# HBT correlation in HIC with a hard process (preliminary)

Workshop on Advances, Innovations, and Future Perspectives in High-Energy Nuclear Physics, Wuhan, China

Weiyao Ke, Central China Normal University In collaboration with Zhong Yang, Xin-Nian Wang, De-Xing Zhu October 23, 2024



## Medium response to the passage of a parton (jet, heavy quark, etc)





A demonstration using linearized MIS-type hydrodynamics

First unambiguous evidence of diffusion wake at HP2024! Compared to simulations.

- $\bullet$  Medium response: a natural consequence of jet quenching  $+$  medium dynamics.
- Will it tell something about the medium dynamics?

## Medium response in the spatial coordinates



- Jet excites QGP at "all" wave-length  $\delta(x) \rightarrow f(k) = 1$ . Small perturbations are characterized by linear **response functions**  $\Leftrightarrow G^{\mu\nu}_{\alpha\beta} = \langle T^{\mu\nu}(t,x) T_{\alpha\beta}(t',x') \rangle.$
- Kineitc v.s. hydro responses give fairly different spatial strucuters. However, we can only access momentum space, e.g., single-hadron spectra.

**Background** is a Bjorken flow. **Response** creates some small perturbations ( $\delta \ll 1$ )

$$
\epsilon = \epsilon_0(\tau) + \delta e(\tau, x, y, \eta_s)
$$

$$
u^{\mu} = u_0^{\mu} + \delta u(\tau, x, y, \eta_s)
$$

$$
\pi^{\mu\nu} = \pi_0^{\mu\nu} + \delta \pi^{\mu\nu}(\tau, x, y, \eta_s)
$$

Study the impact on particle production at freeze out defined by  $e(\tau,x,y,\eta_{\sf s})=e_{\sf f}.$ Linearize the Cooper-Frye formula:



The possible angular structure is quite restricted.

$$
\frac{d\delta N}{dm_{\mathcal{T}}dyd\phi} = \frac{p^{\mathcal{T}}\tau_{\text{frz},0}f_{\text{eq}}(\tilde{p}^{\mathcal{T}})}{(2\pi)^3} (N_0 + \tilde{p}_{\mu}N_1^{\mu} + \tilde{p}_{\mu}\tilde{p}_{\nu}N_2^{\mu\nu} + \tilde{p}_{\mu}\tilde{p}_{\nu}\tilde{p}_{\rho}N_3^{\mu\nu\rho}) + \mathcal{O}(\delta^2), \quad \tilde{p}^{\mu} = \frac{p^{\mu}}{\mathcal{T}}
$$

All information of perturbations are encode in the **spatially-integrated** coefficients  $N$ .

$$
N_0 = \int dx dy d\eta_s \left[ c_s^2 \tilde{\rho}^\tau \delta \tilde{\epsilon} + \delta \tilde{\epsilon} - v^\eta \partial_\eta \delta \tilde{\epsilon} - \tau_f v_\perp \cdot \partial_\perp \delta \tilde{\epsilon} \right],
$$
  
\n
$$
N_1^\mu = \int dx dy d\eta_s \left[ -\delta \tilde{g}^\mu - \tilde{\rho}^\tau \tilde{\pi}^{\mu\nu} \delta \tilde{g}_\nu \right],
$$
  
\n
$$
N_2^{\mu\nu} = \int dx dy d\eta_s \left[ \frac{1}{2} \tilde{\pi}^{\mu\nu} \left[ c_s^2 \tilde{\rho}^\tau \delta \tilde{\epsilon} + \delta \tilde{\epsilon} - v^\eta \partial_\eta \delta \tilde{\epsilon} - \tau_f v_\perp \cdot \partial_\perp \delta \tilde{\epsilon} \right] + \frac{1}{2} \Delta^{\mu\nu}_{\alpha\beta} \delta \tilde{\pi}^{\alpha\beta} \right],
$$
  
\n
$$
N_3^{\mu\nu\rho} = \int dx dy d\eta_s \frac{1}{2} \tilde{\pi}^{\mu\nu} (-\delta \tilde{g}^\rho) \quad \text{where} \quad \delta \tilde{\epsilon} = \frac{\delta \epsilon}{\epsilon + P}, \quad \delta \tilde{g}^\mu = \frac{\delta g^\mu}{\epsilon + P} = \delta u^\mu.
$$

No background radial flow, it will be more interesting with radial flow.



- In this toy example, it seems that interested stuff are being integrated out. Energy-momenutm conservation dominates lowest order of the angular structure.
- So how can we access the spatial information in  $\delta \epsilon(\tau, x, y, \eta_s)$ ?



$$
\frac{d\delta N}{dm_T dy d\phi} = N_0 + \tilde{p}_{\mu} N_1^{\mu} + \tilde{p}_{\mu} \tilde{p}_{\nu} N_2^{\mu\nu} + \tilde{p}_{\mu} \tilde{p}_{\nu} \tilde{p}_{\rho} N_3^{\mu\nu\rho}
$$
  
=  $C_{00} (\tilde{p}^{\tau}) + C_{1m} (\tilde{p}^{\tau}) Y_1^m(\Omega)$   
+  $C_{2m} (\tilde{p}^{\tau}) Y_2^m(\Omega) + C_{3m} (\tilde{p}^{\tau}) Y_3^m(\Omega) + \mathcal{O}(\delta^2)$ 

- In this toy example, it seems that interested stuff are being integrated out. Energy-momenutm conservation dominates lowest order of the angular structure.
- So how can we access the spatial information in  $\delta \epsilon(\tau, x, y, \eta_s)$ ?

## Hanbury-Brown–Twiss (HBT) correlation and spatial information



• Identical bosons are symmetrized

$$
\langle x_1, x_2 | a_{p_1}^{\dagger} a_{p_2}^{\dagger} | 0 \rangle = \frac{e^{ip_1 \cdot x_1 + ip_2 \cdot x_2} + e^{ip_2 \cdot x_1 + ip_1 \cdot x_2}}{\sqrt{2}}
$$

 The intensity to observe one or two identical bosons from a system specified by a density matrix  $\hat{\rho}$ 

$$
n(p) = \text{Tr}\left\{\hat{\rho}\hat{n}_p\right\} = \text{Tr}\left\{\hat{\rho}\hat{a}_{p_1}^{\dagger}\hat{a}_{p_1}\right\}
$$

$$
n(p_1, p_2) = \text{Tr}\left\{\hat{\rho}\hat{n}_{p_1}\hat{n}_{p_2}\right\} = \text{Tr}\left\{\hat{\rho}\hat{a}_{p_1}^{\dagger}\hat{a}_{p_1}\hat{a}_{p_2}^{\dagger}\hat{a}_{p_2}\right\}
$$

#### Earlier focus of HBT measurements



• The two-particle correlation function  $C(q, K)$  contains power spectrum of the Fourier transformed freeze-out surface.

#### Earlier focus of HBT measurements



$$
q^{\mu} = p_1^{\mu} - p_2^{\mu}, \quad K^{\mu} = (K_1^{\mu} + K_2^{\mu})/2,
$$
  

$$
S(q, K) = \int d^4x F(x, K) e^{iqx}
$$

$$
= \underbrace{\int_{\Sigma} K \cdot d^3 \sigma f(x, K) e^{iqx}}_{\text{Production on surface}} \underbrace{+ \cdots}_{\text{decays, etc}}
$$

$$
C\left(\vec{q}, \vec{K}\right) = \frac{N(p_1, p_2)}{N(p_1)N(p_2)} \approx 1 + \frac{|S(q, K)|^2}{|S(0, K)|^2}
$$

- The two-particle correlation function  $C(q, K)$  contains power spectrum of the Fourier transformed freeze-out surface.
- Larger- $q \Leftrightarrow$  finer structures on hypersurface. Study fireball size and inhomogenity. <sup>8</sup>

## With a jet passing through, what may change?



- Tuning q, we scan perturbation with  $\lambda \sim 1/q$  on the hypersurface.
- Compare  $C'(q, K)$  in hard-triggered events vs  $C(q, K)$  in events w/o hard trigger.
- Or, to reduce trigger biases, compare  $C'(q, K)$  in different directions of  $\vec{K}$ .

## An interesting poster from QM2015 in Japan

#### Poster<sup>2</sup> [by Naoto Tanaka, University of Tsukuba](https://utkhii.px.tsukuba.ac.jp/report/2015/QM2015/QMposter_tanaka_final.pdf)

#### **Analysis method**

If jet modification affects medium shape, azimuthally sensitive HBT should have the oscillation with respect to the leading iet axis.

In HBT analysis, momentum range is very low(pr:0.15-2.0 GeV/c). So this analysis will be sensitive not to size of jet itself but to the bulk response and re-distributed hadrons.

Recently non zero jet v<sub>2</sub> is observed<sup>[3]</sup>. Therefore HBT w.r.t. jet axis will also include  $\Psi_2$  HBT signal.

In order to understand jet modification in source shape, Selecting jet axis w.r.t.  $\Psi_2(\Psi_3)$  is important.

\*HBT w.r.t. jet axis 1-4, 2-3 should be symmetric about jet axis



Not sure if this was pursued, and what difficulty was found.

But now, we have (partly) the tools to do theoretical estimations.

## Status of the theoretical tools



- A. Hadrons produced on the bulk hypersurface.
- $B. + C.$  Resonance decay, jet fragmentation.

#### We have three types of two-particles correlations

- $\bullet$   $|A|^2$  Surface-surface correlation, can be treated with Cooper-Frye prescription  $(\checkmark)$ .
- $|B + C|^2$  Jet-jet (like) correlation, and  $\mathfrak{Re}(2A^*(B+C))$  Surface-jet correlation.
	- $\Rightarrow$  Maybe can generalize treatments of resonance decay in [Plumberg, Heinz PRC98(2018)034910] .



W. Chen et al. PRL127(2022)082301, Y. He et al. Phys. Rev. C 91, 054908, S. Cao el al. PRC94(2016)014909, T. Luo et al. PRC109(2024)034919

[W. Chen et al. PLB12(2017)015]

$$
\partial_{\mu} T^{\mu\nu} = J^{\nu}, \quad D_{\tau} \pi^{\mu\nu} = \cdots
$$
  

$$
J^{\nu} = -\frac{\partial}{\partial t} \int f_{\text{hard}}(t, x, p) p^{\nu} \Theta(p \cdot u > E_c) \frac{d^3 p}{(2\pi)^3}
$$
  

$$
(\partial_t + v \cdot \nabla_x) f_{\text{hard}}(t, x, p) = C_{2 \leftrightarrow 2} [f_{\text{hard}}] + \mathcal{R}_{1 \to 2}^{\text{eff}} [f_{\text{hard}}]
$$

### Estimate surface-surface HBT correlation from CoLBT

- CoLBT simulation set up (Z. Yang): controlled deposition of energy momentum along a trajectory  $dp^{\mu}/dt \propto \delta(\eta_s)\delta(x-t)\delta(y)[1, 1, 0, 0].$
- For simplicity, we first neglect viscous correction to the Bose-Einstein distribution function  $f_{\text{BE}}$  at freeze out.
- Consider collinear limit of the pair  $|K| \gg |q|$ , and focus on the region where  $\mathbb{q}^\mu = (0, q_\mathsf{x}, q_\mathsf{y}, 0)$  and  $\vec{\mathcal{K}}_\perp \perp \vec{q}_\perp$ , then

$$
S_{\text{surface}}(q, K) = \frac{g}{(2\pi)^2} \sum_{t_f, \vec{x_f}} K^{\mu} \sigma_{\mu}(t_f, \vec{x_f}) f_{\text{BE}}\left(\frac{K \cdot u_f}{T_{\text{frz}}}\right) J_0\left(q_T | \vec{x_f} - \vec{v_K} t_f|\right)
$$

#### Surface-surface correlation from a single event



$$
R(q, K) = \frac{C_{\text{with jet}}(q, K)}{C_{\text{no jet}}(q, K)}
$$

$$
\pi^+ \pi^+, \quad K = 1.5 \text{ GeV}.
$$

- Central  $Pb+Pb$  and  $O+O$  at LHC.
- Dashed: ideal hydro. Solid: viscous hydro.
- Colors: rotating  $\vec{K}$  from 0 (jet direction) to  $2\pi$ .

$$
S(q, K) = S_{bg} + S_{\text{pert}}
$$
\n
$$
\propto \sum_{t_f, \vec{x_f}} K^{\mu} \sigma_{\mu}(t_f, \vec{x}_f) f_{BE} \left( \frac{K \cdot u}{T_{\text{frz}}}\right) J_0 \left( q_T | \vec{x}_f - \vec{v}_K t_f | \right)
$$
\n
$$
+ \sum_{t_f, \vec{x_f}} K^{\mu} \sigma_{\mu}(t_f, \vec{x}_f) f_{BE}' \frac{K \cdot \delta u}{T_{\text{frz}}} J_0 \left( q_T | \vec{x}_f - \vec{v}_K t_f | \right)
$$
\n
$$
+ \sum_{t_f, \vec{x_f}} K^{\mu} \delta \sigma_{\mu}(t_f, \vec{x}_f) f_{BE} J_0 \left( q_T | \vec{x}_f - \vec{v}_K t_f | \right)
$$

- What determines the sign of correction. K parallel/anti-parallel to flow or freeze-out element corrections.
- What tells the  $q_T$  location of the peak. Inversely related to  $|x - vt|$ , whether the jet's trace on the surface is short or long.

#### How to understand the signal?



16

1. To consistently include jet-jet and jet-surface correlation.<br>0.50  $\leq K_r \leq 1.00$  (GeV) 0.20  $\leq K_r \leq 0.30$  (GeV)



- Pythia8 has a implementation of Bose-Einstein correlation in Lund-string hadronization (Left).
- Need to combine it with space-time information in medium-modified jet shower. Simulations from D.-X. Zhu
- 2. How to reduce signal cancellation when averaging over jet production vertex.
	- Select events that bias jet production location, e.g. [Z. Yang et al. PJC83(2023)652]
	- $\bullet$  Try the back-to-back limit (instead of collinear limit):  $K^0 \neq 0, \vec{K} = 0.$
- Is it possible to use the jet-induced medium response phenomena to study the nature of QGP response?
- We need spatial information. HBT correlation may be useful.
- Preliminary studies using CoLBT reveals interesting structure in surface-surface HBT: interplay of direction of flow, jet, and the direction of the pair.
- Need a careful event selection to preserve a large signal.
- Need to estimate jet-jet and jet-surface HBT correlations.
- Need more discussion and feedback on the feasibility!