

Saturation effects in heavy vector meson photoproduction

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In collaboration with Christophe Royon

Based on [2411.14815](#)

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SURGE collaboration

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Guangzhou, China

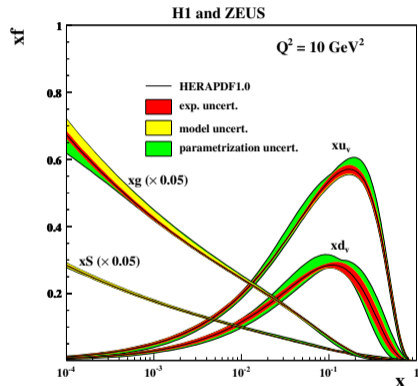


UCLA



Gluon saturation at high energy

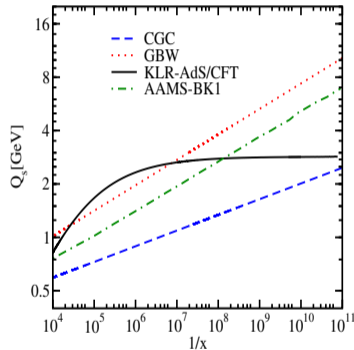
- HERA: rapid growth of gluon distribution at small x
- Growth cannot go on indefinitely: violation of unitarity
- Will eventually be tamed by gluon recombination effects
- Prediction from theory: **gluon saturation**
- Signs of saturation in the experimental data but no definite evidence



H1 and ZEUS (0911.0884)

Searching for saturation

- Saturation effects characterized by the **saturation scale** Q_s^2
- For saturation to be important:
Momentum scale in the process has to be comparable to Q_s^2
- However: Q_s^2 is quite small...
 - Protons: $Q_{s,p}^2 = \mathcal{O}(1 \text{ GeV})$
 - Nuclei: $Q_{s,A}^2 \sim A^{1/3} Q_{s,p}^2$
 \Rightarrow Nuclear enhancement of saturation!
- Energy dependence: $Q_s^2 \sim 1/x^{0.3}$
- Search for saturation:
Need a high energy (small x) and a low momentum scale



Rezaeian (1001.5266)

Exclusive vector meson production

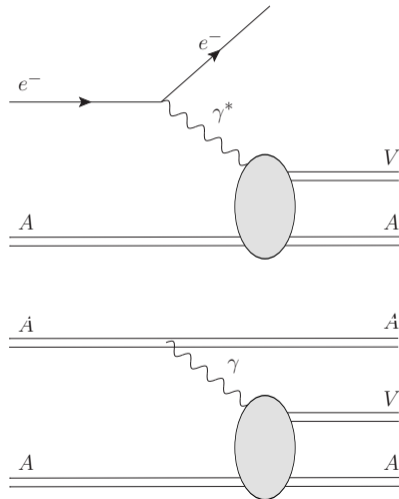
Ryskin, Z.Phys.C 57 (1993) 89-92

$$\frac{d}{dt}\sigma(\gamma^* + A \rightarrow V + A) \sim [xg(x)]^2$$

⇒ Very sensitive to the gluon structure of the target!

- Heavy vector mesons:
Heavy quark mass makes the process perturbative
 - Mass also low enough for saturation!
- Can be measured in:
 - DIS: Electron-ion collisions (HERA, EIC, ...)
 - Ultra-peripheral collisions (LHC, ...)

See Zaochen Ye's talk on Monday

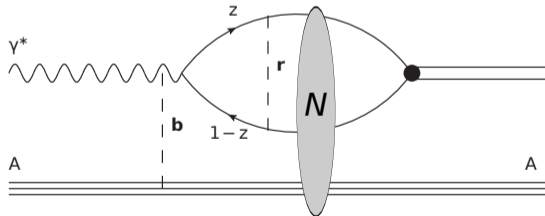


Vector meson production at the leading order in the dipole picture

Invariant amplitude for exclusive vector meson production

$$-i\mathcal{A}^\lambda = 2 \int d^2\mathbf{b} d^2\mathbf{r} \frac{dz}{4\pi} e^{-i\mathbf{b}\cdot\mathbf{\Delta}} \Psi_{\gamma^*}^{q\bar{q}}(\mathbf{r}, z) N(\mathbf{r}, \mathbf{b}, x_{\mathbb{P}}) \Psi_V^{q\bar{q}*}(\mathbf{r}, z), \quad t = -\mathbf{\Delta}^2$$

- $\Psi_{\gamma^*}^{q\bar{q}}$: Photon light-cone wave function
- N : Dipole-target scattering amplitude
- $\Psi_V^{q\bar{q}}$: Vector meson light-cone wave function



Vector meson production at the leading order in the dipole picture

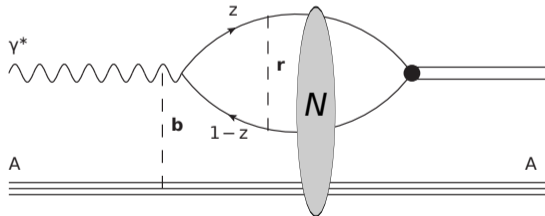
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- Dependence on energy W in the dipole amplitude:

$$x_{\mathbb{P}} = \frac{Q^2 + M_V^2 - t}{W^2 + Q^2 + m_N^2}$$

- $Q^2 = 0$ for photoproduction



Vector meson wave function

- Meson wave function is nonperturbative – has to be modeled
- Various different approaches:
Nonrelativistic QCD, basis light-front quantization...
- We use the Boosted Gaussian that has been found to work well phenomenologically:
[Kowalski, Motyka, Watt \(hep-ph/0606272\)](#)

$$\phi_\lambda(r, z) = \mathcal{N}_\lambda \exp\left(-\frac{m_Q^2 \mathcal{R}^2}{8z(1-z)} - \frac{2z(1-z)r^2}{\mathcal{R}^2} + \frac{m_Q^2 \mathcal{R}^2}{2}\right)$$

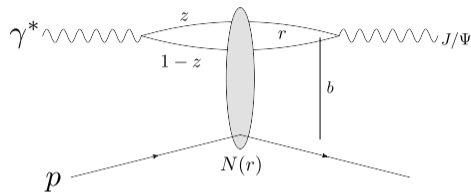
where \mathcal{N}_λ , \mathcal{R} are parameters fixed by normalization and leptonic decay width

Dipole amplitude N

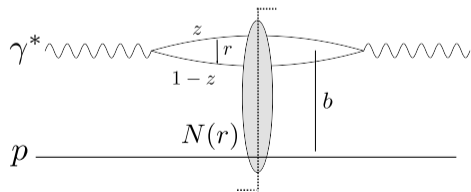
- Describes the interaction with the target
- High energy: **eikonal approximation**

$$N(\mathbf{x}, \mathbf{y}, x_{\mathbb{P}}) = 1 - \frac{1}{N_c} \left\langle \text{Tr} V(\mathbf{x}) V^\dagger(\mathbf{y}) \right\rangle_{x_{\mathbb{P}}}$$

- Universal: appears in different processes
- Energy dependence given by a perturbative **evolution equation**



$$\sigma_{\gamma^* p \rightarrow V p} \sim |N|^2$$



$$\sigma_{\gamma^* p} \sim N$$

Estimating saturation effects

Balitsky–Kovchegov equation

$$\frac{\partial}{\partial \log 1/x} N(\mathbf{x}_0, \mathbf{x}_1) = \frac{N_c \alpha_s}{2\pi^2} \int d^2 \mathbf{x}_2 \frac{\mathbf{x}_{01}^2}{\mathbf{x}_{20}^2 \mathbf{x}_{21}^2} \times [N(\mathbf{x}_0, \mathbf{x}_2) + N(\mathbf{x}_1, \mathbf{x}_2) - N(\mathbf{x}_0, \mathbf{x}_1) - N(\mathbf{x}_0, \mathbf{x}_2)N(\mathbf{x}_1, \mathbf{x}_2)]$$

- Saturation effects introduced by the **nonlinear** term in the BK equation
- Without nonlinear term: BFKL evolution
 - ⇒ Compare BK and BFKL evolutions to estimate saturation effects
- We also include the dependence on the impact parameter (usually neglected)
 - Neglecting it can lead to overestimating saturation effects [JP et al. \(2411.13533\)](#)

Initial condition for the high-energy evolution

- Initial condition chosen as the impact-parameter-dependent McLerran–Venugopalan model JP et al. (2411.13533)

$$\mathbf{r} = \mathbf{x} - \mathbf{y}, \quad \mathbf{b} = \frac{1}{2}(\mathbf{x} + \mathbf{y})$$

$$N(\mathbf{x}, \mathbf{y}) = 1 - \exp\left(-\int d^2\mathbf{z} \kappa T(\mathbf{z}) \left[K_0(m|\mathbf{x} - \mathbf{z}|) - K_0(m|\mathbf{y} - \mathbf{z}|)\right]^2\right)$$

- $T(\mathbf{z})$ = the thickness function describing the shape of the target
- κ = constant describing the strength of the color field
- m = infrared regulator

Proton

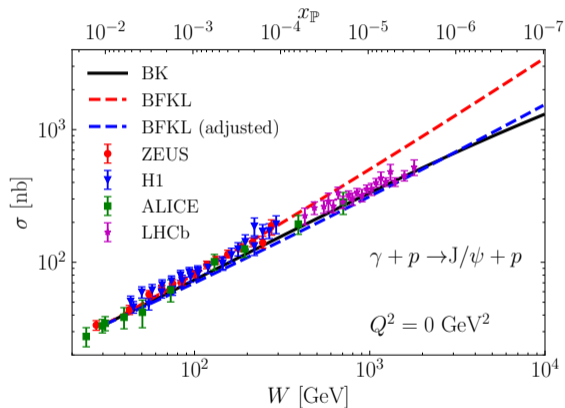
- $T(\mathbf{z})$ = Gaussian
- Parameters fixed by exclusive J/ψ production data

Lead

- From proton to nucleus:
Change only $T(\mathbf{z})$
- $T(\mathbf{z})$ = Woods–Saxon

Exclusive J/ψ production: proton targets

- Slight difference between BFKL and BK in the slope
 - Can be compensated by adjusting α_s for BFKL
- Proton data described well by both BK and BFKL equations



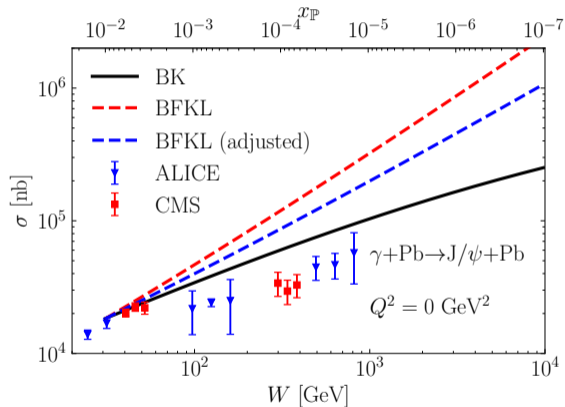
JP, Royon (2411.14815)

Exclusive J/ψ production: nuclear targets

- Fit the model to the proton data
⇒ Predictions for heavy nuclei
- Differences between BK and BFKL:
a factor of 2 for $W \sim 1000$ GeV
- BFKL results linear as predicted
- BK describes the data much better
- Still not exact agreement:
A well-known problem, see e.g.

Mäntysaari, Salazar, Schenke (2312.04194)

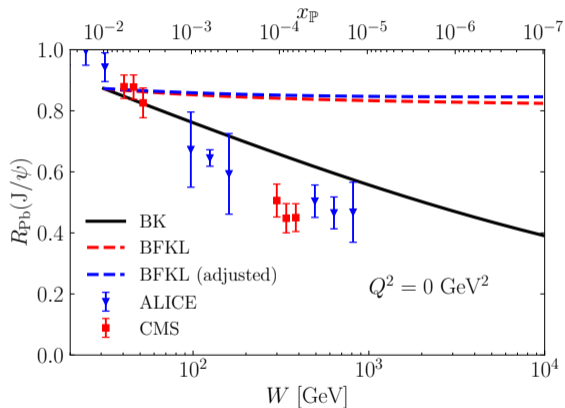
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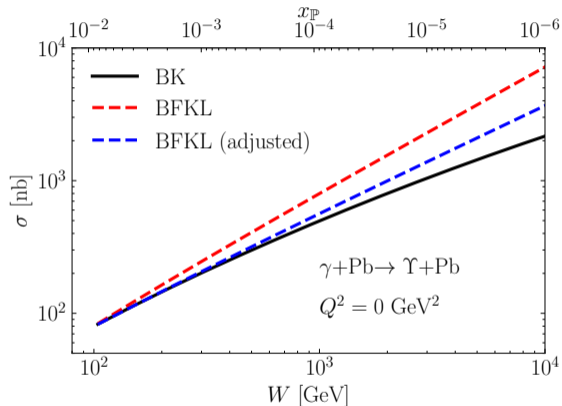
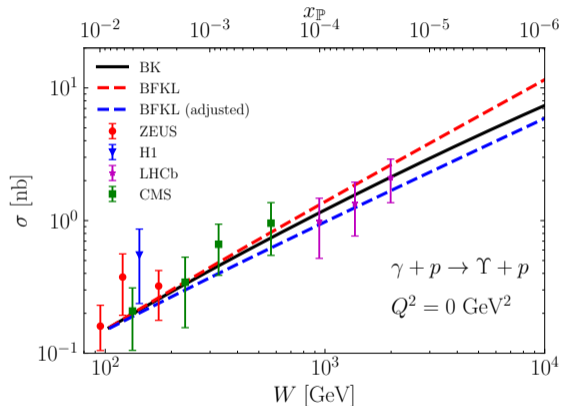
Exclusive J/ψ production: nuclear suppression

- Nuclear suppression factor: $R_A = \sqrt{\sigma_A/\sigma_{IA}}$
where $\sigma_{IA} = \left. \frac{d\sigma^p}{dt} \right|_{t=0} \times \int dt |F_A(t)|^2$
is the impulse approximation
- Without saturation: $R_A \approx 1$
- BFKL essentially constant
 - Linear evolution: energy dependence for protons and nuclei expected to be similar
 - Clear disagreement with the data
 - Saturation provides a natural explanation



JP, Royon (2411.14815)

Υ production



- Saturation effects much smaller \Rightarrow Not expecting sizable differences even at the LHC
- Follows from the large mass of Υ

JP, Royon (2411.14815)

Summary

- Gluon saturation expected at the high-energy limit
- Difficult to measure: need both a high energy and a low momentum scale
- Exclusive heavy vector meson photoproduction is a promising process
 - Diffractive process \Rightarrow sensitive to the gluon density
 - Heavy quark mass \Rightarrow large enough to be perturbative, small enough for saturation
- Compare linear BFKL and nonlinear BK evolution to estimate saturation effects
- Proton targets: no sign of saturation in the current data
- Pb targets: J/ψ **shows a clear preference for BK evolution**
 - Saturation effects already visible in the LHC data?