
SENSITIVITY OF JETS TO INITIAL STAGES

Based on: arXiv:2409.04295 and arXiv:2409.06675

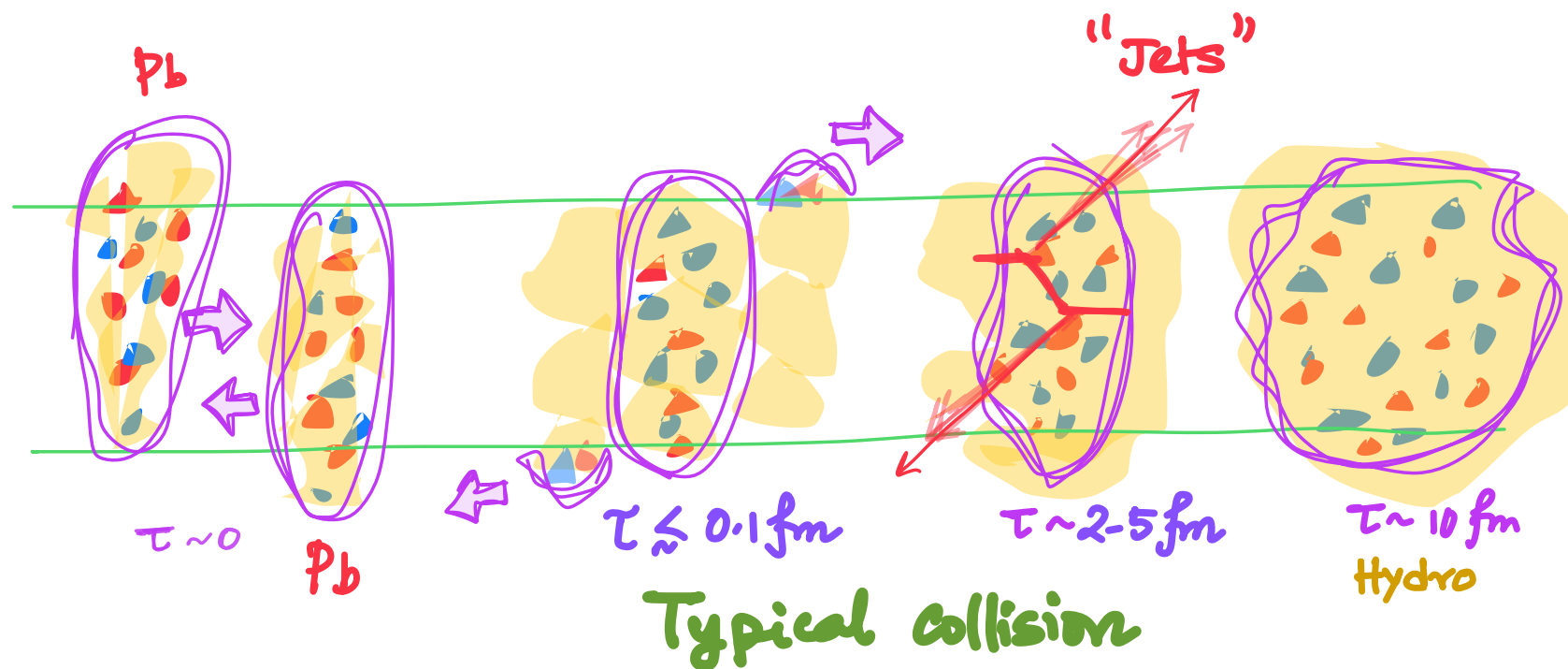
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ATHIC 2025, Gopalpur, India

Jet quenching : quick words



Refer to Introductory slides of N. R. Sahoo, V. Vaidya and S. Oh (Tuesday)

Our field has explored ... and lacked ...

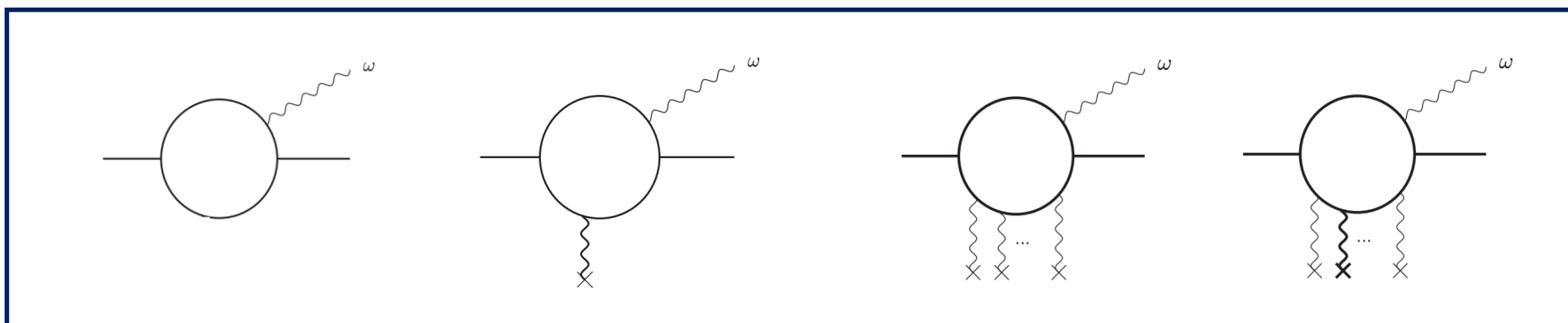
- **QUESTION** : Range of **applicability** of the **hard vs the soft scattering** for jet quenching models.
- A universal splitting rate across all gluon frequencies ?
- **Finite medium size effects** :
 - include realistic medium scenarios (**expanding medium**) relevant for inclusion in phenomenological in-medium parton shower models.
 - validity of the soft multiple and hard scattering not only as a function of energy but also as a **function of the initial quenching time of the medium**.
 - **Rethink** : *Are multiple scatterings important for radiative in-medium parton showers !*

Regimes of phase space emissions

- Various MC in-medium parton showers use two **analytical** approaches :
 - **DILUTE** medium: Single-hard scattering approximation (Opacity expansion).
 - **DENSE** medium: multiple-soft scattering. Harmonic oscillator (HO) approach [1]
- Full **numerical** solutions [2]

Opacity $\chi \equiv L/\lambda \Rightarrow$ denseness of the medium.

- ($L \ll \lambda$) : Medium **DILUTE**, or weakly interacting and ($L \gg \lambda$) : Medium **DENSE**, or strongly interacting



NLO : One hard scattering included (thick wavy line)

+

Multiple soft scattering re-summed to all orders

LO (N=0):
vacuum radiation

NLO (N=1): In medium
Single scattering

LO (HO): Multiple soft
scatterings (wavy vertical lines)

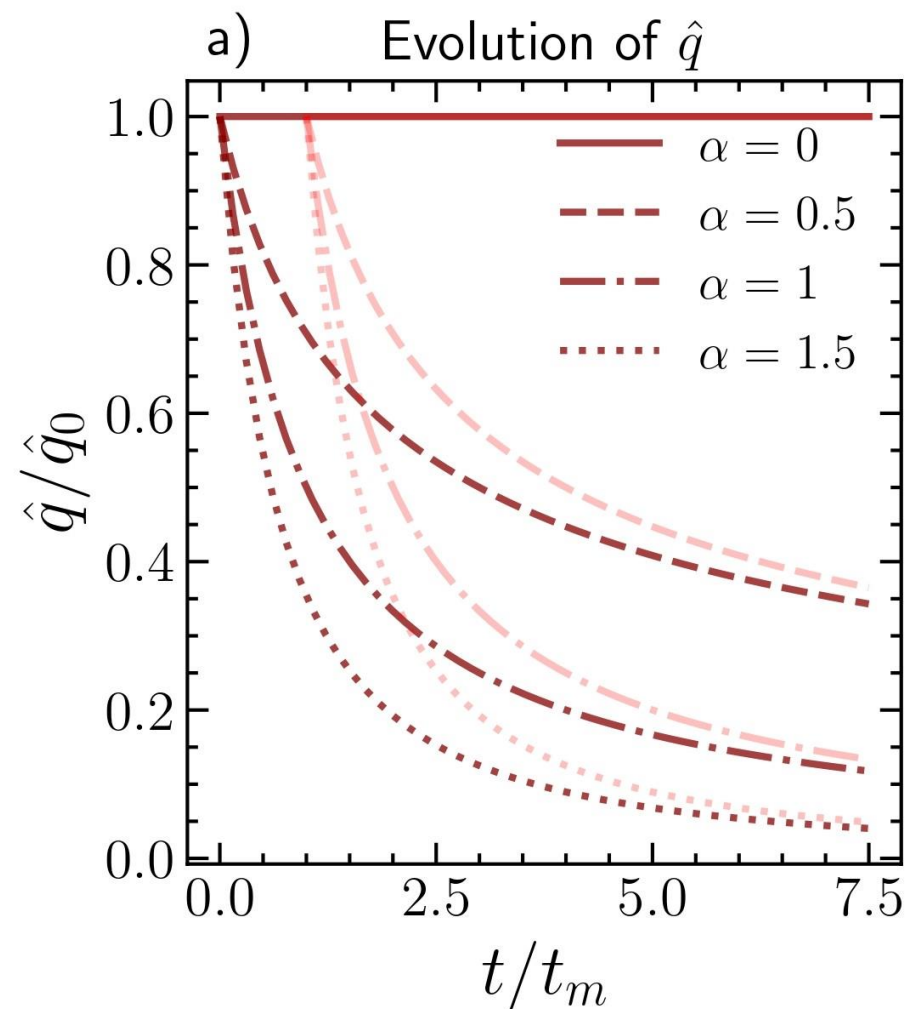
Improved opacity expansion

- Two models for the evolution of quenching parameter \hat{q} :

- model (i) : initially over-occupied system.
- model (ii) : initially under-occupied system.

$$\hat{q}_0(t) = \begin{cases} \hat{q}_0 \left(\frac{t_m}{t+t_m} \right)^\alpha & \text{for model (i), [dark line]} \\ \hat{q}_0 \Theta(t - t_m) \left(\frac{t_m}{t} \right)^\alpha & \text{for model (ii) [light line]} \end{cases}$$

- Previous works have explored the sensitivity of the jets to such expanding model scenarios on the nuclear modification factor and the elliptic flow [Andres 2022, Caucal 2021, SPA 2020-24].



Medium rates : complete analytical rates

- We first derive analytical in-medium emission rates to calculate energy loss that are valid for any medium expansion scenario.

$$\Gamma = \frac{4\alpha_s C_R}{\omega^2} \int_0^t dt_0 \int_{\mathbf{p}, \mathbf{p}_0} \Sigma(\mathbf{p}^2, t) \frac{\mathbf{p} \cdot \mathbf{p}_0}{p^2} \tilde{\mathcal{K}}(\mathbf{p}, t; \mathbf{p}_0, t_0)$$

- Three-point correlator $K(\mathbf{p}; \mathbf{p}_0)$ describes the transverse momentum broadening experienced by the gluon during its formation time.
- We use widely used Gyulassy-Wang (GW) model for medium interaction, with screening mass μ .

$$\Sigma(\mathbf{k}^2, t) = \frac{\hat{q}_0(t)}{\mathbf{k}^2 + \mu^2}$$

Medium rates : complete analytical rates

We will employ three resummation schemes that, in total, cover the whole emission phase space :

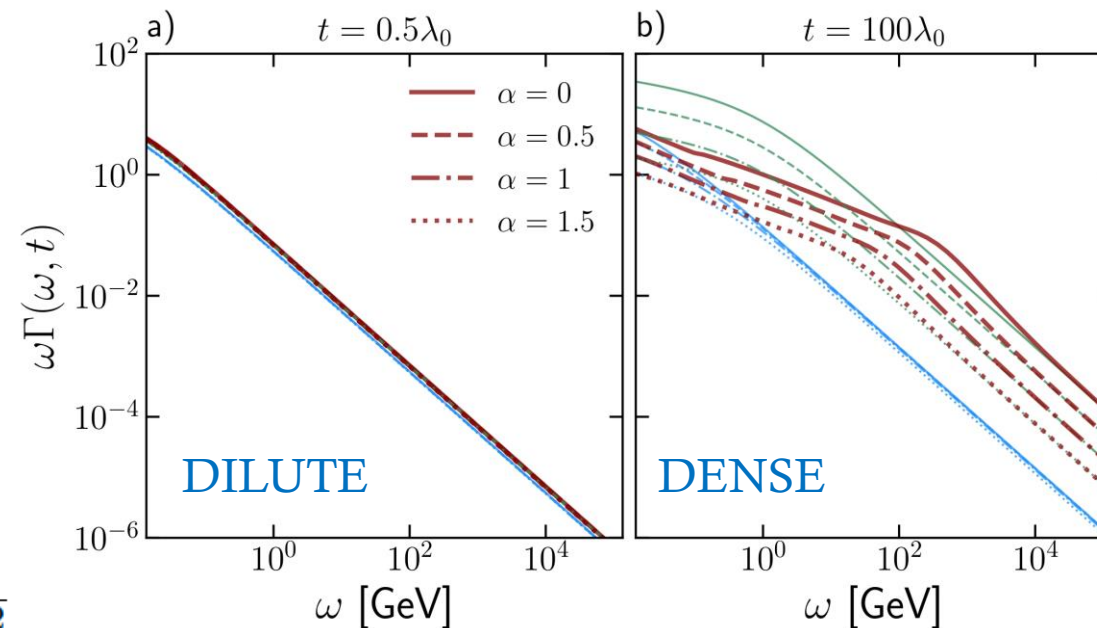
- Opacity expansion (N = 1)/GLV
- Improved opacity expansion (IOE)
- Resummed opacity expansion
- IOE : Leading-log potential split into HO and remainder:

$$v^{\text{LT}}(\mathbf{x}) \equiv v^{\text{HO}}(\mathbf{x}) + \delta v(\mathbf{x}) = \frac{1}{4} \hat{q}_0 \mathbf{x}^2 \log \frac{Q^2}{\mu_*^2} + \frac{1}{4} \hat{q}_0 \mathbf{x}^2 \log \frac{1}{Q^2 \mathbf{x}^2}$$

- Expanding around the harmonic oscillator

$$\mathcal{K}(\mathbf{x}, t_2; \mathbf{y}, t_1) = \mathcal{K}_{\text{HO}}(\mathbf{x}, t_2; \mathbf{y}, t_1) - \int_z \int_{t_1}^{t_2} ds \mathcal{K}_{\text{HO}}(\mathbf{x}, t_2; \mathbf{z}, s) \delta v(\mathbf{z}, s) \mathcal{K}(\mathbf{z}, s; \mathbf{y}, t_1)$$

- **Redefinition: “dressed” transport coefficient**



Q is a separation scale and ensures that the first term dominates.

$$\hat{q} = \frac{\langle k_{\perp}^2 \rangle_{\text{typ}}}{L} = \hat{q}_0 \log \frac{aQ_s^2}{\mu_*^2}$$

Medium rates : novel kinematical conditions

- Ratio of radiative spectrum to NLO in expansion around Harmonic oscillator (LO) gives matching scale Q .

$$\left. \frac{\Gamma_{\text{IOE}}^{(1)}}{\Gamma_{\text{IOE}}^{(0)}} \right|_{\omega \ll \omega_c(t)} = \frac{\hat{q}_0}{\hat{q}} \left[\gamma_E + \frac{\pi}{4} - \frac{3}{2} \ln 2 + \ln \frac{\sqrt{\hat{q}(t)\omega}}{Q^2} \right]$$

- Re-definition of the **scales of the problem** ($\omega_c \gg \omega_{\text{BH}}$) :

$$1 \ll \frac{t}{\lambda_0} \left(\frac{t_m}{t} \right)^\alpha$$

Scattering regimes :

- For sufficiently dilute or small media, the probability of **one scattering** dominates.
- For large or dense media, **multiple scattering** on the background medium
 - An interplay between copious, soft (LPM) interactions and rare, hard momentum exchanges.
 - **Establish whether this regime is relevant for expanding media where the density rapidly decays with proper time ?**

Medium rates : novel kinematical conditions

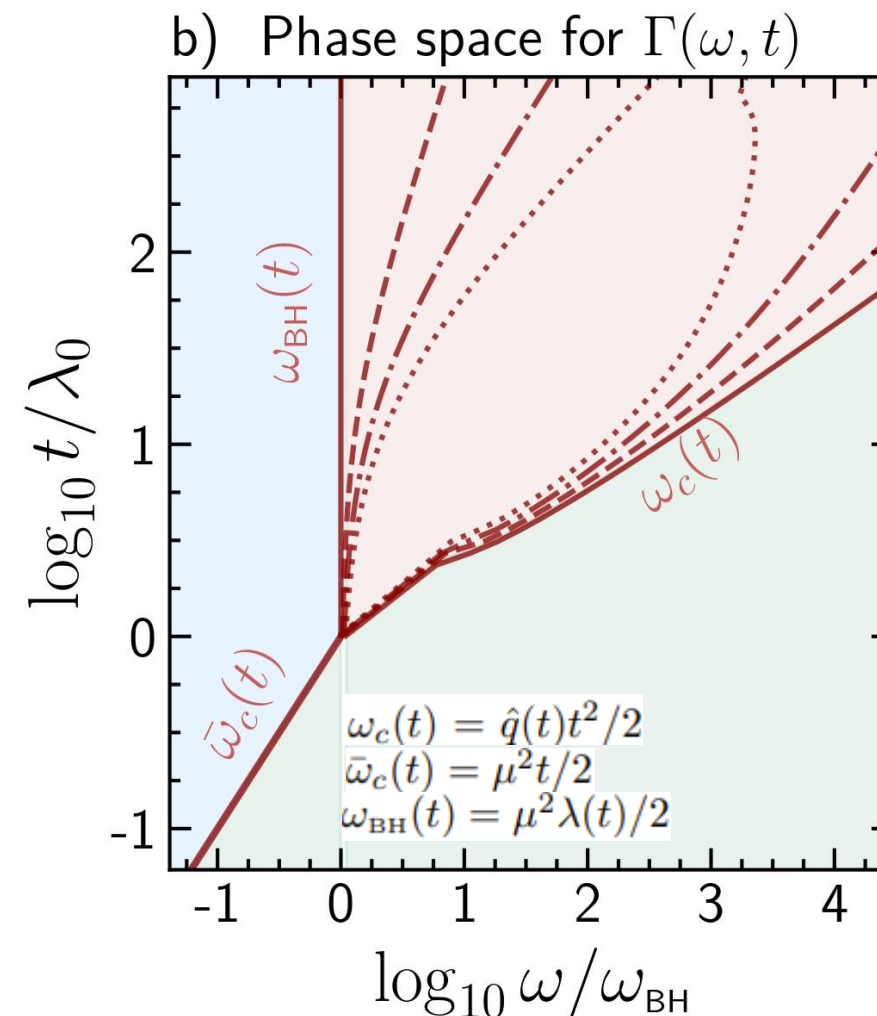
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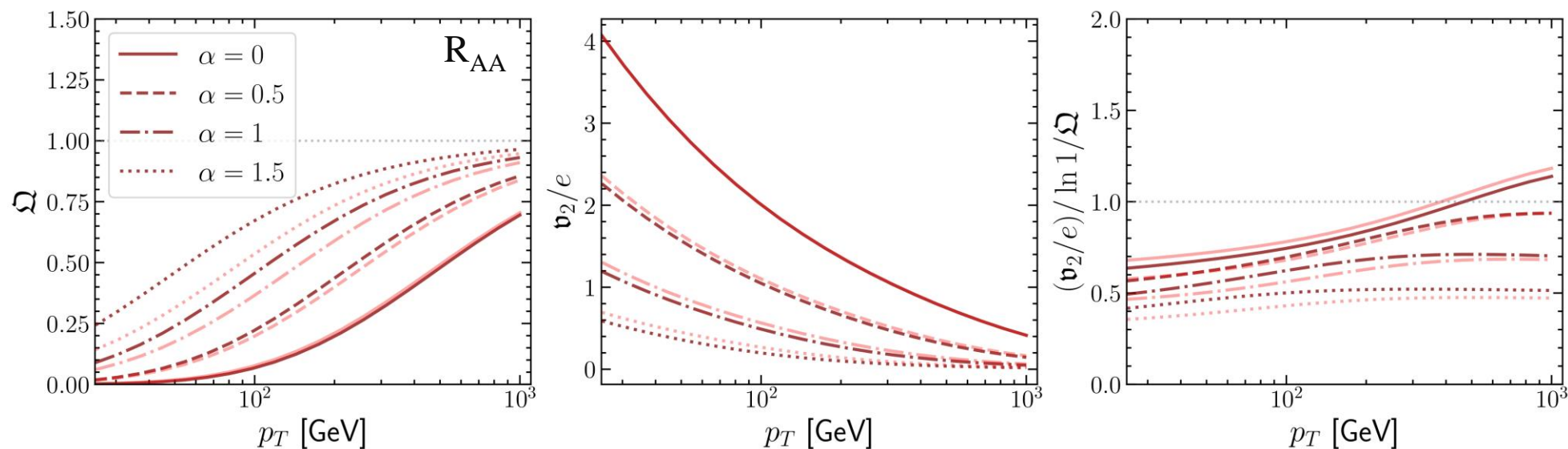
$$1 \ll \frac{t}{\lambda_0} \left(\frac{t_m}{t} \right)^\alpha$$

- Implies** : Medium “hydrodynamization” time should be much bigger than the mean-free-path $t_m \gg \lambda_0$ in order to get contributions from this regime (eg: check Bjorken expansion)



Phenomenological consequence : R_{AA} and v_2

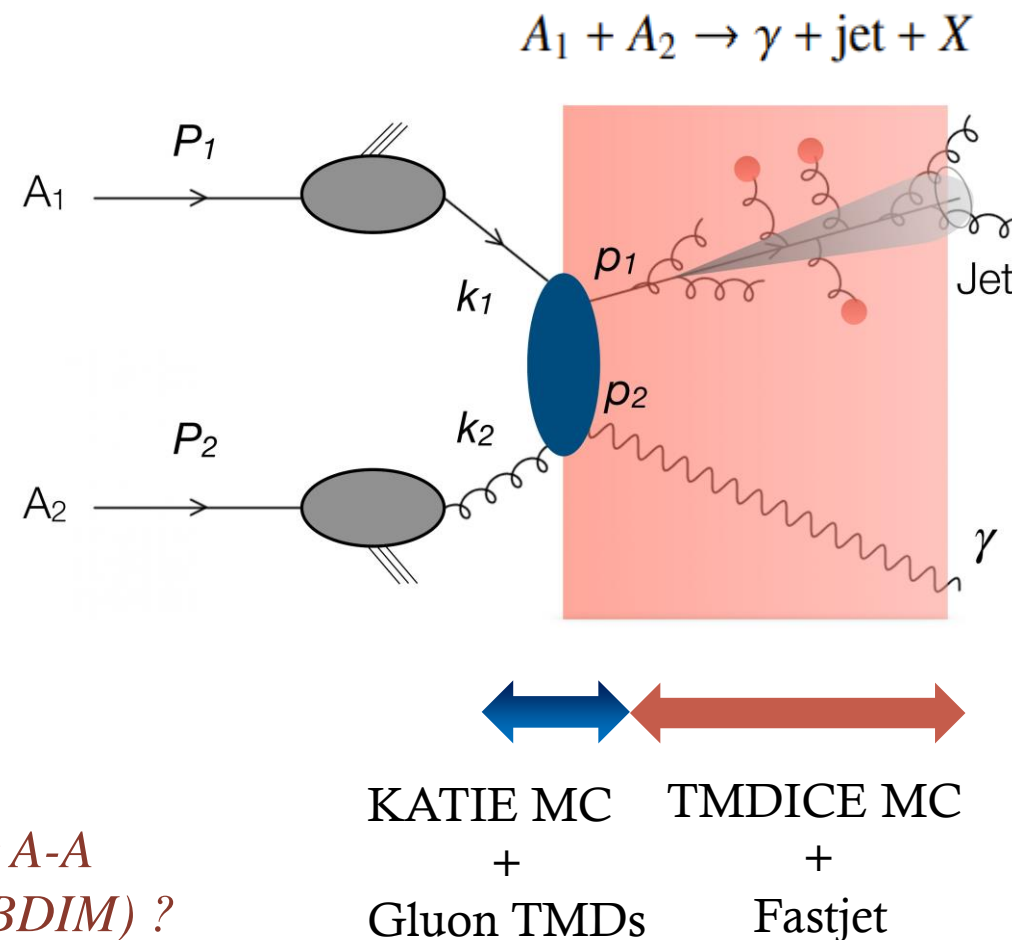
- R_{AA} sensitive to the accumulation of emissions along the entire in-medium path length L .
- v_2 coefficient is directly sensitive to the rate at late times for expanding profiles.
- Assumptions : fully coherent jets, energy loss proceeds as off a single parton.



We demonstrate analytically that a medium evolution, which initially has a small coupling to jets, typically leads to a stronger jet azimuthal asymmetry at the same jet suppression factor.

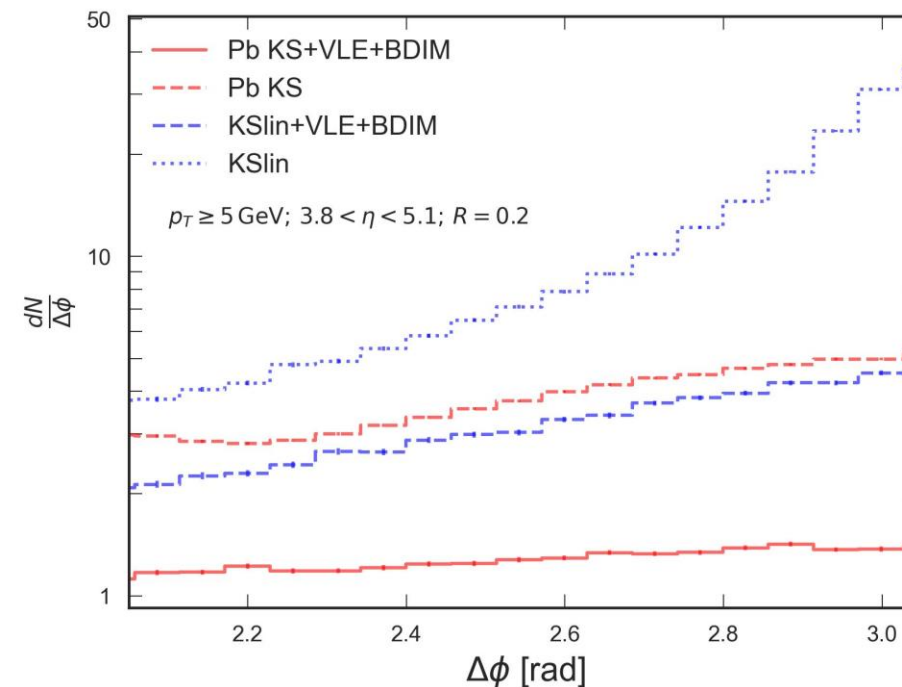
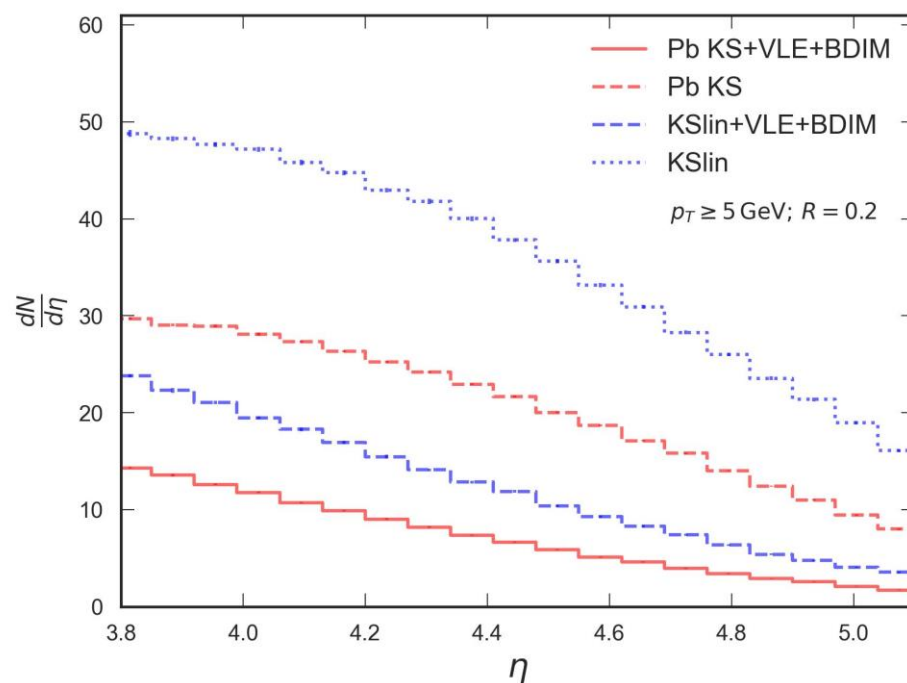
γ -jet at forward rapidity (FOCAL, ALICE)

- Currently, no LHC experiments explicitly studies **saturation physics**, to be observed in processes where longitudinal momentum of target probed at $x < 1E-5$.
- In the LHC jet kinematics, this corresponds to particle production in **a forward rapidity region**.
- Forthcoming **FoCal, ALICE** shall measure jets and photons at $3.4 < \eta < 5.8$.



In pPb collisions, strong saturation effects. True for A-A collisions in spite of medium modifications (VLE + BDIM) ?

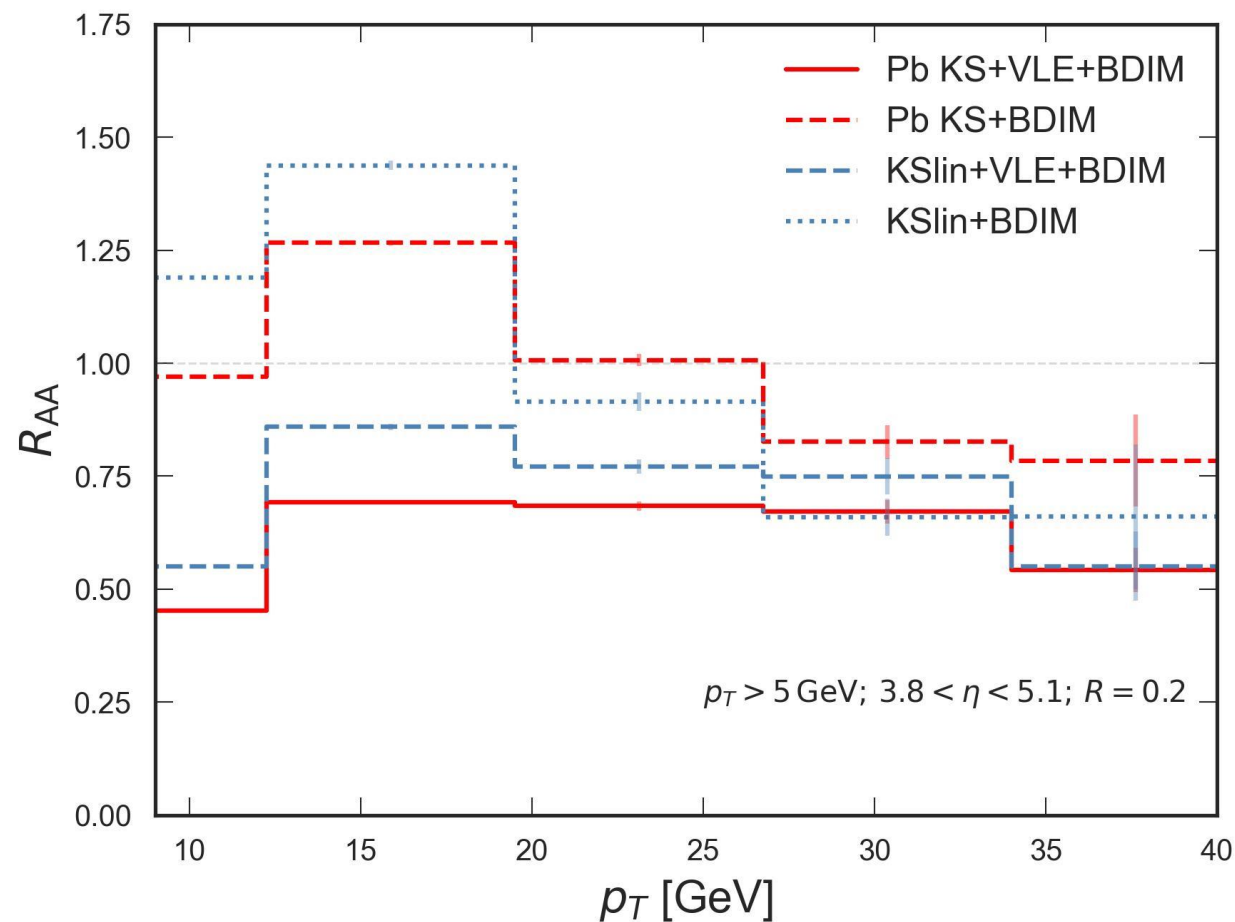
γ -jet rapidity and azimuthal correlations



- Gluon TMD configurations :
 - KSlin: No gluon saturation effects, solution of BFKL equation, DGLAP splitting functions.
 - Pb KS : Gluon saturation, gluon density is a solution of the BK equation, Sudakov effects, DGLAP splitting functions.

γ -jet suppression factor

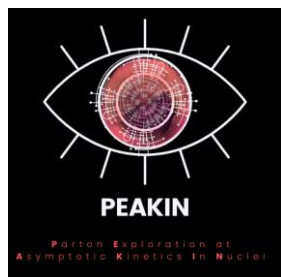
- Inclusion of VLEs strongly increases the jet-suppression at the low- p_T values by an approximate factor of 2 for $p_T \leq 20$ GeV.
- At low enough values of p_T (below $p_T = 20$ GeV, when VLEs are not included, and below $p_T = 35$ GeV, when they are included), a stronger suppression can be observed for the cases with the saturation.
- At high p_T , when VLEs are included, the case with the saturation shows a smaller suppression than the case without it.



VLE = Vacuum like emissions

Future challenges

- Novel incorporation of re-summation techniques for medium induced gluon emissions :
 - Finite size realistic medium effects
 - **Analytical results** : Effective designing of existing Parton shower MCs (faster, precise).
 - **New feature for HIC jet community** : Are multiple scattering important for parton showers ?
- **Future predictions for ALICE** : Description of photon-jet events in forward direction (FOCAL-range) via Monte-Carlo algorithms (saturation + quenching) . Dijets coming soon.
- Utilisation of theoretical developments for upcoming O-O runs (small collision systems) in 2025.



Non linear effects in asymmetric 3-jet (2+1) jet configurations (SPA, Iancu, Tasevsky (in preparation)).

*But this is a story (**EIC**) for another day !*