



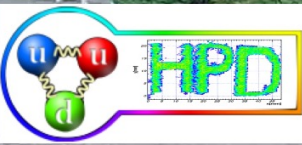
Workshop on Advances, Innovations, and Future Perspectives in High-Energy Nuclear Physics



MINISTERUL CERCETĂRII,
INOVĂRII ȘI DIGITALIZĂRII



QCD Challenges



Mihai Petrovici, Wuhan Workshop, 19-24 October, 2024

Outline:

➤ Introduction

Advances in High Energy Nuclear Physics:

- *Do we see a new state of deconfined matter at LHC ?*
 - $\langle p_T \rangle / [(dN/dy)/S_{\perp}]^{1/2}$ centrality and collision energy dependence
 - $[(dN/dy)/S_{\perp}]^{1/2}$ scaling
 - $\langle dE_T/dy \rangle / \langle dN/dy \rangle - \langle dN/dy \rangle / S_{\perp}$ correlation
 - The slope of $\varepsilon_{Bj} \cdot \tau - \langle dN/dy \rangle / S_{\perp}$ correlation - energy dependence
 - $(dN/dy)_{(strange \text{ and multi strange})} / (dN/dy) - (dN/dy) / S_{\perp}$ correlation
 - collision energy dependence of $(1-RAA) / [(dN/dy)/S_{\perp}]$ for central collisions
- *Similar studies for pp collisions and comparison with Pb-Pb collisions*

Future Perspectives:

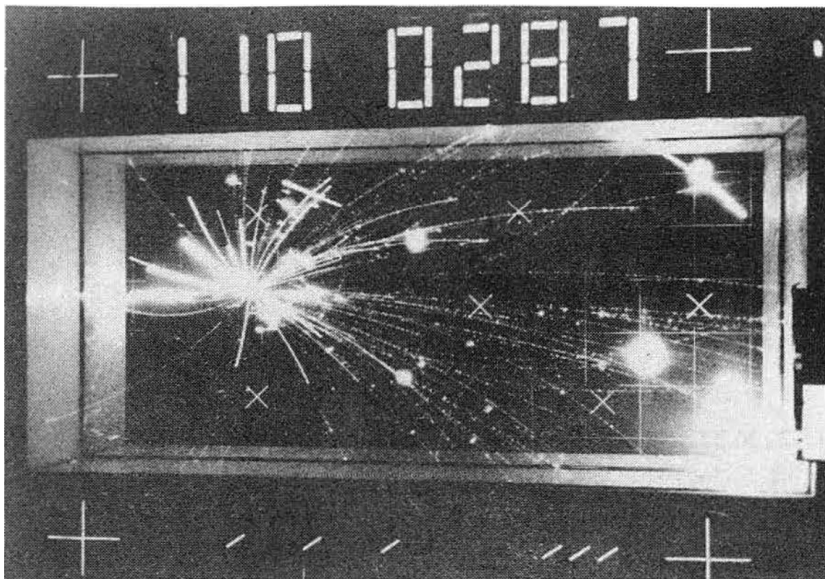
- *pp much larger charged particle multiplicity*
- *Pb-Pb large rapidity PID*
- *What remains to be done at large baryon density ?*

Innovations:

- *A new generation of RPC and TRD with 2D position resolution high granularity, high counting rate, radiation hard*
- *Concluding remark*

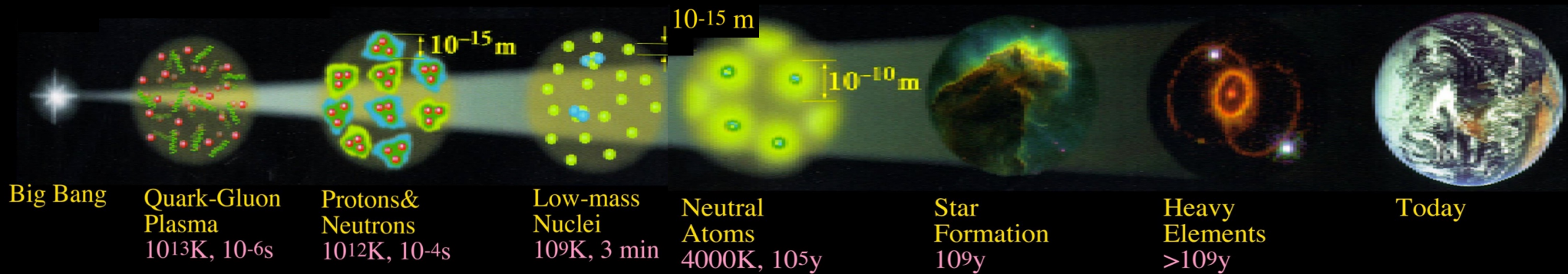
50th anniversary of high energy heavy-ion

- *The high-energy heavy-ion program at LBL has started in summer 1974 (CERN Courier, June 1974)*
- *A University of Frankfurt group has exposed their AgCl detectors to various heavy-ion beams at energies from 250 MeV/A to 2.1 GeV/A. The observed peaks in the angular distributions of light fragments that moved with beam energy in a manner suggestive of these particles arising from shock waves, causing considerable excitement in the nuclear science community.*
- *After being used for several high energy experiments, the LBL streamer chamber used in the collision of 1.8-GeV/nucleon Ar on a lead oxide target, evidenced charged particle multiplicities of over 100 in such reactions.*



<https://escholarship.org/uc/item/8bw3436f>

Could we unravel the History of the Universe



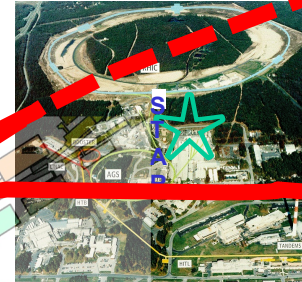
**based on experiments
in terrestrial laboratories ?**

Large scale facilities

LHC: Collider
Pb+Pb @5020GeV/A



RHIC: Collider
Au+Au @ 200GeV/A



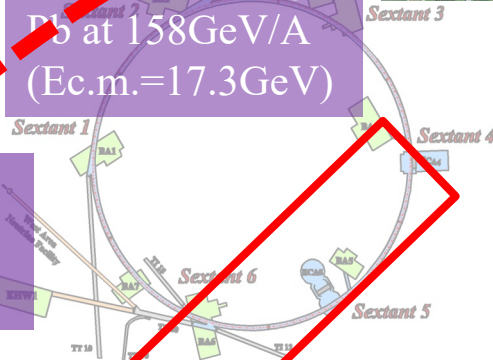
Hotter
Denser
Longer
Bigger



?

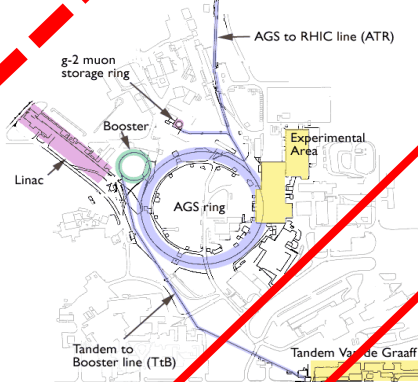
Click on the area of Interest

SPS: Fixed Target
Pb at 158GeV/A
(Ec.m.=17.3GeV)



???

AGS: Fixed Target
Au at 11.7GeV/A
(Ec.m.=4.86GeV)



BES

Bevalac
Fixed Target
1-2GeV/A

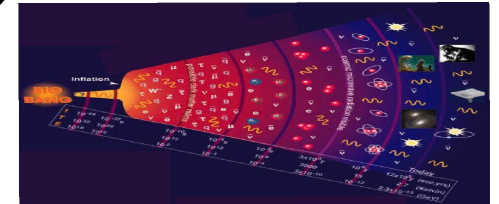
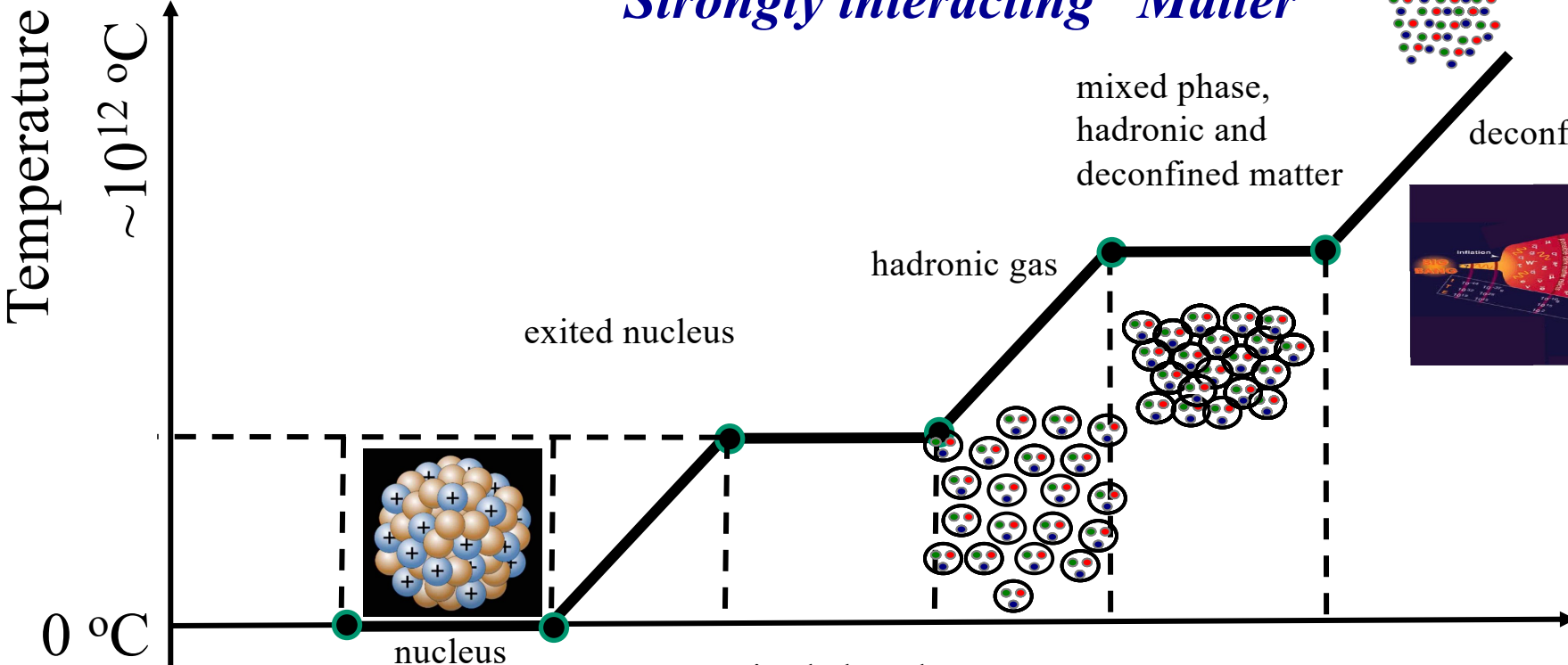


SIS 18



Physics motivation

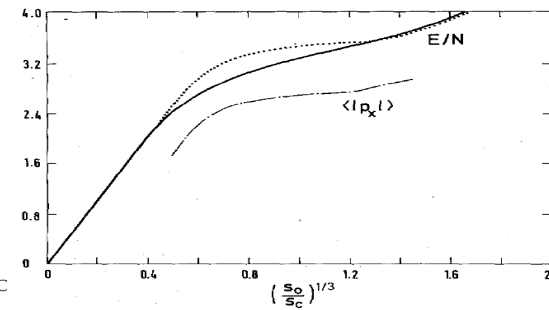
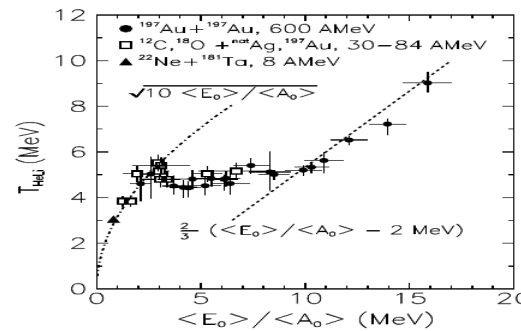
Strongly interacting "Matter"



them.") The elder Bohr, as a young graduate student in 1905, had written a prize-winning paper on the vibration of liquid drops of water. Seventy years later his son is being honored for work growing out of the liquid-drop picture.

mixed phased

thermal energy



*J.Pochodzalla et al.,
 ALADIN Coll.,
 arXiv:[nucl-ex]9607004*

*J.-P. Blaizot and J.-Y. Ollitrault,
 Phys.Lett 191B(1987)21*

Theory predictions

String percolation

T.S.Biro, H.B.Nielsen and J.Knoll, Nucl.Phys. B245(1984)449
J.Dias de Deus and C. Pajares, Phys.Lett. B695(2011)211
I. Bautista et al., Revista Mexicana de Fisica 65(2019)197

$$\frac{dN}{dy} = F(\eta)\bar{N}^s\mu$$

$\eta \equiv (r_0/R)^2\bar{N}^s$ - transverse string density; \bar{N}^s - the average number of strings
 μ - string multiplicity

$$F(\eta) \equiv \sqrt{\frac{1-e^{-\eta}}{\eta}}$$

$$\langle p_T^2 \rangle = \langle p_T^2 \rangle_1 / F(\eta) \quad \langle p_T^2 \rangle_1 - \text{average string transverse momentum}$$

$$\sqrt{\langle p_T^2 \rangle} / \sqrt{\langle dN/dy \rangle / S_{\perp}} \sim 1 / \sqrt{(1-e^{-\eta})}$$

$$\langle p_T^2 \rangle / [(\langle dn/dy \rangle / S_{\perp})] \propto \langle p_T^2 \rangle_1 r_0^2 / \mu (1-e^{-\eta})$$

CGC

Local parton-hadron duality picture
and dimensionality argument

- *Y.L.Dokshitzer, V.A.Khoze and S.Troian, J.Phys.G 17 (1991) 1585*
 - *T. Lappi, Eur.Phys.J. C71 (2011) 1699*
 - *E. Levin and A.H. Rezaeian, Phys.Rev.D 83 (2011)114001*

$$\langle p_T \rangle / \sqrt{\langle dN/dy \rangle / S_{\perp}} \sim \frac{1}{n\sqrt{n}}$$

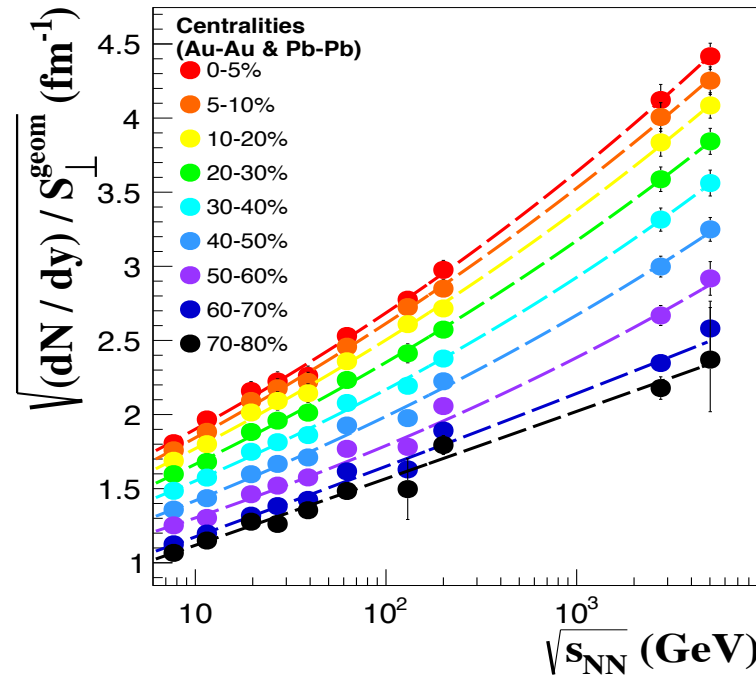
n - no. of charged particles from a gluon fragmentation



$$\langle p_T \rangle / \sqrt{\langle dN/dy \rangle / S_{\perp}}$$

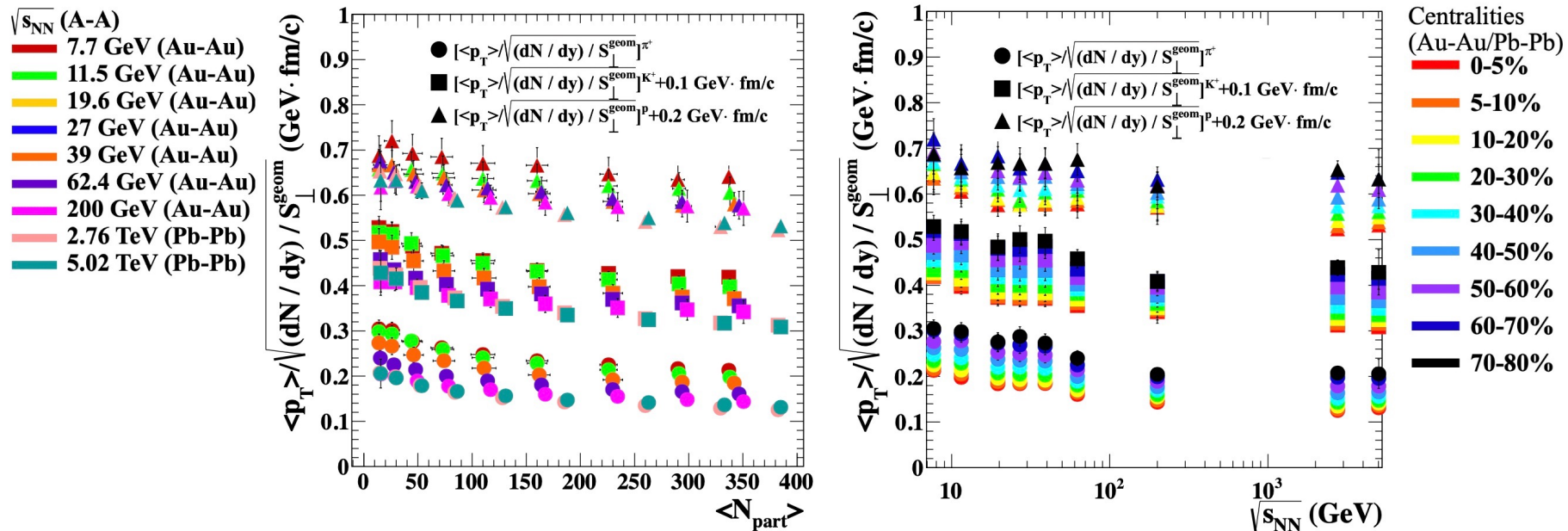
decreases as a function of:

- collision energy
- centrality

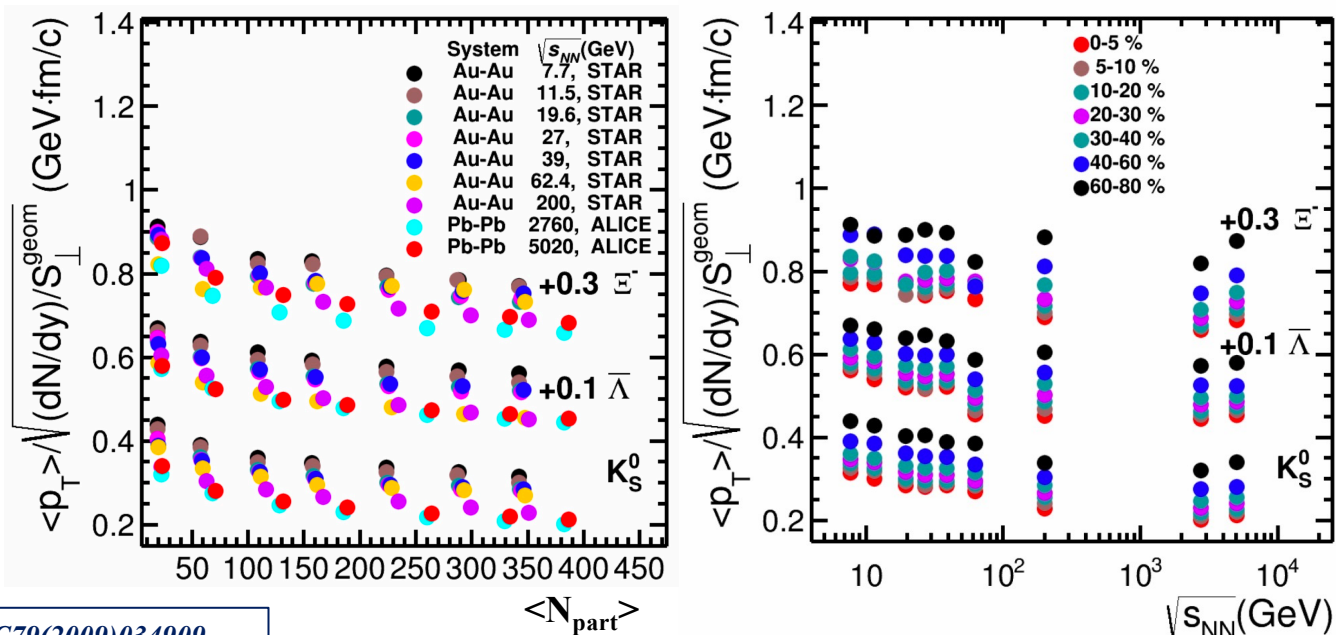


M.Petrovici, A.Lindner and A.Pop, Phys. Rev. C 98(2018)024904

Experimental results



M. Petrovici, A. Lindner and A. Pop, Phys. Rev. C 98(2018)024904



STAR Collaboration, Phys. Rev. C 79(2009)034909

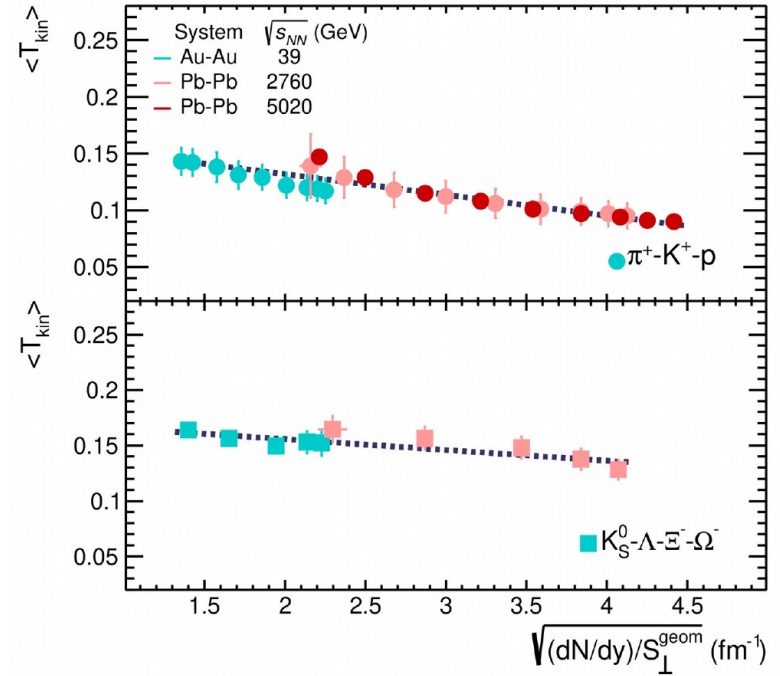
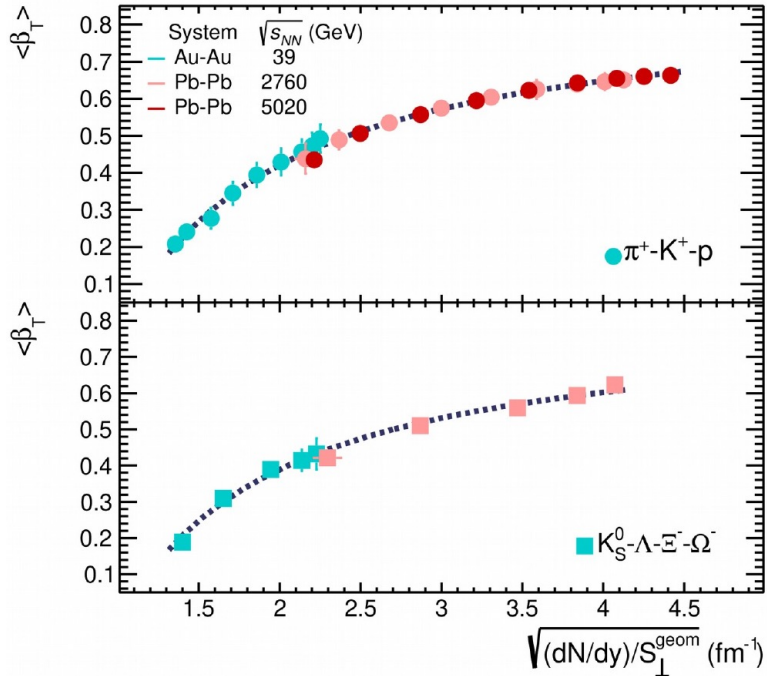
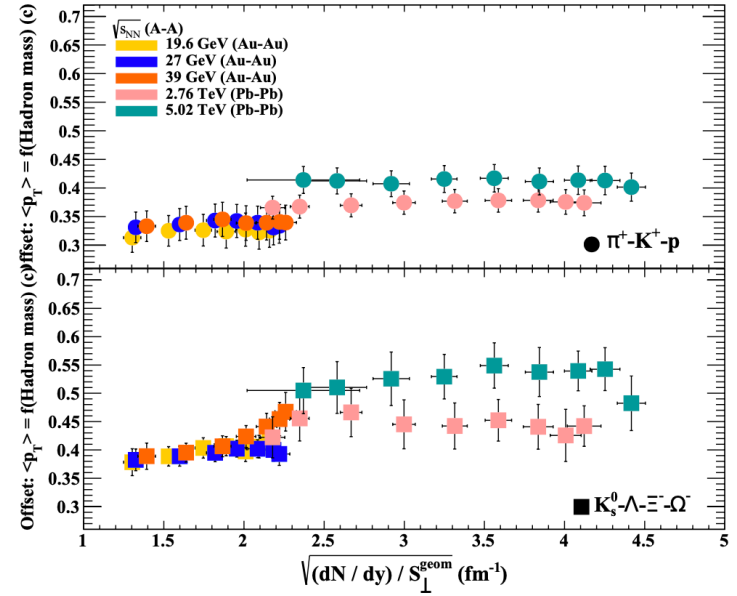
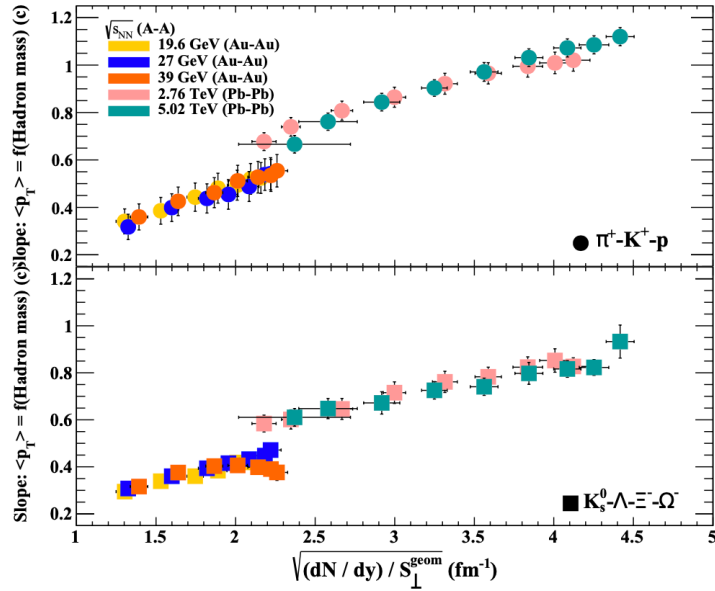
ALICE Collaboration, Phys. Rev. C 88(2013)044910

STAR Collaboration, Phys. Rev. C 96(2017)044904

ALICE Collaboration, Nucl. Phys. A 967(2017)421

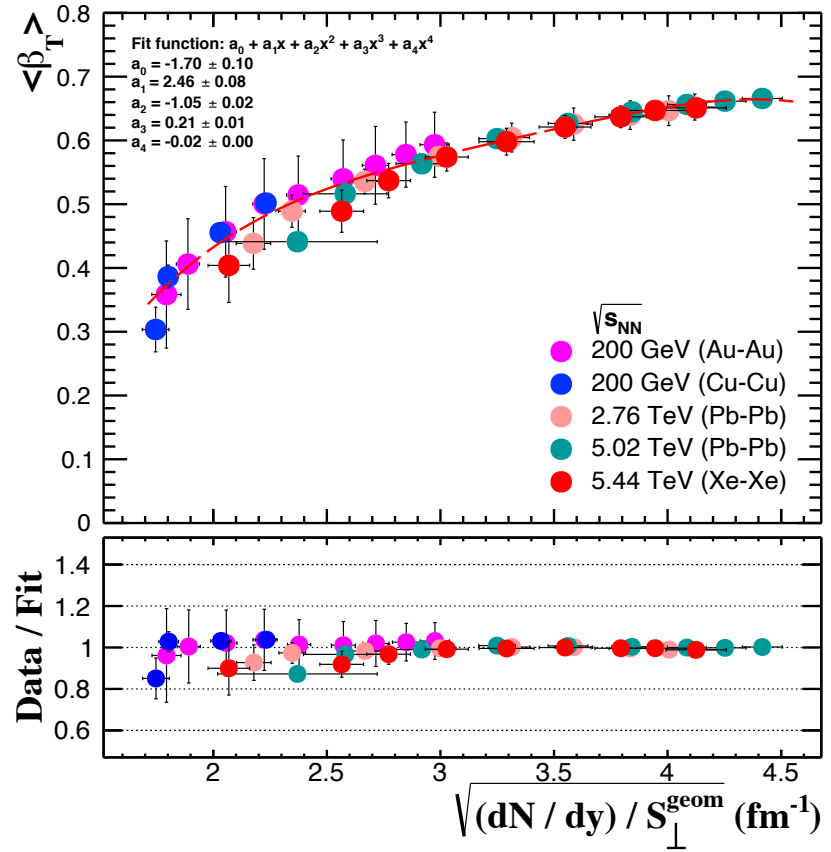
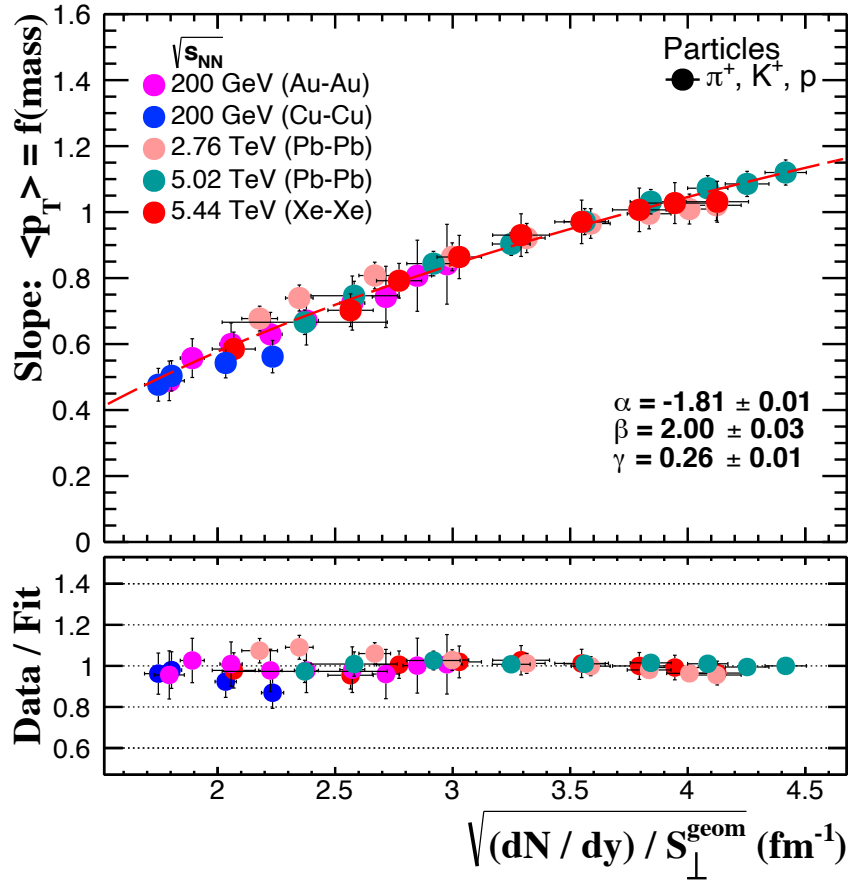
M. Petrovici and A. Pop, EuNPC 2022

$[(dN/dy)/S_{\perp}]^{1/2}$ scaling



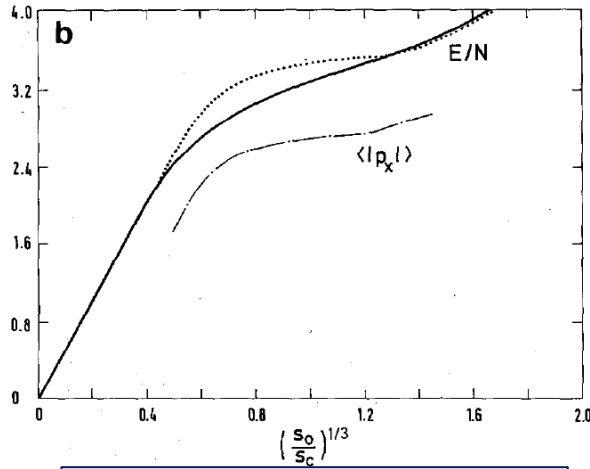
M. Petrovici et al., Phys. Rev. C 98(2018)024904
M. Petrovici and A. Pop, EuNPC 2022
A. Pop and M. Petrovici, arXiv:2402.19115[hep-ph]

$[(dN/dy)/S_{\perp}]^{1/2}$ scaling



M. Petrovici, A. Lindner and A. Pop, AIP Conf.Proc. 2076 (2019) 1, 040001

Energy and Entropy density



J.-P. Blaizot and J.-Y. Ollitrault,
Phys.Lett 191B(1987)21

RHIC BES energies:

$$\frac{dN}{dy} \simeq \frac{3}{2} \frac{dN^{(\pi^+ + \pi^-)}}{dy} + 2 \frac{dN^{(K^+ + K^-, p + \bar{p}, \Xi^- + \bar{\Xi}^+)}}{dy} + \frac{dN^{(\Lambda + \bar{\Lambda})}}{dy}$$

RHIC $\sqrt{s_{NN}}=62.4; 130$ and 200 GeV:

$$\frac{dN}{dy} \simeq \frac{3}{2} \frac{dN^{(\pi^+ + \pi^-)}}{dy} + 2 \frac{dN^{(K^+ + K^-, p + \bar{p}, \Xi^- + \bar{\Xi}^+)}}{dy} + \frac{dN^{(\Lambda + \bar{\Lambda}, \Omega^- + \bar{\Omega}^+)}}{dy}$$

LHC $\sqrt{s_{NN}}=2.76$ and 5.02 TeV:

$$\frac{dN}{dy} \simeq \frac{3}{2} \frac{dN^{(\pi^+ + \pi^-)}}{dy} + 2 \frac{dN^{(p + \bar{p}, \Xi^- + \bar{\Xi}^+)}}{dy} + 4 \frac{dN^{\Sigma^+}}{dy} + \frac{dN^{(K^+ + K^-, K_S^0 + \bar{K}_S^0, \Lambda + \bar{\Lambda}, \Omega^- + \bar{\Omega}^+)}}{dy}$$

for AGS and RHIC energies:

$$\langle m_T \rangle = \sqrt{\langle p_T \rangle^2 + m^2} - m_N \quad \text{- for baryons}$$

$$\langle m_T \rangle = \sqrt{\langle p_T \rangle^2 + m^2} + m_N \quad \text{- for antibaryons}$$

$$\langle m_T \rangle = \sqrt{\langle p_T \rangle^2 + m^2} \quad \text{- for other particles}$$

$$\epsilon = Ts - p$$

- qualitative temperature dependence

of entropy, pressure and energy density

- if p is small, at the transition the entropy density σ increases by the same factor as energy density ϵ

- dn/dy reflects the entropy, created early in the collision mainly through the interaction of the sea gluons of the colliding hadrons

- the entropy being conserved during expansion and hadronization

$$E/N \sim \epsilon/s = E_{fo}/S_{fo}$$

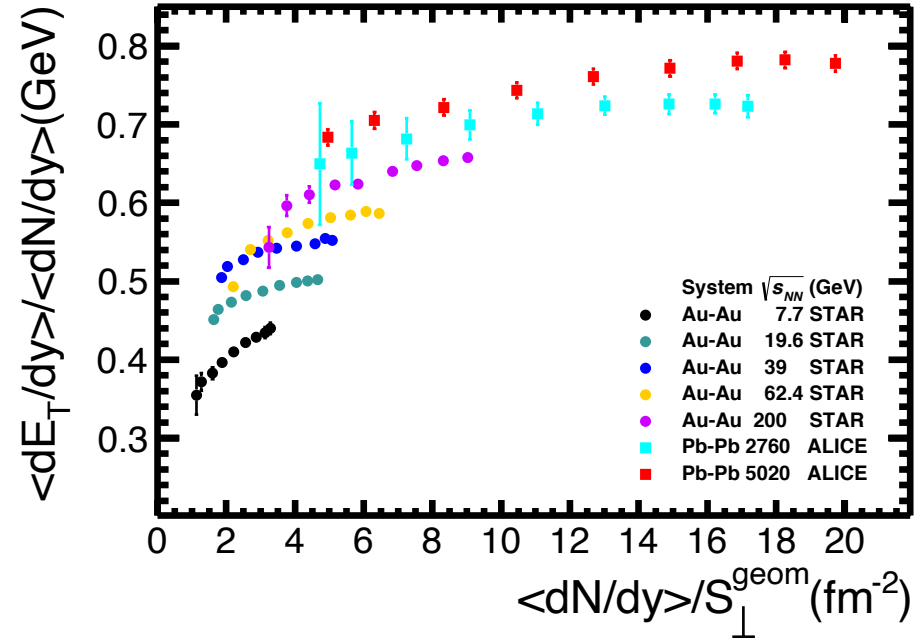
$$s(T_0) \sim (1/R_0^3)(dN/dy)$$

$$\frac{dE_T}{dy} \simeq \frac{3}{2} \left(\langle m_T \rangle \frac{dN}{dy} \right)^{(\pi^+ + \pi^-)} + 2 \left(\langle m_T \rangle \frac{dN}{dy} \right)^{(K^+ + K^-, p + \bar{p}, \Xi^- + \bar{\Xi}^+)} + \left(\langle m_T \rangle \frac{dN}{dy} \right)^{(\Lambda + \bar{\Lambda})}$$

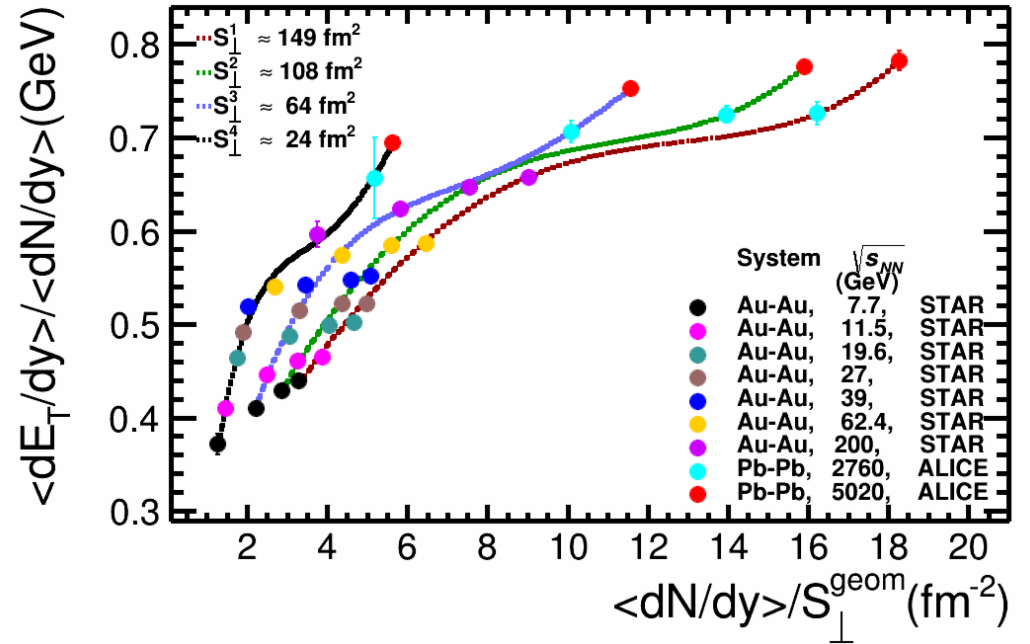
$$\frac{dE_T}{dy} \simeq \frac{3}{2} \left(\langle m_T \rangle \frac{dN}{dy} \right)^{(\pi^+ + \pi^-)} + 2 \left(\langle m_T \rangle \frac{dN}{dy} \right)^{(K^+ + K^-, p + \bar{p}, \Xi^- + \bar{\Xi}^+)} + \left(\langle m_T \rangle \frac{dN}{dy} \right)^{(\Lambda + \bar{\Lambda}, \Omega^- + \bar{\Omega}^+)}$$

$$\frac{dE_T}{dy} \simeq \frac{3}{2} \left(\langle m_T \rangle \frac{dN}{dy} \right)^{(\pi^+ + \pi^-)} + 2 \left(\langle m_T \rangle \frac{dN}{dy} \right)^{(p + \bar{p}, \Xi^- + \bar{\Xi}^+)} + 4 \left(\langle m_T \rangle \frac{dN}{dy} \right)^{\Sigma^+} + \left(\langle m_T \rangle \frac{dN}{dy} \right)^{(K^+ + K^-, K_S^0 + \bar{K}_S^0, \Lambda + \bar{\Lambda}, \Omega^- + \bar{\Omega}^+)}$$

$(dE_T/dy)/(dN/dy) - (dN/dy)/S_{\perp}$ correlation

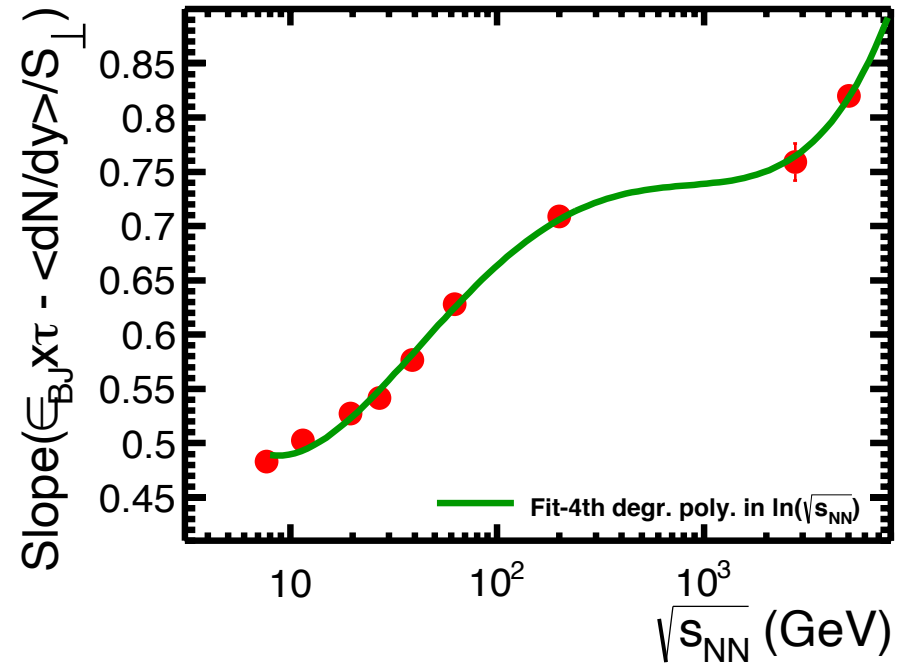
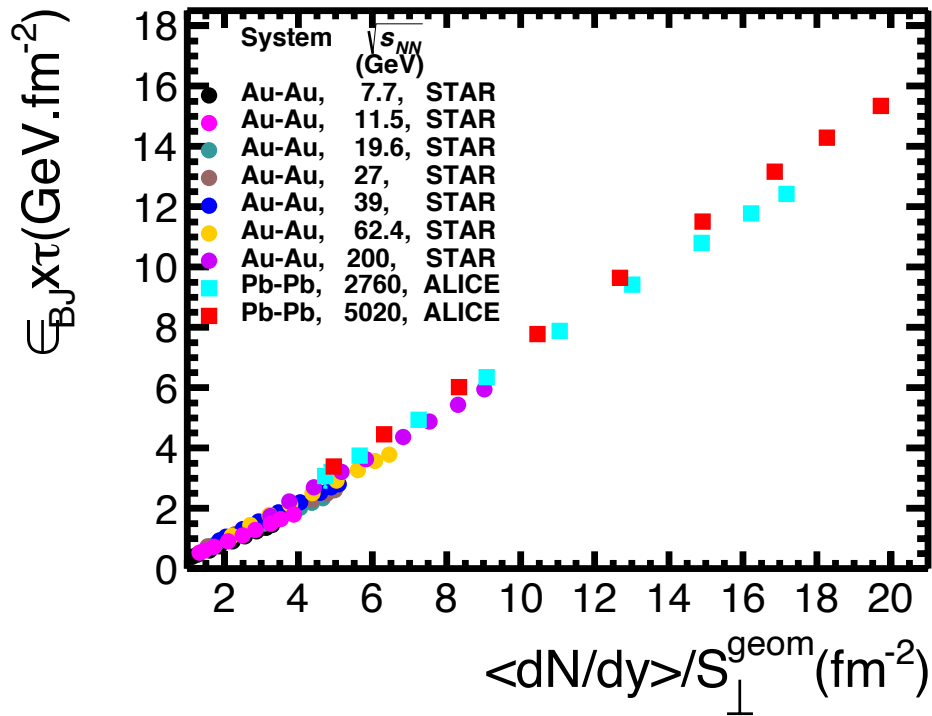


M.Petrovici and A.Pop, Phys.Rev. C107(2023)034913



$\epsilon_{Bj} - (dN/dy)/S_{\perp}$ correlation for A-A - centrality dependence

$$\epsilon_{Bj} \cdot \tau = (dE_T/dy)/S_{\perp}$$

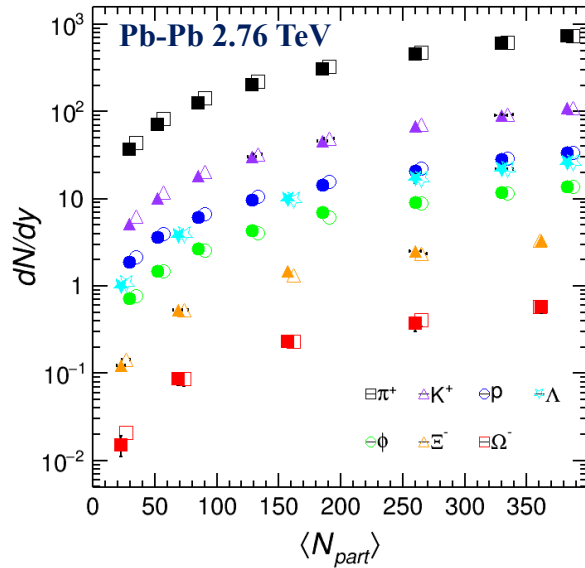


M.Petrovici and A.Pop, Phys.Rev. C107(2023)034913

Strangeness production - smoking gun of deconfinement

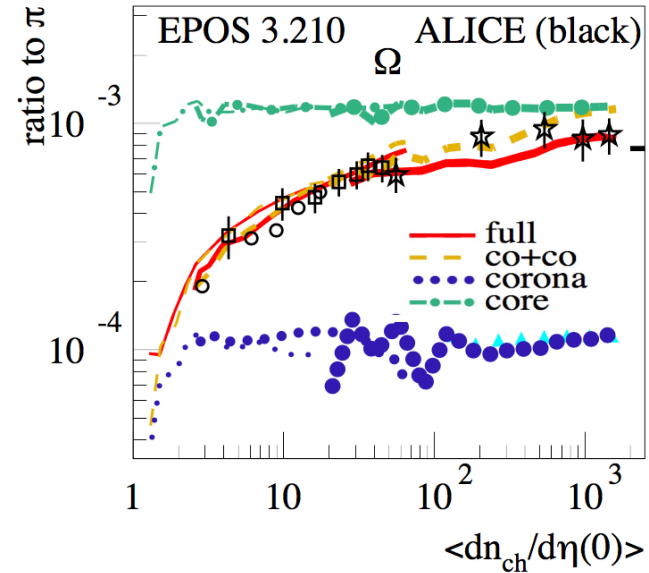
J.Rafelski and B.Muller, Phys.Rev.Lett. 48(1982)1066

$$\left(\frac{dN}{dy}\right)_i^{cen} = N_{part} [(1 - f_{core}) M_i^{ppMB} + f_{core} M_i^{core}] \quad (1)$$



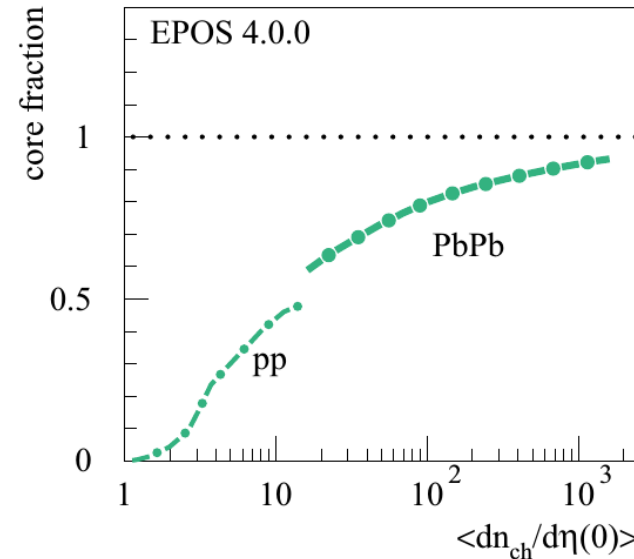
open symbols - Eq.1
full symbols - exp. points

M. Petrovici et al., Phys.Rev. C96(2017)014908



thin lines = pp (7TeV)
intermediate lines = pPb (5TeV)
thick lines = PbPb (2.76TeV)
circles = pp (7TeV)
squares = pPb (5TeV)
stars = PbPb (2.76TeV)

K. Werner, SQM 2017, July 10-15 2017, Utrecht

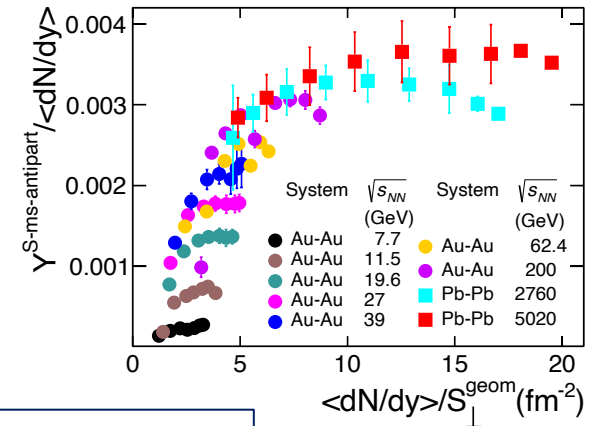
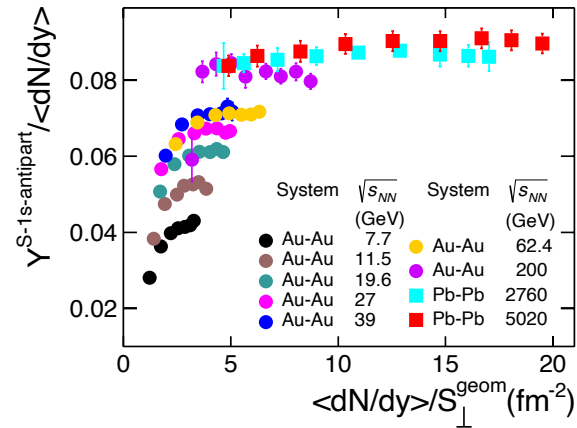
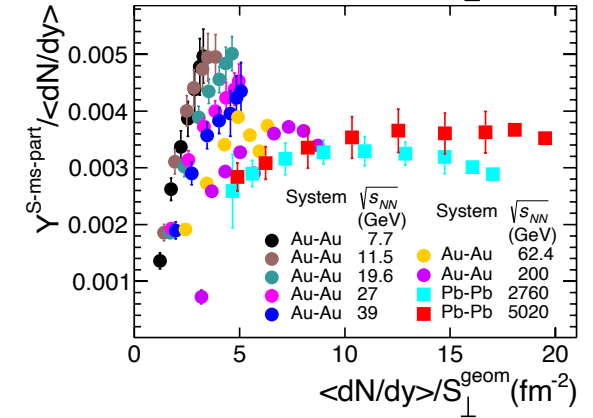
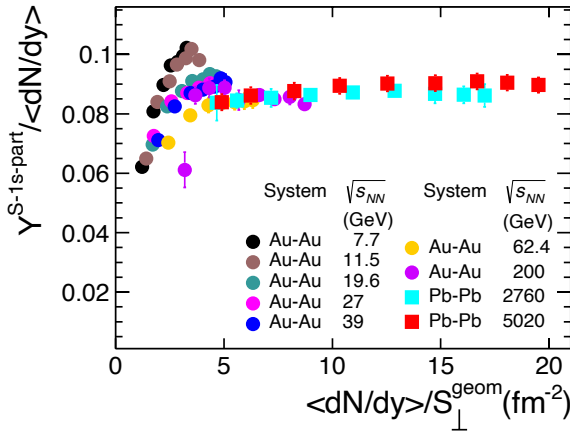
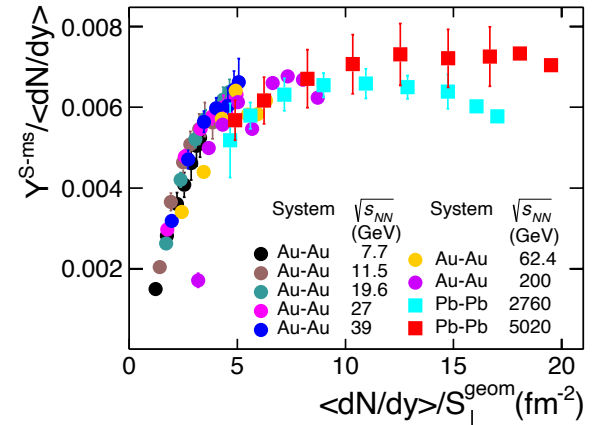
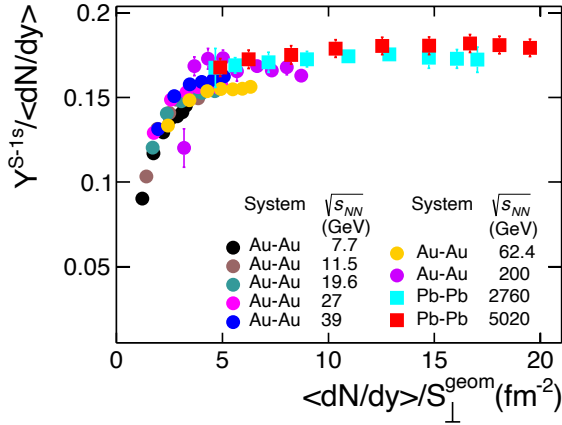


K. Werner, Phys.Rev. C109(2024)014910

$(dN/dy)_{\text{strange and multi strange}} / (dN/dy) - (dN/dy) / S_{\perp}$ correlation

$$Y^{1s} = \frac{dN^{1s}}{dy} = \frac{dN^{(K^++K^-)}}{dy} + 2\frac{dN^{K_s^0}}{dy} + \frac{dN^{(\Lambda+\bar{\Lambda})}}{dy} + 2\frac{dN^{(\Sigma^-+\bar{\Sigma}^+)}}{dy}$$

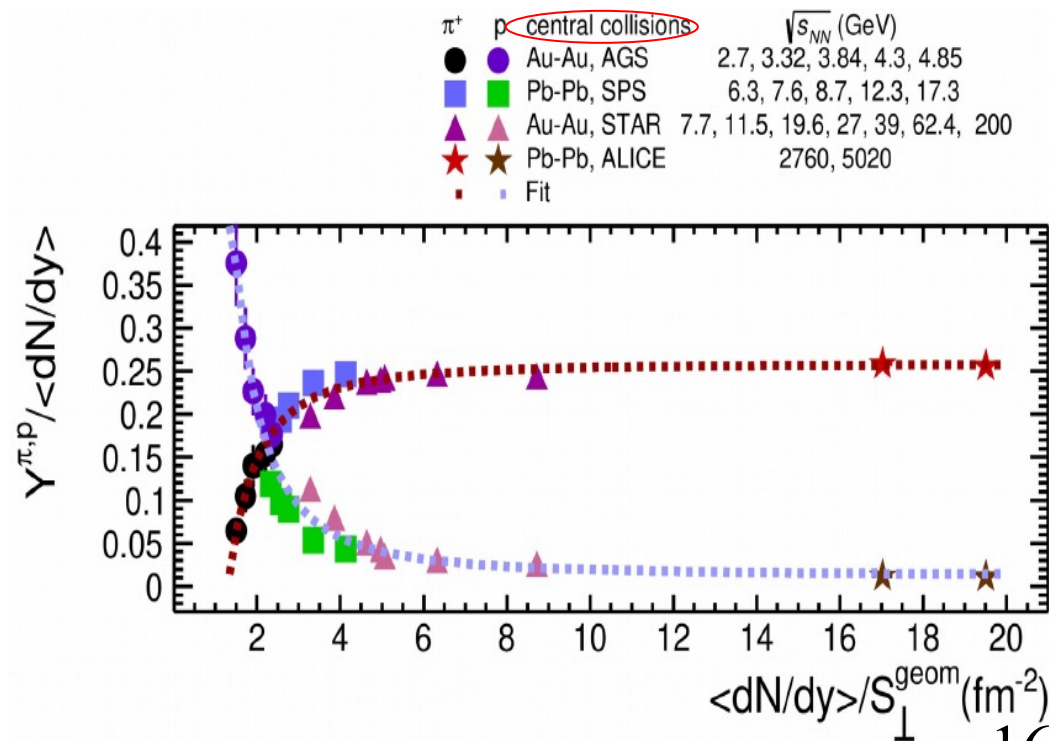
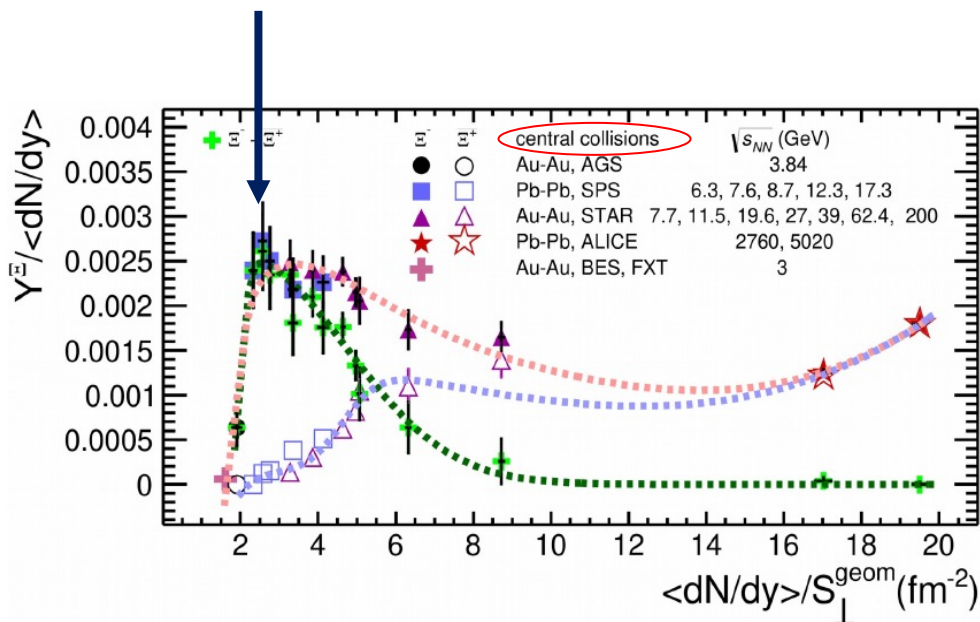
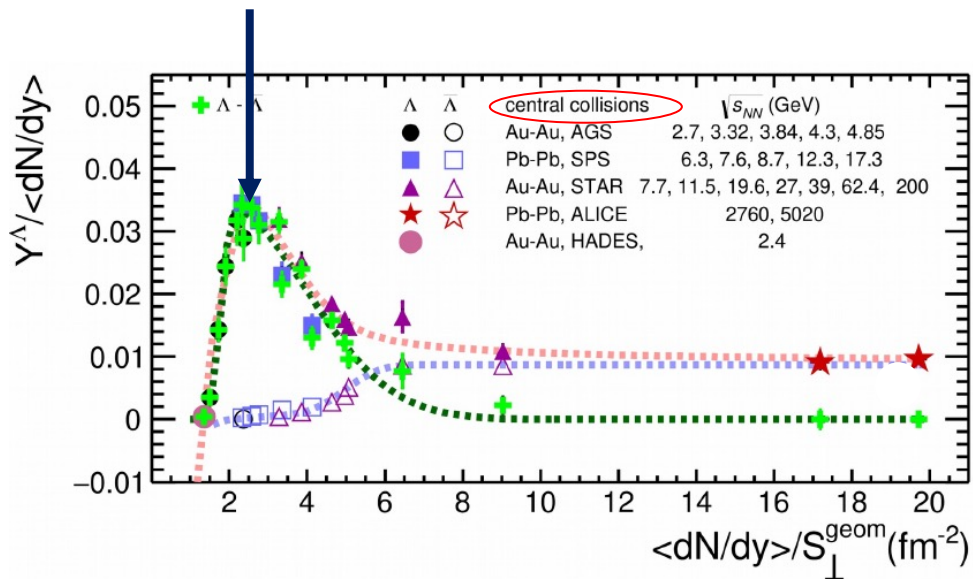
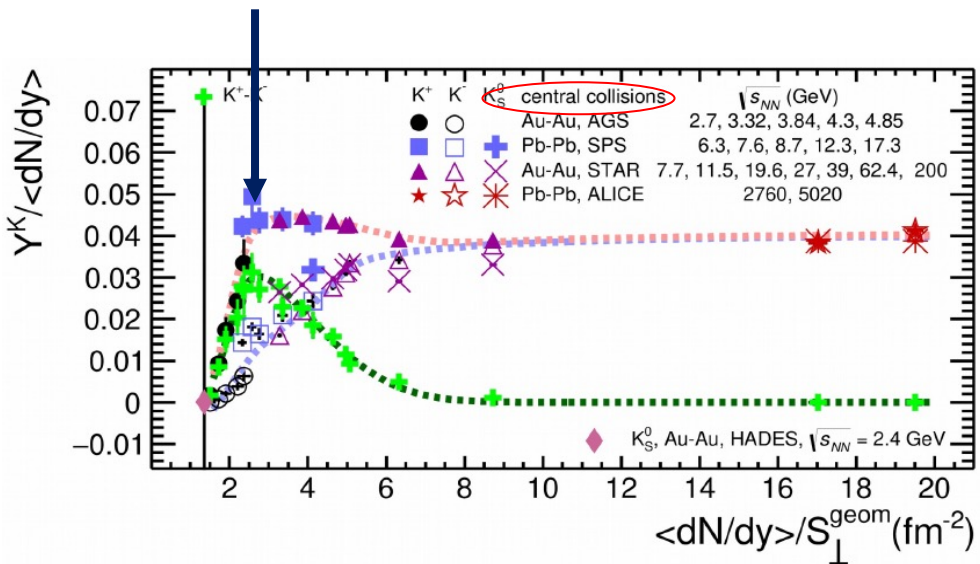
$$Y^{ms} = \frac{dN^{ms}}{dv} = \frac{dN^{(\Omega^-+\bar{\Omega}^+)}}{dv} + 2\frac{dN^{(\Xi^-+\bar{\Xi}^+)}}{dv}$$



M. Petrovici and A. Pop, EuNPC 2022
A. Pop and M. Petrovici, arXiv:2402.19115[hep-ph]

$(dN/dy)_{(strange\ and\ multi\ strange)} / (dN/dy) - (dN/dy) / S_{\perp}$ correlation

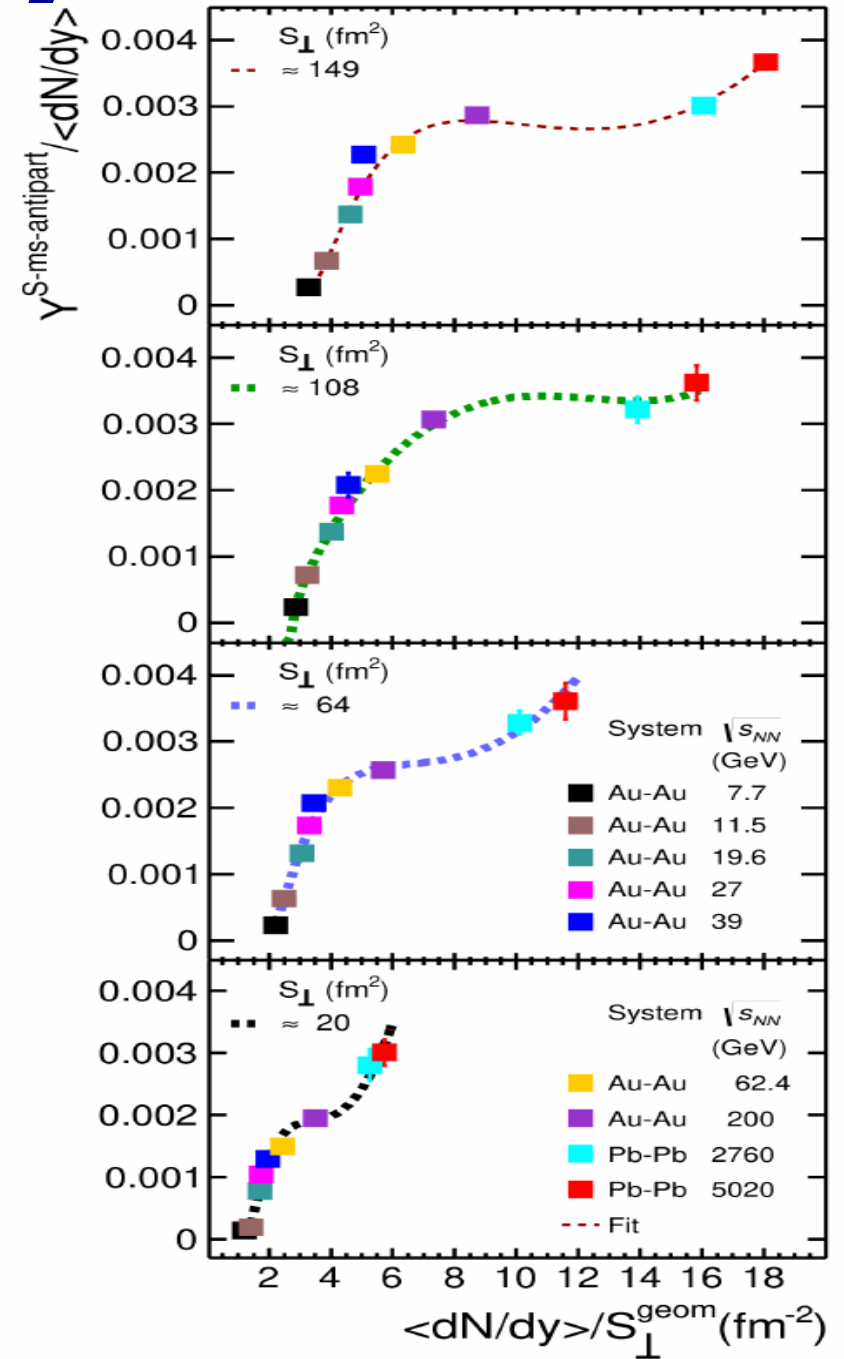
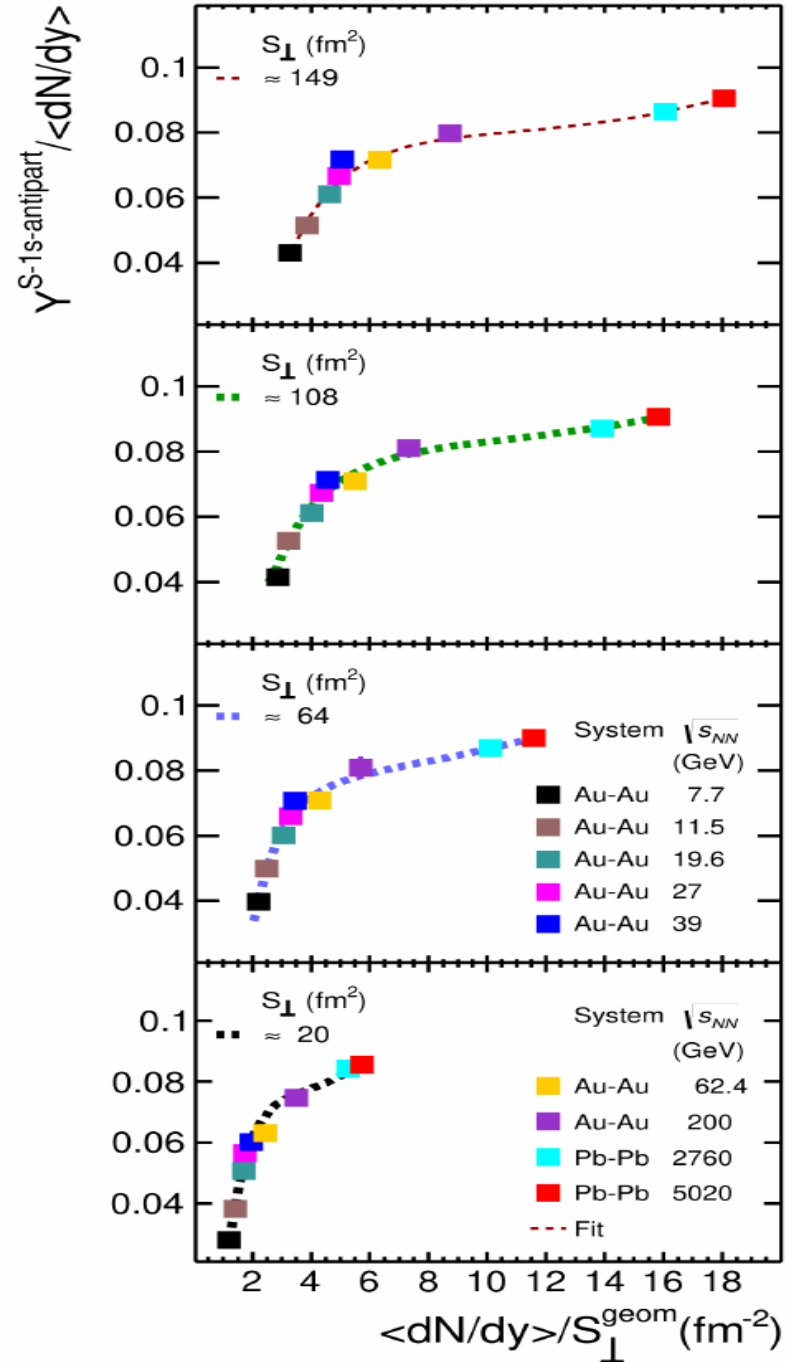
central collisions



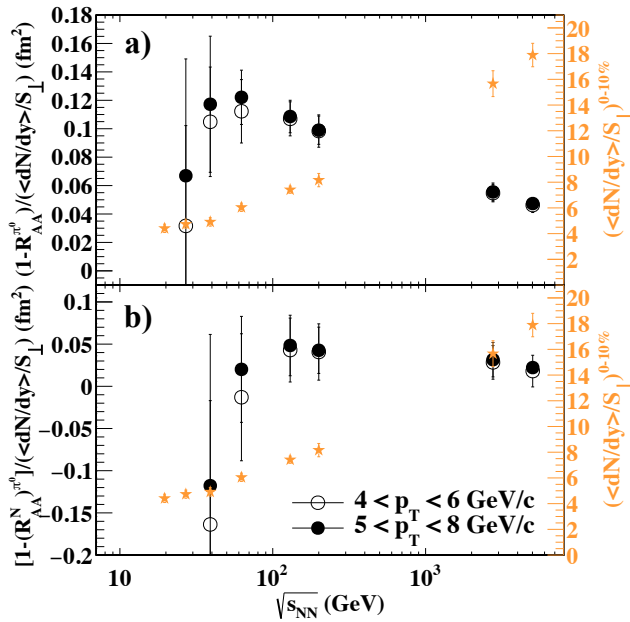
M. Petrovici and A. Pop, EuNPC 2022
 A. Pop and M. Petrovici, arXiv:2402.19115[hep-ph]

$(dN/dy)_{\text{strange and multi strange antihadron}} / (dN/dy) - (dN/dy) / S_{\perp}$ correlation

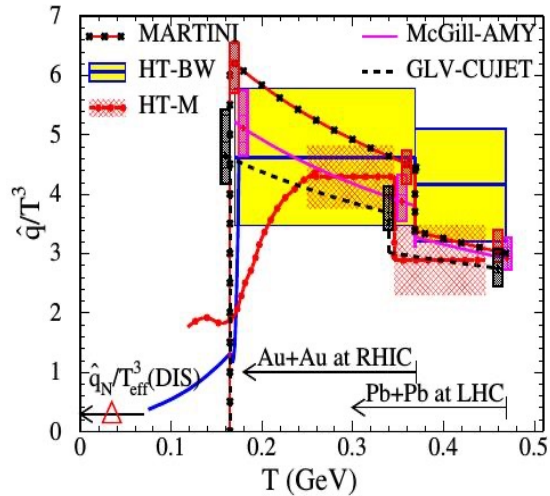
(different S_{\perp})



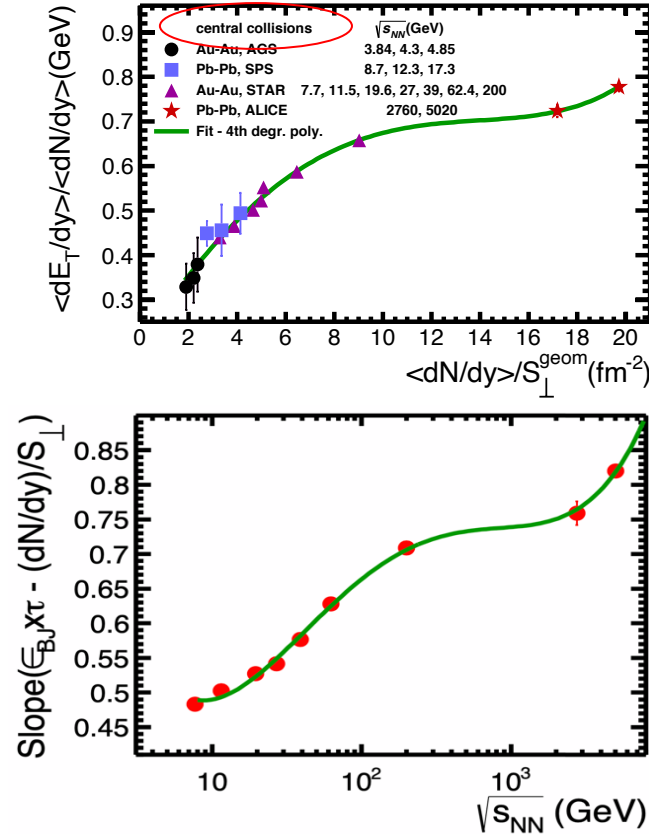
Do we see a new state of deconfined matter at LHC energies?



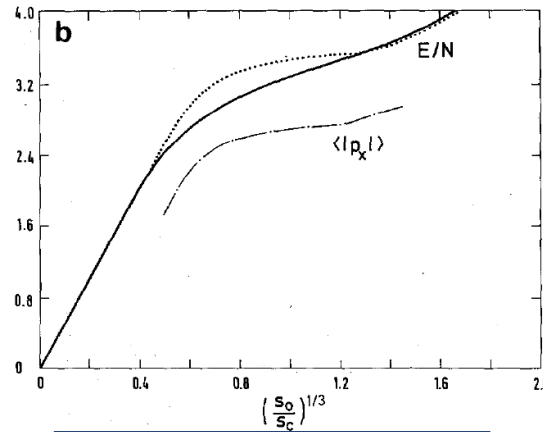
M. Petrovici et al., Phys. Rev. C103(2021)034903



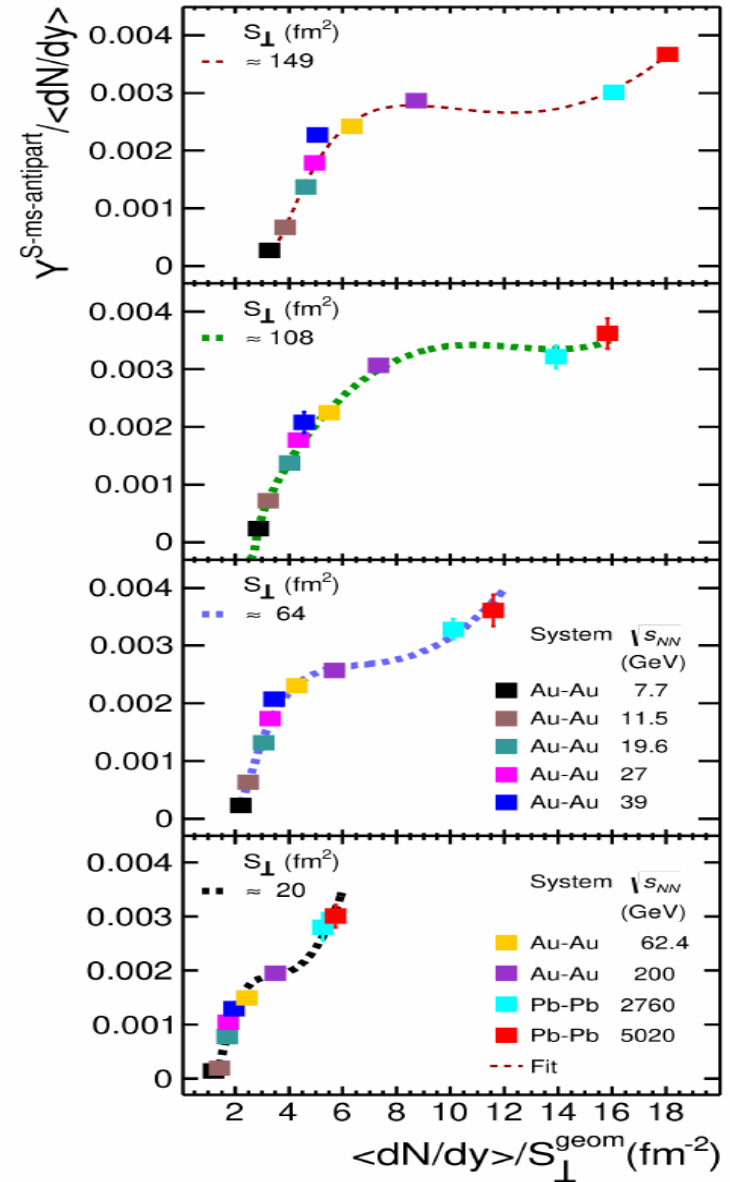
K.M. Burke et al., JET Collaboration, Phys. Rev. C90(2014)014909



M. Petrovici and A. Pop, Phys. Rev. C107(2023)0

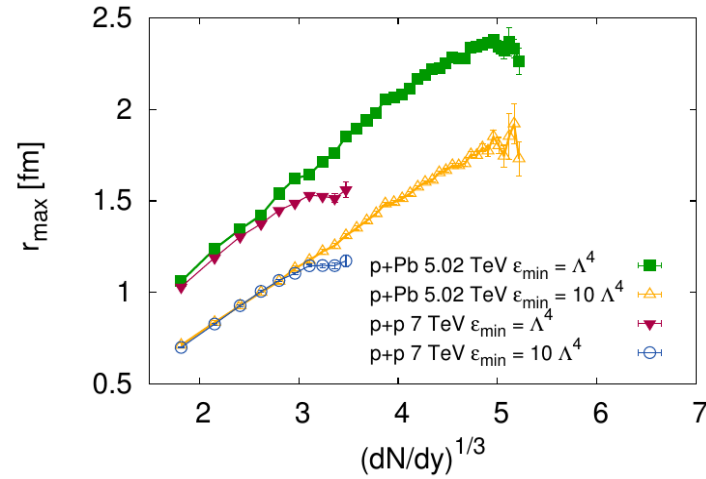


J.-P. Blaizot and J.-Y. Ollitrault, Phys. Lett 191B(1987)21



A. Pop and M. Petrovici, arXiv:2402.19115[hep-ph]

pp vs A-A @ LHC



$R_{pp} = 1 \text{ fm} \cdot f_{pp}$ - maximal radius for which the energy density of the Yang-Mill fields is larger than $\varepsilon = \alpha \Lambda_{QCD}^4$ ($\alpha \in [1, 10]$)

$$S_{\perp}^{pp} = \pi R_{pp}^2$$

$$\alpha=1 \quad f_{pp} = \begin{cases} 0.387 + 0.0335x + 0.274x^2 - 0.0542x^3 & \text{if } x < 3.4 \\ 1.538 & \text{if } x \geq 3.4 \end{cases}$$

$$x = (dN_g/dy)^{1/3}$$

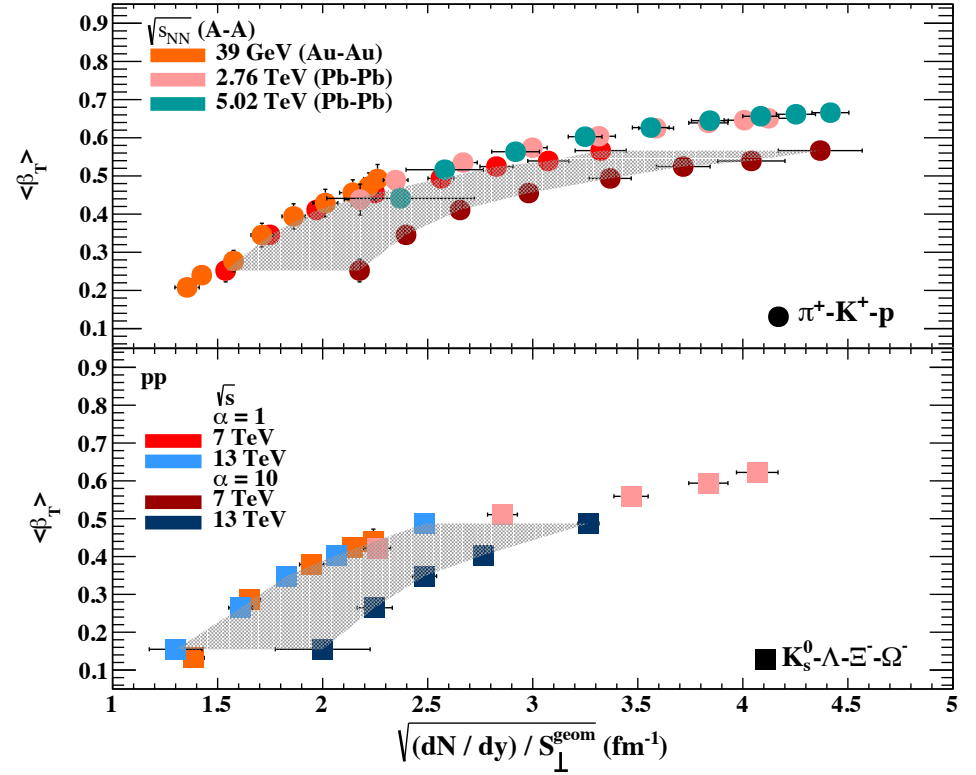
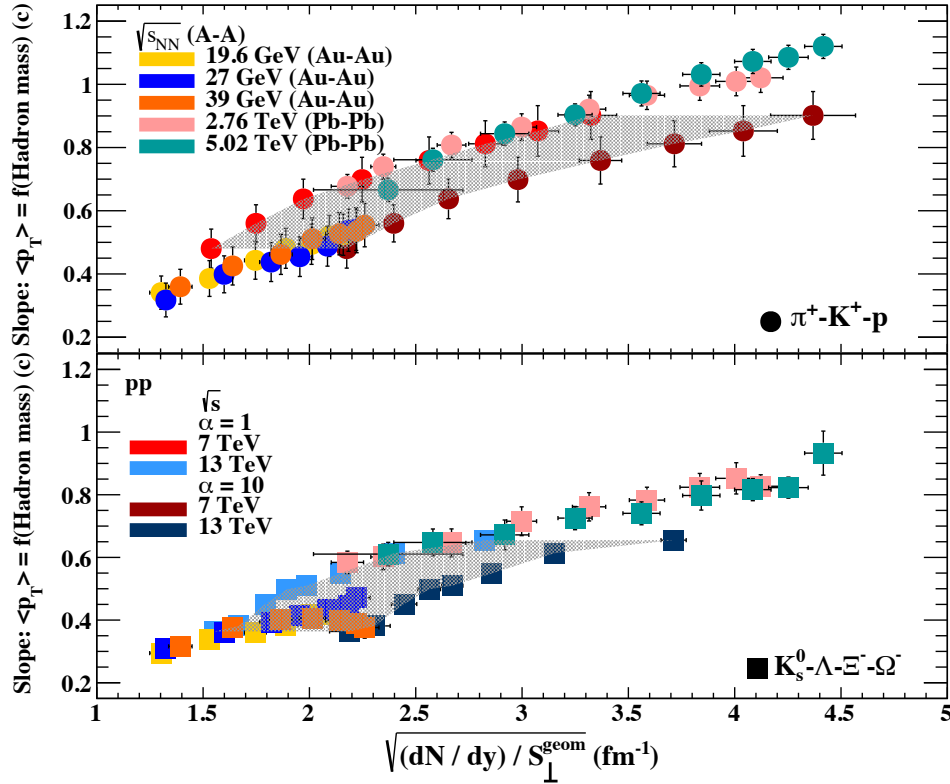
$$dN_g/dy \approx dN/dy$$

A. Bzdak et al., Phys.Rev. C87(2013)064906

McLarren, M. Praszalowicz and B. Schenke, Phys.Rev. C87(2013)064906

A-A vs. pp @ LHC

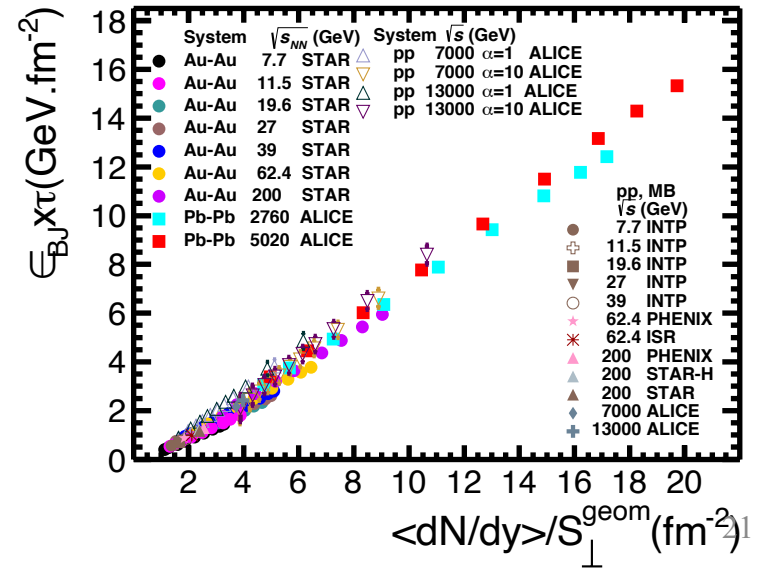
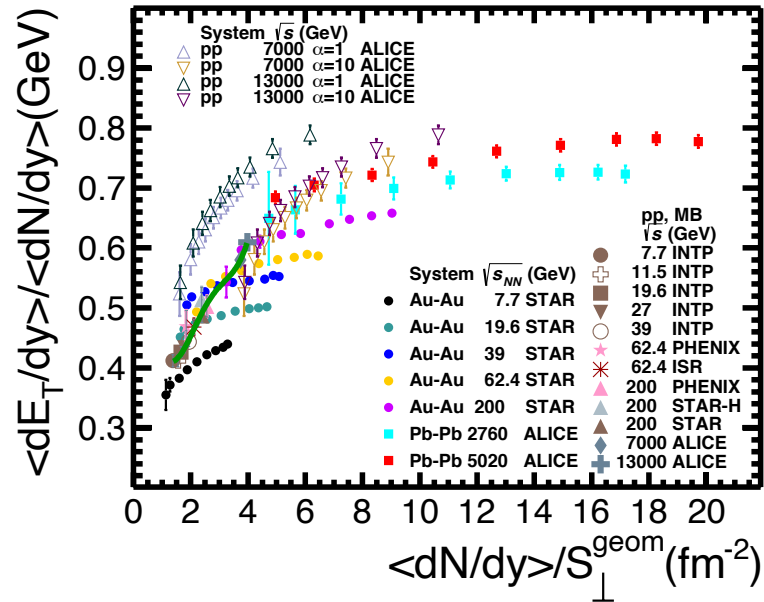
π, K, p



A. Lindner et al., Proceedings of Science (PoS) 380(2021)197
(PANIC2021), <https://pos.sissa.it/380/197/>.

A-A vs. pp @ LHC

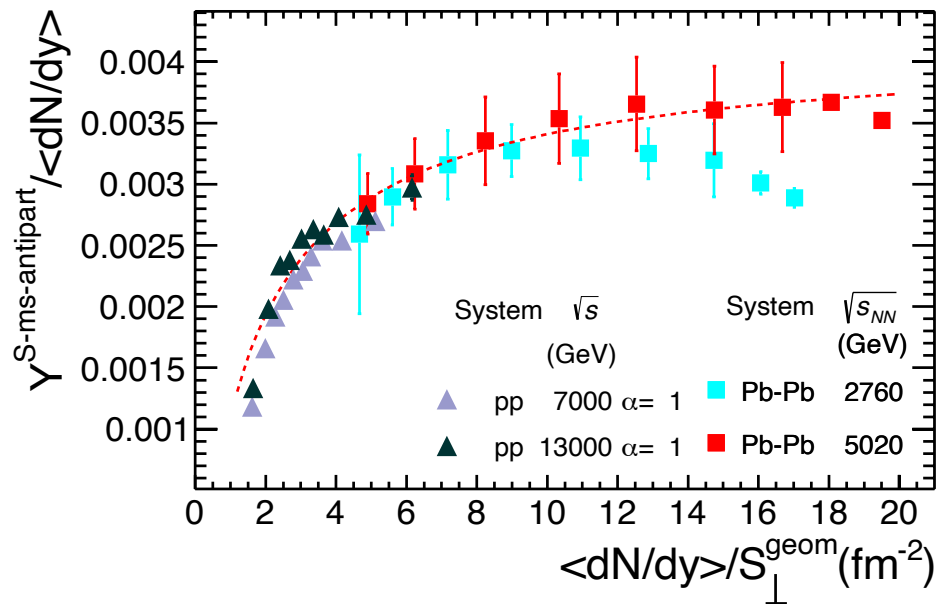
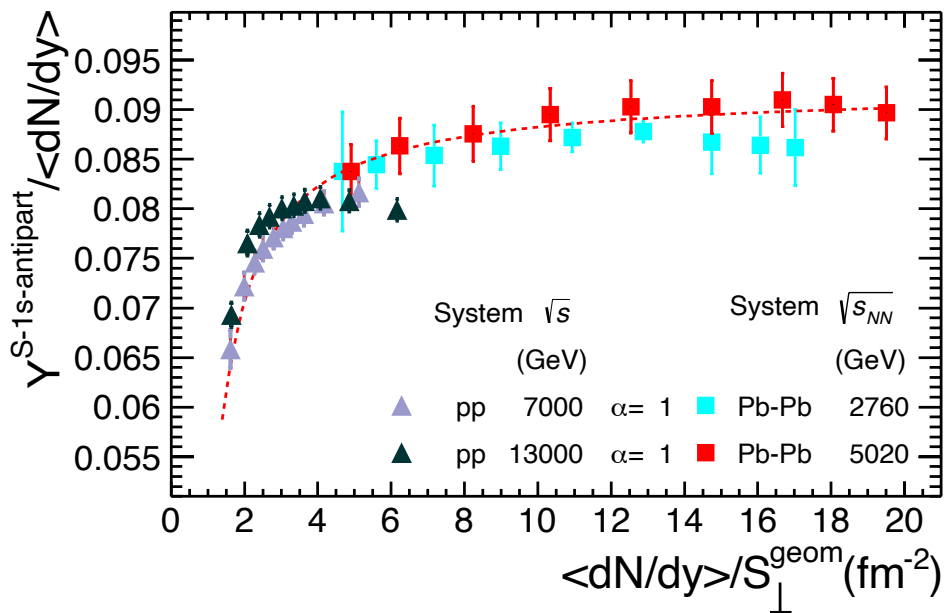
$$(dE_T/dy)/(dN/dy) - (dN/dy)/S_{\perp} \text{ and } \epsilon_{Bj} - (dN/dy)/S_{\perp}$$



M. Petrovici and A. Pop, Phys.Rev. C107(2023)034913

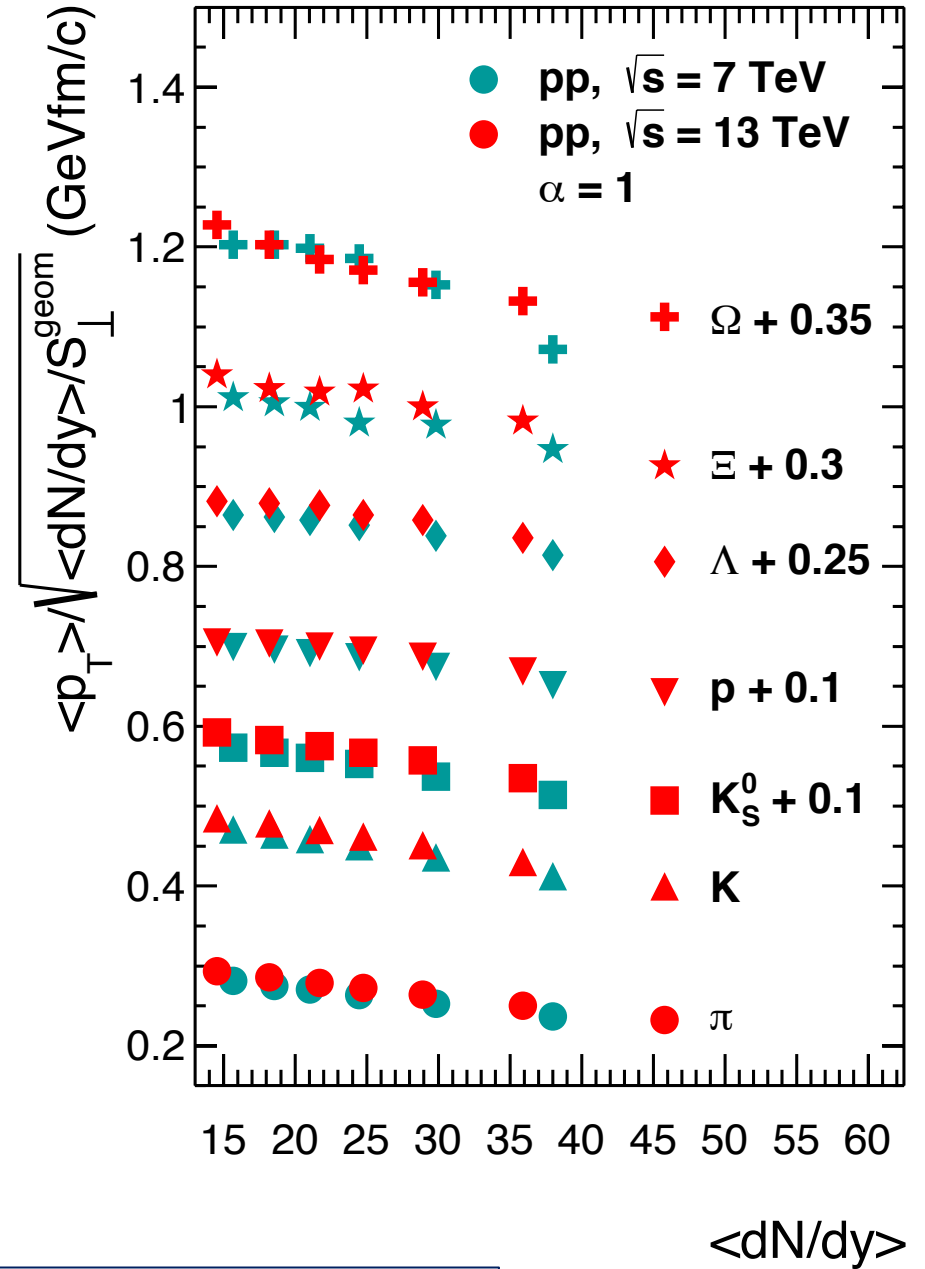
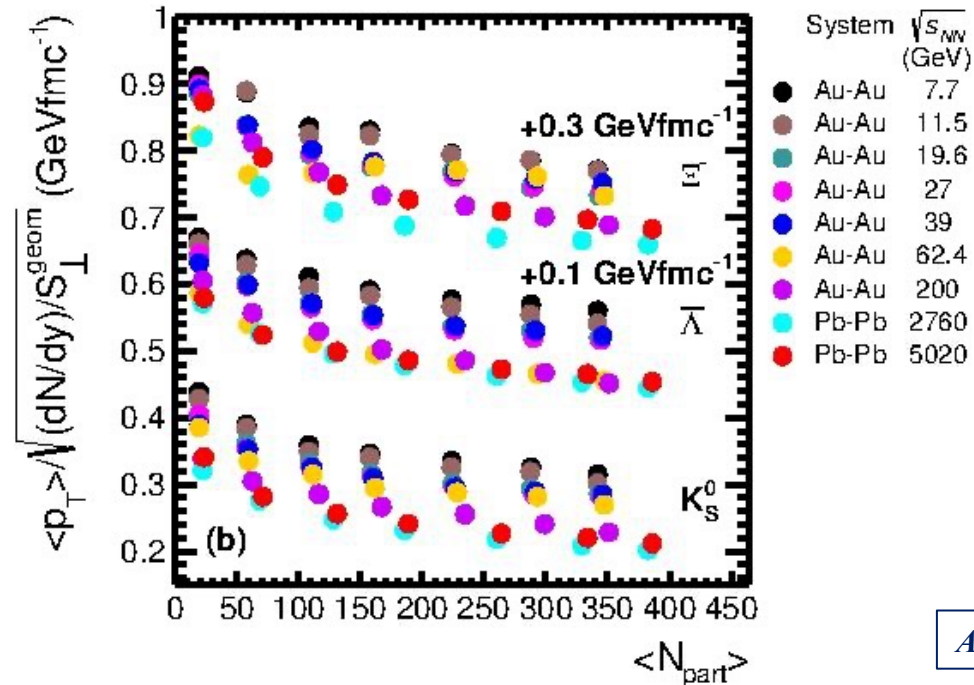
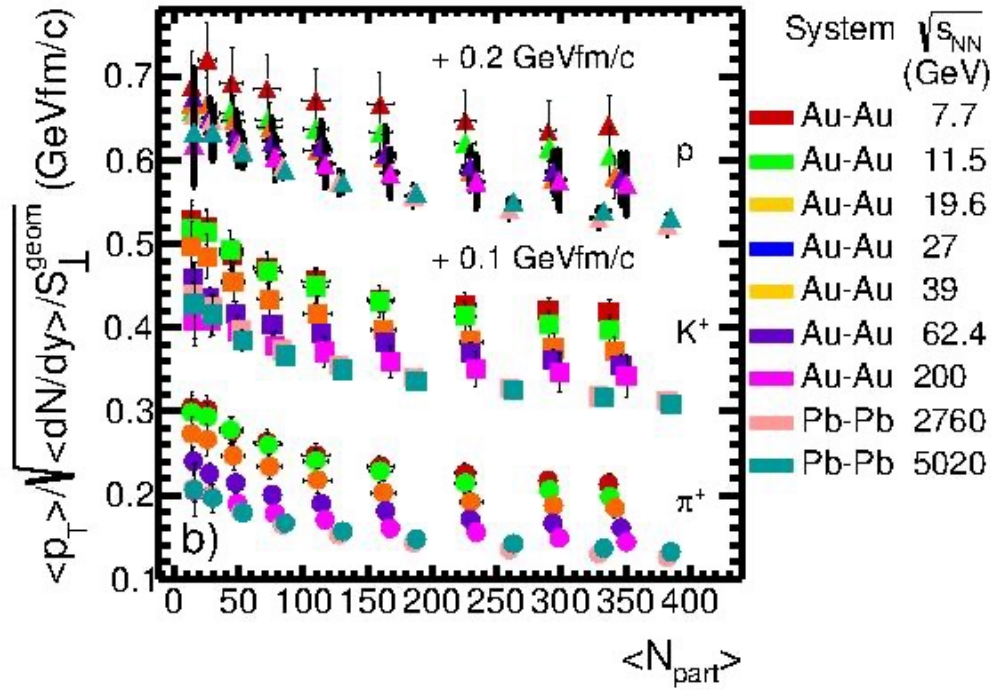
A-A vs. pp @ LHC

$$(dN/dy)^{\text{(strange and multi strange)}} / (dN/dy) - (dN/dy) / S_{\perp}$$



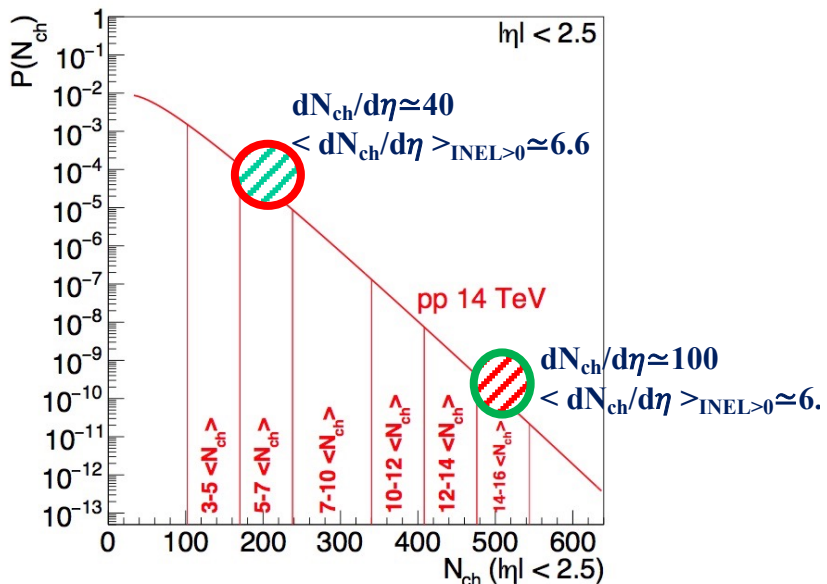
M. Petrovici and A. Pop, EuNPC 2022
A. Pop and M. Petrovici, arXiv:2402.19115[hep-ph]

$$\langle p_T \rangle / [(dN/dy) / S_{\perp}^{geom}]$$

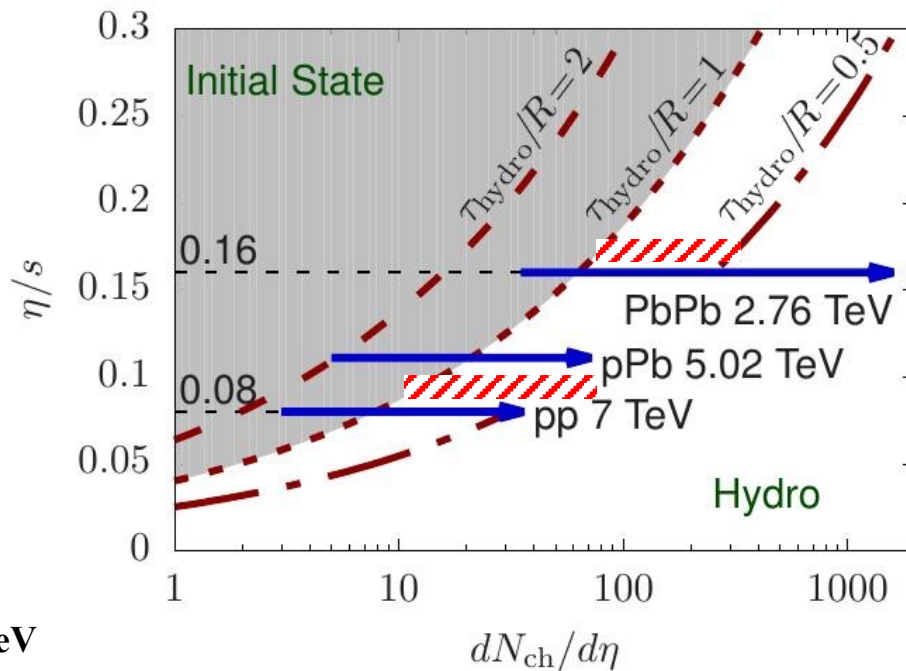


What's next ?

ALICE Coll., arXiv:1812.06772

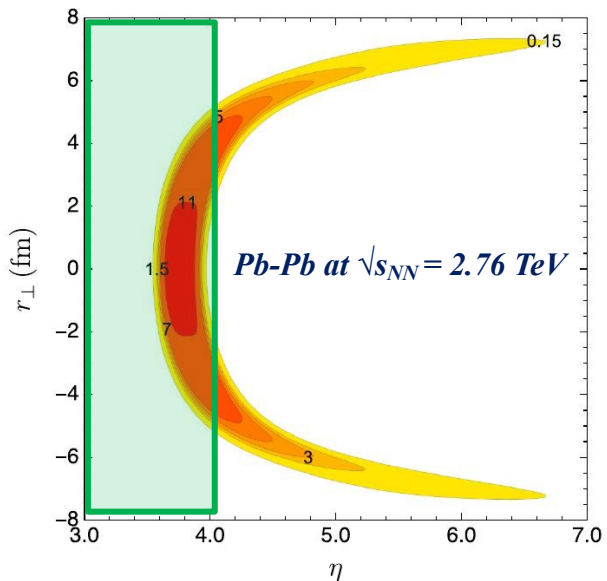
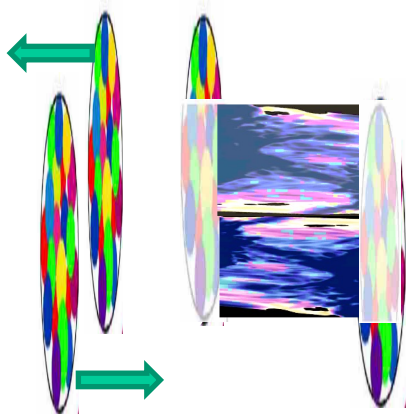


A.Kurkela et al., PoS(Confinement 2018)152



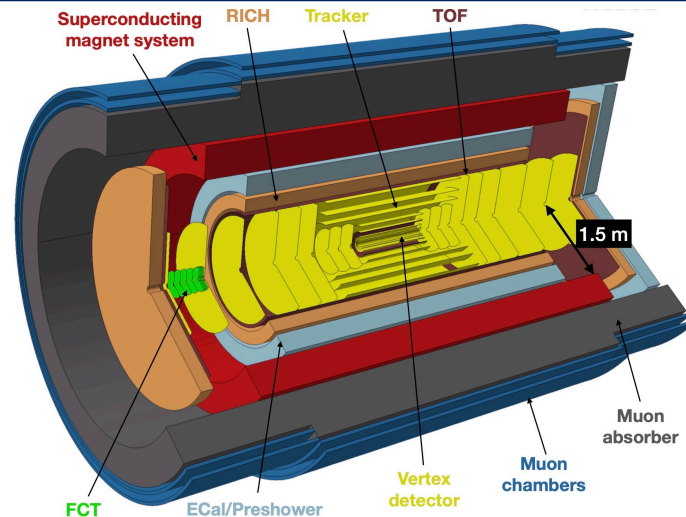
Pb-Pb @ 2-3 energies between 200 GeV and 2.76 TeV

PID @ large rapidity



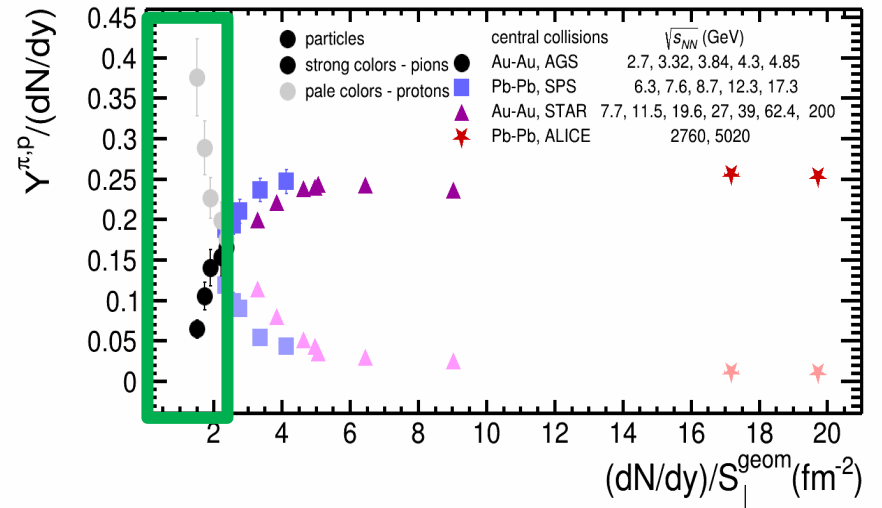
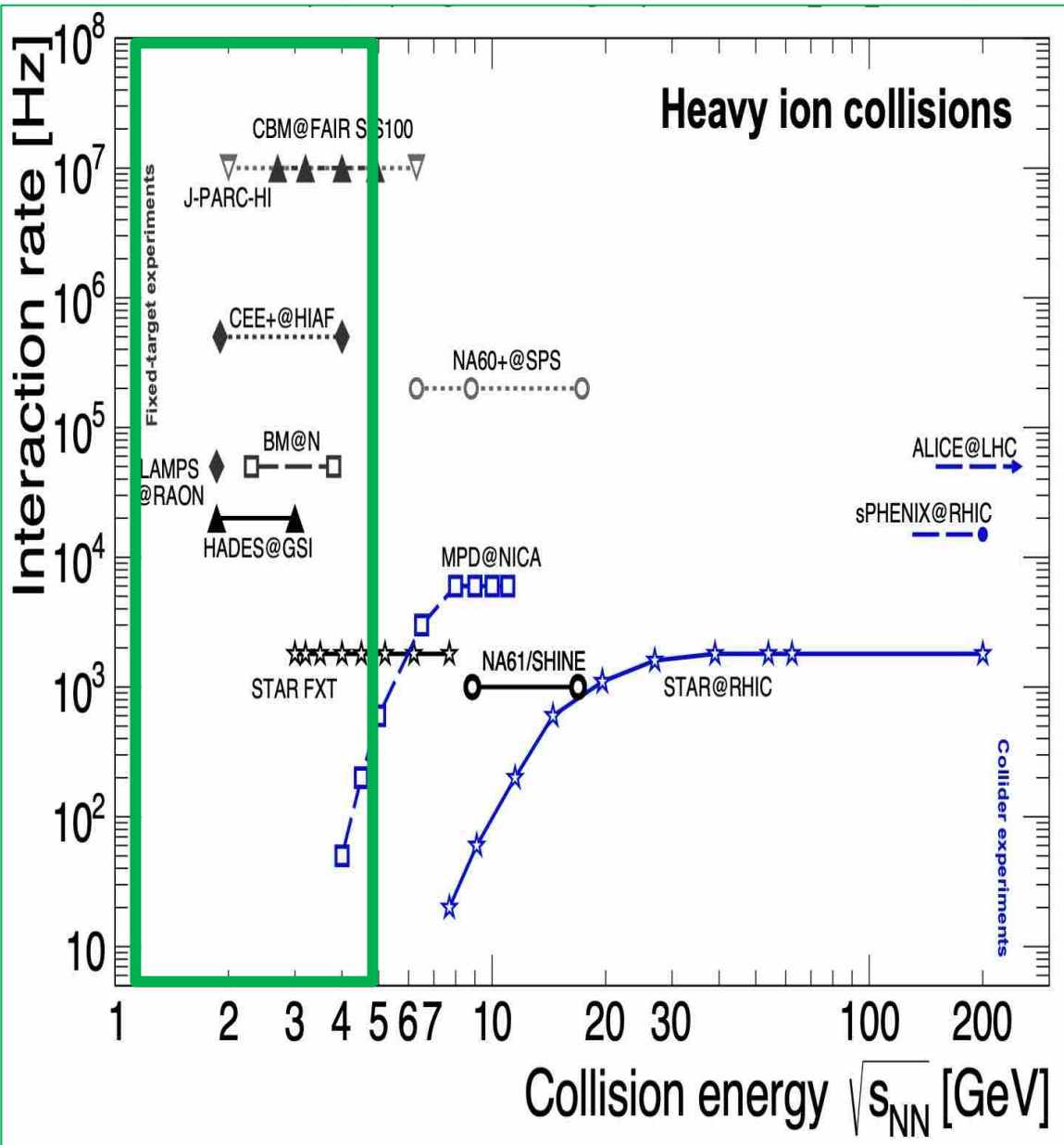
ALICE3

ALICE Collaboration, arXiv:2211.02491v1 [physics.ins-det] 4 Nov 2022



M. Li and J.I. Kapusta, Phys.Rev. C99(2019)014906

What can be studied at baryon densities - similar with those in neutron stars ?



Structure in the speed of sound

NS EoS \rightarrow EoS in HIC

N. Yao et al., arXiv:2311.18819 [nucl-th]

Input: Neutron Star EOS
Symmetry Energy Coefficients

For a given n_B , Convert ε_{NS} to ε_{HIC}

Subtract lepton contribution to P

Obtain $P = -\varepsilon + n_B \mu_B + n_Q \mu_Q$ and $\mu_B = \frac{d\varepsilon}{dn_B}$

Obtain C_s^2 via $C_s^2 = \frac{dP}{d\varepsilon}$

$$\frac{E_{ANM}}{N_B} = \frac{E_{SNM}}{N_B} + E_{\text{sym}} \delta^2 + \mathcal{O}(\delta^4)$$

$$\delta \equiv (n_n - n_p)/(n_n + n_p) = 1 - 2Y_{Q,\text{QCD}}$$

$\delta = 0$ symmetric nuclear matter

$\delta \simeq 1$ neutron star cores (=1 - pure neutron matter (PNM))

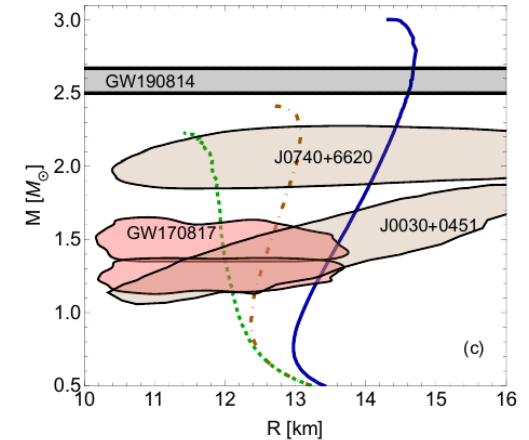
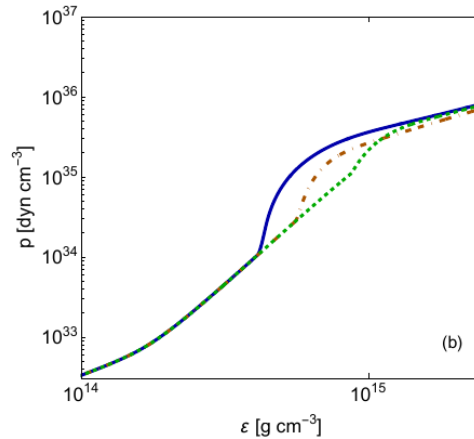
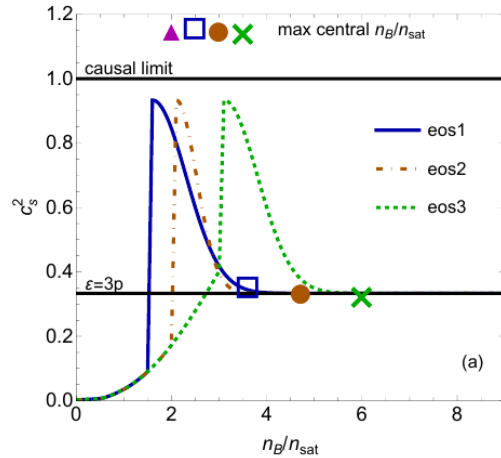
$$\frac{E_{PNM}}{N_B} = \frac{E_{SNM}}{N_B} + E_{\text{sym}}$$

$$\varepsilon_{\text{HIC,asym}} = \varepsilon_{\text{NS,QCD}} - 4n_B \left[E_{\text{sym,sat}} + \frac{L_{\text{sym,sat}}}{3} \left(\frac{n_B}{n_{\text{sat}}} - 1 \right) + \frac{K_{\text{sym,sat}}}{18} \left(\frac{n_B}{n_{\text{sat}}} - 1 \right)^2 + \frac{J_{\text{sym,sat}}}{162} \left(\frac{n_B}{n_{\text{sat}}} - 1 \right)^3 \right] \times \left[\left(Y_{Q,\text{QCD}}^{\text{const}} - Y_{Q,\text{QCD}} \right) + \left(Y_{Q,\text{QCD}}^2 - \left(Y_{Q,\text{QCD}}^{\text{const}} \right)^2 \right) \right]$$

$$\varepsilon + p = n_B \mu_B + n_Q \mu_Q$$

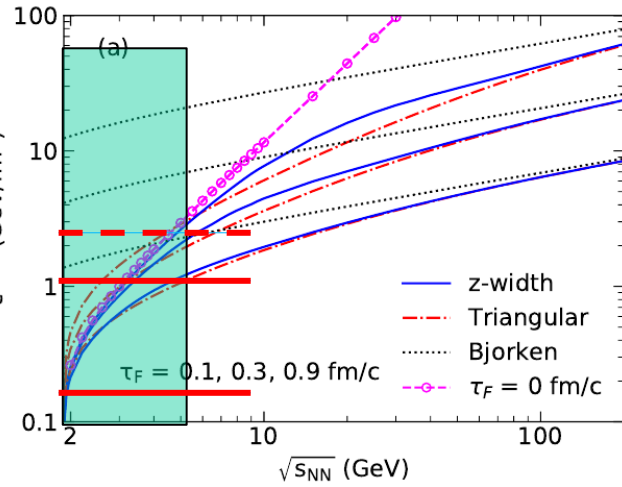
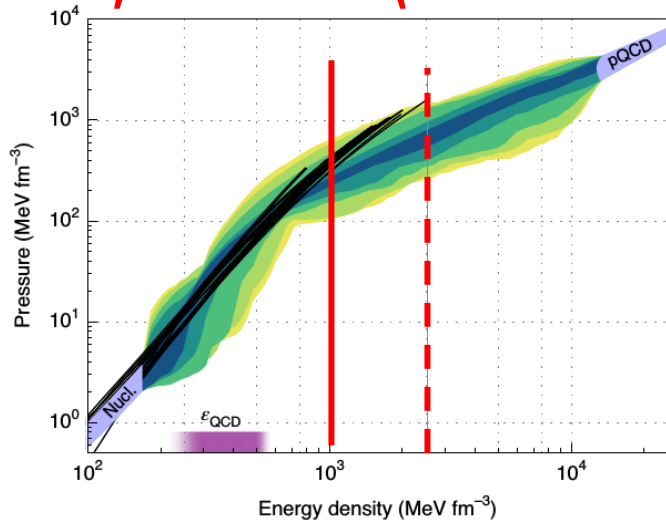
$$p = n_B^2 \frac{d(\varepsilon/n_B)}{dn_B}$$

$$c_s^2 = \left(\frac{dp}{d\varepsilon} \right)_{T=0}$$



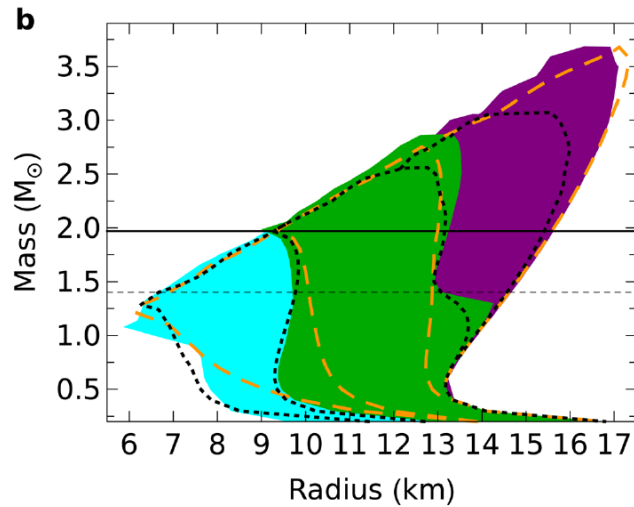
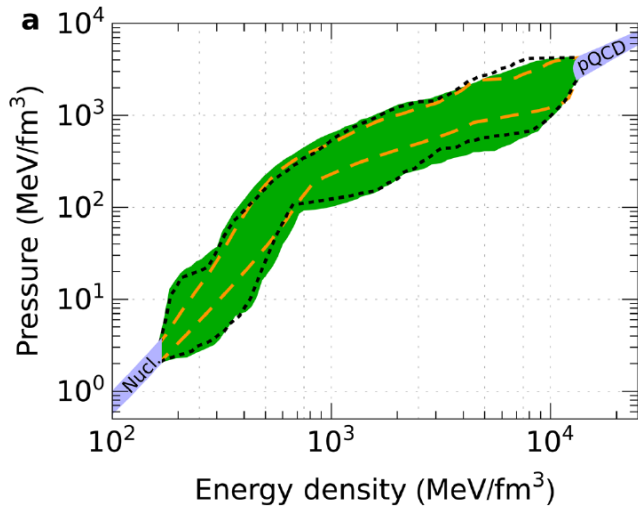
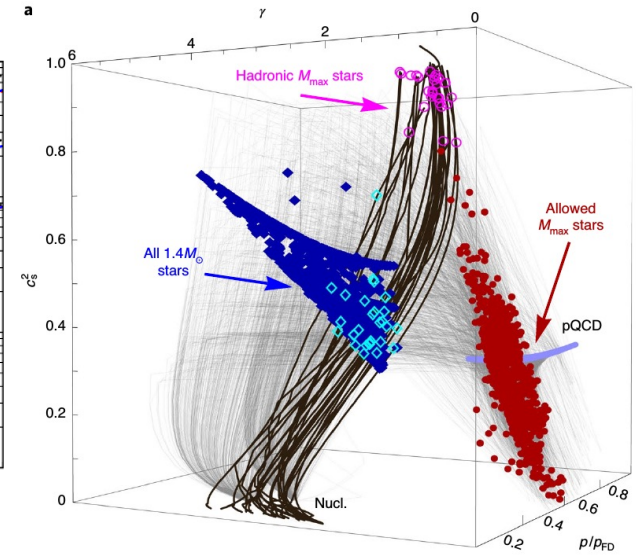
Evidence for quark-matter cores in massive neutron stars

CBM coverage



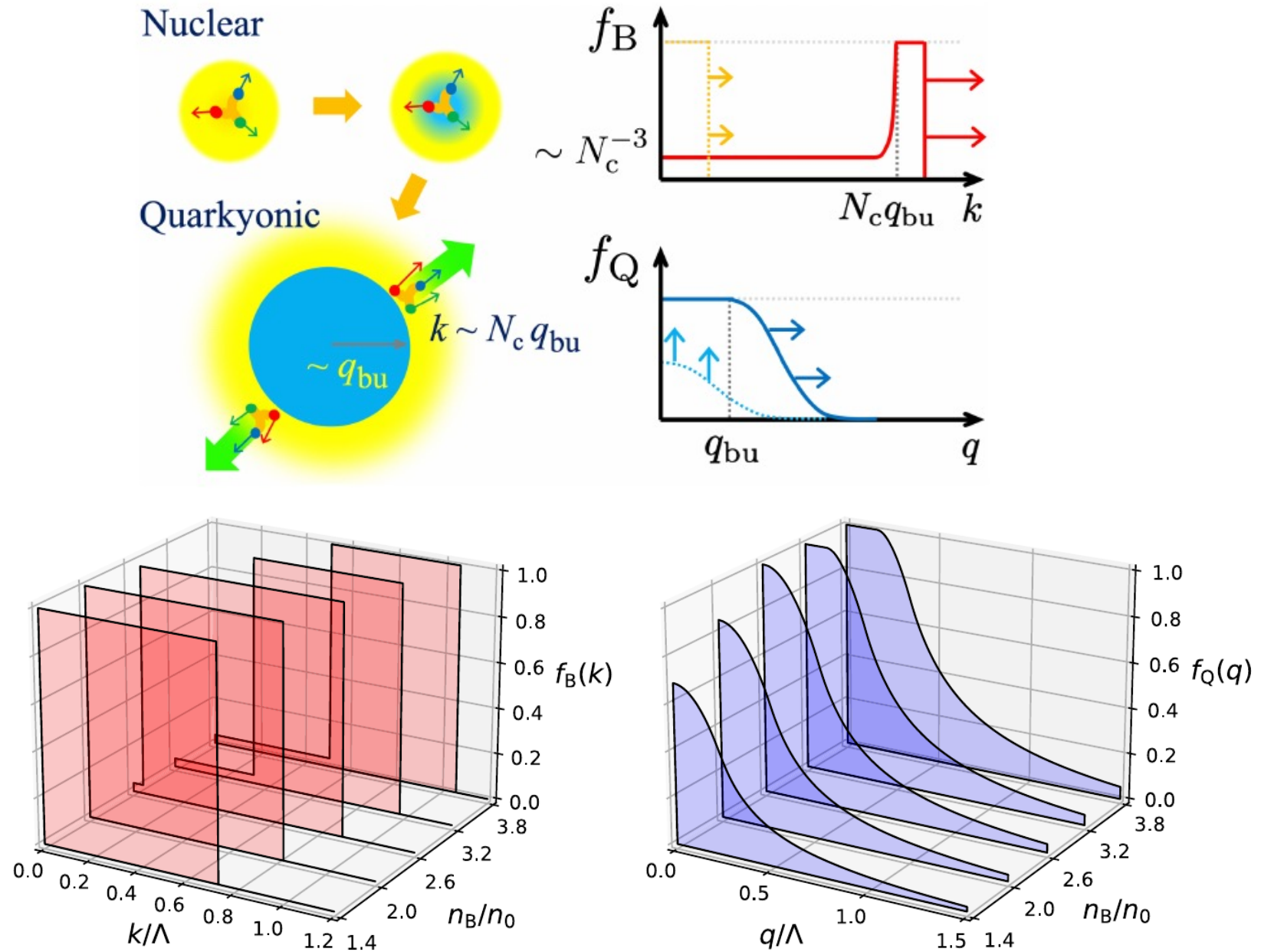
T. Mendenhall and Z.W. Lin, arXiv[nucl-th]2012.13825

$\gamma = c_s^2 \epsilon/P = \epsilon/P(dP/d\epsilon)$
differentiate between quark and hadronic matter
 $\gamma \leq 1.75$ - quark matter



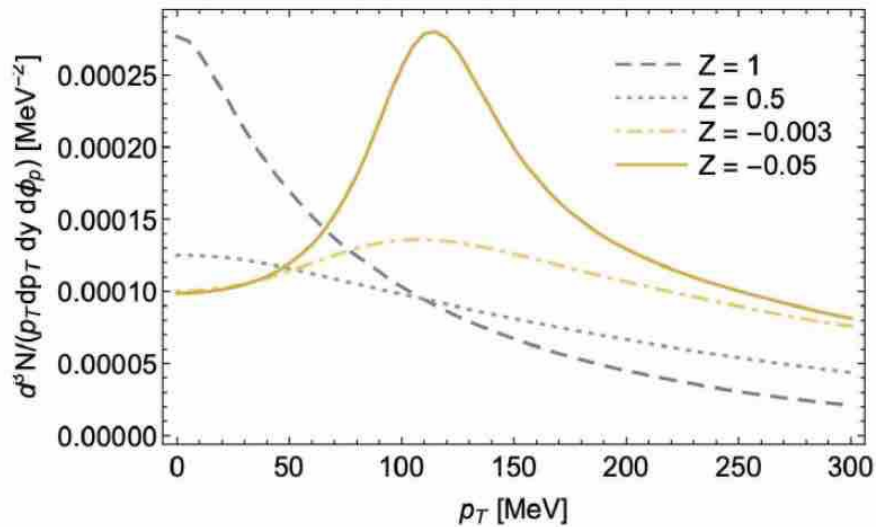
E. Annala et al, Nature Physics 16(2020)907

Quarkyonic Matter



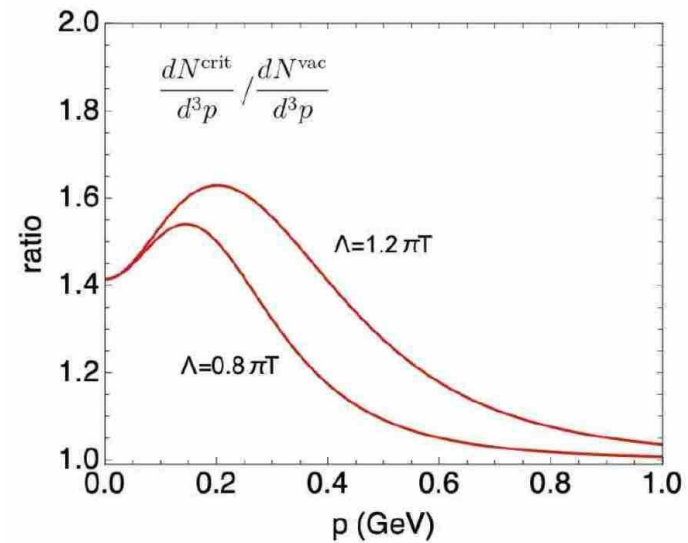
Large Baryon densities

*R.D. Pisarski and F. Rennecke,
arXiv:2103.06890[hep-ph]*



*Model studies suggest that regimes
with periodic spatial modulations
can occur high μ_B*

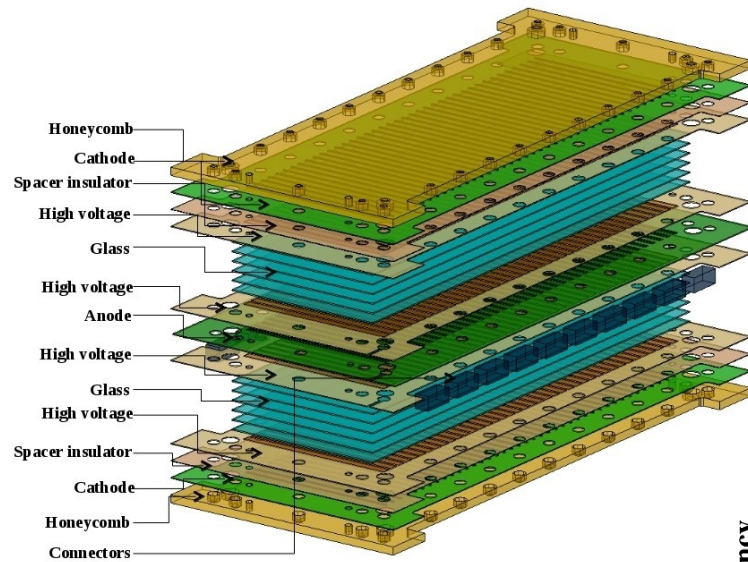
*E. Grossi et al.,
arXiv:2101.10847[nucl-th]*



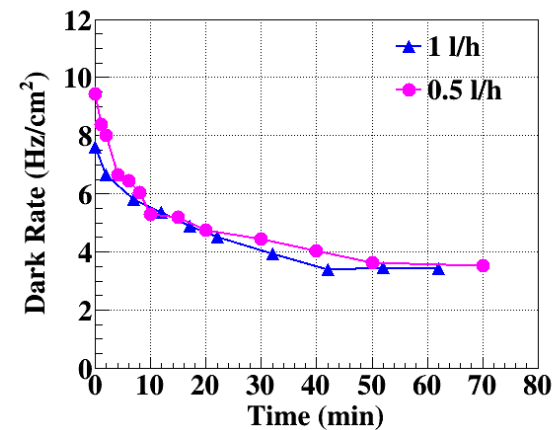
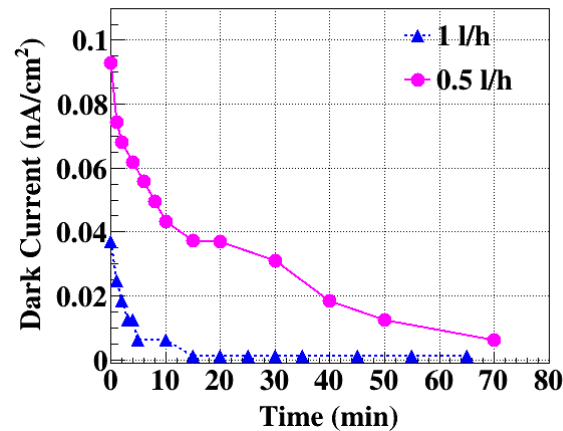
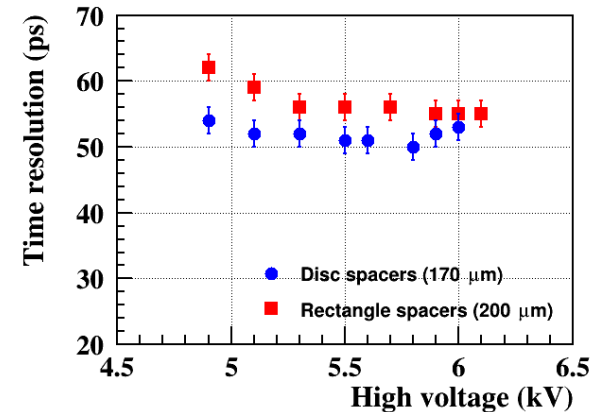
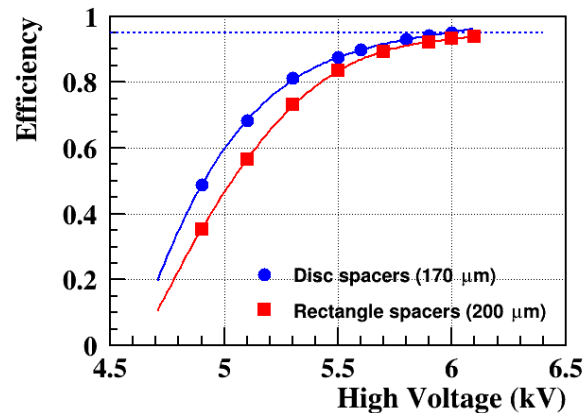
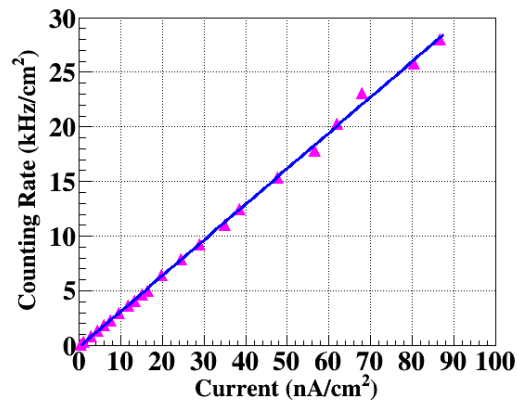
*The enhanced yield of soft pions
near the chiral critical point*

Multi-Strip Multi-Gap Resistive Plate Chambers - MSMGRPC

(Review and references: https://niham.nipne.ro/hpd_courier_no7_August_2024.pdf)

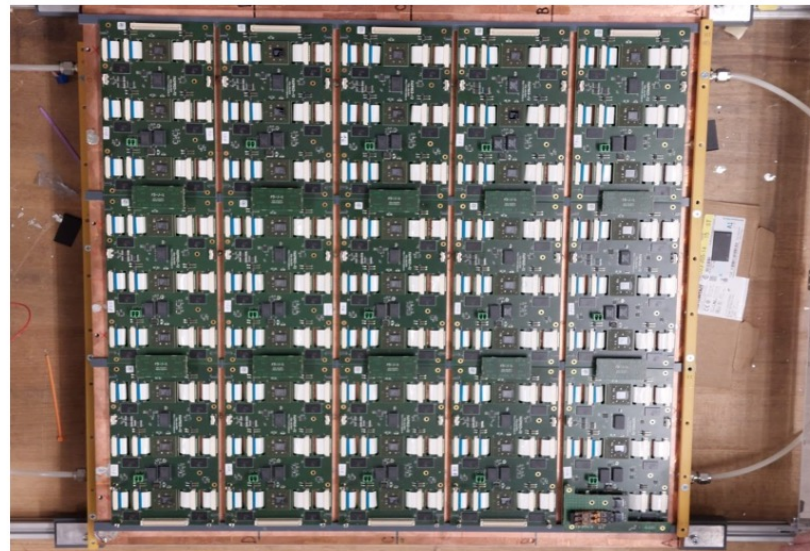
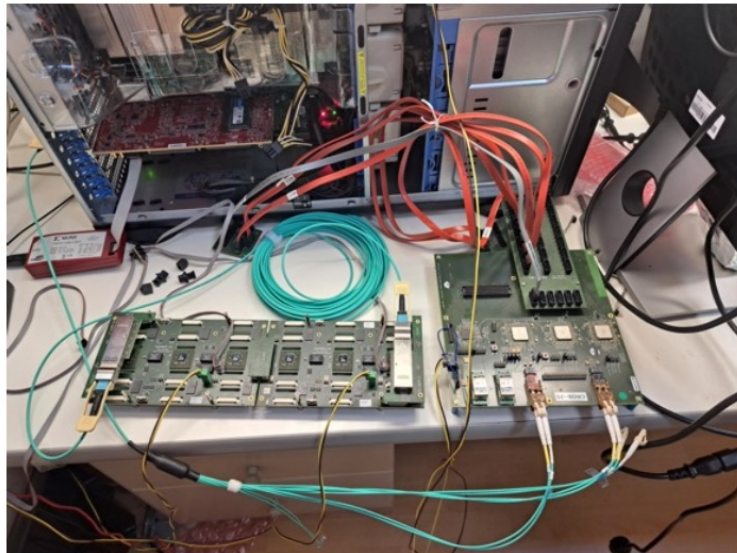
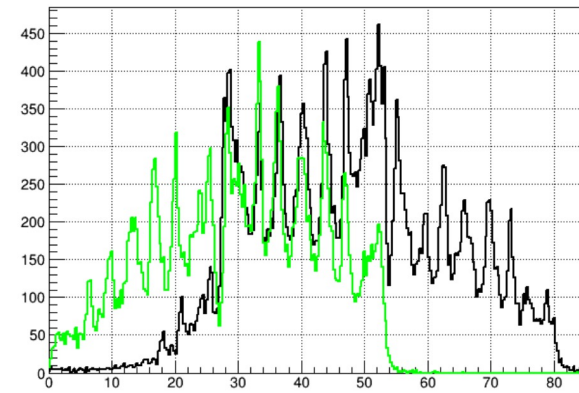
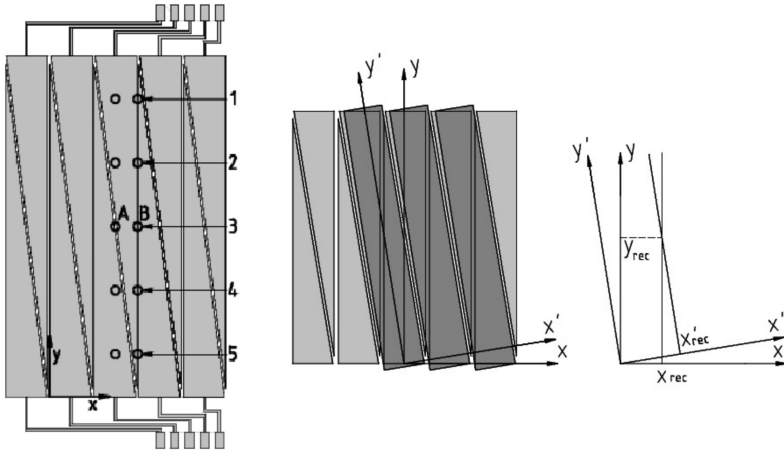


- Symmetric structure: 5 gaps x 2 stacks
- Gas gap thickness: 140 - 200 μm
- Discrete spacers
- Active area: strip length x 9 mm pitch x 32 strips
- Strip length: 56/96/196 mm (MRPC1a/MRPC1b/MRPC1c)
- Resistive electrodes: $\sim 1010 \Omega\text{cm}$, 0.7 mm (Chinese glass)
- Strip structure for Readout & HV electrodes
matched signal transmission line impedance to the input of the FEE *D. Bartos et al., Rom. Journ. of Physics 63, 901 (2018)*
- Differential readout
- Direct flow through the gas gaps



2D-MWPC (TRD)

(Review and references: https://niham.nipne.ro/hpd_courier_no7_August_2024.pdf)



Concluding remark



“We have found it of paramount importance that in order to progress we must recognize the ignorance and leave room for doubt. Scientific knowledge is a body of statements of varying degrees of certainty some most unsure, some nearly sure, none absolutely certain.”

Richard Feynman

Backup slides

