Bootstrapping String-Like Models using Entanglement Minimization and Machine Learning Work with: Debapriyo Chowdhury, Arnab Priya Saha, Aninda Sinha

Based on arXiv:2409.18259 [hep-th]

Cosmological Correlators Worshop, LeCosPA, NTU, Taiwan, 2024

Faizan Bhat, Centre for High-Energy Physics, Indian Institute of Science, Bangalore, India



Open String Amplitude

• Open String Amplitude: Describes the tree-level, two-two scattering of open superstrings ($\alpha' = \ell_s^2 = 1$)

$$\mathcal{M}^{(OS)}(s,t) = \frac{\Gamma(-\alpha's)\Gamma(-\alpha't)}{\Gamma(1-\alpha's-\alpha't)}$$

Duality: For Re(s) < 1,

$$\mathcal{M}^{(OS)}(s,t) = \frac{1}{st} + \sum_{n=1}^{\infty} \frac{(s+1)(s+2)...(s+n-1)}{n!} \frac{1}{t-n}$$

In a (now-very-famous) paper, A.Sinha, A.P.Saha presented

$$\frac{\Gamma(-s)\,\Gamma(-t)}{\Gamma(1-s-t)} = \frac{1}{st} + \sum_{n=1}^{\infty} \frac{1}{n!} \left(\frac{1}{s-n} + \frac{1}{t-n} + \frac{1}{\lambda+n}\right) \\ \times \left(1 - \lambda + \frac{(s+\lambda)(t+\lambda)}{\lambda+n}\right)_{n-1}$$

Field Theory Representations



Figure 1: Field Theory Representation: Poles in all channels + contact terms

Lorentz invariance implies a partial wave expansion of the residues:

$$\operatorname{Res}_{s=n}\mathcal{M}(s,t) = -\pi \sum_{\ell} c_{\ell}^{(n)} \mathcal{C}_{\ell}^{\frac{D-3}{2}} \left(z = 1 + \frac{2t}{s} \right)$$

where $\mathcal{C}_{\ell}^{\frac{D-3}{2}}(z)$ are the Gegenbauer polynomials in D dimensions.

The general formula looks like

$$\mathcal{M}(s,t) = \sum_{n=1}^{\infty} \sum_{\ell=0}^{\ell_{max}} \left[\frac{1}{s-n} + \frac{1}{t-n} + \frac{1}{\lambda+n} \right] c_{\ell}^{(n)} \mathcal{C}_{\ell}^{\left(\frac{D-3}{2}\right)} \left[1 + \frac{2}{n} \left(\frac{(s+\lambda)(t+\lambda)}{\lambda+n} - \lambda \right) \right]$$

- Bootstrap Approach: Find the allowed space of scattering amplitudes by imposing physical constraints.
- Wilson Coefficients: Expansion around s + t = 0 and s t = 0

 $\mathcal{M}_{low}(s,t) = W_{00} + W_{10}(s+t) + W_{01}st + \cdots$

- Q: In the space of consistent scattering amplitudes that satisfy duality, is the open superstring amplitude special?
- A: Yes, it minimizes the total entanglement generated in the scattering process.
- Measured by $\sim -W_{0,0}$. [Aoude, Elor, Remmen, Sumensari]
- Bootstrap Constraints: In D = 10,
 - Crossing Symmetry: $\mathcal{M}(s,t) = \mathcal{M}(t,s)$
 - Analyticity: Only simple poles at $s = n, \forall n \in \mathbb{Z}_{\geq 0}$.
 - Residues at Poles: Polynomials of order $\ell_{max} = n 1$.
 - Fix $W_{1,0} = \zeta(3)$ and $W_{0,1} = \frac{7}{4}\zeta(4)$, i.e. open string values.
 - Unitarity: At tree-level, $c_{\ell}^{(n)} \ge 0$
 - λ -Independence: $\partial_{\lambda}^{k}\mathcal{M}_{\lambda}(s,t) = 0$, for $k \in \mathbb{Z}_{\geq 1}, \lambda \geq -1, (s,t) \in \mathcal{D}_{\lambda}$



Figure 2: Example of \mathcal{D}_{λ} we use for bootstrap

• Maximize $W_{0,0}$ / Minimize Entanglement (via SDPB) ($N_{max} = 30, \epsilon = 10^{-9}, k_{max} = 6, \lambda = 14.6$)



• Maximize $W_{0,0}$ / Minimize Entanglement (via SDPB) ($N_{max} = 30, \epsilon = 10^{-9}, k_{max} = 6, \lambda = 14.6$)



Figure 4: Space of Open-String like theories

Bootstrap using PINNs

- Why do we use PINNs for bootstrap?
- The bootstrap problems we discussed till now and in fact, all bootstrap problems are either Linear optimization problems, or Semi-definite optimization problems.
- These can be handled extremely well via traditional methods.
- Caveat: We used $\epsilon \sim 10^{-9}$ while imposing the constraint

$$-\epsilon \leq \partial_{\lambda}^{k}\mathcal{M}_{\lambda}\left(s,t,c_{\ell}^{(n)}
ight) \leq \epsilon, ext{ for } (s,t) \in \mathcal{D}_{\lambda} ext{ and } 1 \leq k \leq k_{max}$$

However, truncated to some N_{max} , it is not guaranteed that these constraints will be satisfied to some $\epsilon << 1$.

It is more reasonable to impose ratio constraints

$$1 - rac{\mathcal{M}_{\lambda_1}(s,t)}{\mathcal{M}_{\lambda_2}(s,t)} \bigg| \leq \epsilon, \qquad \left| rac{\partial_\lambda^k \mathcal{M}_\lambda(s,t)}{\mathcal{M}_\lambda(s,t)} \right| \leq \epsilon$$

• These are non-linear in the parameters $c_{\ell}^{(n)}$. Traditional methods like SDPB are not useful. This is why we use PINNs.

Bootstrap using PINNs

• Neural networks: Maps with several tunable parameters. In our case,

$$\mathsf{NN}\left(\ell,n,\theta_{j}\right)\equiv c_{\ell}^{(n)}$$

• Neuron Input-Output: $y_j^M = \sigma \left(\sum_{k=1}^{k_{max}} w_{j,k}^M y_k^{M-1} + b_j^M \right)$



Figure 5: Architecture: Input layer with 2 neurons for (ℓ, n) , 2 hidden layers with 64 neurons, output layer wih 1 neuron for $c_{\ell}^{(n)}$. Every neuron has the ReLU activation function $\sigma(x) = \max(0, x)$. Final layer has the SoftPlus activation function $\gamma(x) = \log(1 + e^x)$ for positive $c_{\ell}^{(n)}s$. Total Parameters = [2(64) + 64] + [64(64) + 64] + [64(1) + 1] = 4417.

- Neural networks by minimizing a loss function that measure the violation of constraints.
- We define the following loss function

$$\begin{split} \mathcal{L}\left(\theta_{j}^{M}\right) &= -W_{0,0} + \beta_{1} \left(W_{10} - \left(-\zeta(3)\right)\right)^{2} + \beta_{2} \left(W_{01} - \frac{7\zeta(4)}{4}\right)^{2} \\ &+ \beta_{3} \frac{1}{N} \sum_{s,t \in \mathcal{D}_{\lambda}} \tilde{\mathcal{L}}(s,t). \end{split}$$

• Hyperparameters β_i set the tolerance for constraint violation. Bigger β_i means smaller tolerance $\implies \min(EPM) = \min(\mathcal{L}(\theta_i^M))$

Bootstrap using PINNs

- We implement PINN using the Python library PyTorch.
- Case 1: When $\tilde{\mathcal{L}}(s,t) = \left(1 \frac{\mathcal{M}_{\lambda_1}(s,t)}{\mathcal{M}_{\lambda_2}(s,t)}\right)^2$,
 - $\mathcal{D}_{\lambda} = \{(s,t) \mid -5.5 \le s \le 5.5, -0.2 \le t \le 0.2, \Delta_s = 1, \Delta_t = 0.4\}$
 - Leading Regge Trajectory:

	$c_0^{(1)}$	$c_1^{(2)}$	$c_2^{(3)}$	$c_{3}^{(4)}$	$c_{4}^{(5)}$	$c_{5}^{(6)}$
Open String	1	0.0714	0.0119	0.00289	0.000867	0.000300
PINN	0.999	0.0715	0.0121	0.00300	0.000922	0.000332

• Case 2: When
$$\tilde{\mathcal{L}}(s,t) = \sum_{k=1}^{k_{max}} \left(\frac{1}{\mathcal{M}_{\lambda}(s,t)} \frac{\partial^k \mathcal{M}_{\lambda}(s,t)}{\partial \lambda^k} \right)^2$$
,

- $\mathcal{D}_{\lambda} = \{(s, t) | 0.4 \le s \le 10.4, t = 10.1, \Delta_s = 1\}$
- Leading Regge Trajectory:

	$c_{0}^{(1)}$	$c_1^{(2)}$	$c_{2}^{(3)}$	$c_{3}^{(4)}$	$C_{4}^{(5)}$	$c_{5}^{(6)}$
Open String	1	0.0714	0.0119	0.00289	0.000867	0.000300
PINN	0.998	0.0702	0.0124	0.00341	0.000990	0.000339

• For open string, at (s, t) = (10.4, 10.1), $\mathcal{M}(s, t) \approx 1.34 \times 10^5$ and $\partial_{\lambda}\mathcal{M}_{\lambda}(s, t) \approx -2.51$. \implies Only PINN method can work!.

Summary

- We present a new way to set up the numerical S-Matrix bootstrap using a parametric crossing symmetric dispersion relation λ – CSDR.
- We maximize $W_{0,0}$ / minimize the first finite moment of the entangling power (EPM) and find that the optimal solution is an excellent approximation to the open superstring amplitude.
- We initiate the use of Physics-Informed Neural Networks for the bootstrap to perform non-linear, constrained optimization.
- We also study closed string-like amplitudes and find Dual resonance models there also minimize EPM.

THANK YOU