

LHCb

Wat doen wij?



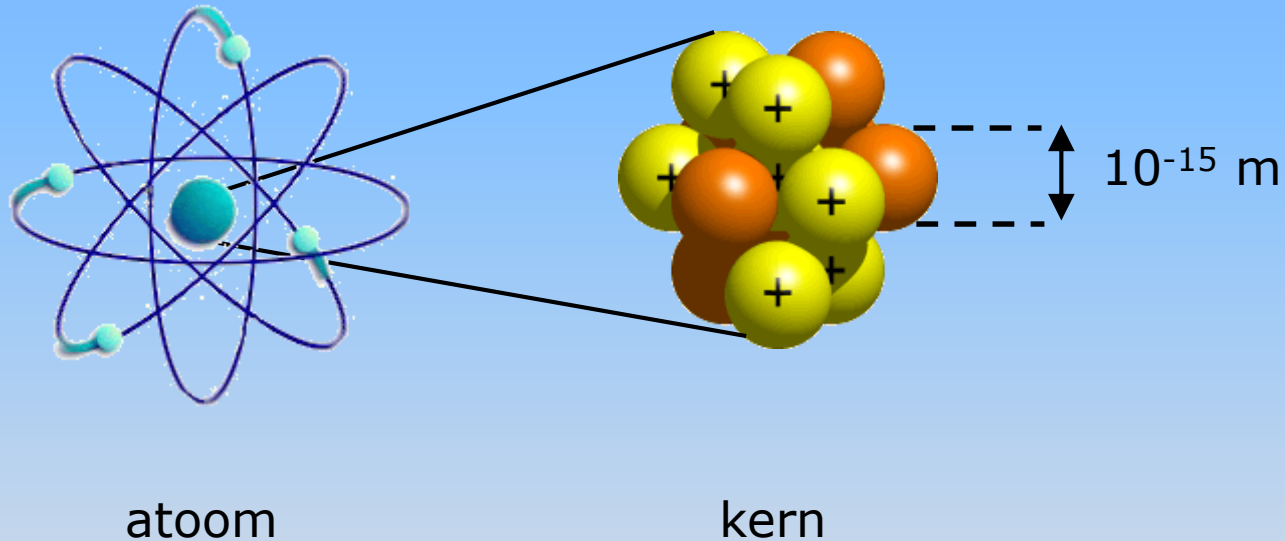
Niels Tuning – 9 okt 2024

LHCb

- Waarom deeltjesfysica?
- Waarom LHCb?
- Resultaten
- Higgs en LHCb

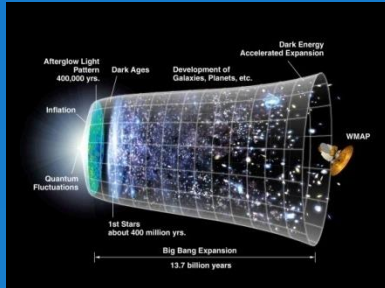
Deeltjesfysica

Bestudeert de natuur op afstanden $< 10^{-15}$ m



Quantum theorie beschrijft alle metingen tot 10^{-18} m
(Ter vergelijik: 10^{+18} m = 100 lichtjaar)

Machten van tien ...



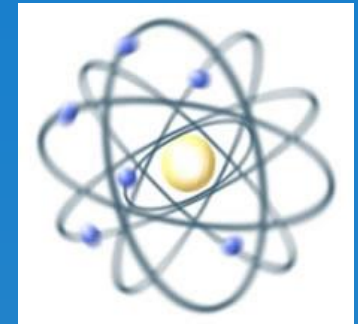
Heelal
 10^{26} m

Spin
 10^{-2} m



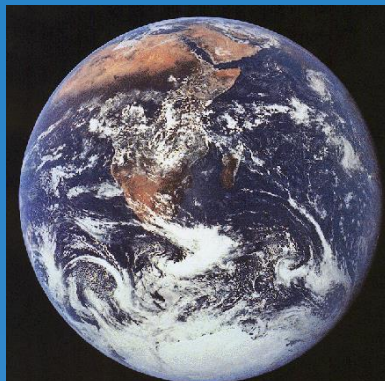
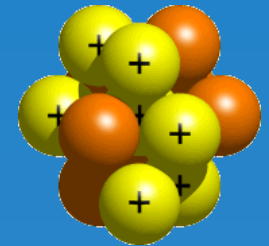
Melkweg
 10^{21} m

Atoom
 10^{-10} m



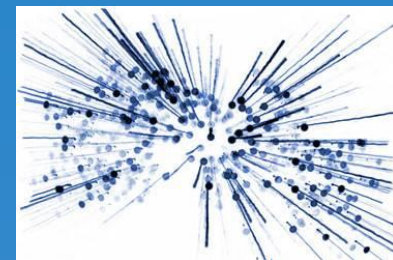
Zonnestelsel
 10^{13} m

Kern
 10^{-15} m

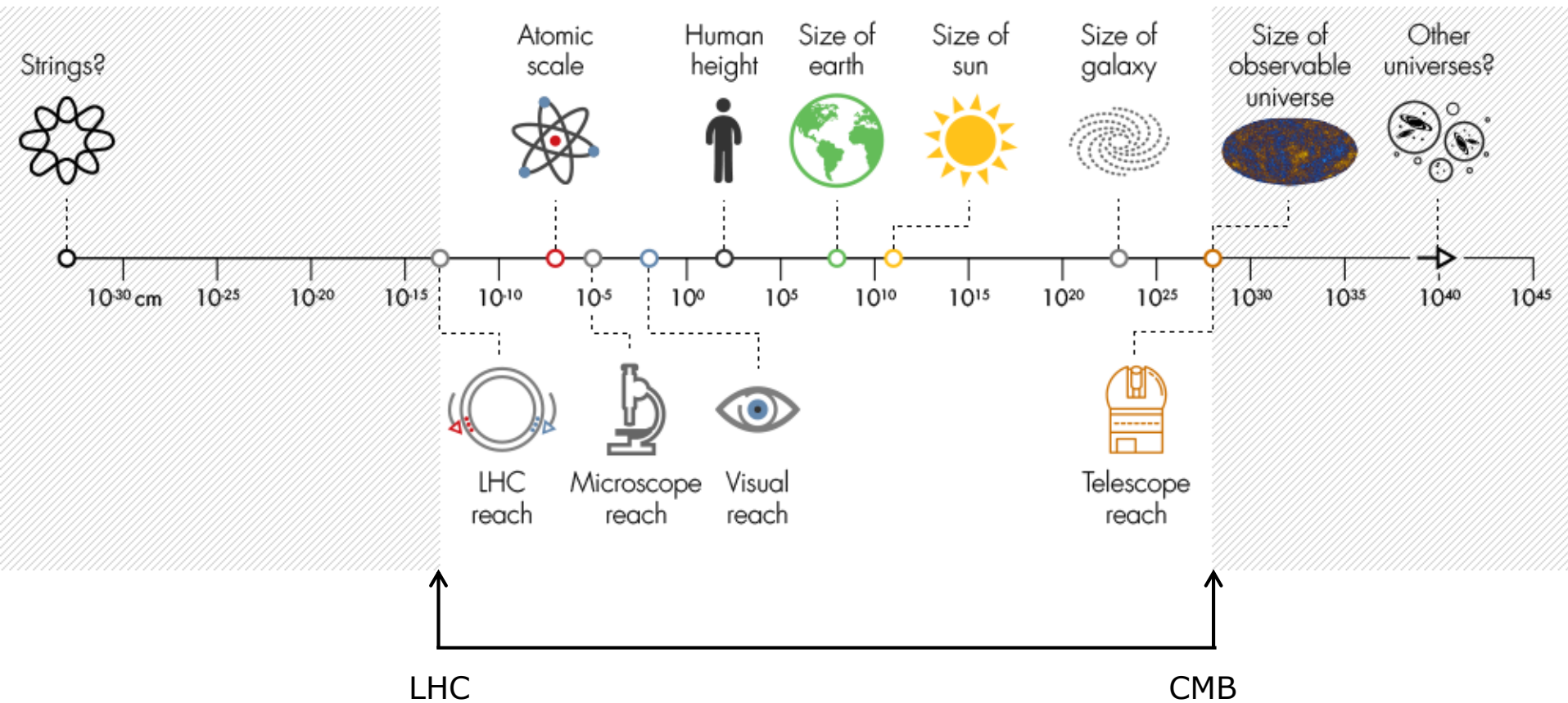


Aarde
 10^7 m

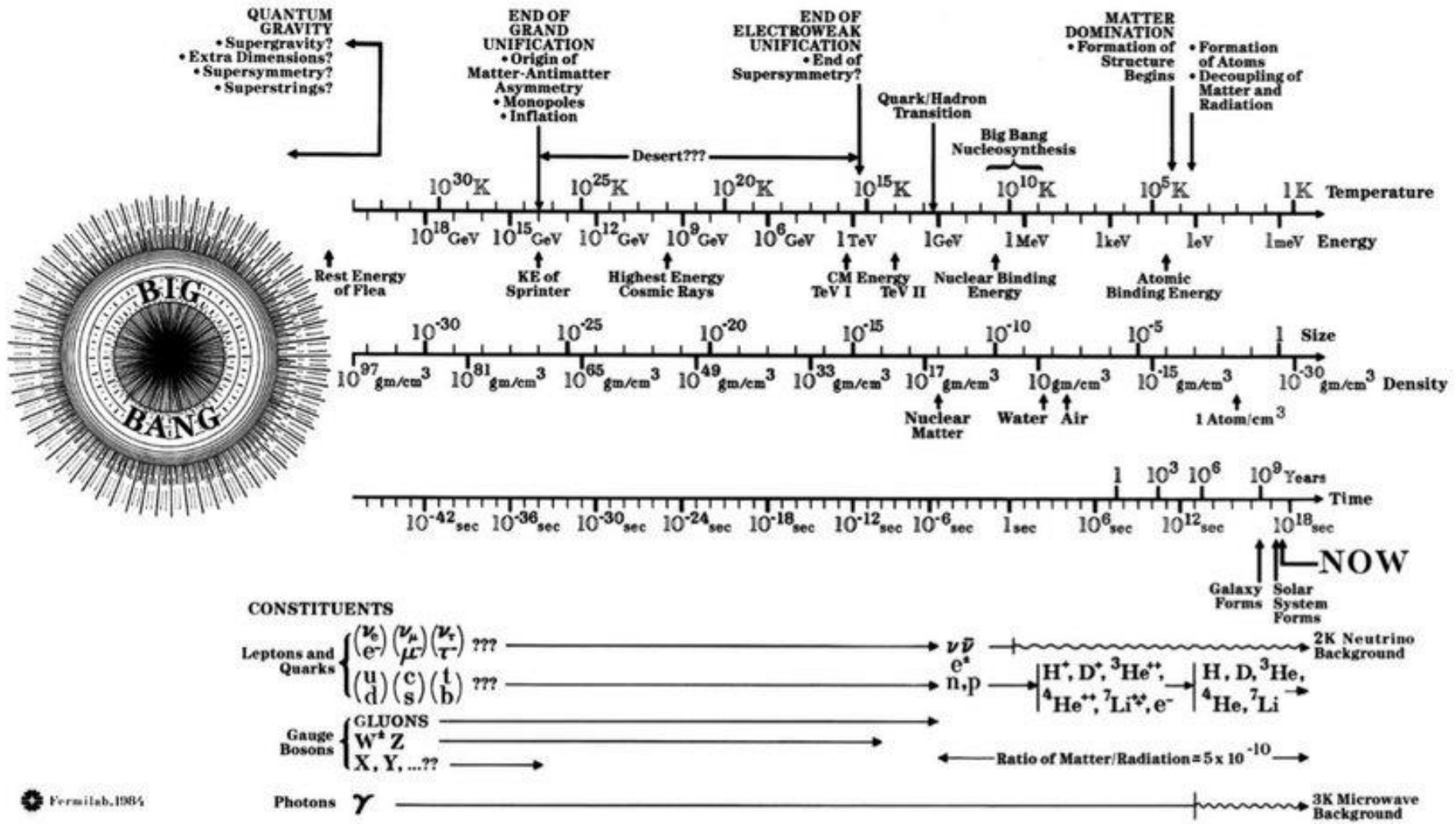
Botsingen
 10^{-18} m



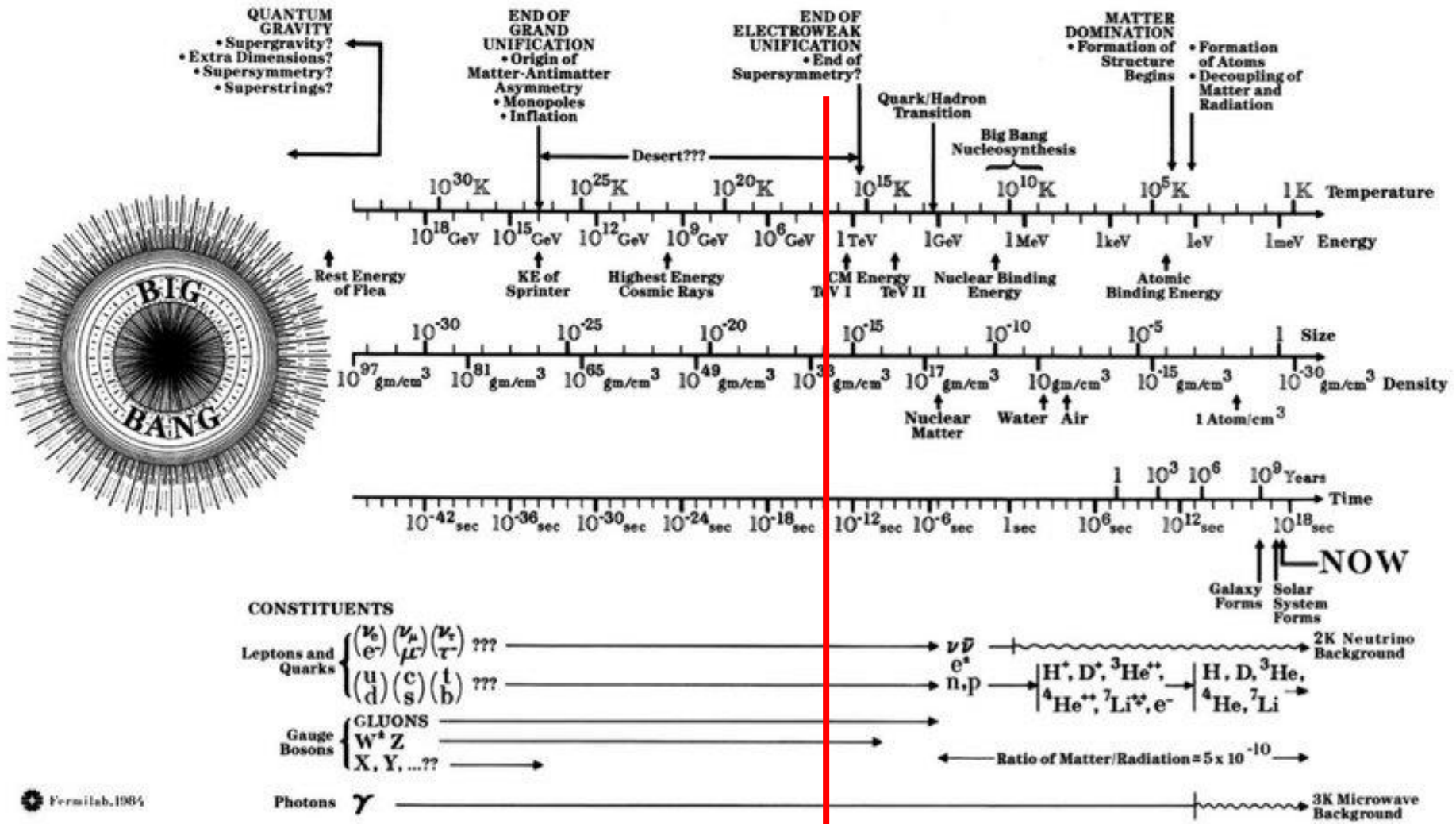
Schaal



Complete History of the Universe



Complete History of the Universe



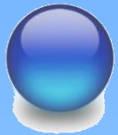
U bent hier: \blacklozenge $10^{-16} \text{ m}, 10^{-14} \text{ s}, 10^{16} \text{ K}$

De stand van zaken in 2022

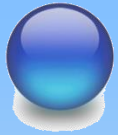


[http:// pdg.lbl.gov](http://pdg.lbl.gov)

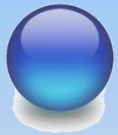
De elementaire deeltjes



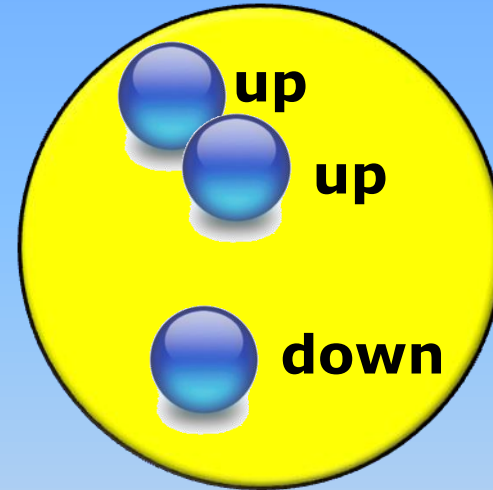
up



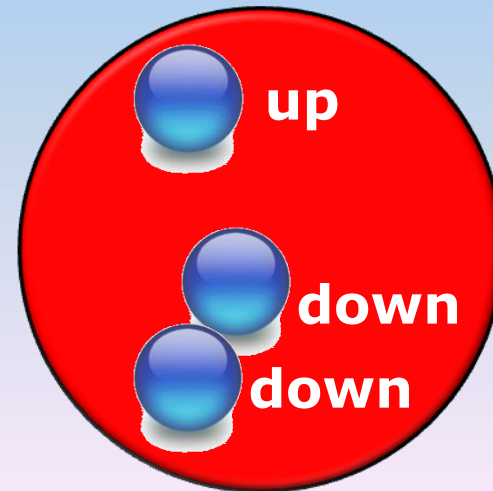
down



elektron

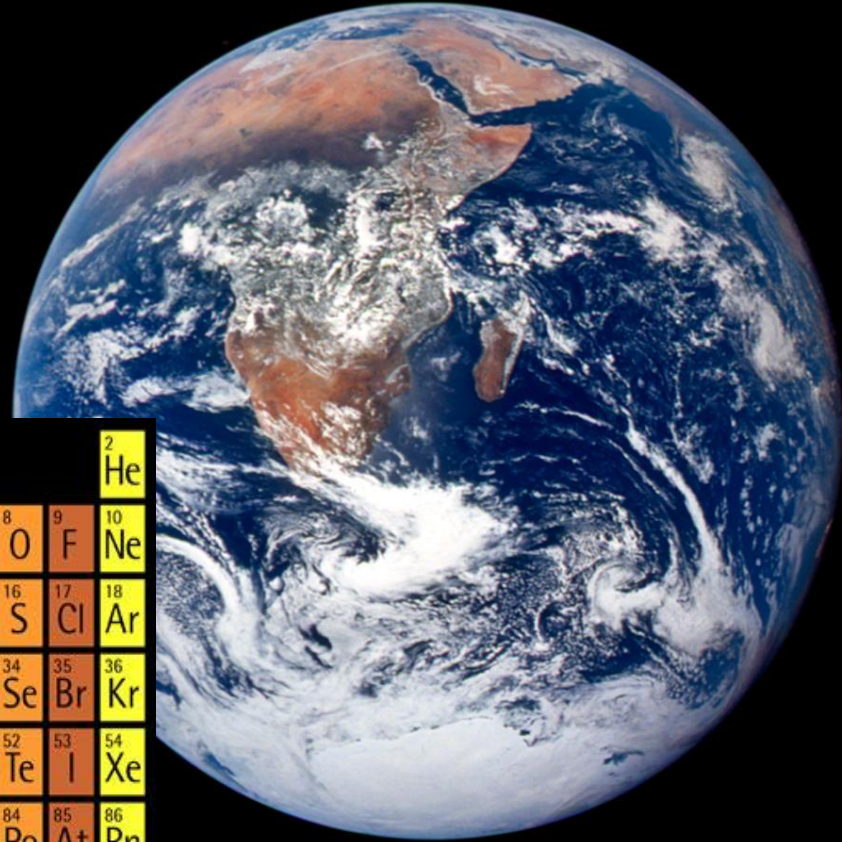
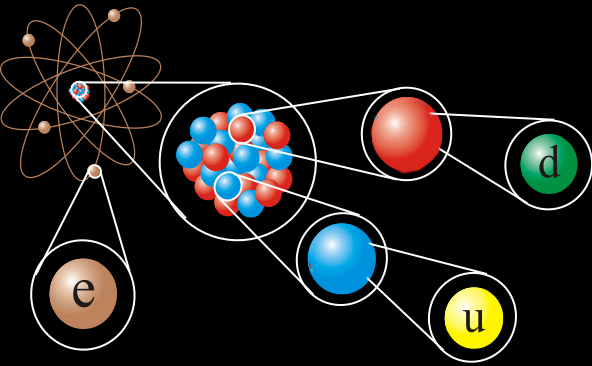


Proton



Neutron

Wat kan je maken van deze 3 bouwstenen?



periodiek systeem
van Mendeleev

1 H																	2 He
3 Li	4 Be											5 B	6 C	7 N	8 O	9 F	10 Ne
11 Na	12 Mg											13 Al	14 Si	15 P	16 S	17 Cl	18 Ar
19 K	20 Ca	21 Sc	22 Ti	23 V	24 Cr	25 Mn	26 Fe	27 Co	28 Ni	29 Cu	30 Zn	31 Ga	32 Ge	33 As	34 Se	35 Br	36 Kr
37 Rb	38 Sr	39 Y	40 Zr	41 Nb	42 Mo	43 Tc	44 Ru	45 Rh	46 Pd	47 Ag	48 Cd	49 In	50 Sn	51 Sb	52 Te	53 I	54 Xe
55 Cs	56 Ba	71 Lu	72 Hf	73 Ta	74 W	75 Re	76 Os	77 Ir	78 Pt	79 Au	80 Hg	81 Tl	82 Pb	83 Bi	84 Po	85 At	86 Rn
87 Fr	88 Ra	103 Lr	104 Rf	105 Db	106 Sg	107 Bh	108 Hs	109 Mt									
		57 La	58 Ce	59 Pr	60 Nd	61 Pm	62 Sm	63 Eu	64 Gd	65 Tb	66 Dy	67 Ho	68 Er	69 Tm	70 Yb		
		89 Ac	90 Th	91 Pa	92 U	93 Np	94 Pu	95 Am	96 Cm	97 Bk	98 Cf	99 Es	100 Fm	101 Md	102 No		

Alles!

De elementaire deeltjes

Niet één serie, maar drie!

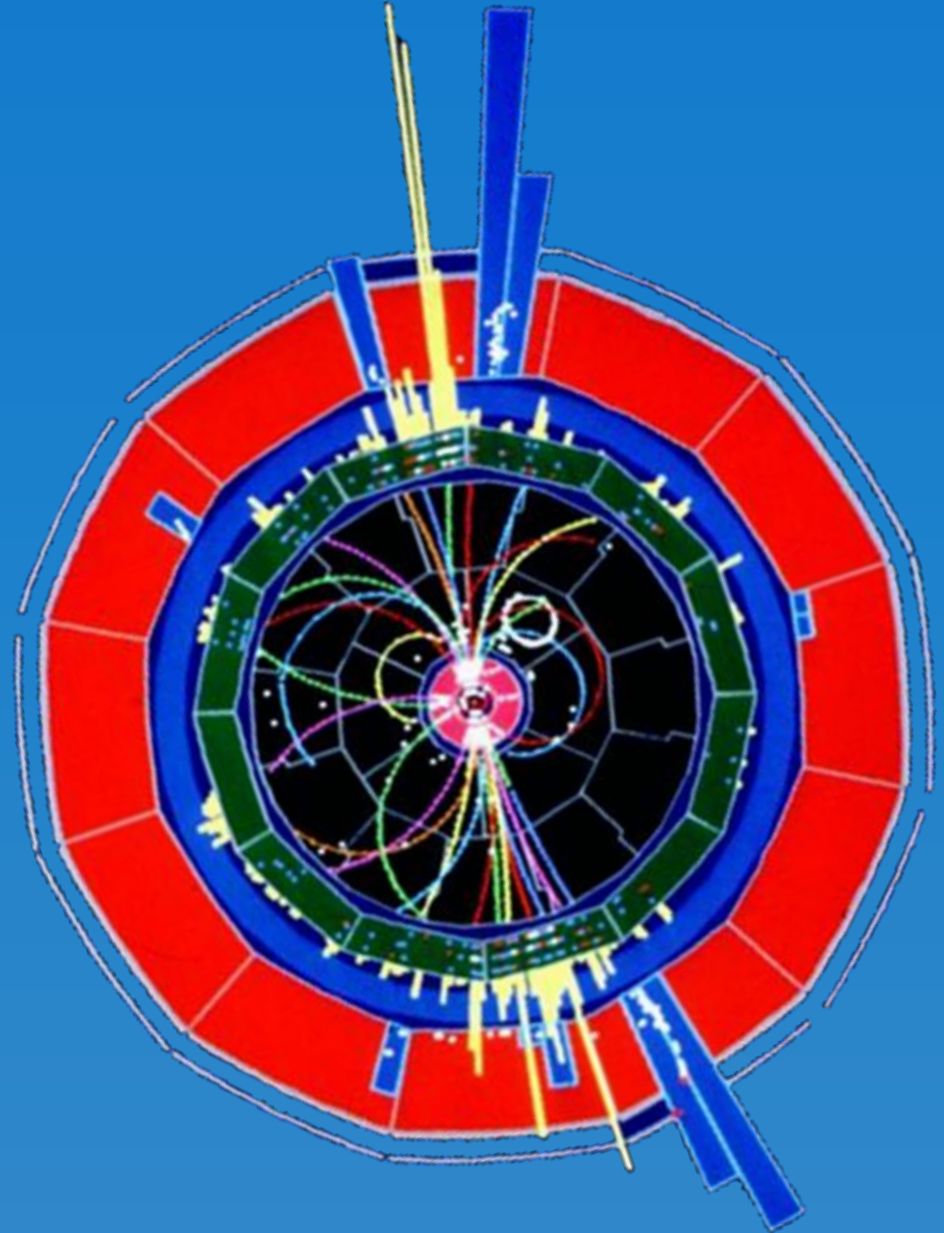
I II III

quarks

u	c <i>(1976)</i>	t <i>(1995)</i>
d	s <i>(1947)</i>	b <i>(1978)</i>

leptons

e <i>(1895)</i>	μ <i>(1936)</i>	τ <i>(1973)</i>
ν_e <i>(1956)</i>	ν_μ <i>(1963)</i>	ν_τ <i>(2000)</i>



De elementaire deeltjes

Generatie:

I II III Lading

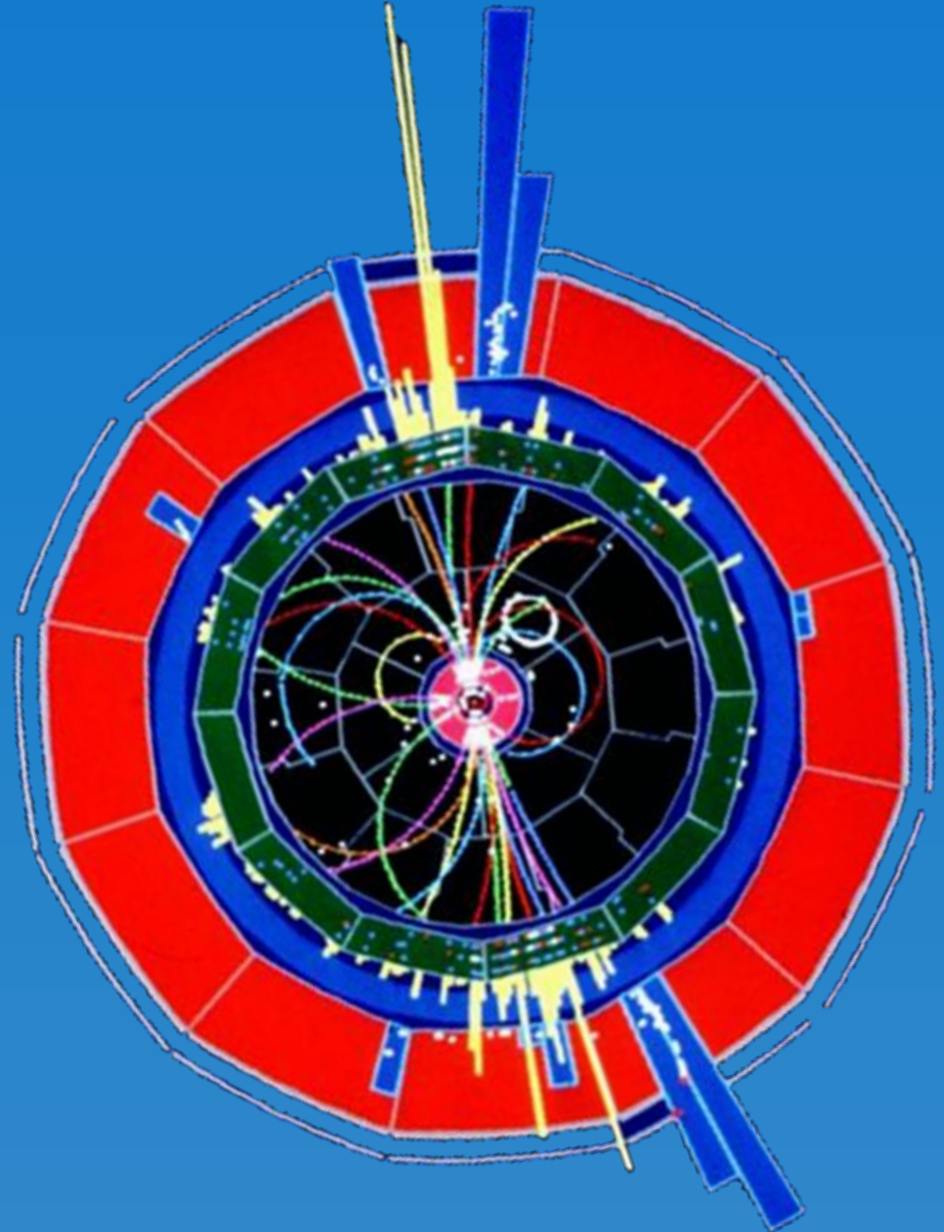
quarks

u (1976)	c (1976)	t (1995)	+2/3 e
d (1947)	s (1947)	b (1978)	-1/3 e

leptons

e (1895)	μ (1936)	τ (1973)	-1 e
ν_e (1956)	ν_μ (1963)	ν_τ (2000)	0 e

Materie

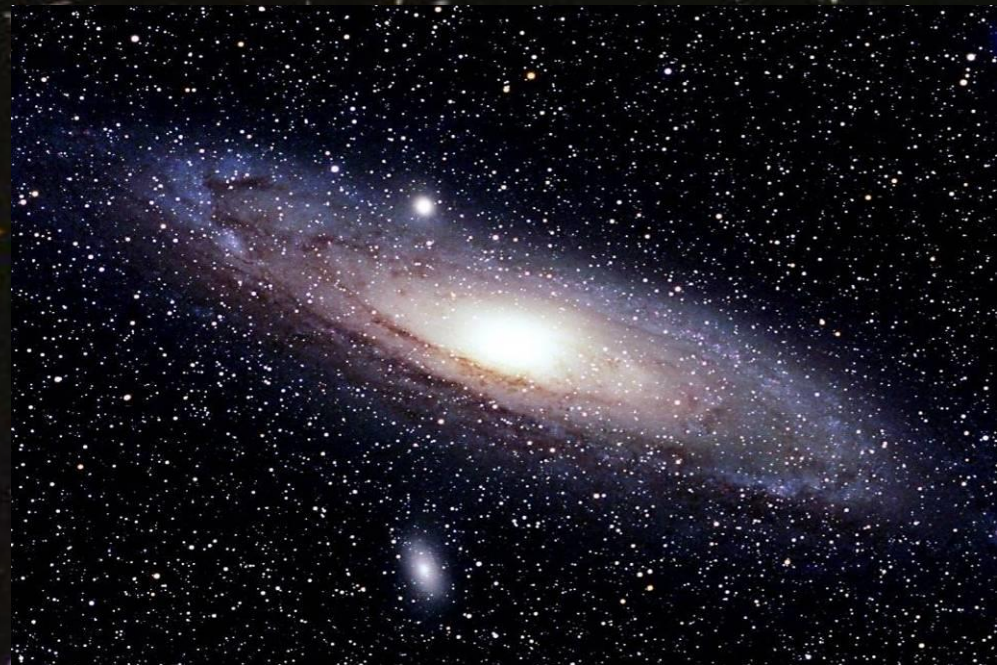


Is dit alles?

Generatie:

	I	II	III	<u>Lading</u>
quarks	u (1976)	c (1976)	t (1995)	+2/3 e
	d (1947)	s (1947)	b (1978)	-1/3 e
leptons	e (1895)	μ (1936)	τ (1973)	-1 e
	ν_e (1956)	ν_μ (1963)	ν_τ (2000)	0 e

Materie



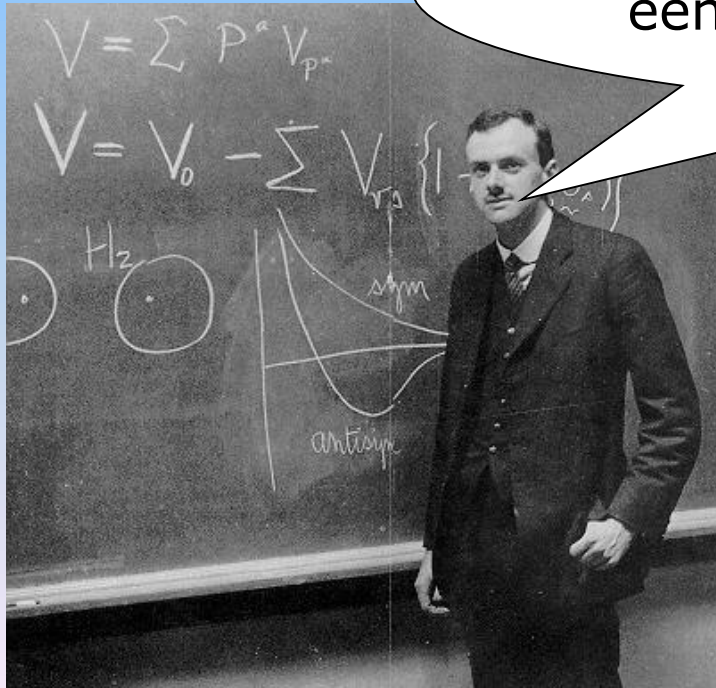
Anti-materie

Revoluties begin vorige eeuw:

- Relativiteitstheorie
- Quantum Mechanica

Paul Dirac (1928): relativistische quantum theorie!

Voor elk materiedeeltje bestaat een anti-materiedeeltje!



Anti-materie deeltje:

- Zelfde massa
- Tegenovergestelde lading

De elementaire deeltjes

	I	II	III	<u>Lading</u>
quarks	u (1976)	c (1976)	t (1995)	+2/3 e
	d	s (1947)	b (1978)	-1/3 e
leptons	e (1895)	μ (1936)	τ (1973)	-1 e
	ν_e (1956)	ν_μ (1963)	ν_τ (2000)	0 e

Materie

De elementaire deeltjes

	I	II	III	<u>Lading</u>
quarks	u	c <i>(1976)</i>	t <i>(1995)</i>	+2/3 e
	d	s <i>(1947)</i>	b <i>(1978)</i>	-1/3 e

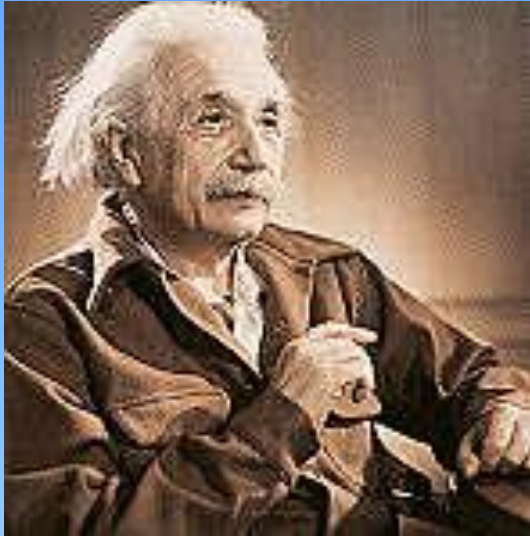
leptons	e <i>(1895)</i>	μ <i>(1936)</i>	τ <i>(1973)</i>	-1 e
	ν_e <i>(1956)</i>	ν_μ <i>(1963)</i>	ν_τ <i>(2000)</i>	0 e

Materie

<u>Lading</u>	I	II	III
-2/3 e	ū	c̄	t̄
+1/3 e	d̄	s̄	b̄
+1 e	ē	μ̄	τ̄
0 e	ν̄_e	ν̄_μ	ν̄_τ

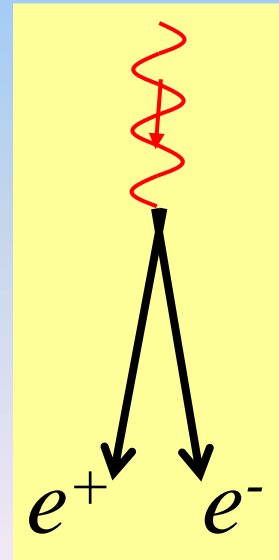
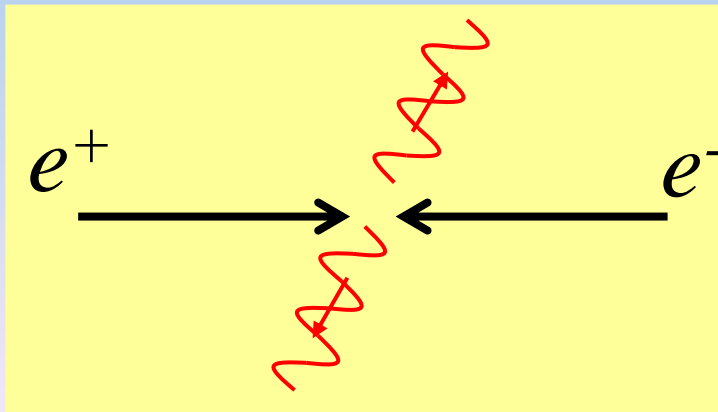
Anti-materie

Hoe maak je anti-materie??

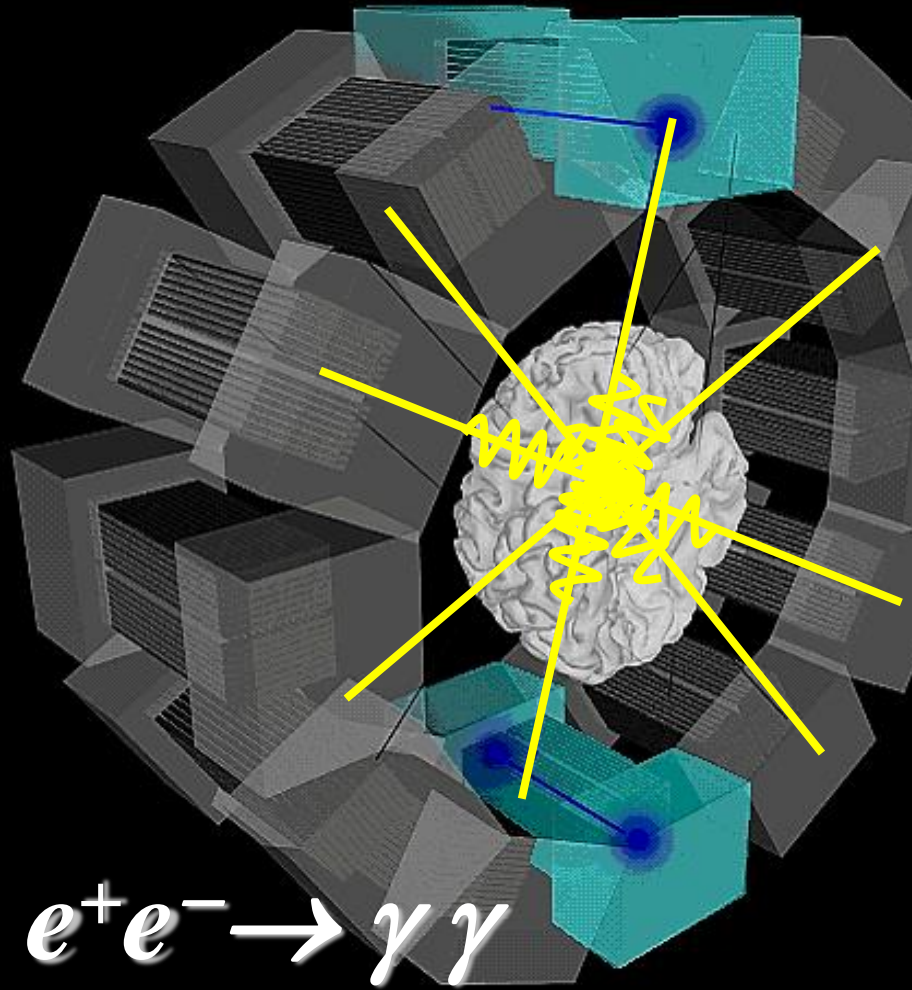
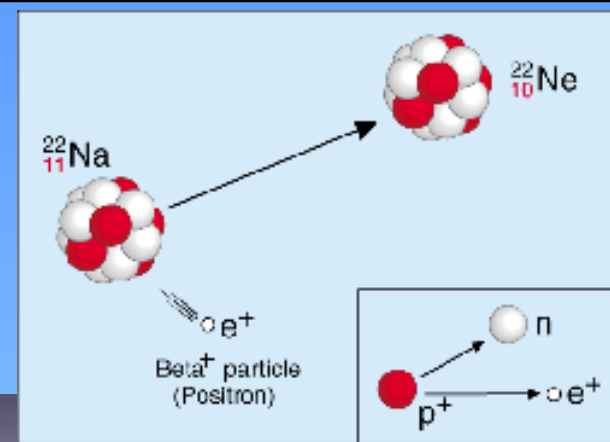


Albert Einstein:
 $E=mc^2$

materie + antimaterie = licht !
(en vice versa)



Anti-materie in ziekenhuizen: de PET-scan



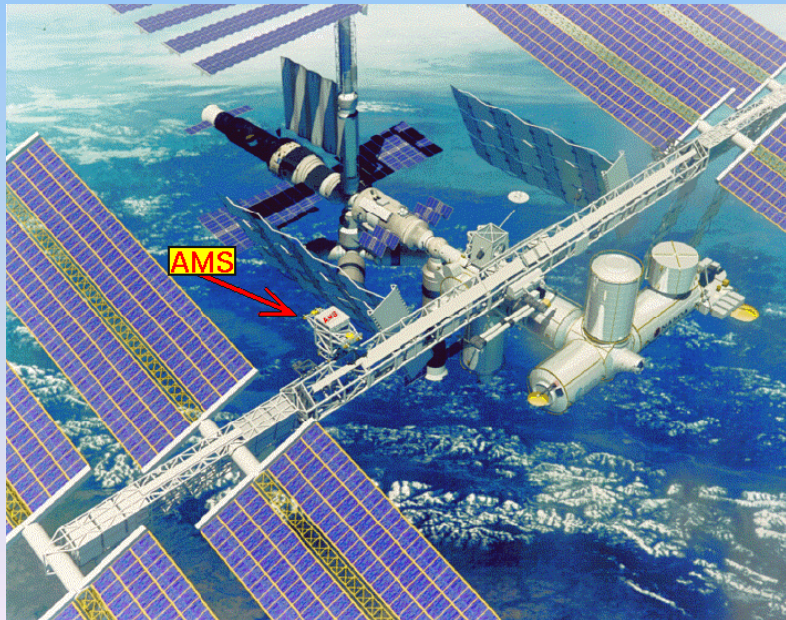
Wat snappen we nog niet:



I. Wat snappen we nog niet? “*Anti-materie*”

Waar is de anti-materie gebleven?

*Geen anti-materie
met satellieten*



*Geen anti-materie
sterrenstelsels*



II. Wat snappen we nog niet? “Higgs”

(Gedeeltelijk beantwoord op 4 juli 2012 !)

Massa van deeltjes



Bijzondere voorspelling:

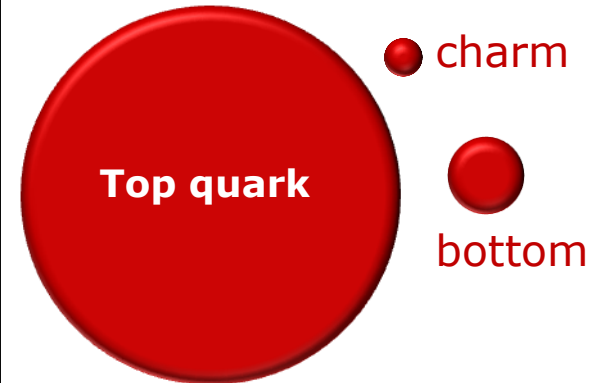
Het Higgs boson:

zorgt ervoor dat deeltjes massa kunnen hebben in de theorie

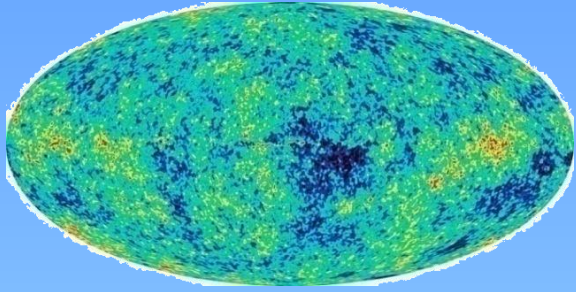
Neutrino's

- Elektron
- Muon
- Tau

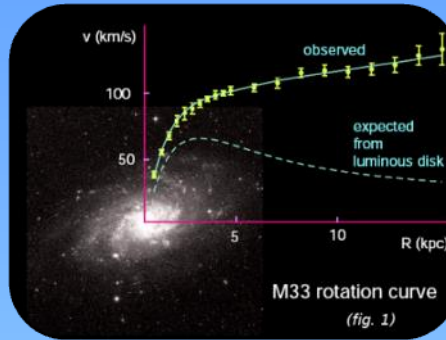
- up, down, strange



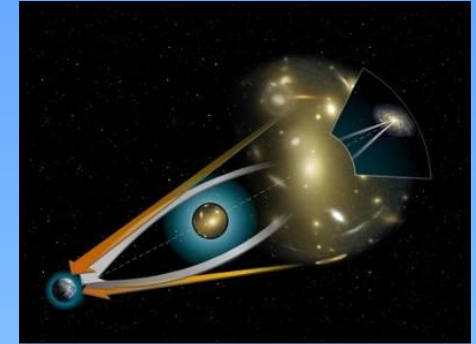
III. Wat snappen we nog niet? “Donkere materie”



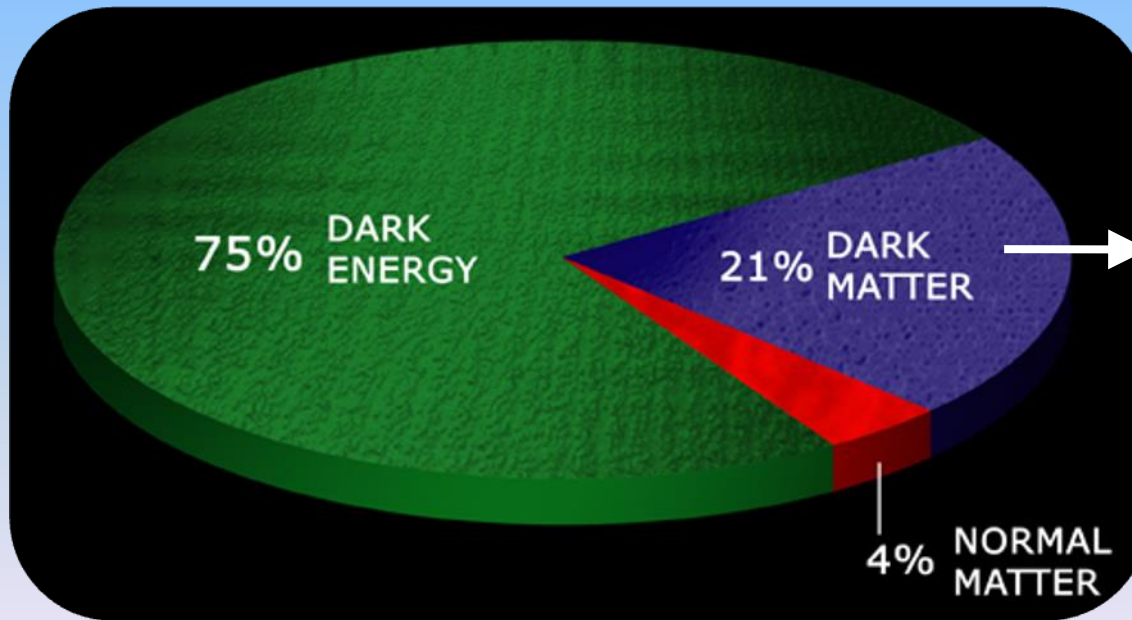
Temperatuurfluctuaties
structuur van het heelal



Rotatie-curves



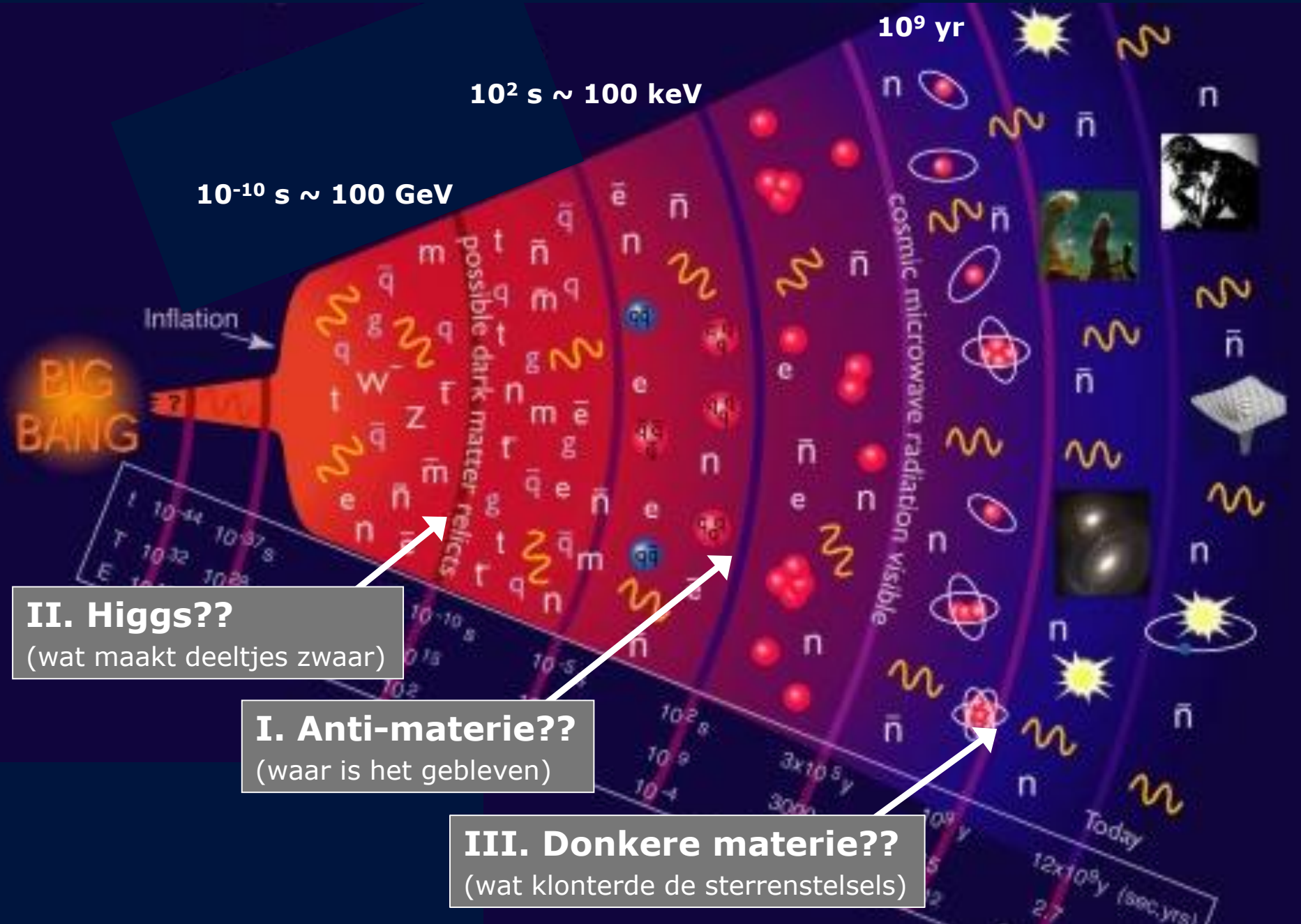
Gravitationele lens



Wat is de
donkere materie ?

We hebben al die tijd maar 4% van het heelal bestudeerd!

Wat snappen we niet? Drie Grote Vragen



II. Higgs??
(wat maakt deeltjes zwaar)

I. Anti-materie??
(waar is het gebleven)

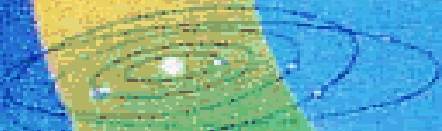
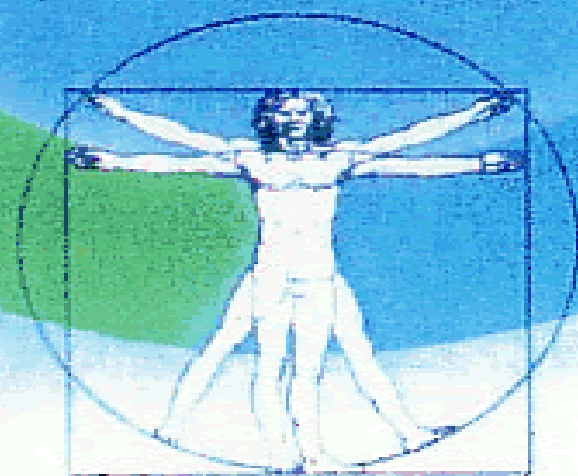
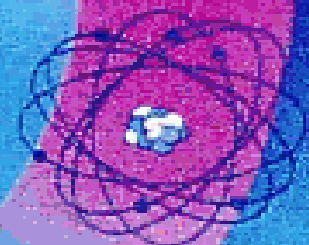
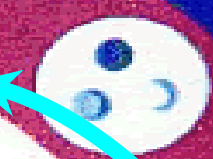
III. Donkere materie??
(wat klonterde de sterrenstelsels)

Astronomie

Deeltjes
fysica



Fundamenteel
(nieuwsgierigheid gedreven)
onderzoek



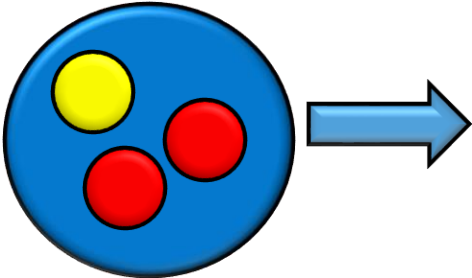




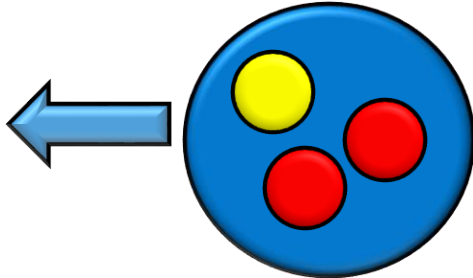
Klassiek botsen

Quantummechanisch botsen

proton



proton



Wat verwacht je ?

$$\begin{aligned}
 & -\frac{1}{2}\partial_\nu g_\mu^a \partial_\nu g_\mu^a - g_s f^{abc} \partial_\mu g_\nu^a g_\mu^b g_\nu^c - \frac{1}{4g_s^2} f^{abc} f^{ade} g_\mu^b g_\nu^c g_\mu^d g_\nu^e + \\
 & \frac{1}{2} i g_s^2 (\bar{q}^c \gamma^\mu q_j^c) g_\mu^a + G^a \partial^2 G^a + g_s f^{abc} \partial_\mu G^a G^b g_\mu^c - \partial_\nu W_\mu^+ \partial_\nu W_\mu^- - \\
 & M^2 W_\mu^+ W_\mu^- - \frac{1}{2} \partial_\nu Z_\mu^0 \partial_\nu Z_\mu^0 - \frac{1}{2c_w^2} M^2 Z_\mu^0 Z_\mu^0 - \frac{1}{2} \partial_\mu A_\nu \partial_\mu A_\nu - \frac{1}{2} \partial_\mu H \partial_\mu H - \\
 & \frac{1}{2} m_h^2 H^2 - \partial_\mu \phi^+ \partial_\mu \phi^- - M^2 \phi^+ \phi^- - \frac{1}{2} \partial_\mu \phi^0 \partial_\mu \phi^0 - \frac{1}{2c_w^2} M \phi^0 \phi^0 - \beta_h \left[\frac{2M^2}{g^2} + \right. \\
 & \left. \frac{2M}{g} H + \frac{1}{2} (H^2 + \phi^0 \phi^0 + 2\phi^+ \phi^-) \right] + \frac{2M^4}{g^2} \alpha_h - i g c_w [\partial_\nu Z_\mu^0 (W_\mu^+ W_\nu^- - \\
 & W_\nu^+ W_\mu^-) - Z_\nu^0 (W_\mu^+ \partial_\nu W_\mu^- - W_\nu^- \partial_\mu W_\mu^+) + Z_\mu^0 (W_\nu^+ \partial_\nu W_\mu^- - \\
 & W_\nu^- \partial_\mu W_\mu^+)] - i g s_w [\partial_\nu A_\mu (W_\mu^+ W_\nu^- - W_\nu^+ W_\mu^-) - A_\nu (W_\mu^+ \partial_\nu W_\mu^- - \\
 & W_\nu^- \partial_\mu W_\mu^+) + A_\mu (W_\nu^+ \partial_\nu W_\mu^- - W_\nu^- \partial_\mu W_\mu^+)] - \frac{1}{2} g^2 W_\mu^+ W_\mu^- W_\nu^+ W_\nu^- + \\
 & \frac{1}{2} g^2 W_\mu^+ W_\nu^- W_\mu^+ W_\nu^- + g^2 c_w^2 (Z_\mu^0 W_\mu^+ Z_\nu^0 W_\nu^- - Z_\mu^0 Z_\nu^0 W_\mu^+ W_\nu^-) + \\
 & g^2 s_w^2 (A_\mu W_\mu^+ A_\nu W_\nu^- - A_\mu A_\nu W_\mu^+ W_\nu^-) + g^2 s_w c_w [A_\mu Z_\nu^0 (W_\mu^+ W_\nu^- - \\
 & W_\nu^+ W_\mu^-) - 2A_\mu Z_\mu^0 W_\nu^+ W_\nu^-] - g \alpha [H^3 + H \phi^0 \phi^0 + 2H \phi^+ \phi^-] - \\
 & \frac{1}{8} g^2 \alpha_h [H^4 + (\phi^0)^4 + 4(\phi^+ \phi^-)^2 + 4(\phi^0)^2 \phi^+ \phi^- + 4H^2 \phi^+ \phi^- + 2(\phi^0)^2 H^2] - \\
 & g M W_\mu^+ W_\mu^- H - \frac{1}{2} g \frac{M}{c_w^2} Z_\mu^0 Z_\mu^0 H - \frac{1}{2} i g [W_\mu^+ (\phi^0 \partial_\mu \phi^- - \phi^- \partial_\mu \phi^0) - \\
 & W_\mu^- (\phi^0 \partial_\mu \phi^+ - \phi^+ \partial_\mu \phi^0)] + \frac{1}{2} g [W_\mu^+ (H \partial_\mu \phi^- - \phi^- \partial_\mu H) - W_\mu^- (H \partial_\mu \phi^+ - \\
 & \phi^+ \partial_\mu H)] + \frac{1}{2} g \frac{1}{c_w} (Z_\mu^0 (H \partial_\mu \phi^0 - \phi^0 \partial_\mu H) - i g \frac{s_w}{c_w} M Z_\mu^0 (W_\mu^+ \phi^- - W_\mu^- \phi^+) + \\
 & i g s_w M A_\mu (W_\mu^+ \phi^- - W_\mu^- \phi^+) - i g \frac{1-2c_w^2}{2c_w} Z_\mu^0 (\phi^+ \partial_\mu \phi^- - \phi^- \partial_\mu \phi^+) + \\
 & i g s_w A_\mu (\phi^+ \partial_\mu \phi^- - \phi^- \partial_\mu \phi^+) - \frac{1}{4} g^2 W_\mu^+ W_\mu^- [H^2 + (\phi^0)^2 + 2\phi^+ \phi^-] - \\
 & \frac{1}{4} g^2 \frac{1}{c_w^2} Z_\mu^0 Z_\mu^0 [H^2 + (\phi^0)^2 + 2(2s_w^2 - 1)^2 \phi^+ \phi^-] - \frac{1}{2} g^2 \frac{s_w^2}{c_w} Z_\mu^0 \phi^0 (W_\mu^+ \phi^- + \\
 & W_\mu^- \phi^+) - \frac{1}{2} i g^2 \frac{s_w^2}{c_w} Z_\mu^0 H (W_\mu^+ \phi^- - W_\mu^- \phi^+) + \frac{1}{2} g^2 s_w A_\mu \phi^0 (W_\mu^+ \phi^- + \\
 & W_\mu^- \phi^+) + \frac{1}{2} i g^2 s_w A_\mu H (W_\mu^+ \phi^- - W_\mu^- \phi^+) - g^2 \frac{s_w}{c_w} (2c_w^2 - 1) Z_\mu^0 A_\mu \phi^+ \phi^- - \\
 & g^1 s_w^2 A_\mu A_\mu \phi^+ \phi^- - \bar{e}^\lambda (\gamma \partial + m_e^\lambda) e^\lambda - \bar{\nu}^\lambda \gamma \partial \nu^\lambda - \bar{u}_j^\lambda (\gamma \partial + m_u^\lambda) u_j^\lambda - \\
 & \bar{d}_j^\lambda (\gamma \partial + m_d^\lambda) d_j^\lambda + i g s_w A_\mu [-\bar{e}^\lambda \gamma^\mu e^\lambda] + \frac{2}{3} (\bar{u}_j^\lambda \gamma^\mu u_j^\lambda) - \frac{1}{3} (\bar{d}_j^\lambda \gamma^\mu d_j^\lambda) + \\
 & \frac{i g}{4c_w} Z_\mu^0 [(\bar{\nu}^\lambda \gamma^\mu (1 + \gamma^5) \nu^\lambda) + (\bar{e}^\lambda \gamma^\mu (4s_w^2 - 1 - \gamma^5) e^\lambda) + (\bar{u}_j^\lambda \gamma^\mu (\frac{4}{3}s_w^2 - \\
 & 1 - \gamma^5) u_j^\lambda) + (\bar{d}_j^\lambda \gamma^\mu (1 - \frac{8}{3}s_w^2 - \gamma^5) d_j^\lambda)] + \frac{i g}{2\sqrt{2}} W_\mu^+ [(\bar{\nu}^\lambda \gamma^\mu (1 + \gamma^5) e^\lambda) + \\
 & (\bar{u}_j^\lambda \gamma^\mu (1 + \gamma^5) C_{\lambda\kappa} d_j^\kappa)] + \frac{i g}{2\sqrt{2}} W_\mu^- [(\bar{e}^\lambda \gamma^\mu (1 + \gamma^5) \nu^\lambda) + (\bar{d}_j^\kappa C_{\lambda\kappa}^\dagger \gamma^\mu (1 + \\
 & \gamma^5) u_j^\lambda)] + \frac{i g}{2\sqrt{2}} \frac{m_\lambda^2}{M} [-\phi^+ (\bar{\nu}^\lambda (1 - \gamma^5) e^\lambda) + \phi^- (\bar{e}^\lambda (1 + \gamma^5) \nu^\lambda)] - \\
 & \frac{g}{2} \frac{m_\lambda^2}{M} [H (\bar{e}^\lambda e^\lambda) + i \phi^0 (\bar{e}^\lambda \gamma^5 e^\lambda)] + \frac{i g}{2M\sqrt{2}} \phi^+ [-m_d^\lambda (\bar{u}_j^\lambda C_{\lambda\kappa} (1 - \gamma^5) d_j^\kappa) + \\
 & m_u^\lambda (\bar{u}_j^\lambda C_{\lambda\kappa} (1 + \gamma^5) d_j^\kappa)] + \frac{i g}{2M\sqrt{2}} \phi^- [m_d^\lambda (\bar{d}_j^\kappa C_{\lambda\kappa}^\dagger (1 + \gamma^5) u_j^\kappa) - m_u^\lambda (\bar{d}_j^\kappa C_{\lambda\kappa}^\dagger (1 - \\
 & \gamma^5) u_j^\kappa)] - \frac{g}{2} \frac{m_\lambda^2}{M} H (\bar{u}_j^\lambda u_j^\lambda) - \frac{g}{2} \frac{m_\lambda^2}{M} H (\bar{d}_j^\lambda d_j^\lambda) + \frac{i g}{2} \frac{m_\lambda^2}{M} \phi^0 (\bar{u}_j^\lambda \gamma^5 u_j^\lambda) - \\
 & \frac{i g}{2} \frac{m_\lambda^2}{M} \phi^0 (\bar{d}_j^\lambda \gamma^5 d_j^\lambda) + \bar{X}^+ (\partial^2 - M^2) X^+ + \bar{X}^- (\partial^2 - M^2) X^- + \bar{X}^0 (\partial^2 - \\
 & \frac{M^2}{c_w^2}) X^0 + \bar{Y} \partial^2 Y + i g c_w W_\mu^+ (\partial_\mu \bar{X}^0 X^- - \partial_\mu \bar{X}^+ X^0) + i g s_w W_\mu^+ (\partial_\mu \bar{Y} X^- - \\
 & \partial_\mu \bar{X}^+ Y) + i g c_w W_\mu^- (\partial_\mu \bar{X}^- X^0 - \partial_\mu \bar{X}^0 X^+) + i g s_w W_\mu^- (\partial_\mu \bar{X}^- Y - \\
 & \partial_\mu \bar{Y} X^+) + i g c_w Z_\mu^0 (\partial_\mu \bar{X}^+ X^+ - \partial_\mu \bar{X}^- X^-) + i g s_w A_\mu (\partial_\mu \bar{X}^+ X^+ - \\
 & \partial_\mu \bar{X}^- X^-) - \frac{1}{2} g M [\bar{X}^+ X^+ H + \bar{X}^- X^- H + \frac{1}{c_w^2} \bar{X}^0 X^0 H] + \\
 & \frac{1-2c_w^2}{2c_w} i g M [\bar{X}^+ X^0 \phi^+ - \bar{X}^- X^0 \phi^-] + \frac{1}{2c_w} i g M [\bar{X}^0 X^- \phi^+ - \bar{X}^0 X^+ \phi^-] + \\
 & i g M s_w [\bar{X}^0 X^- \phi^+ - \bar{X}^0 X^+ \phi^-] + \frac{1}{2} i g M [\bar{X}^+ X^+ \phi^0 - \bar{X}^- X^- \phi^0]
 \end{aligned}$$

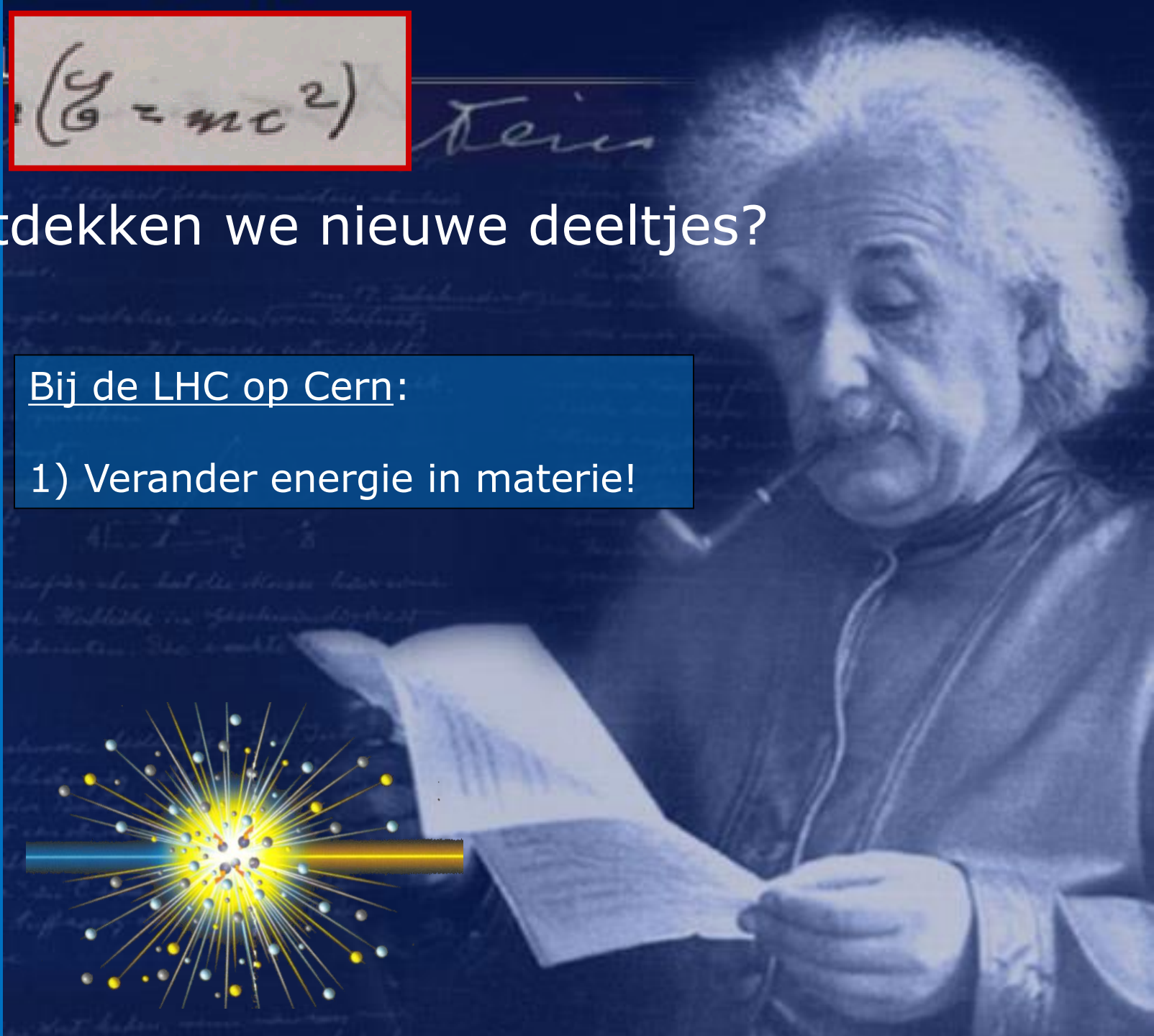
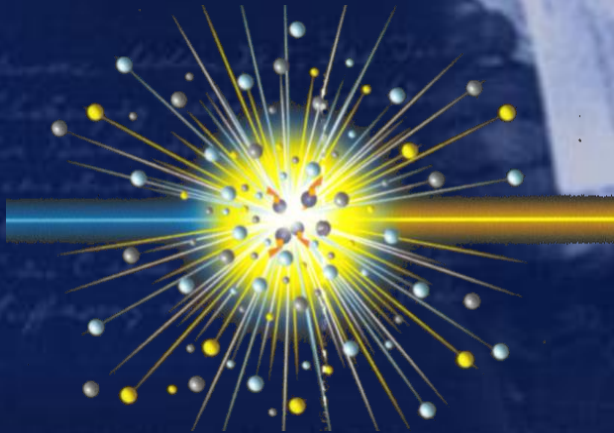
Al 40 jaar bestaan er precieze wiskundige voorspellingen!

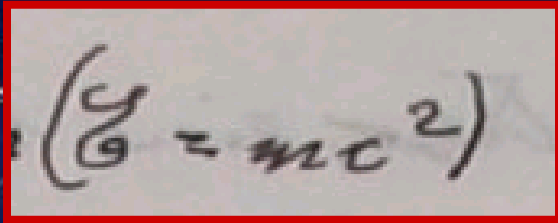
$$E = mc^2$$

Hoe ontdekken we nieuwe deeltjes?

Bij de LHC op Cern:

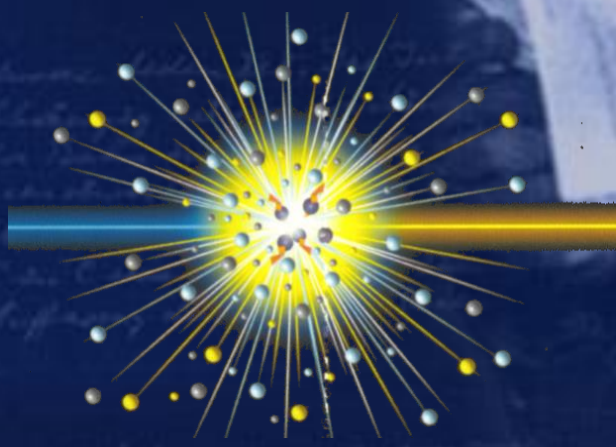
1) Verander energie in materie!




$$E = mc^2$$

Hoe ontdekken we nieuwe deeltjes?

- Bij de LHC op Cern:
- 1) Verander energie in materie!
 - 2) Nieuwe deeltjes veranderen voorspellingen





LHCb

ATLAS

CMS

ALICE

2) Nieuwe deeltjes
veranderen voorspellingen

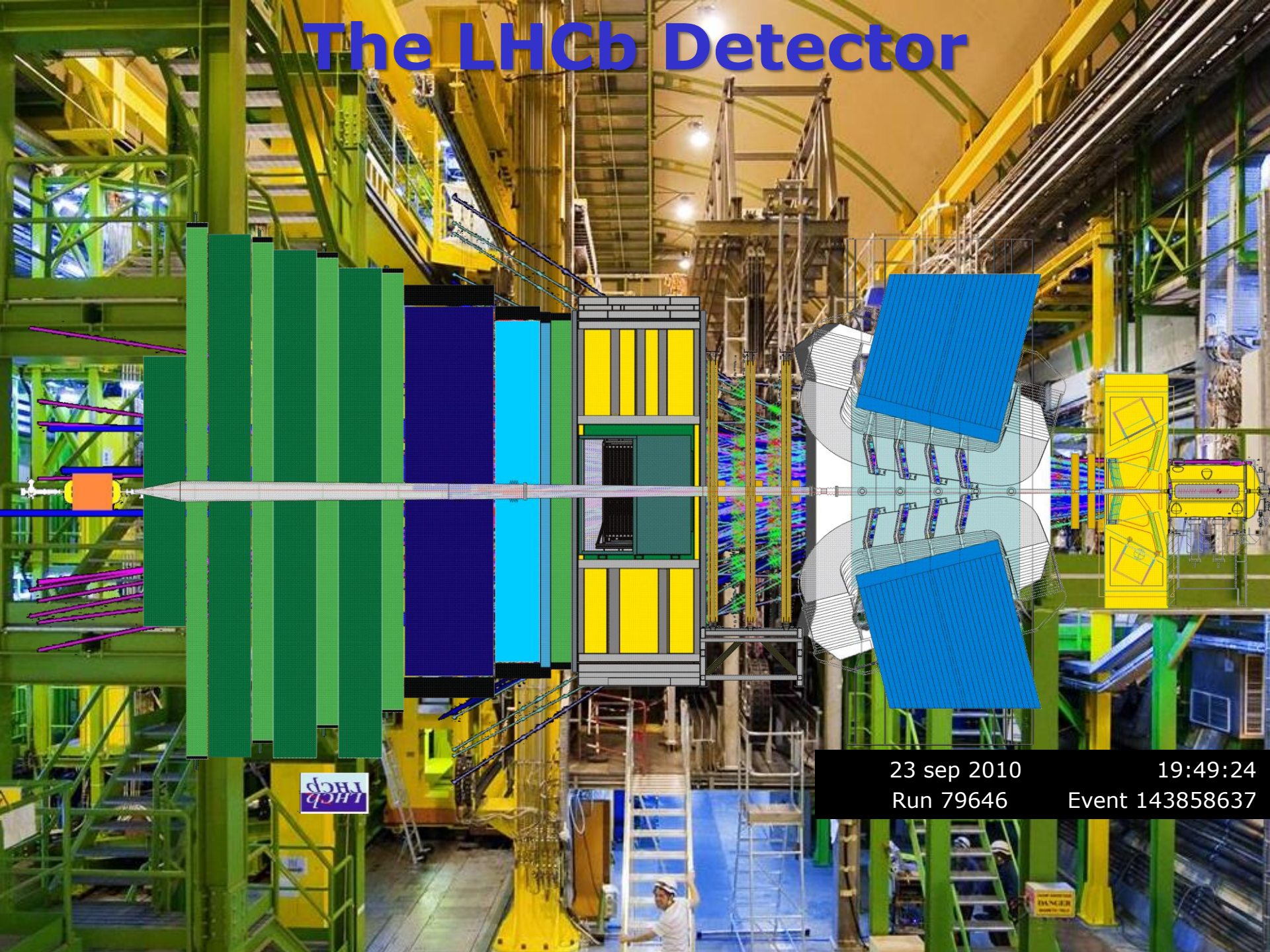
1) Verander energie in
materie

LHCb

ATLAS



The LHCb Detector



23 sep 2010

19:49:24

Run 79646

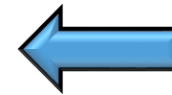
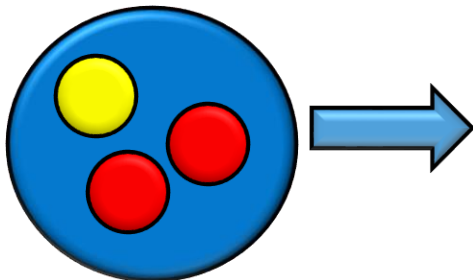
Event 143858637

DANGER

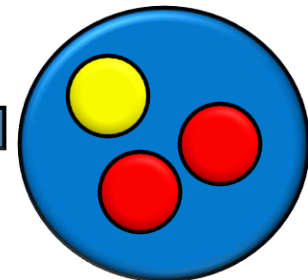
$$\begin{aligned}
& -\frac{1}{2}\partial_\nu g_\mu^a \partial_\nu g_\mu^a - g_a f^{abc} \partial_\mu g_\nu^b g_\mu^c - \frac{1}{2}g^2 f^{abc} f^{cde} g_\mu^b g_\nu^d g_\mu^e + \\
& \frac{1}{2}ig^2(\bar{q}^i \gamma_\mu q^j)g_\mu^a + G^a \partial^2 G^a + g_a f^{abc} \partial_\mu^i G^a G^b g_\mu^c - \partial_\nu W_\mu^+ \partial_\nu W_\mu^- - \\
& M^2 W_\mu^+ W_\mu^- - \frac{1}{2}\partial_\nu Z_\mu^0 \partial_\nu Z_\mu^0 - \frac{1}{2c_w} M^2 Z_\mu^0 Z_\mu^0 - \frac{1}{2}\partial_\mu A_\nu \partial_\mu A_\nu - \frac{1}{2}\partial_\mu H \partial_\mu H - \\
& \frac{1}{2}m_h^2 H^2 - \partial_\mu \phi^+ \partial_\mu \phi^- - M^2 \phi^+ \phi^- - \frac{1}{2}\partial_\mu \phi^0 \partial_\mu \phi^0 - \frac{1}{2c_w} M \phi^0 \phi^0 - \beta_h \left[\frac{2M^2}{g^2} + \right. \\
& \left. \frac{2M}{g} H + \frac{1}{2}(H^2 + \phi^0 \phi^0 + 2\phi^+ \phi^-) \right] + \frac{2M}{g^2} \alpha_h - ig_{c_w} [\partial_\nu Z_\mu^0 (W_\mu^+ W_\nu^- - \\
& W_\nu^+ W_\mu^-) - Z_\mu^0 (W_\mu^+ \partial_\nu W_\nu^- - W_\nu^- \partial_\mu W_\mu^+) + Z_\mu^0 (W_\nu^+ \partial_\nu W_\mu^- - \\
& W_\nu^- \partial_\mu W_\mu^+)] - ig_{s_w} [\partial_\nu A_\mu (W_\mu^+ W_\nu^- - W_\nu^- W_\mu^+) - A_\nu (W_\mu^+ \partial_\nu W_\mu^- - \\
& W_\nu^- \partial_\mu W_\mu^+) + A_\mu (W_\mu^+ \partial_\nu W_\nu^- - W_\nu^- \partial_\mu W_\mu^+)] - \frac{1}{2}g^2 W_\mu^+ W_\nu^- W_\mu^+ W_\nu^- + \\
& \frac{1}{2}g^2 W_\mu^+ W_\nu^- W_\mu^+ W_\nu^- + g^2 c_w^2 (Z_\mu^0 W_\nu^+ Z_\mu^0 W_\nu^- - Z_\mu^0 Z_\nu^0 W_\mu^+ W_\nu^-) + \\
& g^2 s_w^2 (A_\mu W_\nu^+ A_\nu W_\mu^- - A_\mu A_\nu W_\nu^+ W_\mu^-) + g^2 s_w c_w [A_\mu Z_\nu^0 (W_\mu^+ W_\nu^- - \\
& W_\nu^- W_\mu^+) - 2A_\nu Z_\mu^0 W_\mu^+ W_\nu^-] - g\alpha [H^3 + H\phi^0 \phi^0 + 2H\phi^+ \phi^-] - \\
& \frac{1}{4}g^2 \alpha_h [H^4 + (\phi^0)^4 + 4(\phi^+ \phi^-)^2 + 4(\phi^0)^2 \phi^+ \phi^- + 4H^2 \phi^+ \phi^- + 2(\phi^0)^2 H^2] - \\
& gM W_\mu^+ W_\nu^- H - \frac{1}{2}g \frac{M}{c_w} Z_\mu^0 Z_\nu^0 H - \frac{1}{2}ig [W_\mu^+ (\phi^0 \partial_\nu \phi^- - \phi^- \partial_\nu \phi^0) - \\
& W_\nu^- (\phi^0 \partial_\mu \phi^+ - \phi^+ \partial_\mu \phi^0)] + \frac{1}{2}ig [W_\mu^+ (H \partial_\nu \phi^- - \phi^- \partial_\nu H) - W_\nu^- (H \partial_\mu \phi^+ - \\
& \phi^+ \partial_\mu H)] + \frac{1}{2}g \frac{1}{c_w} (Z_\mu^0 (H \partial_\nu \phi^0 - \phi^0 \partial_\nu H) - ig \frac{c_w^2}{c_w} M Z_\mu^0 (W_\mu^+ \phi^- - W_\nu^- \phi^+) + \\
& ig_{s_w} M A_\mu (W_\mu^+ \phi^- - W_\nu^- \phi^+) - ig \frac{2c_w}{c_w} Z_\mu^0 (\phi^+ \partial_\nu \phi^- - \phi^- \partial_\nu \phi^+) + \\
& ig_{s_w} A_\mu (\phi^+ \partial_\nu \phi^- - \phi^- \partial_\nu \phi^+) - \frac{1}{4}g^2 W_\mu^+ W_\nu^- [H^2 + (\phi^0)^2 + 2\phi^+ \phi^-] - \\
& \frac{1}{4}g^2 \frac{1}{2} Z_\mu^0 Z_\nu^0 [H^2 + (\phi^0)^2 + 2(2s_w^2 - 1)^2 \phi^+ \phi^-] - \frac{1}{2}g^2 \frac{c_w^2}{c_w} Z_\mu^0 \phi^0 (W_\mu^+ \phi^- + \\
& W_\nu^- \phi^+) - \frac{1}{2}ig^2 \frac{c_w^2}{c_w} Z_\mu^0 H (W_\mu^+ \phi^- - W_\nu^- \phi^+) + \frac{1}{2}g^2 s_w A_\mu \phi^0 (W_\mu^+ \phi^- + \\
& W_\nu^- \phi^+) + \frac{1}{2}ig^2 s_w A_\mu H (W_\mu^+ \phi^- - W_\nu^- \phi^+) - g^2 c_w (2c_w^2 - 1) Z_\mu^0 A_\nu \phi^+ \phi^- - \\
& g^1 s_w^2 A_\mu A_\nu \phi^+ \phi^- - \bar{e}^\lambda (\gamma \partial + m_e^\lambda) e^\lambda - \bar{\nu}^\lambda \gamma \partial \nu^\lambda - \bar{u}_j^\lambda (\gamma \partial + m_u^\lambda) u_j^\lambda - \\
& \bar{d}_j^\lambda (\gamma \partial + m_d^\lambda) d_j^\lambda + ig_{s_w} A_\mu [-(\bar{e}^\lambda \gamma^\mu e^\lambda) + \frac{2}{3}(\bar{u}_j^\lambda \gamma^\mu u_j^\lambda) - \frac{1}{3}(\bar{d}_j^\lambda \gamma^\mu d_j^\lambda)] + \\
& \frac{ig}{4c_w} Z_\mu^0 [(\bar{\nu}^\lambda \gamma^\mu (1 + \gamma^5) \nu^\lambda) + (\bar{e}^\lambda \gamma^\mu (4s_w^2 - 1 - \gamma^5) e^\lambda) + (\bar{u}_j^\lambda \gamma^\mu (\frac{4}{3}s_w^2 - \\
& 1 - \gamma^5) u_j^\lambda) + (\bar{d}_j^\lambda \gamma^\mu (1 - \frac{8}{3}s_w^2 - \gamma^5) d_j^\lambda)] + \frac{ig}{2\sqrt{2}} W_\mu^+ [(\bar{\nu}^\lambda \gamma^\mu (1 + \gamma^5) e^\lambda) + \\
& (\bar{u}_j^\lambda \gamma^\mu (1 + \gamma^5) C_{\lambda\kappa} d_j^\kappa)] + \frac{ig}{2\sqrt{2}} W_\mu^- [(\bar{e}^\lambda \gamma^\mu (1 + \gamma^5) \nu^\lambda) + (\bar{d}_j^\kappa C_{\lambda\kappa} \gamma^\mu (1 + \\
& \gamma^5) u_j^\lambda)] + \frac{ig}{2\sqrt{2}} \frac{m_h^2}{M} [-\phi^+ (\bar{\nu}^\lambda (1 - \gamma^5) e^\lambda) + \phi^- (\bar{e}^\lambda (1 + \gamma^5) \nu^\lambda)] - \\
& \frac{g}{2} \frac{m_h^2}{M} [H (\bar{e}^\lambda e^\lambda) + i\phi^0 (\bar{e}^\lambda \gamma^5 e^\lambda)] + \frac{ig}{2M\sqrt{2}} \phi^+ [-m_h^2 (\bar{u}_j^\lambda C_{\lambda\kappa} (1 - \gamma^5) d_j^\kappa) + \\
& m_h^2 (\bar{u}_j^\lambda C_{\lambda\kappa} (1 + \gamma^5) d_j^\kappa)] + \frac{ig}{2M\sqrt{2}} \phi^- [m_h^2 (\bar{d}_j^\kappa C_{\lambda\kappa}^1 (1 + \gamma^5) u_j^\lambda) - m_h^2 (\bar{d}_j^\kappa C_{\lambda\kappa}^1 (1 - \\
& \gamma^5) u_j^\lambda)] - \frac{g}{2} \frac{m_h^2}{M} H (\bar{u}_j^\lambda u_j^\lambda) - \frac{g}{2} \frac{m_h^2}{M} H (\bar{d}_j^\lambda d_j^\lambda) + \frac{ig}{2} \frac{m_h^2}{M} \phi^0 (\bar{u}_j^\lambda \gamma^5 u_j^\lambda) - \\
& \frac{ig}{2} \frac{m_h^2}{M} \phi^0 (\bar{d}_j^\lambda \gamma^5 d_j^\lambda) + \bar{X}^+ (\partial^2 - M^2) X^+ + \bar{X}^- (\partial^2 - M^2) X^- + \bar{X}^0 (\partial^2 - \\
& \frac{M^2}{c_w^2}) X^0 + \bar{Y} \partial^2 Y + ig_{c_w} W_\mu^+ (\partial_\mu \bar{X}^0 X^- - \partial_\nu \bar{X}^+ X^0) + ig_{s_w} W_\mu^+ (\partial_\mu \bar{Y} X^- - \\
& \partial_\nu \bar{X}^+ Y) + ig_{c_w} W_\nu^- (\partial_\mu \bar{X}^- X^0 - \partial_\nu \bar{X}^0 X^+) + ig_{s_w} W_\nu^- (\partial_\mu \bar{X}^- Y - \\
& \partial_\nu \bar{Y} X^+) + ig_{c_w} Z_\mu^0 (\partial_\mu \bar{X}^+ X^+ - \partial_\nu \bar{X}^- X^-) + ig_{s_w} A_\mu (\partial_\mu \bar{X}^+ X^+ - \\
& \partial_\nu \bar{X}^- X^-) - \frac{1}{2}gM [\bar{X}^+ X^+ H + \bar{X}^- X^- H + \frac{1}{c_w} \bar{X}^0 X^0 H] + \\
& \frac{1-2c_w^2}{2c_w} igM [\bar{X}^+ X^0 \phi^+ - \bar{X}^- X^0 \phi^-] + \frac{1}{2c_w} igM [\bar{X}^0 X^- \phi^+ - \bar{X}^0 X^+ \phi^-] + \\
& igM s_w [\bar{X}^0 X^- \phi^+ - \bar{X}^0 X^+ \phi^-] + \frac{1}{2}igM [\bar{X}^+ X^+ \phi^0 - \bar{X}^- X^- \phi^0]
\end{aligned}$$

Quantummechanisch botsen

proton



proton



LHCb in getallen

120,000 B events per sec

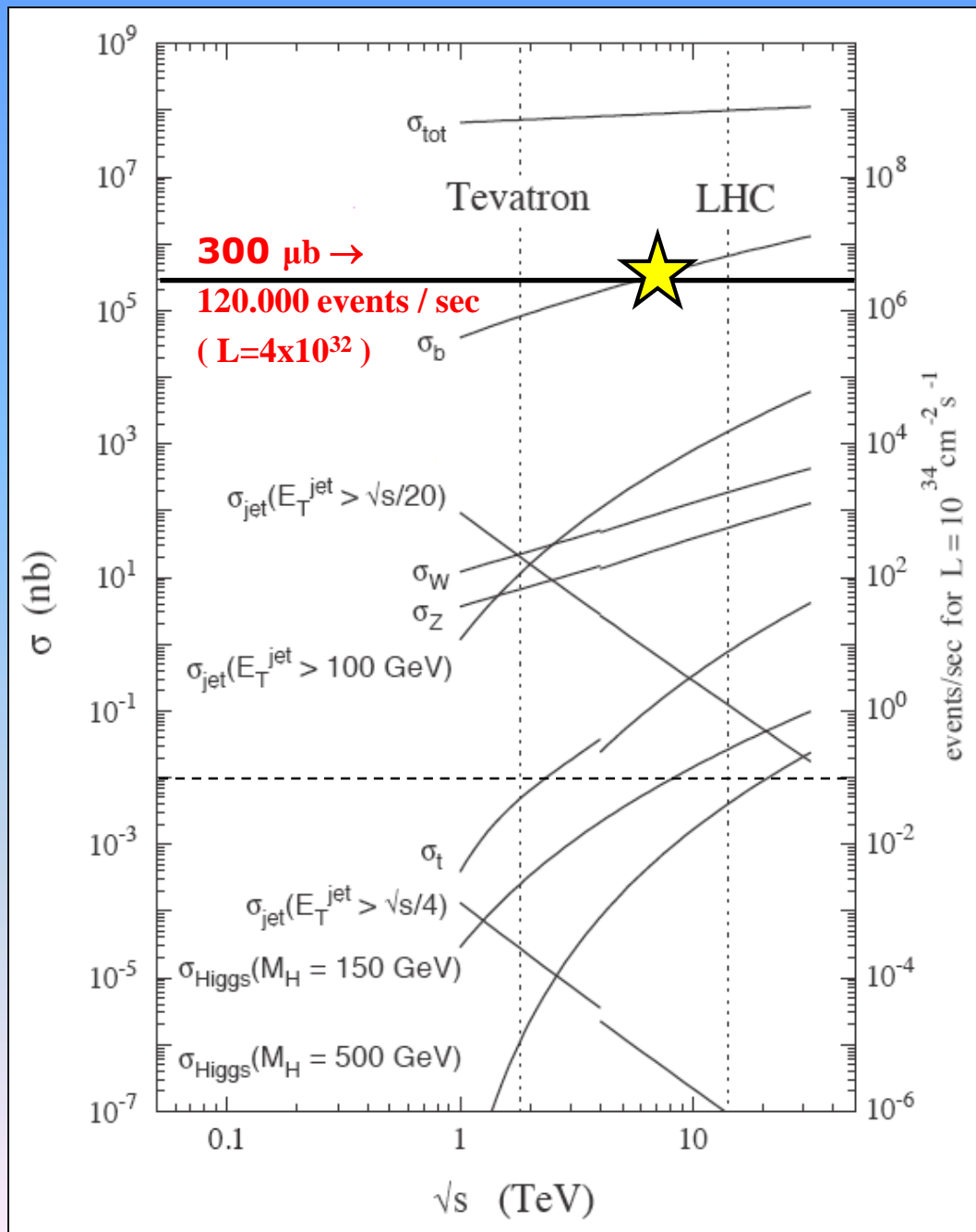
(ter vergelijking: in ATLAS : 1 Higgs in 100 sec)

10^{12} B events in 2022

(ter vergelijking: Babar heeft in totaal 10^9 B events)

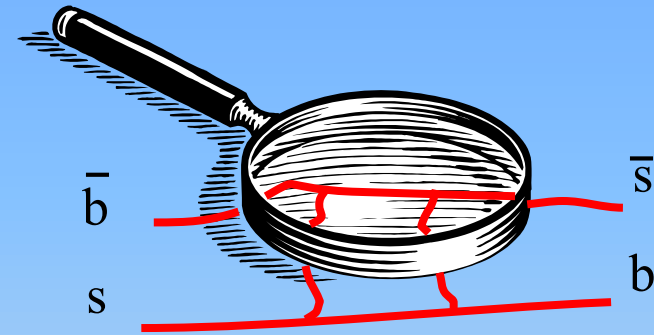
5 kHz naar tape

(ter vergelijking: ATLAS schrijft 200 Hz weg)

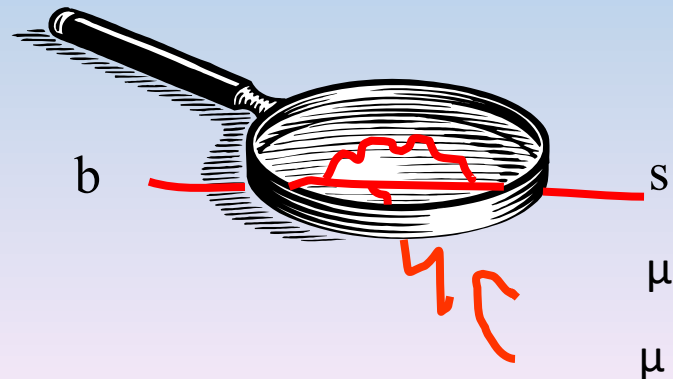


LHCb: bestuderen van B deeltje

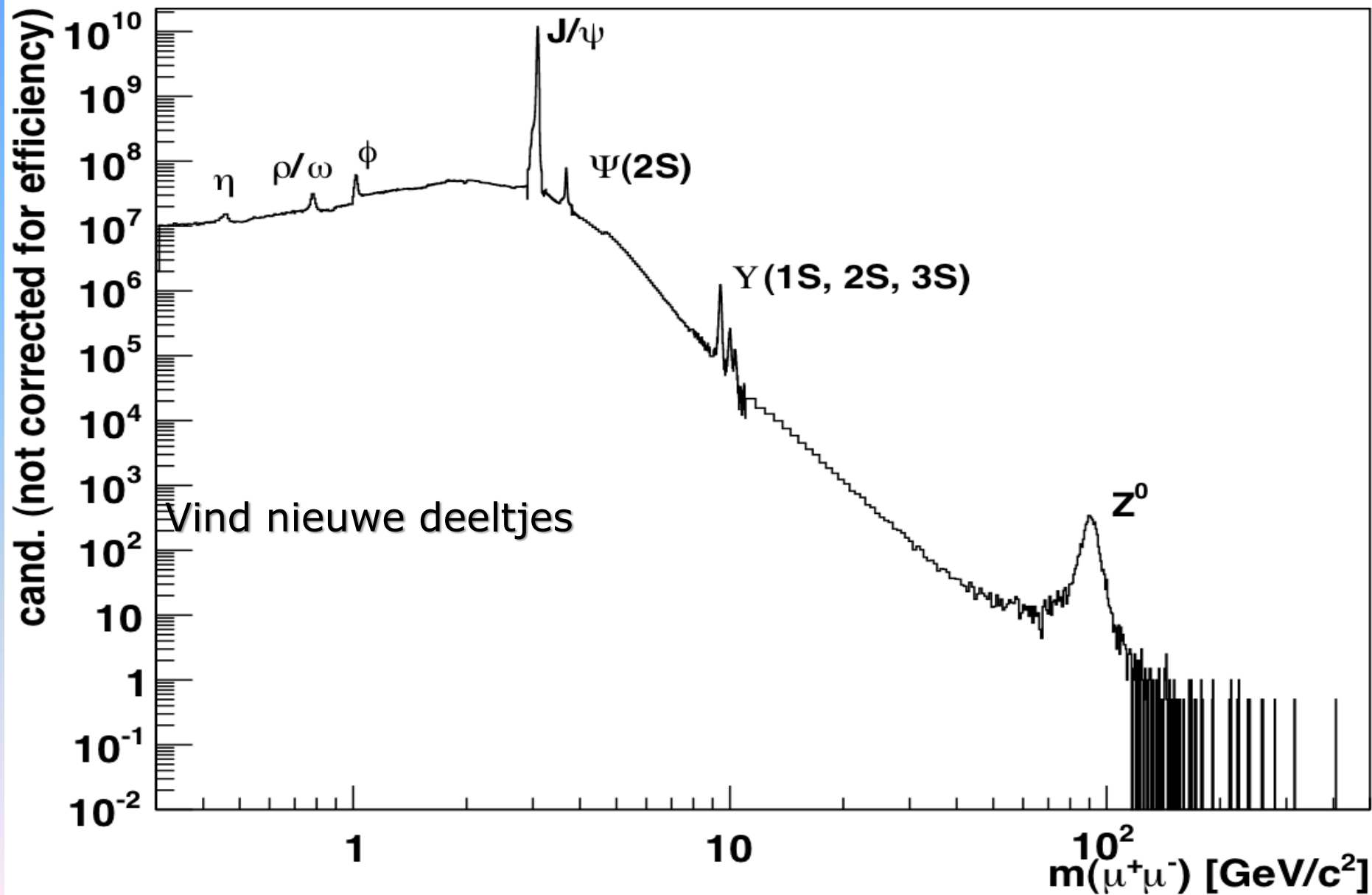
1) Vind verschillen tussen materie en anti-materie



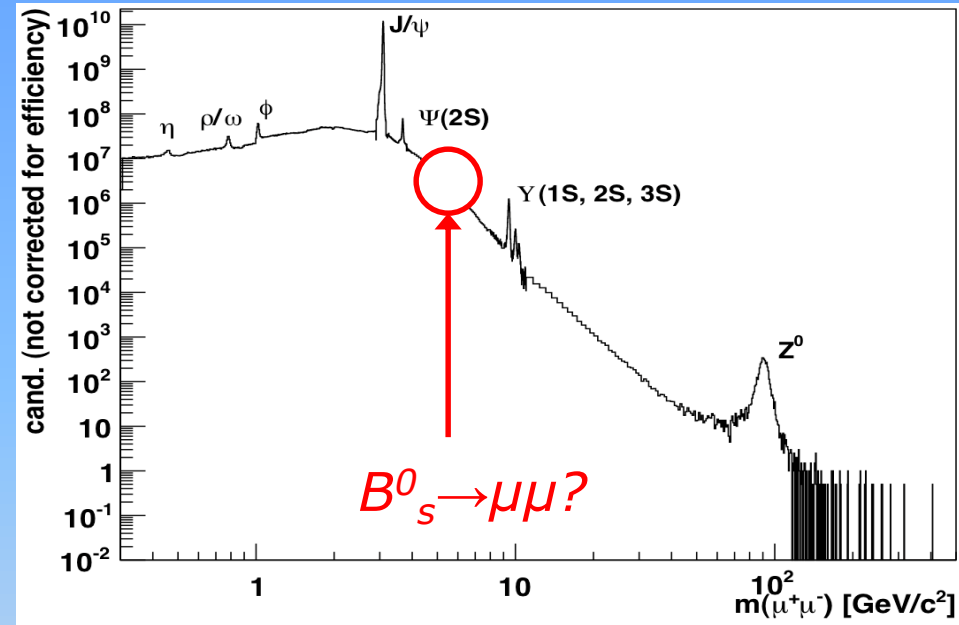
2) Vind nieuwe deeltjes



LHCb: bestuderen van B deeltje

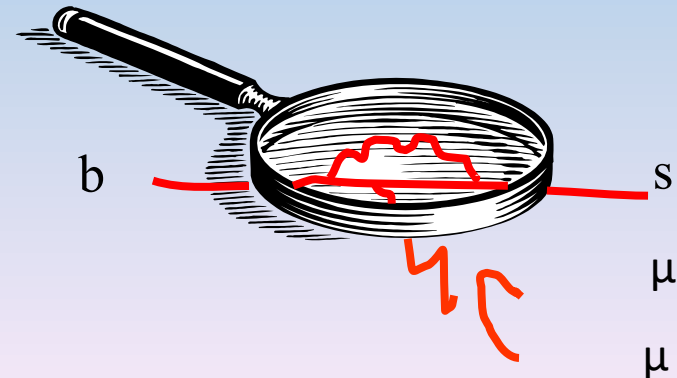
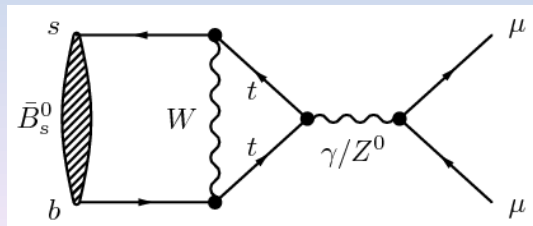


LHCb: bestuderen van B deeltje

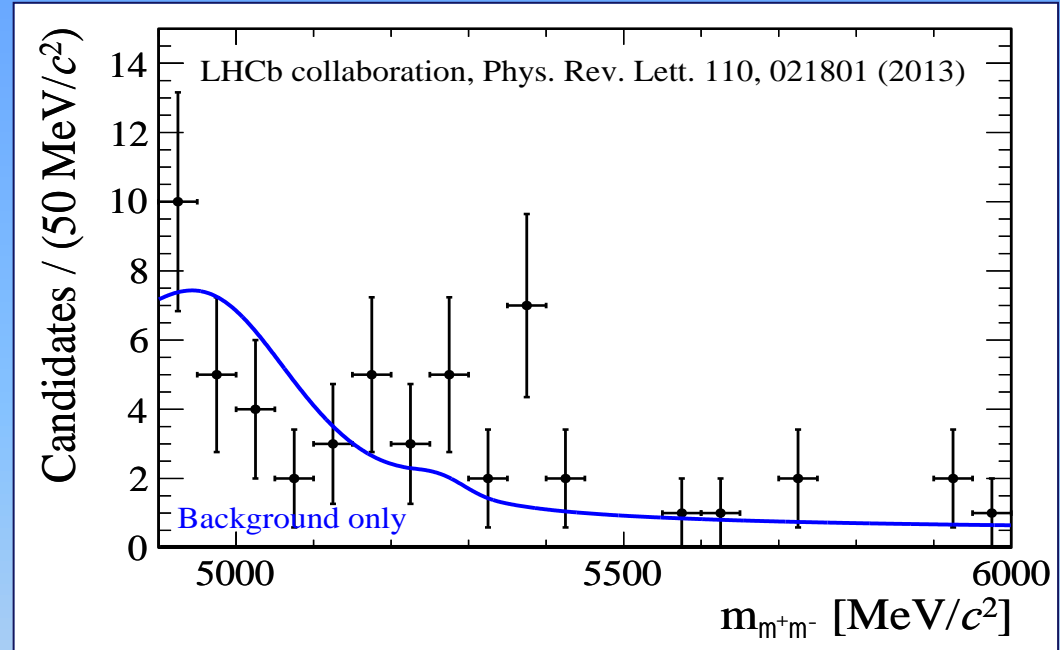


2) Vind nieuwe deeltjes

$$B^0_s \rightarrow \mu\mu$$

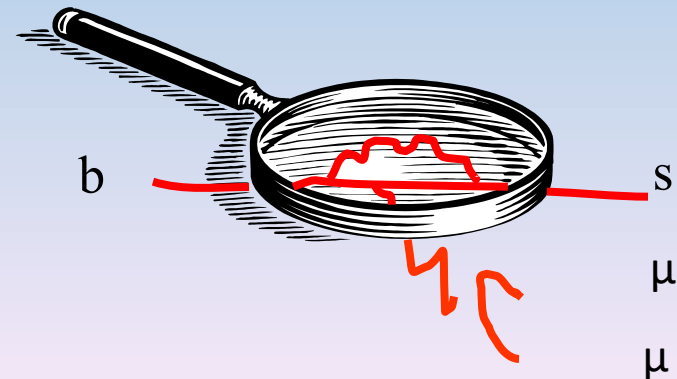


LHCb: bestuderen van B deeltje

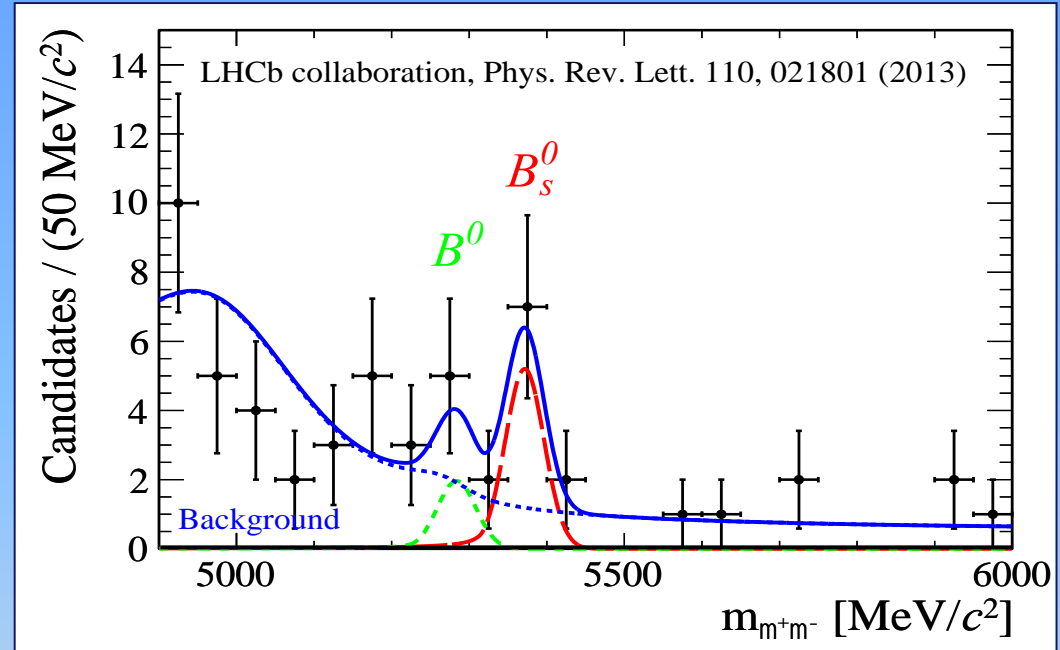


Slechts 3 op de miljard B deeltjes vervalst naar 2 muonen

Bestaan er nieuwe deeltjes?

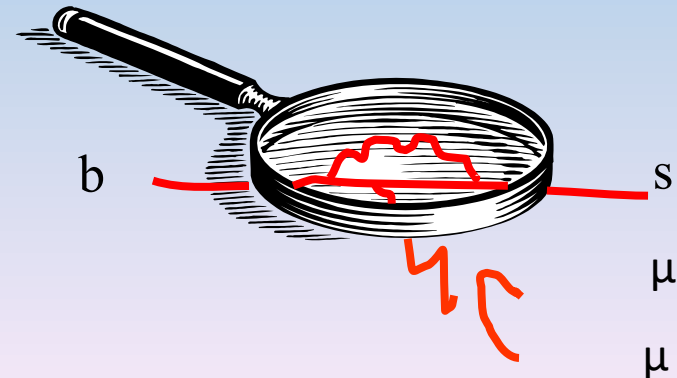


LHCb: bestuderen van B deeltje

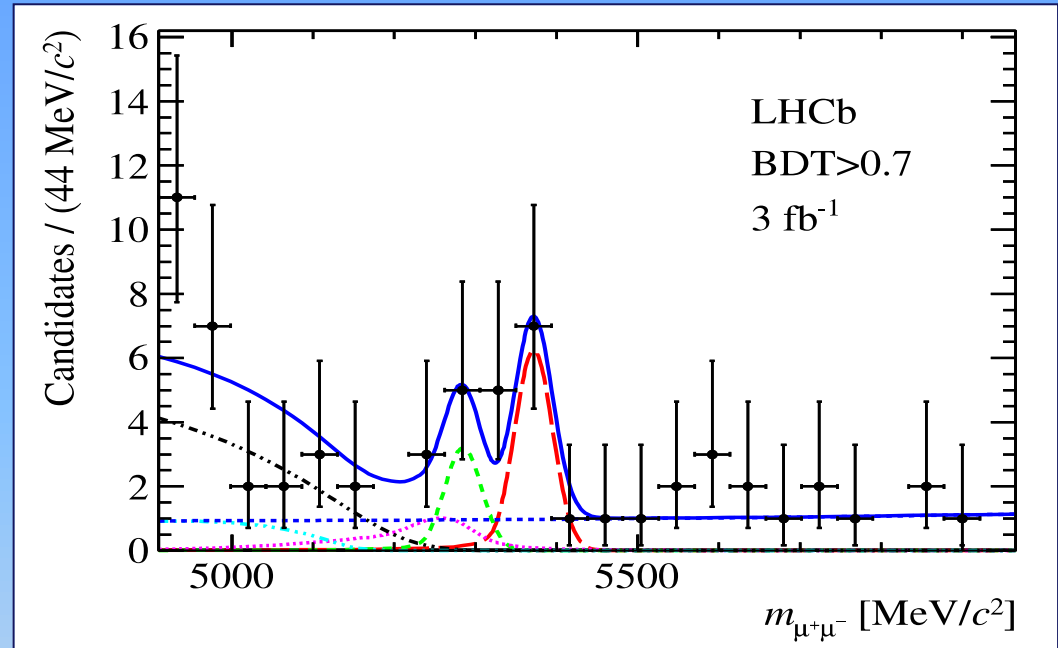


Slechts 3 op de miljard B deeltjes vervalst naar 2 muonen

Bestaan er nieuwe deeltjes?

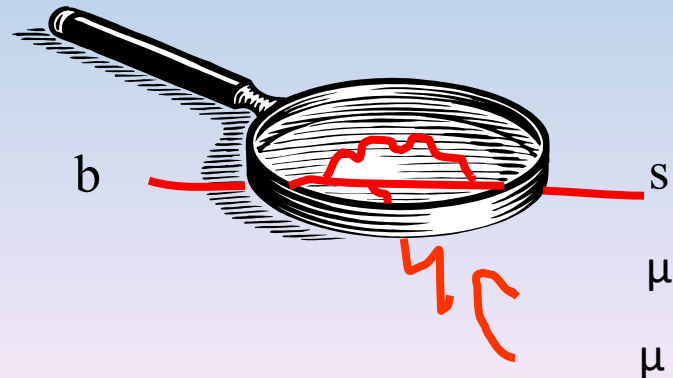


LHCb: bestuderen van B deeltje

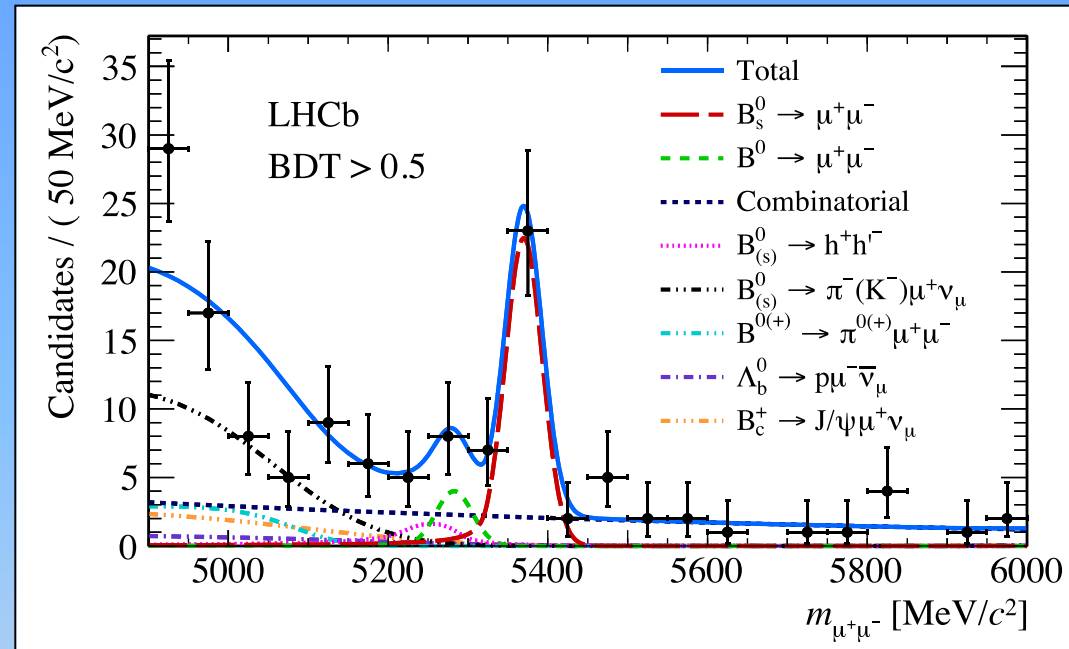


Slechts 3 op de miljard B deeltjes vervalst naar 2 muonen

Bestaan er nieuwe deeltjes?

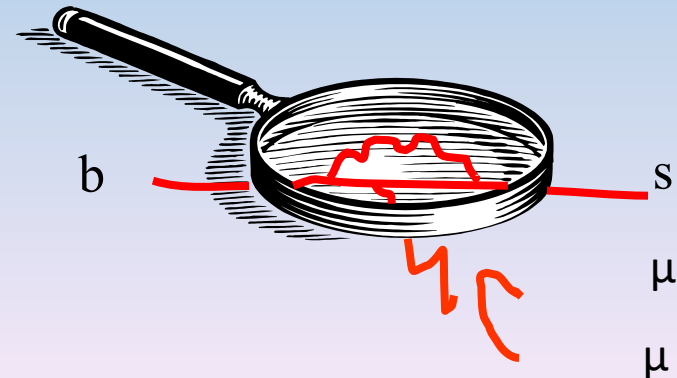


LHCb: bestuderen van B deeltje

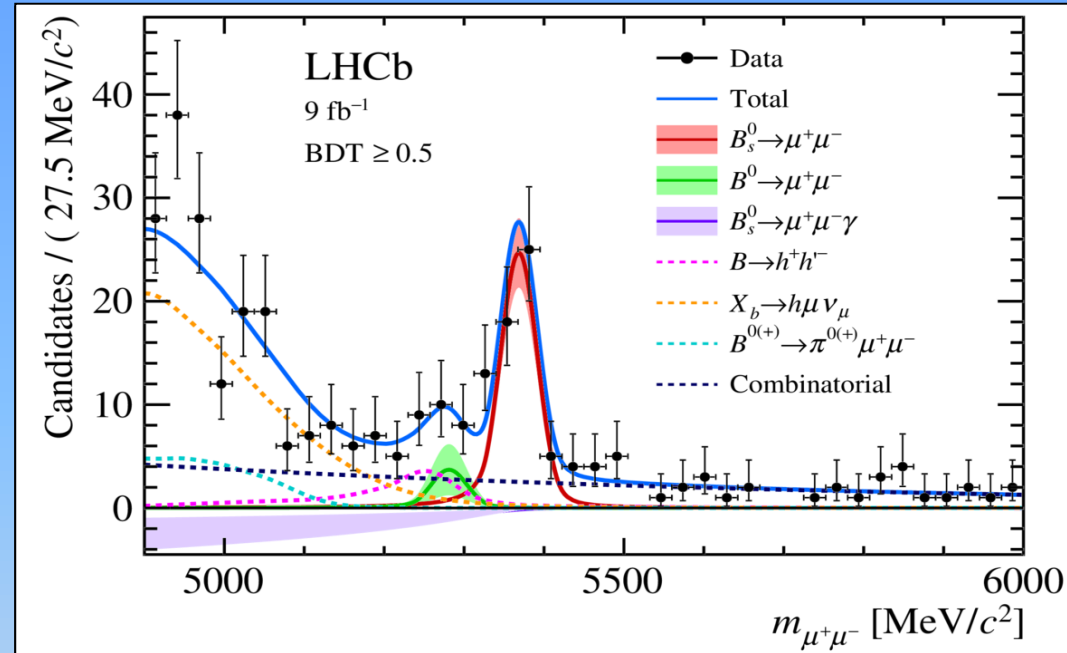


Slechts 3 op de miljard B deeltjes vervalst naar 2 muonen

Bestaan er nieuwe deeltjes?



LHCb: bestuderen van B deeltje



Phys. Rev. Lett. 128, (2022) 041801

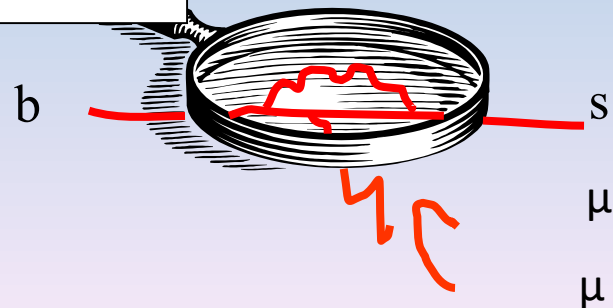
$$\mathcal{B}(B_s^0 \rightarrow \mu^+ \mu^-) = (3.09^{+0.46+0.15}_{-0.43-0.11}) \times 10^{-9}$$

$$\mathcal{B}(B^0 \rightarrow \mu^+ \mu^-) = (1.2^{+0.8}_{-0.7} \pm 0.1) \times 10^{-10}$$

Theory:

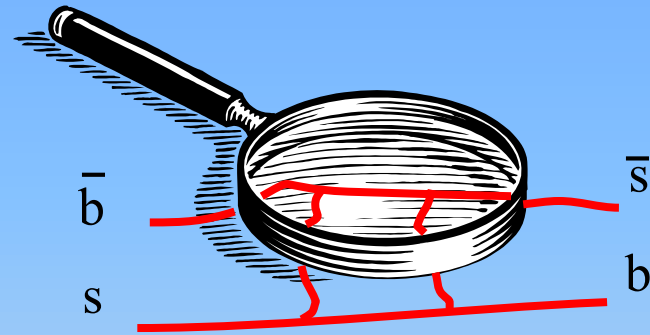
$$\mathcal{B}(B_s^0 \rightarrow m^+ m^-) = (3.65 \pm 0.23) \times 10^{-9}$$

$$\mathcal{B}(B^0 \rightarrow m^+ m^-) = (1.06 \pm 0.09) \times 10^{-10}$$

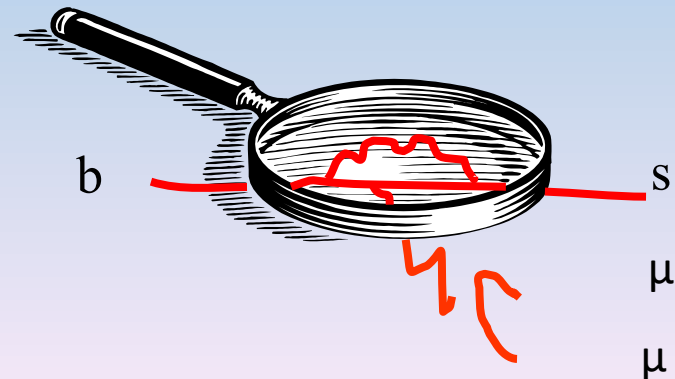


LHCb: bestuderen van B deeltje

1) Vind verschillen tussen materie en anti-materie



2) Vind nieuwe deeltjes



LHCb: highlights

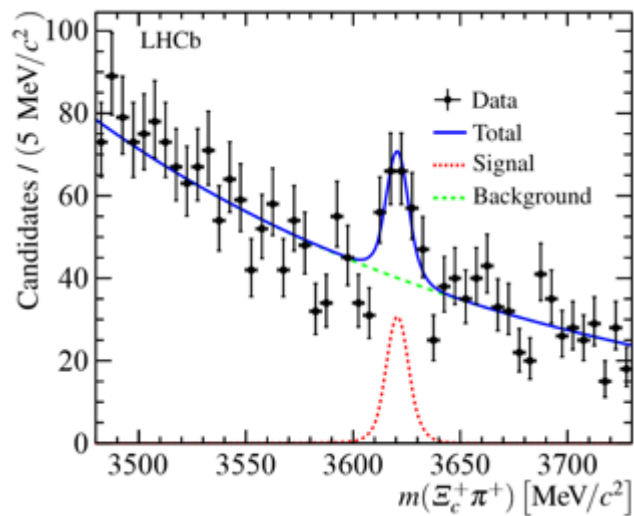
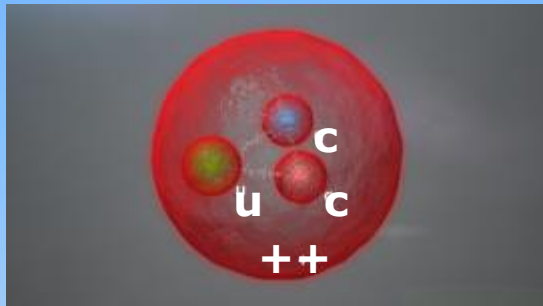
- 1) Nieuwe 'gewone' hadronen
- 2) Nieuwe 'exotische' hadronen: Tetraquark en pentaquark
- 3) Ontdekking 'CP schending' B_s
- 4) Ontdekking 'CP schending' charm

Hot topic:

- 5) Verschil electron, muon, tau??

LHCb: nieuwe 'gewone' hadronen

(ccu): Ξ_{cc}^{++}



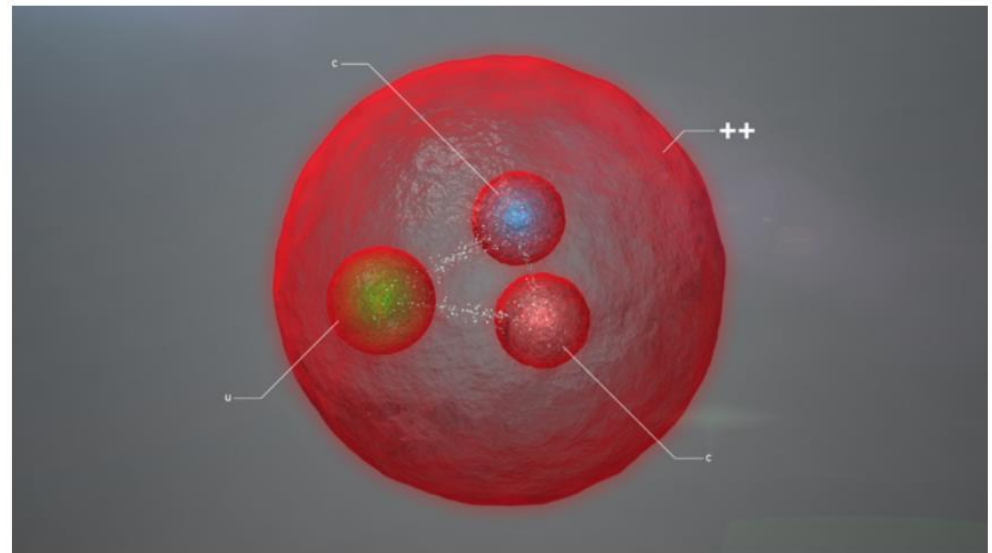
NewScientist
IDEEEN DIE DE WERELD VERANDEREN

Nieuw zwaar deeltje legt sterke kernkracht op de pijnbank

12 juli 2017



Jacob Aron en Leah Crane



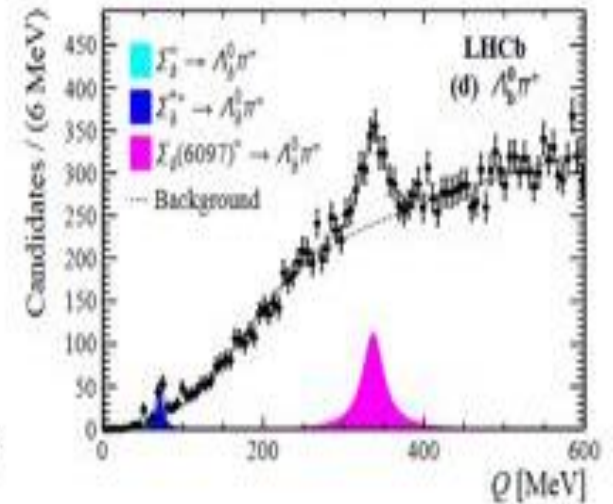
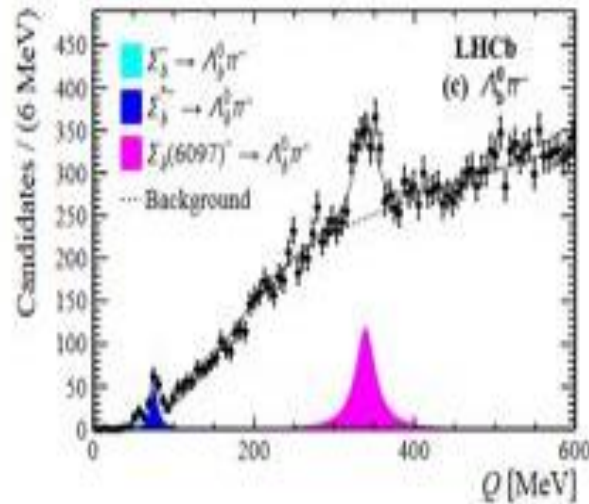
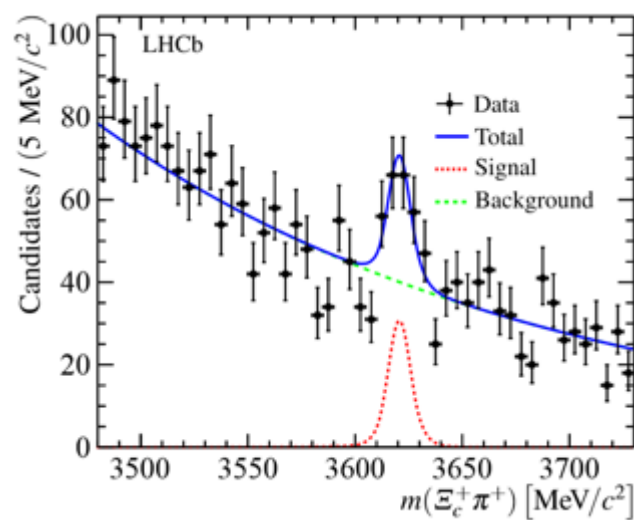
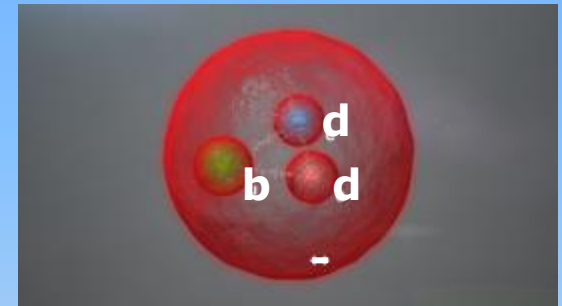
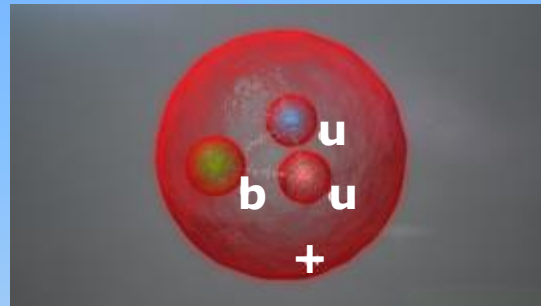
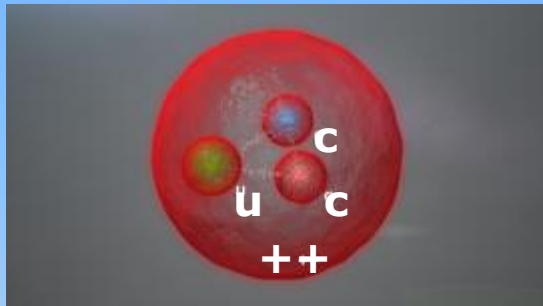
Het nieuwe deeltje bevat twee charm-quarks in het midden en een up-quark daaromheen. Beeld: Daniel Dominguez/CERN.

LHCb: nieuwe 'gewone' hadronen

(ccu): Ξ_{cc}^{++}

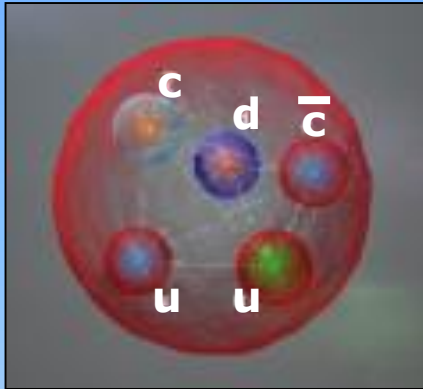
(buu): $\Sigma_b(6097)^+$

(bdd): $\Sigma_b(6097)^-$

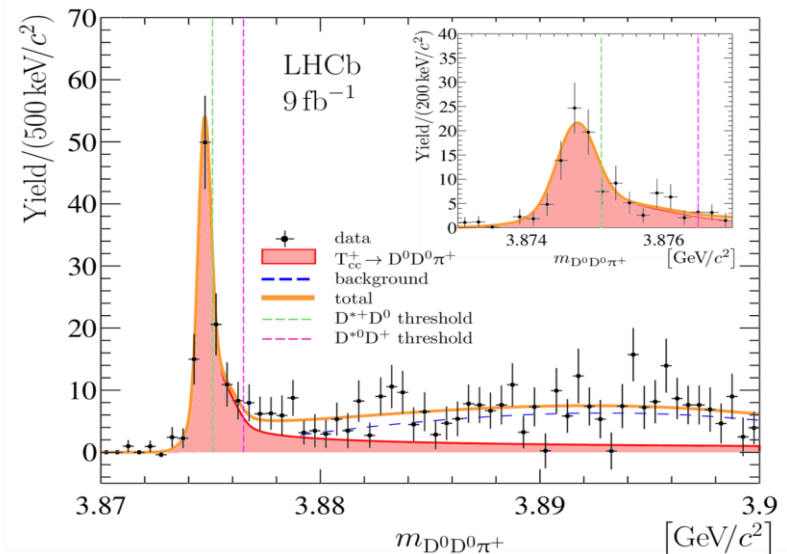
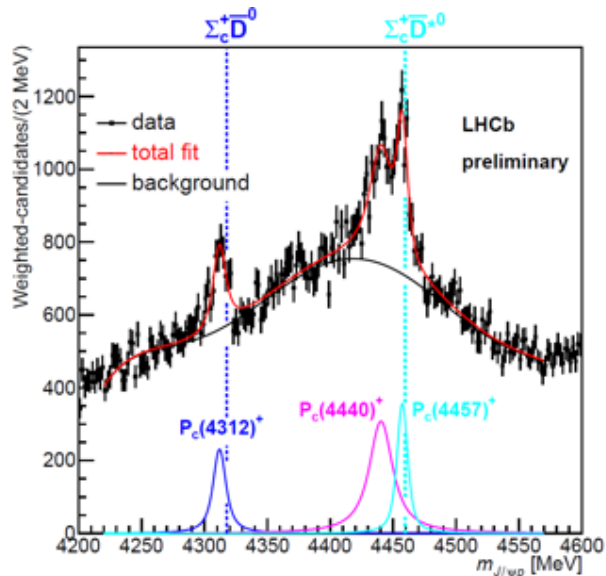
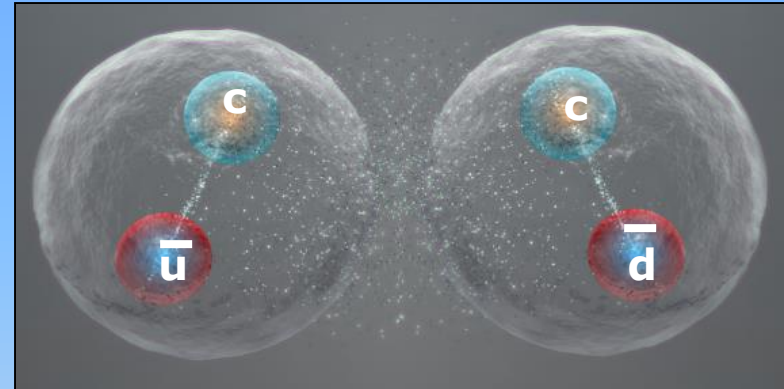


LHCb: nieuwe 'exotische' hadronen

$(c\bar{c}duu)$: $\mathbf{P}_c(4312)^+$

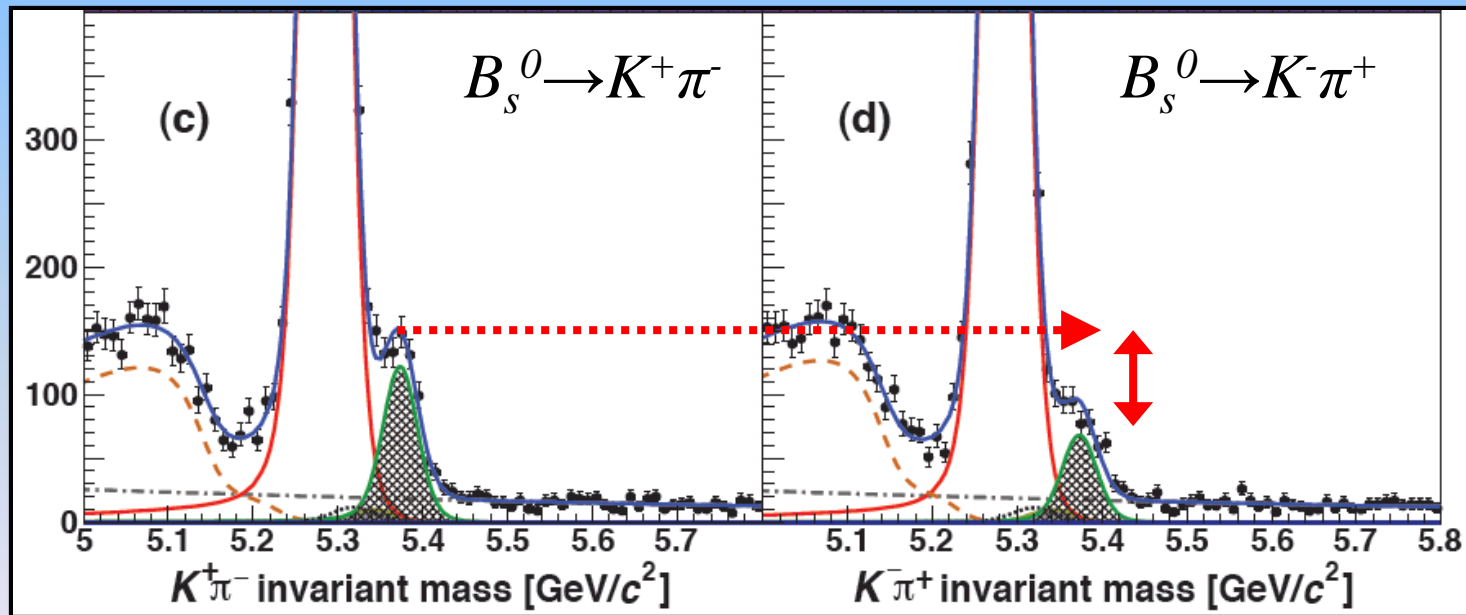
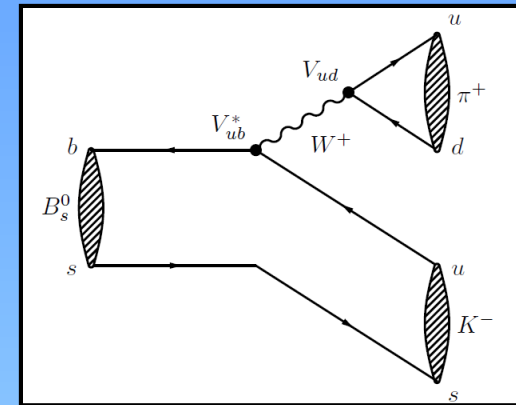


$(c\bar{u} c\bar{d})$: $\mathbf{T}_{cc}^+(3875)$



LHCb: antimaterie verschil in B_s^0

CP schending in B_s^0



LHCb: antimaterie verschil in charm

“CP schending”

$D^0 \rightarrow K^+ K^-$ same as

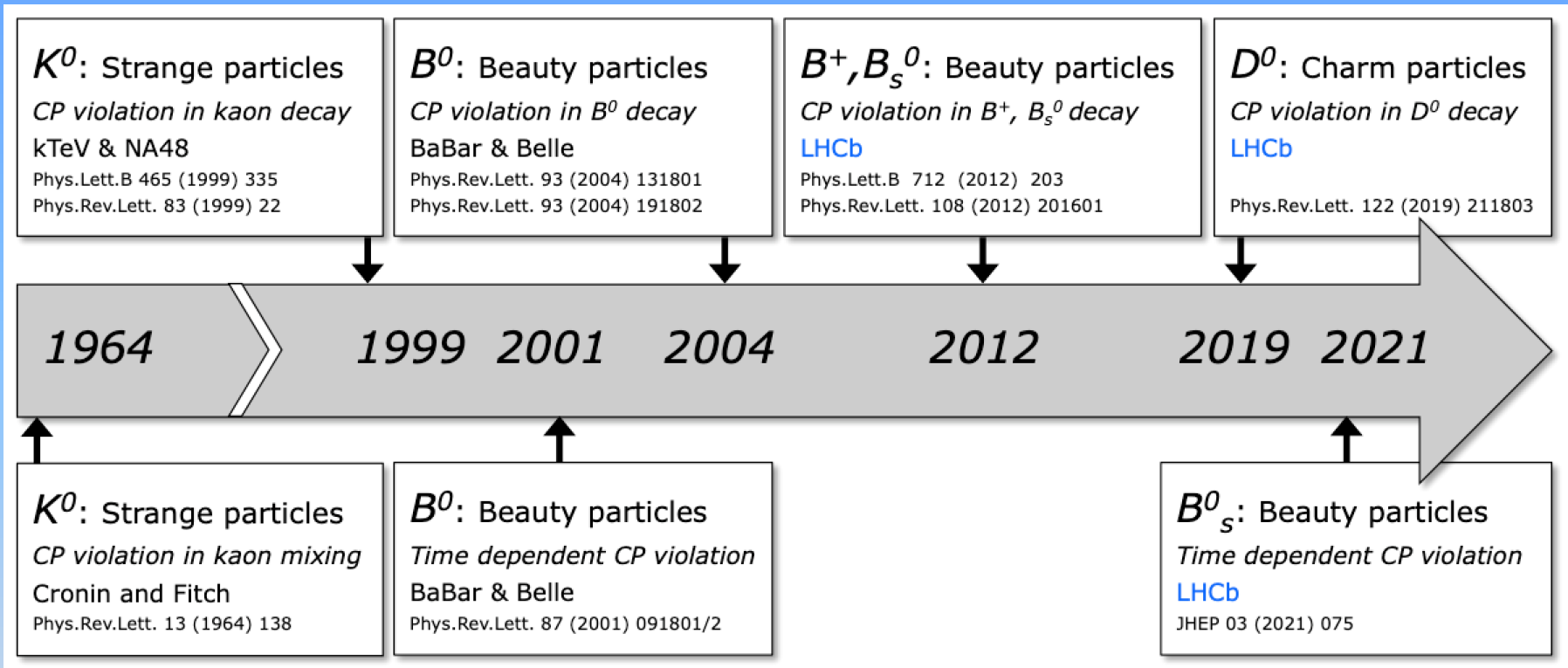
$\bar{D}^0 \rightarrow K^+ K^-$??

at least it is different compared to
 $D^0 \rightarrow \pi^+ \pi^- \dots$:

$$\Delta A_{CP} = (-15.4 \pm 2.9) \times 10^{-4}$$



LHCb: antimaterie verschillen



- | | | |
|------|---------------------------------------|-------------------|
| (ds) | 1964: CP schending met K^0 | (Nobelprijs 1980) |
| (bd) | 2000: CP schending met B^0 | (Nobelprijs 2008) |
| (bs) | 2012: CP schending met B_s^0 | (LHCb) |
| (cu) | 2019: CP schending met D^0 | (LHCb) |

LHCb: highlights

- 1) Nieuwe 'gewone' hadronen
- 2) Nieuwe 'exotische' hadronen: Tetraquark en pentaquark
- 3) Ontdekking 'CP schending' B_s
- 4) Ontdekking 'CP schending' charm

Hot topic:

- 5) Verschil electron, muon, tau?

LHCb: hot topic

NewScientist
IDEEËN DIE DE WERELD VERANDEREN

Cern vindt nieuwe hint voor scheurtjes in standaardmodel

19 april 2017



George van Hal

de Volkskrant

CERN is 'voorzichtig opgewonden' over subtiele verschillen in deeltjeswereld

Een gevoel van 'voorzichtige opwinding' heeft zich meester gemaakt van deeltjesfysici van CERN in Genève. Dinsdag maakte de LHCb-detector daar bekend subtiele verschillen te zien tussen bepaalde deeltjes. De gangbare deeltjestheorie neemt aan dat deeltjes in essentie identiek zijn.

Martijn van Calmthout 19 april 2017, 21:29



Wellicht is de deeltjeswereld niet zo democratisch als vooraf gedacht werd. Beeld epa



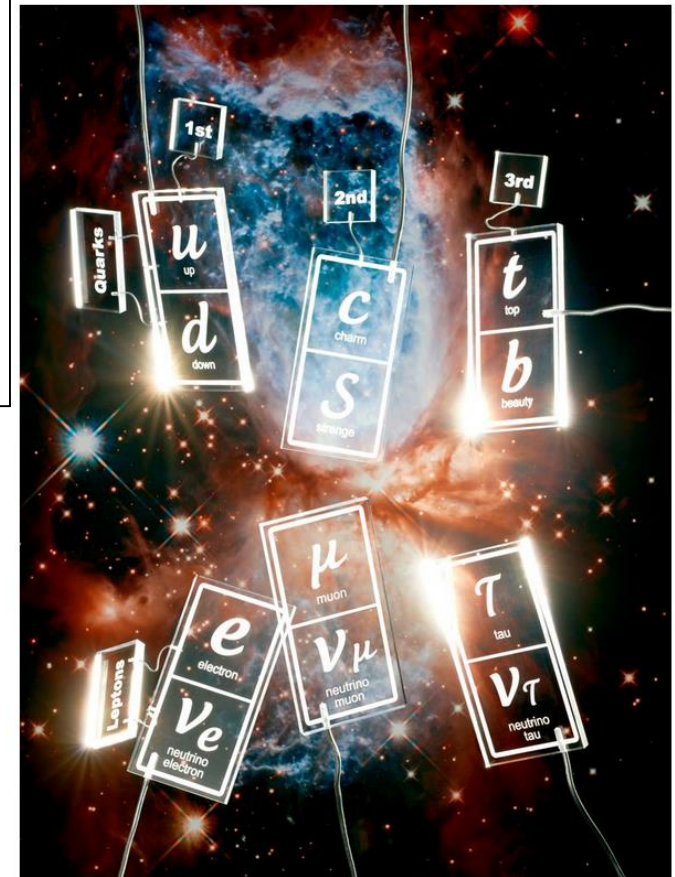
mogelijke hint dat er meer is dan alleen het rdmodel. Beeld: Cern.

de Volkskrant

Moeder aller deeltjes: de zoektocht naar de leptoquark

Is het fundamenteelste deeltje in het universum altijd over het hoofd gezien? Komende week kan de wereld opgeschud worden, als natuurkundigen in Seoul hun resultaten bekendmaken. Leptoquark, onthoud dat woord.

Martijn van Calmthout 29 juni 2018, 11:25



Beeld Rein Janssen

LHCb: hot topic

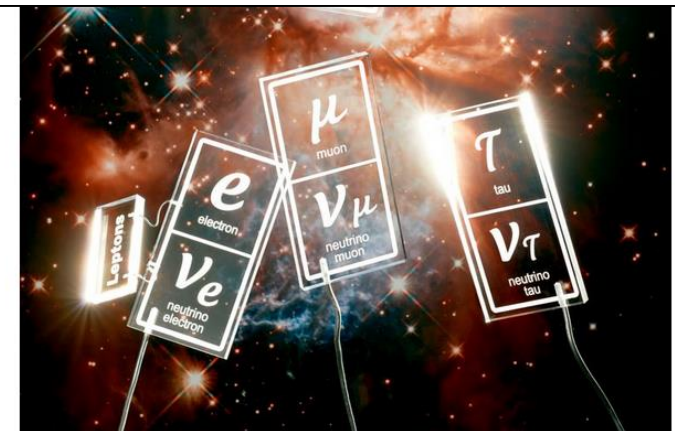
de Volkskrant

Moeder aller deeltjes: de zoektocht naar de leptoquark

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Martijn van Calmthout 29 juni 2018, 11:25

Maar de LHCb-metingen geven al jaren kleine hints dat er iets mis is met deze keurige lepton-universaliteit. En dat elektronen en muonen ergens diep van binnen toch net iets anders met quarks omgaan.



Beeld Rein Janssen

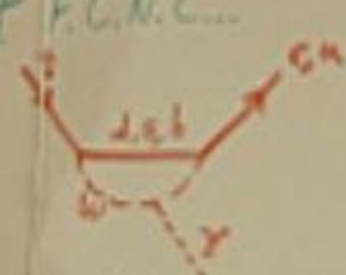
$t \rightarrow W^+ b$

$$BR(t \rightarrow W^+ b) = \frac{\Gamma(t \rightarrow W^+ b)}{\Gamma(t \rightarrow W^+ b) + \Gamma(t \rightarrow W^+ c) + \Gamma(t \rightarrow W^+ s)}$$

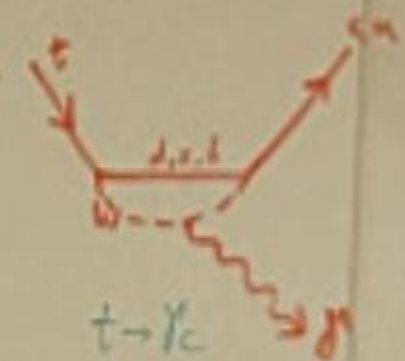


$$= \frac{|V_{cb}|^2}{|V_{cb}|^2 + |V_{cb}|^2 + |V_{cb}|^2}$$
$$\approx \frac{(0.9945)^2}{(0.9945)^2 + (0.0077)^2 + (0.07745)^2}$$
$$= 99.82\%$$

but F.C.N.C...



$t \rightarrow Z c$
 $t \rightarrow Z b$



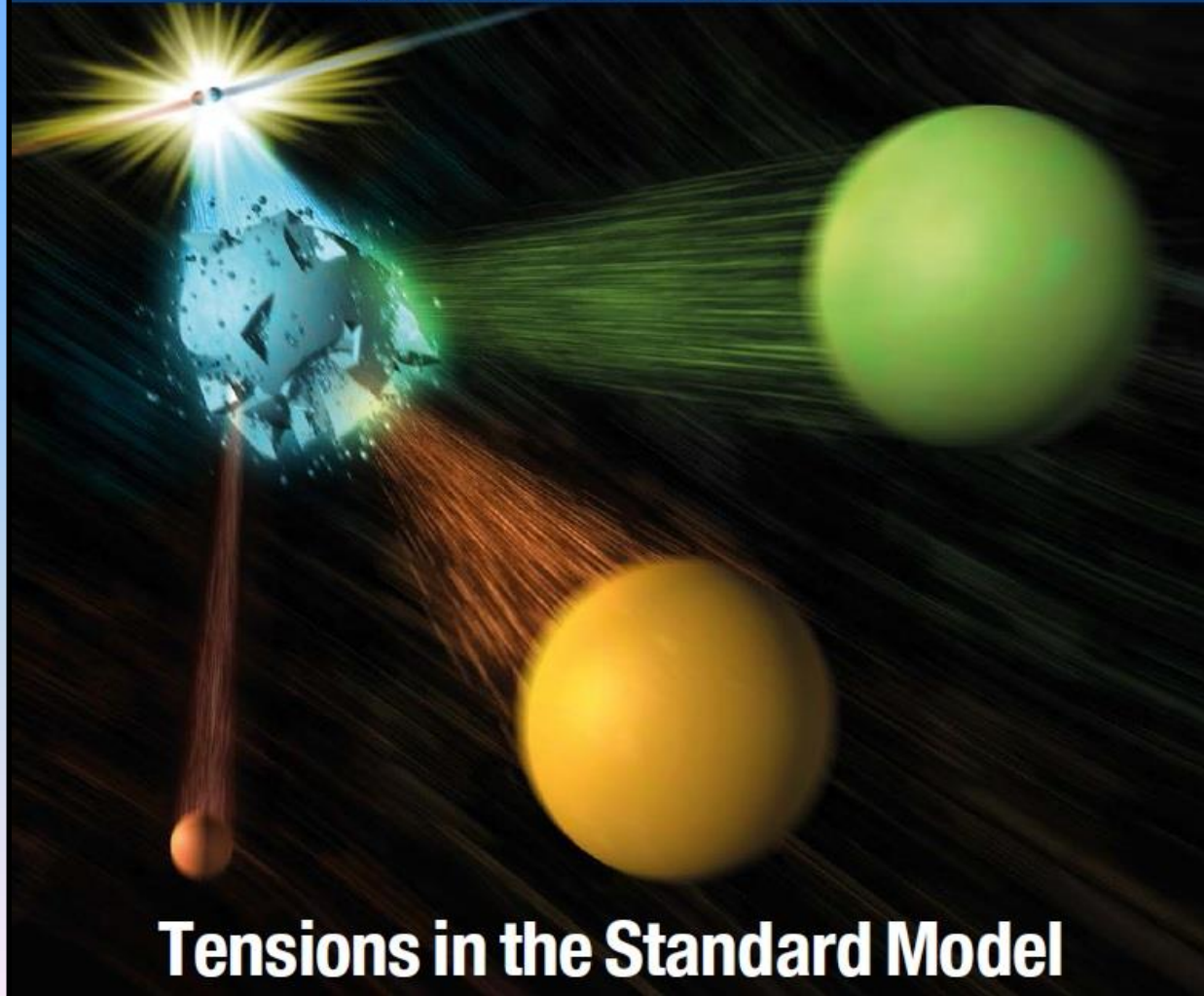
$t \rightarrow Y c$
 $t \rightarrow Y b$

$$U_{CKM} = \begin{pmatrix} c_{12}c_{13} & & \\ -s_{12}c_{13} & -c_{12}s_{13} & s_{13}e^{i\delta} \\ s_{12}c_{13} & c_{12}s_{13} & s_{13}e^{i\delta} \end{pmatrix}$$

INTERNATIONAL JOURNAL OF HIGH-ENERGY PHYSICS

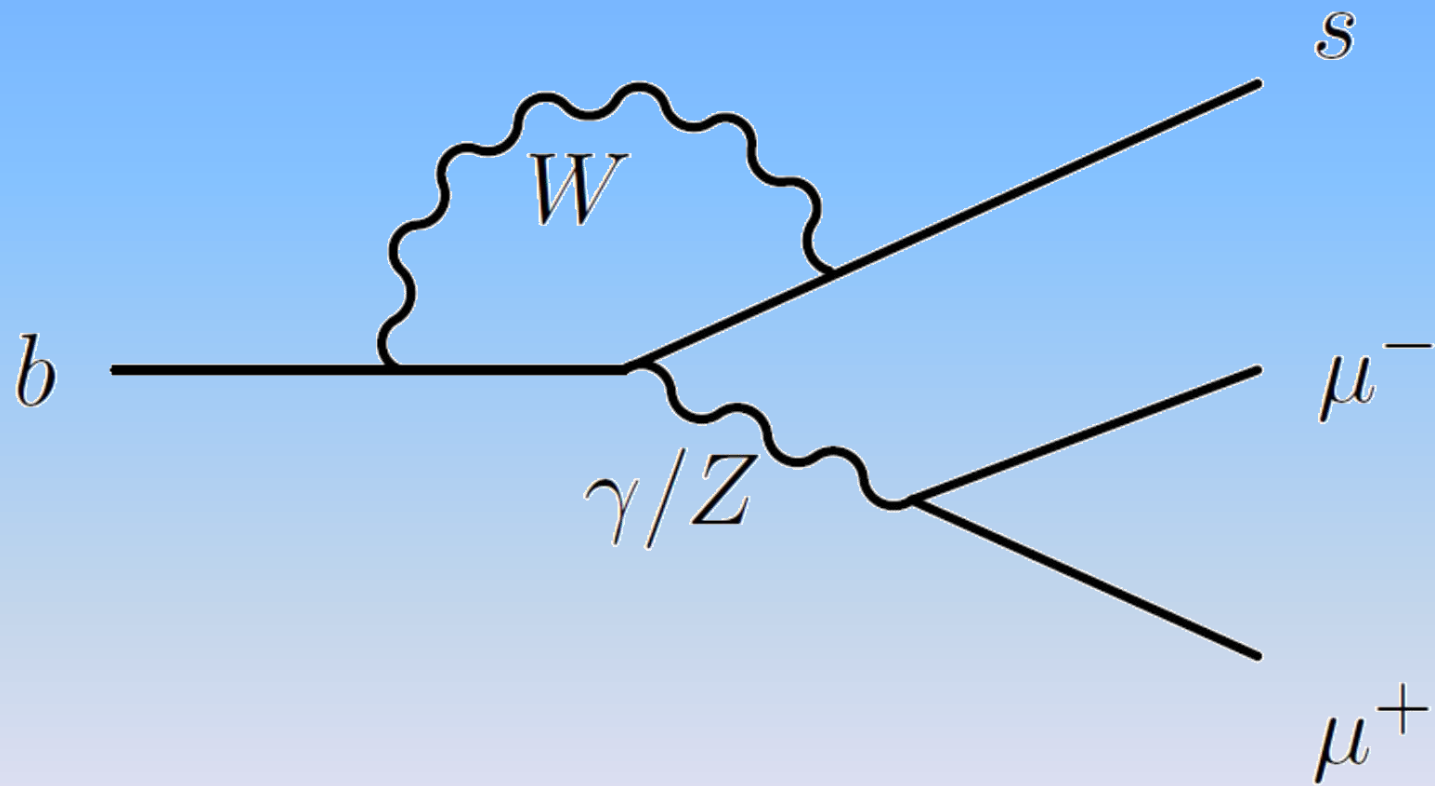
CERN COURIER

VOLUME 55 NUMBER 9 NOVEMBER 2015

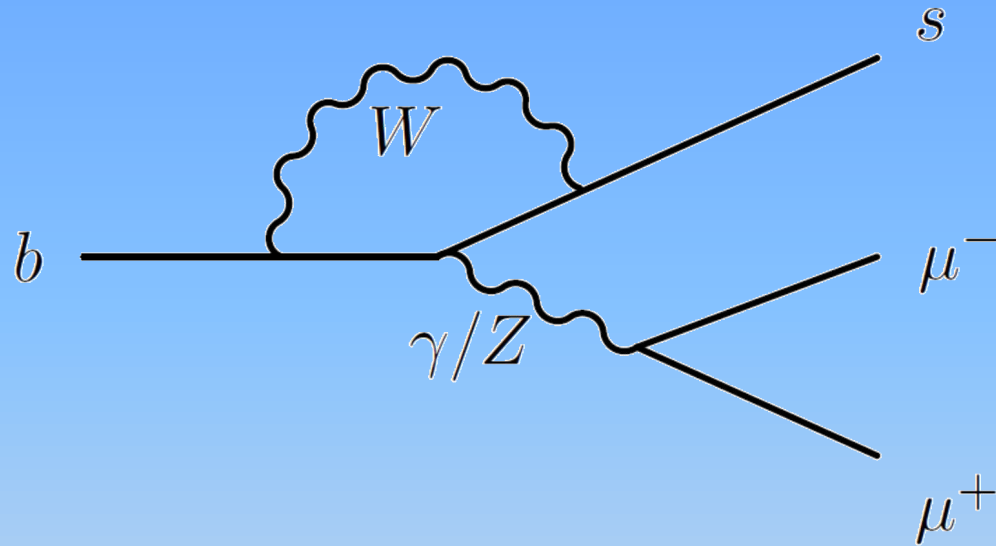


Tensions in the Standard Model

LHCb: hot topic



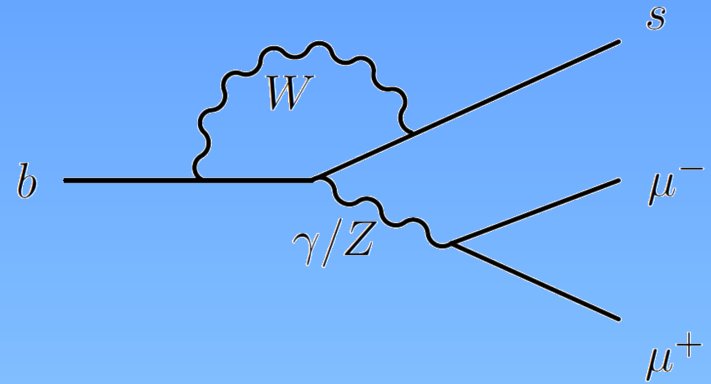
LHCb: hot topic



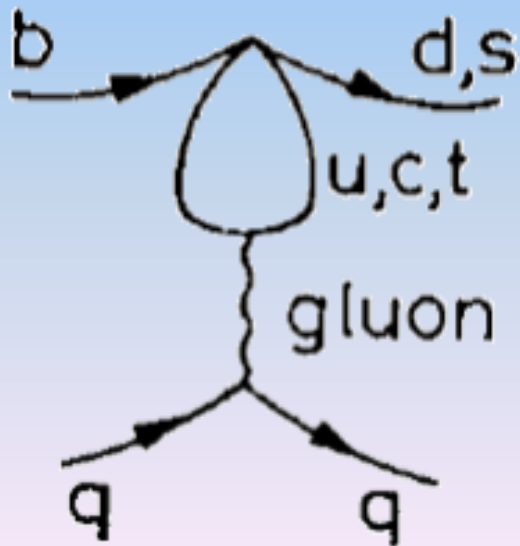
Flavour changing neutral current electroweak penguin

FCNC EWP

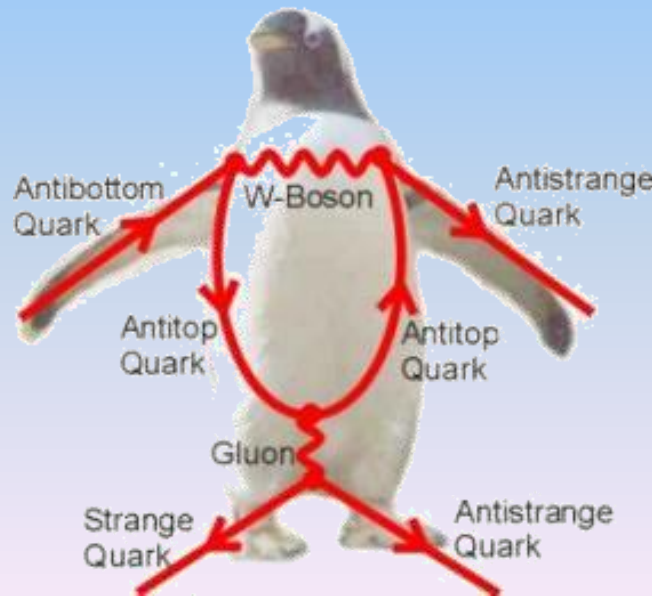
LHCb: hot topic



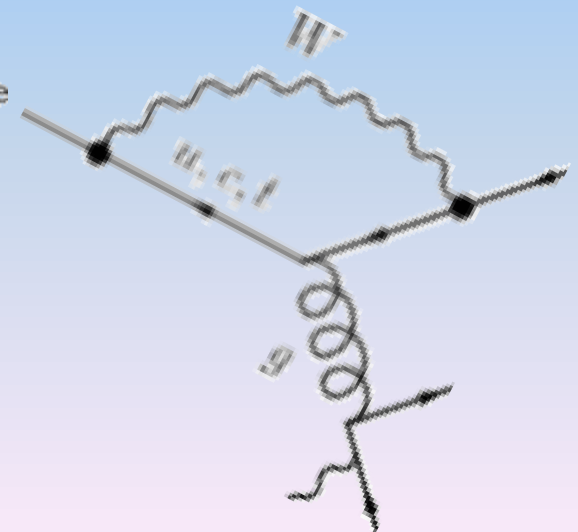
The original penguin:



A real penguin:

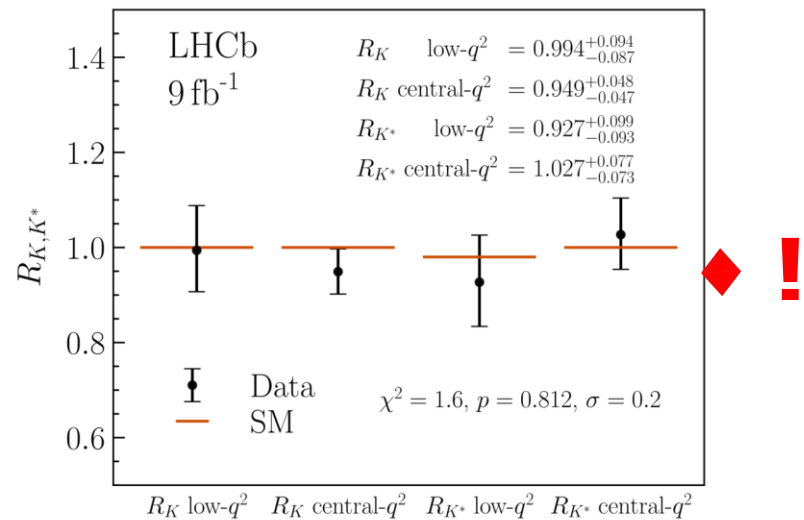
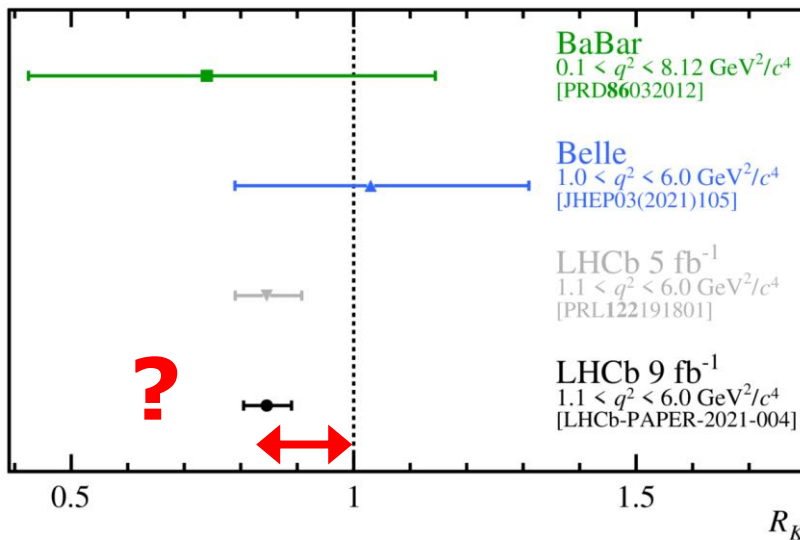
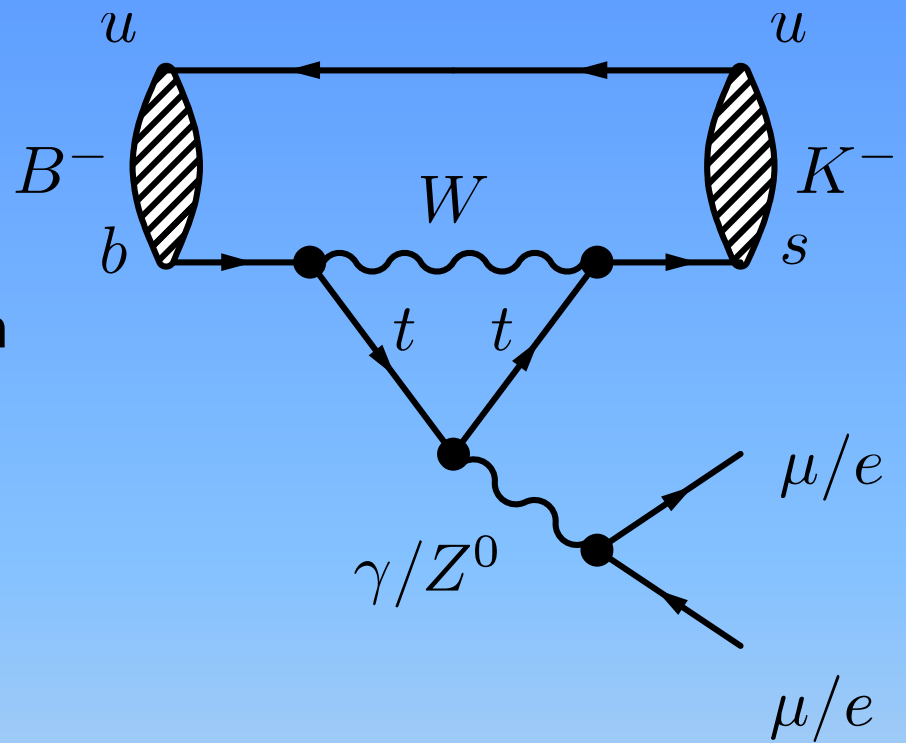


Our penguin:



LHCb: hot topic

Electronen en muonen gedragen zich anders?

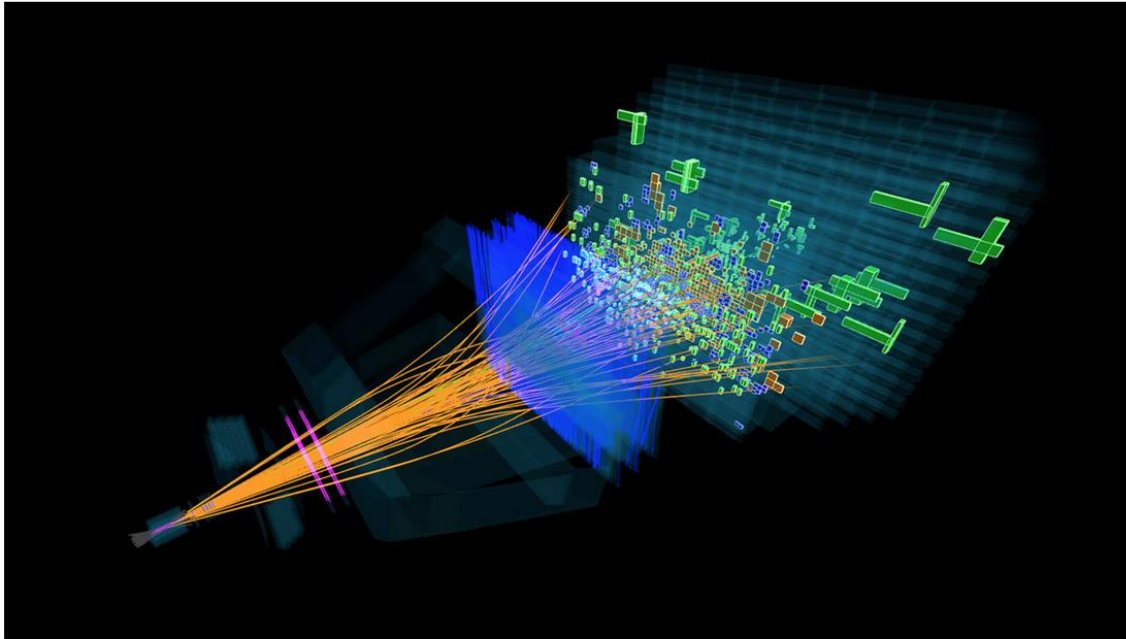


Muon gedraagt zich toch best normaal – of is er meer aan de hand?

Nieuws

Jean-Paul Keulen 29-12-2022 15:00:00

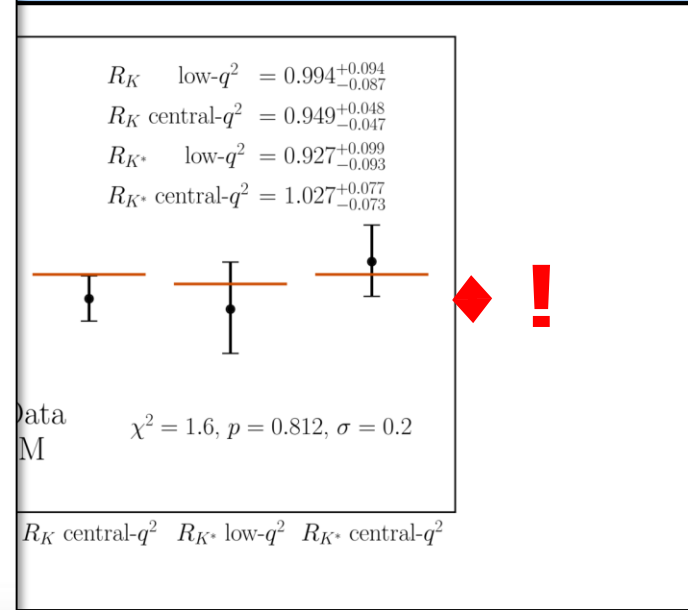
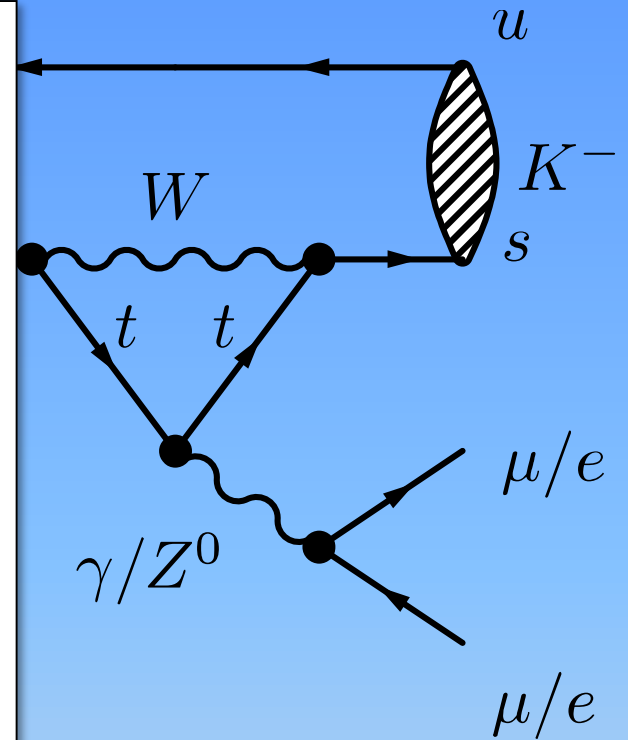
Deel dit artikel: [f](#) [t](#) [p](#) [e](#)



Vervallend B-deeltje in LHCb. Beeld: LHCb/CERN

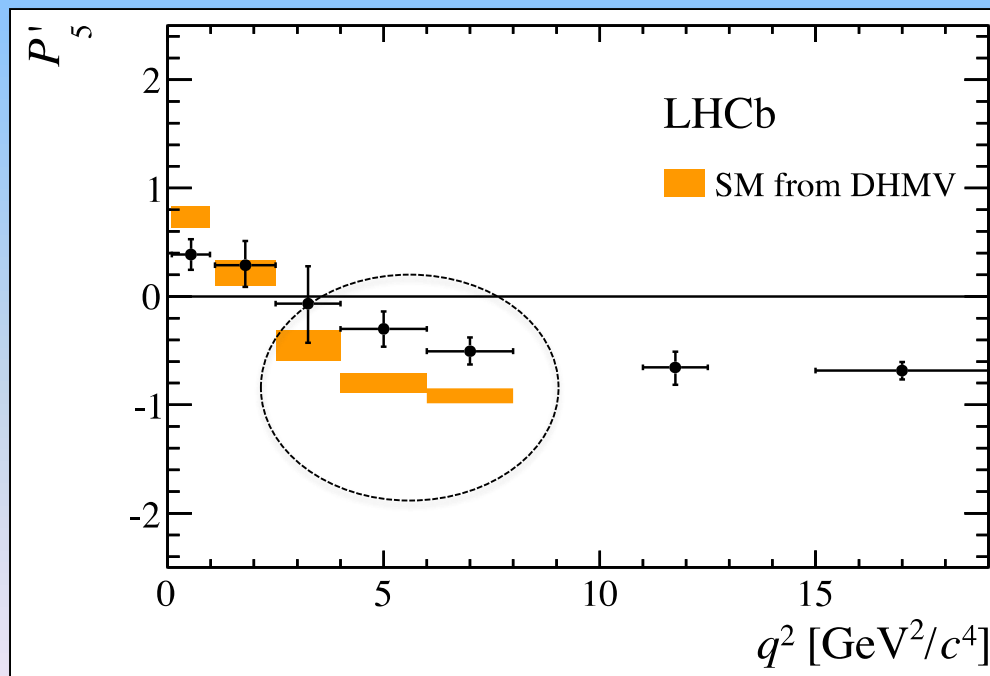
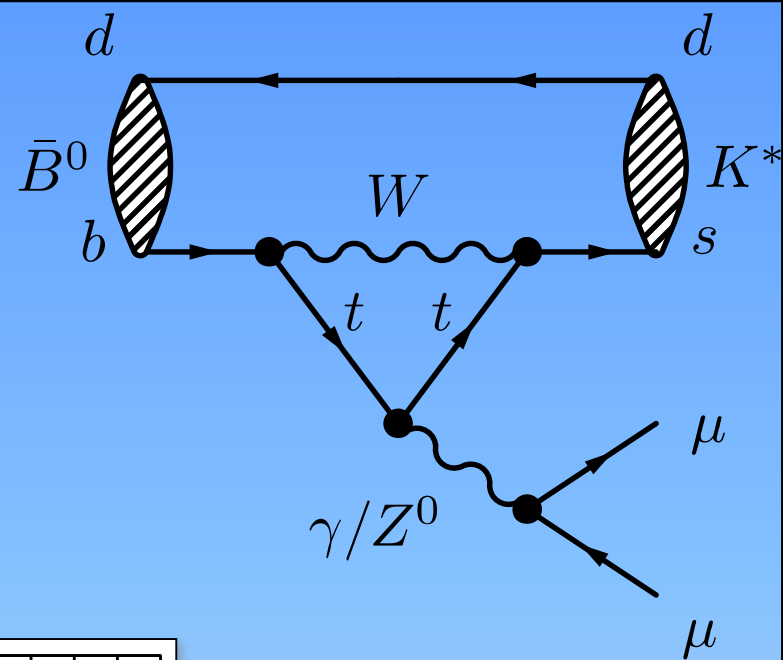
Een verrassend resultaat uit de deeltjesfysica lijkt te zijn afgeserveerd. Toch zijn er nog openstaande vragen.

Het gold als een van de interessantste resultaten binnen de deeltjesfysica sinds de ontdekking van het higgsdeeltje: het feit dat er bij het verval van bepaalde deeltjes **minder vaak muonen ontstaan dan je zou verwachten**. Zou die afwijking van onze huidige deeltjestheorie, blootgelegd met het deeltjesexperiment **LHCb**, wijzen op het bestaan van nieuwe deeltjes of nieuwe natuurkrachten?



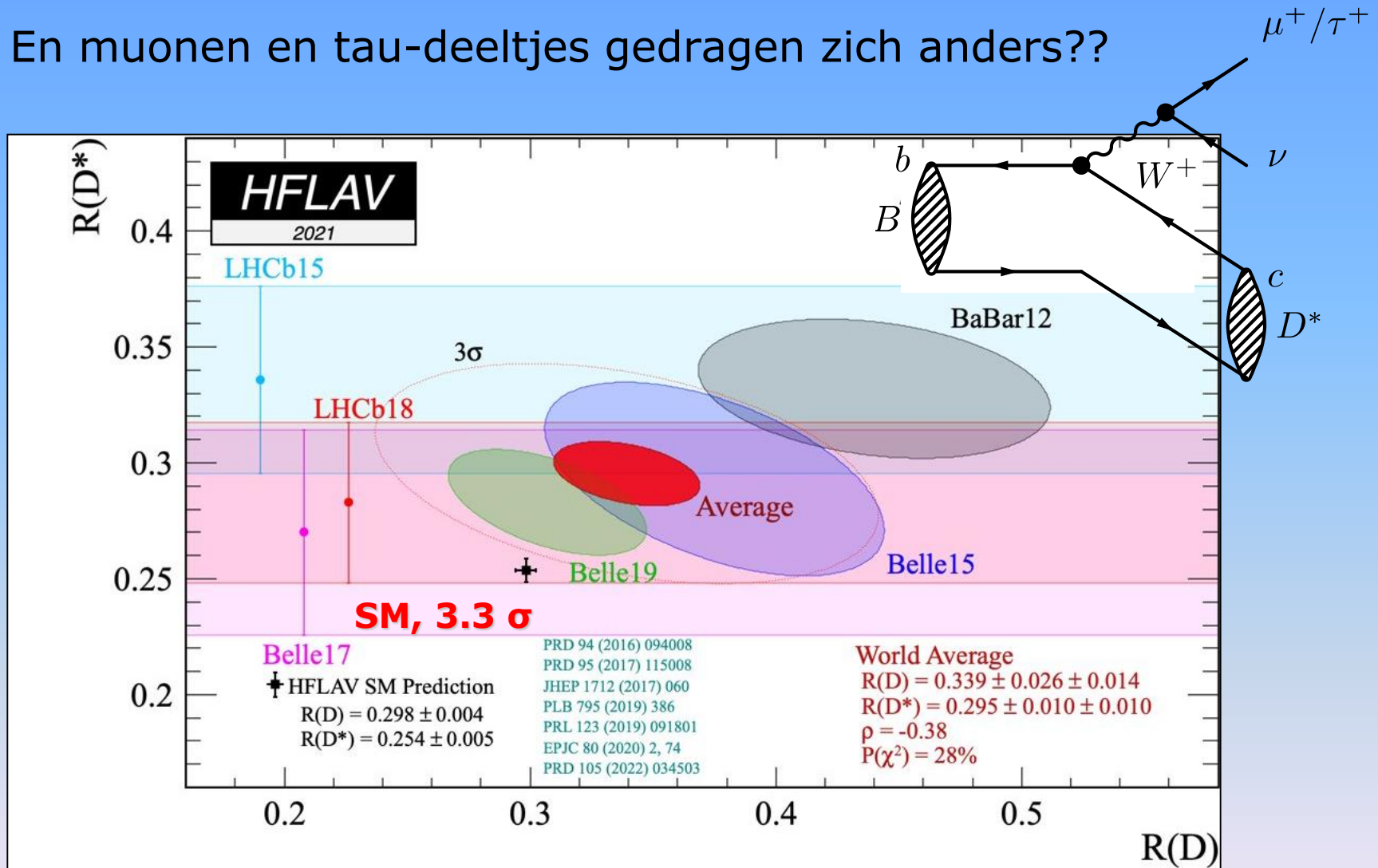
LHCb: hot topic

Ook hoekverdeling is anders...



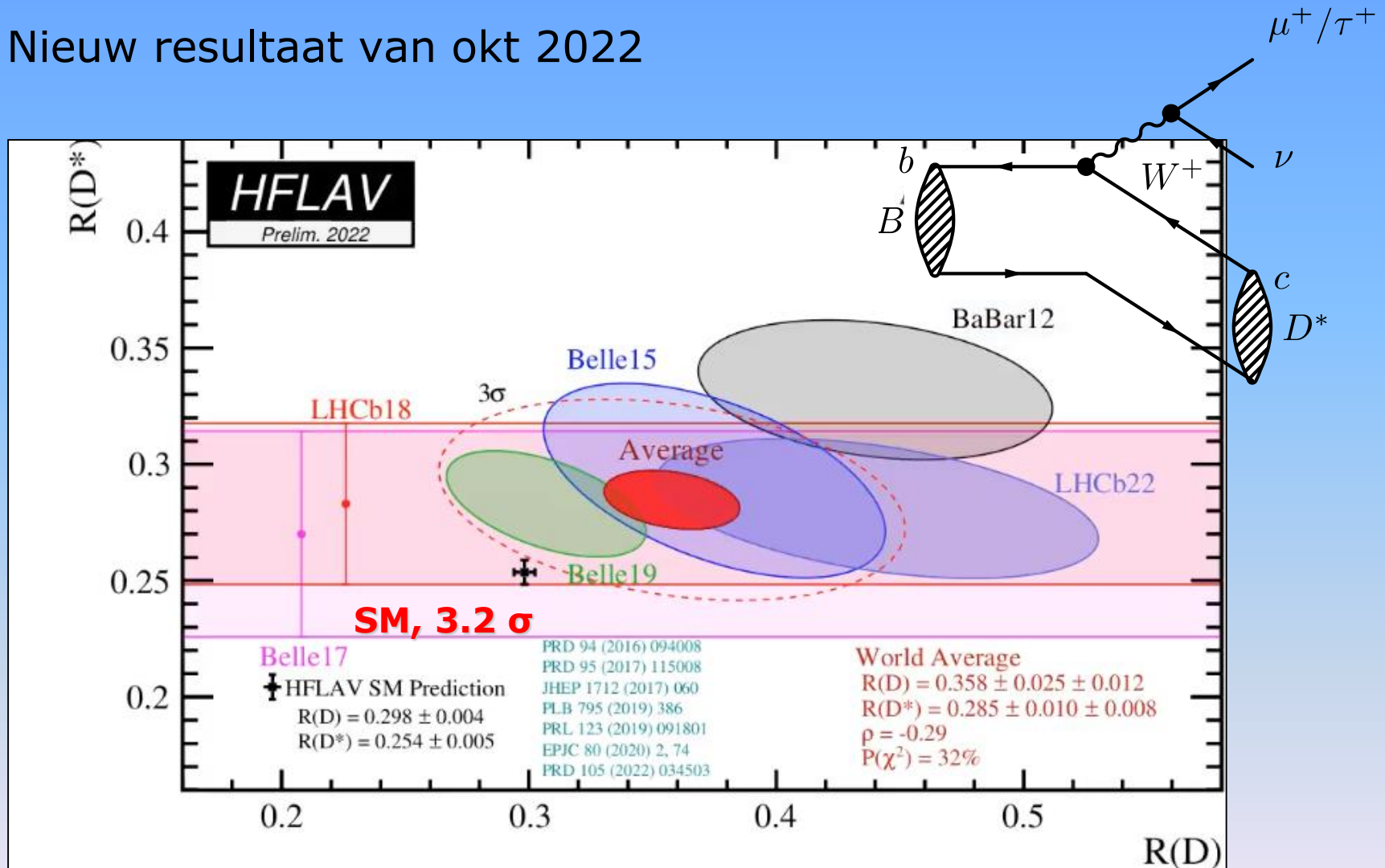
LHCb: hot topic

En muonen en tau-deeltjes gedragen zich anders??



LHCb: hot topic

Nieuw resultaat van okt 2022



LHCb: wat kan het zijn?

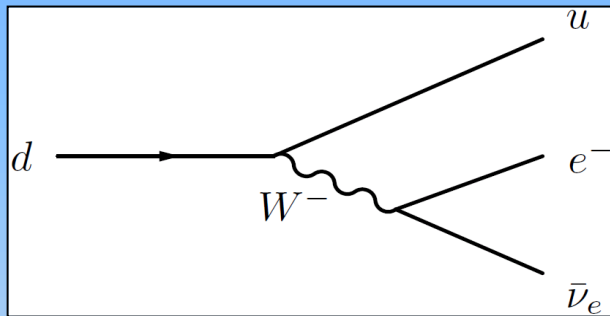
de Volkskrant

Moeder aller deeltjes: de zoektocht naar de leptoquark

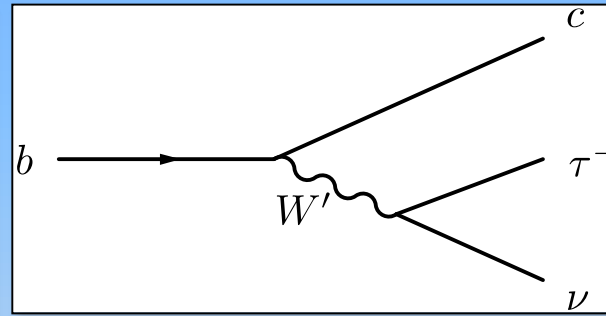
Is het fundamenteelste deeltje in het universum altijd over het hoofd gezien? Komende week kan de wereld opgeschud worden, als natuurkundigen in Seoul hun resultaten bekendmaken. Leptoquark, onthoud dat woord.

Martijn van Calmthout 29 juni 2018, 11:25

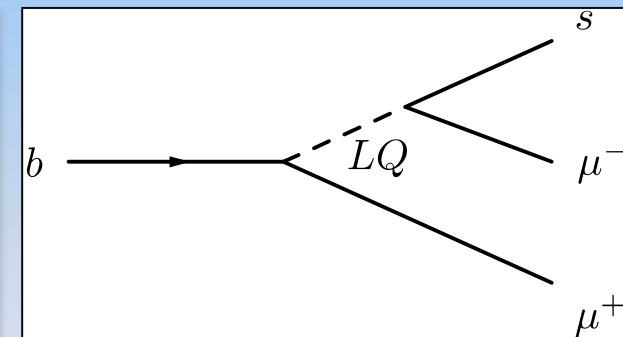
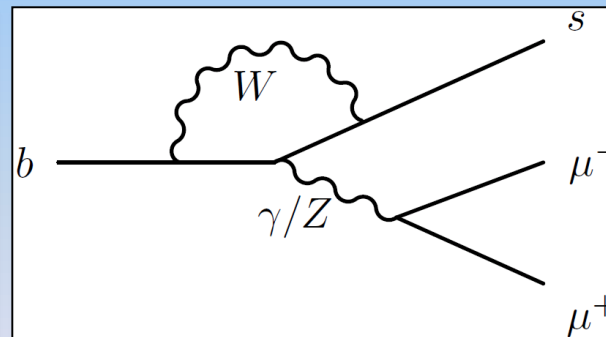
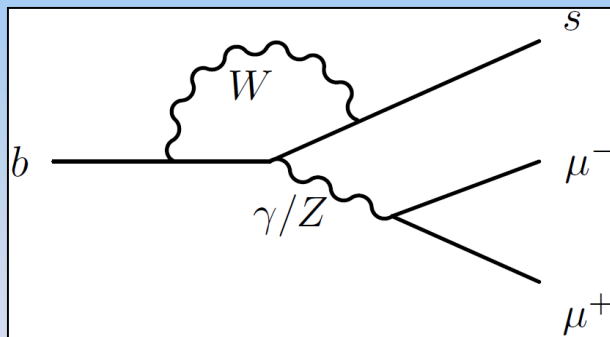
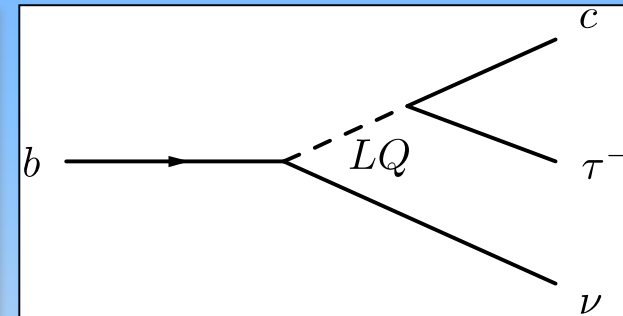
SM



SU(2)'



Leptoquark



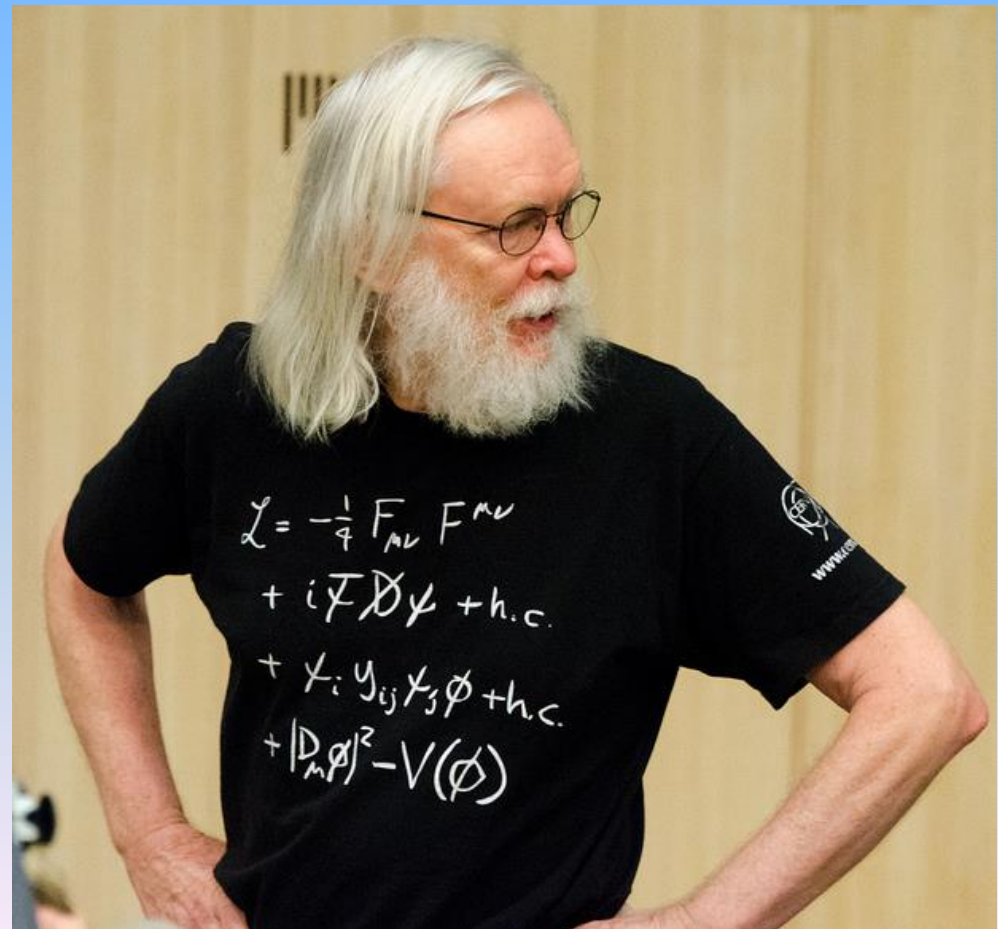
Leptoquark, onthoud dat woord.

M. Van Calmthout

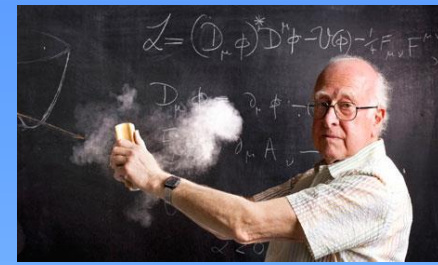
Dank!

Higgs en LHCb

Waarom is de Higgs zo bijzonder?



Waarom is de Higgs zo bijzonder?



Higgs heeft unieke rol in de wereld van elementaire deeltjes

ψ : "normale" deeltjes

ϕ : Higgs

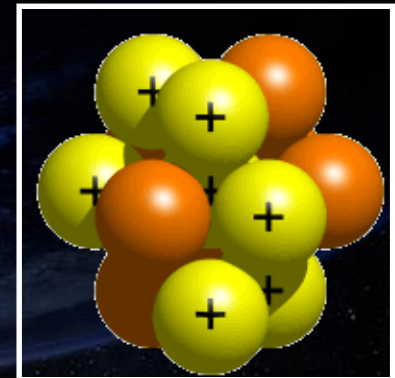
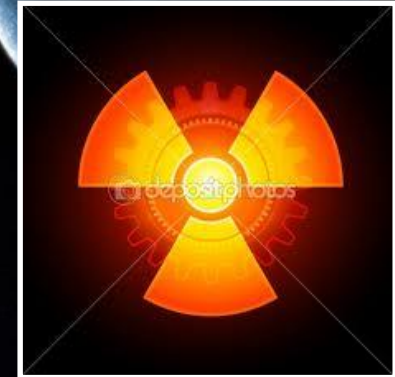
De helft van het T-shirt gaat over Higgs!

$$\begin{aligned}\mathcal{L} = & -\frac{1}{4} F_{\mu\nu} F^{\mu\nu} \\ & + i\bar{\psi} \not{D} \psi \\ & + \chi_i y_{ij} \chi_j \phi \\ & + |D_\mu \phi|^2 - V(\phi)\end{aligned}$$

Quarks	u up	c charm	t top
	d down	s strange	b bottom

Leptons	ν_e electron neutrino	ν_μ muon neutrino	ν_τ tau neutrino
	e electron	μ muon	τ tau

γ photon	Force carriers
Z Z boson	
W W boson	
g gluon	



12 deeltjes

4 krachten

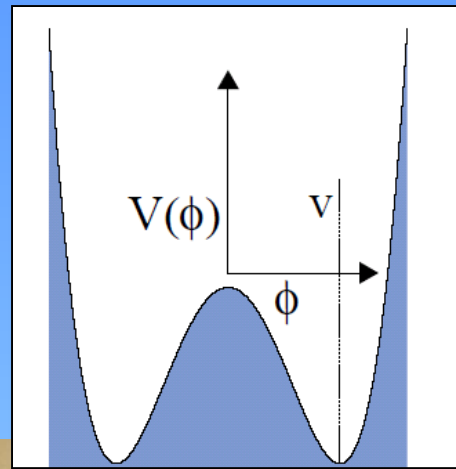
+ Higgs



Higgs en LHCb?

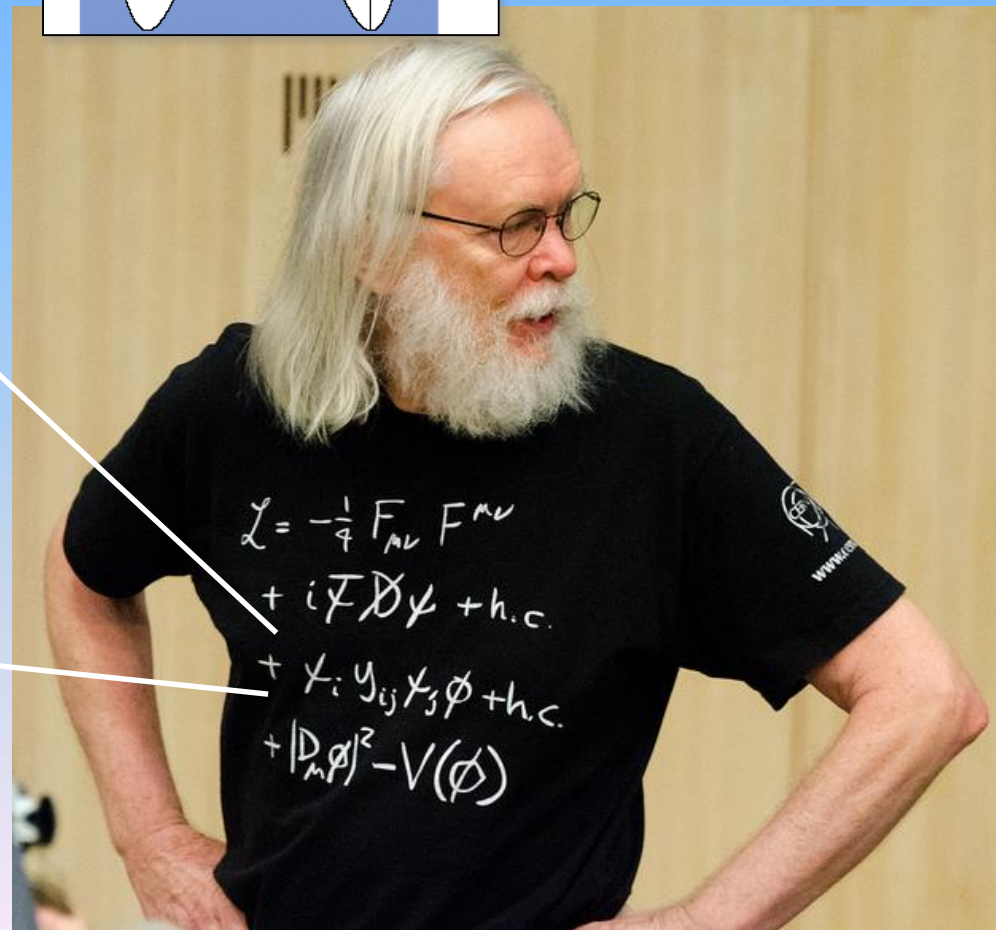
$$Y_{ij} y_i y_j f \rightarrow Y_{ij} y_i y_j (v + H) / \sqrt{2}$$

m: $Y_{ij} v$



$$\phi = \frac{v + h}{\sqrt{2}} e^{i\chi/v}$$

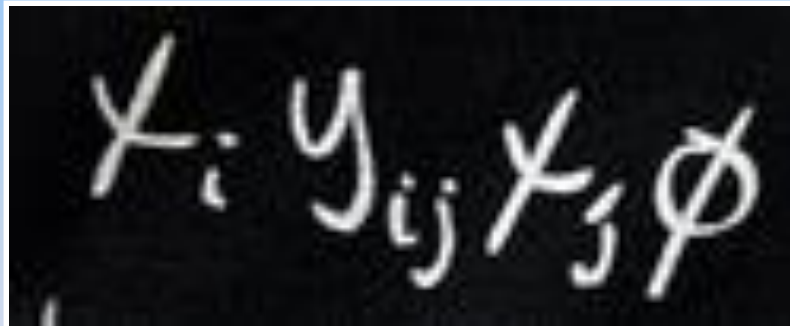
$$\mathcal{L} = y_{ij} \psi_i \psi_j \phi$$

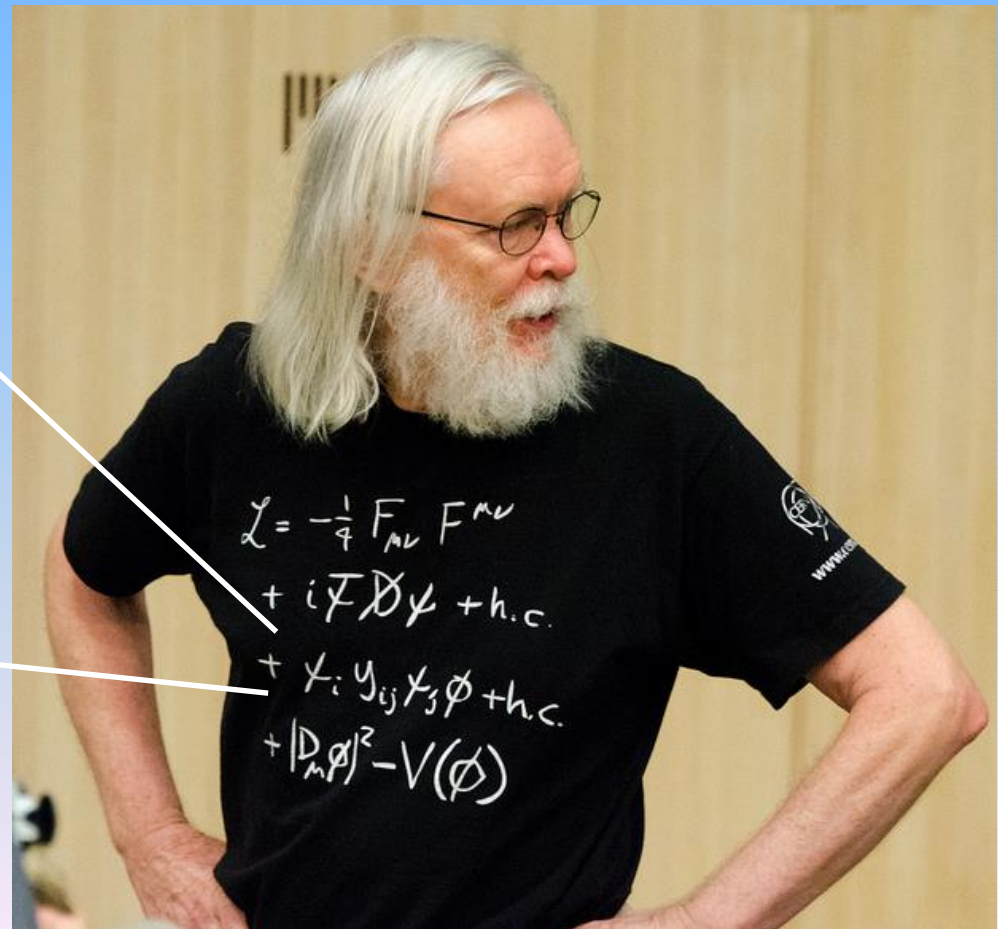


Higgs en LHCb?

Ψ : quarks

Y_{ij} : koppeling tussen verschillende quarks i, j

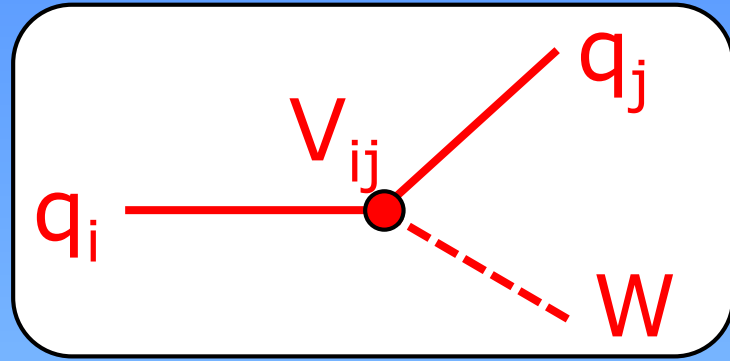

$$\chi_i y_{ij} \chi_j \phi$$



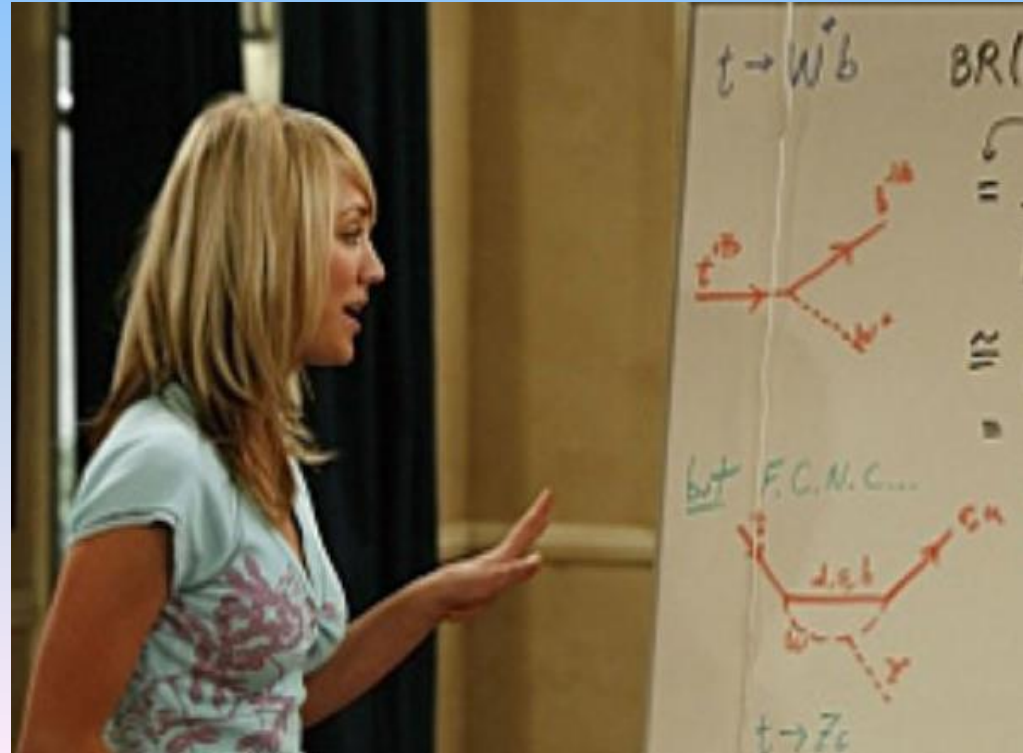
Higgs en LHCb?

Ψ : quarks

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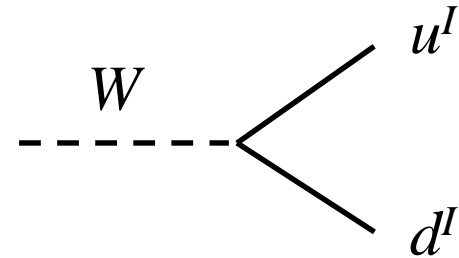
$$\mathcal{L} = \bar{\psi}_i Y_{ij} \psi_j + \dots$$



$$\mathcal{L}_{SM} = \mathcal{L}_{Kinetic} + \mathcal{L}_{Higgs} + \mathcal{L}_{Yukawa}$$

$$-\mathcal{L}_{Yuk} = Y_{ij}^d (\bar{u}_L^I, \bar{d}_L^I)_i \begin{pmatrix} \varphi^+ \\ \varphi^0 \end{pmatrix} d_{Rj}^I + \dots$$

$$\mathcal{L}_{Kinetic} = \frac{g}{\sqrt{2}} \bar{u}_{Li}^I \gamma^\mu W_\mu^- d_{Li}^I + \frac{g}{\sqrt{2}} \bar{d}_{Li}^I \gamma^\mu W_\mu^+ u_{Li}^I + \dots$$



- Diagonalize Yukawa matrix Y_{ij}
 - Mass terms
 - Quarks rotate
 - Off diagonal terms in charged current couplings

$$\begin{pmatrix} d^I \\ s^I \\ b^I \end{pmatrix} \rightarrow V_{CKM} \begin{pmatrix} d \\ s \\ b \end{pmatrix}$$

$$-\mathcal{L}_{Mass} = (\bar{d}, \bar{s}, \bar{b})_L \begin{pmatrix} m_d & & \\ & m_s & \\ & & m_b \end{pmatrix} \begin{pmatrix} d \\ s \\ b \end{pmatrix}_R + (\bar{u}, \bar{c}, \bar{t})_L \begin{pmatrix} m_u & & \\ & m_c & \\ & & m_t \end{pmatrix} \begin{pmatrix} u \\ c \\ t \end{pmatrix}_R + \dots$$

$$\mathcal{L}_{CKM} = \frac{g}{\sqrt{2}} \bar{u}_i \gamma^\mu W_\mu^- V_{ij} (1 - \gamma^5) d_j + \frac{g}{\sqrt{2}} \bar{d}_j \gamma^\mu W_\mu^+ V_{ij}^* (1 - \gamma^5) u_i + \dots$$

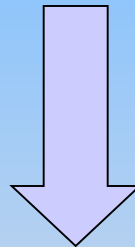
$$\mathcal{L}_{SM} = \mathcal{L}_{CKM} + \mathcal{L}_{Higgs} + \mathcal{L}_{Mass}$$

What do we know about the CKM matrix?

Magnitudes of elements have been measured over time

- Result of a *large* number of measurements and calculations

$$\begin{pmatrix} d' \\ s' \\ b' \end{pmatrix} = \begin{pmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{pmatrix} \begin{pmatrix} d \\ s \\ b \end{pmatrix}$$

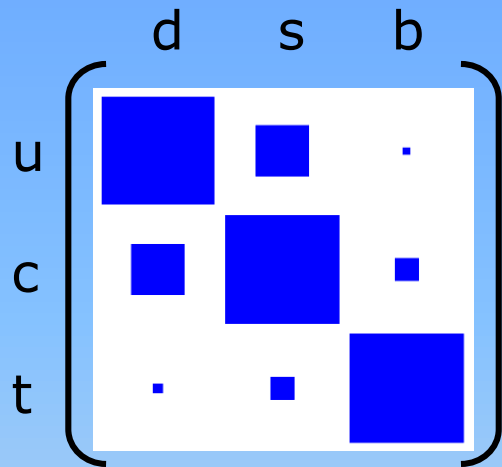


$$\begin{pmatrix} |V_{ud}| & |V_{us}| & |V_{ub}| \\ |V_{cd}| & |V_{cs}| & |V_{cb}| \\ |V_{td}| & |V_{ts}| & |V_{tb}| \end{pmatrix} = \begin{pmatrix} 0.97419 & 0.2257 & 0.00359 \\ 0.2256 & 0.97334 & 0.0415 \\ 0.00874 & 0.0407 & 0.999133 \end{pmatrix} \pm \begin{pmatrix} 0.00022 & 0.0010 & 0.00016 \\ 0.0010 & 0.00023 & 0.0011 \\ 0.00037 & 0.0010 & 0.000044 \end{pmatrix}$$

•Magnitude of elements shown only, no information of phase

Higgs en LHCb?

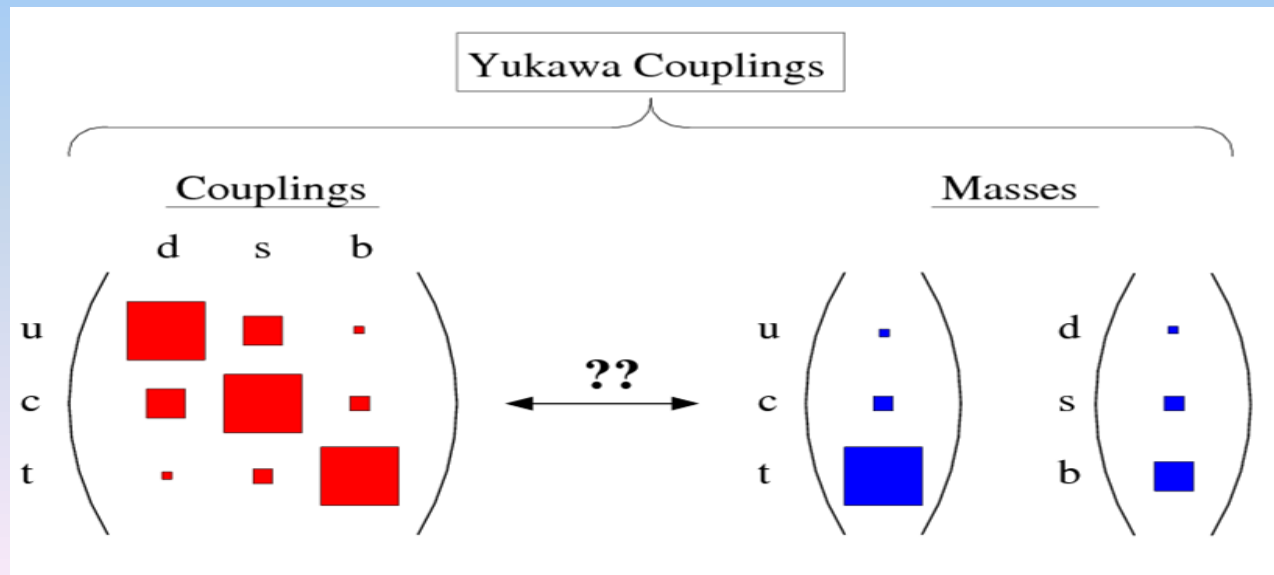
Quark koppelingen:



Waarom dit patroon in quark koppelingen?

Waarom dit patroon in quark massa's?

→ **Is er een verband?**



Intermezzo: How about the leptons?

the equivalent of the CKM matrix

– Pontecorvo-Maki-Nakagawa-Sakata matrix

$$\begin{bmatrix} \nu_e \\ \nu_\mu \\ \nu_\tau \end{bmatrix} = \begin{bmatrix} U_{e1} & U_{e2} & U_{e3} \\ U_{\mu1} & U_{\mu2} & U_{\mu3} \\ U_{\tau1} & U_{\tau2} & U_{\tau3} \end{bmatrix} \begin{bmatrix} \nu_1 \\ \nu_2 \\ \nu_3 \end{bmatrix} \quad \bullet \text{vs} \quad \begin{bmatrix} |d'\rangle \\ |s'\rangle \\ |b'\rangle \end{bmatrix} = \begin{bmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{bmatrix} \begin{bmatrix} |d\rangle \\ |s\rangle \\ |b\rangle \end{bmatrix}.$$

a completely different hierarchy!

$$U_{MNSP} \approx \begin{pmatrix} 0.85 & 0.53 & 0 \\ -0.37 & 0.60 & 0.71 \\ -0.37 & 0.60 & -0.71 \end{pmatrix}$$

$$V_{CKM} = \begin{pmatrix} 0.97428 & 0.2253 & 0.00347 \\ 0.2252 & 0.97345 & 0.0410 \\ 0.00862 & 0.0403 & 0.999152 \end{pmatrix}$$

Intermezzo: How about the leptons?

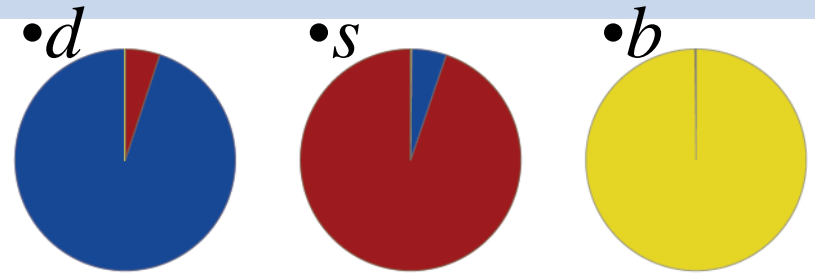
the equivalent of the CKM matrix

– Pontecorvo-Maki-Nakagawa-Sakata matrix

$$\begin{bmatrix} \nu_e \\ \nu_\mu \\ \nu_\tau \end{bmatrix} = \begin{bmatrix} U_{e1} & U_{e2} & U_{e3} \\ U_{\mu1} & U_{\mu2} & U_{\mu3} \\ U_{\tau1} & U_{\tau2} & U_{\tau3} \end{bmatrix} \begin{bmatrix} \nu_1 \\ \nu_2 \\ \nu_3 \end{bmatrix} \quad \bullet\text{vs} \quad \begin{bmatrix} |d'\rangle \\ |s'\rangle \\ |b'\rangle \end{bmatrix} = \begin{bmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{bmatrix} \begin{bmatrix} |d\rangle \\ |s\rangle \\ |b\rangle \end{bmatrix}.$$

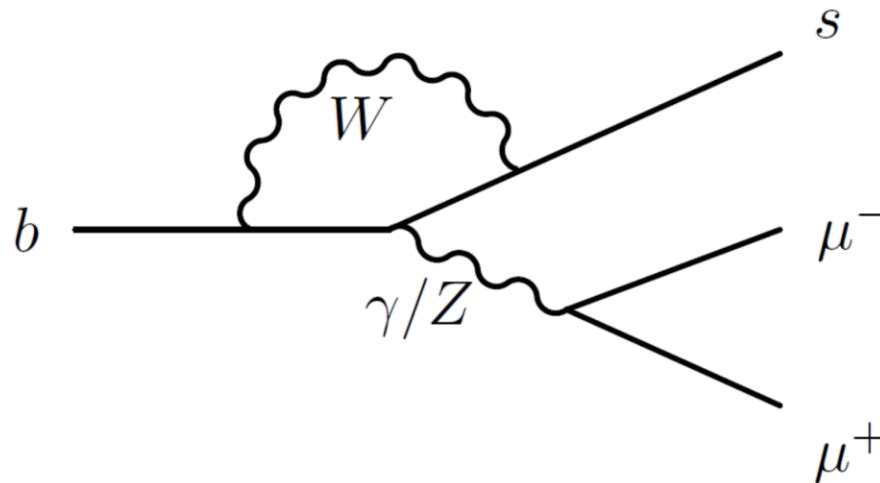
a completely different hi

$$\begin{pmatrix} |U_{e1}|^2 & |U_{e2}|^2 & |U_{e3}|^2 \\ |U_{\mu1}|^2 & |U_{\mu2}|^2 & |U_{\mu3}|^2 \\ |U_{\tau1}|^2 & |U_{\tau2}|^2 & |U_{\tau3}|^2 \end{pmatrix} \approx \begin{pmatrix} \frac{2}{3} & \frac{1}{3} & 0 \\ \frac{1}{6} & \frac{1}{3} & \frac{1}{2} \\ \frac{1}{6} & \frac{1}{3} & \frac{1}{2} \end{pmatrix}$$



Take home message

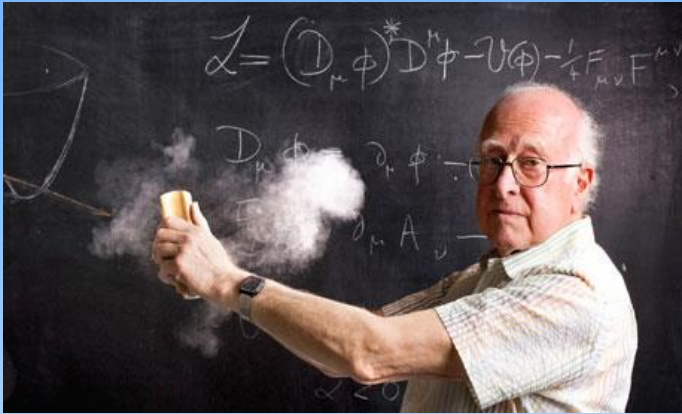
- 1) LHCb zoekt verschillen tussen materie en antimaterie
- 2) LHCb kan zeer zware deeltjes vinden (maar alleen *virtueel*)
- 3) Nieuwe deeltjes helpen om grote vragen te beantwoorden



Einde

Wat snappen we nog niet?

Massa van deeltjes



Bijzondere voorspelling:

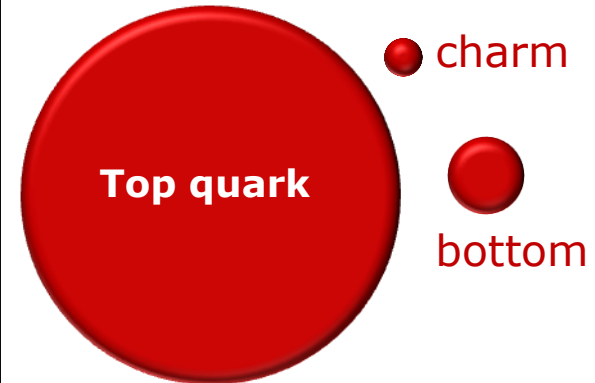
Het Higgs boson:

zorgt ervoor dat deeltjes massa kunnen hebben in de theorie

Neutrino's

- Elektron
- Muon
- Tau

● up, down, strange



Wat is massa ?? Anno 1687

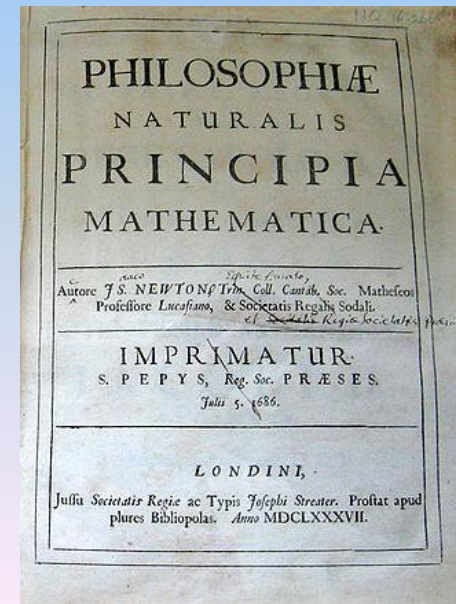
Massa is de 'wisselkoers' tussen kracht en versnelling:

$$F = m \times a$$

Beschrijft niet wat massa **is** ...



Newton



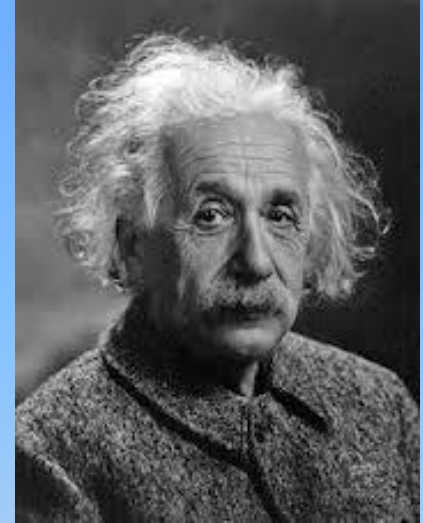
Wat is massa ?? Anno 1905

Massa is energie

$$E = m \times c^2$$

Beschrijft wel wat massa *is* !

Maar niet waar het vandaan komt ...



Einstein

13. *Ist die Trägheit eines Körpers von seinem Energieinhalt abhängig?*
von A. Einstein.

Die Resultate einer jüngst in diesen Annalen von mir publizierten elektrodynamischen Untersuchung¹⁾ führen zu einer sehr interessanten Folgerung, die hier abgeleitet werden soll. Ich legte dort die Maxwell-Hertz'schen Gleichungen für den leeren Raum nebst dem Maxwellschen Ausdruck für die elektromagnetische Energie des Raumes zugrunde und außerdem das Prinzip:

Die Gesetze, nach denen sich die Zustände der physikalischen Systeme ändern, sind unabhängig davon, auf welches von zwei relativ zueinander in gleichförmiger Parallel-Translationsbewegung befindlichen Koordinatensystemen diese Zustandsänderungen bezogen werden (Relativitätsprinzip).

Gestützt auf diese Grundlagen²⁾ leitete ich unter anderem das nachfolgende Resultat ab (l. c. § 8):

Ein System von ebenen Lichtwellen besitze, auf das Koordinatensystem (x, y, z) bezogen, die Energie l ; die Strahlrichtung (Wellennormale) bilde den Winkel φ mit der x -Achse des Systems. Führt man ein neues, gegen das System (x, y, z) in gleichförmiger Paralleltranslation begriffenes Koordinatensystem (ξ, η, ζ) ein, dessen Ursprung sich mit der Geschwindigkeit v längs der x -Achse bewegt, so besitzt die genannte Lichtmenge — im System (ξ, η, ζ) gemessen — die Energie:

$$l^* = l \frac{1 - \frac{v}{V} \cos \varphi}{\sqrt{1 - \left(\frac{v}{V}\right)^2}}$$

wobei V die Lichtgeschwindigkeit bedeutet. Von diesem Resultat machen wir im folgenden Gebrauch.

1) A. Einstein, Ann. d. Phys. 17, p. 891. 1905.

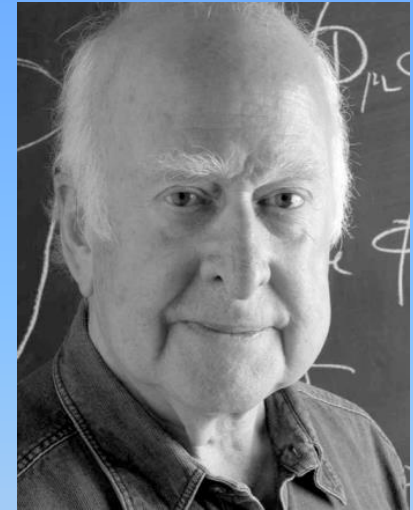
2) Das dort benutzte Prinzip der Konstanz der Lichtgeschwindigkeit ist natürlich in den Maxwellschen Gleichungen enthalten.

Wat is massa ?? Anno 1964

Massa van elementaire deeltjes komt door
 “wrijving” met alomtegenwoordig ‘Higgs veld’

$$m: \psi\psi H$$

Huh?



Higgs

BROKEN SYMMETRIES, MASSLESS PARTICLES AND GAUGE FIELDS

D. M. HIGGS
School of Mathematical Physics, University of Edinburgh, Scotland

Received 21 July 1964

Recently a number of people have discussed the Goldstone theorem [1, 2]. This says solutions of a Lorentz-invariant theory which violate an internal symmetry operation of that theory must contain a massless scalar particle. Kibble and Lee [3] showed that this theorem does not necessarily apply to non-relativistic theories and argued that their considerations would apply equally well to Lorentz-invariant field theories. Gilbert [4], however, gave a proof that the failure of the Goldstone theorem in the non-relativistic case is of a type which cannot exist when Lorentz invariance is imposed on a theory. The purpose of this note is to show that Gilbert's argument fails for an important class of field theories, that is which the conserved currents are coupled to gauge fields.

Following the procedure used by Utiyama [5], let us consider a theory of two fermionic scalar fields

$\psi_1(x), \psi_2(x)$ which is invariant under the phase transformations

$$\psi_1 \rightarrow e^{i\alpha} \psi_1, \quad \psi_2 \rightarrow e^{i\beta} \psi_2 \quad (1)$$

There then is a conserved current J_μ such that

$$\partial_\mu J^\mu = 0 \quad (2)$$

We assume that the Lagrangian is such that symmetry is broken by the spontaneous of the vacuum expectation value of ψ_1 . Goldstone's theorem is proved by showing that the Fourier transform of $\langle \psi_1(x), \psi_1(0) \rangle$ contains a term

$$\frac{1}{k^2} \delta(k^2) \delta^4(k) \quad (3)$$

where $\delta(k^2)$ is the distribution, as a consequence of Lorentz invariance, the conservation law (2). Kibble and Lee [3] avoided this result in the non-relativistic case by showing that the more general form of the Fourier transform is now, in Gilbert's notation,

$$F(k) = \frac{1}{k^2} \delta(k^2) + \frac{1}{k^2} \delta(k^2) + \frac{1}{k^2} \delta(k^2) \quad (4)$$

where $\delta(k^2)$ may be taken as $\delta(k^2) \delta^4(k)$, (5) plus an a priori Lorentz frame. The conservation law then reduces eq. (3) to the less general

obtained in order to define a radiation gauge in which the vector gauge fields are well defined operators. Such theories are nevertheless Lorentz-covariant, as has been shown by Dirac [6]. (This has, of course, long been known of the special such theory, quantum electrodynamics.) There seems to be no reason why the vector $\delta(k^2)$ should not appear in the Fourier transform under consideration.

It is characteristic of gauge theories that the conservation laws hold in the strong sense, as a consequence of field equations of the form

$$\partial_\mu \psi = 0 \quad (6)$$

Except in the case of Abelian gauge theories, the fields A_μ, V_μ are not simply the gauge field variables A_μ, V_μ , the former additional terms with coefficients of the structure constants of the group as functions of A_μ, V_μ must be given by eq. (6). Applying eq. (6) to the commutator given as the Fourier transform of $\langle \psi_1(x), \psi_1(0) \rangle$ the single term $\frac{1}{k^2} \delta(k^2) \delta^4(k)$, etc. the term thus cancelled both Goldstone's zero-mass bosons and the

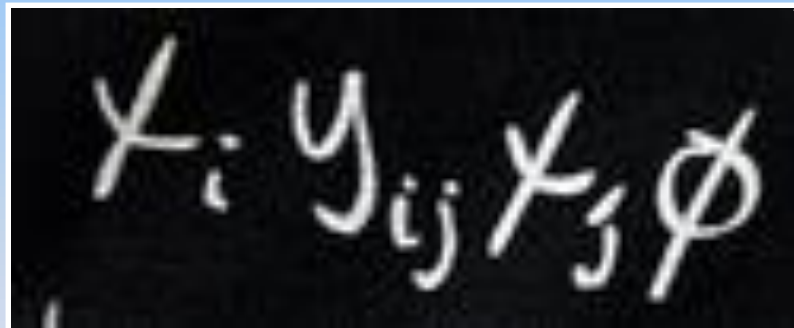
Wat is massa ?? Anno 1964

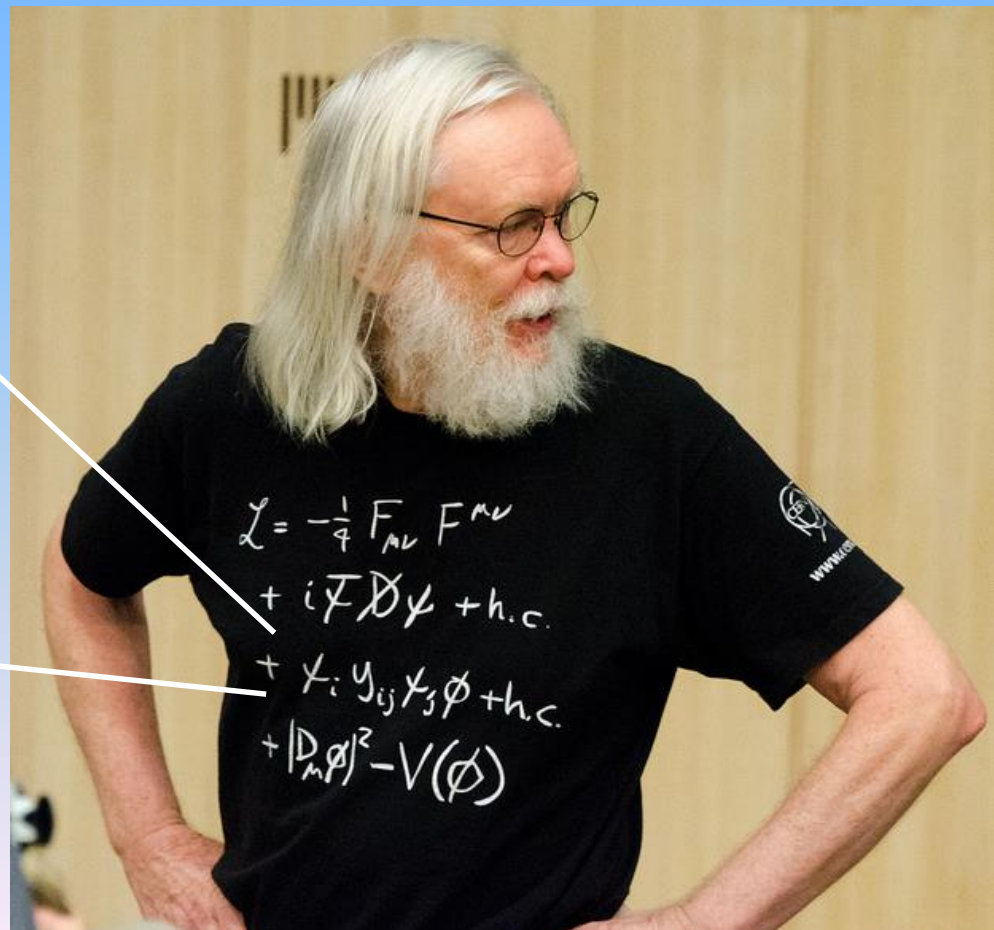
Massa van elementaire deeltjes komt door

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$$m: \psi\psi H$$




$$\mathcal{L} = -\frac{1}{4} F_{\mu\nu} F^{\mu\nu} + i\bar{\psi}\not{D}\psi + h.c. + \chi_i y_{ij} \chi_j \phi + h.c. + |D_\mu \phi|^2 - V(\phi)$$



Nee, dit is niet Pierre, dit is John Ellis

Modelleren van interactie

Standaard Model Lagrangiaan

$$\begin{aligned} & -\frac{1}{2}\partial_\nu g_\mu^a \partial_\nu g_\mu^a - g_s f^{abc} \partial_\mu g_\mu^a g_\mu^b g_\mu^c - \frac{1}{4}g_s^2 f^{abc} f^{ade} g_\mu^a g_\mu^b g_\mu^c g_\mu^d g_\mu^e + \frac{1}{2}ig_\mu^2 (g_\mu^a \gamma^\mu g_\mu^a) g_\mu^2 + \\ & G^a \partial^2 G^a + g_s f^{abc} \partial_\mu G^a G^b \partial_\mu W_\mu^c - \partial_\mu W_\mu^+ \partial_\nu W_\nu^- - M^2 W_\mu^+ W_\mu^- - \frac{1}{2}\partial_\mu Z_\mu^0 Z_\mu^0 - \frac{1}{2}g^2 M^2 Z_\mu^0 Z_\mu^0 - \\ & \frac{1}{2}\partial_\mu A_\nu \partial_\mu A_\nu - \frac{1}{2}\partial_\mu H \partial_\mu H - \frac{1}{2}m_H^2 H^2 - \partial_\mu \phi^+ \partial_\mu \phi^- - M^2 \phi^+ \phi^- - \frac{1}{2}\partial_\mu \phi^0 \partial_\mu \phi^0 - \\ & \frac{1}{2c_w^2} M \phi^0 \phi^0 - \beta_h \left[\frac{2M^2}{g^2} + \frac{2M}{g} H + \frac{1}{2} (H^2 + \phi^0 \phi^0 + 2\phi^+ \phi^-) \right] + \frac{2M^2}{g^2} \alpha_h - ig_{cw} [\partial_\mu Z_\mu^0 (W_\mu^+ W_\mu^- \\ & W_\mu^+ W_\mu^-) - Z_\mu^0 (W_\mu^+ \partial_\mu W_\mu^- - W_\mu^- \partial_\mu W_\mu^+) + Z_\mu^0 (W_\mu^+ \partial_\mu W_\mu^- - W_\mu^- \partial_\mu W_\mu^+)] - ig_{sw} [\partial_\mu A_\nu (W_\mu^+ W_\nu^- \\ & W_\nu^+ W_\mu^-) - A_\nu (W_\mu^+ \partial_\nu W_\nu^- - W_\nu^- \partial_\nu W_\mu^+) + A_\nu (W_\nu^+ \partial_\mu W_\mu^- - W_\mu^- \partial_\mu W_\nu^+)] - \frac{1}{2}g^2 W_\mu^+ W_\mu^- W_\mu^+ W_\mu^- + \\ & \frac{1}{2}g^2 W_\mu^+ W_\mu^- W_\mu^+ W_\mu^- + g^2 c_w^2 (Z_\mu^0 W_\mu^+ Z_\mu^0 W_\mu^- - Z_\mu^0 Z_\mu^0 W_\mu^+ W_\mu^-) + g^2 s_w^2 (A_\mu W_\mu^+ A_\mu W_\mu^- - \\ & A_\mu A_\mu W_\mu^+ W_\mu^-) + g^2 s_w c_w [A_\mu Z_\mu^0 (W_\mu^+ W_\mu^- - W_\mu^- W_\mu^+) - 2A_\mu Z_\mu^0 W_\mu^+ W_\mu^-] - g\alpha [H^3 + \\ & H \phi^0 \phi^0 + 2H \phi^+ \phi^-] - \frac{1}{2}g^2 \alpha_h [H^4 + (\phi^0)^4 + 4(\phi^+ \phi^-)^2 + 4(\phi^0)^2 \phi^+ \phi^- + 4H^2 \phi^+ \phi^- + \\ & 2(\phi^0)^2 H^2] - gM W_\mu^+ W_\mu^- H - \frac{1}{2}g \frac{M^2}{c_w} Z_\mu^0 Z_\mu^0 H - \frac{1}{2}ig [W_\mu^+ (\phi^0 \partial_\mu \phi^- - \phi^- \partial_\mu \phi^0) - W_\mu^- (\phi^0 \partial_\mu \phi^+ - \\ & \phi^+ \partial_\mu \phi^0)] + \frac{1}{2}ig [W_\mu^+ (H \partial_\mu \phi^- - \phi^- \partial_\mu H) - W_\mu^- (H \partial_\mu \phi^+ - \phi^+ \partial_\mu H)] + \frac{1}{2}g \frac{1}{c_w} [Z_\mu^0 (H \partial_\mu \phi^0 - \\ & \phi^0 \partial_\mu H) - ig \frac{s_w}{c_w} M Z_\mu^0 (W_\mu^+ \phi^- - W_\mu^- \phi^+) + ig_{sw} M A_\mu (W_\mu^+ \phi^- - W_\mu^- \phi^+) - ig \frac{1-2c_w^2}{2c_w} Z_\mu^0 (\phi^+ \partial_\mu \phi^- - \\ & \phi^- \partial_\mu \phi^+) + ig_{sw} A_\mu (\phi^+ \partial_\mu \phi^- - \phi^- \partial_\mu \phi^+) - \frac{1}{4}g^2 W_\mu^+ W_\mu^- [H^2 + (\phi^0)^2 + 2\phi^+ \phi^-] - \\ & \frac{1}{4}g^2 \frac{1}{c_w} Z_\mu^0 Z_\mu^0 [H^2 + (\phi^0)^2 + 2(2s_w^2 - 1)^2 \phi^+ \phi^-] - \frac{1}{2}g \frac{s_w^2}{c_w} Z_\mu^0 \phi^0 (W_\mu^+ \phi^- + W_\mu^- \phi^+) - \\ & \frac{1}{2}ig^2 \frac{s_w^2}{c_w} Z_\mu^0 H (W_\mu^+ \phi^- - W_\mu^- \phi^+) + \frac{1}{2}g^2 s_w A_\mu \phi^0 (W_\mu^+ \phi^- + W_\mu^- \phi^+) + \frac{1}{2}ig^2 s_w A_\mu H (W_\mu^+ \phi^- - \\ & W_\mu^- \phi^+) - g^2 \frac{s_w}{c_w} (2c_w^2 - 1) Z_\mu^0 A_\mu \phi^+ \phi^- - g^2 s_w^2 A_\mu A_\mu \phi^+ \phi^- - e^2 (\gamma \partial + m_\Delta^2) e^\lambda - \\ & \bar{\nu}^\lambda \gamma \partial \nu^\lambda - \bar{u}_j^2 (\gamma \partial + m_\Delta^2) u_j^2 - \bar{d}_j^2 (\gamma \partial + m_\Delta^2) d_j^2 + ig_{sw} A_\mu [-(e^\lambda \gamma^\mu e^\lambda) + \frac{2}{3}(\bar{u}_j^2 \gamma^\mu u_j^2) - \\ & \frac{1}{3}(\bar{d}_j^2 \gamma^\mu d_j^2)] + \frac{ig}{4c_w} Z_\mu^0 [(\bar{\nu}^\lambda \gamma^\mu (1 + \gamma^5) \nu^\lambda) + (e^\lambda \gamma^\mu (4s_w^2 - 1 - \gamma^5) e^\lambda) + (\bar{u}_j^2 \gamma^\mu (\frac{2}{3}s_w^2 - \\ & 1 - \gamma^5) u_j^2) + (\bar{d}_j^2 \gamma^\mu (1 - \frac{2}{3}s_w^2 - \gamma^5) d_j^2)] + \frac{ig}{2\sqrt{2}} W_\mu^+ [(\bar{\nu}^\lambda \gamma^\mu (1 + \gamma^5) e^\lambda) + (\bar{u}_j^2 \gamma^\mu (1 + \\ & \gamma^5) C_{\lambda\alpha} d_j^2) + \frac{ig}{2\sqrt{2}} W_\mu^- [(\bar{e}^\lambda \gamma^\mu (1 + \gamma^5) \nu^\lambda) + (\bar{d}_j^2 C_{\lambda\alpha} \gamma^\mu (1 + \gamma^5) u_j^2)] + \frac{ig}{2\sqrt{2}} \frac{m_\Delta^2}{M} [-\phi^+ (\bar{\nu}^\lambda (1 - \\ & \gamma^5) e^\lambda) + \phi^- (\bar{e}^\lambda (1 + \gamma^5) \nu^\lambda)] - \frac{g}{2} \frac{m_\Delta^2}{M} [H (\bar{e}^\lambda e^\lambda) + i\phi^0 (\bar{e}^\lambda \gamma^5 e^\lambda)] + \frac{ig}{2M\sqrt{2}} \phi^+ [-m_\Delta^2 (\bar{u}_j^2 C_{\lambda\alpha} (1 - \\ & \gamma^5) d_j^2) + m_\Delta^2 (\bar{u}_j^2 C_{\lambda\alpha} (1 + \gamma^5) d_j^2)] + \frac{ig}{2M\sqrt{2}} \phi^- [m_\Delta^2 (\bar{d}_j^2 C_{\lambda\alpha} (1 + \gamma^5) u_j^2) - m_\Delta^2 (\bar{d}_j^2 C_{\lambda\alpha} (1 - \\ & \gamma^5) u_j^2)] - \frac{g}{2} \frac{m_\Delta^2}{M} H (\bar{u}_j^2 u_j^2) - \frac{g}{2} \frac{m_\Delta^2}{M} H (\bar{d}_j^2 d_j^2) + \frac{ig}{2} \frac{m_\Delta^2}{M} \phi^0 (\bar{u}_j^2 \gamma^5 u_j^2) - \frac{ig}{2} \frac{m_\Delta^2}{M} \phi^0 (\bar{d}_j^2 \gamma^5 d_j^2) + \\ & \bar{X} + (\partial^2 - M^2) X + \bar{X} - (\partial^2 - M^2) X - \bar{X}^0 (\partial^2 - \frac{M^2}{c_w^2}) X^0 + Y \partial Y + ig_{cw} W_\mu^+ (\partial_\mu \bar{X}^0 X - \\ & \partial_\mu \bar{X}^+ X^0) + ig_{sw} W_\mu^+ (\partial_\nu \bar{Y} X^- - \partial_\nu \bar{X}^+ Y) + ig_{cw} W_\mu^- (\partial_\mu \bar{X}^- X^0 - \partial_\mu \bar{X}^0 X^+) + \\ & ig_{sw} W_\mu^- (\partial_\mu \bar{X}^- Y - \partial_\mu \bar{Y} X^+) + ig_{cw} Z_\mu^0 (\partial_\mu \bar{X}^+ X^+ - \partial_\mu \bar{X}^- X^-) + ig_{sw} A_\mu (\partial_\mu \bar{X}^+ X^+ - \\ & \partial_\mu \bar{X}^- X^-) - \frac{1}{2}gM [\bar{X}^+ X^+ H + \bar{X}^- X^- H + \frac{1}{c_w} \bar{X}^0 X^0 H] + \frac{1-2c_w^2}{2c_w} igM [\bar{X}^+ X^0 \phi^+ - \\ & \bar{X}^- X^0 \phi^-] + \frac{1}{2c_w} igM [\bar{X}^0 X^- \phi^+ - \bar{X}^0 X^+ \phi^-] + igM s_w [\bar{X}^0 X^- \phi^+ - \bar{X}^0 X^+ \phi^-] + \\ & \frac{1}{2}igM [\bar{X}^+ X^+ \phi^0 - \bar{X}^- X^- \phi^0] \end{aligned}$$

Bladmuziek (J.S. Bach) bladmuziek

Allegro moderato. (♩ = 96)

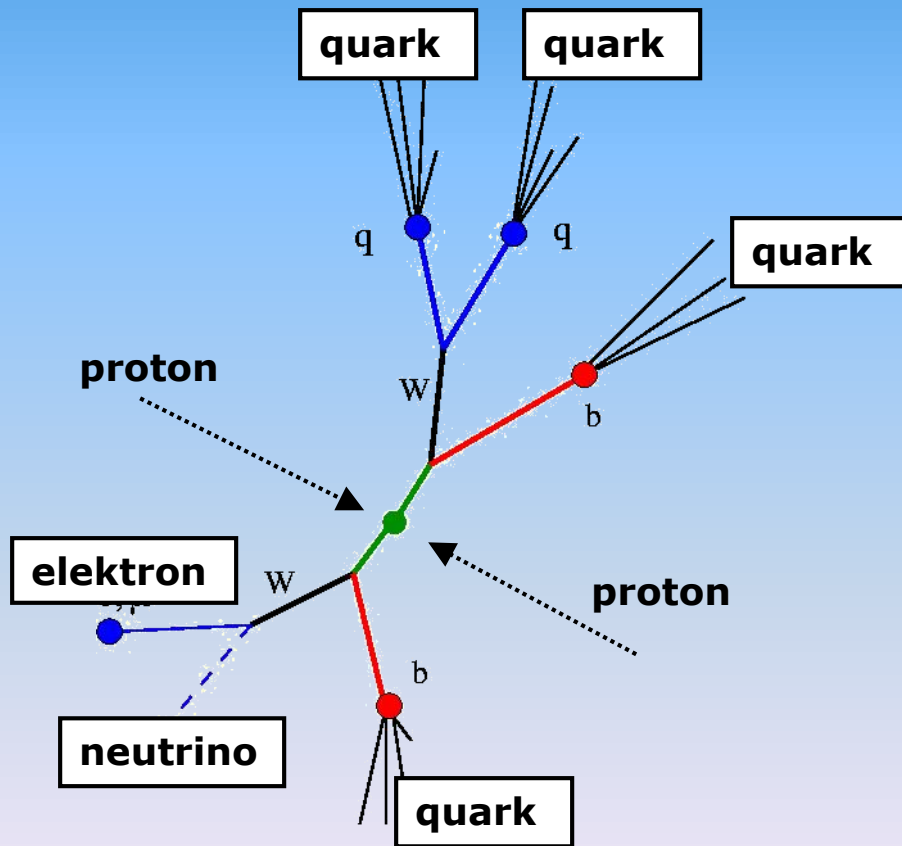
legato

cresc.

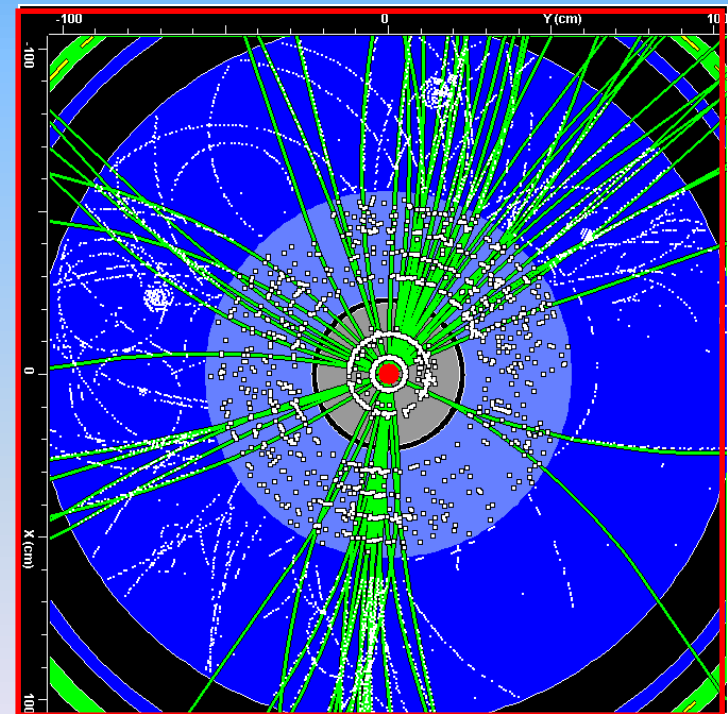
dimin.

poco a poco cresc.

Hoe zien die botsingen er nou uit ?

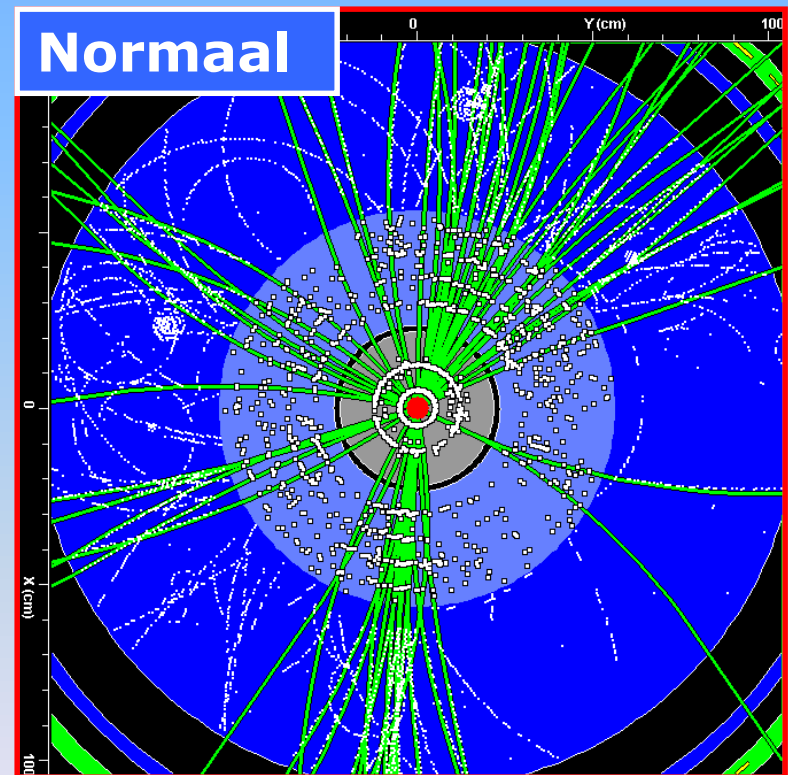
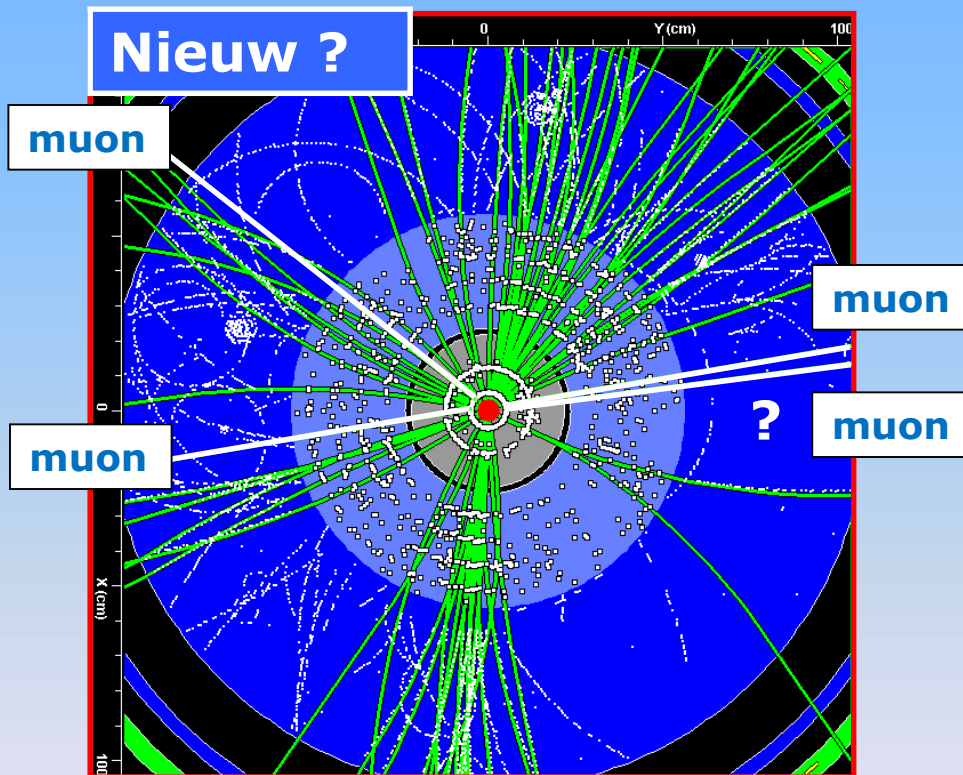


Simulatie top quark productie



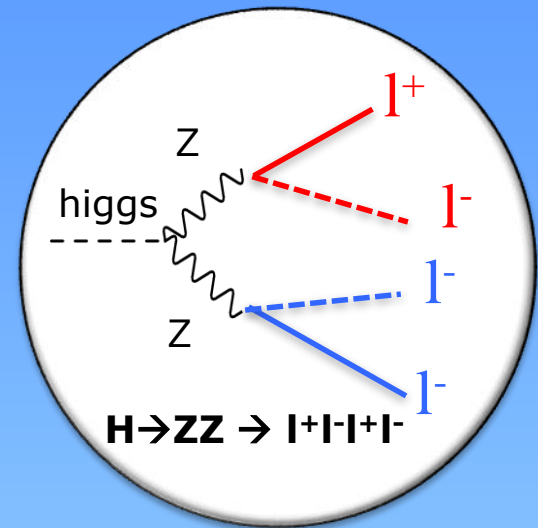
Hoe ontdek je nou nieuwe dingen

Nieuwe afstandschaal EN nieuwe detector



Higgs \rightarrow ZZ \rightarrow 4 leptonen

klein aantal schitterende botsingen



120.000 Higgs bosonen

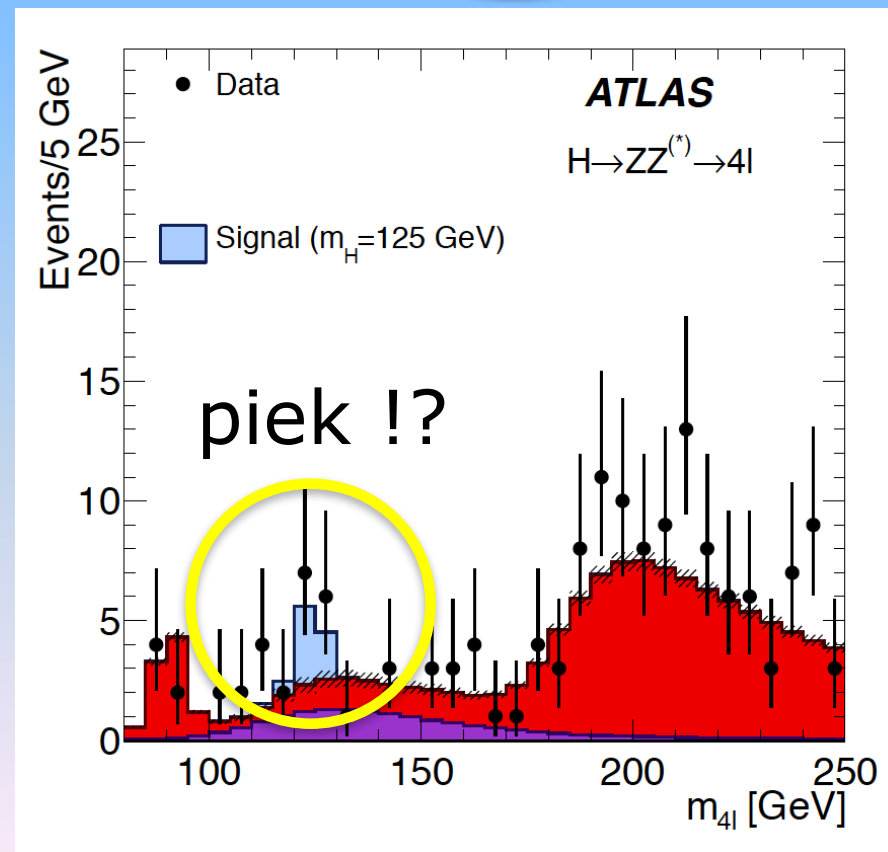


- Maar 1 op de 1000 Higgs bosonen vervalt naar 4 leptonen
- 50% kans dat ATLAS detector ze allemaal goed terugvindt



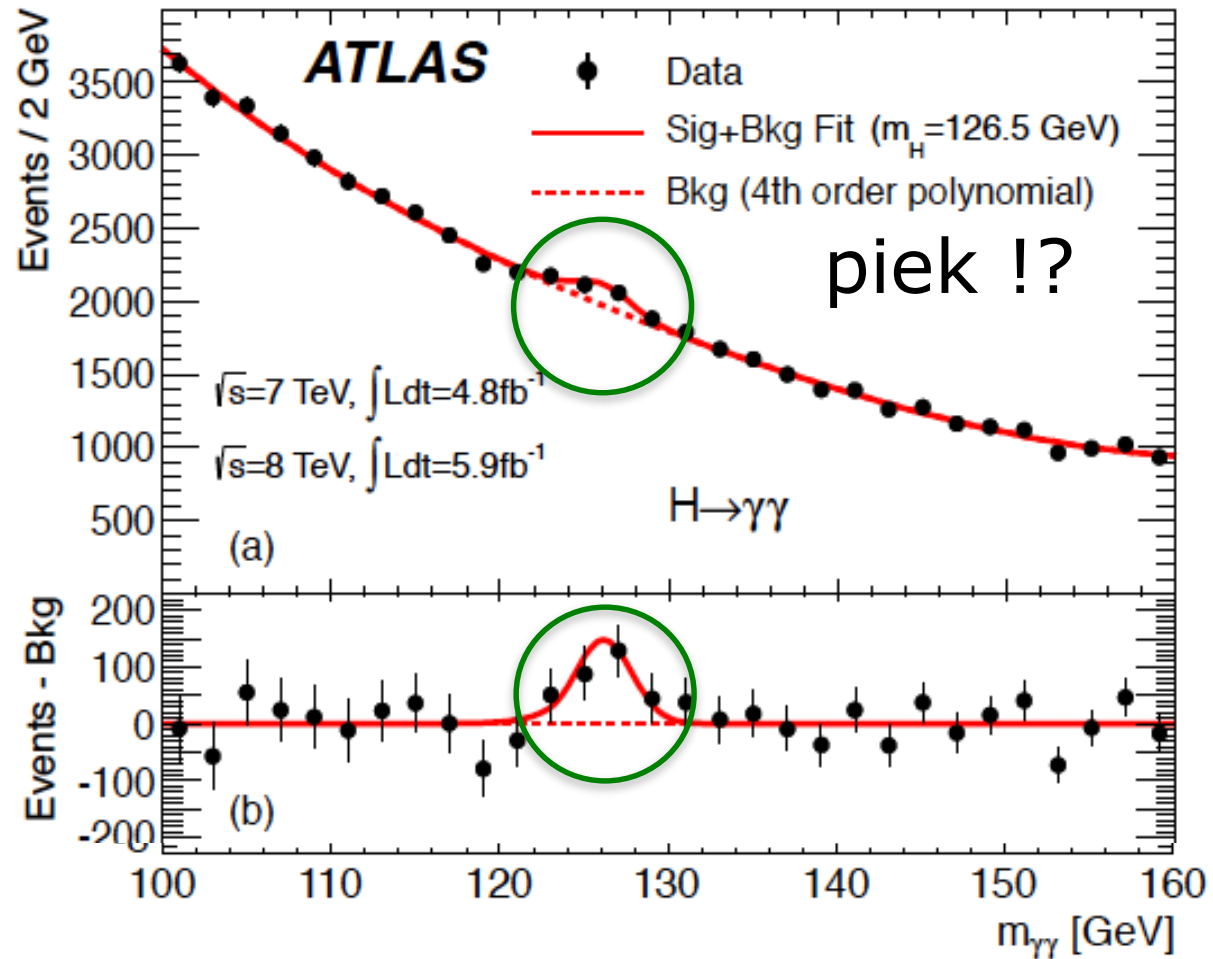
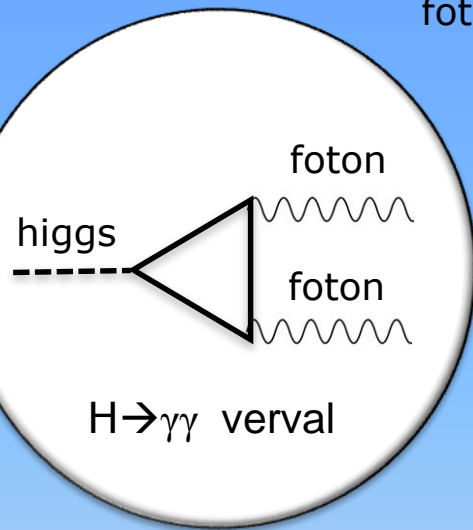
60 (Higgs \rightarrow 4 lepton) events

'overig'	52 events
Met Higgs	68 events



Higgs \rightarrow 2 fotonen

foton



Interpretatie overschot in ATLAS



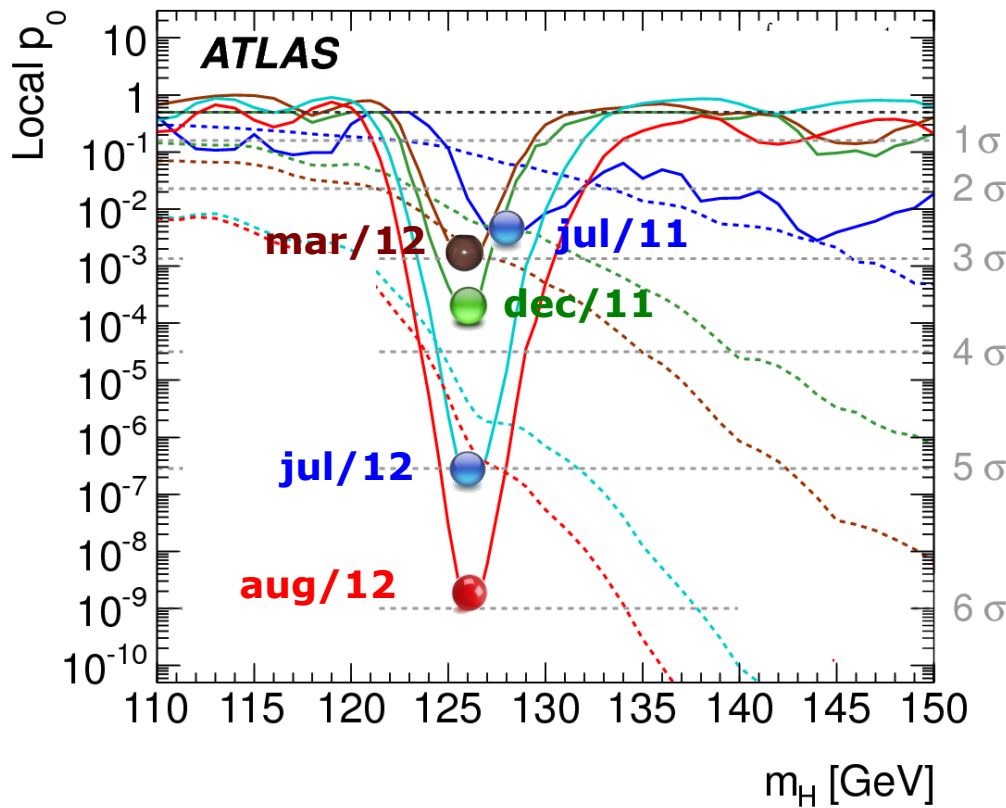
Claim pas ontdekking als:

Kans op toevallige fluctuatie zoals geobserveerd kleiner dan 1 op 1 miljoen

8 keer 6 gooien achter elkaar

Een ontdekking in slow-motion

Time-line higgs ontdekking



EUROPEAN ORGANISATION FOR NUCLEAR RESEARCH (CERN)



CERN-PH-EP-2012-218
Accepted by: Physics Letters B

Observation of a New Particle in the Search for the Standard Model Higgs Boson with the ATLAS Detector at the LHC

The ATLAS Collaboration

This paper is dedicated to the memory of our ATLAS colleagues who did not live to see the full impact and significance of their contributions to the experiment.

Abstract

A search for the Standard Model Higgs boson in proton-proton collisions with the ATLAS detector at the LHC is presented. The datasets used correspond to integrated luminosities of approximately 4.8 fb^{-1} collected at $\sqrt{s} = 7 \text{ TeV}$ in 2011 and 5.8 fb^{-1} at $\sqrt{s} = 8 \text{ TeV}$ in 2012. Individual searches in the channels $H \rightarrow ZZ^{(0)} \rightarrow 4\ell$, $H \rightarrow \gamma\gamma$ and $H \rightarrow WW^{(0)} \rightarrow \ell\nu\ell\nu$ in the 8 TeV data are combined with previously published results of searches for $H \rightarrow ZZ^{(0)}$, $WW^{(0)}$, $b\bar{b}$ and $\tau^+\tau^-$ in the 7 TeV data and results from improved analyses of the $H \rightarrow ZZ^{(0)} \rightarrow 4\ell$ and $H \rightarrow \gamma\gamma$ channels in the 7 TeV data. Clear evidence for the production of a neutral boson with a measured mass of $126.0 \pm 0.4 (\text{stat}) \pm 0.4 (\text{sys}) \text{ GeV}$ is presented. This observation, which has a significance of 5.9 standard deviations, corresponding to a background fluctuation probability of 1.7×10^{-4} , is compatible with the production and decay of the Standard Model Higgs boson.

arXiv:1207.7214v2 [hep-ex] 31 Aug 2012

Ontdekking van het Higgs deeltje op 4 juli 2012



Nog een paar 'kleine' dingetjes:

1

Waar is alle anti-materie gebleven ?

2

80% van de materie in het heelal is onbekend
→ donkere materie

3

Higgs boson en quark koppelingen?
(wat is het verband) ?

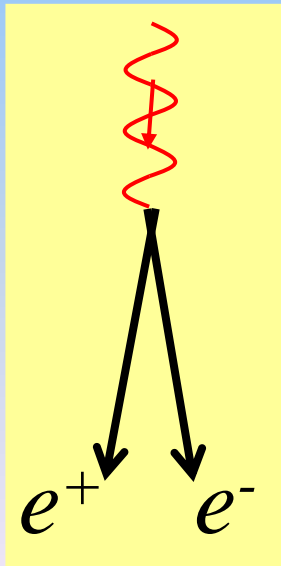
- waarom past gravitatie niet in SM, extra dimensies, waarom 3 families, fermionen fundamentele deeltjes, supersymmetrie, protonen stabiel, kwantisatie elektrische lading, exploderende quantumcorrecties, kleine neutrino massa's, string theorie, ...

Higgs en het Universum

Higgs: Deeltje? Veld?

Deeltje

Foton (lichtdeeltje)



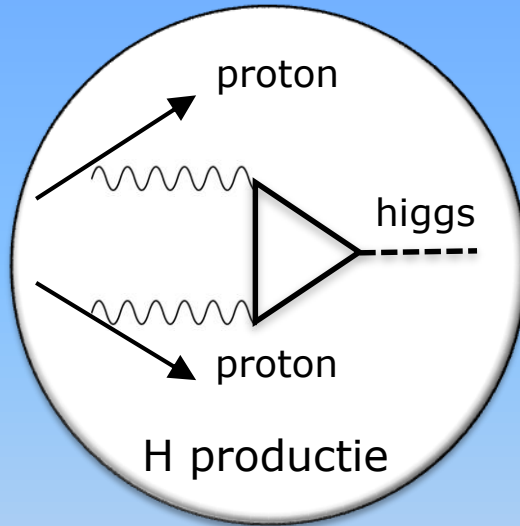
Veld

Elektrisch veld

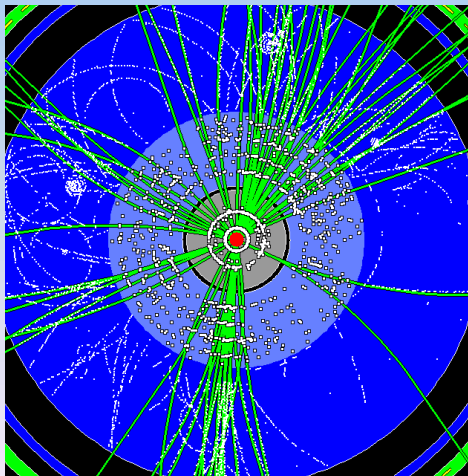
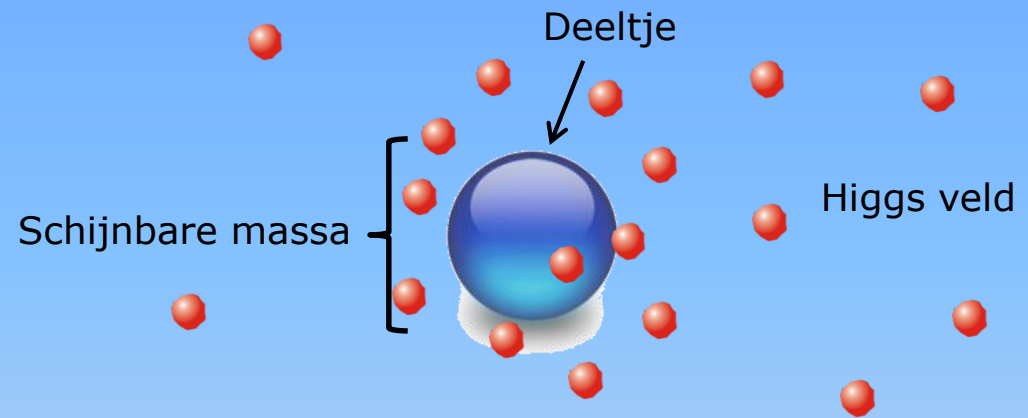


Waarom is de Higgs zo bijzonder?

Deeltje



Veld



Alsof de vis het water heeft ontdekt...

Het Higgs veld – kun je het zien?

Het Higgs veld is uniform – als het meer in deze foto

Het maken van een Higgs deeltje is als een rimpeling op het meer

De theorie van Higgs:

als het veld bestaat is er
ook een bijbehorend deeltje

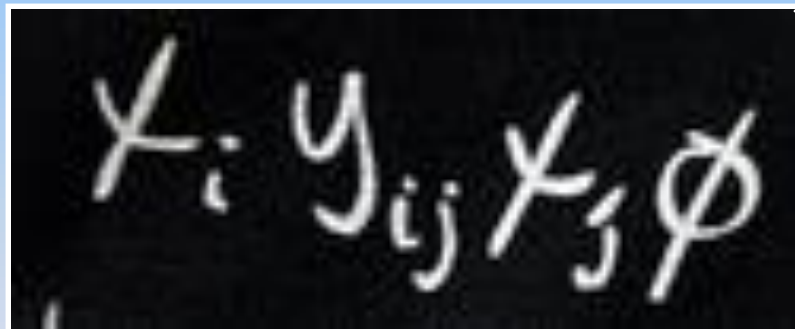


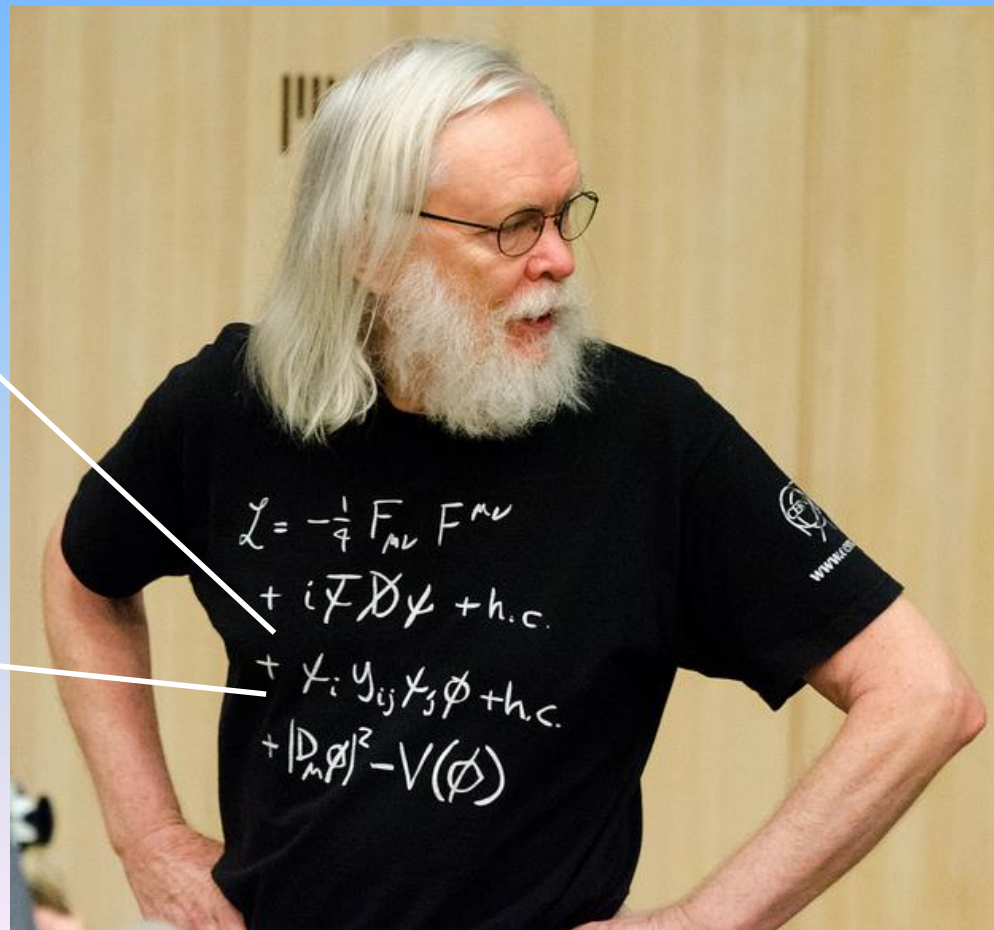
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$$m: \psi\psi H$$




$$\mathcal{L} = y_{ij} \psi_j \phi + h.c.$$



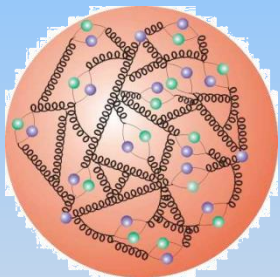
Nee, dit is niet Pierre, dit is John Ellis

Wat is massa ?

Massa van elementaire deeltjes komt door
"wrijving" met alomtegenwoordig 'Higgs veld'

Einstein:

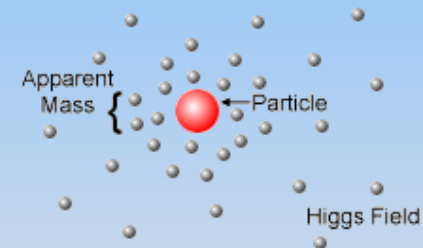
proton massa =
bindings energie



Elementair deeltje
in lege ruimte:
geen rust-energie=
geen massa

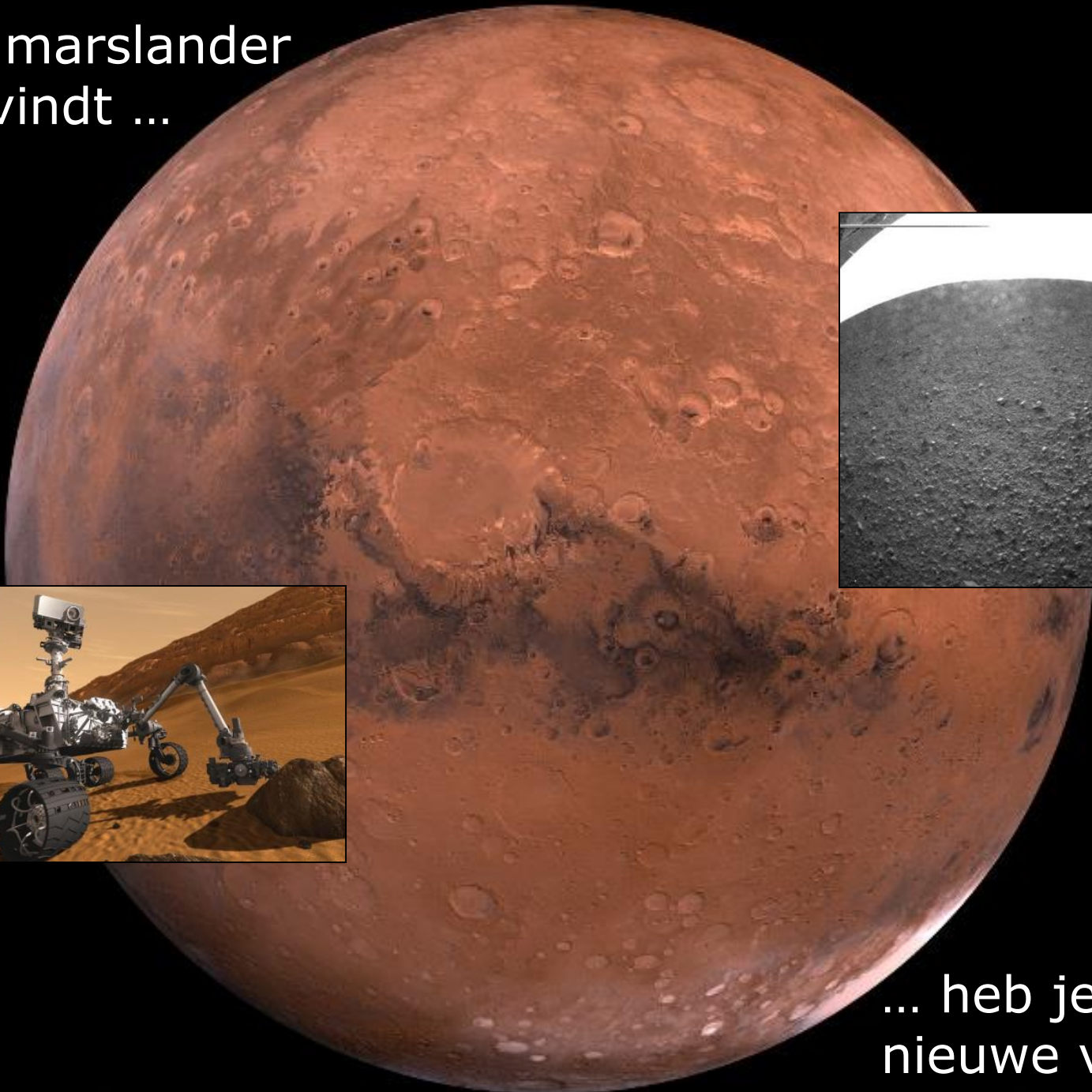


Elementair deeltje
in Higgs veld:
rust energie =
interactie met Higgs veld
= massa



Revolutionair – met spectaculaire consequenties :
de ruimte is *niet* leeg, maar gevuld met soort 'ether'

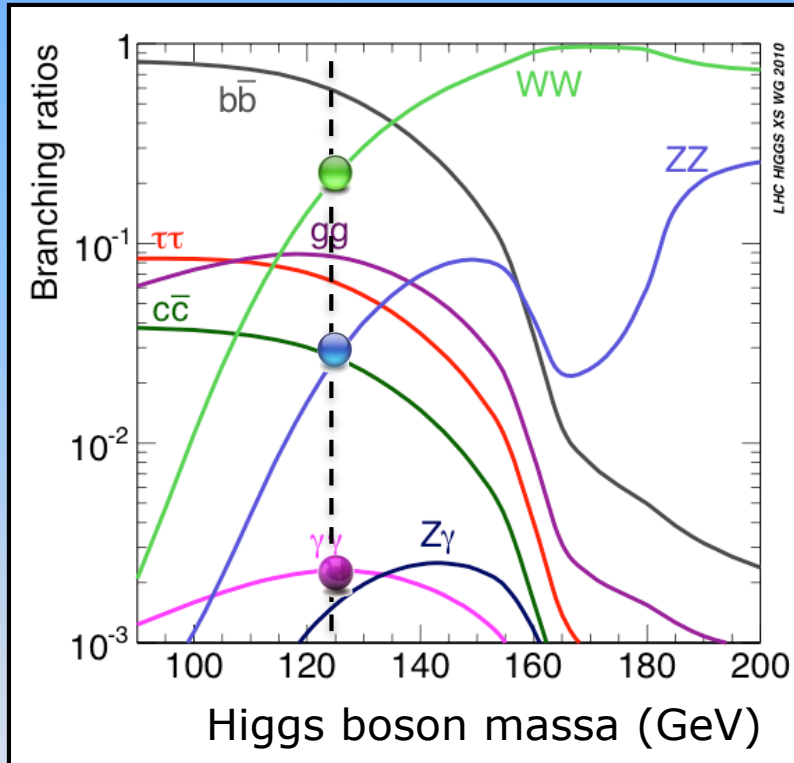
Als de marslander
leven vindt ...



... heb je 1000
nieuwe vragen

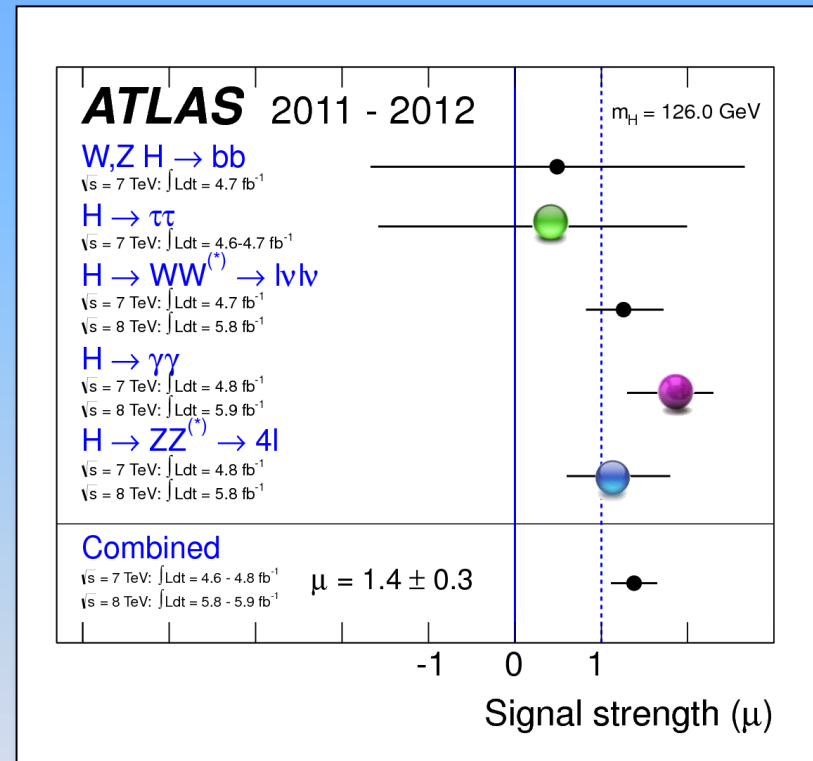
Kloppen Higgs' eigenschappen ?

$m_h = 125 \text{ GeV}$



voorspelling

Standaard Model



meting

Kloppen Higgs' eigenschappen ?

Zijn er nog meer soorten Higgs deeltjes ?

$$\begin{aligned}\mathcal{L}_{\text{susy}} = & -\frac{g^2}{8} (H_u^\dagger \sigma^a H_u + H_d^\dagger \sigma^a H_d)^2 - \frac{g'^2}{8} (H_u^\dagger H_u - H_d^\dagger H_d)^2 \\ & + \lambda_{ij}^u H_u^T \epsilon \bar{u}_i q_j - \lambda_{ij}^d H_d^T \epsilon \bar{d}_i q_j - \lambda_{ij}^e H_e^T \epsilon \bar{e}_i \ell_j \\ & - \frac{H_u^\dagger}{\sqrt{2}} (g \sigma^a \bar{W}^a + g' \bar{B}) \bar{H}_u - \frac{H_d^\dagger}{\sqrt{2}} (g \sigma^a \bar{W}^a - g' \bar{B}) \bar{H}_d + \text{h.c.}\end{aligned}$$

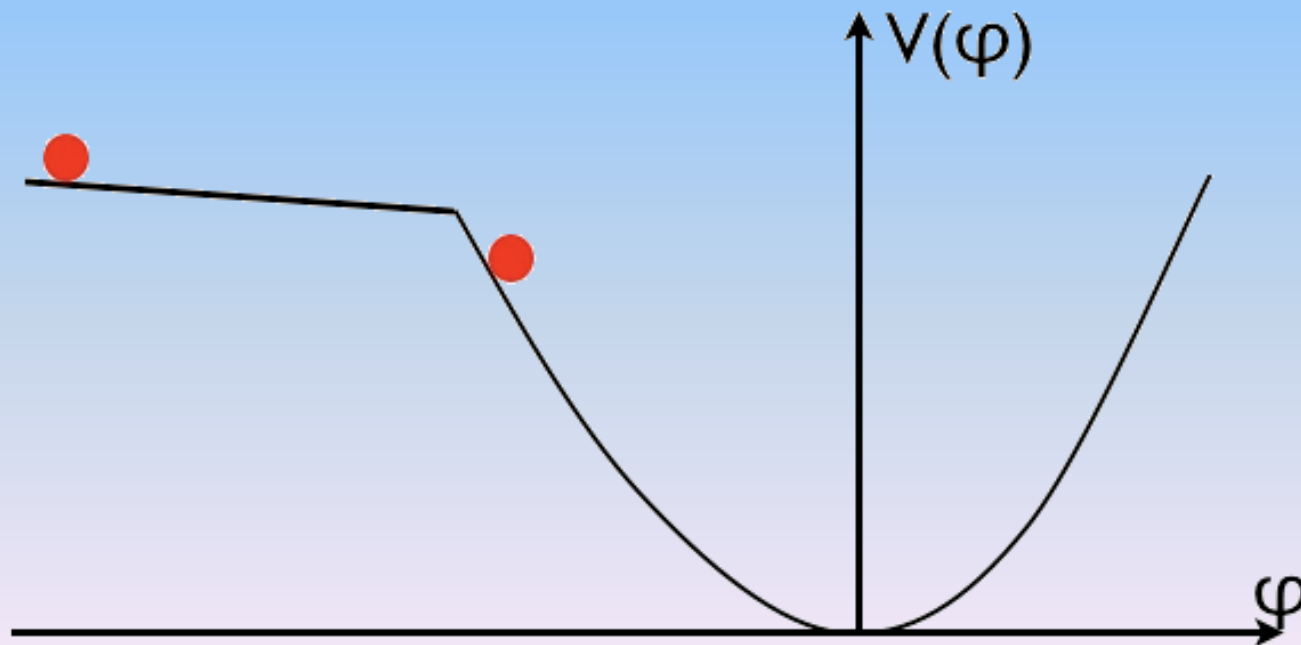
Een stap verder...



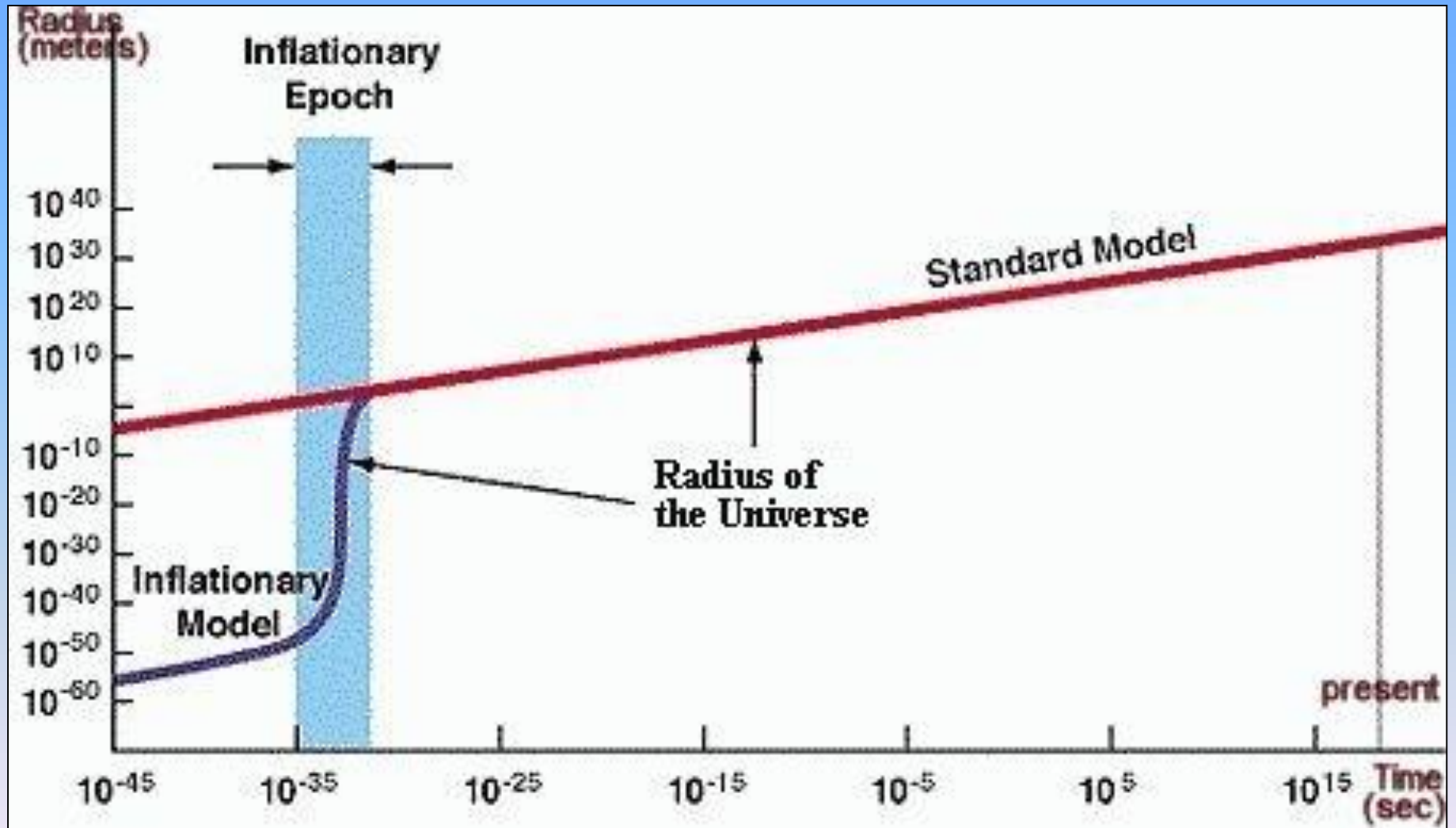
Een ander veld: de Big Bang

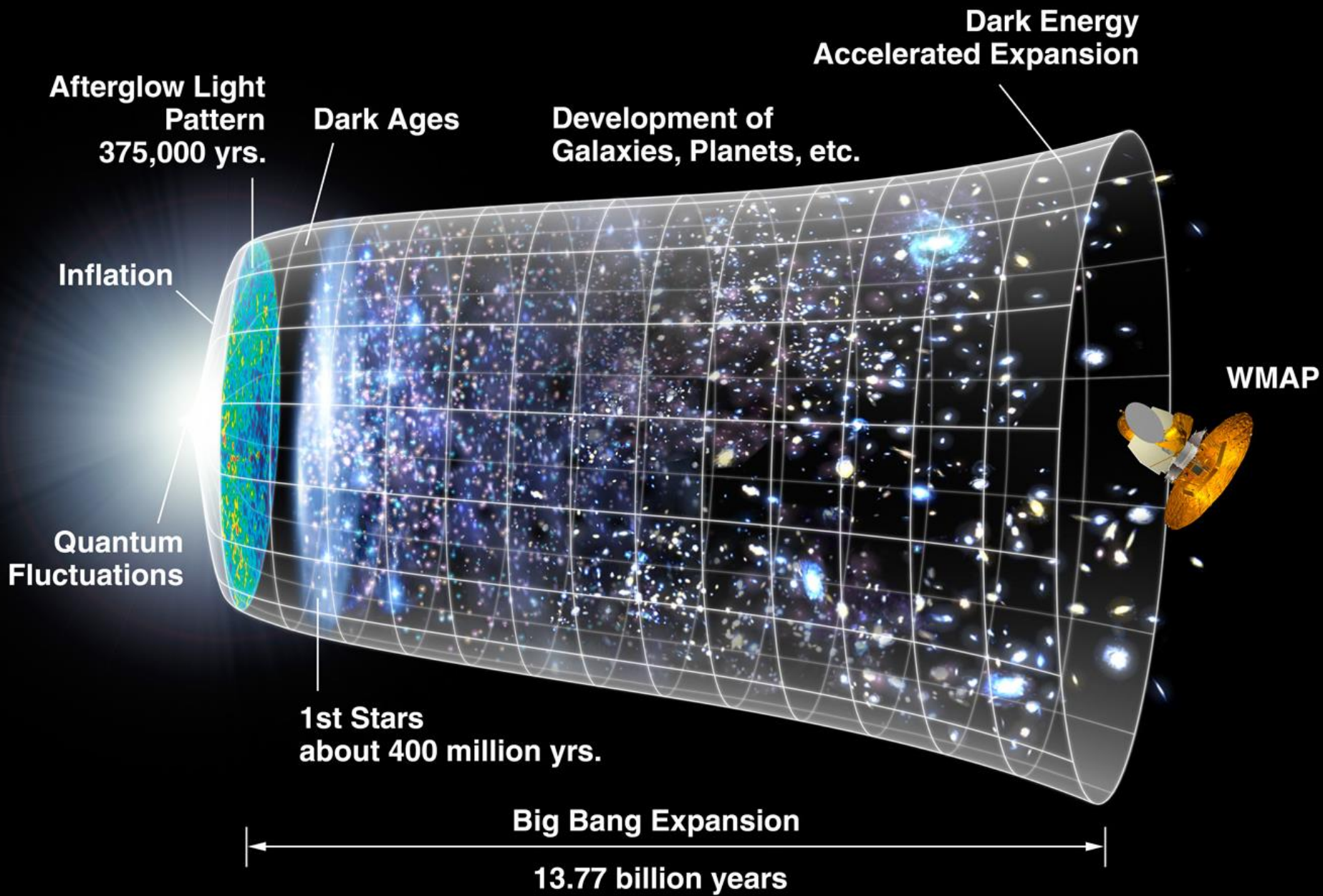
Een van Higgs' eigenschappen komt overeen met een ander veld...

Het *inflaton* dat de het heelal opblies tussen 10^{-33} en 10^{-32} seconde na de Big Bang



Een ander veld: de Big Bang





Higgs:

1

Een Higgs deeltje gevonden in Geneve

2

Heelal gevuld met Higgs veld

3

Kloppen zijn eigenschappen?

Nog een paar 'kleine' dingetjes:

4

80% van de materie in het heelal is onbekend
→ donkere materie

5

Waar is alle anti-materie gebleven ?

6

Higgs boson (hoe krijgen deeltjes massa) ?

- waarom past gravitatie niet in SM, extra dimensies, waarom 3 families, fermionen fundamentele deeltjes, supersymmetrie, protonen stabiel, kwantisatie elektrische lading, exploderende quantumcorrecties, kleine neutrino massa's, string theorie, ...

EINDE

Wat is het nut van dit onderzoek?

Fundamenteel onderzoek

- Kan leiden tot verrassingen,
 - Soms zelfs nuttig...
 - Maar per definitie van te voren onbekend



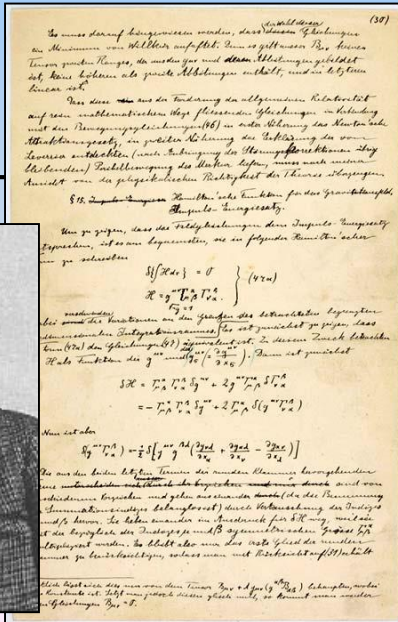
“Oneindig veel toegepast onderzoek aan de kaars zou ons nooit het elektrische licht hebben gebracht.”



Wat is het nut van dit onderzoek?

Fundamenteel onderzoek

- Kan leiden tot verrassingen,
 - Soms zelfs nuttig...
 - Maar per definitie van te voren onbekend

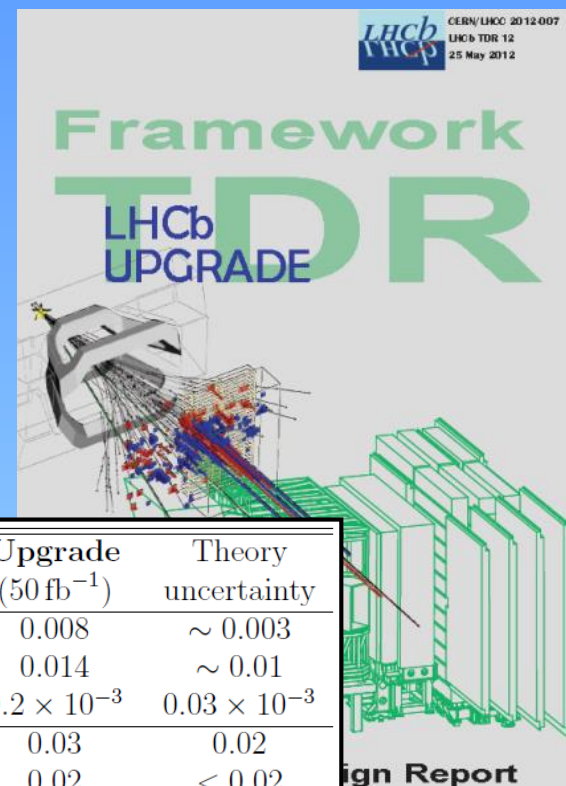


“Zonder relativiteitstheorie zit de GPS er 10km/dag naast!”



LHCb: Hoe verder?

- Preciezer! → Upgrade (2018)



Type	Observable	Current precision	LHCb 2018	Upgrade (50 fb ⁻¹)	Theory uncertainty
B_s^0 mixing	$2\beta_s (B_s^0 \rightarrow J/\psi \phi)$	0.10 [30]	0.025	0.008	~ 0.003
	$2\beta_s (B_s^0 \rightarrow J/\psi f_0(980))$	0.17 [32]	0.045	0.014	~ 0.01
	a_{s1}^s	6.4×10^{-3} [63]	0.6×10^{-3}	0.2×10^{-3}	0.03×10^{-3}
Gluonic penguins	$2\beta_s^{\text{eff}}(B_s^0 \rightarrow \phi\phi)$	–	0.17	0.03	0.02
	$2\beta_s^{\text{eff}}(B_s^0 \rightarrow K^{*0}\bar{K}^{*0})$	–	0.13	0.02	< 0.02
	$2\beta_s^{\text{eff}}(B^0 \rightarrow \phi K_S^0)$	0.17 [63]	0.30	0.05	0.02
Right-handed currents	$2\beta_s^{\text{eff}}(B_s^0 \rightarrow \phi\gamma)$	–	0.09	0.02	< 0.01
	$\tau^{\text{eff}}(B_s^0 \rightarrow \phi\gamma)/\tau_{B_s^0}$	–	5 %	1 %	0.2 %
Electroweak penguins	$S_3(B^0 \rightarrow K^{*0}\mu^+\mu^-; 1 < q^2 < 6 \text{ GeV}^2/c^4)$	0.08 [64]	0.025	0.008	0.02
	$s_0 A_{\text{FB}}(B^0 \rightarrow K^{*0}\mu^+\mu^-)$	25 % [64]	6 %	2 %	7 %
	$A_{\text{I}}(K\mu^+\mu^-; 1 < q^2 < 6 \text{ GeV}^2/c^4)$	0.25 [9]	0.08	0.025	~ 0.02
	$\mathcal{B}(B^+ \rightarrow \pi^+\mu^+\mu^-)/\mathcal{B}(B^+ \rightarrow K^+\mu^+\mu^-)$	25 % [29]	8 %	2.5 %	$\sim 10\%$
Higgs penguins	$\mathcal{B}(B_s^0 \rightarrow \mu^+\mu^-)$	1.5×10^{-9} [4]	0.5×10^{-9}	0.15×10^{-9}	0.3×10^{-9}
	$\mathcal{B}(B^0 \rightarrow \mu^+\mu^-)/\mathcal{B}(B_s^0 \rightarrow \mu^+\mu^-)$	–	$\sim 100\%$	$\sim 35\%$	$\sim 5\%$
Unitarity triangle angles	$\gamma (B \rightarrow D^{(*)}K^{(*)})$	$\sim 10\text{--}12^\circ$ [40, 41]	4°	0.9°	negligible
	$\gamma (B_s^0 \rightarrow D_s K)$	–	11°	2.0°	negligible
	$\beta (B^0 \rightarrow J/\psi K_S^0)$	0.8° [63]	0.6°	0.2°	negligible
Charm	A_Γ	2.3×10^{-3} [63]	0.40×10^{-3}	0.07×10^{-3}	–
CP violation	ΔA_{CP}	2.1×10^{-3} [8]	0.65×10^{-3}	0.12×10^{-3}	–

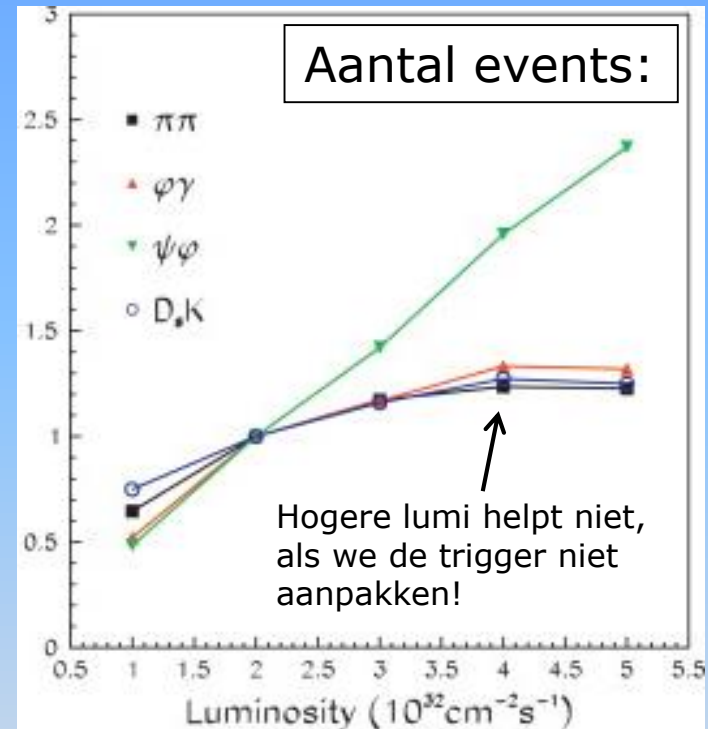
LHCb: Upgrade - Trigger

- Precisie meting → Meer luminositeit
- Meer luminositeit → ~~Hogere trigger rate~~
- Meer luminositeit → Hogere threshold
- Hogere threshold → ~~Minder events ...~~

Oplossing:

Slimmere trigger → *alle* events naar CPU farm:

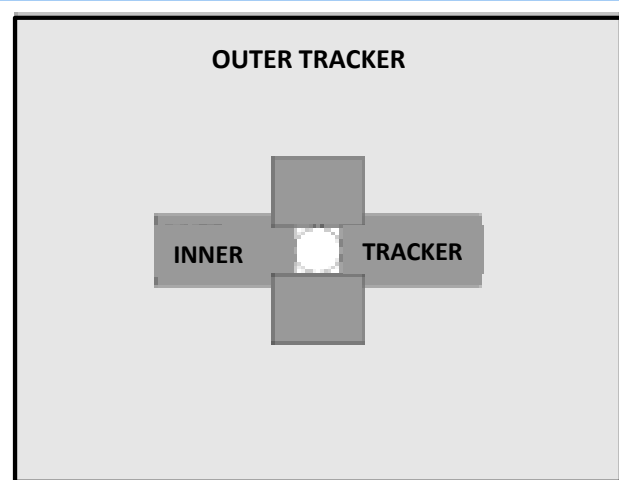
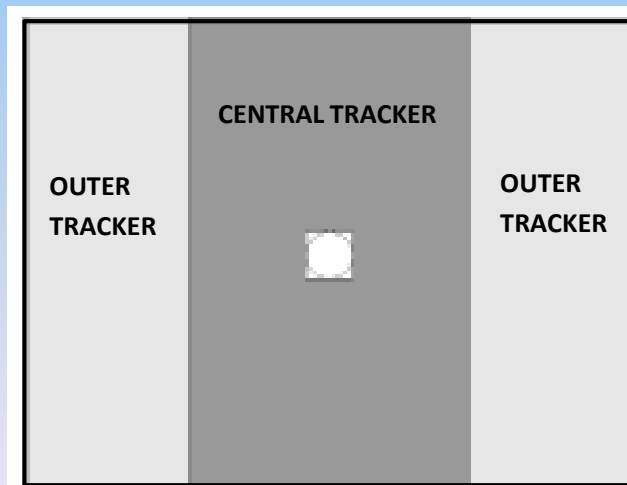
➤ Readout @40 MHz, niet 1 MHz ...



LHCb: Upgrade - Detectors

- Precisie meting → Meer luminositeit
- Meer luminositeit → Hogere particle rate
- Hogere particle rate → **Occupancy te hoog** in de Outer Tracker

➤ 2 opties:

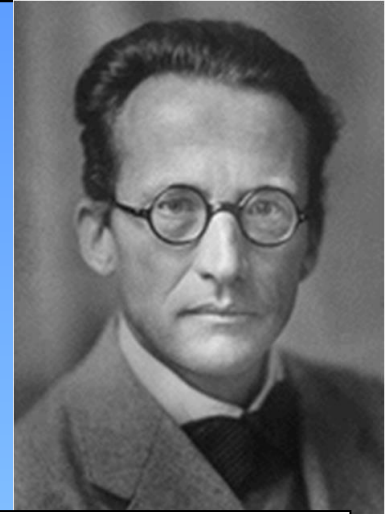


Beslissing in 2013

1) Inner Tracker wordt Scintil. Fiber,
Outer Tracker wordt minder

2) Inner Tracker wordt groter,
Outer Tracker wordt kleiner

Schrödinger



Klassiek verband tussen E and p:

$$E = \frac{\vec{p}^2}{2m}$$

Quantum mechanische substitutie:
(operator acting on wave function ψ)

$$E \rightarrow i\frac{\partial}{\partial t} \quad \text{and} \quad \vec{p} \rightarrow -i\vec{\nabla}$$

Schrodinger vergelijking:

$$i\frac{\partial}{\partial t} \psi = \frac{-1}{2m} \nabla^2 \psi$$

Oplossing:

$$\psi = N e^{i(\vec{p}\vec{x} - Et)}$$

•(show it is a solution)

Klein-Gordon



Relativistisch verband tussen E and p:

$$E^2 = \vec{p}^2 + m^2$$

Quantum mechanische substitutie:
(operator acting on wave function ψ)

$$E \rightarrow i \frac{\partial}{\partial t} \quad \text{and} \quad \vec{p} \rightarrow -i \vec{\nabla}$$

Klein Gordon vergelijking:

$$-\frac{\partial^2}{\partial t^2} \phi = -\nabla^2 \phi + m^2 \phi$$

or :

$$(\square + m^2) \phi(x) = 0$$

or :

$$(\partial_\mu \partial^\mu + m^2) \phi(x) = 0$$

Oplossing:

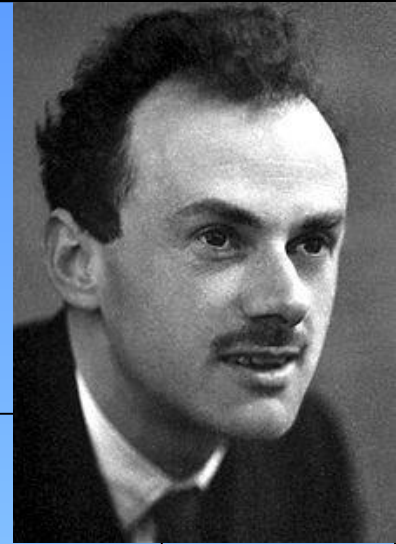
$$\phi(x) = N e^{-ip_\mu x^\mu}$$

$$E^2 = \vec{p}^2 + m^2$$

Maar: negatieve energie oplossing?

$$E = \pm \sqrt{\vec{p}^2 + m^2}$$

Dirac



Paul Dirac zocht een vergelijking, die

- relativistisch correct is,
- Maar linear in d/dt om negatieve energie te vermijden
- (en lineair in d/dx (or ∇) vanwege Lorentz covariantie)

Hij vond een vergelijking, die

- spin-1/2 deeltjes bleek te beschrijven en
- het bestaan van anti-deeltjes voorspelde

Dirac

➤ How to find that relativistic, linear equation ??

Write Hamiltonian in general form,

$$H\psi = (\vec{\alpha} \cdot \vec{p} + \beta m) \psi$$

but when squared, it must satisfy:

$$H^2\psi = (\vec{p}^2 + m^2) \psi$$

Let's find α_i and β !

$$\begin{aligned} H^2\psi &= (\alpha_i p_i + \beta m)^2 \psi && \text{with : } i = 1, 2, 3 \\ &= \left(\underbrace{\alpha_i^2}_{=1} p_i^2 + \underbrace{(\alpha_i \alpha_j + \alpha_j \alpha_i)}_{=0 \quad i>j} p_i p_j + \underbrace{(\alpha_i \beta + \beta \alpha_i)}_{=0} p_i m + \underbrace{\beta^2}_{=1} m^2 \right) \psi \end{aligned}$$

So, α_i and β must satisfy:

- $\alpha_1^2 = \alpha_2^2 = \alpha_3^2 = \beta^2$
- $\alpha_1, \alpha_2, \alpha_3, \beta$ anti-commute with each other
- (not a unique choice!)

Dirac

$$H\psi = (\vec{\alpha} \cdot \vec{p} + \beta m) \psi$$

➤ What are α and β ??

- The lowest dimensional matrix that has the desired behaviour is 4x4 !?

$$\vec{\alpha} = \begin{pmatrix} 0 & \vec{\sigma} \\ \vec{\sigma} & 0 \end{pmatrix} \quad ; \quad \beta = \begin{pmatrix} I & 0 \\ 0 & -I \end{pmatrix}$$

- Often used
- Pauli-Dirac representation:

- with:
$$\sigma_1 = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \quad ; \quad \sigma_2 = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix} \quad ; \quad \sigma_3 = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$$

So, α_i and β must satisfy:

- $\alpha_1^2 = \alpha_2^2 = \alpha_3^2 = \beta^2$
- $\alpha_1, \alpha_2, \alpha_3, \beta$ anti-commute with each other
- (not a unique choice!)

Dirac

$$H\psi = (\vec{\alpha} \cdot \vec{p} + \beta m) \psi$$

Usual substitution:

$$H \rightarrow i\frac{\partial}{\partial t}, \vec{p} \rightarrow -i\vec{\nabla}$$

Leads to:

$$i\frac{\partial}{\partial t}\psi = (-i\vec{\alpha} \cdot \vec{\nabla} + \beta m) \psi$$

Multiply by β :

$$\left(i\beta\frac{\partial}{\partial t}\psi + i\beta\alpha_1\frac{\partial}{\partial x} + i\beta\alpha_2\frac{\partial}{\partial y} + i\beta\alpha_3\frac{\partial}{\partial z} \right) \psi - m\psi = 0 \quad \bullet(\beta^2=1)$$

Gives the famous Dirac equation:

$$(i\gamma^\mu \partial_\mu - m) \psi = 0$$

with : $\gamma^\mu = (\beta, \beta\vec{\alpha}) \equiv$ Dirac γ -matrices

$$\text{for each } j=1,2,3,4 \quad : \quad \sum_{k=1}^4 \left[\sum_{\mu=0}^3 i(\gamma^\mu)_{jk} \partial_\mu - m\delta_{jk} \right] (\psi_k) = 0$$

Dirac

$$H\psi = (\vec{\alpha} \cdot \vec{p} + \beta m) \psi$$

The famous Dirac equation:

$$(i\gamma^\mu \partial_\mu - m) \psi = 0$$

with : $\gamma^\mu = (\beta, \beta\vec{\alpha}) \equiv$ Dirac γ -matrices

R.I.P. :



Dirac vergelijking

Schrödinger equation

- Time-dependence of wave function

$$E = \frac{\vec{p}^2}{2m}$$

$$i \frac{\partial}{\partial t} \psi = \frac{-1}{2m} \nabla^2 \psi$$

Klein-Gordon equation

- Relativistic equation of motion of scalar particles

$$E^2 = \vec{p}^2 + m^2$$

$$-\frac{\partial^2}{\partial t^2} \phi = -\nabla^2 \phi + m^2 \phi$$

Dirac equation

- Relativistically correct, and linear
- Equation of motion for spin-1/2 particles
- Prediction of anti-matter

$$(i\gamma^\mu \partial_\mu - m) \psi = 0$$



$$\psi = \begin{pmatrix} \psi_1 \\ \psi_2 \\ \psi_3 \\ \psi_4 \end{pmatrix}$$