

Viscous effects of a hot QGP medium in time dependent magnetic field

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Introduction- Shear and bulk viscosity and importance of time dependent EM fields

Shear viscosity results

Phenomenologically significant quantities



The energy-momentum tensor of the QGP fluid can be expressed as

$$T^{\mu\nu} = \int \frac{d^3p}{(2\pi)^3 p^0} \, p^{\mu} p^{\nu} f(x,p),$$

The spatial components of the dissipative part of the energy momentum tensor can be related to the shear and bulk viscous coefficients in the following manner,

$$\Delta T^{ij} = -\eta W^{ij} - \zeta \delta^{ij} \partial_I u^I, \qquad (1)$$

and the distribution function, $f(x, p) = f_0(x, p) + \delta f(x, p)$





Figure: Evolution of magnetic field in vacuum (blue) and in medium (red).

Heavy-ion collision experiments at RHIC and LHC have shown the presence of magnetic fields. The form of the decay of the magnetic field can be represented as,

$$eB(t) = eB_0 Exp[-rac{t}{ au_B}]$$

In the presence of external EM fields, the away from equilibrium

distribution function is, $\delta f(t) \propto E(t), B(t)$

Writing the Boltzmann equation out into its components, we have

$$\frac{p^{\mu}}{p^{0}}\frac{\partial f_{0}}{\partial x^{\mu}} + \frac{\partial(\delta f(t))}{\partial t} + qE(t) \cdot \frac{\partial f_{0}}{\partial p} + q(v \times B(t))\frac{\partial(\delta f(t))}{\partial p} = -\frac{\delta f(t)}{\tau_{R}}$$

using the following Ansatz for δf

$$\delta f(t) = -\tau_R q E(t) \cdot \frac{\partial f_0}{\partial p} - \Gamma(t) \cdot \frac{\partial f_0}{\partial p},$$

To obtain the shear viscosity, we compare with Eq.1, and read off the coefficient of $-W^{ij}$. The expression reads

$$\eta(t) = -\left[\sum_{f} g_{f}\left(\frac{2}{15}\right) \int \frac{d^{3}p}{(2\pi)^{3}} \frac{p^{4}}{2\epsilon_{f}^{2}} \left(\frac{\partial f^{0}}{\partial \epsilon_{f}}\right) \right. \\ \left. \left(e^{\int (-1/\tau_{R} - iqB(t)/\epsilon)dt} \int dt \, e^{-\int (-1/\tau_{R} - iqB(t)/\epsilon)dt} + e^{\int (-1/\tau_{R} + iqB(t)/\epsilon)dt} \int dt \, e^{-\int (-1/\tau_{R} + iqB(t)/\epsilon)dt}\right) + G.C \right]$$



Consider the magnetic field, $eB(t) = eB_0 Exp[-\frac{t}{\tau_B}]$, in the limiting case we can see that $\lim_{\tau_B \to \infty} eB(t) = eB_0$. Then the factor,

$$\lim_{\tau_B \to \infty} e^{\int (-1/\tau_R - iqB(t)/\epsilon)dt} \int dt \, e^{-\int (-1/\tau - iqB(t)/\epsilon)dt} = \frac{e^{(1/\tau_R - i\omega)t}e^{-(1/\tau_R - i\omega)t}}{(1/\tau_R - i\omega)}$$

$$=rac{1}{(1/ au_{R}-i\omega)}$$

here, $\omega=qB_0/\epsilon.$ Combining the two terms,

$$\frac{1}{(1/\tau_R - i\omega)} + \frac{1}{(1/\tau_R + i\omega)} = \frac{2\tau_R}{(1 + \omega^2 \tau_R^2)},$$

and hence we reduce to the equation in literature for the constant field case,

$$\eta = -\left[\sum_{f} g_{f}\left(\frac{2}{15}\right) \int \frac{d^{3}p}{(2\pi)^{3}} \frac{p^{4}}{2\epsilon_{f}^{2}} \left(\frac{\partial f^{0}}{\partial \epsilon_{f}}\right) \left[\frac{2\tau_{R}}{(1+\omega^{2}\tau_{R}^{2})}\right] + G.C\right]$$



Shear viscosity Vs Temperature



Figure: Temperature dependence of shear viscosity and its dependency on the decay rate of the magnetic field, τ_B at t = 5 fm, $qB_0 = 0.01 \text{ GeV}^2$ and $\mu = 0.1 \text{ GeV}$.



Shear viscosity Vs Time



Figure: Time dependence of shear viscosity and its dependency on the decay rate of the magnetic field, τ_B at T = 0.2 GeV, $qB_0 = 0.01$ GeV² and $\mu = 0.1$ GeV.



Shear viscosity as a function of qB_0 and τ_B



Figure: Decay rate, τ_B and amplitude qB_0 dependency of the shear viscosity at T = 0.2 GeV, t = 5 fm and $\mu = 0.1$ GeV.



The Shear viscosity to entropy density ratio



Figure: Shear viscosity to entropy density as a function of temperature and its dependency on the decay rate of the magnetic field, τ_B at t = 5 fm, $qB_0 = 0.01$ GeV² and $\mu = 0.1$ GeV.



Thermalisation time

The thermalisation time, τ_{th} can be modelled as, $\tau_{th} = 4\pi\eta/Ts$,



Figure: Thermalistaion time, τ_{th} dependency on the decay rate, τ_B , and amplitude, qB_0 , at T = 0.2 GeV, t = 5 fm and $\mu = 0.1$ GeV.

