Heavy quark potential in anisotropic medium

Arghya Mukherjee

Ramakrishna Mission Residential College (Autonomous), Narendrapur, Kolkata

ATHIC 2025

In collaboration with: Dr. Margaret E. Carrington & Dr. Gabor Kunstatter



OUTLINE:

- \diamond Motivation
- $\diamond\,$ Overview of the formalism
- \diamond Results
- ♦ Conclusions

MOTIVATION:

 \diamond A model of anisotropic expansion: (PhysRevD.68.036004)

$$n_{\xi}(k) = n_{\rm eq} \left(\sqrt{k^2 + \xi_9 \left(\vec{k} \cdot \hat{n}_3 \right)^2} \right) \,. \tag{1}$$

where $n_{\rm eq}(k)$ represents the FD or BE distribution function.

- ◊ Have been used in several scenarios: (dispersion, instability, aHydro, HQ dynamics and so on).
- ◊ Our interest here is the HQ potential (real part) in anisotropic medium.

MOTIVATION:

$\diamond~{\rm HQ}$ potential:

$$V(\vec{r}) = -g^2 C_F \int \frac{d^3 p}{(2\pi)^3} \left(e^{i\vec{p}\cdot\vec{r}} - 1 \right) \mathcal{D}_{00}(p^0 \to 0, \vec{p})$$
(2)

where g is the strong coupling constant and $C_F = \frac{4}{3}$. \diamond In RTF:

$$\mathcal{D}_{\mu\nu}^{-1}(p_0, \vec{p}) = (\mathcal{D}_0)_{\mu\nu}^{-1}(p_0, \vec{p}) + \Pi_{\mu\nu}(p_0, \vec{p}).$$
(3)

♦ where $(\mathcal{D}_0)^{-1}_{\mu\nu}(p_0, \vec{p})$ is the inverse of the free gluon propagator

$$\mathcal{D}_{0}^{\mu\nu}(p_{0},\vec{p}) = -\frac{g^{\mu\nu}}{P^{2}} + (1-\chi)\frac{P^{\mu}P^{\nu}}{P^{4}}.$$
 (4)

♦ The anisotropic distribution function appears within $\Pi_{\mu\nu}$.

MOTIVATION:

♦ Phys. Lett. B 662, 37 (2008) :



◇ There is stronger attraction on distance scales on the order of the inverse Debye mass for quark pairs aligned along the direction of anisotropy than for transverse alignment.

FORMALISM:



FORMALISM:

♦ Generalization of the one parameter model:

$$n_{\xi}(k) = n_{\text{eq}} \left(\sqrt{k^2 + \xi_9 \left(\vec{k} \cdot \hat{n}_3 \right)^2} \right) = n_{\text{eq}} \left(k \sqrt{1 + \xi_9 \left(\vec{v} \cdot \hat{n}_3 \right)^2} \right).$$
(5)

can be achieved by defining:

$$n(\vec{k}) = C_{\xi} n(kH_{\xi}(\vec{v})) \tag{6}$$

 $\diamond\,$ We restrict ourselves to:

$$H_{\xi}^{2}(\vec{v}) = (1+\xi_{0}) + \xi_{2}(\vec{n}_{1}\cdot\vec{v})^{2} + \xi_{9}(\vec{n}_{3}\cdot\vec{v})^{2} + \xi_{6}(\vec{n}_{1}\cdot\vec{v})(\vec{n}_{3}\cdot\vec{v}) + \xi_{4}(\vec{n}_{1}\cdot\vec{v})^{4} + \xi_{8}(\vec{n}_{1}\cdot\vec{v})^{3}(\vec{n}_{3}\cdot\vec{v}) + \xi_{11}(\vec{n}_{1}\cdot\vec{v})^{2}(\vec{n}_{3}\cdot\vec{v})^{2} + \xi_{13}(\vec{n}_{1}\cdot\vec{v})(\vec{n}_{3}\cdot\vec{v})^{3} + \xi_{14}(\vec{n}_{3}\cdot\vec{v})^{4}.$$
(7)

7/9

RESULTS: PHYSREVC.110.035203



- ♦ The scaled potential for three values of θ at $\phi = \pi/4$ and three values of ϕ at $\theta = \pi/2$.
- $\diamond \ (\xi_0,\xi_2,\xi_4,\xi_6,\xi_8,\xi_9,\xi_{11},\xi_{13},\xi_{14}) = (5,8,37,-21,21,8,15,-28,40).$

CONCLUSIONS:

- ♦ General anisotropic scenario is expected to provide more realistic description of the evolving medium.
- ♦ Directional dependence of the potential is more complex.
- ◊ Imaginary part of the potential is an interesting future direction.

Thank You