



ATHIC 2025



# Experimental overview on the open heavy flavor measurement in relativistic heavy-ion collisions



KoALICE



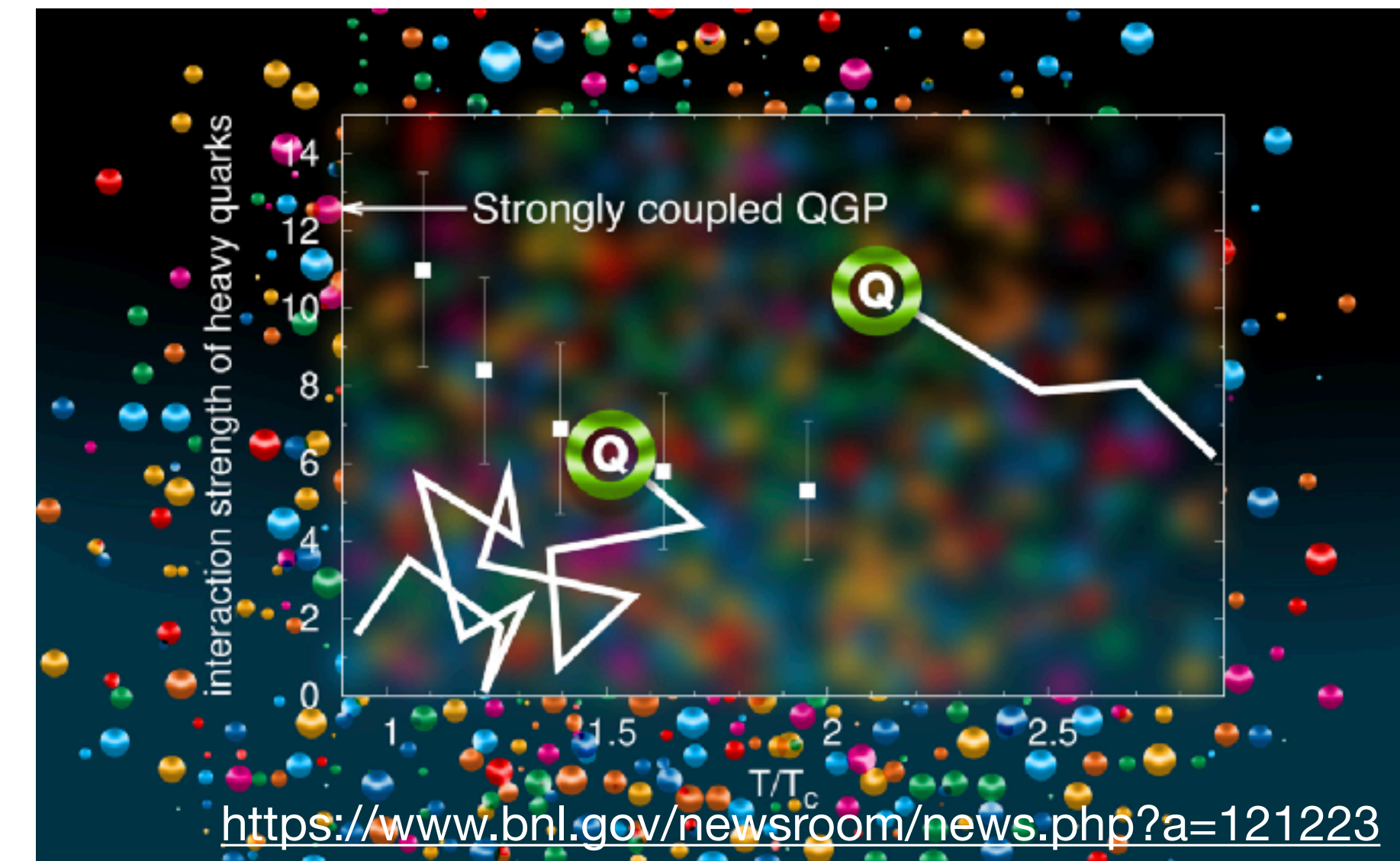
MinJung Kweon

Inha University

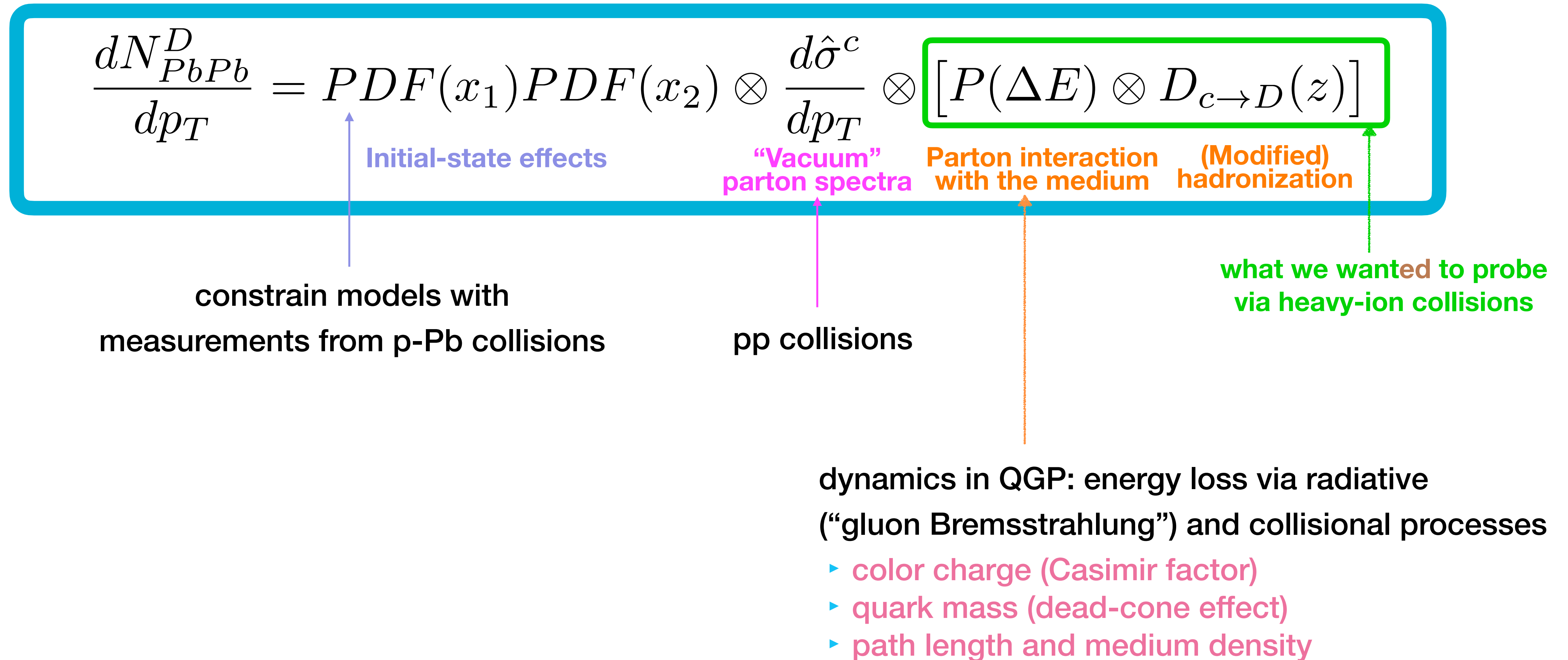
ATHIC 2025, 2025. 1.14

# Suitable probes to understand QCD medium, heavy quarks

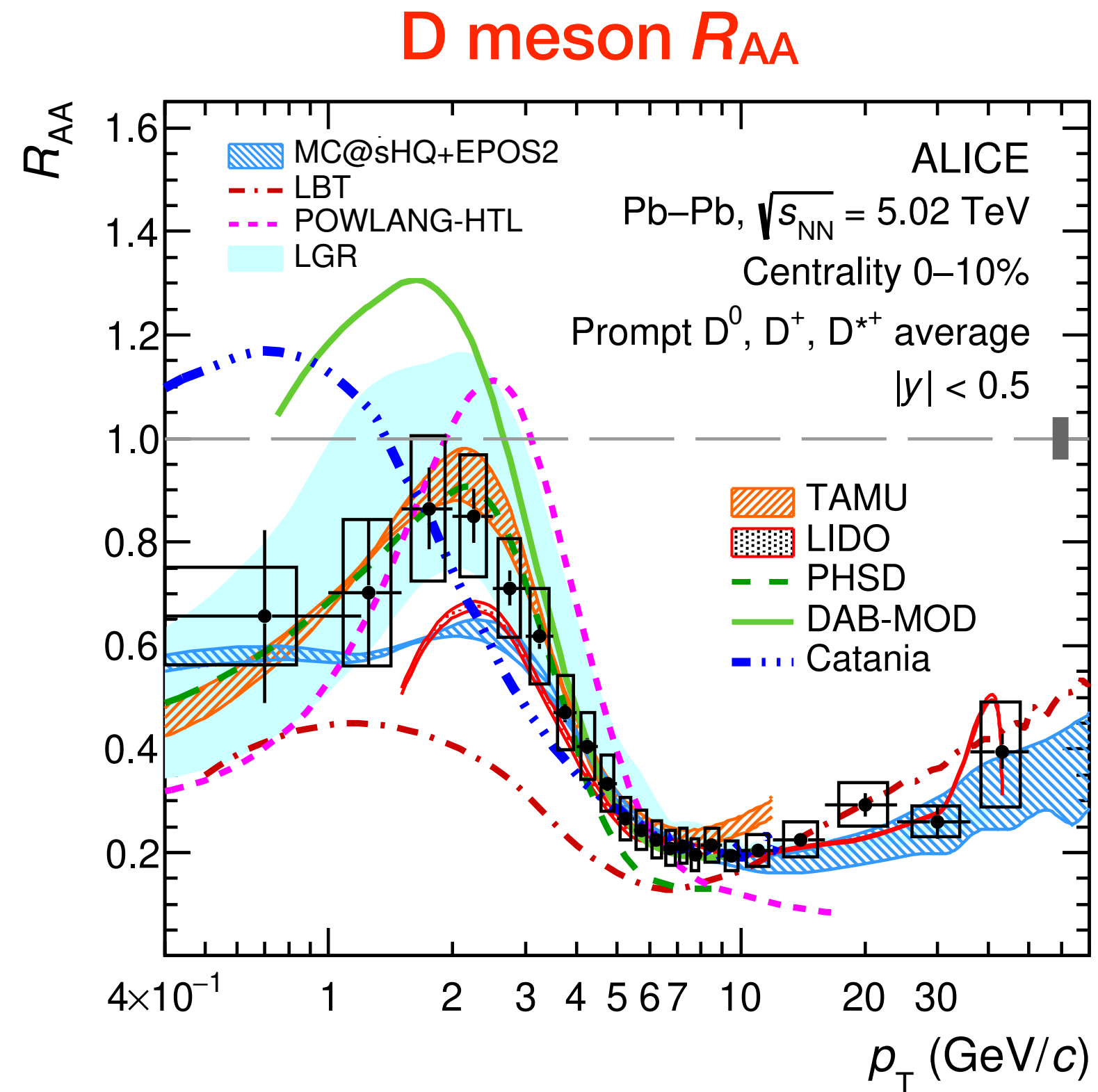
- Heavy-quark (HQ) mass is much larger than the nonperturbative QCD scale → produced mainly in initial hard scatterings (reasonably well described by perturbative QCD)
- **However, the story is not easy since what we measure is ...**
  - dynamics of heavy quarks from their creation at the onset of a heavy-ion collision through their evolution in the QCD medium until their detection as heavy hadrons
- We need a comprehensive description of the initial production of the heavy quarks, their interactions with the QGP, hadronization, and the interactions of heavy hadrons in the hadronic phase → rather complex to describe using first-principles QCD!



# Heavy flavour production in medium: what we measure



# Heavy flavour production in medium: where to see, what we see



- Very strong coupling between heavy quarks and medium
- Strong suppressions are fairly described by transport models
  - nPDF changes the total yields
  - Suppression at low  $p_T$  via collisional energy loss
  - Suppression at high  $p_T$  via radiative energy loss ( $\Delta E/E$  decreases)
  - Push by radial flow (low  $p_T$  to high  $p_T$ )
  - Hadronization picks light quark kinematics
- However, describing full  $p_T$  is still challenging...

ALI-PUB-501952

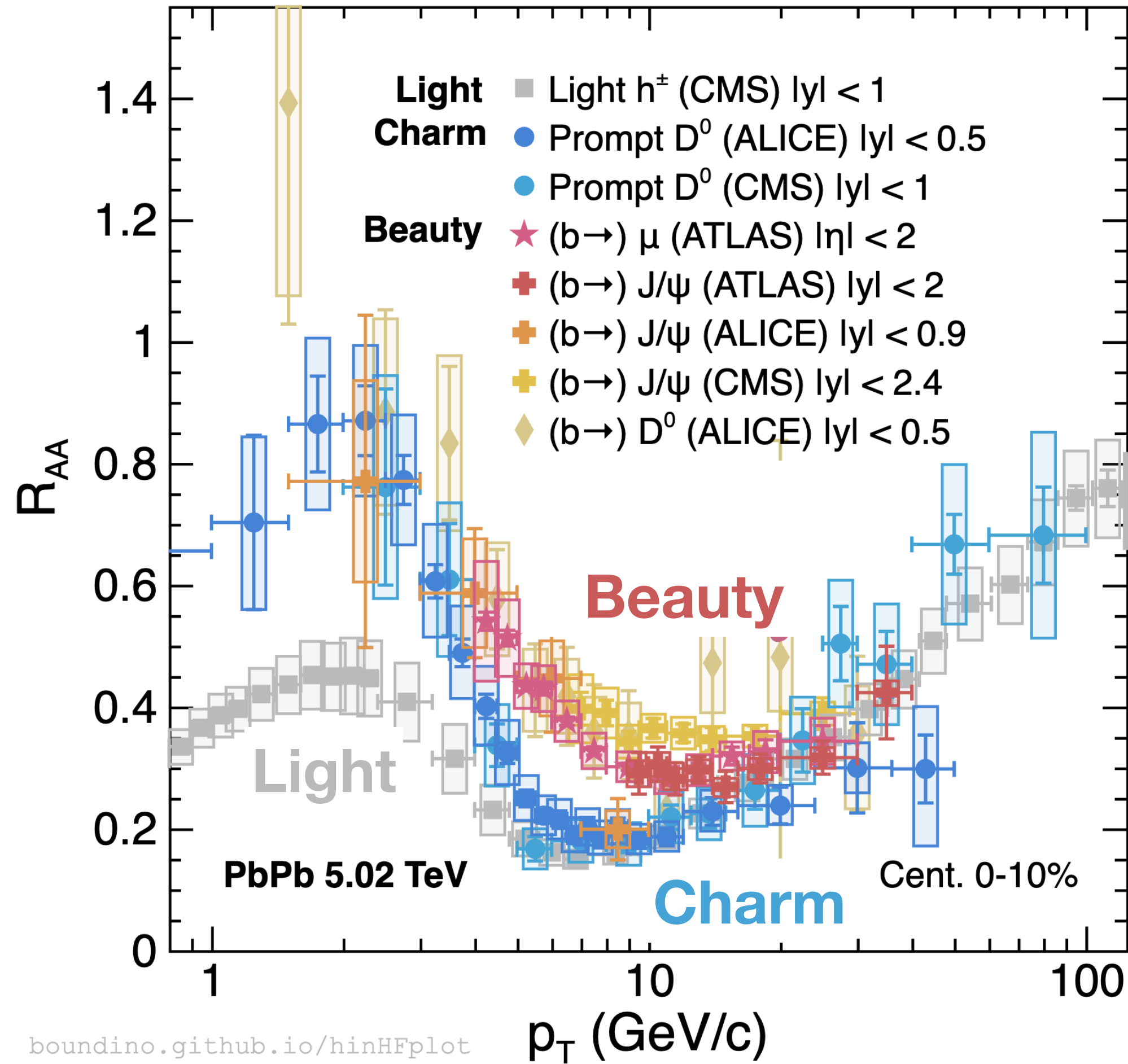
collective flow, hadronisation, nuclear PDF

collisional E loss

radiative E loss

JHEP 01 (2022) 174

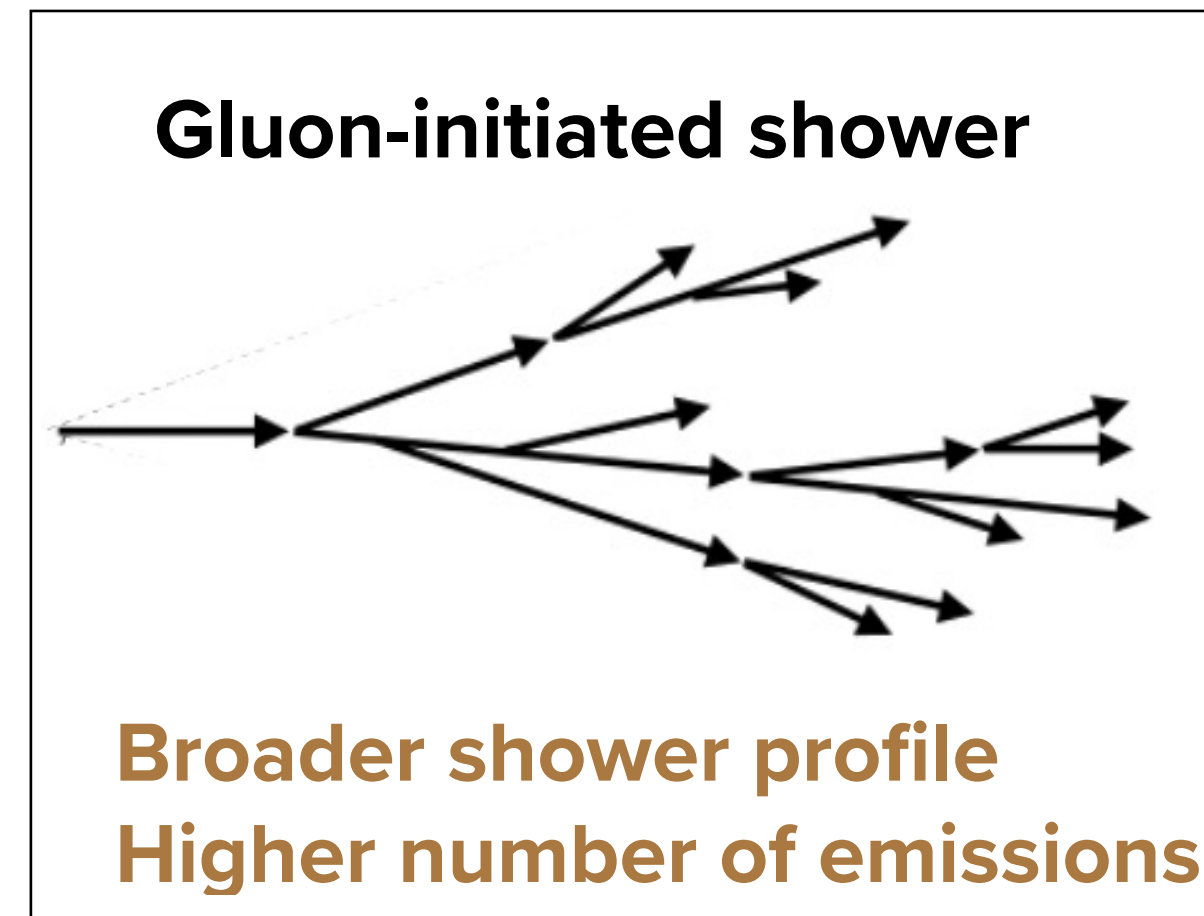
# Mass and color charge dependence of energy loss: via $R_{AA}$



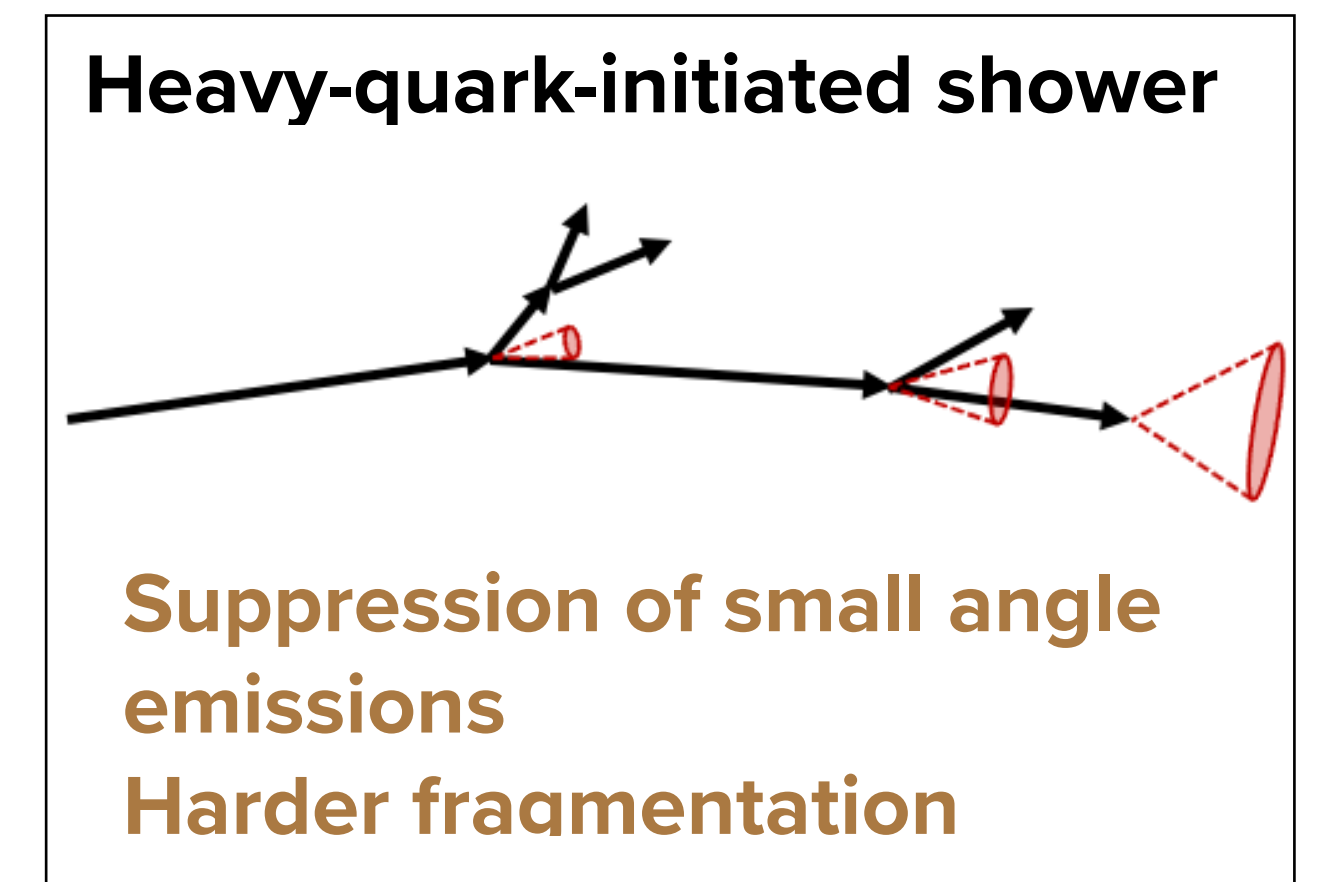
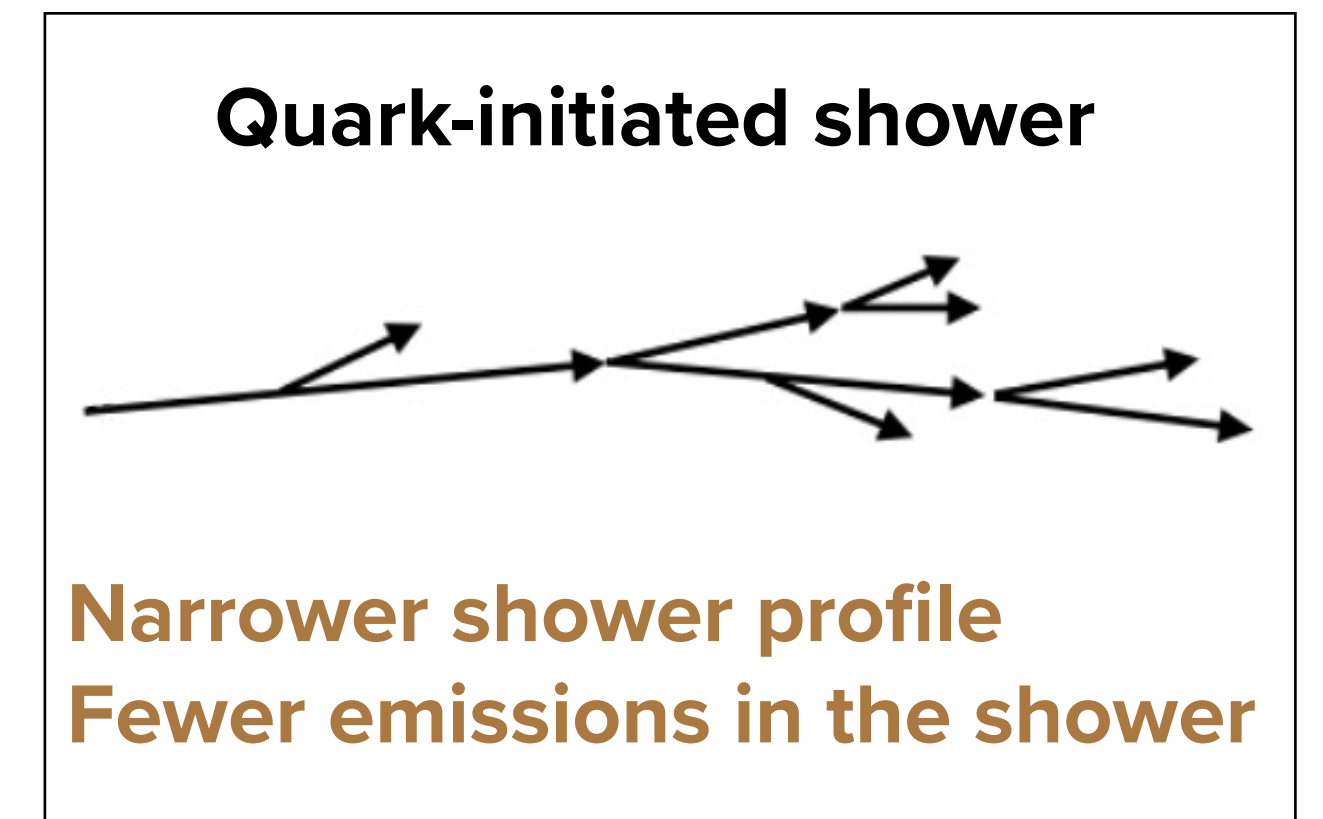
ATLAS PLB 829 (2022) 137077, ATLAS EPJC 78 (2018) 762  
 ALICE HEP 02 (2024) 066, ALICE JHEP 12 (2022) 126  
 CMS EPJC 78 (2018) 509

- Mass and color charge dependent suppression

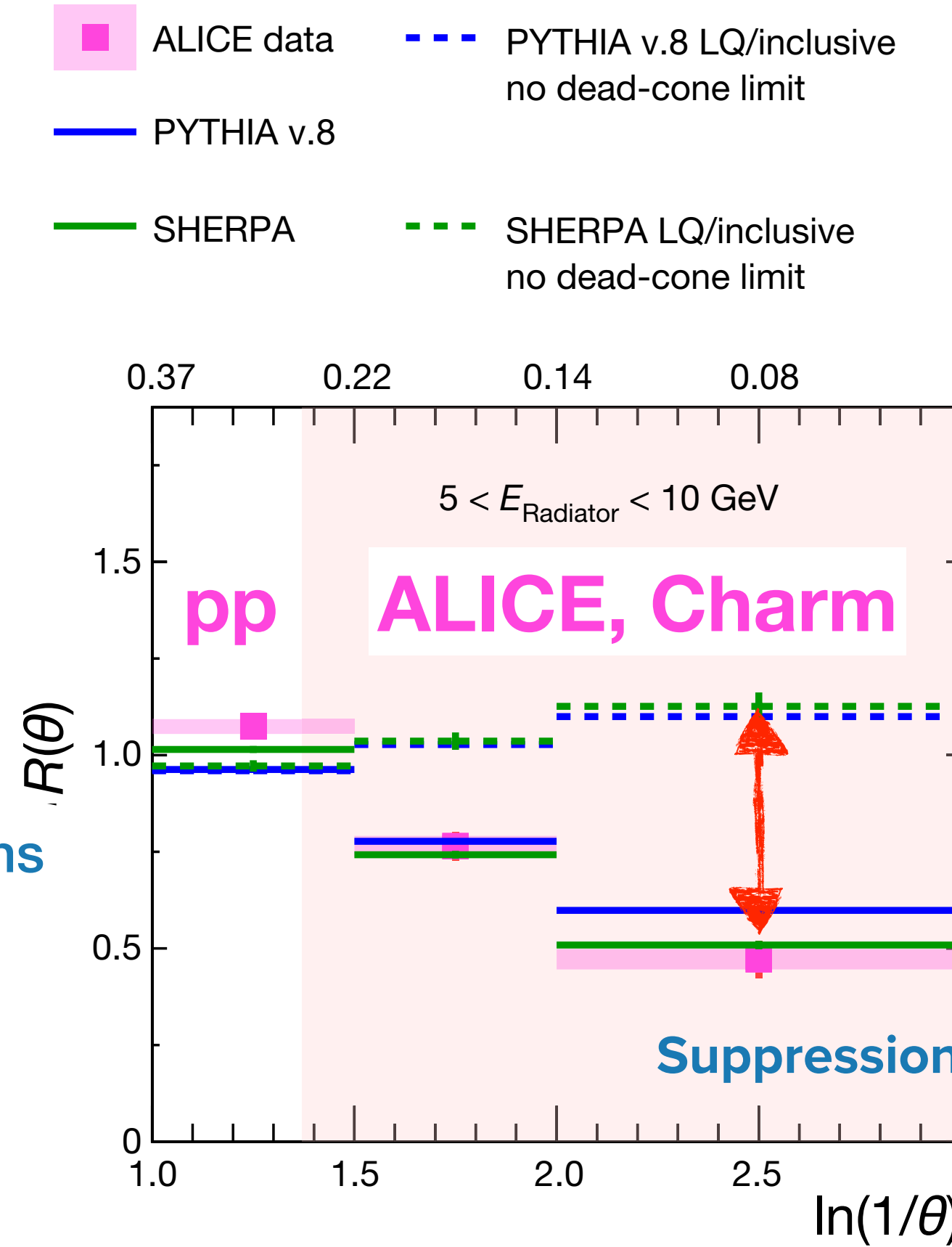
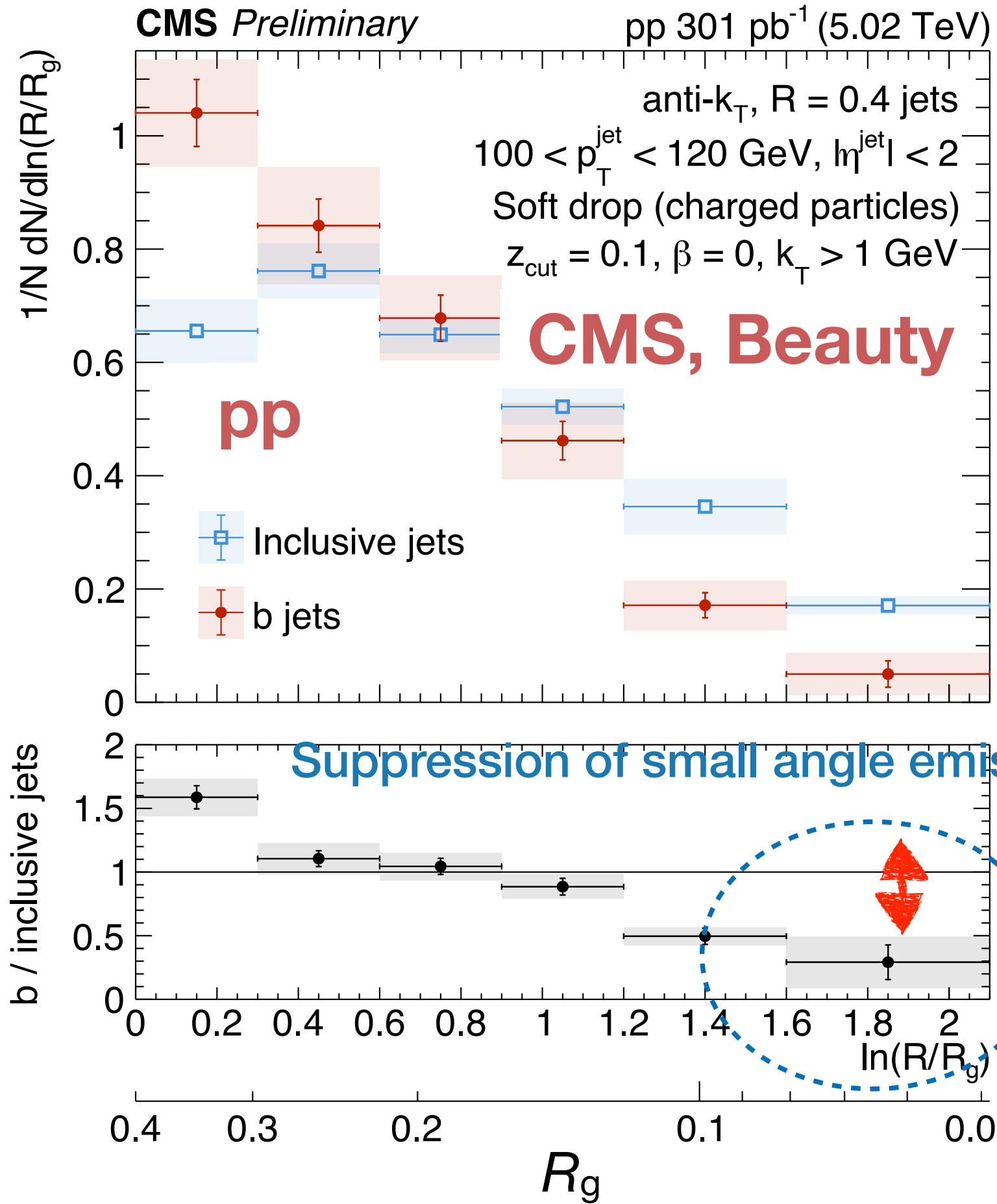
In terms of energy loss  
 Gluons  $>$  Light Quarks (u, d, s)  $>$  Charm Quarks (c)  $>$  Beauty Quarks (b)



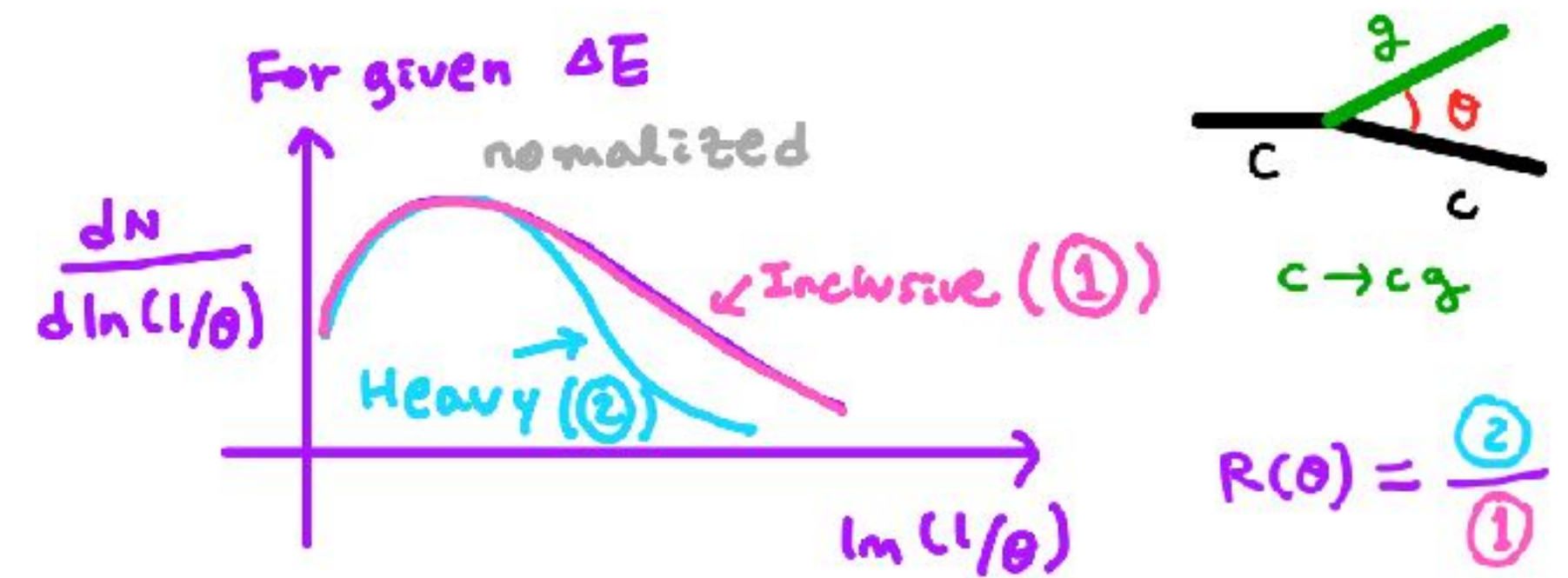
relevant to light quark suppression



# Mass dependence of energy loss: via dead cone



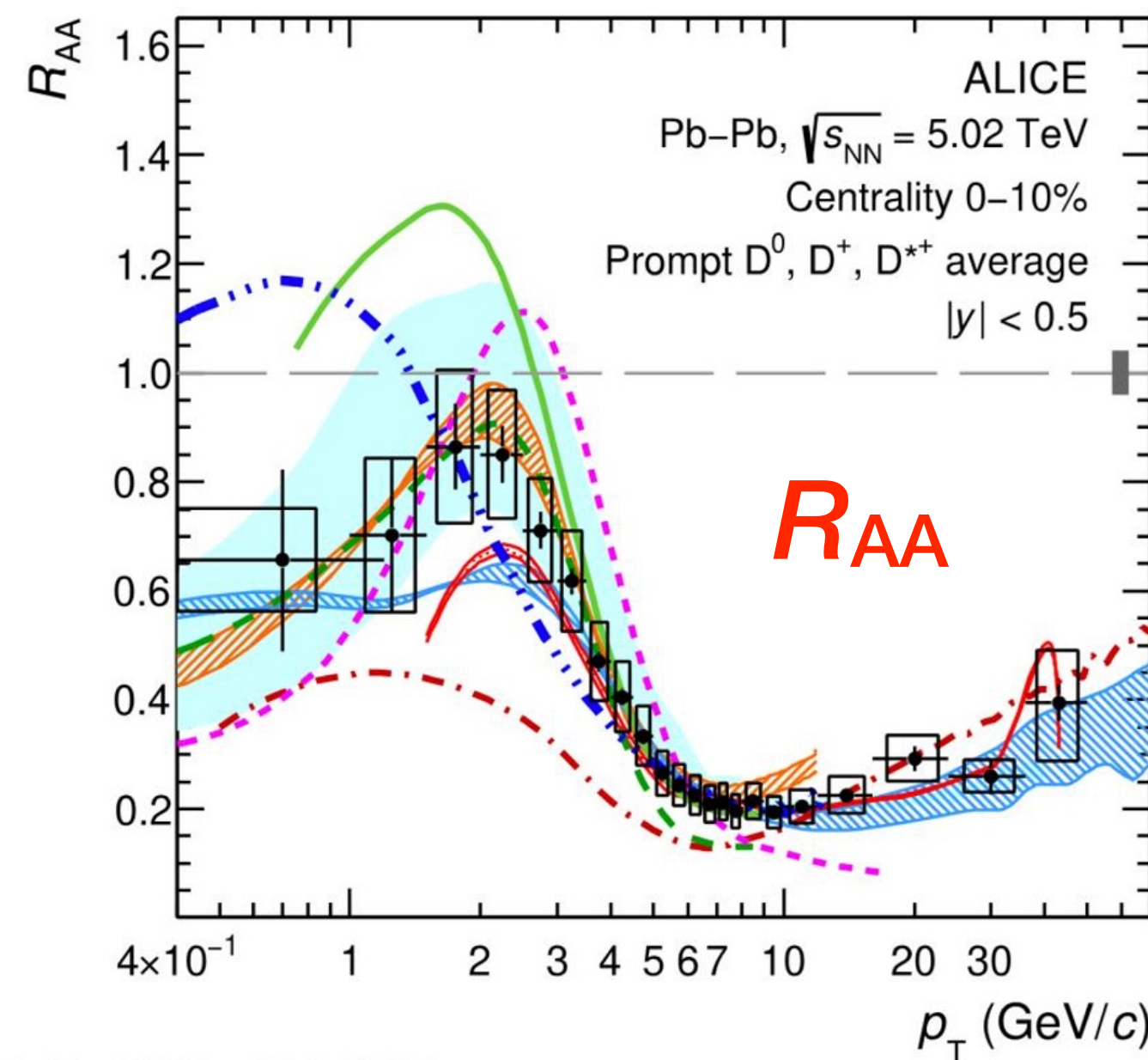
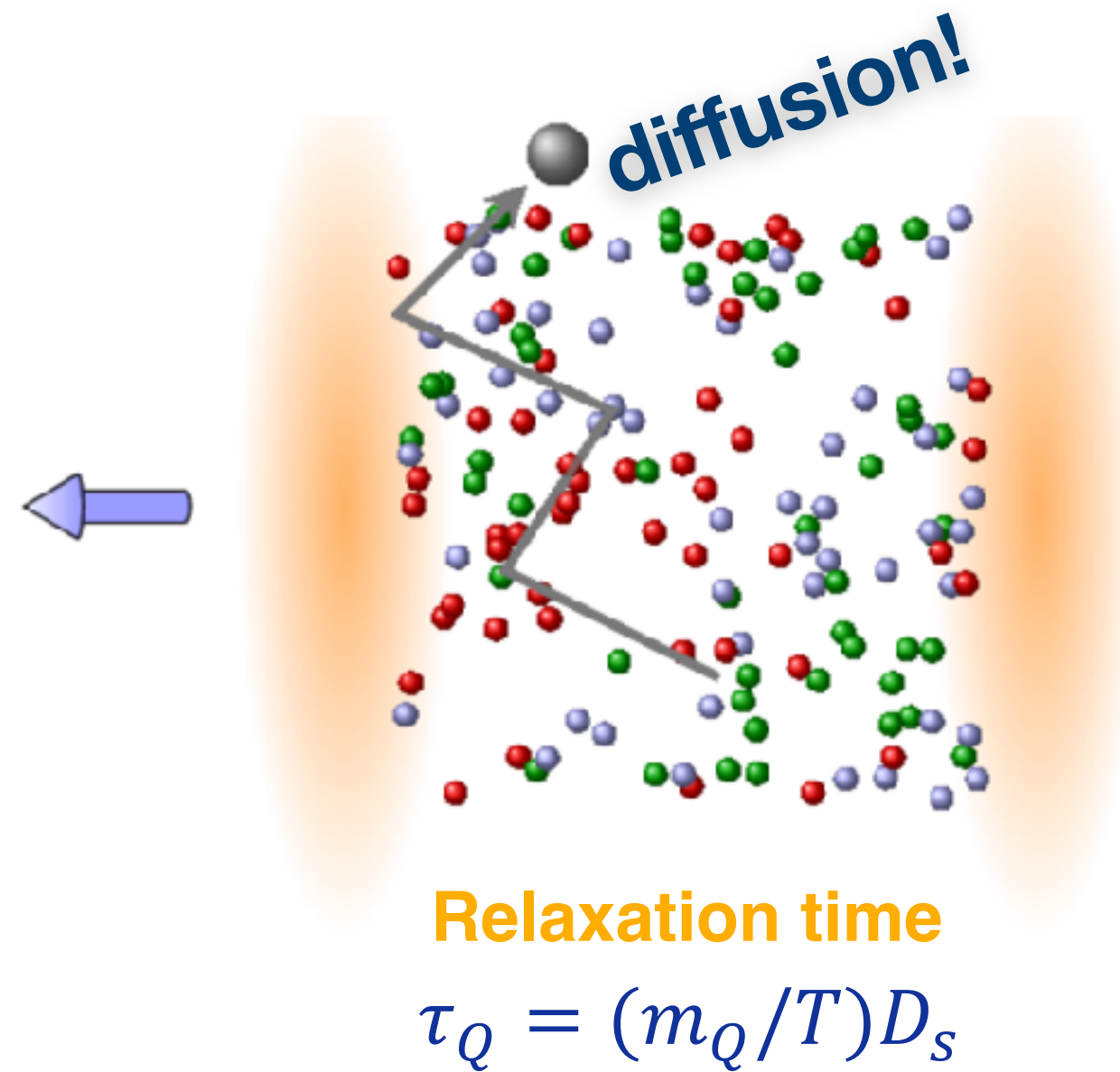
The ratio of the splitting angle ( $\theta$ ) distributions for D<sup>0</sup>-meson tagged jets & inclusive jets



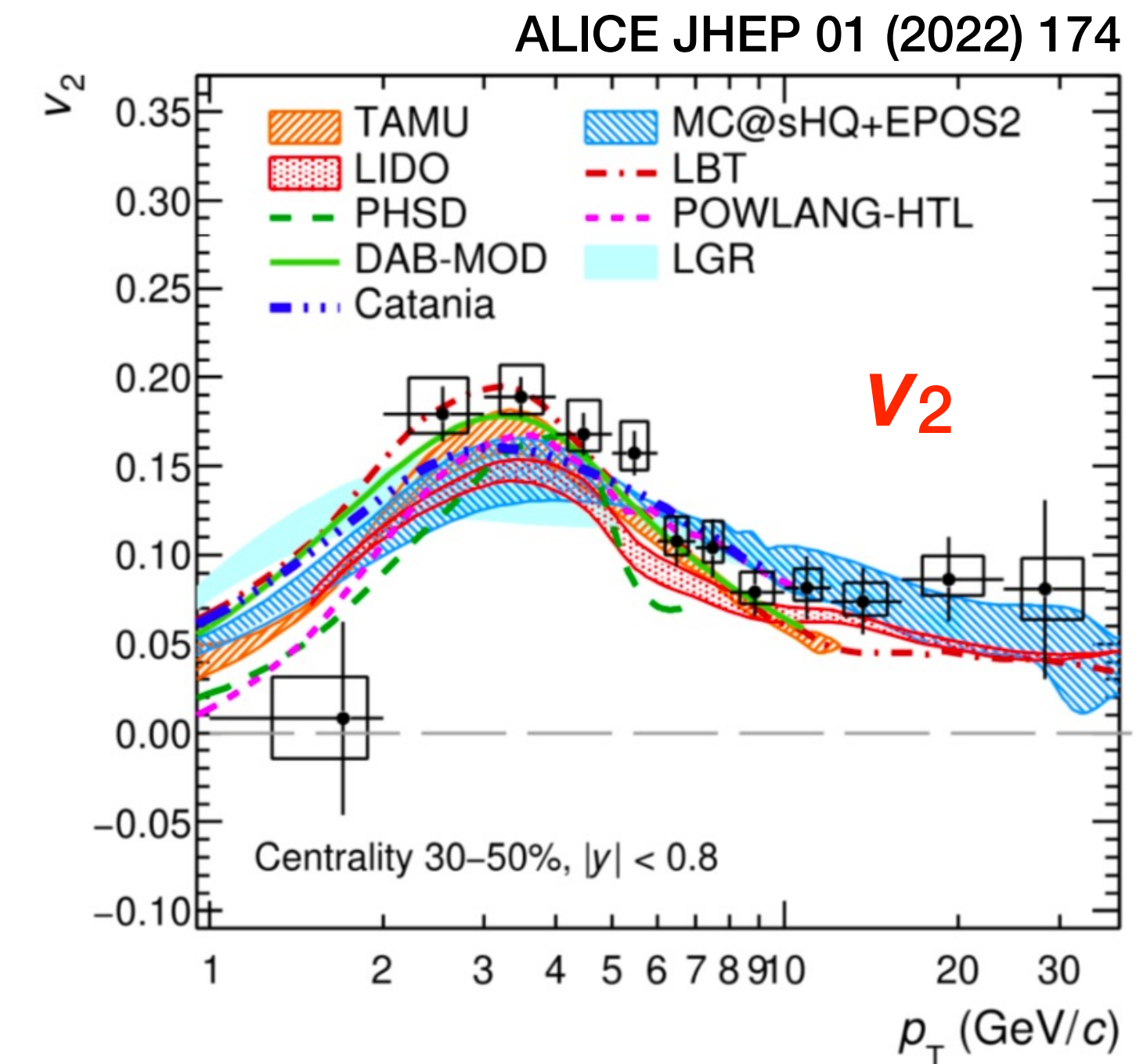
- Direct observation of dead cone in pp collisions
- The medium-induced radiation in a QGP can alter the same observables → to be understood to separate the contributions of parton mass effects and QGP-induced modifications.

# Charm quark transport in the medium

- Low- $p_T$  region provides insight into the heavy-quark interactions with the medium (by diffusion, analogous to a ‘Brownian motion’)
- Charm quark interacts with the medium via collisional and radiative processes in heavy-ion collisions → comparison of  $R_{AA}$  and  $v_2$  with transport models → constraint on the diffusion coefficient



ALI-PUB-498687

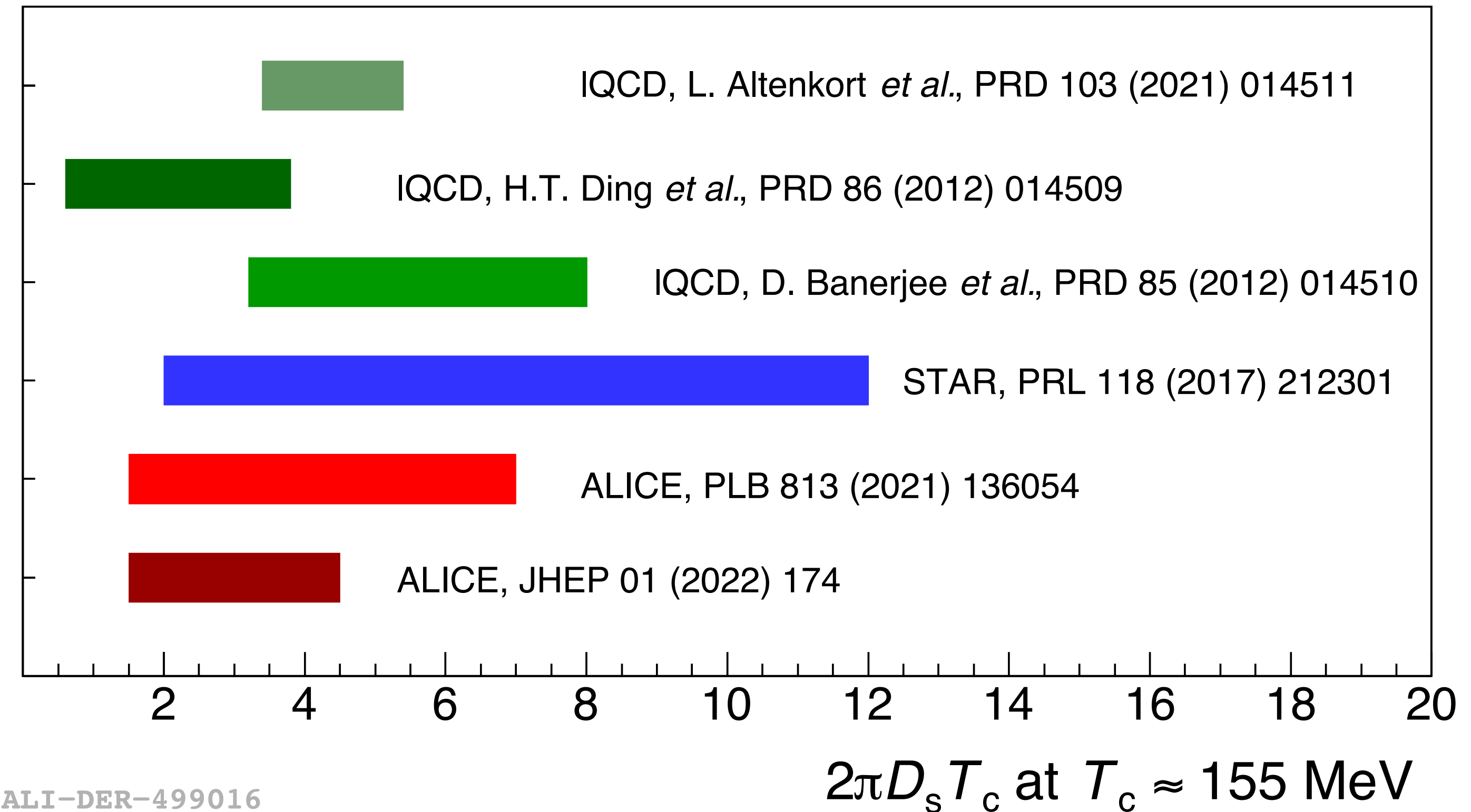
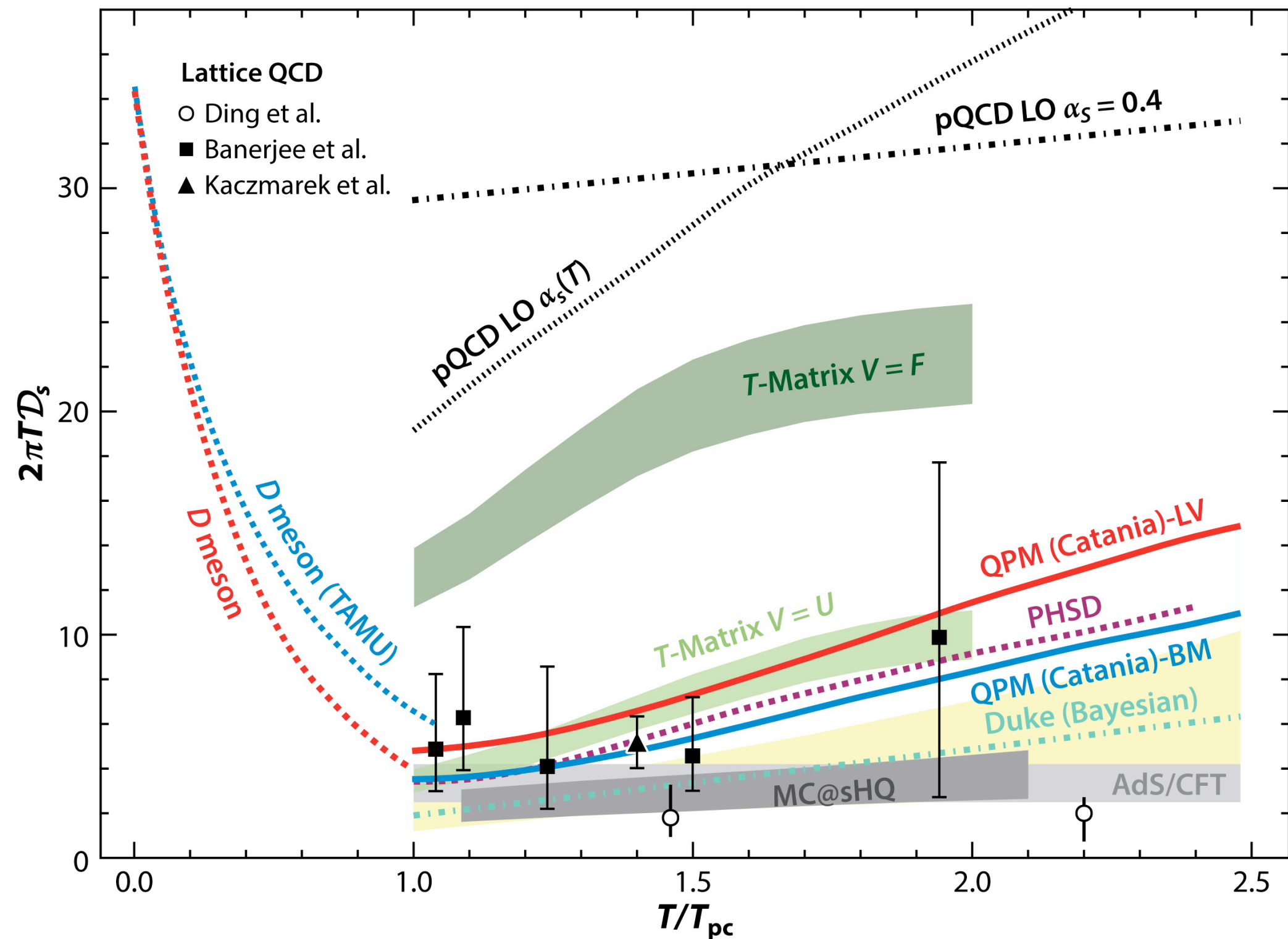


At low- $p_T$ , also shadowing & bulk evolution of the medium

- Charm quarks are **thermalised** with medium (having different thermalization timescales due to their masses) → collective motion (slower thermalization leading to smaller  $v_2$ )

# Quantitative information via spacial diffusion coefficient

X. Dong et al., Ann.Rev.Nucl.Part.Sci. 69 (2019) 417



ALI-DER-499016

$2\pi D_s T_c$  at  $T_c \approx 155$  MeV

## Relaxation time

$$\tau_Q = (m_Q/T)D_s \quad 1.5 < 2\pi D_s T_c < 4.5 \rightarrow \text{direct access to heavy-flavour relaxation time: } \tau_{\text{charm}} \sim 3\text{-}8 \text{ fm}/c$$

Note: hadronization is hard to control in the model...



# Heavy flavour production in medium: hadronization

Going back to the original assumption...

$$\frac{dN_{PbPb}^D}{dp_T} = \underbrace{PDF(x_1)PDF(x_2)}_{\text{Initial-state effects}} \otimes \underbrace{\frac{d\hat{\sigma}^c}{dp_T}}_{\text{"Vacuum" parton spectra}} \otimes \underbrace{[P(\Delta E) \otimes D_{c \rightarrow D}(z)]}_{\substack{\text{Parton interaction} \\ \text{with the medium} \quad \text{(Modified)} \\ \text{hadronization}}}$$

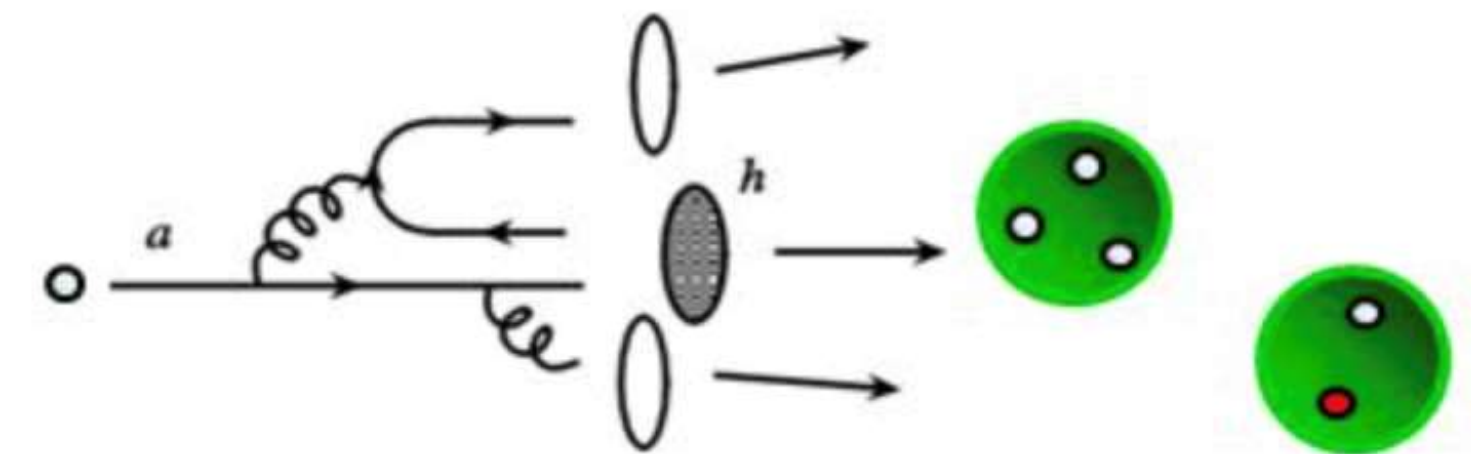
- Fragmentation functions  $D(z)$  are phenomenological functions to parameterize the non-perturbative parton-to-hadron transition

- $z$  = fraction of the parton momentum taken by the hadron  $h$
  - Do not specify the hadronisation mechanism

- Parametrized on data and assumed to be **“universal”**

- In A-A collisions:

- Energy-loss of hard-scattered partons while traversing the QGP
  - Modified fragmentation function  $D(z)$  by “rescaling” the variable  $z$



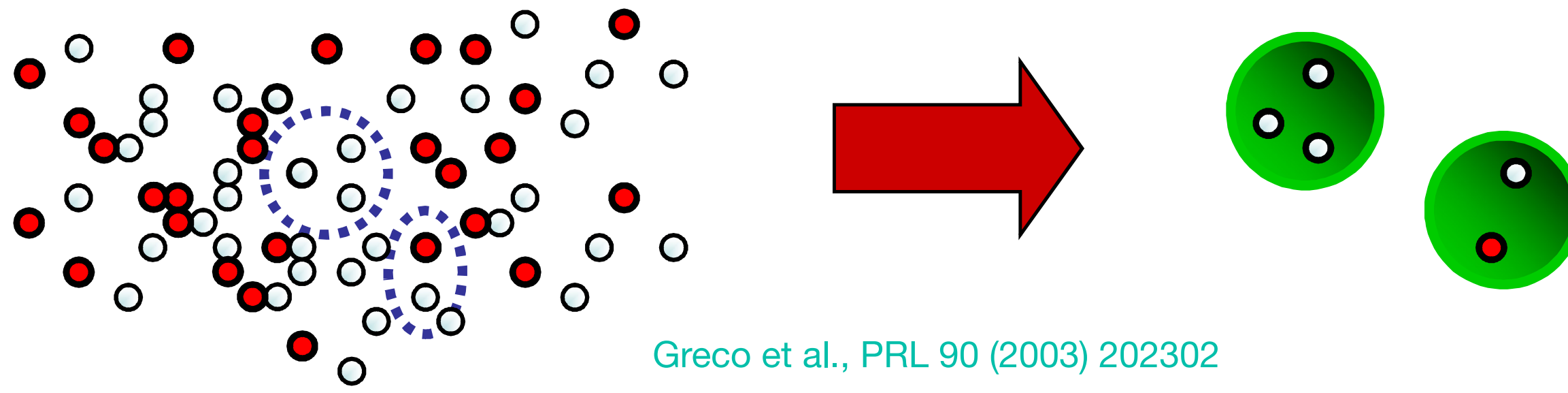
extracted from experiments such as  $e^+e^-$  collisions, ex. Peterson

$$D_{Q \rightarrow H}(z) \propto \frac{1}{z \left[ 1 - \frac{1}{z} - \frac{\epsilon}{1-z} \right]^2}$$

$$\epsilon = m_q^2 / m_Q^2$$

# Hadronization in medium; expectation

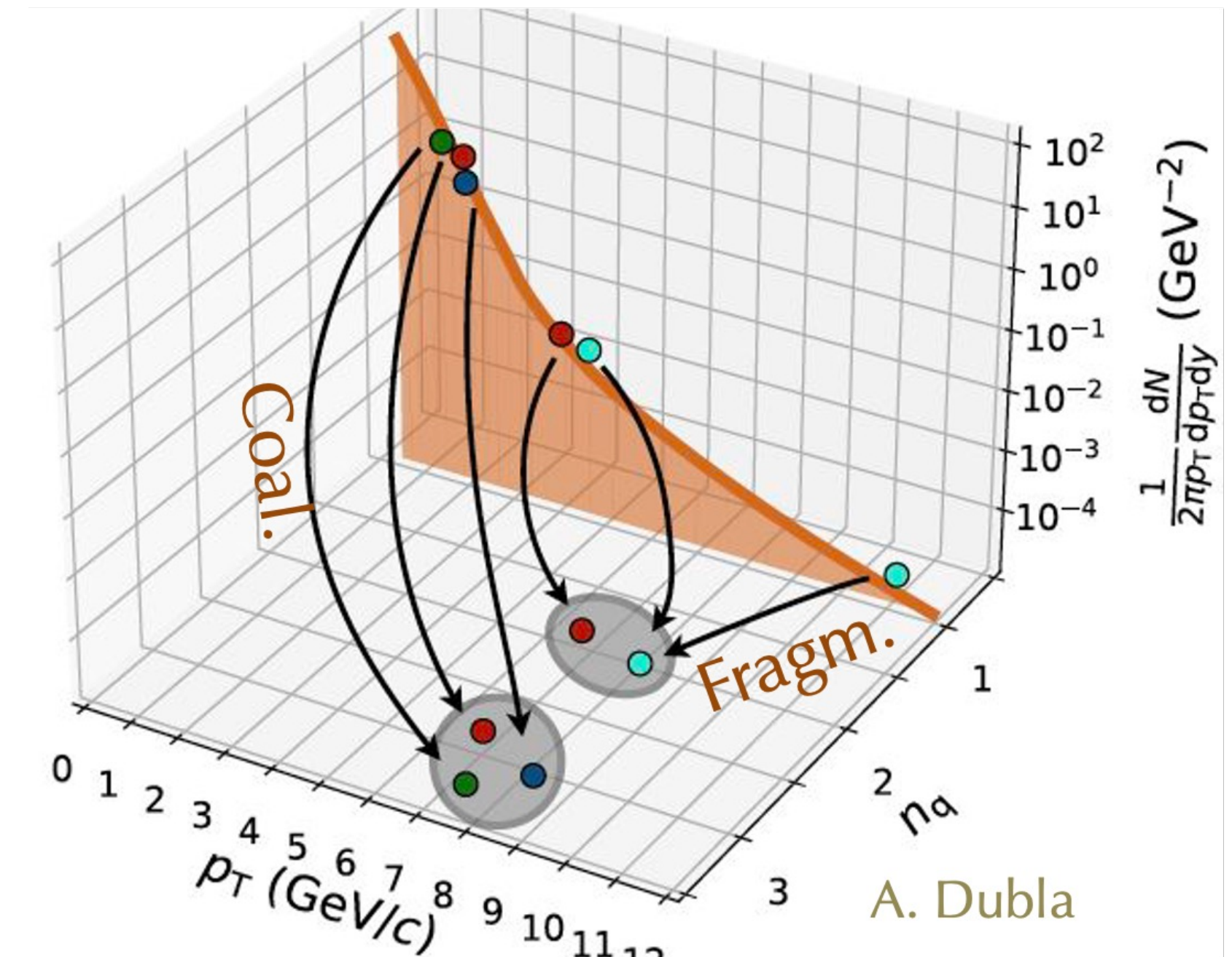
- Phase space at the hadronization is filled with partons
  - Single parton description may not be valid anymore
  - No need to create  $q\bar{q}$  pairs via splitting / string breaking
  - Partons that are “close” to each other in phase space (position and momentum) can simply recombine into hadrons



Greco et al., PRL 90 (2003) 202302  
 Fries et al., PRL 90 (2003) 202303  
 Hwa, Yang, PRC 67 (2003) 034902

- **Recombination vs. fragmentation:**

- Competing mechanisms
- Recombination naturally enhances baryon/meson ratios at intermediate  $p_T$

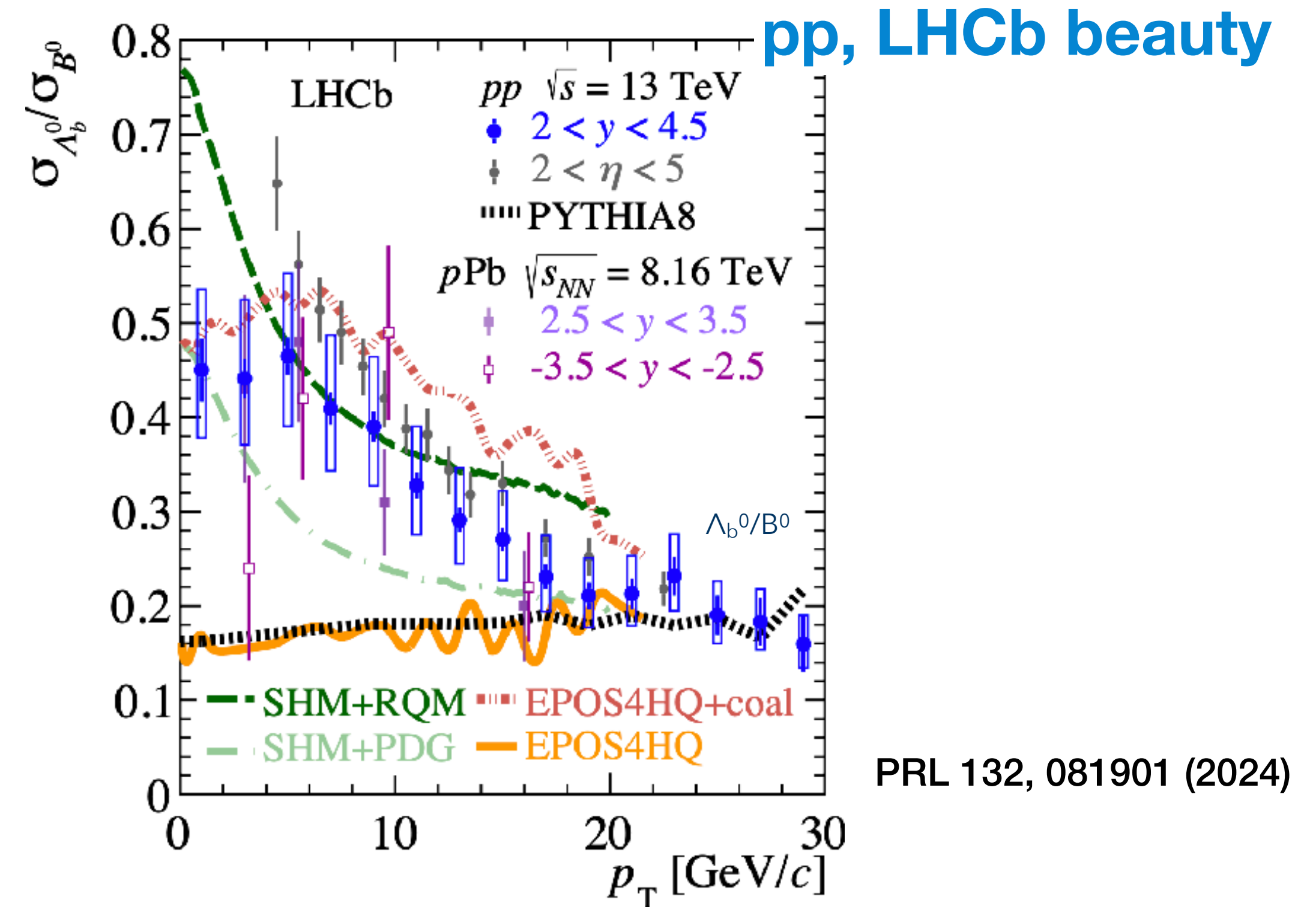
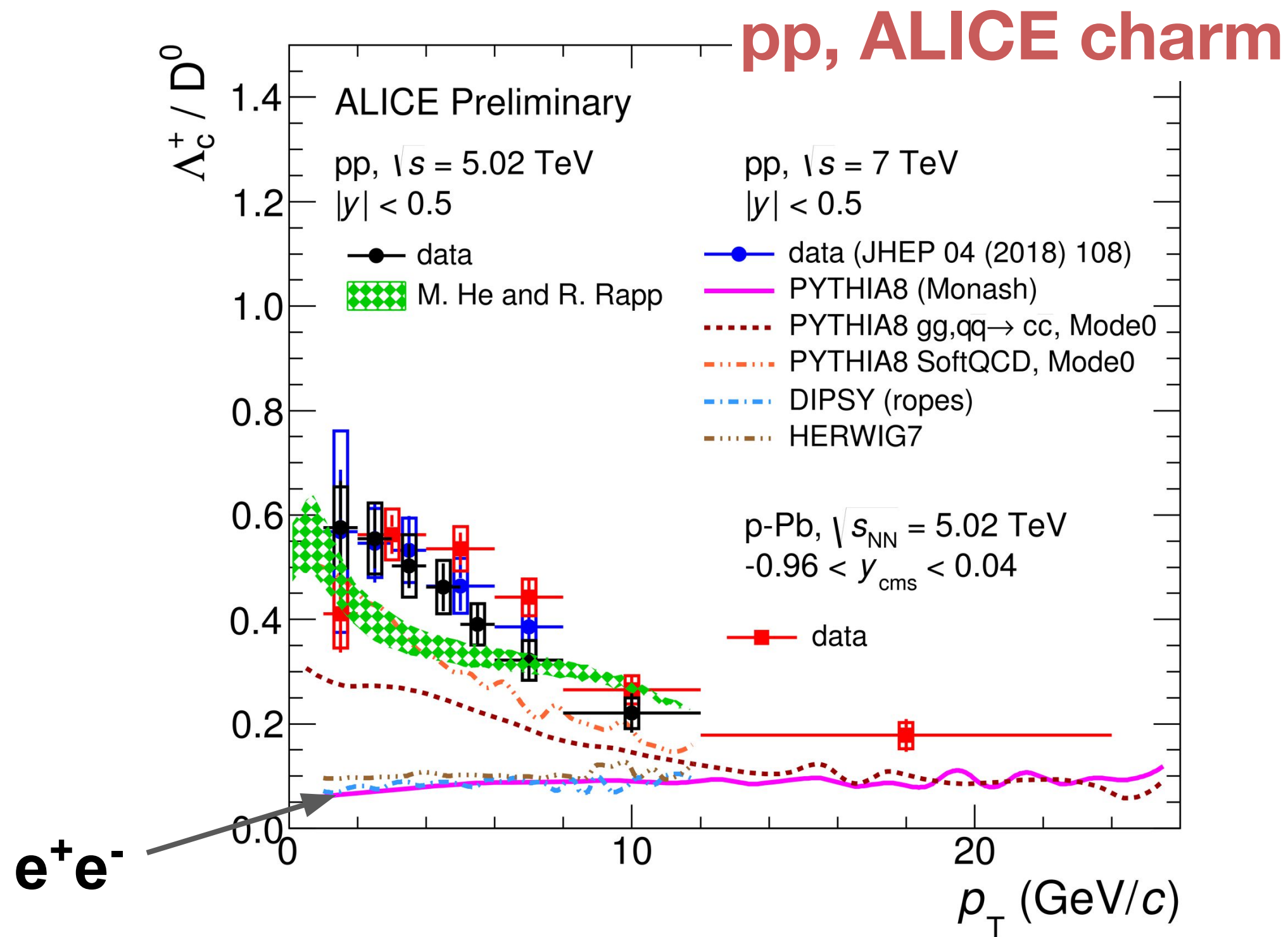


A. Dubla

# Hadronization in vacuum; observation

“Naive expectation: ratios of particle-species yields independent from collision system”

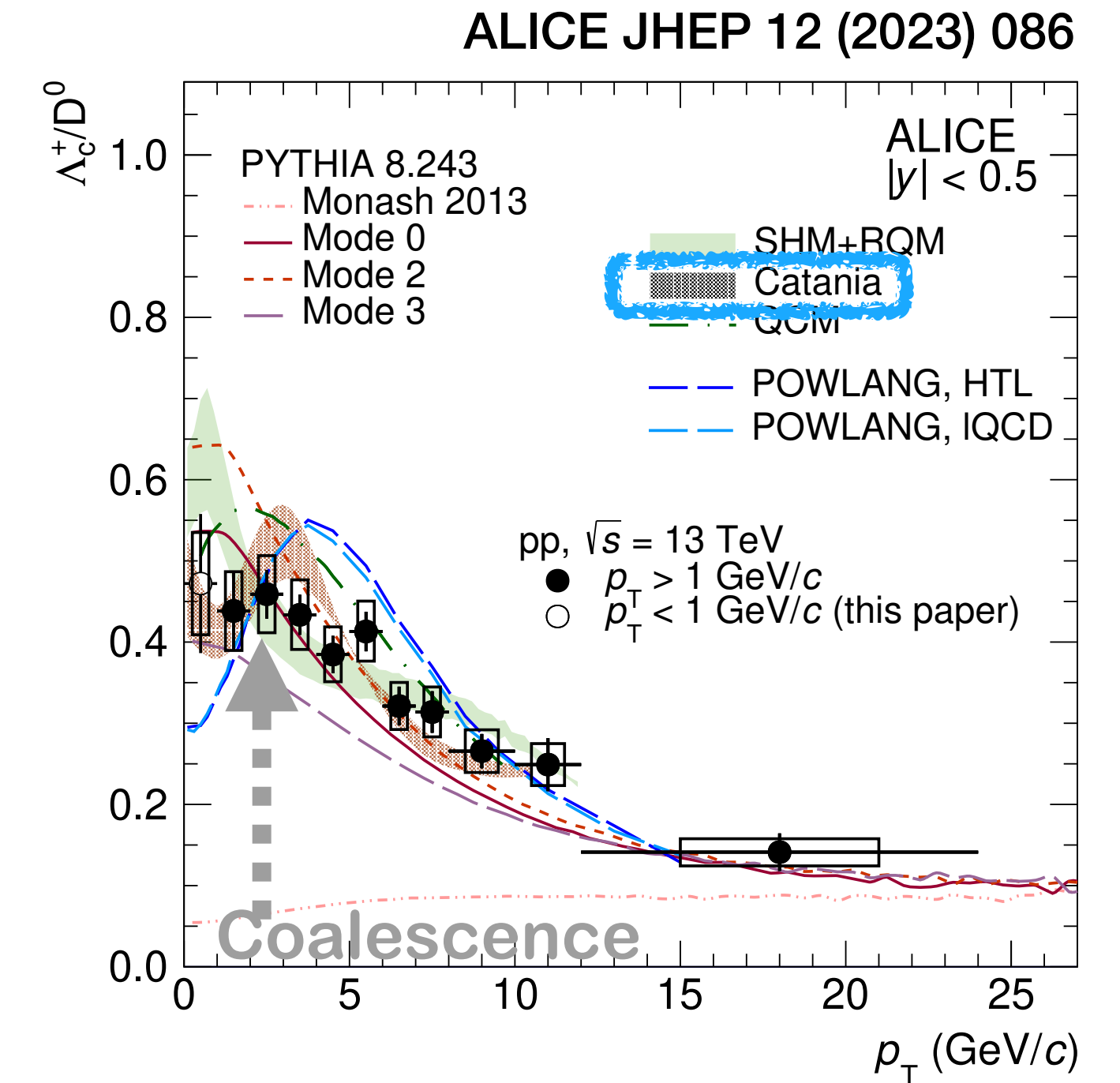
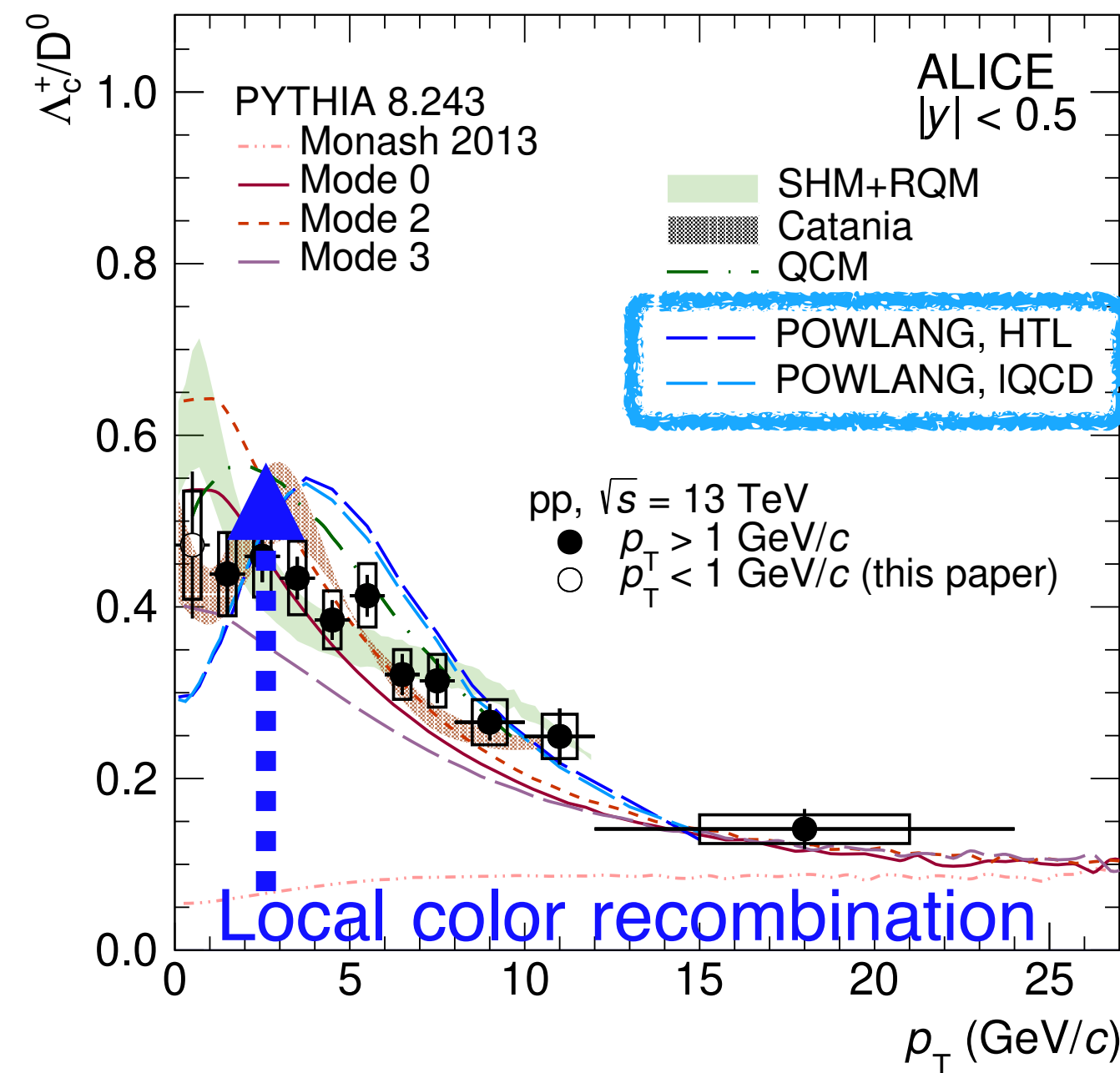
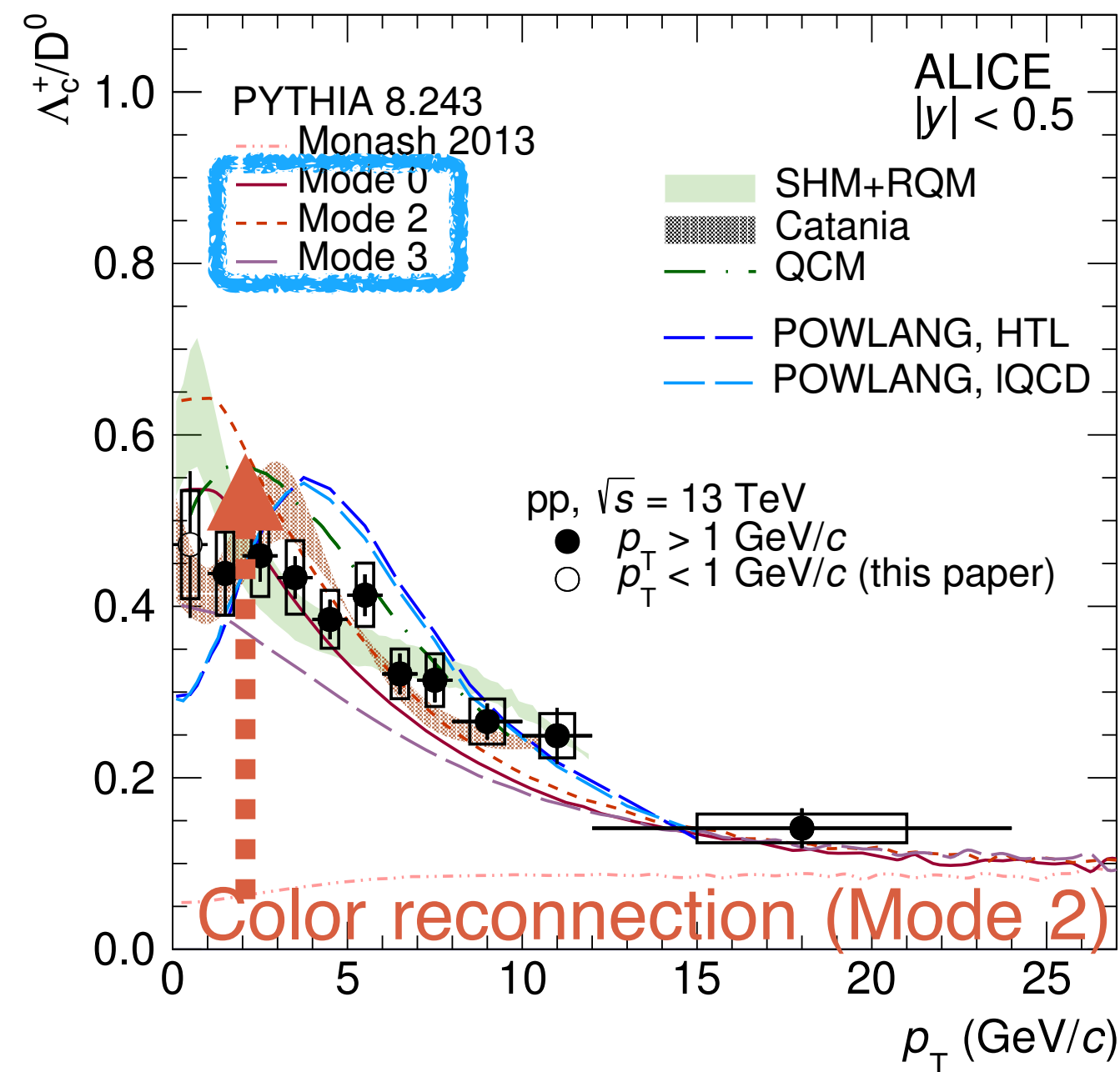
Surprises:  $\Lambda_c/D^0 \sim 0.5$  (at intermediate  $p_T$ ) not only in AA but even in pp  $\rightarrow$  strong enhancement wrt  $e^+e^-$



$\Lambda_b/B^0 \rightarrow$  Similar trend in charm and beauty sectors!

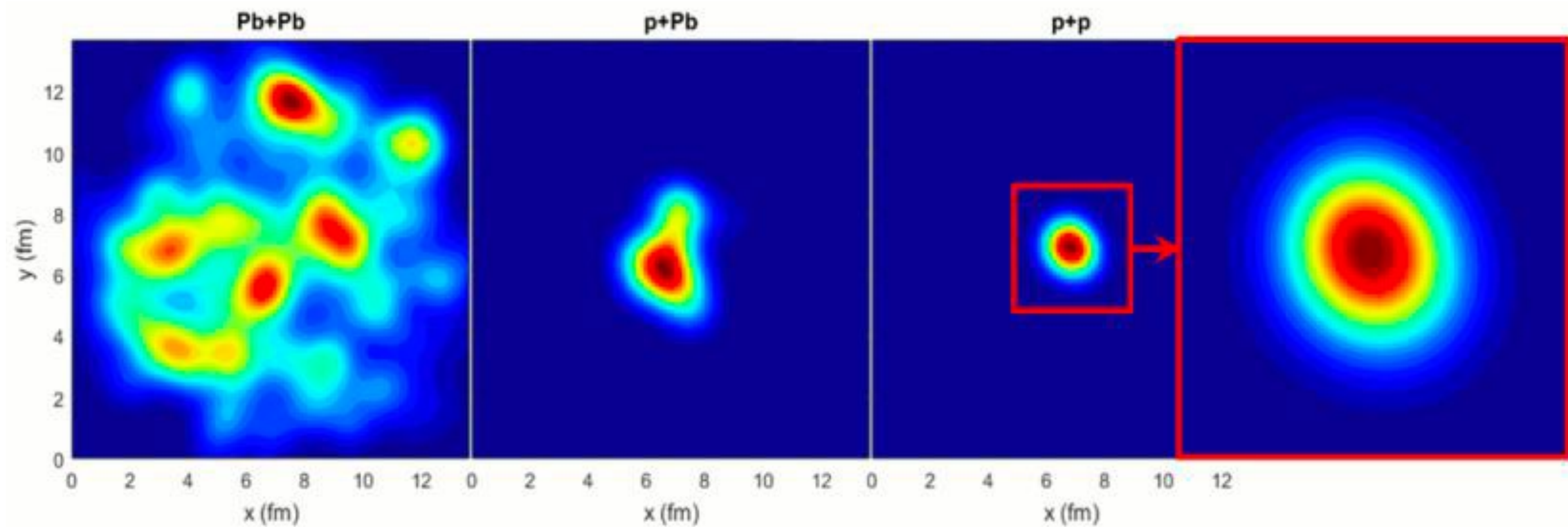
# OK, how do we explain? Heavy flavour baryon enhancement impact...

- Heavy-flavour hadronization stimulated the model developments
  - PYTHIA with Color Reconnection (CR) beyond Leading Color (LC) in pp
  - Catania: Coalescence+Fragmentation approach applied to pp
  - Local color recombination: POWLANG in AA and in pp
  - Inclusion of heavy-flavour Coalescence+Fragmentation in EPOS (pp & AA)



- Different hadronization mechanisms proposed!

# As an example, in Catania, coalescence + fragmentation in pp



R. D. Weller, P. Romatschke, PLB 774 (2017) 351-356

Vincenzo Greco's expression in his SQM talk!

**Daring** to assume a small fireball according to **viscous hydro** applied to pp as in AA, but **size, time, flow given by hydro for pp**

**p+p @ 5 TeV**

- $t_{pp} = 1.7 \text{ fm}/c$
  - $\beta_0 = 0.4$
  - $R = 2.5 \text{ fm}$
  - $V \sim 30 \text{ fm}^3$
- +  $f_c(p)$  from **FONNL distribution**

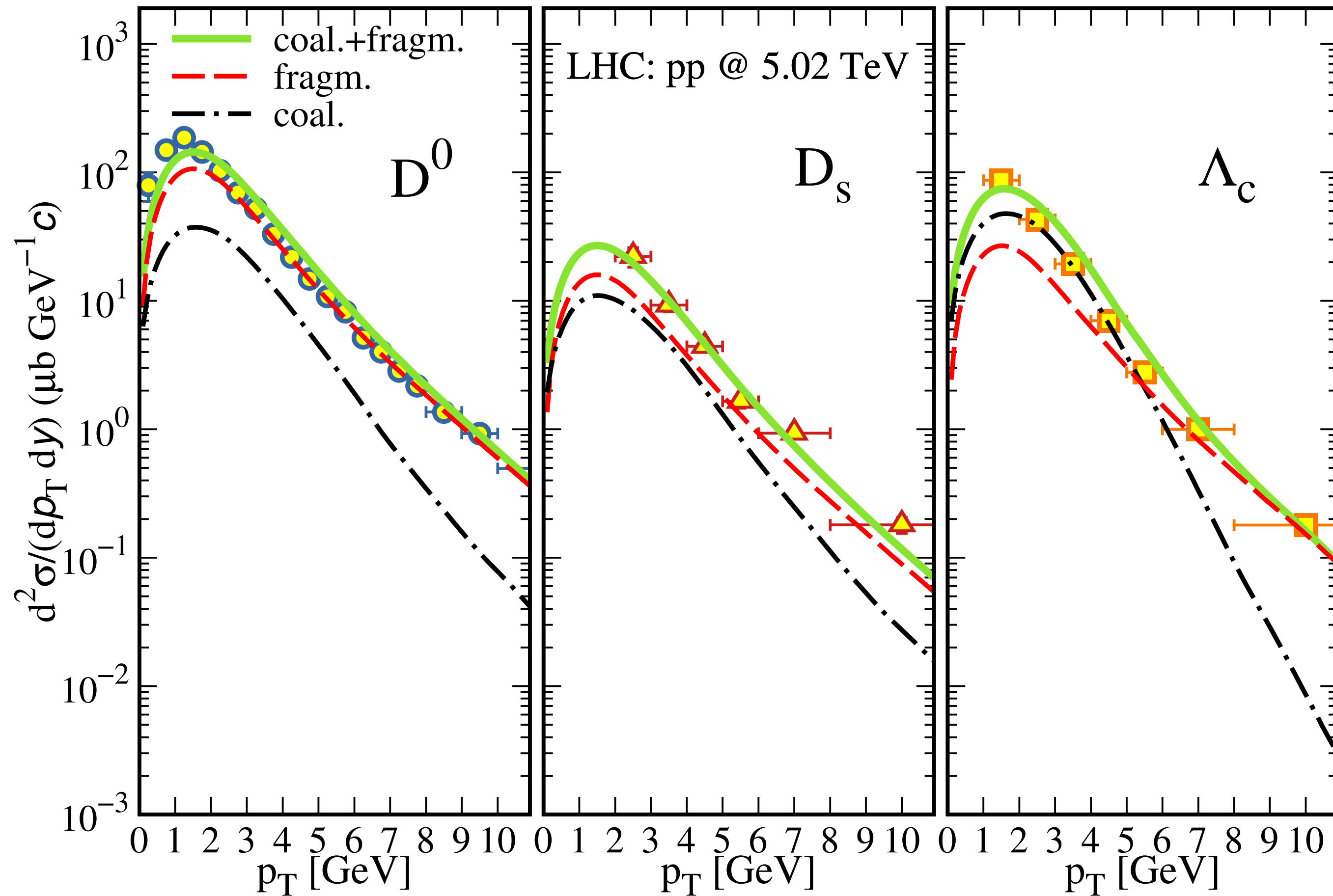
$$f_q(p) \sim \frac{dN_{q,\bar{q}}}{d^2p_T} \sim \exp\left(-\frac{\gamma_T(m_T - p_T \cdot \beta_T \mp \mu_q)}{T}\right)$$

+ same Wigner function widths  $\sigma_{r,i}$  of hadrons in AA

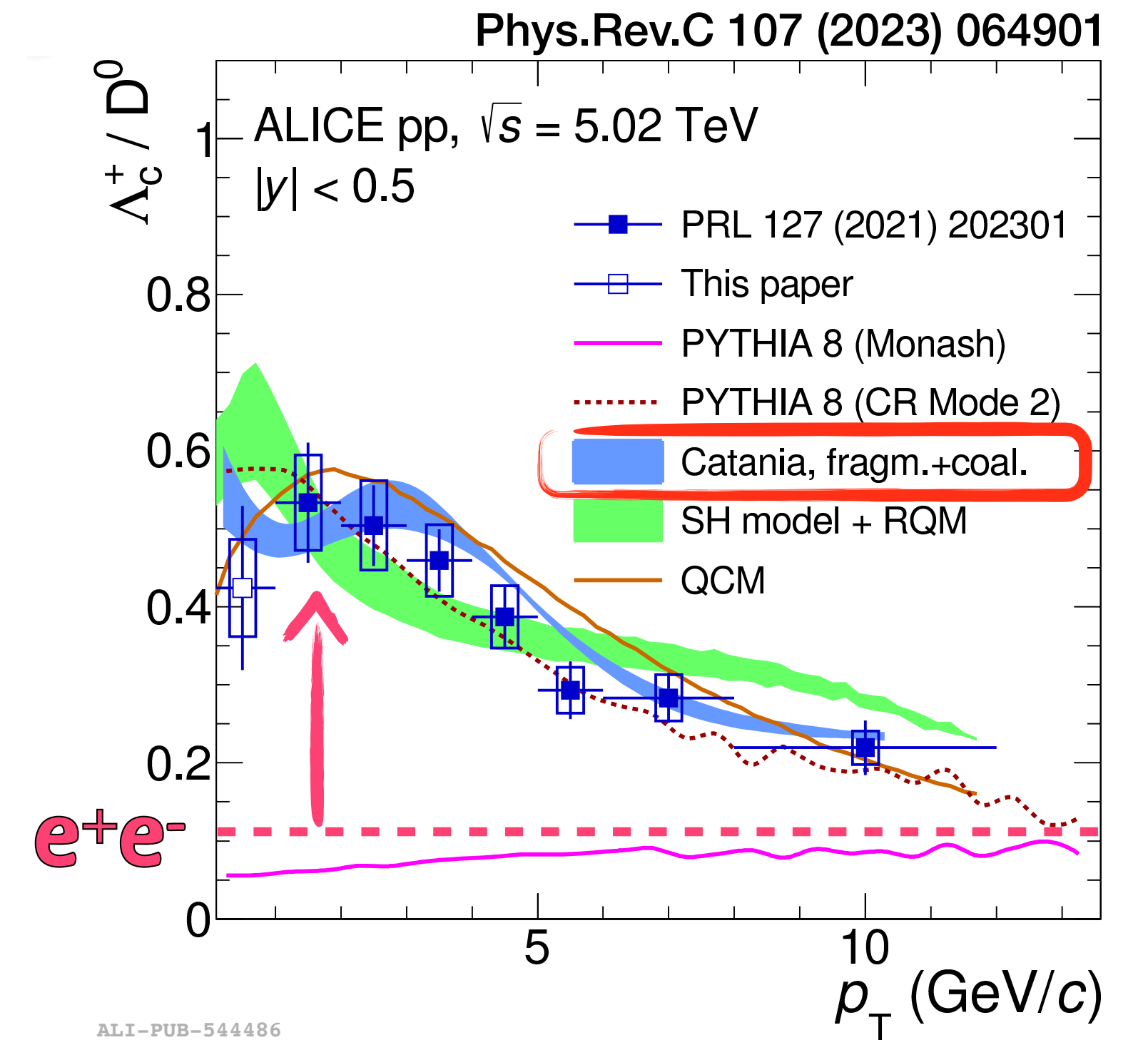
$$f_H(x_i, p_i) = \prod_{i=1}^{N_q-1} 8 \exp\left(-\frac{x_{r,i}^2}{\sigma_{r,i}^2} - p_{r,i}^2 \sigma_{r,i}^2\right)$$

# Coalescence in pp vs $p_T$ in Catania

Phys. Lett. B 821 (2021) 136622

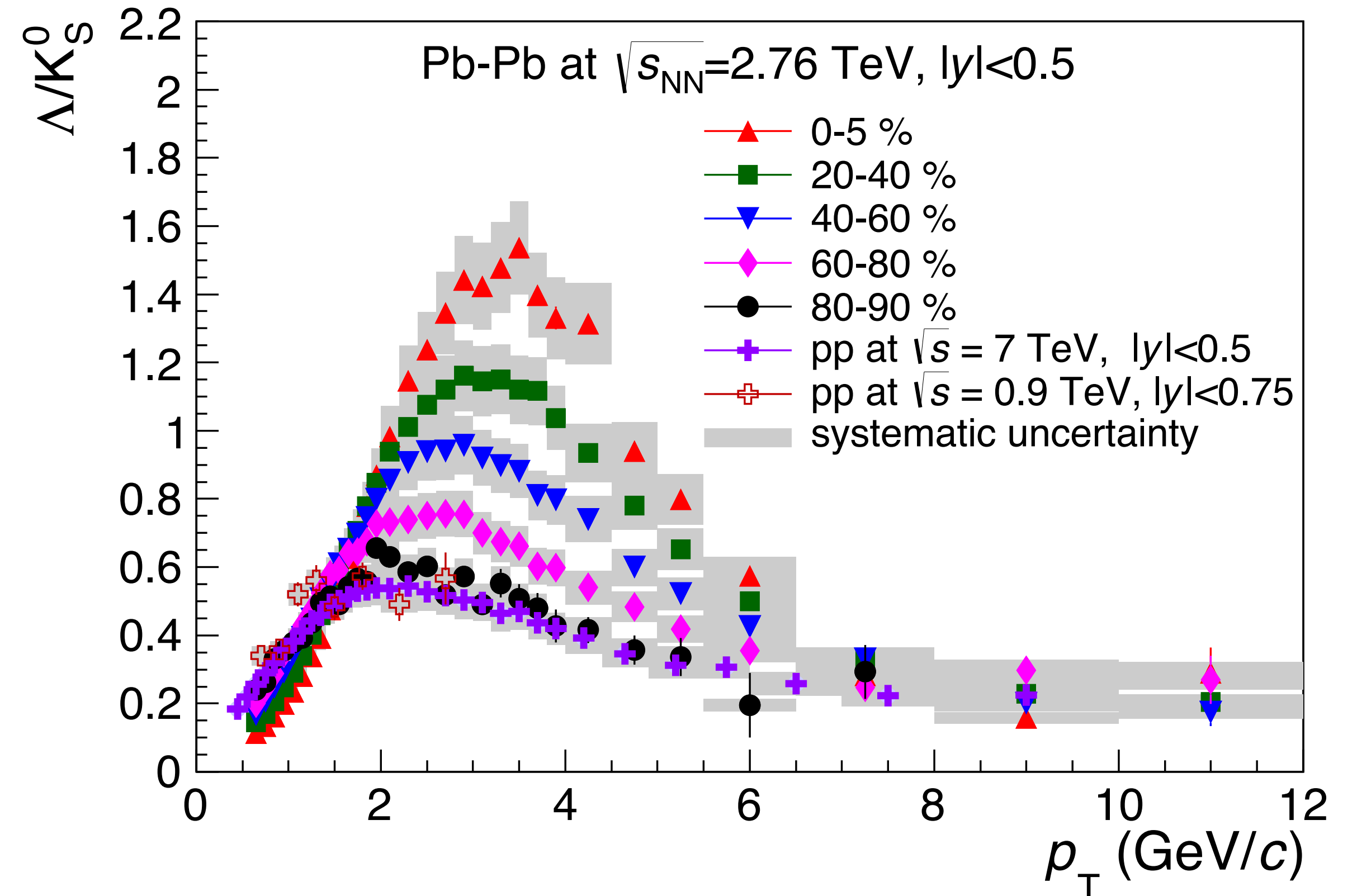
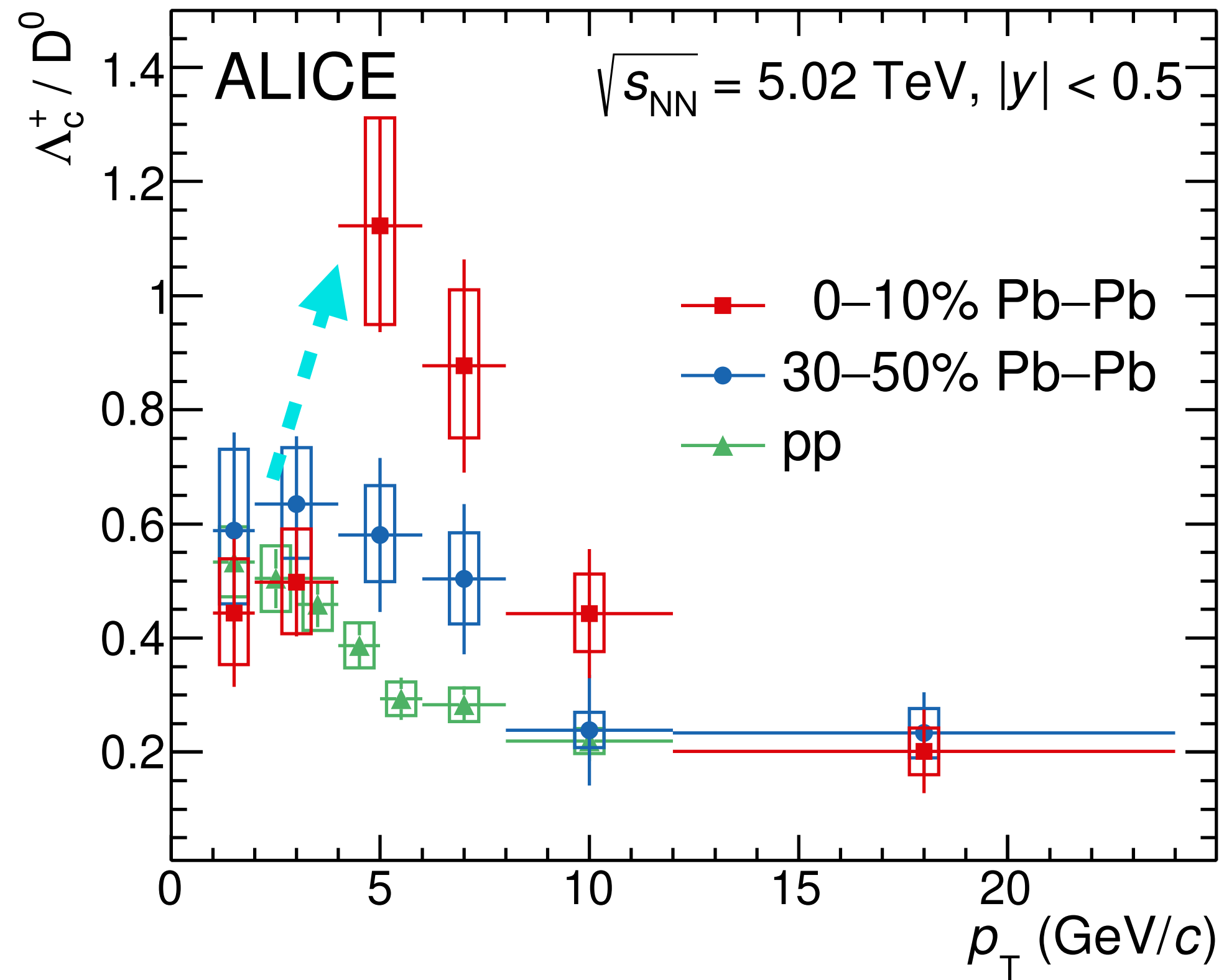


- All the coalescence does not affect significantly  $D^0$ , but is dominant for baryons  $\Lambda_c$  and  $\Xi_c$



# How about in Pb–Pb?

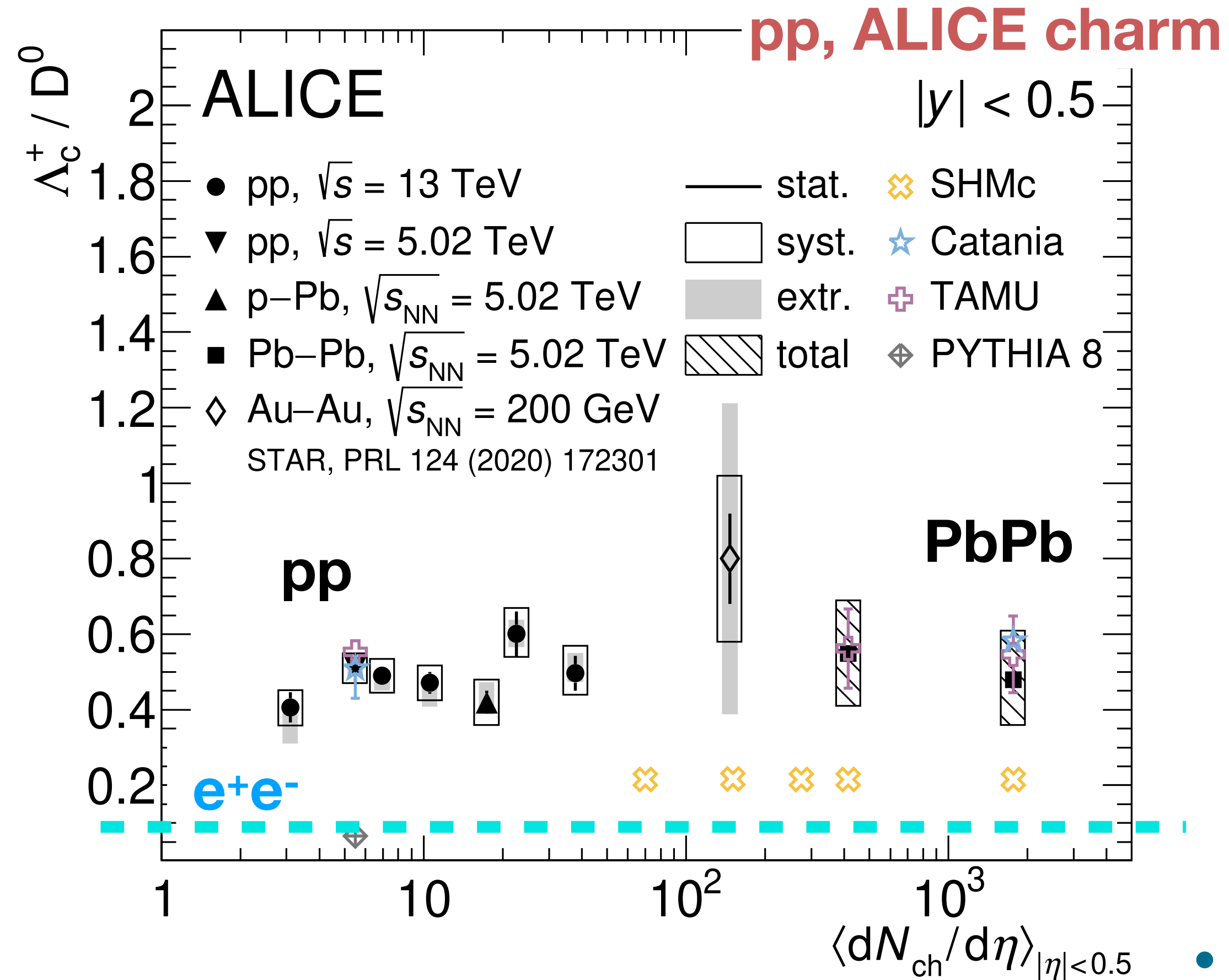
Phys. Lett. B 839 (2023) 137796



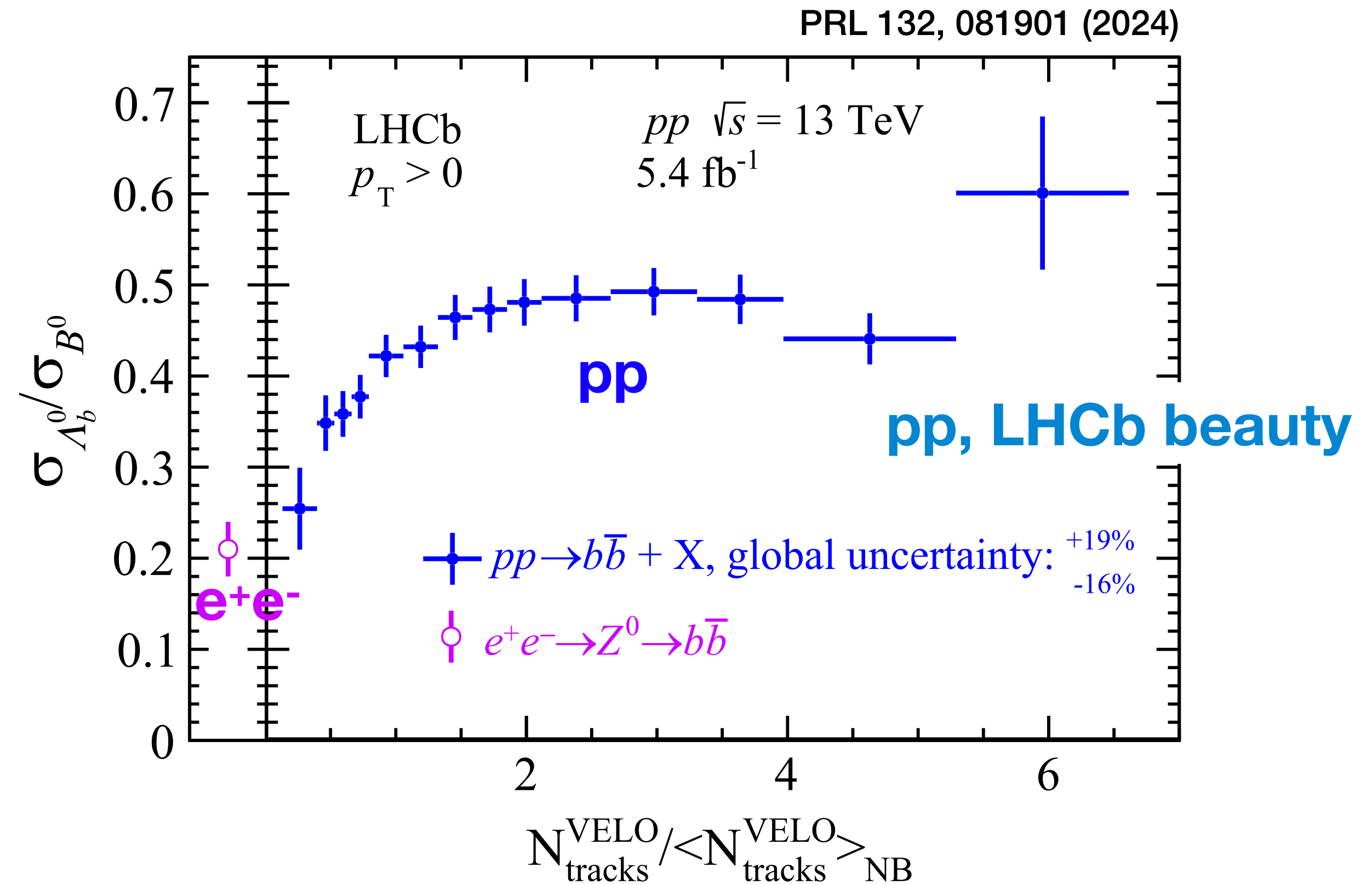
- Ratio increases from pp to mid-central and central Pb-Pb at intermediate  $p_T$
- Trend qualitatively similar to what is observed for  $\Lambda / K_s^0$  ratios

# $p_T$ -Integrated yield ratio in pp collisions

## $p_T$ -integrated $\Lambda_c^+/D^0$ ratios



## $p_T$ -integrated $\Lambda_b^+/B^0$ ratios

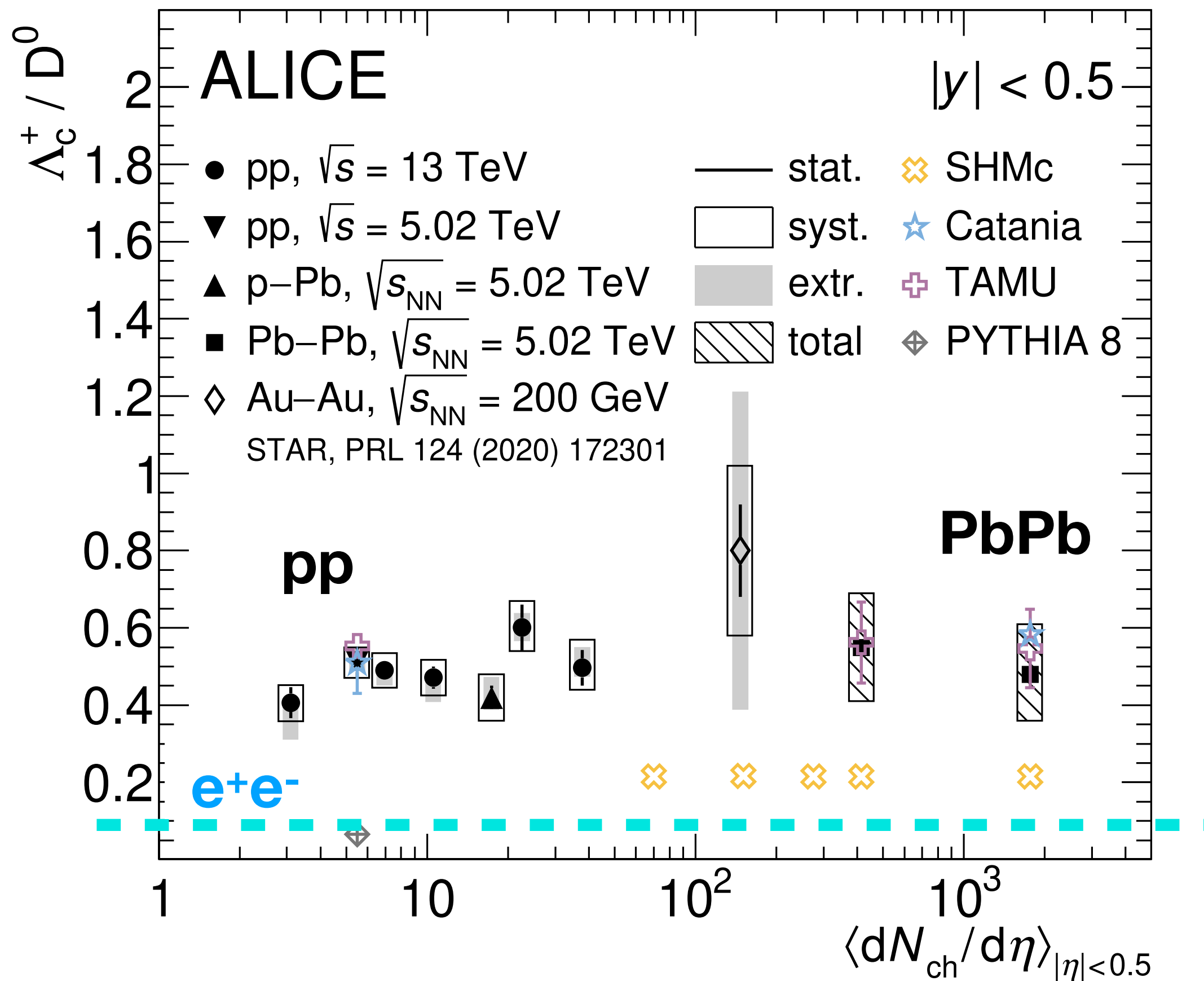


- $p_T$ -integrated yield ratio is saturated in all hadronic collision systems
- Then, enhancement is due to  $p_T$  redistributions?
- Similar to b sector

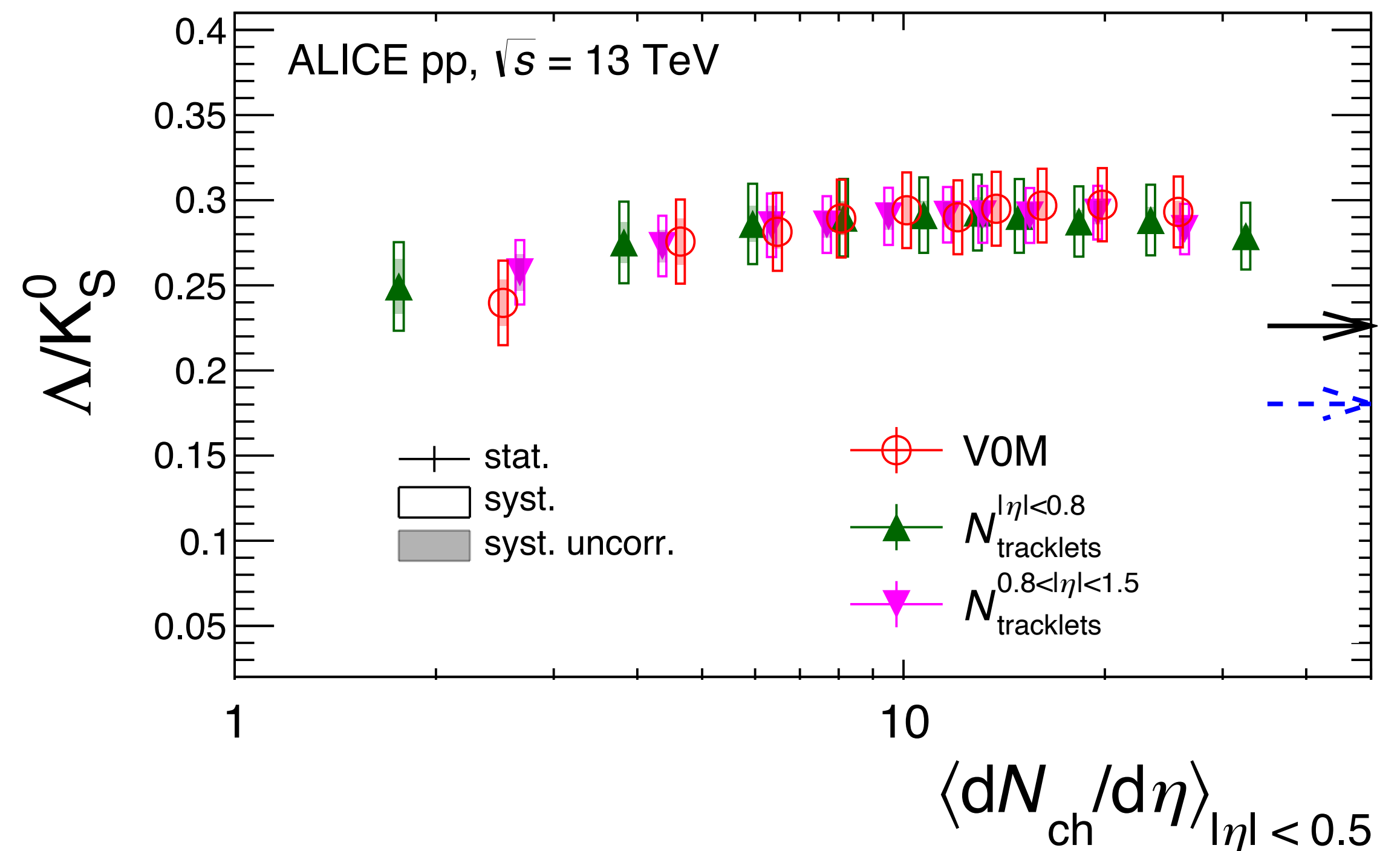


# Where does the $p_T$ differential enhancement come from?

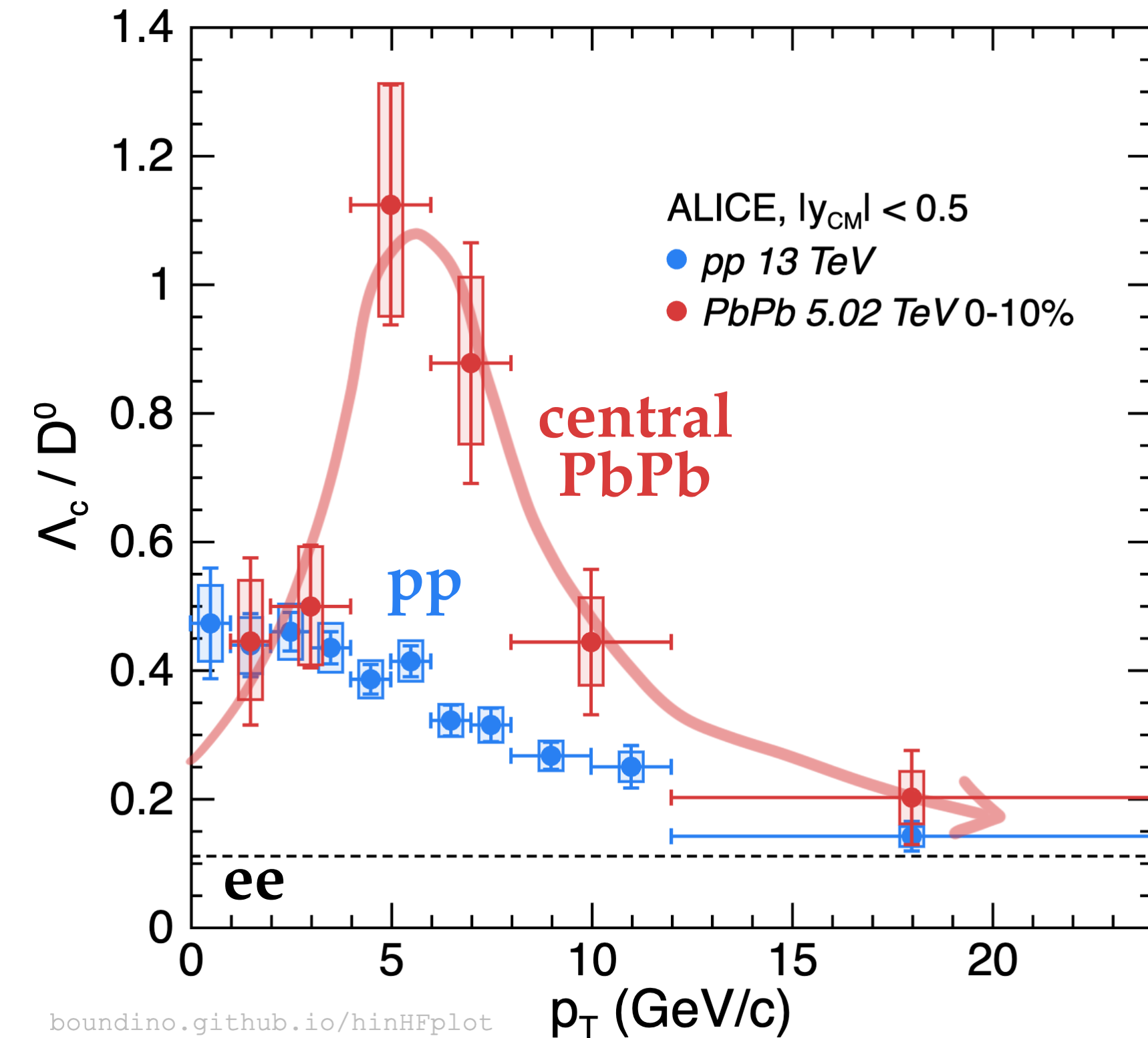
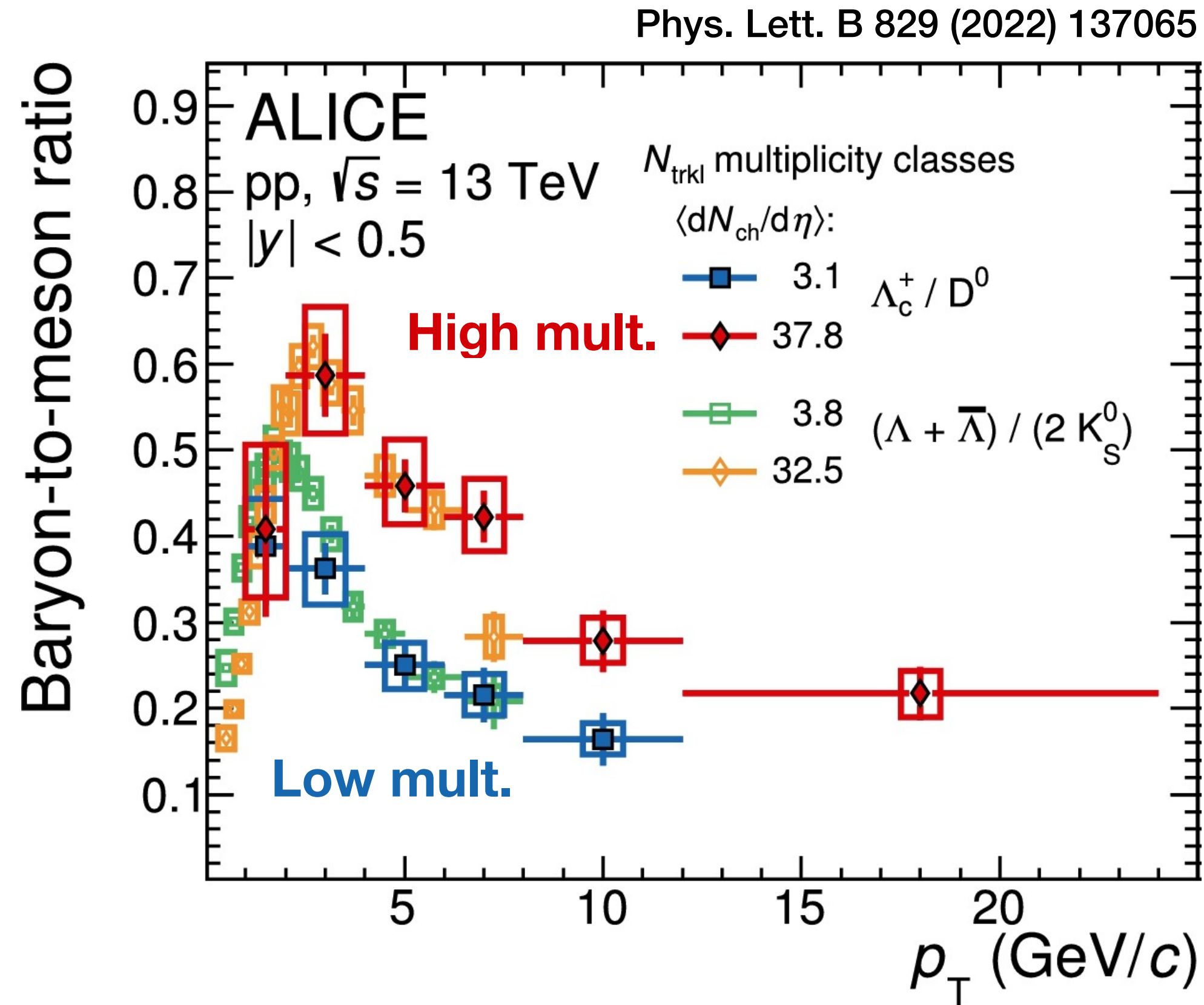
## $p_T$ -integrated $\Lambda_c^+/D^0$ ratios



- Due to **different  $p_T$  redistribution** for baryons and mesons rather than multiplicity dependence in hadronization process itself?
- Modified mechanism of hadronization **in all hadronic collision systems** with respect to charm fragmentation tuned on e<sup>+</sup>e<sup>-</sup> and e-p measurements?

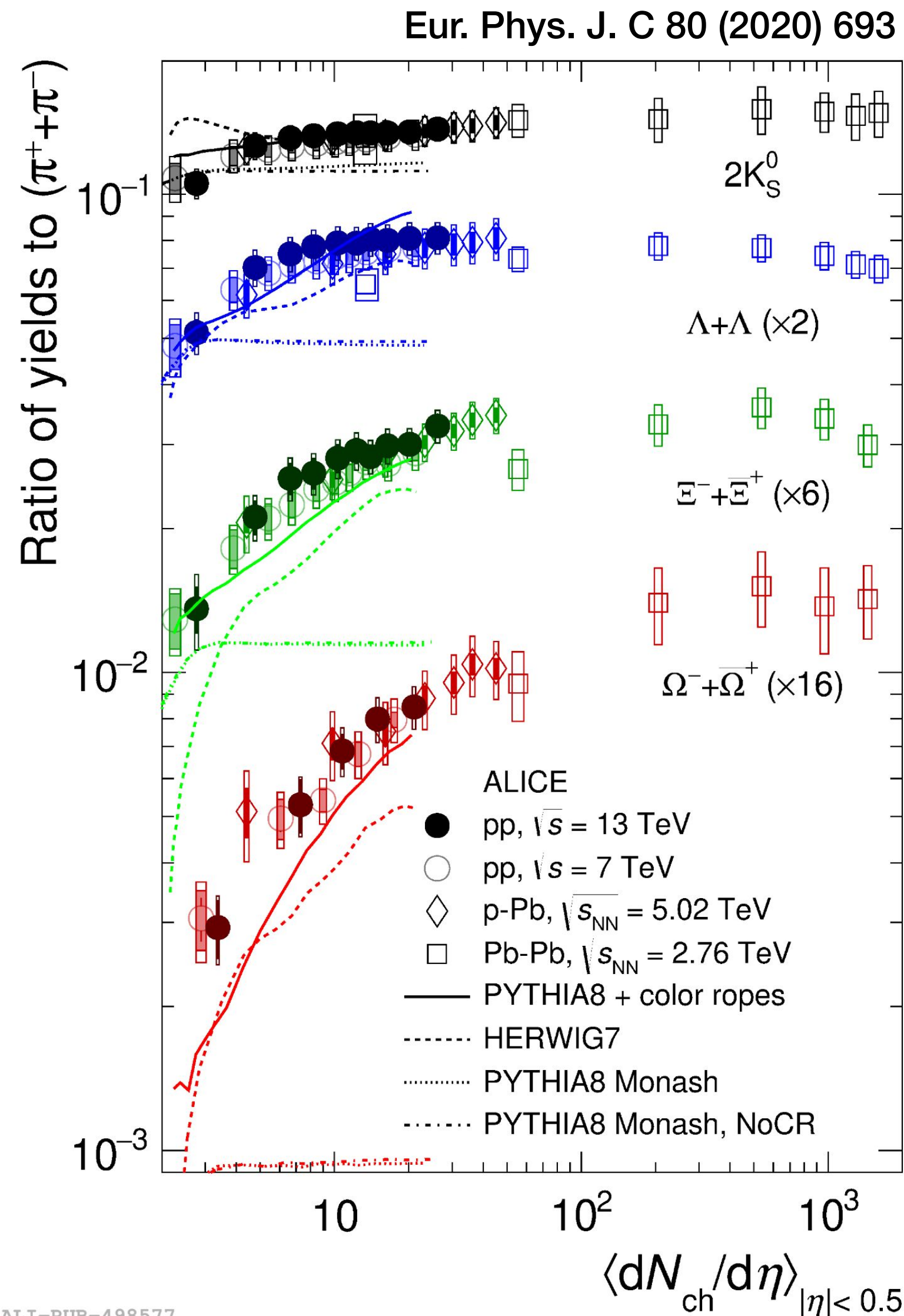


# $p_T$ redistribution



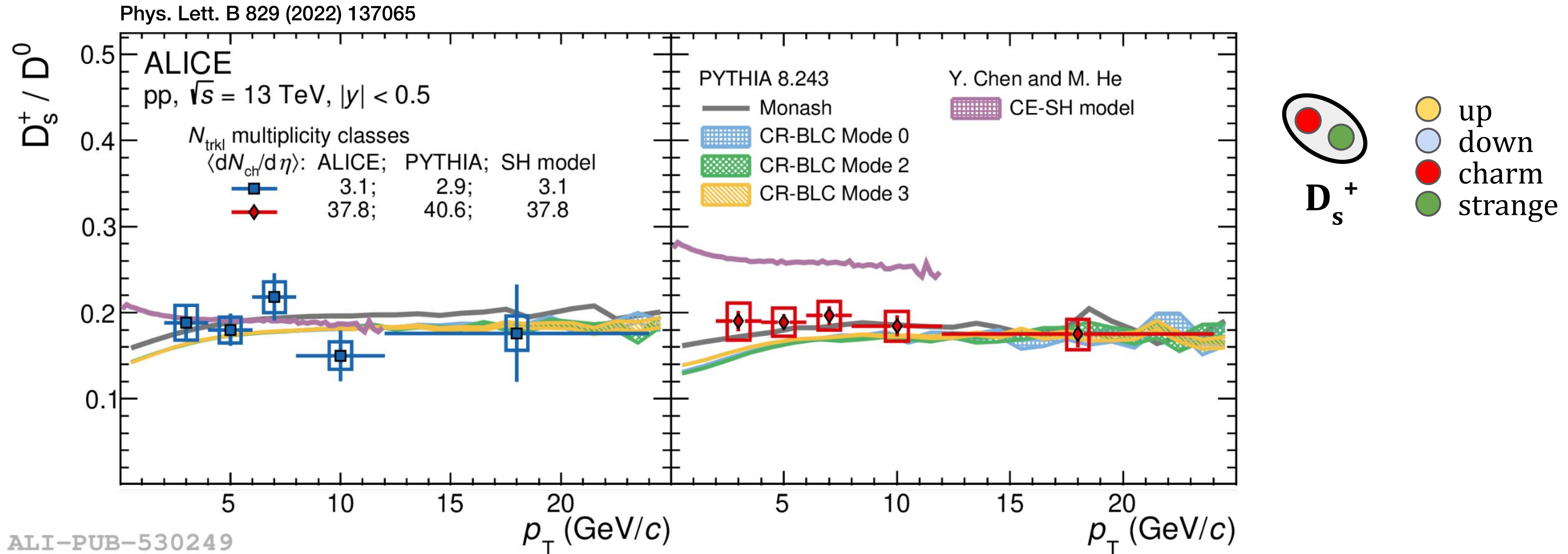
- Charm baryons/meson like for strangeness!
- Common mechanism for light- and charm-baryon formation in hadronic collisions? (unlikely) or coincidence in a redistribution?
- Shape changes dramatically in central PbPb → Strongest radial flow?

# Role of strangeness in heavy-quark hadronization



- **Strangeness enhancement:** yield-ratio between (multi)strange hadrons and pion larger **in heavy-ion collisions** than minimum-bias pp collisions
- Smooth increase vs. event **multiplicity**, without a clear **collision-system dependence**
- What do we learn from **strange heavy hadron** ( $D^0, \Lambda_c^+, \Xi_c^0, \dots$ ) production about heavy-quark **hadronization**
  - ➔ **evolve vs. event multiplicity?**
  - ➔ **sensitive to QGP-induced effects** (e.g. strangeness enhancement, coalescence,  $E$ -loss, flow, ...)?

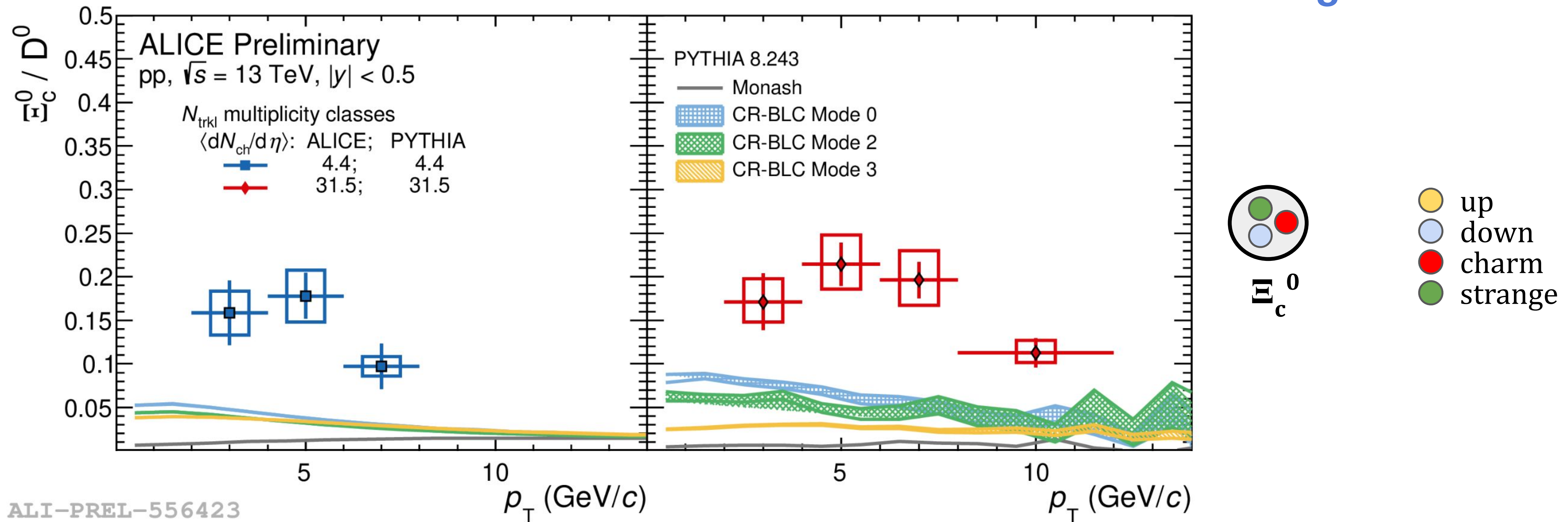
# $D_s^+$ : strange charm meson



- $D_s^+/D^0$  ratio are independent of  $p_T$
- No strong multiplicity dependency
- Comparable with measurement at  $e^+e^-$  and  $e-p$  collisions

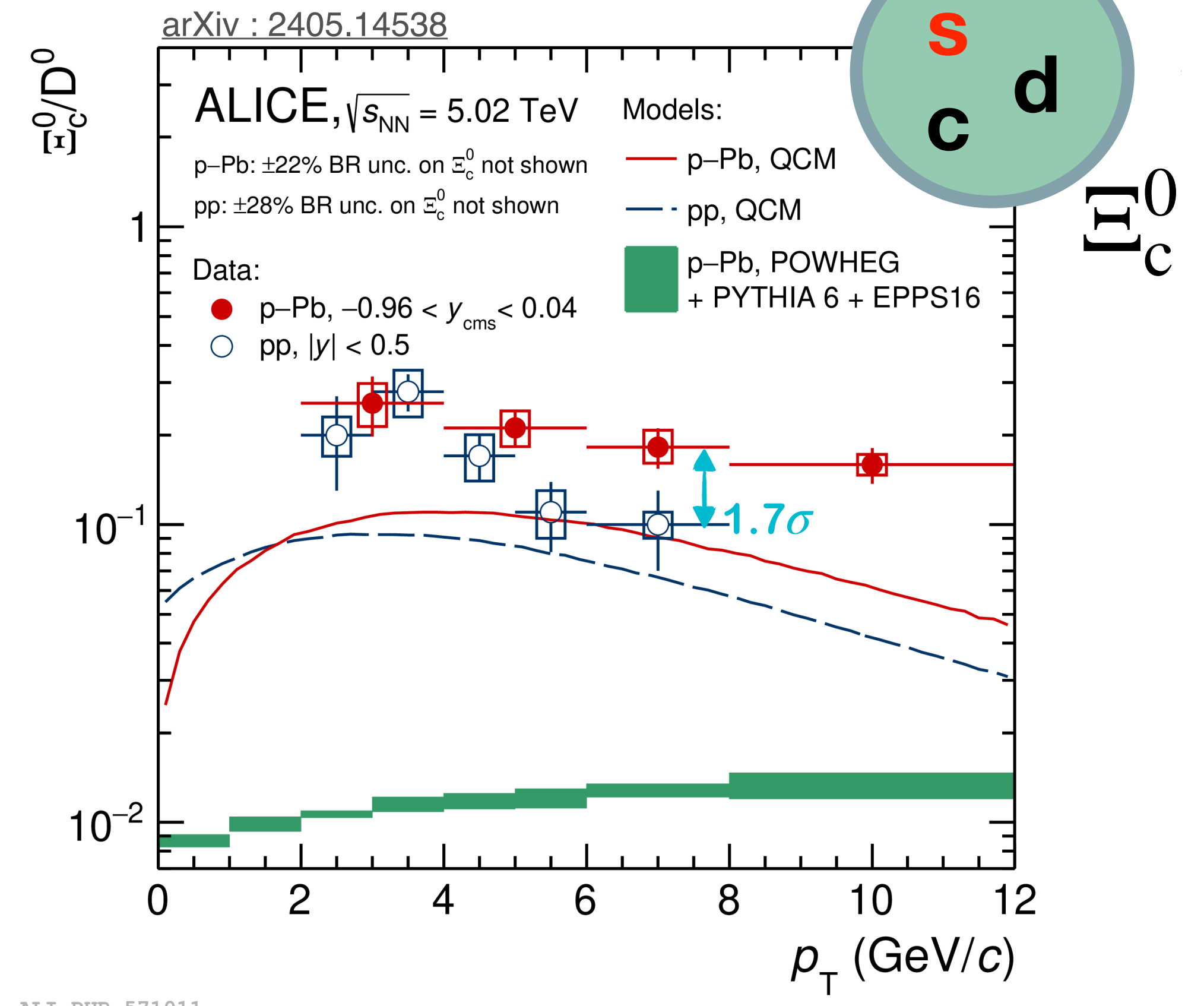
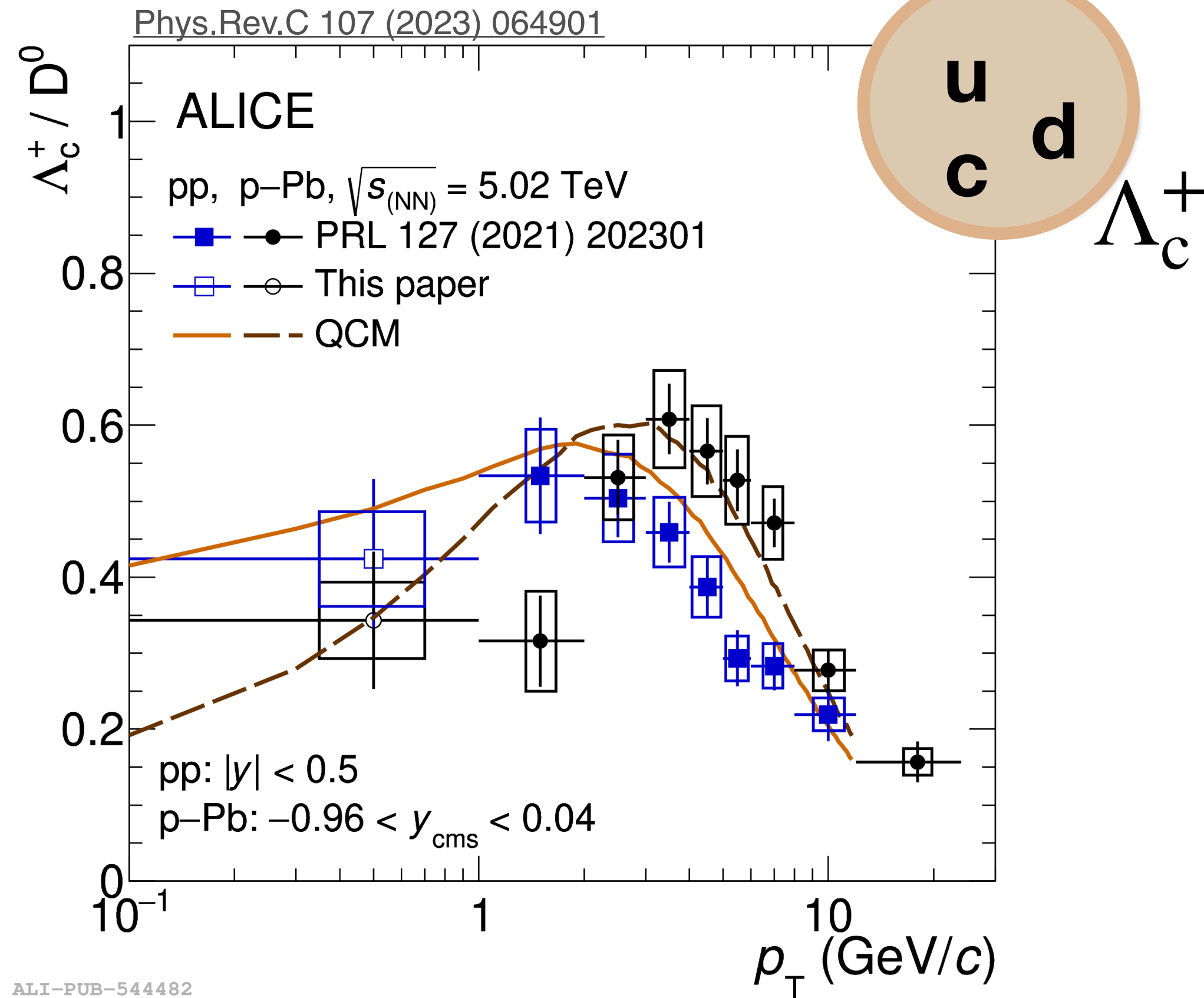
# What about strangeness charm baryon?

More challenges!



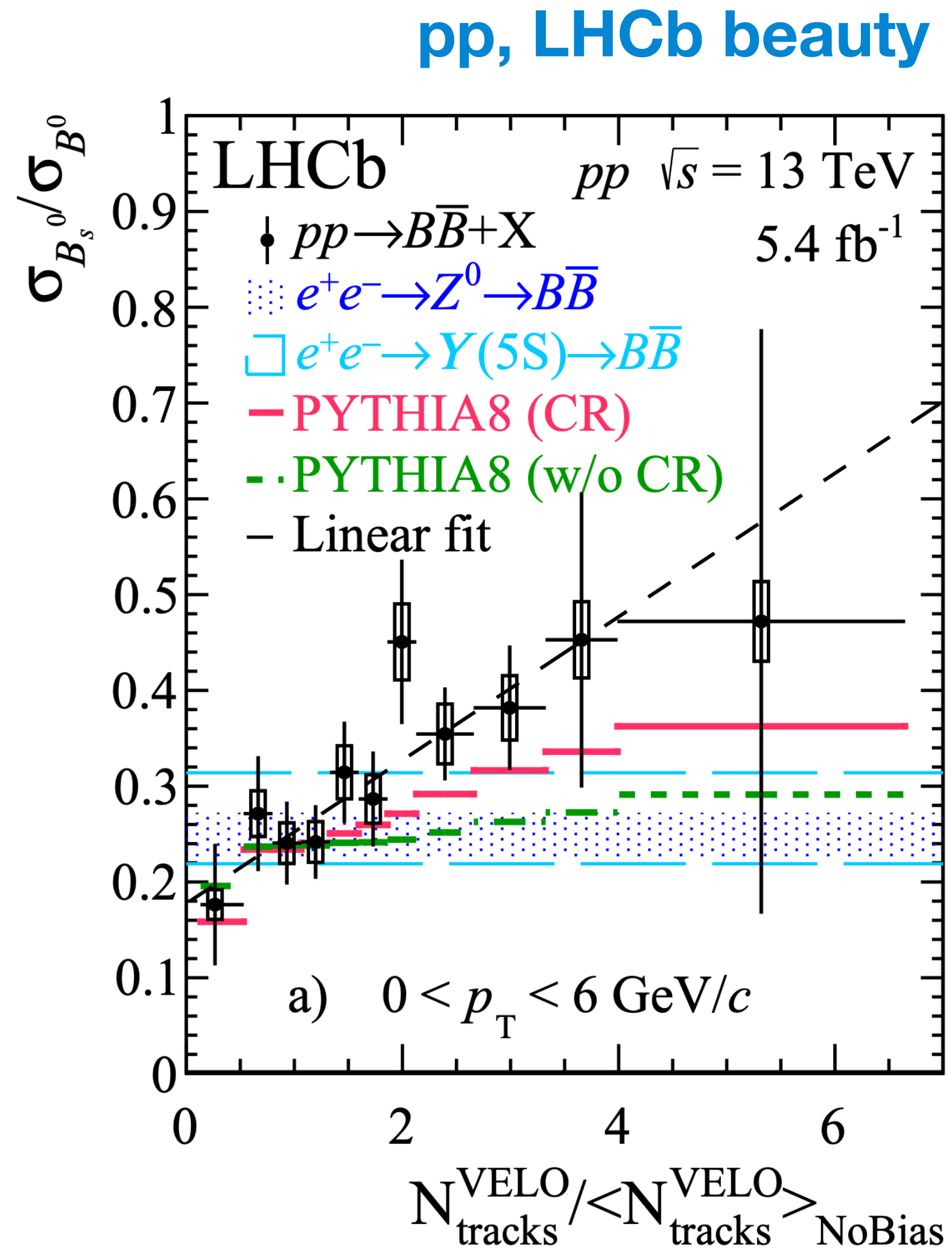
- **Strong  $p_T$  dependence**
- **Enhancement** compared to the measurement in  $e^+e^-$  and  $e^-p$  collisions

# Challenging models

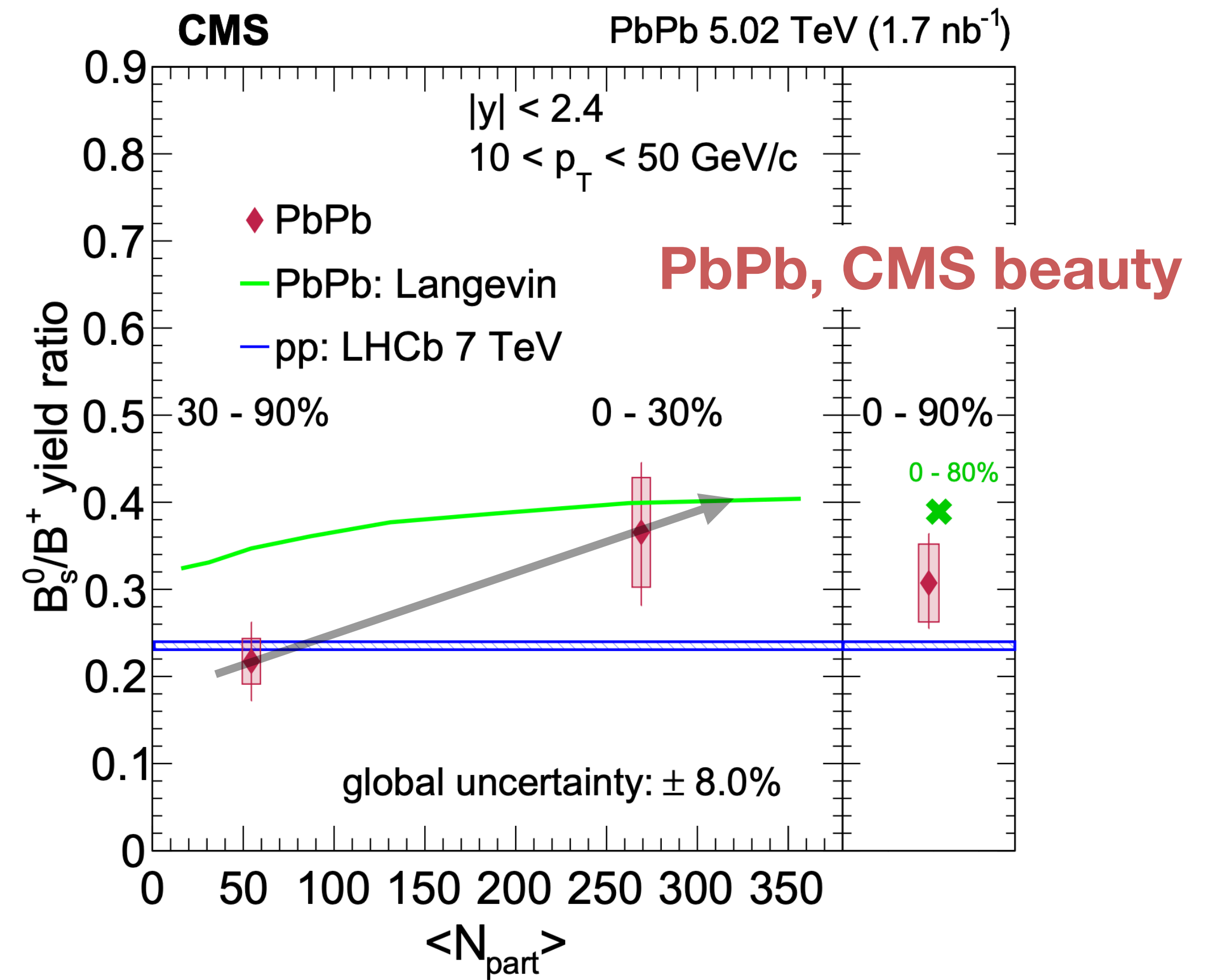


- Shift of distribution peak towards higher  $p_T$  could be attributed to radial flow
- QCM describes the magnitude of the ratio for  $\Lambda_c^+/D^0$ , but underestimate for  $\Xi_c^0$

# Is beauty different?



$B_s^0/B^+$

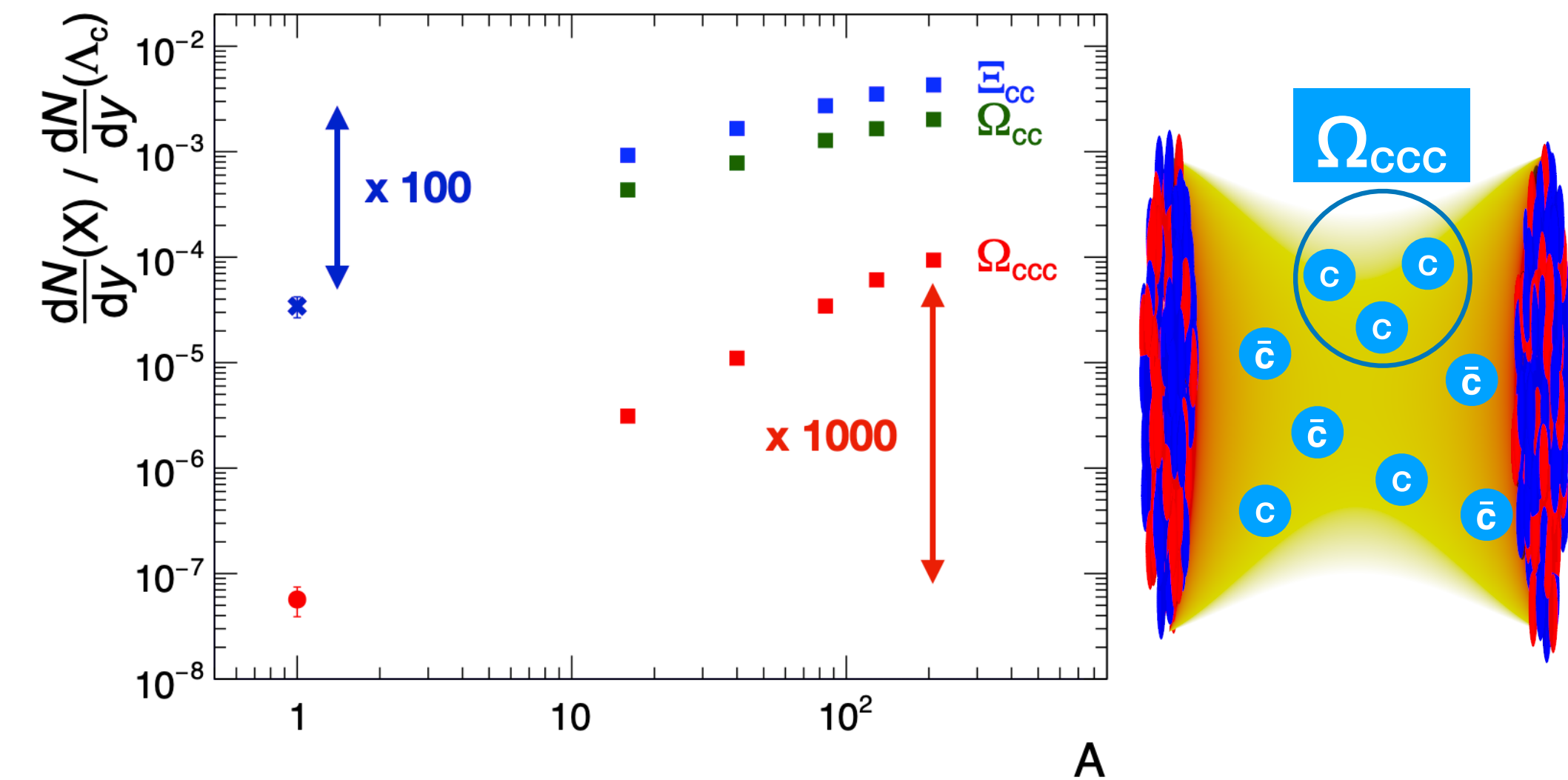


- Hint of different behavior of beauty
- Need precision measurement

# What is obvious?, what is vague, what is unknown, ...

- Enhancement due to **different  $p_T$  redistribution** for baryons and mesons rather than multiplicity dependence in hadronization process itself?
- **At least in the market, coalescence** → a common framework for heavy-flavor hadronization from pp to AA?
- Other approaches such as PYTHIA-CR, POWLANG-LCN, ... point also to
  - In medium local recombination
  - Large evolution from  $e^+e^-$  to pp while reshuffling in  $p_T$  from pp to AA

- **Need more differential observables and precision measurement**
  - Rapidity evolution
  - Extend to bottom
  - Effect on the other observables (ex.  $v_2$ )
  - ....



→ Very large **enhancement** predicted by Statistical hadronization model in Pb-Pb

⇒ **Require new detector ALICE 3!**







Thank you for your attention!

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# Extra Slides

# Heavy-flavour hadronization modelings

## PYTHIA 8

Hadronization via **fragmentation**, color reconnection between partons from different multiparton interactions



### Monash tune

(tuned to  $e^+e^-$  measurements)

[Eur.Phys.J. C 74 \(2014\) 3024](#)

### Mode 2

the **junction** topology leads to an increase of baryon production

[JHEP 08 \(2015\) 003](#)

## SHM + RQM

[Phys.Lett.B 795 \(2019\) 117-121](#)

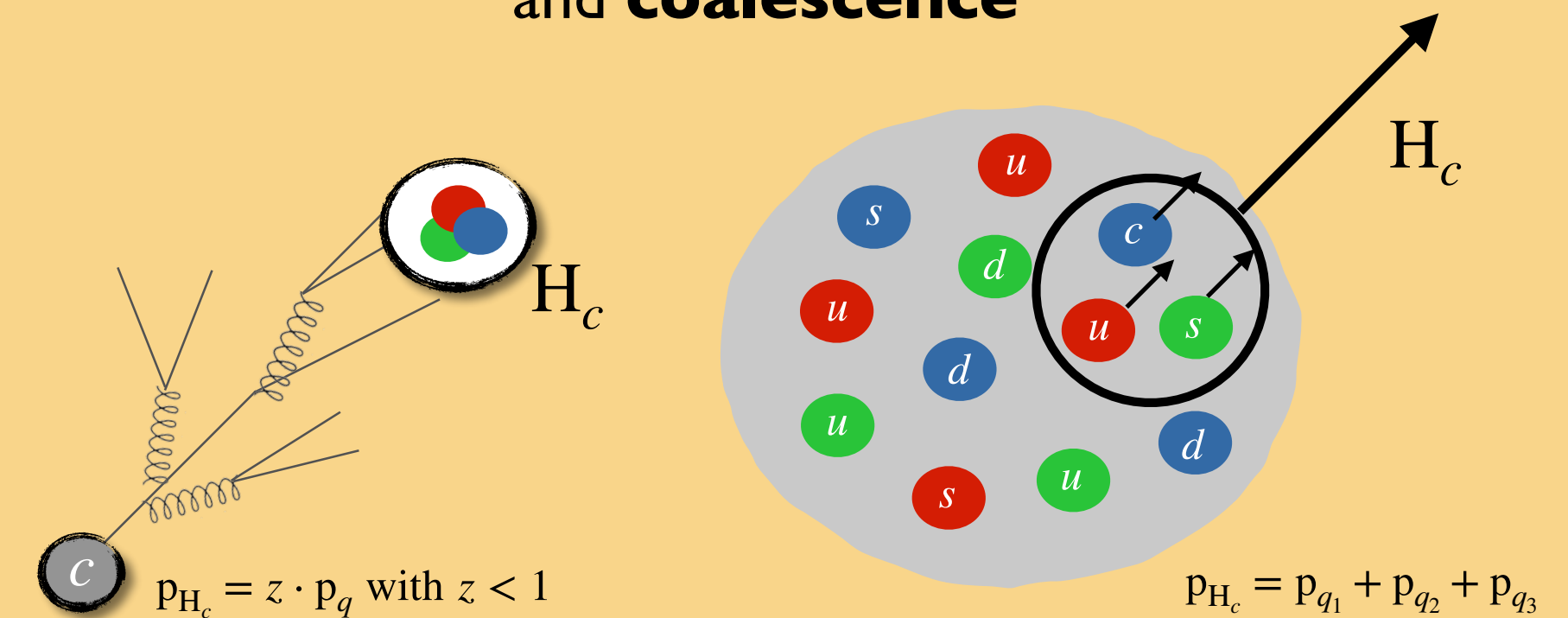
[Phys.Rev.D. 84 \(2011\) 014025](#)

- Complexity of hadronization process replaced by **statistical weights** governed by hadron mass
- Feed-down from largely **augmented set of charm baryon states** beyond the ones currently listed in the PDG, as predicted by Relativistic Quark Model

## CATANIA

[Phys.Lett.B 821 \(2021\) 136622](#)

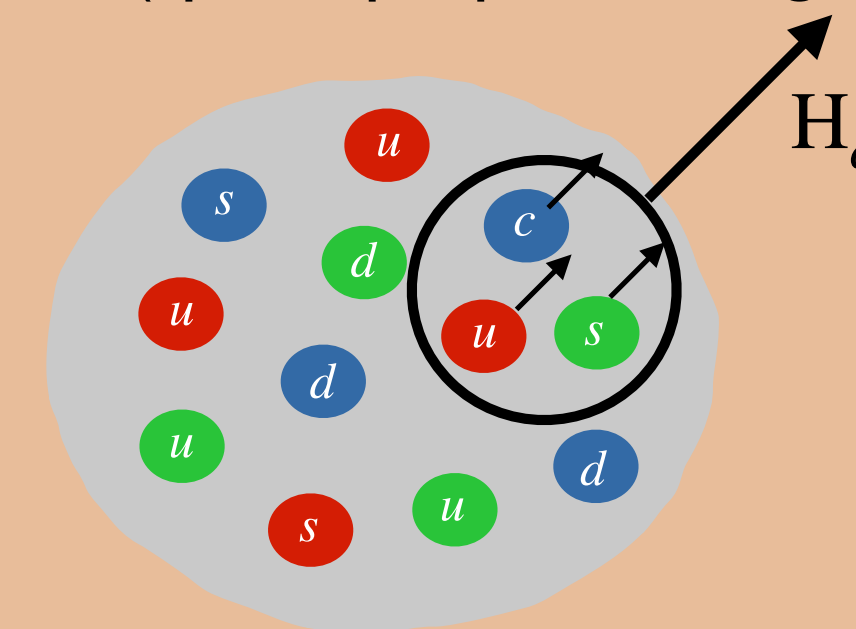
Hadronization via both **fragmentation** and **coalescence**



## QCM

[Eur.Phys.J.C 78 \(2018\) 344](#)

Quark (re-)Combination Mechanism  
**equal-velocity combination** of charm quark and light quarks (spatial properties neglected)



# Fragmentation

Independent fragmentation of partons into hadrons is the **standard way** to describe hadronization in elementary collision systems (pp, e<sup>+</sup>e<sup>-</sup>)

$$E \frac{d\sigma_H}{d^3 P_H} = E_p \frac{d\sigma_i}{d^3 p_i} \otimes \mathcal{D}_{i \rightarrow H}(z) \quad z = P_H/p$$

$\mathcal{D}(z)$  is **non-perturbative** quantity but it is considered to be universal and **usually extracted from experiments** such as e<sup>+</sup>e<sup>-</sup> collisions.

ex. Peterson

$$\mathcal{D}_{Q \rightarrow H}(z) \propto \frac{1}{z \left[ 1 - \frac{1}{z} - \frac{\epsilon}{1-z} \right]^2} \quad \epsilon = m_q^2/m_Q^2$$

ex. in PYTHIA with a modified Lund string fragmentation function

$$\mathcal{D}_{Q \rightarrow H} \propto \frac{1}{z^{1+rbm_Q^2}} z^{a_\alpha} \left( \frac{1-z}{z} \right)^{a_\beta} \exp \left( -\frac{bm_T^2}{z} \right)$$

# Question on the universality

## Fragmentation Issues

### Fragmentation Function (FF):

provides information about the energy fraction which is transferred from quark to a given meson (the larger  $m_Q$  the harder the fragmentation function)

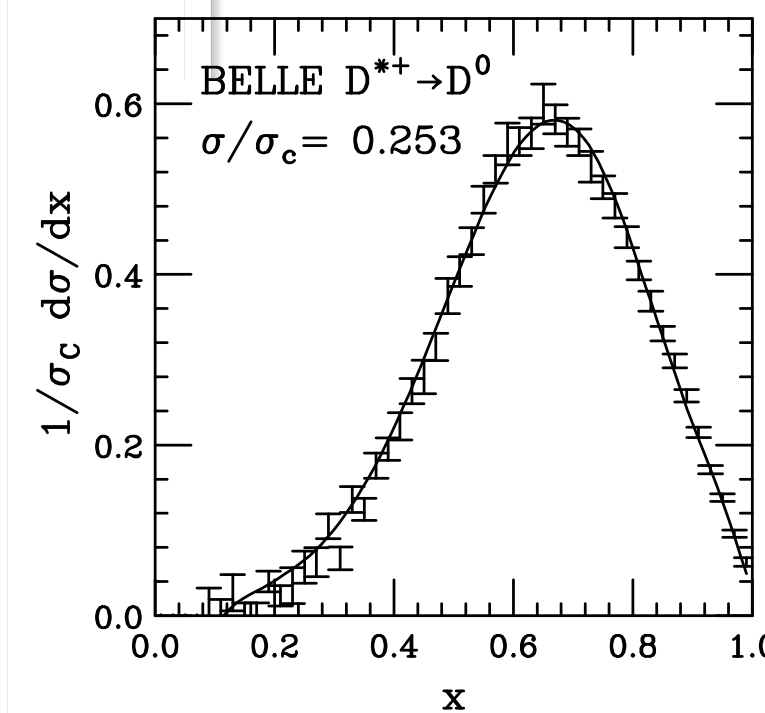
### Questions to be answered:

▷ what's the **proper parametrization** of non-perturbative frag. function?

- Peterson:  $f(z) \propto 1/[z(1 - \frac{1}{z} - \frac{\epsilon}{(1-z)})^2]$
- Kartvelishvili:  $f(z) \propto z^\alpha(1 - z)$
- Lund symmetric:  $f(z) \propto \frac{1}{z}(1 - z)^a \exp(-\frac{bm_t^2}{z})$
- Bowler:  $f(z) \propto \frac{1}{z^{1+rbm_t^2}}(1 - z)^a \exp(-\frac{bm_t^2}{z})$

▷ is fragmentation function **universal**?  
(i.e. are FF portable from  $e^+e^-$  to  $ep$  and  $pp$ ?)

- ▷ different observable definitions
- ▷ different center of mass energies, thus different pert. components as well  
⇒ **Direct shape comparison impossible!**



Fit to BELLE data  
(Cacciari, Nason, Oleari)

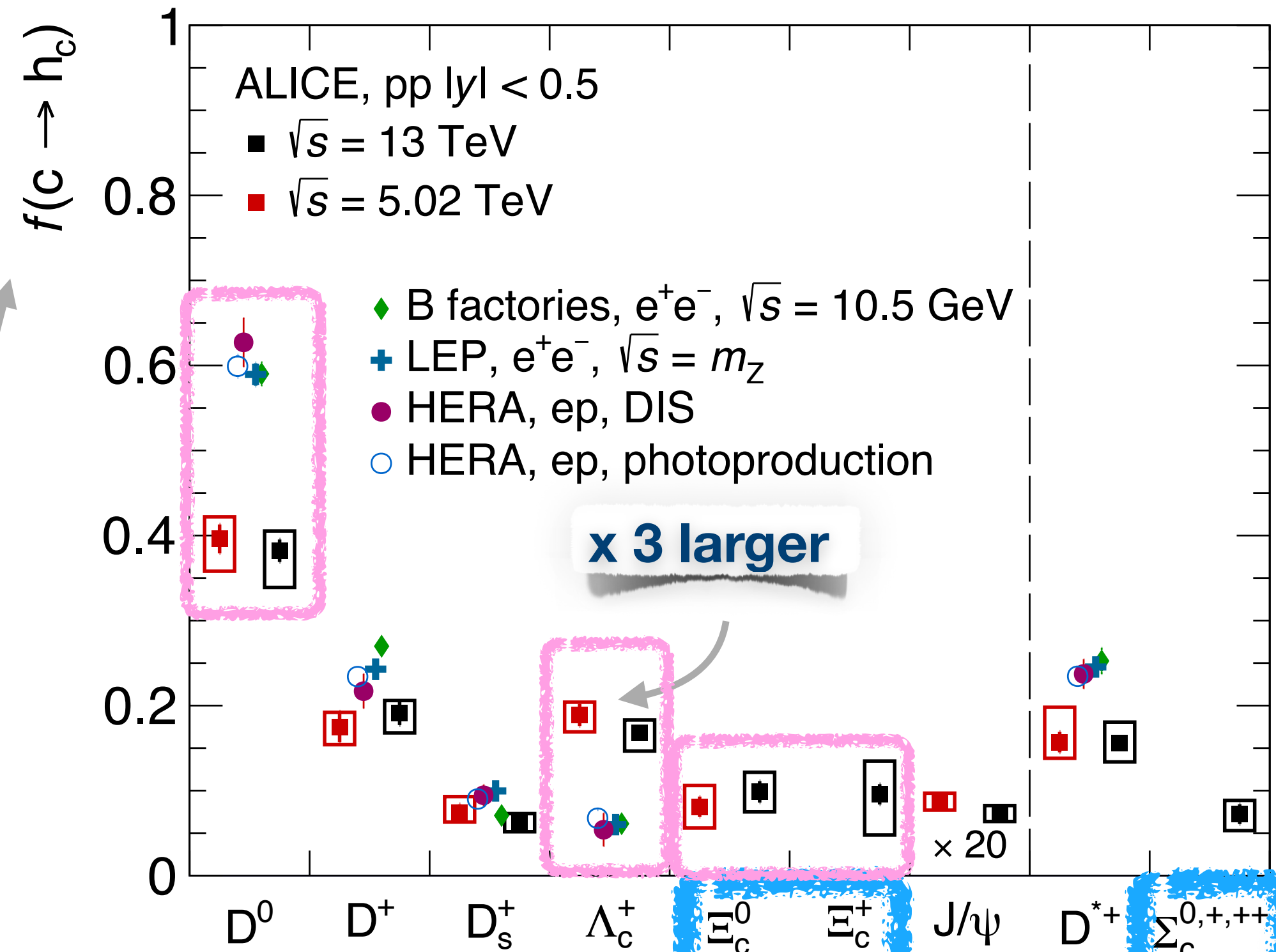
▷ **Fitted parametrization:**  $f(x) \propto \delta(1 - x) + \frac{c}{N_{a,b}}(1 - x)^a x^b$

▷ **ALEPH:**  $a = 2.4 \pm 1.2, b = 13.9 \pm 5.7, c = 5.9 \pm 1.7$

▷ **CLEO/BELLE:**  $a = 1.8 \pm 0.2, b = 11.3 \pm 0.6, c = 2.46 \pm 0.07$

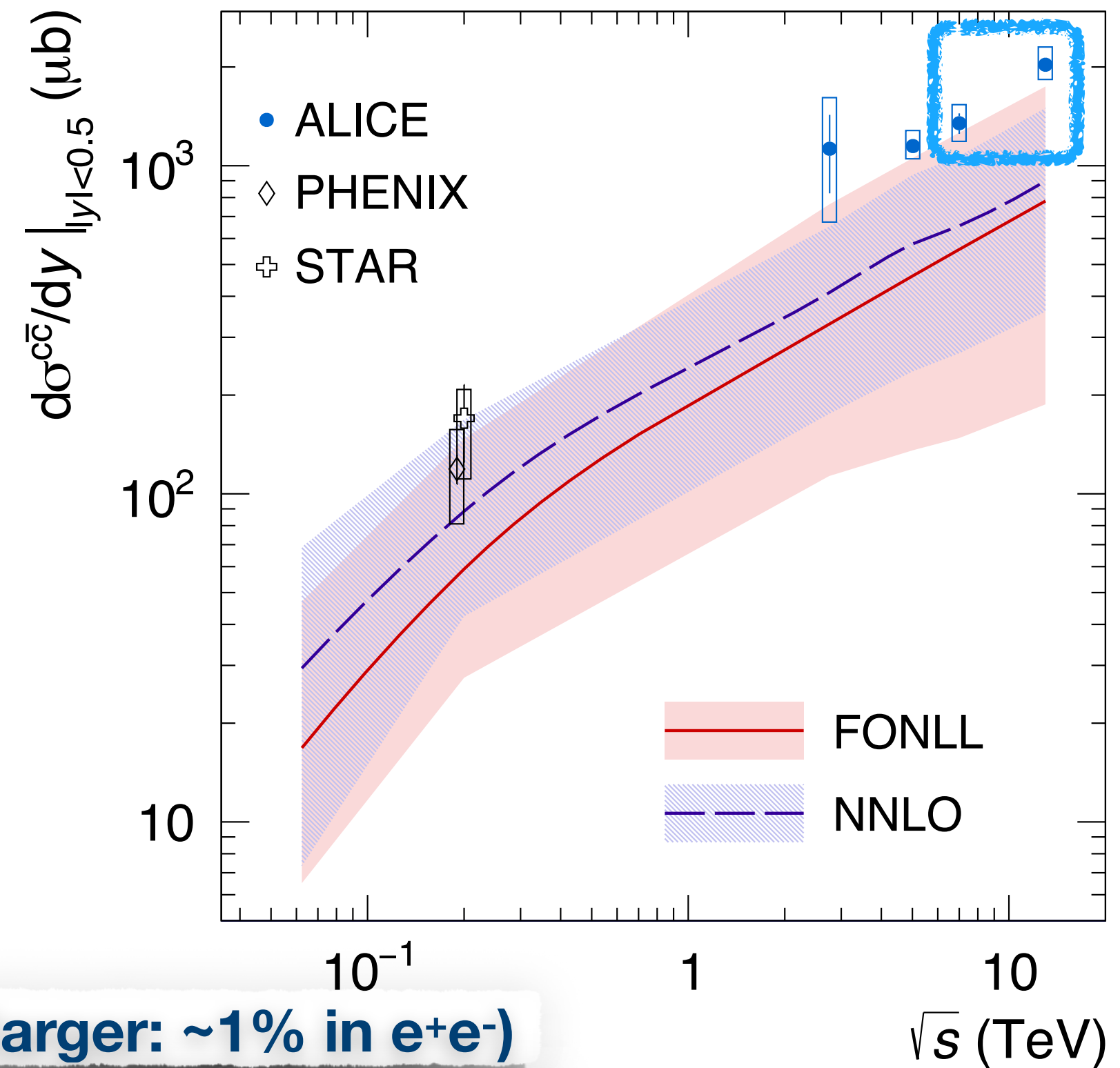
**Fits not in agreement! Does universality of  $FF_{np}$  not hold?**

# Charm-quark fragmentation fraction



10% of total charm cross section (considered negligible in  $e^+e^-$ )

$\Sigma_c^0$ : Larger feed-down to  $\Lambda_c^+$  (40%, 17% in  $e^+e^-$ )



Used the sum of the  $p_T$ -integrated cross sections of  $D^0$ ,  $D^+$ ,  $D_s^+$ ,  $J/\psi$ ,  $\Lambda_c^+$ ,  $\Xi_c^0$ ,  $\Xi_c^+$

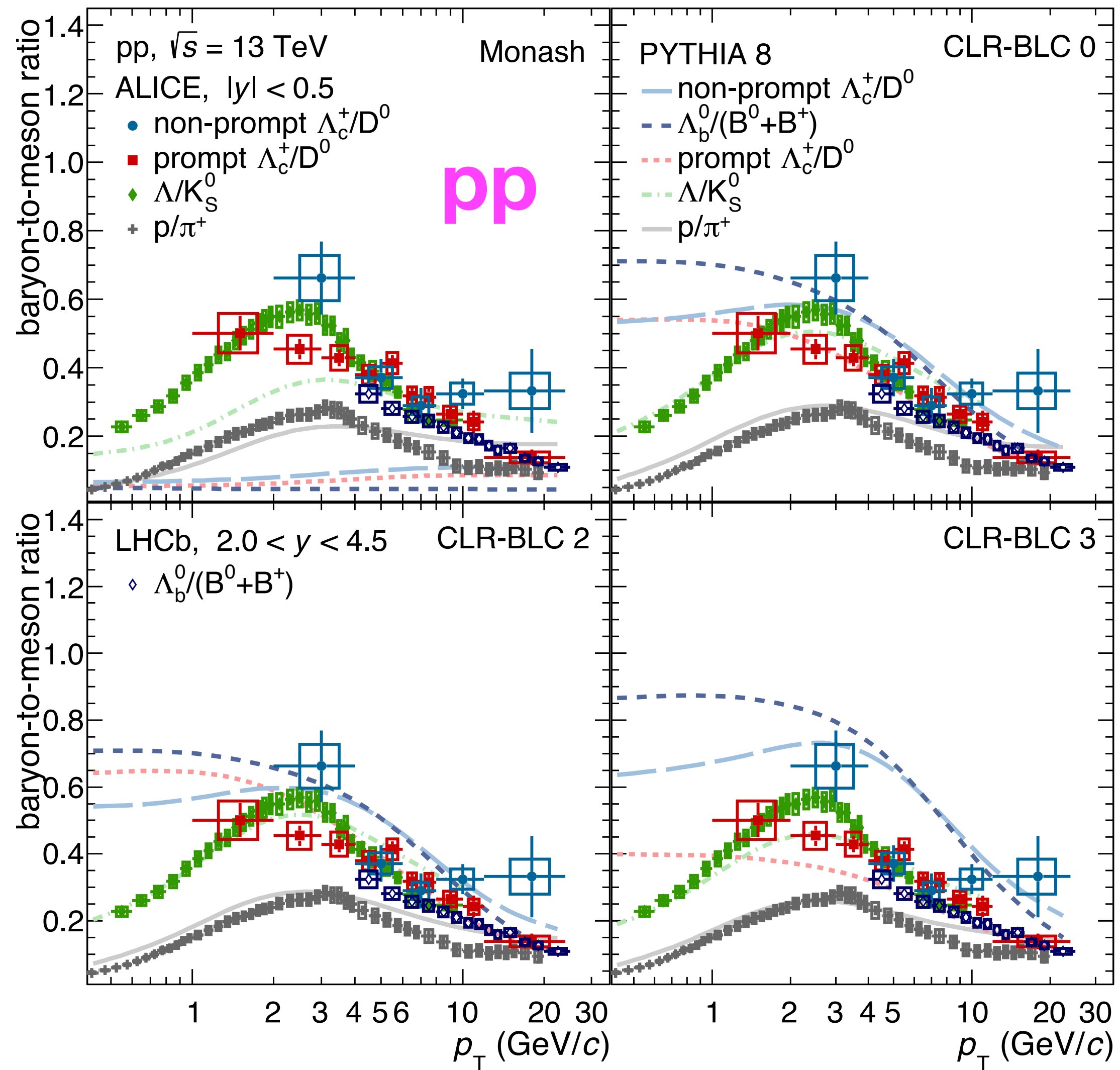
JHEP 12 (2023) 086

Normalized by the sum of the  $p_T$ -integrated cross sections of  $D^0$ ,  $D^+$ ,  $D_s^+$ ,  $J/\psi$ ,  $\Lambda_c^+$ ,  $\Xi_c^0$ ,  $\Xi_c^+$

**Conclusion: baryon enhancement at the LHC with respect to  $e^+e^-$  collisions is caused by different hadronisation mechanisms at play in the parton-rich environment produced in pp collisions**

# Baryon to meson ratios of different flavors

Phys. Rev. D 108, 112003 (2023)



- All the measurements for beauty, charm, and strange hadrons show a similar trend as a function of  $p_T$  and are compatible within the uncertainties

→ Similar baryon-formation mechanism among light, strange, charm and beauty hadrons?

- non-prompt  $\Lambda_c^+/D^0$
- prompt  $\Lambda_c^+/D^0$
- ◆  $\Lambda/K_S^0$
- +  $\rho/\pi^+$
- ◆  $\Lambda_b^0/(B^0+B^+)$

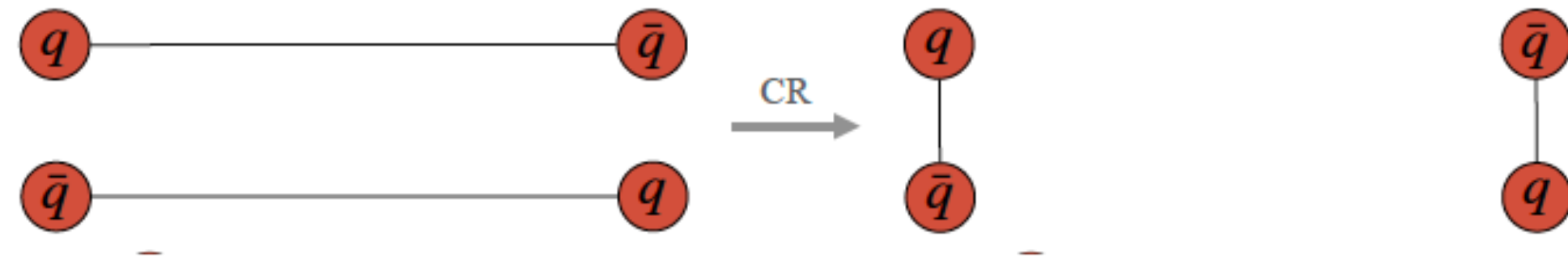
Note: for LHCb, different normalization & should consider decay kinematics (for the other case)

\* These three tunes are characterized by different constraints on the time dilation and causality

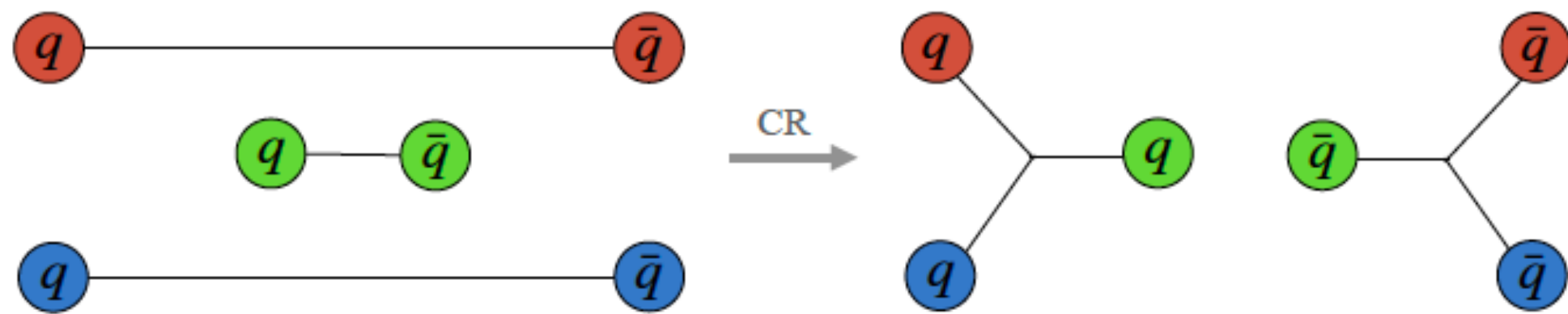


# PYTHIA Color Reconnection

Altmann et al., arXiv 2405.19137



(a) Dipole-type reconnection.



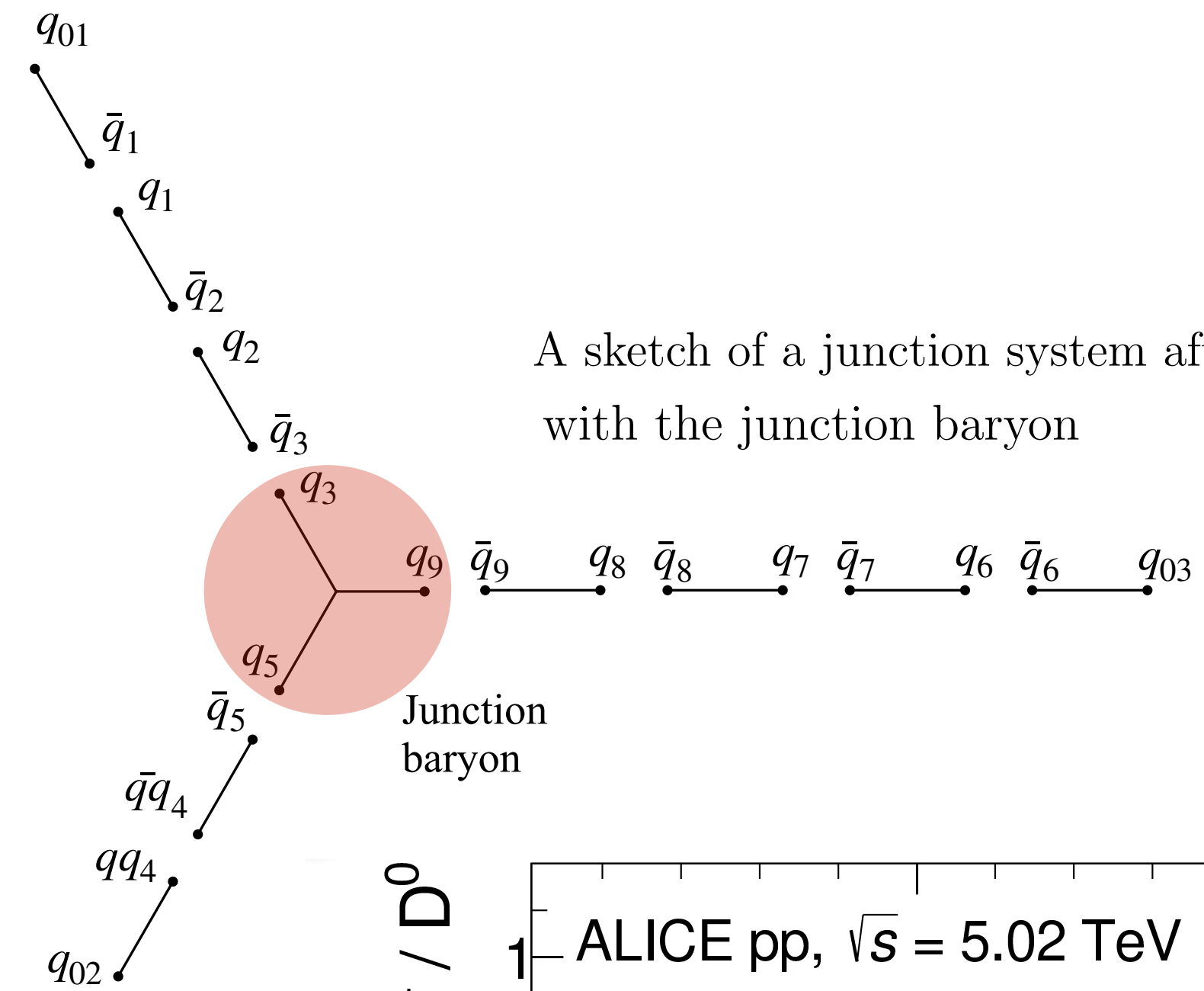
(b) Junction reconnection.

- When string color reconnection is switched-on in pp:

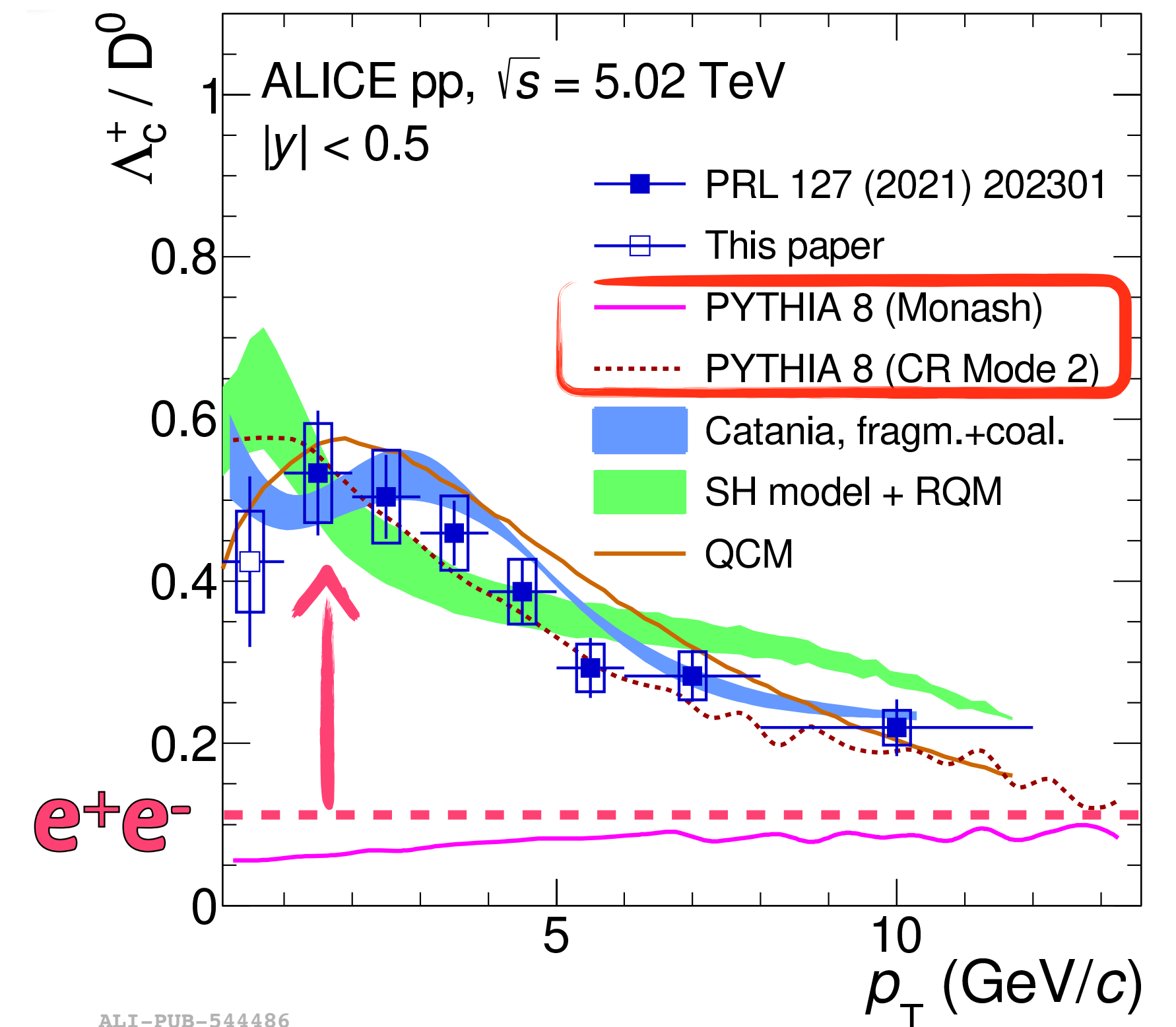
→ Very large baryon  $\Lambda_c$  enhancement

→ not that relevant for D

Not so different qualitatively wrt Coalescence and POWLANG Local color recombination



A sketch of a junction system after fragmentation, with the junction baryon

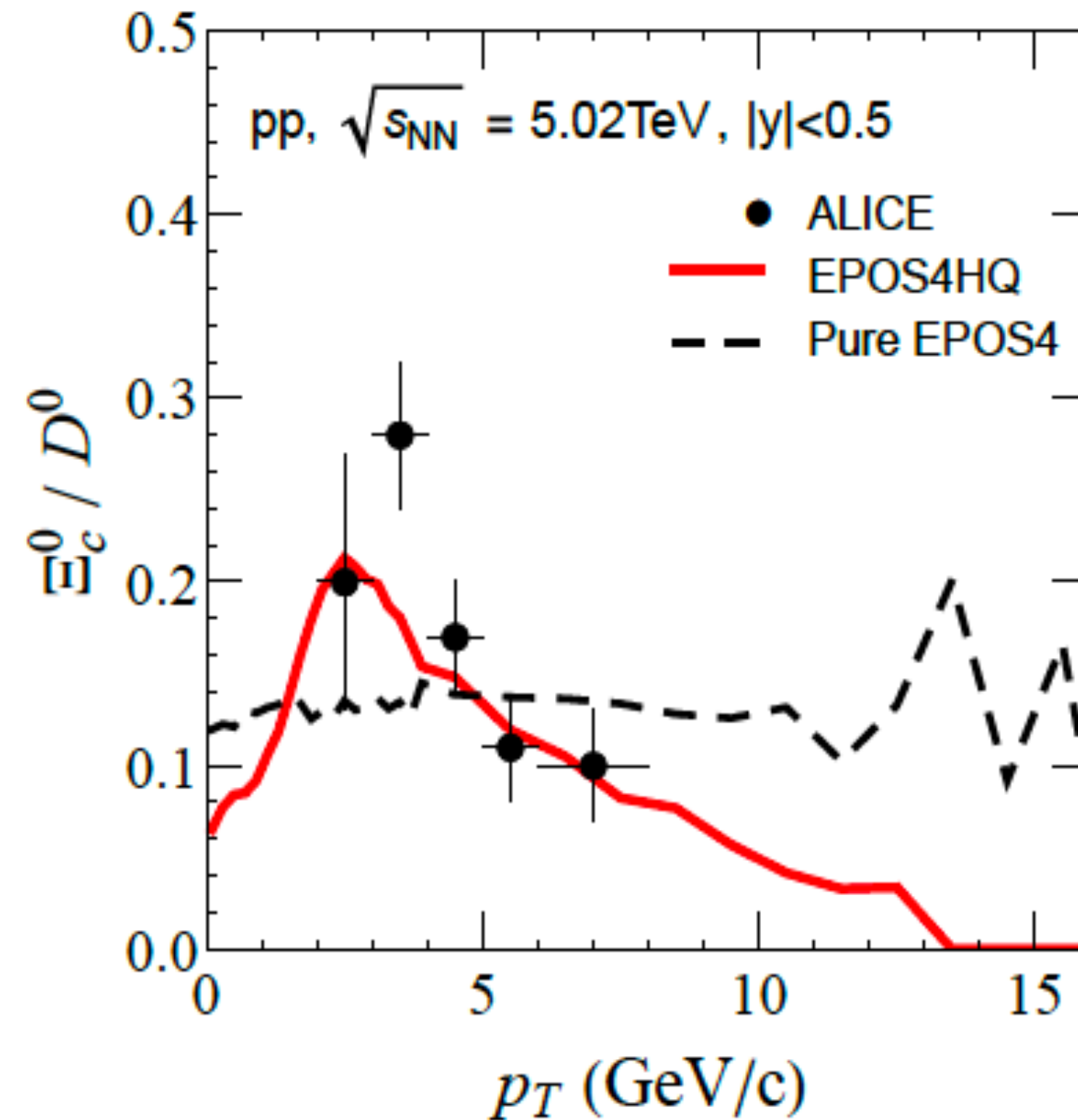
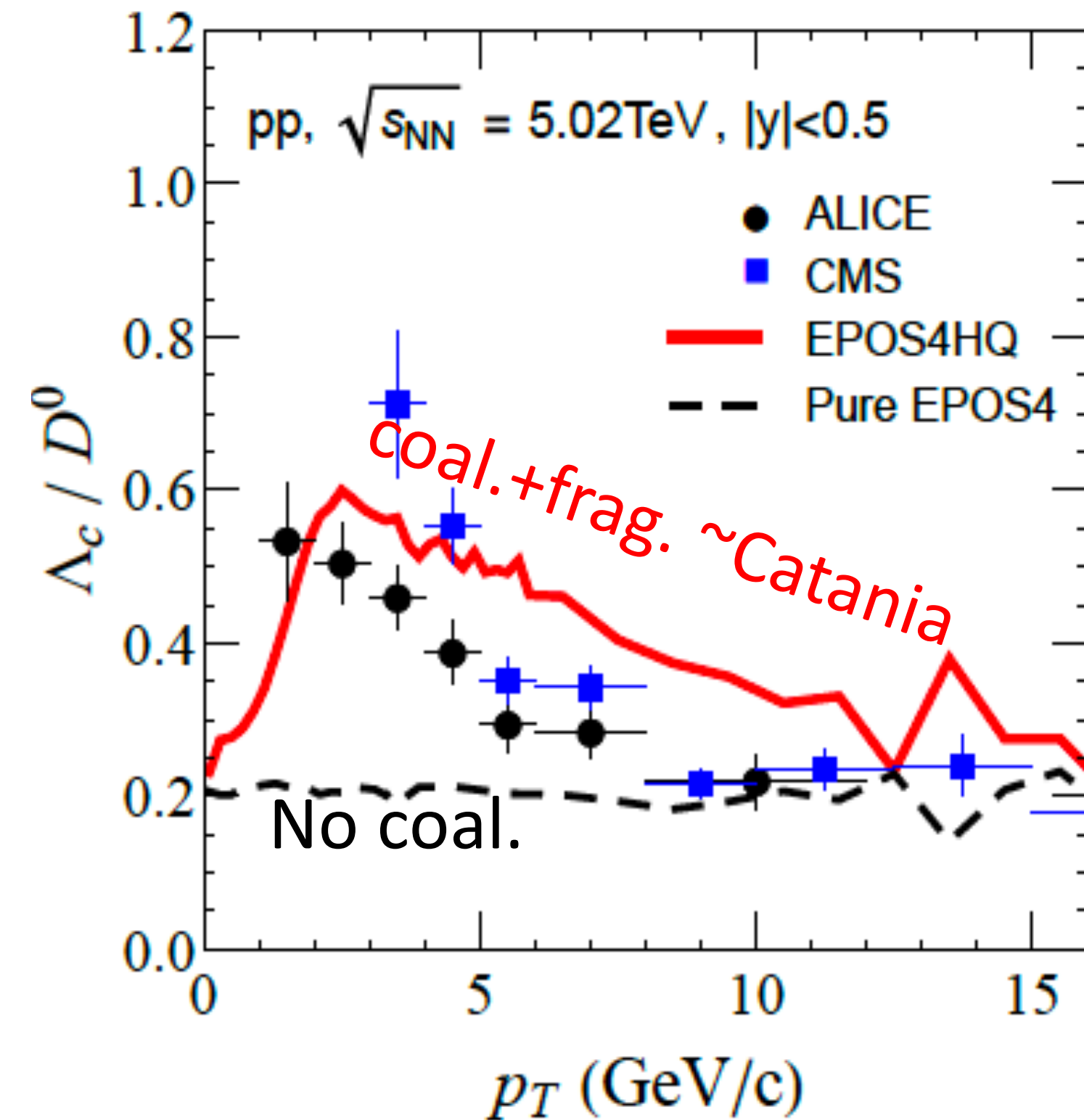


# Many models in market enhancing baryon production

- Coalescence [+Fragmentations] model:

→ Catania, Coal-TAMU(KO), Ko-Cao, CCNU-Duke, [QCM], PHSD, RRM-TAMU, Nantes-EPOS4HQ,...

J.Zhao et al., PRD109 (2024)



Ex) EPOS4HQ

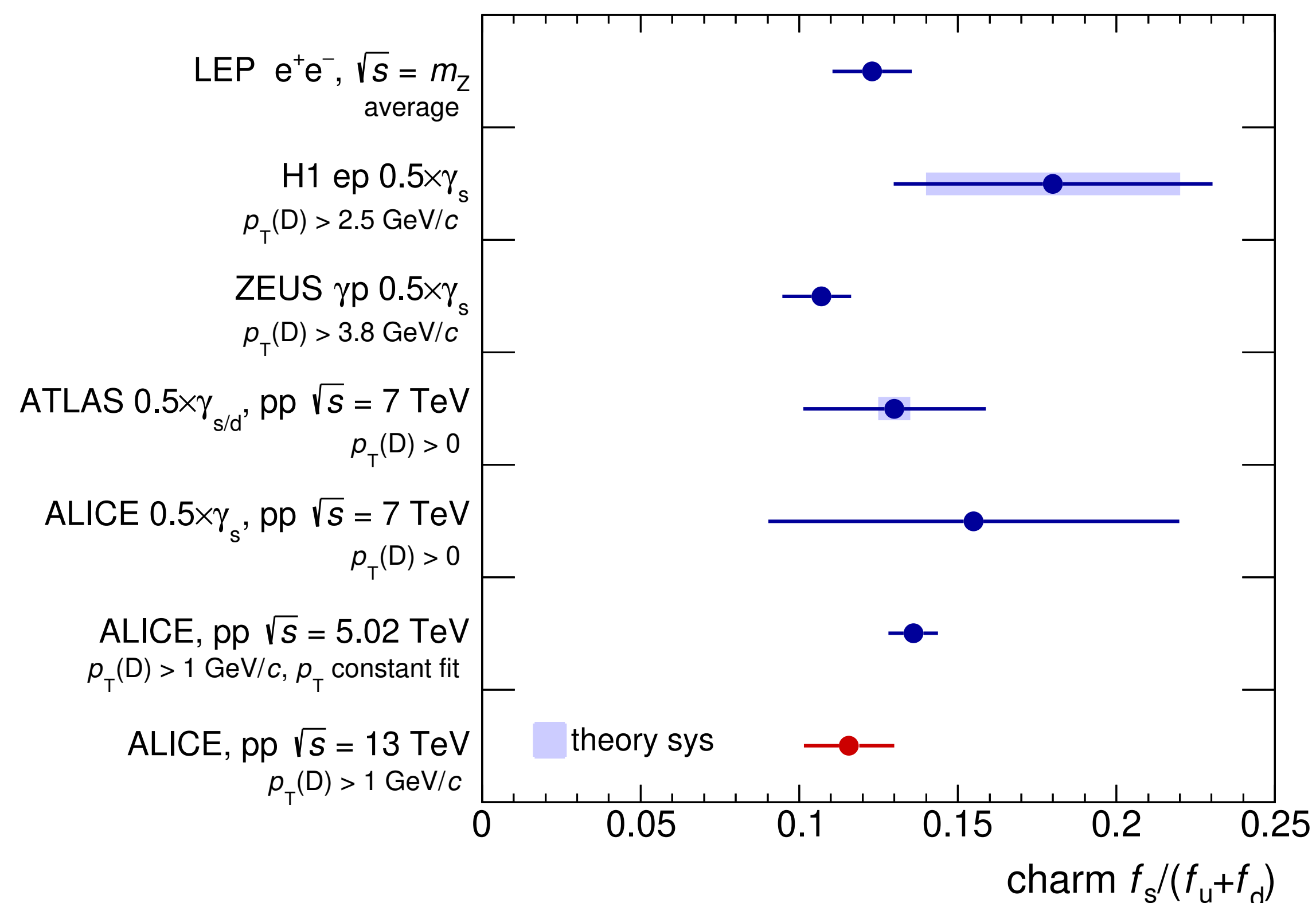
→ To describe HF spectra & ratios needs  
Coalescence in phase space ~Catania

Only difference wrt Catania:

- Assume RQM states like in SHM

# Charm-quark fragmentation-fraction ratio

## Strange to non-strange charm-meson production ratio



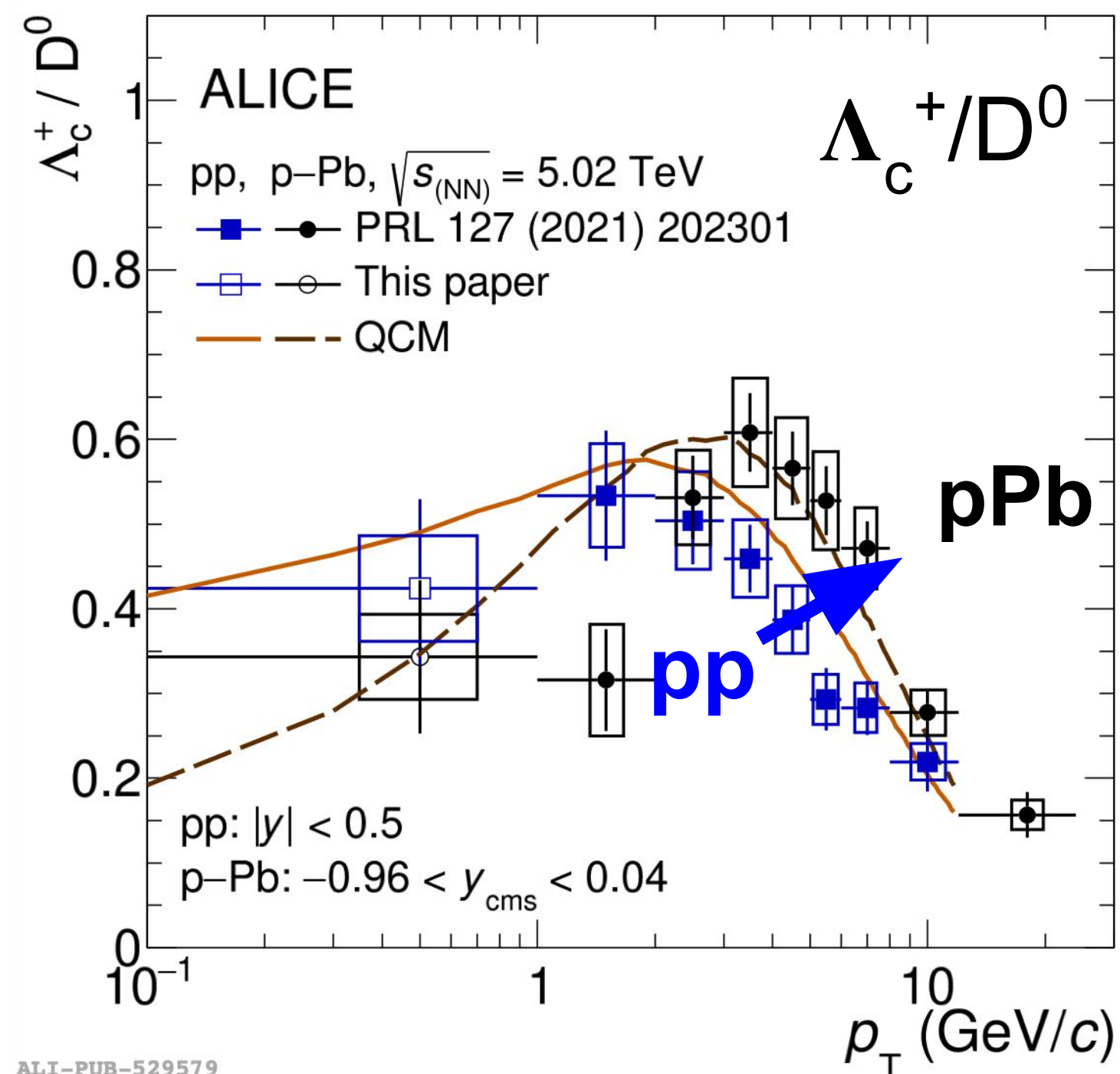
	$d\sigma/dy _{ y <0.5} (\mu\text{b}), p_T > 0$			
$D^0$	$749 \pm 27$ (stat.)	$^{+48}_{-50}$ (syst.)	$\pm 12$ (lumi.)	$\pm 6$ (BR)
$D^+$	$375 \pm 32$ (stat.)	$^{+35}_{-35}$ (syst.)	$\pm 6$ (lumi.)	$\pm 6$ (BR)
$D_s^+$	$120 \pm 11$ (stat.)	$^{+12}_{-13}$ (syst.)	$^{+25}_{-10}$ (extrap.)	$\pm 2$ (lumi.) $\pm 3$ (BR)
$\Lambda_c^+$	$329 \pm 15$ (stat.)	$^{+28}_{-29}$ (syst.)	$\pm 5$ (lumi.)	$\pm 15$ (BR)
$\Xi_c^0$ [52]	$194 \pm 27$ (stat.)	$^{+46}_{-46}$ (syst.)	$^{+18}_{-12}$ (extrap.)	$\pm 3$ (lumi.)
$\Xi_c^+$	$187 \pm 25$ (stat.)	$^{+19}_{-19}$ (syst.)	$^{+13}_{-59}$ (extrap.)	$\pm 3$ (lumi.) $\pm 82$ (BR)
$J/\psi$ [84]	$7.29 \pm 0.27$ (stat.)	$^{+0.52}_{-0.52}$ (syst.)	$^{+0.04}_{-0.01}$ (extrap.)	
$D^{*+}$	$306 \pm 26$ (stat.)	$^{+33}_{-34}$ (syst.)	$^{+48}_{-17}$ (extrap.)	$\pm 5$ (lumi.) $\pm 3$ (BR)
$\Sigma_c^{0,+,++}$	$142 \pm 22$ (stat.)	$^{+24}_{-24}$ (syst.)	$^{+24}_{-32}$ (extrap.)	$\pm 2$ (lumi.) $\pm 6$ (BR)

$f_x$ : probability for a charm quark to hadronize with another quark of flavour  $x$

$\Rightarrow D_s^+/D^0+D^+$

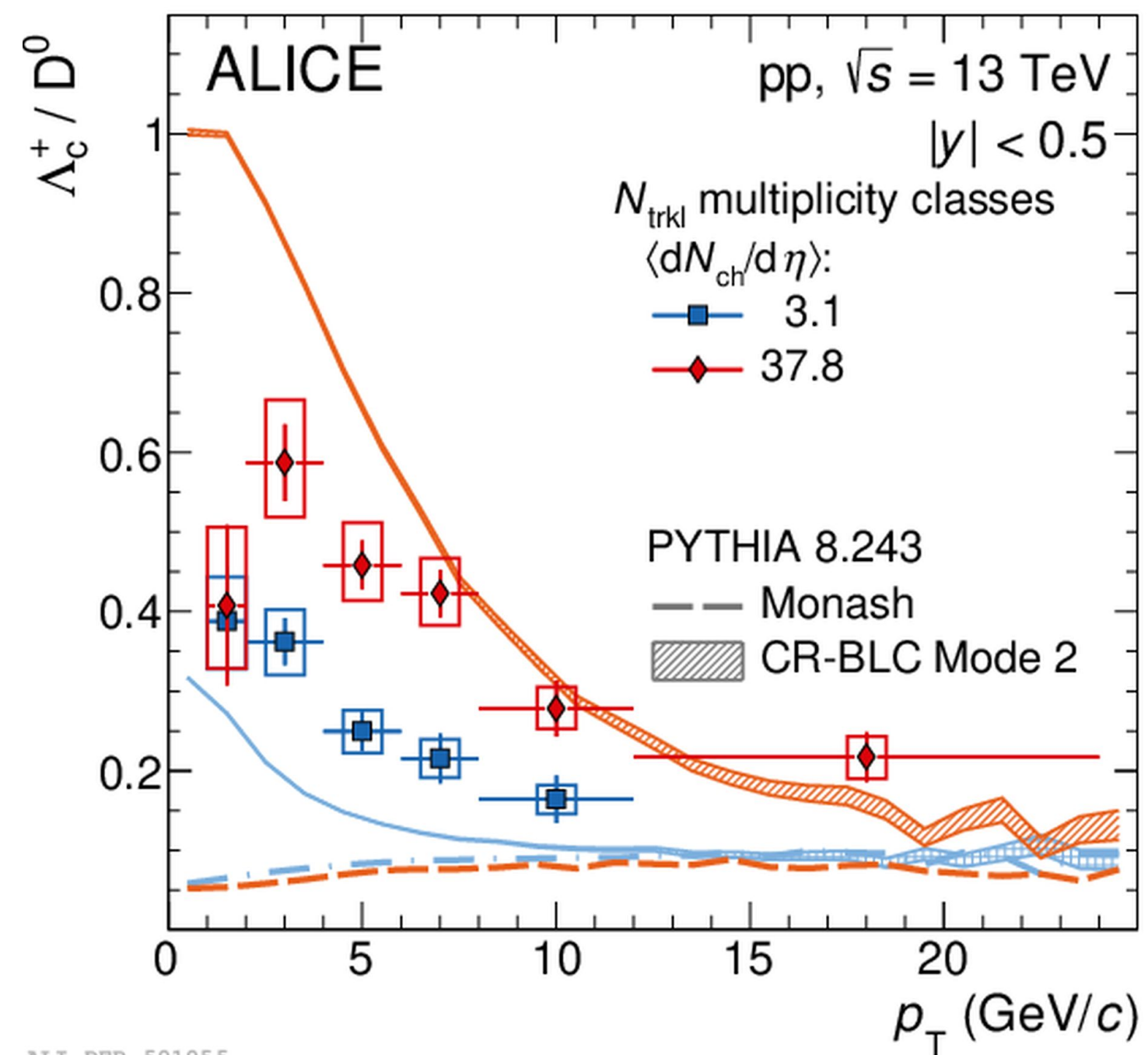
Production of **prompt strange D mesons / prompt non-strange D mesons** in  $e^+e^-$ , ep and pp collisions doesn't show any significant dependence of the collision system & energy!

PRC 104 054905 (2021)



ALI-PUB-529579

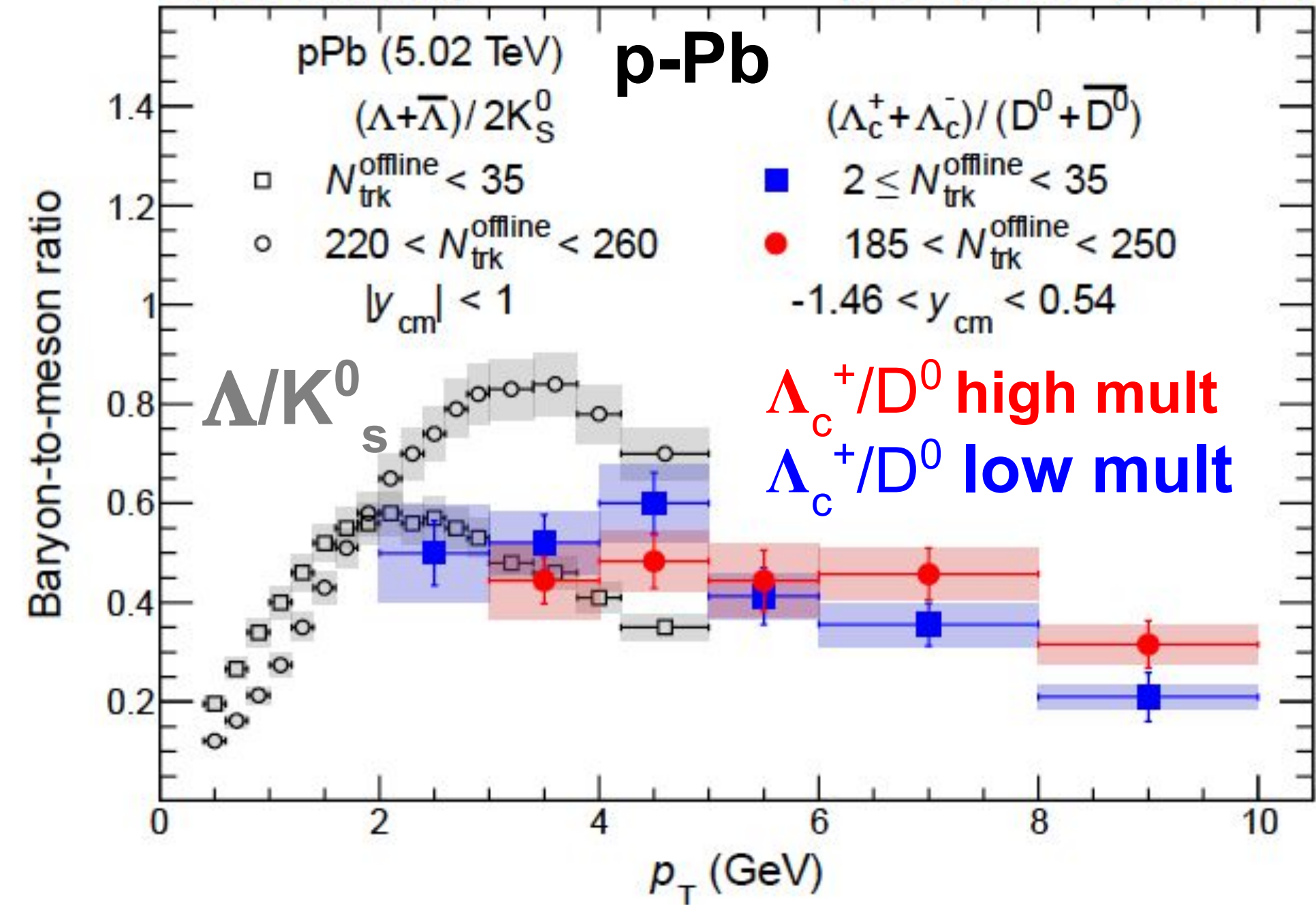
PLB 829 (2022) 137065



ALI-DER-501055

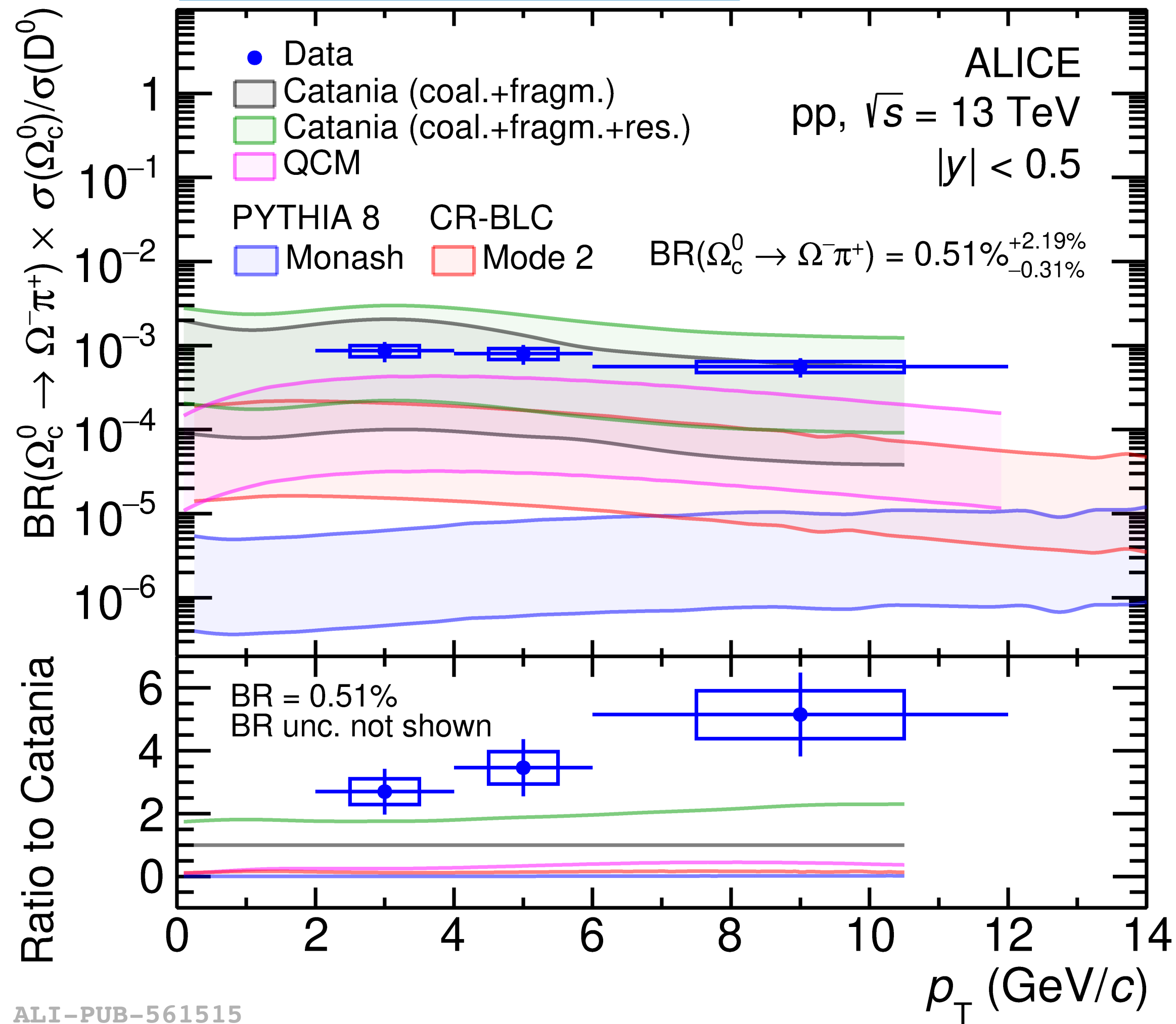
CMS Preliminary

pPb 97.8 nb<sup>-1</sup> (8.16 TeV)



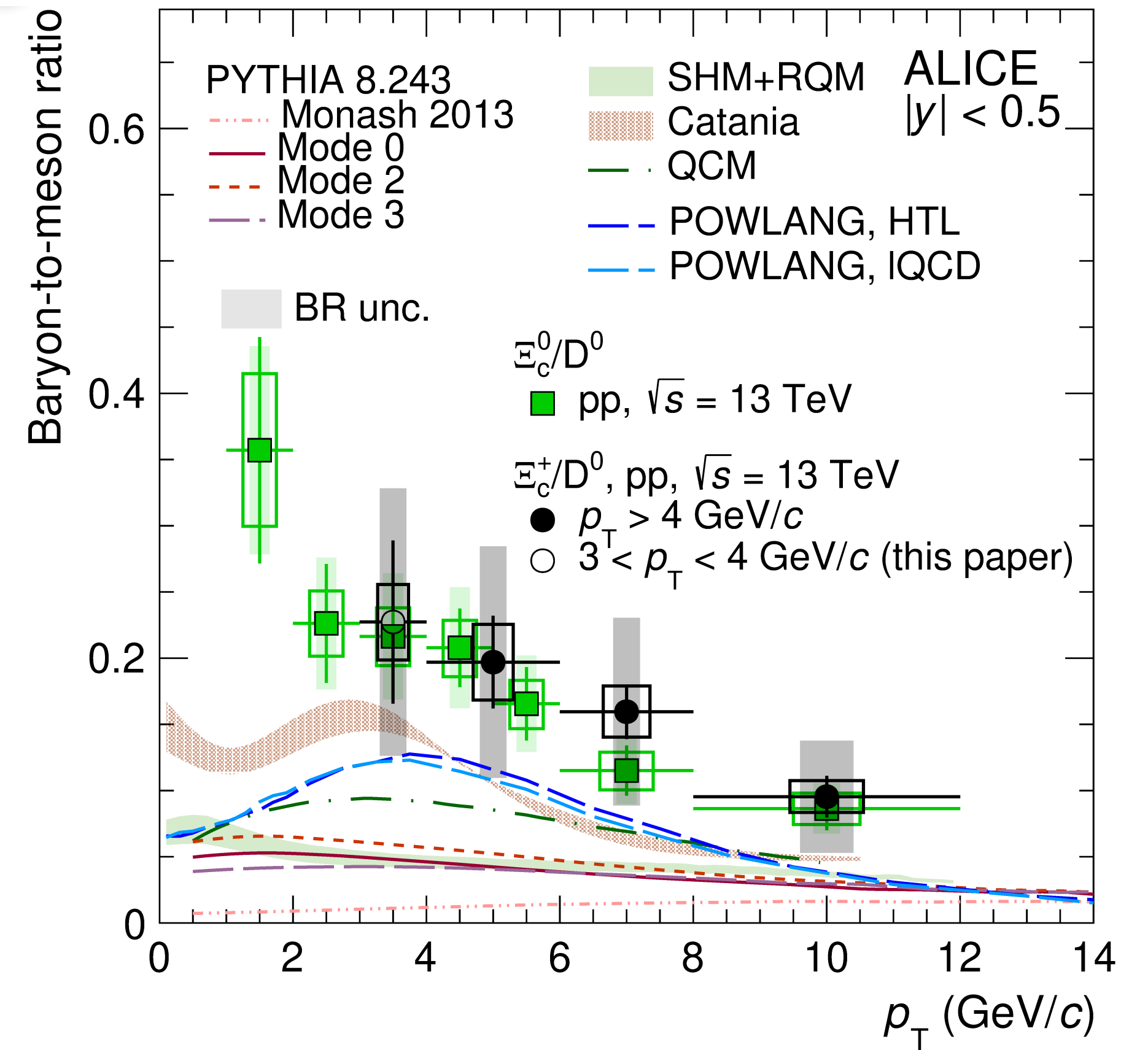
# With 2 strangeness

Phys.Lett.B 846 (2023) 137625



ALI-PUB-561515

JHEP 12 (2023) 086



ALI-PUB-567881