

# **Emergent Cosmology from Quantum Gravity**

(In collaboration with P. Fischer, A. Jercher, T. Ladstätter, D. Oriti, A. Pithis)

## Luca Marchetti

Atlantic General Relativity 24 University of New Brunswick, Fredericton, NB 19 June 2024

Department of Mathematics and Statistics UNB Fredericton





Homogeneous Inhomogeneities: initial conditions

- ► Nature of dark matter?
- Nature of dark energy?
- Singularity resolution?

Inhomogeneities: dynamics

- ► Impact of singularity resolution on pert.?
- ▶ Is the evolution of pert. modified by QG effects?
- **>** ...

Homogeneous sector

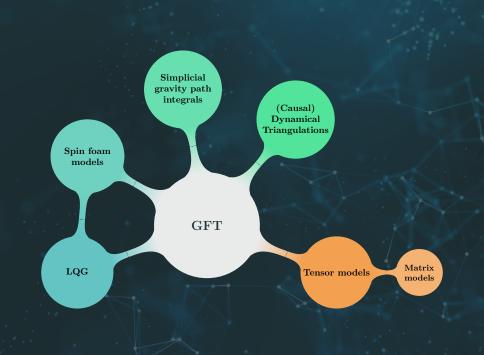
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- Quantum-geometric inflation?
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## The GFT approach to quantum gravity



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- ► Take seriously the idea of a microscopic structure of spacetime.
- ► Access to powerful field theoretic methods (Fock space, RG...)!

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### **Group Field Theory Quanta**

- ▶ GFT quanta are atoms of quantum 3-space, i.e. tetrahedra.
- ▶ Quantum numbers associated with fields discretized over a tetrahedron ( $g_a$  = gravitational,  $\Phi$  = scalar fields).
- ▶ GFT field domain is a group manifold *G*: not spacetime!

$$\varphi^{\dagger}(g_a, \mathbf{\Phi}) |0\rangle = \mathbf{g}_3$$

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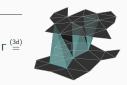
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## **Group Field Theory Processes**

- GFT Feynman diagrams Γ (QG processes) are associated with 4d triangulated (pseudo-)manifolds.
- $ightharpoonup Z_{GFT} = discrete matter-gravity path-integral.$
- ▶ GFT interactions are combinatorially non-local.



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Homogeneous Inhomogeneities: sector initial conditions ► Nature of dark matter?

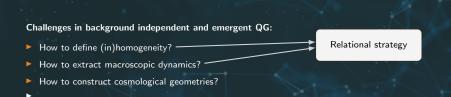
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## Challenges in background independent and emergent QG:

- How to define (in)homogeneity?
- How to extract macroscopic dynamics?
- How to construct cosmological geometries?

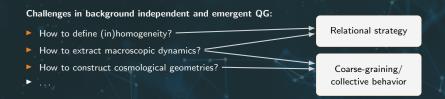


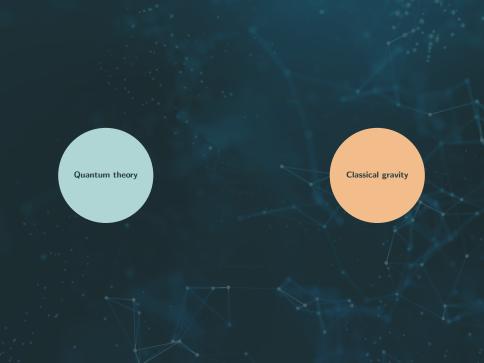
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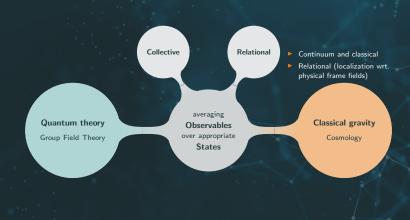


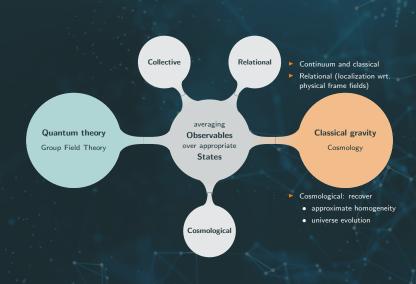






Collective ► Continuum and classical averaging Classical gravity Quantum theory Observables over appropriate Group Field Theory Cosmology States









notation: 
$$\varphi \cdot \psi = \int_{\Omega} d\Omega \ \varphi \psi$$

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- Look for states that can accommodate an infinite number of quanta

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Macroscopic dynamics = mean-field approximation

- GFT dynamics is captured by quantum equations of motion, or Schwinger-Dyson (SD) equations.
- Simplest SD equations are the averaged quantum equations of motion:  $\langle \delta S_{\text{GFT}}/\delta \hat{\varphi} \rangle_{\psi} = 0$ .

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$$\hat{N} = \varphi^{\dagger} \cdot \hat{\varphi}, \qquad \hat{\Phi} = \hat{\varphi}^{\dagger} \cdot (\Phi \hat{\varphi}), \qquad \hat{V} = \hat{\varphi}^{\dagger} \cdot V[\hat{\varphi}].$$

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Relational peaking

- Constructing relational observables in GFT is not easy (QFT with no continuum intuition).
- Relational localization implemented at an effective level on observable averages.

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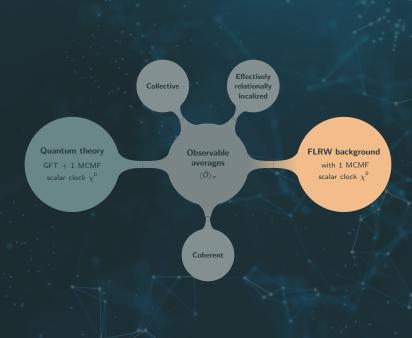
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$$\begin{split} \sigma_{\rm x} &= \left( {\rm fixed~peaking~function~} \eta_{\rm x} \right) \times \left( {\rm reduced~wavefunction~} \tilde{\sigma} \right), \\ {\rm E.g.:~} \chi^{\mu} \ {\rm frame~and~} |\sigma\rangle : \\ \mathcal{O}_{\sigma}({\rm x}) &\equiv \langle \hat{\mathcal{O}} \rangle_{\sigma_{\rm x}} \simeq \mathcal{O}[\tilde{\sigma}] |_{\chi^{\mu} = {\rm x}^{\mu}} \,, \qquad \langle \hat{\chi}^{\mu} \rangle_{\sigma_{\rm x}} \simeq {\rm x}^{\mu} \end{split}$$





# Effective FLRW cosmological dynamics

## Mean-field approximation

- ▶ Homogeneity:  $\tilde{\sigma}$  depends only on MCMF clock  $\chi^0$ .
- ▶ Isotropy:  $\tilde{\sigma}_{\upsilon} \equiv \rho_{\upsilon} e^{i\theta_{\upsilon}}$  ( $\upsilon_{\mathsf{EPRL}} \in \mathbb{N}/2$ ,  $\upsilon_{\mathsf{BC}} \in \mathbb{R}$ ).
- ► Mesoscopic regime: negligible interactions.

$$0 = \tilde{\sigma}_{\upsilon}^{\prime\prime} - 2i\tilde{\pi}_{0}\tilde{\sigma}_{\upsilon}^{\prime} - E_{\upsilon}^{2}\tilde{\sigma},$$

$$V_{\sigma}(x^{0}) = \sum_{\upsilon} V_{\upsilon} |\tilde{\sigma}_{\upsilon}|^{2} (x^{0}).$$

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#### Effective volume dynamics

$$\left(\frac{V_\sigma'}{3V_\sigma}\right)^2 = \left(\frac{2\sum_v V_v \rho_v \operatorname{sgn}(\rho_v') \sqrt{\mathcal{E}_v - Q_v^2/\rho_v^2 + \mu_v^2 \rho_v^2}}{3\sum_v V_v \rho_v^2}\right)^2, \quad \frac{V_\sigma''}{V_\sigma} = \frac{2\sum_v V_v \left[\mathcal{E}_v + 2\mu_v^2 \rho_v^2\right]}{\sum_v V_v \rho_v^2}$$

Classical limit

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Large number of quanta (large volume and late times)

- Volume quantum fluctuations under control.
- If  $\mu_v^2$  is mildly dependent on v (or one  $v_o$  is dominating) and equal to  $3\pi G$

$$(V_{\sigma}'/3V_{\sigma})^2 \simeq 4\pi G/3 \longrightarrow \mathsf{flat}\;\mathsf{FLRW}$$

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  - $(V'_{\sigma}/3V_{\sigma})^2 \simeq 4\pi G/3 \longrightarrow \text{flat FLRW}$

- ✓ Matching with GR (in harmonic gauge)!
- $\checkmark x^0 = \langle \hat{\chi}^0 \rangle_{\sigma_{..0}}$ , clock quantum fluct.  $\simeq 0$ .
- $\label{eq:continuous} \checkmark \ \langle \hat{\Pi}^0 \rangle_{\sigma_{\chi^0}} = \langle \hat{H}_\sigma \rangle_{\sigma_{\chi^0}} \mbox{ (higher moments} \simeq 0).$

Effective relational framework reliable!

## Mean-field approximation

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Singularity res. into quantum bounce?

- $\triangleright$   $x^0$  may not coincide with  $\langle \hat{\chi}^0 \rangle_{\sigma_{x,0}}$  anymore!
- Clock quantum fluctuations may be large!
- $\wedge$   $\langle \hat{\Pi}^0 \rangle_{\sigma_{\cdot,0}} \neq \langle \hat{H}_{\sigma} \rangle_{\sigma_{\cdot,0}}$  (higher moments  $\neq$  0).

Effective rel. framework may break down!

# Emergent inflation and phantom dark energys

Tensor (modulus) Cellular (phase)

 $\mbox{notation:} \ \varphi \cdot \psi = \!\! \int_{\Omega} \! \mathrm{d}\Omega \, \varphi \, \psi$  Cellular (phase)

### Tensor (modulus)

$$\mathsf{Tr}^{(m)}_{\mathcal{V}_{\gamma_I}}[\varphi,\bar{\varphi}] \sim \mathcal{V}^{(m)}_{\gamma_I} \cdot \bar{\varphi}^{(l+1)/2} \cdot \varphi^{(l+1)/2}$$

- ► Highly symmetric, studied in renormalization.
- $\blacktriangleright$  Modulus-only dependence after  $\sigma\textsc{-isotropy}.$

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# $\text{notation: } \varphi \, \cdot \, \psi \, = \!\! \int_{\Omega} \! \mathrm{d}\Omega \, \varphi \, \psi$

#### Cellular (phase)

$$\mathrm{Tr}_{\mathcal{V}_{\gamma_{I}}}^{(p)}[\varphi,\bar{\varphi}] \sim \mathcal{V}_{\gamma_{I}}^{(p)} \cdot \varphi^{I+1}$$

- Admit a more clear geometric interpretation.
- Modulus&phase dependence after  $\sigma$ -isotropy.

$$\mathsf{Tr}_{\mathcal{V}_{\gamma_l}}^{(m)}[\varphi,\bar{\varphi}] \sim \mathcal{V}_{\gamma_l}^{(m)} \cdot \bar{\varphi}^{(l+1)/2} \cdot \varphi^{(l+1)/2}$$

- ► Highly symmetric, studied in renormalization.
- $\blacktriangleright \ \ \mathsf{Modulus}\text{-}\mathsf{only}\ \mathsf{dependence}\ \mathsf{after}\ \sigma\text{-}\mathsf{isotropy}.$

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#### Phantom Dark Energy

lacktriangledown Consider l=5 modulus interactions at very late times, but include a subdominant spin  $\upsilon'$ :

$$w = 3 - 2(VV'')/(V')^2 \simeq -1 - b/V, \qquad b > 0.$$

Universe effectively dominated by (non-pathologic) emergent phantom dark energy.

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**Emergent Inflation** 

PRELIMINARY

$$0 = \tilde{\sigma}_{\upsilon}^{\prime\prime} - 2i\tilde{\pi}_{0}\tilde{\sigma}_{\upsilon}^{\prime} - E_{\upsilon}^{2}\tilde{\sigma}_{\upsilon} - \lambda_{\upsilon}\bar{\tilde{\sigma}}_{\upsilon}^{\prime}$$

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#### Slow-roll phase

- ▶ Initial conditions: I = 5, Re[ $\lambda_v \bar{\tilde{\sigma}}_v^I$ ] close to a maximum,  $\theta_n$  determines slow-roll.
- Long-lasting quasi-deSitter phase!

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#### Graceful exit and fast oscillations

- ► Natural slow-roll breakdown: fast oscillations.
- Interactions washed away on average.
- Graceful exit: matter (clock) dominated phase!



# Cosmic inhomogeneities from quantum gravity entanglement

# Setting

#### Classical

- 4 MCMF reference fields  $(\chi^0, \chi^i)$ ,
- $\begin{tabular}{ll} $\mathbf{I}$ MCMF matter field $\phi$ dominating the energy-momentum budget and slightly relationally inhomogeneous wrt. $\chi^i$. \end{tabular}$

#### Quantum

Beyond condensates: time- and spacelike tetrahedra.

Inhomogeneities = Quantum Entanglement

$$|\Delta;x\rangle = \mathcal{N}_{\Delta} e^{\hat{\sigma}\otimes\mathbb{I}_{-}+\mathbb{I}_{+}\otimes\hat{\tau}+\widehat{\delta\Phi}\otimes\mathbb{I}_{-}+\widehat{\delta\Psi}+\mathbb{I}_{+}\otimes\widehat{\delta\Xi}}\,|0\rangle\,.$$

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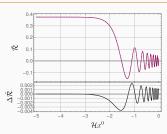
$$|\Delta;x\rangle = \mathcal{N}_{\Delta} e^{\hat{\sigma}\otimes\mathbb{I}_{-}+\mathbb{I}_{+}\otimes\hat{\tau}+\widehat{\delta\Phi}\otimes\mathbb{I}_{-}+\widehat{\delta\Psi}+\mathbb{I}_{+}\otimes\widehat{\delta\Xi}} |0\rangle \ .$$

#### Classical dynamics with trans-Planckian QG effects

 $\begin{tabular}{ll} \hline & Matter $\delta\phi_{\sf GFT}$ and "curvature-like" (isotropic) \\ & {\sf pert.} \ \tilde{\mathcal{R}}$ emerge from to two-body relational \\ & {\sf nearest-neighbor QG}$ correlations $(\delta\widehat{\Phi}, \widehat{\delta\Psi}, \widehat{\delta\Xi})$. \\ \end{tabular}$ 

$$\begin{split} \delta\phi_{\mathsf{GFT}}^{\prime\prime\prime} + k^2 a^4 \delta\phi_{\mathsf{GFT}} &= \left(\frac{a^2 k}{M_{\mathrm{pl}}}\right) j_{\phi}[\bar{\phi}] \,, \\ \tilde{\mathcal{R}}_{\mathsf{GFT}}^{\prime\prime\prime} + k^2 a^4 \tilde{\mathcal{R}}_{\mathsf{GFT}} &= \left(\frac{a^2 k}{M_{\mathrm{pl}}}\right) j_{\bar{\mathcal{R}}}[\bar{\phi}] \,, \end{split}$$

- Trans-Planckian QG corrections to the dynamics of scalar isotropic perturbations.
- ✓ Remarkable agreement with GR at larger scales.



Top:  $\tilde{\mathcal{R}}_{GFT}$  (blue) and  $\tilde{\mathcal{R}}_{GR}$  (dashed red) for  $k/M_{Pl}=10^2$ . Bottom: their difference  $\Delta \tilde{\mathcal{R}}$ .

