

Effect of Coriolis Force on Diffusion of D Meson

Ashutosh Dwibedi

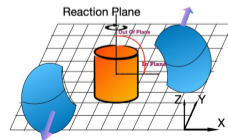
Collaborators: Nandita Padhan, Dipannita Das,
Arghya Chatterjee, Sudipan De, Sabyasachi Ghosh

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Indian Institute of Technology Bhilai

ashutoshd@iitbhilai.ac.in

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Motivation: off-central HIC and vorticity

- In off-central HIC, a huge magnetic field as well as angular velocity can be created.

Becattini et al. Phys. Rev. C 77, 024906, Jiang et al. Phys. Rev. C 94, 044910

$$L_0 = A b \frac{\sqrt{S_{NN}}}{2}$$

- For Au-Au collision at $\sqrt{S_{NN}} = 200$ GeV, $b = 5$ fm, $A = 197 \implies L_0 \sim 10^5 \hbar$
- Most of the angular momentum gets carried away by the spectators. Imagining only 10% or 0.1 part of this is deposited in the reaction zone,
 $\implies L = 0.1 L_0 \sim 10^4 \hbar$.
- Average angular velocity $\Omega_{\max} = 0.06 - 0.07 \text{ fm}^{-1}$

Literatures on the rotating nuclear matter

- Studies on the thermodynamics of the QCD matter rotating with a **global angular velocity $\vec{\Omega}$**

NJL Model: Wang et al. Phys. Rev. D 99, 016018

Quark-meson model: Chen et al. Phys. Rev. D 108, 054006

linear sigma model: Singha et al. Phys. Rev. D 110, 094053

lattice QCD: Yamamoto et al. Phys. Rev. Lett. 111, 081601

- Studies on transport properties like, shear viscosity, electrical conductivity, etc

Dwibedi et al., Phys. Rev. C 109, 034913, Phys. Rev. C 109, 034914, Phys. Rev. C 110, 024904

- Question: What happens to a heavy quark/ meson diffusing in the rotating QCD matter?

Model description: The BTE in RTA

- D meson diffusing under the background of rotating nuclear matter

$$p^\mu \frac{\partial f^0 + \delta f}{\partial x^\mu} - \Gamma_{\mu\lambda}^\alpha p^\mu p^\lambda \frac{\partial f^0 + \delta f}{\partial p^\alpha} = -(u^\alpha p_\alpha) \frac{\delta f}{\tau_c}, \quad (1)$$

eqbm. distribution = $f^0 = \frac{1}{e^{(p^\alpha u_\alpha - \mu_D)/T} - 1}$,

$u^\alpha = (\frac{1}{\sqrt{g_{00}}}, 0)$, $g_{\mu\nu}$ = metric tensor

Effective force $\equiv F^\alpha = -\frac{1}{m_D} \Gamma_{\mu\lambda}^\alpha p^\mu p^\lambda = \frac{dp^\alpha}{d\tau}$, τ = proper time

deviation from eqbm. = δf

relaxation time of D meson = τ_c

Model description: heavy meson current and spatial diffusion coefficients

- Heavy meson current J^i coming from the variation of its chemical potential μ_D ,

$$J^i = \int \frac{d^3\vec{p}}{(2\pi)^3} \frac{p^i}{\rho_0} \delta f, \quad (2)$$

$$J^i = -\sigma^{ij} \nabla_j \mu_D, \quad (3)$$

σ^{ij} is heavy meson conductivity.

- Solving BTE for δf leads to J^i (equivalently σ^{ij}). Spatial diffusion coefficients can be evaluated using the relation

$$D^{ij} = \frac{\sigma^{ij}}{\chi}, \quad \chi = \text{susceptibility}. \quad (4)$$

- Everything boils down to evaluate δf from BTE. Any force term in BTE?

Model description: Forces in rotating frame and expression of diffusion coefficients

- The relativistic EOM in the rotating frame,

$$\frac{d\vec{p}}{dt} = \gamma_v m_D (\vec{\Omega} \times \vec{r}) \times \vec{\Omega} + 2\gamma_v m_D (\vec{v} \times \vec{\Omega}), \quad (5)$$

$$\text{Centrifugal force} = \gamma_v m_D (\vec{\Omega} \times \vec{r}) \times \vec{\Omega},$$

$$\text{Coriolis force} = 2\gamma_v m_D (\vec{v} \times \vec{\Omega})$$

- Solving BTE with the inclusion of Coriolis force we get,

$$D^{\parallel} = \frac{\int \frac{d^3 p}{(2\pi)^3} \tau_c \times \frac{p^2}{E^2} f_0 (1 + f_0)}{3 \int \frac{d^3 p}{(2\pi)^3} f_0 (1 + f_0)}, \quad D^{\perp} = \frac{\int \frac{d^3 p}{(2\pi)^3} \tau_c^{\perp} \times \frac{p^2}{E^2} f_0 (1 + f_0)}{3 \int \frac{d^3 p}{(2\pi)^3} f_0 (1 + f_0)},$$
$$D^{\times} = \frac{\int \frac{d^3 p}{(2\pi)^3} \tau_c^{\times} \times \frac{p^2}{E^2} f_0 (1 + f_0)}{3 \int \frac{d^3 p}{(2\pi)^3} f_0 (1 + f_0)}, \text{ defined relative to } \vec{\Omega} \quad (6)$$

Model description: Determination of τ_C

- Focus on the hadronic temperature domain

The rotating hadron gas background $\xrightarrow{\text{modeled}}$ The ideal hadron resonance gas (HRG) model.

τ_C of D meson-hadron scatterings $\xrightarrow{\text{modeled}}$ hard-sphere type scatterings,

$$\tau_C = 1 / (n_{\text{HRG}} v_{\text{av}} \pi a^2), \quad (7)$$

$$n_{\text{HRG}} = \sum_B g_B \int_0^\infty \frac{d^3 p}{(2\pi)^3} \frac{1}{e^{E/T} + 1} + \sum_M g_M \int_0^\infty \frac{d^3 p}{(2\pi)^3} \frac{1}{e^{E/T} - 1}, \quad (8)$$

- a = scattering length, v_{av} = thermal average velocity of D meson
- We will calibrate τ_C by changing a to cover the existing model calculations of D^{\parallel} (\equiv diffusion coefficient at $\vec{\Omega} = 0$).

Result and Conclusion: variation of anisotropic diffusion coefficients

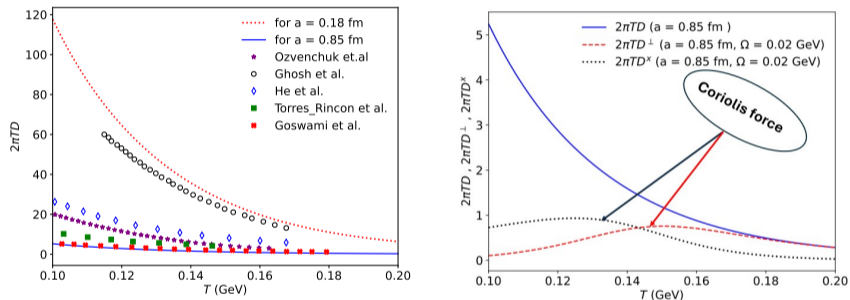


Figure: Calibration of relaxation time of D meson (left) and variation of anisotropic diffusion coefficients (right)

- Anisotropic spatial diffusion coefficients ($D^{\perp,\times}$) can differ significantly from D^{\parallel} at low temperatures. This can affect the experimental observables like R_{AA} and v_2 of D meson.

References

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Thank You!