

The role of Entanglement and CCR in QED scattering Processes

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*M. Blasone, G. Lambiase and B. M., “Entanglement distribution in Bhabha scattering with entangled spectator particle”, Phys. Rev. D 109 (2024).

†M. Blasone, S. De Siena, G. Lambiase, C. Matrella and B. M., “Complete complementarity relations in tree level QED processes”, [arXiv:2402.09195 [quant-ph]].

Summary

1. Entanglement in QED scattering processes
2. Complete Complementarity Relations in QED scattering processes

Motivations

- Investigation of fundamental properties of (elementary) particles via quantum correlations[‡]
- Entanglement in relativistic systems
- A possible resource for quantum information
- Entanglement in high energy processes as a probe for new physics beyond Standard Model

[‡]In this conference see related talks by **P. Lamba, E.M. Sessolo, K. Sakurai**

Entanglement in QED scattering processes

Entanglement generation in QED processes*

- How entanglement (in helicity) is generated at the fundamental level in QED

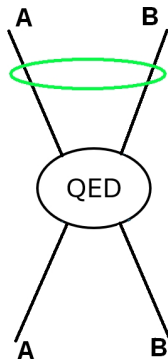
$$e^-e^+ \rightarrow \mu^-\mu^+$$

$$e^-e^+ \rightarrow e^-e^+$$

$$e^-\mu^- \rightarrow e^-\mu^-$$

$$e^-\gamma \rightarrow e^-\gamma$$

...



*A. Cervera-Lierta et al., SciPost Phys. (2017).
S. Fedida, A. Serafini, Phys.Rev.D (2023).

Entanglement Generation in QED processes

Concurrence:

$$C(\rho) = \max(0, \lambda_1 - \lambda_2 - \lambda_3 - \lambda_4)$$

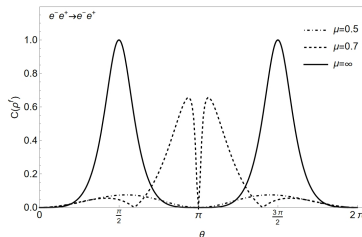
$$\mathcal{R} = \rho_{s_1 s_2} \tilde{\rho}_{s_1 s_2}$$

$$\tilde{\rho}_{s_1 s_2} = (\sigma_y \otimes \sigma_y) \rho_{s_1 s_2}^* (\sigma_y \otimes \sigma_y)$$

Initial and final helicity states:

$$|i\rangle = |R\rangle_A \otimes |L\rangle_B$$

$$|f\rangle = \sum_{r,s=R,L} \mathcal{M}(RL; rs) |r\rangle_A |s\rangle_B$$

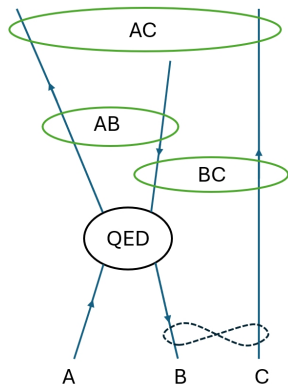


QED processes with spectator particles[†]

Initial and final states

$$|i\rangle = |R\rangle_A \otimes \left(\cos \eta |R\rangle_B |R\rangle_C + e^{i\beta} \sin \eta |L\rangle_B |L\rangle_C \right)$$

$$|f\rangle = \sum_{r,s=R,L} \left[\cos \eta \mathcal{M}(RR; rs) |r\rangle_A |s\rangle_B |R\rangle_C + e^{i\beta} \sin \eta \mathcal{M}(RL; rs) |r\rangle_A |s\rangle_B |L\rangle_C \right]$$



[†]J.B. Araujo et al., Phys. Rev.D (2019).

M. Blasone, G. Lambiase and B. M., Phys. Rev. D 109 (2024).

Entanglement distribution in Bhabha scattering [‡]

- In the same framework we have studied the generation and distribution of the entanglement in Bhabha scattering for the three bipartite channel AB, AC and BC.
- For incoming momenta of the order of the mass, the entanglement has a non trivial distribution in the three output channels.
- In the relativistic regime, analytic expressions for the concurrences:

$$\lim_{\mu \rightarrow \infty} C(\rho_{AB}^f) = \frac{2 \sin^2 \eta \sin^4(\theta/2) \cos^4(\theta/2)}{1 - (1 - \frac{1}{8} \sin^2 \theta) \sin^2 \theta \sin^2 \eta}$$

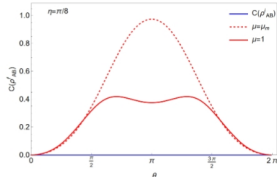
$$\lim_{\mu \rightarrow \infty} C(\rho_{AC}^f) = \frac{\sin(2\eta) \sin^4(\theta/2)}{1 - (1 - \frac{1}{8} \sin^2 \theta) \sin^2 \theta \sin^2 \eta}$$

$$\lim_{\mu \rightarrow \infty} C(\rho_{BC}^f) = \frac{\sin(2\eta) \cos^4(\theta/2)}{1 - (1 - \frac{1}{8} \sin^2 \theta) \sin^2 \theta \sin^2 \eta}$$

[‡]M. Blasone, G. Lambiase and B. M., Phys. Rev. D 109 (2024).

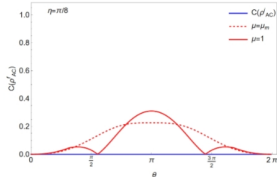
Entanglement distribution ($\eta = \pi/8$)

AB



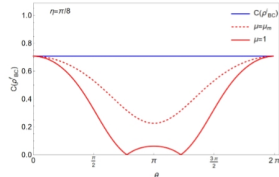
(a)

AC

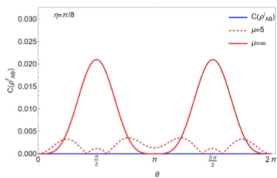


(b)

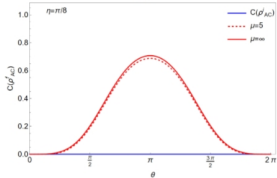
BC



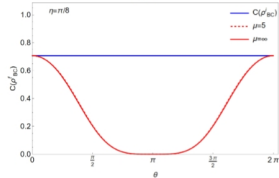
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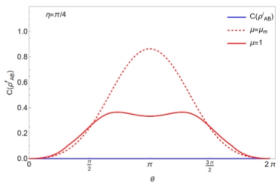


(f)

$$\mu = \frac{|\vec{p}|}{m_e}, \quad \mu_m = \frac{1}{2} \sqrt{-3 + \sqrt{17}}$$

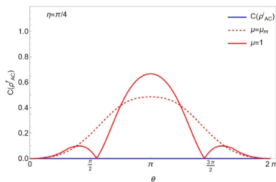
Entanglement distribution ($\eta = \pi/4$)

AB



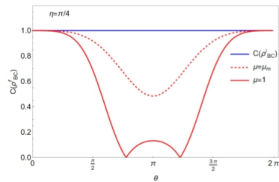
(a)

AC

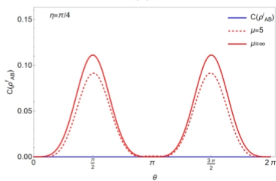


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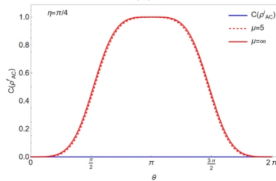
BC



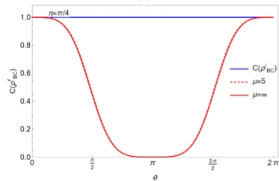
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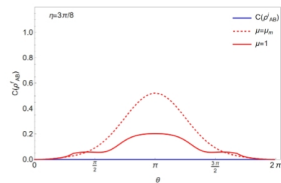
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Entanglement distribution ($\eta = 3\pi/8$)

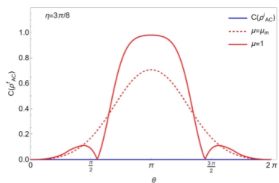
AB

AC

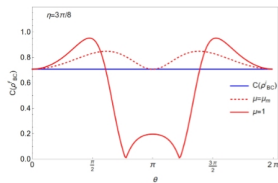
BC



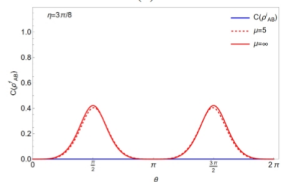
(a)



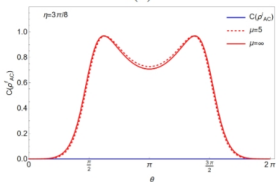
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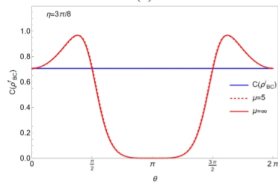
(c)



(d)



(e)



(f)

- The distribution and generation of entanglement tend to concentrate in some bipartitions.
- Complete transfer from BC channel to the AC channel in $\theta = \pi$ in relativistic regim: interaction as a *quantum gate* \Rightarrow possible application in quantum information protocols.
- For $\frac{\pi}{4} < \eta < \frac{3\pi}{4}$ the interaction generates entanglement in the BC channel, where it was present before the interaction.

Complete Complementarity Relations in QED scattering processes

Complementarity

Complementarity[§]: the two aspect of duality, corpuscular and undulatory, are equally real but mutually exclusive.

The first quantitative version[¶] of the wave-particle duality was summarized by a simple complementarity relation

$$P^2 + V^2 \leq 1.$$

where P is the **predictability**, a measure of path information and V is the **visibility** of the interference pattern.

[§]N. Bohr, *The quantum postulate and the recent development of atomic theory*, Nature (1928)

[¶]W. K. Wootters, W. H. Zurek, Phys. Rev. D. (1979),
B-G. Englert, Phys. Rev. Lett. (1996)

Complete Complementarity Relations

For bipartite states[‡] we have to consider a **triatlity relation** formed by two quantities generating *local*, single-partite realities and the correlation between subsystems, which generates an exclusive bipartite *non-local* reality:

$$P_k^2 + V_k^2 + C^2 = 1, \quad k = A, B$$

where P_k and V_k are the predictability and visibility for the single-partite systems. C is a measure of entanglement.

[‡]M. Jakob, J. A. Bergou, Opt. Commun. (2010)

CCR for general bipartite states

We can write a general bipartite state in terms of two qubits as

$$|\Theta\rangle = a|00\rangle + b|01\rangle + c|10\rangle + d|11\rangle,$$

and express the three terms of the previous triality relation as:

$$C(\Theta) = 2|ad - bc|$$

$$P_A = |(|c|^2 + |d|^2) - (|a|^2 + |b|^2)|$$

$$P_B = |(|b|^2 + |d|^2) - (|a|^2 + |c|^2)|$$

$$V_A = 2|ac^* + bd^*|$$

$$V_B = 2|ab^* + cd^*|$$

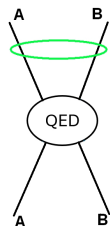
CCR for Bhabha scattering: initial states

We have analyzed C , V and P before and after the QED interaction for three different states:

$$|i\rangle_I = |R\rangle_A |L\rangle_B$$

$$|i\rangle_{II} = (\cos \alpha |R\rangle_A + e^{i\xi} \sin \alpha |L\rangle_A) \otimes (\cos \beta |R\rangle_B + e^{i\eta} \sin \beta |L\rangle_B)$$

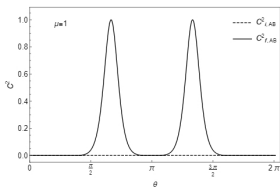
$$|i\rangle_{III} = \cos \alpha |R\rangle_A |R\rangle_B + e^{i\xi} \sin \alpha \cos \beta |R\rangle_A |L\rangle_B \\ + e^{i\eta} \sin \alpha \sin \beta \cos \chi |L\rangle_A |R\rangle_B + e^{i\tau} \sin \alpha \sin \beta \sin \chi |L\rangle_A |L\rangle_B$$



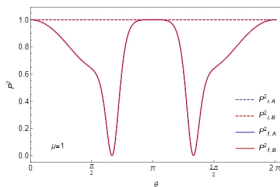
‡M. Blasone, S. De Siena, G. Lambiase, C. Matrella and B. M., [arXiv:2402.09195 [quant-ph]].

Examples: initial factorized state $|i\rangle_I$

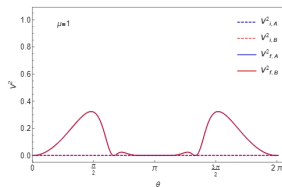
$$P_i = 1, \quad V_i = 0, \quad C = 0.$$



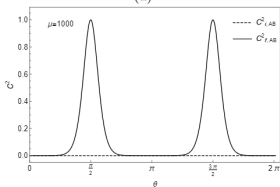
(a)



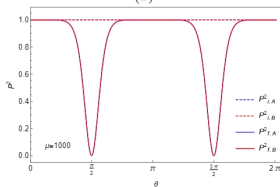
(b)



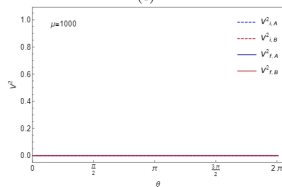
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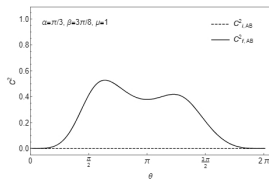
(e)



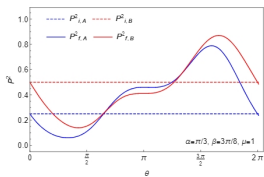
(f)

Examples: initial factorized state $|i\rangle_{II}$

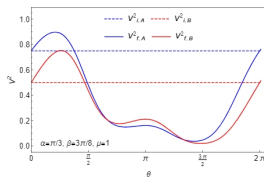
$$P_i \neq 0, \quad V_i \neq 0, \quad C = 0.$$



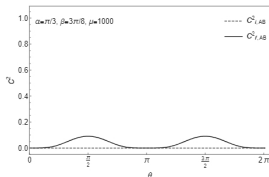
(a)



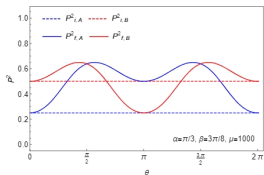
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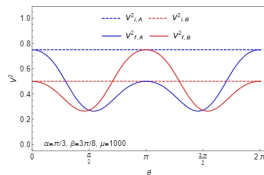
(c)



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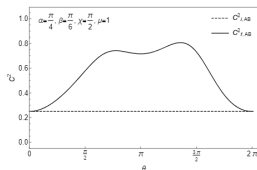
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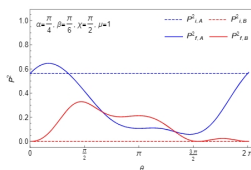
(f)

Examples: initial entangled states $|i\rangle_{III}$

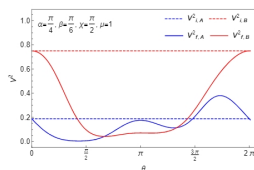
Entanglophilus regime



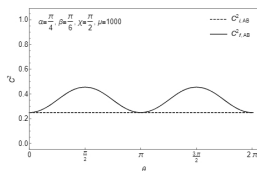
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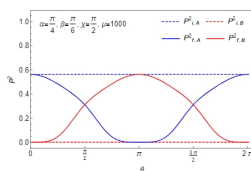
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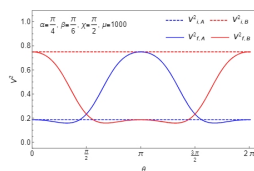
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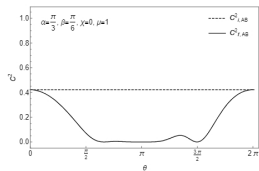
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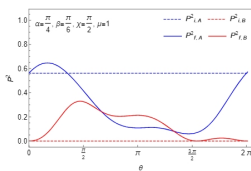
(f)

Examples: initial entangled states III

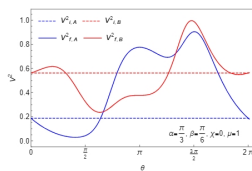
Entanglophobous regime



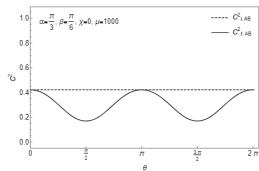
(a)



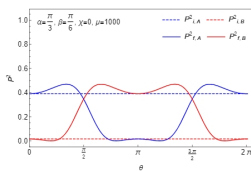
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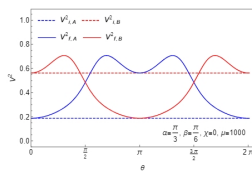
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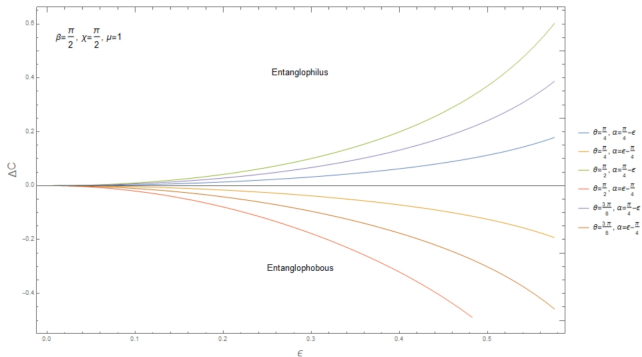


(f)

Entanglophilus, Entanglophobous and Mixed regimes

Entanglophilus, Entanglophobous from Bell states

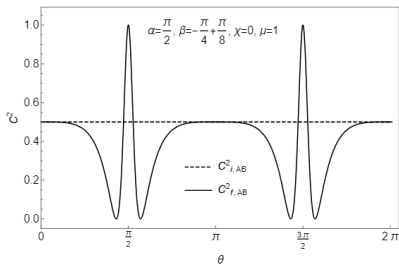
$$\Delta C \equiv (C_f - C_i)/C_i \quad \Phi_{\alpha}^{\pm} = \cos \alpha |RR\rangle \pm \sin \alpha |LL\rangle \quad \Phi^{\pm} = \frac{1}{\sqrt{2}}(|RR\rangle \pm |LL\rangle)$$



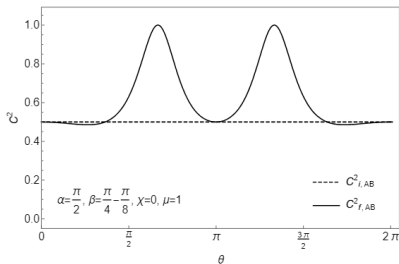
Entanglophilus, Entanglophobous and Mixed regimes

Mixed regimes from Bell states

$$\Psi_{\alpha}^{\pm} = \cos \alpha |RL\rangle \pm \sin \alpha |LR\rangle \quad \Psi^{\pm} = \frac{1}{\sqrt{2}}(|RL\rangle \pm |LR\rangle)$$



(a)



(b)

Maximal entanglement conservation

$$P_i = 0 \quad V_i = 0 \quad C = 1$$

- Processes preserving maximal entanglement

- $e^-e^+ \rightarrow e^-e^+$
- $e^-e^- \rightarrow e^-e^-$
- $e^-e^+ \rightarrow \mu^-\mu^+$
- $e^-\mu^- \rightarrow e^-\mu^-$

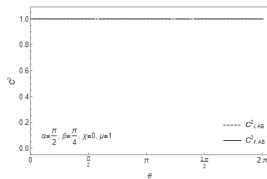
An Example: Bhabha scattering

$$\Phi^+ \rightarrow \Phi^+$$

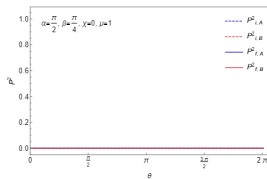
$$\Phi^- \rightarrow \cos s_1 \Phi^- + \sin s_1 \Psi^+$$

$$\Psi^+ \rightarrow \cos s_2 \Phi^- + \sin s_2 \Psi^+$$

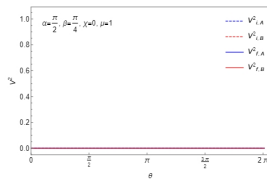
$$\Psi^- \rightarrow \Psi^-$$



(a)



(b)



(c)

QED scattering with photons

- $e^-e^+ \rightarrow 2\gamma$: maximal entanglement is not conserved for all the initial maximally entangled states
- $e^-\gamma \rightarrow e^-\gamma$: maximal entanglement is never achieved

	$e^-e^+ \rightarrow \gamma\gamma$		*		Compton	
In. State	Fin. State	Conc.	*	In. State	Fin. State	Conc.
Φ^+	Φ^-	1	*	Φ^+	GC	< 1
Φ^-	$\cos r \Phi^+ + \sin r \Psi^+$	$\sin(2r)$	*	Φ^-	GC	< 1
Ψ^+	Ψ^+	1	*	Ψ^+	GC	< 1
Ψ^-	Ψ^-	1	*	Ψ^-	GC	< 1

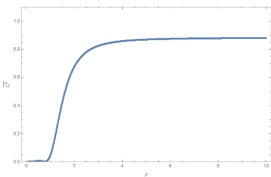
Discussion of results

- Can we consider and quantify the QED scattering processes as a good resource in order to generate entangled particles?
- How can we interpret these results?

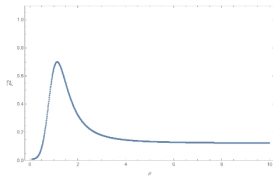
Entanglement resource

Weighted CCR averages for initial state $|RL\rangle$

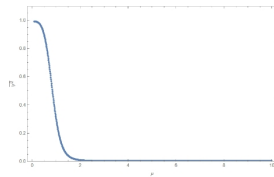
$$\frac{d\sigma}{d\Omega} = \frac{|\mathcal{M}|^2}{64\pi^2 E_{CM}^2} \quad \overline{Q^2} = \frac{1}{\mathbf{N}} \sum_{rs} \int_{\mathcal{D}} |M(a, b; r, s)|^2 Q^2(\theta) d\theta \quad Q = P, V, C$$



(a)

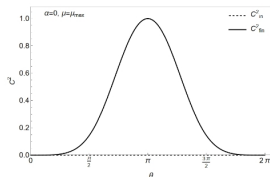


(b)

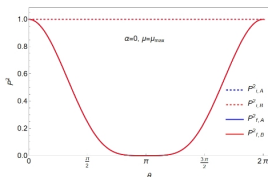


(c)

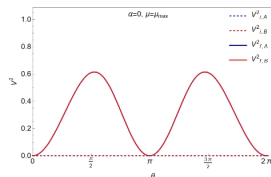
CCR vs Probabilities for initial state $|RR\rangle$



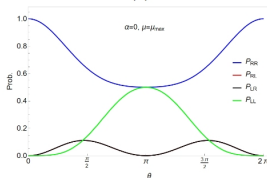
(a)



(b)

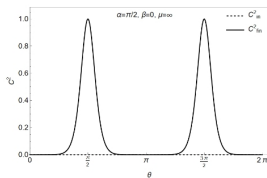


(c)

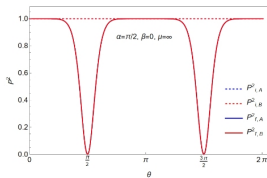


(d)

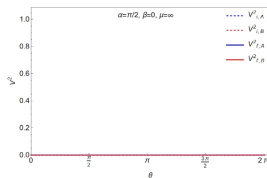
CCR vs Probabilities for initial state $|RL\rangle$



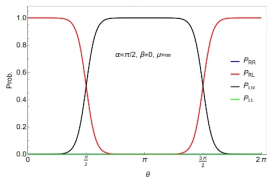
(a)



(b)



(c)



(d)

Conclusions

- We have studied entanglement generation and distribution in QED processes with and without entangled spectator particle
- With entangled spectator particle we observe a non trivial distribution of initial entanglement in the various bipartitions
- For some values of parameters we find a complete transfer of entanglement from one channel to another (quantum gate)
- By means of CCR we have widely characterized entanglement generation in elementary scattering process at tree level.
- We find a regime in which entanglement increases with respect to its initial value and another in which it decreases
- Remarkably, the maximal entanglement (Bell state) is totally preserved for processes that involve 1/2-spin massive particle

Outlook

- Extension of the analysis to other interactions
- 1 loop correction**
- Analysis of the same processes in other reference frames
- Possible implementations of our results in quantum information protocols
- Possible identification of new physics BSM

K. Kowalska and E. M. Sessolo, Entanglement in flavored scalar scattering, JHEP **07 (2024), 156.