# HBT correlation in HIC with a hard process (preliminary)

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## Medium response to the passage of a parton (jet, heavy quark, etc)





A demonstration using linearized MIS-type hydrodynamics First unambiguous evidence of diffusion wake at HP2024! Compared to simulations.

- Medium response: a natural consequence of jet quenching + medium dynamics.
- Will it tell something about the medium dynamics?

## Medium response in the spatial coordinates



- Jet excites QGP at "all" wave-length δ(x) → f(k) = 1. Small perturbations are characterized by linear response functions ⇔ G<sup>µν</sup><sub>αβ</sub> = ⟨T<sup>µν</sup>(t, x)T<sub>αβ</sub>(t', x')⟩.
- Kineitc v.s. hydro responses give fairly different spatial strucuters. However, we can only access momentum space, e.g., single-hadron spectra.

**Background** is a Bjorken flow. **Response** creates some small perturbations ( $\delta \ll 1$ )

$$\epsilon = \epsilon_0(\tau) + \delta e(\tau, x, y, \eta_s)$$
$$u^{\mu} = u_0^{\mu} + \delta u(\tau, x, y, \eta_s)$$
$$\pi^{\mu\nu} = \pi_0^{\mu\nu} + \delta \pi^{\mu\nu}(\tau, x, y, \eta_s)$$

Study the impact on particle production at freeze out defined by  $e(\tau, x, y, \eta_s) = e_f$ . Linearize the Cooper-Frye formula:



## How are information encoded in one-particle angular distribution? (toy study)

The possible angular structure is quite restricted.

$$\frac{d\delta N}{dm_T dy d\phi} = \frac{p^{\tau} \tau_{\rm frz,0} f_{eq}(\tilde{p}^{\tau})}{(2\pi)^3} \left(N_0 + \tilde{p}_{\mu} N_1^{\mu} + \tilde{p}_{\mu} \tilde{p}_{\nu} N_2^{\mu\nu} + \tilde{p}_{\mu} \tilde{p}_{\nu} \tilde{p}_{\rho} N_3^{\mu\nu\rho}\right) + \mathcal{O}\left(\delta^2\right), \quad \tilde{p}^{\mu} = \frac{p^{\mu}}{T}$$

All information of perturbations are encode in the spatially-integrated coefficients N.

$$\begin{split} \mathsf{N}_{0} &= \int d\mathsf{x} d\mathsf{y} d\eta_{s} \left[ c_{s}^{2} \tilde{p}^{\tau} \delta \tilde{\epsilon} + \delta \tilde{\epsilon} - \mathsf{v}^{\eta} \partial_{\eta} \delta \tilde{\epsilon} - \tau_{f} \mathsf{v}_{\perp} \cdot \partial_{\perp} \delta \tilde{\epsilon} \right], \\ \mathsf{N}_{1}^{\mu} &= \int d\mathsf{x} d\mathsf{y} d\eta_{s} \left[ -\delta \tilde{g}^{\mu} - \tilde{p}^{\tau} \tilde{\pi}^{\mu\nu} \delta \tilde{g}_{\nu} \right], \\ \mathsf{N}_{2}^{\mu\nu} &= \int d\mathsf{x} d\mathsf{y} d\eta_{s} \left[ \frac{1}{2} \tilde{\pi}^{\mu\nu} \left[ c_{s}^{2} \tilde{p}^{\tau} \delta \tilde{\epsilon} + \delta \tilde{\epsilon} - \mathsf{v}^{\eta} \partial_{\eta} \delta \tilde{\epsilon} - \tau_{f} \mathsf{v}_{\perp} \cdot \partial_{\perp} \delta \tilde{\epsilon} \right] + \frac{1}{2} \Delta_{\alpha\beta}^{\mu\nu} \delta \tilde{\pi}^{\alpha\beta} \right], \\ \mathsf{N}_{3}^{\mu\nu\rho} &= \int d\mathsf{x} d\mathsf{y} d\eta_{s} \frac{1}{2} \tilde{\pi}^{\mu\nu} (-\delta \tilde{g}^{\rho}) \quad \text{where} \quad \delta \tilde{\epsilon} = \frac{\delta \epsilon}{\epsilon + P}, \quad \delta \tilde{g}^{\mu} = \frac{\delta g^{\mu}}{\epsilon + P} = \delta u^{\mu}. \end{split}$$

No background radial flow, it will be more interesting with radial flow.



- In this toy example, it seems that interested stuff are being integrated out. Energy-momenutm conservation dominates lowest order of the angular structure.
- So how can we access the spatial information in  $\delta \epsilon(\tau, x, y, \eta_s)$ ?



$$\frac{d\delta N}{dm_{T} dy d\phi} = N_{0} + \tilde{p}_{\mu} N_{1}^{\mu} + \tilde{p}_{\mu} \tilde{p}_{\nu} N_{2}^{\mu\nu} + \tilde{p}_{\mu} \tilde{p}_{\nu} \tilde{p}_{\rho} N_{3}^{\mu\nu\rho}$$
$$= C_{00} \left(\tilde{p}^{\tau}\right) + C_{1m} \left(\tilde{p}^{\tau}\right) Y_{1}^{m}(\Omega)$$
$$+ C_{2m} \left(\tilde{p}^{\tau}\right) Y_{2}^{m}(\Omega) + C_{3m} \left(\tilde{p}^{\tau}\right) Y_{3}^{m}(\Omega) + \mathcal{O}(\delta^{2})$$

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## Hanbury-Brown–Twiss (HBT) correlation and spatial information



• Identical bosons are symmetrized

$$\langle x_1, x_2 | a^{\dagger}_{p_1} a^{\dagger}_{p_2} | 0 
angle = rac{e^{i p_1 \cdot x_1 + i p_2 \cdot x_2} + e^{i p_2 \cdot x_1 + i p_1 \cdot x_2}}{\sqrt{2}}$$

• The intensity to observe one or two identical bosons from a system specified by a density matrix  $\hat{\rho}$ 

$$\begin{split} n(p) &= \operatorname{Tr} \left\{ \hat{\rho} \hat{n}_{p} \right\} = \operatorname{Tr} \left\{ \hat{\rho} \hat{a}_{p_{1}}^{\dagger} \hat{a}_{p_{1}} \right\} \\ n(p_{1}, p_{2}) &= \operatorname{Tr} \left\{ \hat{\rho} \hat{n}_{p_{1}} \hat{n}_{p_{2}} \right\} = \operatorname{Tr} \left\{ \hat{\rho} \hat{a}_{p_{1}}^{\dagger} \hat{a}_{p_{1}} \hat{a}_{p_{2}}^{\dagger} \hat{a}_{p_{2}} \right\} \end{split}$$

### Earlier focus of HBT measurements



• The two-particle correlation function C(q, K) contains power spectrum of the Fourier transformed freeze-out surface.

## Earlier focus of HBT measurements



$$q^{\mu} = p_{1}^{\mu} - p_{2}^{\mu}, \quad K^{\mu} = (K_{1}^{\mu} + K_{2}^{\mu})/2$$

$$S(q, K) = \int d^{4}x F(x, K) e^{iqx}$$

$$= \underbrace{\int_{\Sigma} K \cdot d^{3}\sigma f(x, K) e^{iqx}}_{\text{Production on surface}} \underbrace{+\cdots}_{\text{decays, etc}}$$

$$C\left(\vec{q}, \vec{K}\right) = \frac{N(p_{1}, p_{2})}{N(p_{1})N(p_{2})} \approx 1 + \frac{|S(q, K)|^{2}}{|S(0, K)|^{2}}$$

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- The two-particle correlation function C(q, K) contains power spectrum of the Fourier transformed freeze-out surface.
- Larger- $q \Leftrightarrow$  finer structures on hypersurface. Study fireball size and inhomogenity.

## With a jet passing through, what may change?



- Tuning q, we scan perturbation with  $\lambda \sim 1/q$  on the hypersurface.
- Compare C'(q, K) in hard-triggered events vs C(q, K) in events w/o hard trigger.
- Or, to reduce trigger biases, compare C'(q, K) in different directions of  $\vec{K}$ .

## An interesting poster from QM2015 in Japan

#### Posterc³ by Naoto Tanaka, University of Tsukuba

#### **Analysis method**

If jet modification affects medium shape, azimuthally sensitive HBT should have the oscillation with respect to the leading jet axis.

In HBT analysis, momentum range is very low(p::0.15-2.0 GeV/c). So this analysis will be sensitive not to size of jet itself but to the bulk response and re-distributed hadrons.

Recently **non zero jet v**\_2 is observed<sup>[3]</sup>. Therefore HBT w.r.t. jet axis will also include  $\Psi_2$  HBT signal.

In order to understand jet modification in source shape, Selecting jet axis w.r.t.  $\Psi_2(\Psi_3)$  is important.

\*HBT w.r.t. jet axis ①-④, ②-③ should be symmetric about jet axis



Not sure if this was pursued, and what difficulty was found.

But now, we have (partly) the tools to do theoretical estimations.

## Status of the theoretical tools



- A. Hadrons produced on the bulk hypersurface.
- B. + C. Resonance decay, jet fragmentation.

#### We have three types of two-particles correlations

- |A|<sup>2</sup> Surface-surface correlation, can be treated with Cooper-Frye prescription (√).
- $|B + C|^2$  Jet-jet (like) correlation, and  $\mathfrak{Re}(2A^*(B + C))$  Surface-jet correlation.
  - $\Rightarrow$  Maybe can generalize treatments of resonance

decay in [Plumberg, Heinz PRC98(2018)034910] .



W. Chen et al. PRL127(2022)082301,
Y. He et al. Phys. Rev. C 91, 054908,
S. Cao el al. PRC94(2016)014909,
T. Luo et al. PRC109(2024)034919

[W. Chen et al. PLB12(2017)015]

$$\partial_{\mu} T^{\mu\nu} = J^{\nu}, \quad D_{\tau} \pi^{\mu\nu} = \cdots$$
  
 $J^{\nu} = -\frac{\partial}{\partial t} \int f_{hard}(t, x, p) p^{\nu} \Theta(p \cdot u > E_c) \frac{d^3 p}{(2\pi)^3}$   
 $(\partial_t + v \cdot \nabla_x) f_{hard}(t, x, p) = C_{2\leftrightarrow 2} [f_{hard}] + \mathcal{R}_{1\rightarrow 2}^{eff} [f_{hard}]$ 

## Estimate surface-surface HBT correlation from CoLBT

- CoLBT simulation set up (Z. Yang): controlled deposition of energy momentum along a trajectory  $dp^{\mu}/dt \propto \delta(\eta_s)\delta(x-t)\delta(y)[1,1,0,0]$ .
- For simplicity, we first neglect viscous correction to the Bose-Einstein distribution function  $f_{\rm BE}$  at freeze out.
- Consider collinear limit of the pair  $|K| \gg |q|$ , and focus on the region where  $q^{\mu} = (0, q_x, q_y, 0)$  and  $\vec{K}_{\perp} \perp \vec{q}_{\perp}$ , then

$$S_{\rm surface}(q,K) = \frac{g}{(2\pi)^2} \sum_{t_f,\vec{x_f}} K^{\mu} \sigma_{\mu}(t_f,\vec{x_f}) f_{\rm BE}\left(\frac{K \cdot u_f}{T_{\rm frz}}\right) J_0\left(q_T |\vec{x_f} - \vec{v_K} t_f|\right)$$

### Surface-surface correlation from a single event



$$egin{aligned} R(q, \mathcal{K}) &= rac{C_{ ext{with jet}}(q, \mathcal{K})}{C_{ ext{no jet}}(q, \mathcal{K})} \ \pi^+ \pi^+, \quad \mathcal{K} &= 1.5 ext{ GeV}. \end{aligned}$$

- Central Pb+Pb and O+O at LHC.
- Dashed: ideal hydro. Solid: viscous hydro.
- Colors: rotating  $\vec{K}$  from 0 (jet direction) to  $2\pi$ .

 $S(q,K)=S_{bg}+S_{
m pert}$ 

$$\begin{split} &\propto \sum_{t_f, \vec{x_f}} \mathcal{K}^{\mu} \sigma_{\mu}(t_f, \vec{x}_f) f_{BE} \left( \frac{\mathcal{K} \cdot u}{\mathcal{T}_{\mathrm{frz}}} \right) J_0 \left( q_T \left| \vec{x}_f - \vec{v}_{\mathcal{K}} t_f \right| \right. \\ &+ \sum_{t_f, \vec{x_f}} \mathcal{K}^{\mu} \sigma_{\mu}(t_f, \vec{x}_f) f_{BE}' \frac{\mathcal{K} \cdot \delta u}{\mathcal{T}_{\mathrm{frz}}} J_0 \left( q_T \left| \vec{x}_f - \vec{v}_{\mathcal{K}} t_f \right| \right) \\ &+ \sum_{t_f, \vec{x_f}} \mathcal{K}^{\mu} \delta \sigma_{\mu}(t_f, \vec{x}_f) f_{BE} J_0 \left( q_T \left| \vec{x}_f - \vec{v}_{\mathcal{K}} t_f \right| \right) \end{split}$$

- What determines the sign of correction. *K* parallel/anti-parallel to flow or freeze-out element corrections.
- What tells the q<sub>T</sub> location of the peak. Inversely related to |x vt|, whether the jet's trace on the surface is short or long.

#### How to understand the signal?



1. To consistently include jet-jet and jet-surface correlation.



- Pythia8 has a implementation of Bose-Einstein correlation in Lund-string hadronization (Left).
- Need to combine it with space-time information in medium-modified jet shower. Simulations from D.-X. Zhu
- 2. How to reduce signal cancellation when averaging over jet production vertex.
  - Select events that bias jet production location, e.g. [Z. Yang et al. PJC83(2023)652]
  - Try the back-to-back limit (instead of collinear limit):  $K^0 \neq 0, \vec{K} = 0$ .

- Is it possible to use the jet-induced medium response phenomena to study the nature of QGP response?
- We need spatial information. HBT correlation may be useful.
- Preliminary studies using CoLBT reveals interesting structure in surface-surface HBT: interplay of direction of flow, jet, and the direction of the pair.
- Need a careful event selection to preserve a large signal.
- Need to estimate jet-jet and jet-surface HBT correlations.
- Need more discussion and feedback on the feasibility!