# Light and heavy flavor jet quenching in relativistic heavy-ion collisions



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## Jet quenching in high-energy nuclear collisions



### **Nuclear modification factor**

 $R_{AA} \equiv rac{d^2 N^{AA}/dy dp_{\perp}}{d^2 N^{pp}/dy dp_{\perp} imes \langle N^{AA}_{coll} 
angle}$ 



[Mueller et al., Ann. Rev. Nucl. Part. Sci. 62, 361 (2012)]

## Jets tagged with heavy quarks





- Produced from initial hard scatterings
- Serve as an ideal probe of the QGP properties
- Provide a unique opportunity for studying the flavor dependence of parton splitting (dead cone effect)

## Searches for the flavor dependence of parton splitting

### Hadron *R*<sub>AA</sub> (parton energy loss)



No clear separation between charged hadrons, D, and B, except at very low  $p_{T}$ 

### Goals:

- Understand flavor hierarchies embedded in both hadrons and jets Use hadron and jet observables to probe the QGP properties

### Distribution of splitting angles in pp



Clear suppression of splitting at small  $\theta$ in *D*-jets *vs*. inclusive jets

## **Theoretical framework of jet quenching**



$$d\sigma_{h} = \sum_{abjd} f_{alp} \otimes f_{blp} \otimes d\sigma_{ab \to jd} \otimes D_{hlj}$$

- $f_{a/p}, f_{b/p} \rightarrow f_{a/A}, f_{b/B}$ : cold nuclear matter (initial state) effect, e.g., shadowing, Cronin, ..., measured in pA collisions



$$d\tilde{\sigma}_{h} = \sum_{abjd} f_{a/A} \otimes f_{b/B} \otimes d\sigma_{ab \rightarrow jd} \otimes \tilde{D}_{h/j}$$

•  $D_{h/i} \rightarrow D_{h/i}$ : medium modified fragmentation function, hot nuclear matter (final state) effect • Factorization assumption:  $\tilde{D}_{h/j} = \sum_{i'} P_{j \to j'} \otimes D_{h/j'}$ , nuclear modification of parton j

## Parton transport inside the QGP

### Linear Boltzmann Transport (LBT)

 $p_a \cdot \partial f_a(x_a, p_a) = E_a(\mathscr{C}_a^{\text{el}} + \mathscr{C}_a^{\text{inel}})$ 

## Parton transport inside the QGP

### Linear Boltzmann Transport (LBT)

**Elastic scattering (** $ab \rightarrow cd$ **)** 

$$\mathscr{C}_{a}^{\text{el}} = \sum_{b,c,d} \int \prod_{i=b,c,d} \frac{d[p_i]}{2E_a} (\gamma_d f_c f_d - \gamma_b f_a f_b) \cdot (2\pi)^4 \delta^4 (p_a + p_b - p_c - p_d) \left| \mathscr{M}_{ab \to cd} \right|^2$$

 $p_a \cdot \partial f_a(x_a, p_a) = E_a(\mathscr{C}_a^{\text{el}} + \mathscr{C}_a^{\text{inel}})$ 

 $2 \rightarrow 2$  scattering matrices

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loss term: scattering rate (for Monte-Carlo simulation)

$$\Gamma_a^{\text{el}}(\mathbf{p}_a, T) = \sum_{b,c,d} \frac{\gamma_b}{2E_a} \int \prod_{i=b,c,d} d[p_i] f_d$$

 $p_a \cdot \partial f_a(x_a, p_a) = E_a(\mathscr{C}_a^{\text{el}} + \mathscr{C}_a^{\text{inel}})$ 

 $2 \rightarrow 2$  scattering matrices

 $f_b \cdot (2\pi)^4 \delta^{(4)}(p_a + p_b - p_c - p_d) |\mathcal{M}_{ab \to cd}|^2$ 

## **Inelastic scattering**





• Medium information absorbed in  $\hat{q} \equiv d \langle p_{\perp}^2 \rangle / dt$ 

[ Majumder PRD 85 (2012); Zhang, Wang and Wang, PRL 93 (2004)]

• Higher-twist formalism: collinear expansion (  $\langle k_{\perp}^2 \rangle \ll l_{\perp}^2 \ll Q^2$  )

$$\frac{1}{4}\sin^2\left(\frac{t-t_i}{2\tau_f}\right)$$

## Flavor hierarchy in hadron suppression





## Flavor hierarchy in hadron suppression

### NLO initial production and fragmentation + Boltzmann transport + hydrodynamic medium for QGP

### 1.2 CMS 0-10% ALICE 0-5% 0.8 R 44 0.6 0.4 100 10 p<sub>T</sub> (GeV)

### charged hadron

- g-initiated h & D  $R_{AA} < q$ -initiated h & D  $R_{AA} \left[ \Delta E_q > \Delta E_{q/c} \right]$



•  $R_{AA}(c > D) > R_{AA}(q > h) [\Delta E_q > \Delta E_c], R_{AA}(q > D) < R_{AA}(q > h) [different FFs] => R_{AA}(h) \approx R_{AA}(D)$ Signature of flavor hierarchy of parton ΔE offset by NLO production/fragmentation in hadron R<sub>AA</sub>



## Flavor hierarchy in hadron suppression



- starting from  $p_T \sim 8 \text{ GeV}$
- confirmation from future precision measurement

• A simultaneous description of charged hadron, D meson, B meson, B-decay D meson R<sub>AA</sub>'s

• Predict R<sub>AA</sub> separation between B and h / D below 40 GeV, but similar values above – wait for



## Extraction of parton energy loss from hadron RAA

- Mean  $p_T$  loss:  $\langle \Delta p_T^J \rangle = C_i \beta_g p_T^\gamma \log(p_T)$ 
  - $\beta_g$ : overall magnitude for g
  - $C_i$ : flavor dependence
  - $\gamma$ :  $p_T$  dependence
- $p_{T}$  loss distribution:

$$W_{AA}(x) = \frac{\alpha^{\alpha} x^{\alpha - 1} e^{-\alpha x}}{\Gamma(\alpha)}$$

$$x \equiv \Delta p_{\rm T} / \langle \Delta p_{\rm T} \rangle$$



**Bayesian calibration to data** 

### NLO initial production and fragmentation + Parametrized parton energy loss inside the QGP

[Xing, SC, Qin, Phys. Lett. B 850 (2024) 138523]



**Constraints on parameters** 



## **Extraction of parton energy loss from hadron RAA**

### Average energy loss



•  $\Delta E_g > \Delta E_q \sim \Delta E_c > \Delta E_b$ 

- Flavor hierarchy of parton energy loss is encoded in the hadron R<sub>AA</sub> data
- between parton energy loss and NLO production and fragmentation

### **Energy loss distribution**



More stringent test on QCD calculation

• No obvious hierarchy for the hadron  $R_{AA}$  data themselves, due to the interplay

## From hadrons to full jets



- Jet partons and medium background cannot be cleanly separated in reality •
- Jet-medium interactions: medium modification of jets + medium response

Medium response (energy deposition + depletion) is naturally included in all jet observables

## Jet R<sub>AA</sub> and v<sub>2</sub>

### RAA



- Including medium response reduces jet energy loss and thus increases the jet RAA



• With  $R_{AA}$  fixed, including medium response (coupled to medium flow) increases the jet  $v_2$ 

### Jet substructure



### Transverse (*r*) distribution: jet shape



[Tachibana, Chang, Qin, Phys. Rev. C 95] (2017) 044909 ]

### Longitudinal (z) distribution: jet fragmentation function



[Chen, SC, Luo, Pang, Wang, Phys. Lett. B 777 (2018) 86-90 ]



## Search for unique signatures of medium response

### **Energy suppression in diffusion wake**



- SC, Luo, Pang, Wang, PLB 777 (2018) 86, Yang, Luo, Chen, Pang, Wang, PRL 130 (2023) 052301 ]
- Confirmed by recent CMS data [CMS-PAS-HIN-23-006]

Energy suppression predicted in the backward direction of jets at  $1 < p_T^h < 2$  GeV [Chen,

## Search for unique signatures of medium response

### Hadron chemistry around quenched jets

### **Baryon enhancement**



- Larger quark density and strangeness density in QGP than in vacuum jets
- Stronger enhancement at larger distance from the jet axis

### Strangeness enhancement



Enhanced baryon-to-meson ratio and strangeness around jets in AA vs. pp collisions

[Luo, Mao, Qin, Wang, Zhang, PLB 837 (2023) 137638; Chen, SC, Luo, Pang, Wang, NPA 1005 (2021) 121934]

## A novel observable: energy-energy correlator (EEC)



[Komiske et. al., PRL 130 (2023) 051901]

- EEC:  $\frac{d\Sigma}{dR_L} = \int d\vec{n}_1 d\vec{n}_2 \frac{\langle \mathscr{E}(\vec{n}_1)\mathscr{E}(\vec{n}_2) \rangle}{Q^2} \delta(\Delta R_{12} - R_L)$
- EEC of jets presents a clear angular scale separation between perturbative and nonperturbative (e.g. hadronization) regions



[Craft et. al., arXiv:2210.09311]

 $\mathscr{E}$ : energy flow in a given direction,  $\Delta R_{12} = \sqrt{\Delta \phi_{12}^2 + \Delta \eta_{12}^2}$  : relative angle, Q: hard scale

• EEC can also reveal the flavor dependence of splitting angles of partons in pp collisions Implement a first realistic calculation on light and heavy flavor jet EEC in AA collisions

## Light vs. heavy flavor jet EEC in pp collisions



• Jet in pp: Pythia 8 simulation

• EEC analysis (*i*, *j* denote jet constituents)

$$\frac{d\Sigma(\theta)}{d\theta} = \frac{1}{\Delta\theta} \sum_{\substack{|\theta_{ij} - \theta| < \Delta\theta/2}} \frac{p_{\mathrm{T},i}(\vec{n}_i) p_{\mathrm{T},j}(\vec{n}_j)}{p_{\mathrm{T},j\text{et}}^2}$$

• Flavor (mass) dependence:

- Overall magnitude: charged jet > D-jet > B-jet
- Typical (peak) angle: charged jet < D-jet < B-jet

Suppression of splitting within  $\theta_0 \sim m_Q/E_Q$  in vacuum

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• Jet energy dependence

• Higher  $p_T \rightarrow \Sigma$  peaks at smaller  $\theta$ 

 $p_{\rm T}\theta_{\rm peak}$  ~ transition scale between pert. and non-pert.

## EEC of partons developed from a single quark



- Single quark  $\rightarrow$  LBT + static medium  $\rightarrow$  EEC of daughter partons
- Flavor (mass) hierarchy of EEC:

  - Clear strong suppression of  $\Sigma$  below  $\theta_0 \sim m_Q/E_{\rm initial}$
- Contributions form medium response and gluon emission show similar hierarchies

Xing, SC, Qin, Wang, arXiv:2409.12843

• Magnitude: charged > D > B-jet; peak position: charged < D < B-jet (similar to vacuum jets)

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## Light vs. heavy flavor jet EEC in central PbPb collisions



Nuclear modification (AA - pp) — Pythia + LBT in hydro

- and large  $\theta$
- tagged with heavier mesons



• General features: suppression at intermediate  $\theta$ , enhancement at small  $\theta$  (except for *B*-jet)

• Flavor hierarchy: weaker nuclear modification (both suppression and enhancement) for jets

## **Different contributions to medium modification on EEC**



- Medium response enhances EEC at large  $\theta$

S: shower partons inherited from Pythia S+R: add medium-induced gluons S+R+M: further add medium response

• Jet energy loss causes suppression over the entire  $\theta$  region - Medium-induced gluon emission enhances EEC at small  $\theta$ 

## Effects of trigger bias on the jet EEC

### $p_{T}$ trigger in both pp and AA



- AA jets with trigger bias originate from pp jets with higher  $p_T$  and initial virtuality scale → Stronger but narrower vacuum splittings
- Can be tested using  $\gamma$ -jets

 $p_{T}$  trigger only in pp (no trigger bias in AA)



 $\rightarrow$  Enhances EEC at small  $\theta$ , reduces the suppression/enhancement at intermediate/large  $\theta$ 

## **Constraints on jet transport coefficient inside the QGP**



[JET, Phys. Rev.C 90 (2014) 1, 014909]

 $\hat{q} \equiv d\langle k_{\perp}^2 \rangle / dt \sim \langle F^{ai+}(0)F_i^{a+}(y^-) \rangle$ 

- QGP is much more opaque than cold nuclear matter to jet propagation
- Recent developments on  $\hat{q}$  extraction:

Multistage jet evolution model with Bayesian analysis

[JETSCAPE, Phys. Rev. C 104 (2021) 1, 024905]

Information field based global interference

[Xie et al., Phys. Rev. C 108 (2023) 1, L011901]

## Probing the equation of state of the QGP

[F.-L. Liu, X.-Y. Wu, SC, G-Y. Qin, X.-N. Wang, Phys. Lett. B 848 (2024) 138355]

### Transport

 $p_a \cdot \partial f_a(x_a, p_a) = E_a(\mathscr{C}_a^{el} + \mathscr{C}_a^{inel})$ 

### **Strong coupling strength** g(E,T)

### **Equation of state**

 $P_{qp}(m_u, m_d, \dots, T) = \sum_{i=u,d,s,g} d_i \int \frac{d^3p}{(2\pi)^3} \frac{\left|\vec{p}\right|^2}{3E_i(p)} f_i(p) - B(T)$  $= \Sigma_i P_{kin}^i(m_i, T) - B(T)$  $\epsilon = TdP(T)/dT - P(T), \quad s = (\epsilon + P)/T$ 

### **Thermal mass of partons**

$$\begin{split} m_g^2 &= \frac{1}{6} g^2 \left[ (N_c + \frac{1}{2} n_f) T^2 + \frac{N_c}{2\pi^2} \Sigma_q \mu_q^2 \right] \\ m_{u,d}^2 &= \frac{N_c^2 - 1}{8N_c} g^2 \left[ T^2 + \frac{\mu_{u,d}^2}{\pi^2} \right] \\ m_s^2 - m_{0s}^2 &= \frac{N_c^2 - 1}{8N_c} g^2 \left[ T^2 + \frac{\mu_s^2}{\pi^2} \right] \end{split}$$

**Strategy:** Fit g from comparing transport model to data Calculate EoS from g

## **EoS of QGP and diffusion coefficient of heavy quarks**

### **Equation of state**



- Agreement with the lattice data

### **Diffusion coefficient**



### Simultaneous constraint on QGP properties and transport properties of hard probes

## **Constraints on the E&M field with the D meson v**<sub>1</sub>





## **Constraints on the E&M field with the D meson v**<sub>1</sub>



- Strong E&M field dominates at the LHC energy
- [Jiang, SC, Xing, Wu, Yang, Zhang, Phys. Rev. C 105 (2022) 5, 054907]



 Tilted geometry w.r.t. the beam direction dominates at the RHIC energy • Sensitivity of the D meson v<sub>1</sub> to different E&M evolution profiles at the LHC

## **Summary**

### Probing strongly interacting matter using energetic hadrons and jets

- Flavor hierarchy of parton energy loss is encoded in the hadron R<sub>AA</sub>, though not explicit due to the interplay between energy loss and NLO contributions
- The jet EEC is an excellent observable to study the dead cone effect on parton splitting (magnitude and peak position of EEC) in both pp and AA collisions
- Jet and heavy flavor observables can constrain various QGP properties



Thank. you!