Exploring the flow harmonic correlations via multi-particle Symmetric and Asymmetric Cumulants in Au+Au collisions at $\sqrt{s_{NN}}$ = 200 GeV

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Outline

Introduction

2 Symmetric and Asymmetric Cumulants of Flow Amplitudes

Error estimation

Framework

Results





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Anisotropic flow is a sensitive probe both of initial conditions in heavy-ion collisions, and of QGPs transport properties (e.g. of its shear viscosity)



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^ahttps://cerncourier.com/a/going-with-the-flow

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Anisotropi	c Flow				

Fourier series to describe anisotropic emission of particles in the plane transverse to the beam direction after every heavy-ion collision. ¹

$$E\frac{d^{3}N}{dp^{3}} = \frac{d^{2}N}{2\pi p_{t}dp_{t}dy} \left[1 + 2\sum_{n=1}^{\infty} v_{n} \cos\left[n(\phi - \psi_{n})\right] \right]$$

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- v_n : flow amplitudes
- ψ_n : symmetry planes

Anisotropic flow is quantified with v_n and ψ_n .

- *v*₁ : directed flow
- v_2 : elliptic flow
- v_3 : triangular flow
- *v*₄ : quadrangular flow, etc.

¹Poskanzer, Voloshin, Phys.Rev.C 58 (1998) 1671-1678

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Traditional	flow observables				

• v_n are insensitivite to temperature dependence of η/s^2



²H. Niemi, K. J. Eskola, R. Paatelainen, Phys. Rev. C 93, 024907 (2016)

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Multi-par	ticle correlations				
Two-narticle	correlations				



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Decomposition of the two-particle distribution into uncorrelated and correlated components. $^{\rm 3}$

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³Borghini, Dinh, Ollitrault, PhysRevC.63.054906

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Multi-par	ticle correlations				
Three-partic	le correlations				

$$f(\mathbf{p}_{1}, \mathbf{p}_{2}, \mathbf{p}_{3}) = f(\mathbf{p}_{1})f(\mathbf{p}_{2})f(\mathbf{p}_{3}) + f_{c}(\mathbf{p}_{1}, \mathbf{p}_{2})f(\mathbf{p}_{3}) + f_{c}(\mathbf{p}_{1}, \mathbf{p}_{3})f(\mathbf{p}_{2}) + f_{c}(\mathbf{p}_{2}, \mathbf{p}_{3})f(\mathbf{p}_{1}) + f_{c}(\mathbf{p}_{1}, \mathbf{p}_{2}), \mathbf{p}_{3})$$

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Decomposition of the three-particle distribution.⁴

⁴Borghini, Dinh, Ollitrault, PhysRevC.63.054906

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Multi-par Four-particle	ticle correlations				
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Decomposition of four-particle distribution.⁵

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$$\begin{aligned} & \mathcal{E}_{n}\{4\} & \equiv & \langle \langle e^{in(\varphi_{1}+\varphi_{2}-\varphi_{3}-\varphi_{4})} \rangle \rangle \\ & - & \langle \langle e^{in(\varphi_{1}-\varphi_{3})} \rangle \rangle \langle \langle e^{in(\varphi_{2}-\varphi_{4})} \rangle \rangle \\ & - & \langle \langle e^{in(\varphi_{1}-\varphi_{4})} \rangle \rangle \langle \langle e^{in(\varphi_{2}-\varphi_{3})} \rangle \rangle. \end{aligned}$$

$$c_n{4} \equiv \langle v_n^4 \rangle - 2 \langle v_n^2 \rangle^2$$

$$v_n{4} = \sqrt[4]{-c_n{4}}$$

⁵Borghini, Dinh, Ollitrault, PhysRevC.63.054906

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Cumulants	definitions				

The Symmetric Cumulants, defined below, can probe the genuine correlations between different flow harmonics $^{\rm 6}$

$$SC(m,n) \equiv \langle v_m^2 v_n^2 \rangle_c = \langle v_m^2 v_n^2 \rangle - \langle v_m^2 \rangle \langle v_n^2 \rangle$$
$$= \langle \langle e^{i(m\varphi_1 + n\varphi_2 - m\varphi_3 - n\varphi_4)} \rangle \rangle - \langle \langle e^{i(m\varphi_1 - n\varphi_3)} \rangle \rangle \langle \langle e^{i(m\varphi_2 - n\varphi_4)} \rangle \rangle$$

The Asymmetric Cumulants, which are more generalized, probe the genuine correlations between the different moments of different flow harmonics: ⁷

$$\begin{aligned} \mathrm{AC}_{2,1}(m,n) &\equiv \langle (v_m^2)^2 v_n^2 \rangle_c \equiv \langle v_m^4 v_n^2 \rangle_c, \\ &= \langle v_m^4 v_n^2 \rangle - \langle v_m^4 \rangle \langle v_n^2 \rangle - 2 \langle v_m^2 v_n^2 \rangle \langle v_m^2 \rangle + 2 \langle v_m^2 \rangle^2 \langle v_n^2 \rangle, \\ \mathrm{AC}_{3,1}(m,n) &\equiv \langle (v_m^2)^3 v_n^2 \rangle_c \equiv \langle v_m^6 v_n^2 \rangle_c \end{aligned}$$

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 6 Bilandzic, Christensen, Gulbrandsen, Hansen, Zhou, Phys.Rev.C 89 (2014) 6, 064904 7 Bilandzic, Lesch, Mordasini, Taghavi, Phys.Rev.C 105 (2022) 2, 024912 $$<\!\!\!\!\!\!\!\!\!\!\!\!\!\!\!$

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To eliminate the effect of the magnitudes of v_m and v_n on the cumulants, we define the normalized cumulants. This enables us to compare data and model calculations in a quantitative way and compare the fluctuations of the initial and final states. The normalized symmetric and asymmetric cumulants are defined as: ⁸

$$NSC(m, n) = \frac{SC(m, n)}{\langle v_m^2 \rangle \langle v_n^2 \rangle}$$

NAC_{2,1}(*m*, *n*) =
$$\frac{AC_{2,1}(m, n)}{\langle v_m^2 \rangle^2 \langle v_n^2 \rangle}$$

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⁸Taghavi, Eur.Phys.J.C 81 (2021) 7, 652

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The Bootst	rap method: The calculation of sa	mpling variar	ice		

Let \hat{O} be the observable on which we intend to find the standard error.

- Given a parent sample of size *n*, we construct *B* independent bootstrap samples $X_1^*, X_2^*, \ldots, X_B^*$, each with *n* data points randomly drawn with replacement.
- We evaluate the observable for each bootstrap sample.

$$\hat{O}_b^* = \hat{O}(X_b^*), \qquad b = 1, 2, ..., B.$$

• The sampling variance of the observable can then be calculated as follows:

$$Var(\hat{O}) = \frac{1}{B-1} \sum_{b=1}^{B} (\hat{O}_{b}^{*} - \bar{\hat{O}}),$$

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where,
$$\overline{\hat{O}} = \frac{1}{B} \sum_{b=1}^{B} \hat{O}_{b}^{*}$$
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Sensitivity to hydrodynamic transport coefficients



- The large sensitivity of symmetric cumulants primarily arises from the characteristics of anisotropic flow, as the symmetric cumulants involve higher powers of flow coefficients.
- NSC(2,3) and NAC_{2,1}(2,3) are insensitive to the hydro model parameters. Reliable for constraining the initial state of the system's evolution.
- NSC(2, 4) and NAC_{2,1}(2, 4) show considerable sensitivity to these parameters, allowing them to effectively constrain hydrodynamic models.

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Sensitivity to hadronic interactions



- The large sensitivity of symmetric cumulants primarily arises from the characteristics of anisotropic flow, as the symmetric cumulants involve higher powers of flow coefficients.
- NSC(2,3) and NAC_{2,1}(2,3) are insensitive to late-stage hadronic interactions. Reliable for constraining the initial state of the systems evolution.
- NSC(2, 4) and NAC_{2,1}(2, 4) show considerable sensitivity to these stages, allowing them to effectively constrain transport models.

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Sensitivity to hadronic interactions



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Centrality dependence of *SC*(*m*, *n*)



- (Normalized) Symmetric cumulant vs centrality from four-particle correlations compared with STAR data ¹⁰ in Au+Au collisions at $\sqrt{s_{NN}}$ = 200 GeV. SC(2,3) is consistently negative while as SC(2,4) is consistently positive.
- Note that we have shown the Glauber results for NSC(2, 4) to highlight that in non-central collisions v_4 is mainly driven by ε_2^2 and not ε_4 .

¹⁰Adam et al., Phys.Lett.B 783 (2018) 459-465

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Centrality dependence of $AC_{2,1}(m, n)$



- (Normalized) Asymmetric cumulant vs centrality from six-particle correlations in Au+Au collisions at $\sqrt{s_{NN}} = 200$ GeV. AC_{2,1}(2,3) is consistently negative while as AC_{2,1}(2,4) is consistently positive.
- Note that we have shown the Glauber results for NAC_{2,1}(2,4) to highlight that in non-central collisions v_4 is mainly driven by ε_2^2 and not ε_4 .

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$\sqrt{s_{NN}}$ dependence



- The values of SC(*m*, *n*) and AC_{2,1}(*m*, *n*) at ALICE ¹¹ ¹² are larger-by factors ranging from 3 to 8-compared to those at STAR ¹³.
- Important to note the similarities in the magnitude and centrality dependence of the normalized cumulants. Minimal energy dependence for the normalized cumulants.

 ¹¹Acharya et. al., Phys.Rev.C 108 (2023) 5, 055203
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- v_n are insensitivite to temperature dependence of η/s .
- NSC(2, 3) and NAC_{2,1}(2, 3) are to both the hydro model parameters (shear and bulk viscosity) and late-stage hadronic interactions. Reliable for constraining the initial state of the systems evolution.
- NSC(2, 4) and NAC_{2,1}(2, 4) show considerable sensitivity to these stages, allowing them to effectively constrain hydrodynamic and transport models.
- The values of SC(*m*, *n*) and AC_{2,1}(*m*, *n*) at ALICE are larger-by factors ranging from 3 to 8-compared to those at STAR.
- Important to note the similarities in the magnitude and centrality dependence of the normalized cumulants. Minimal energy dependence for the normalized cumulants.

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Thank you for your attention!

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