

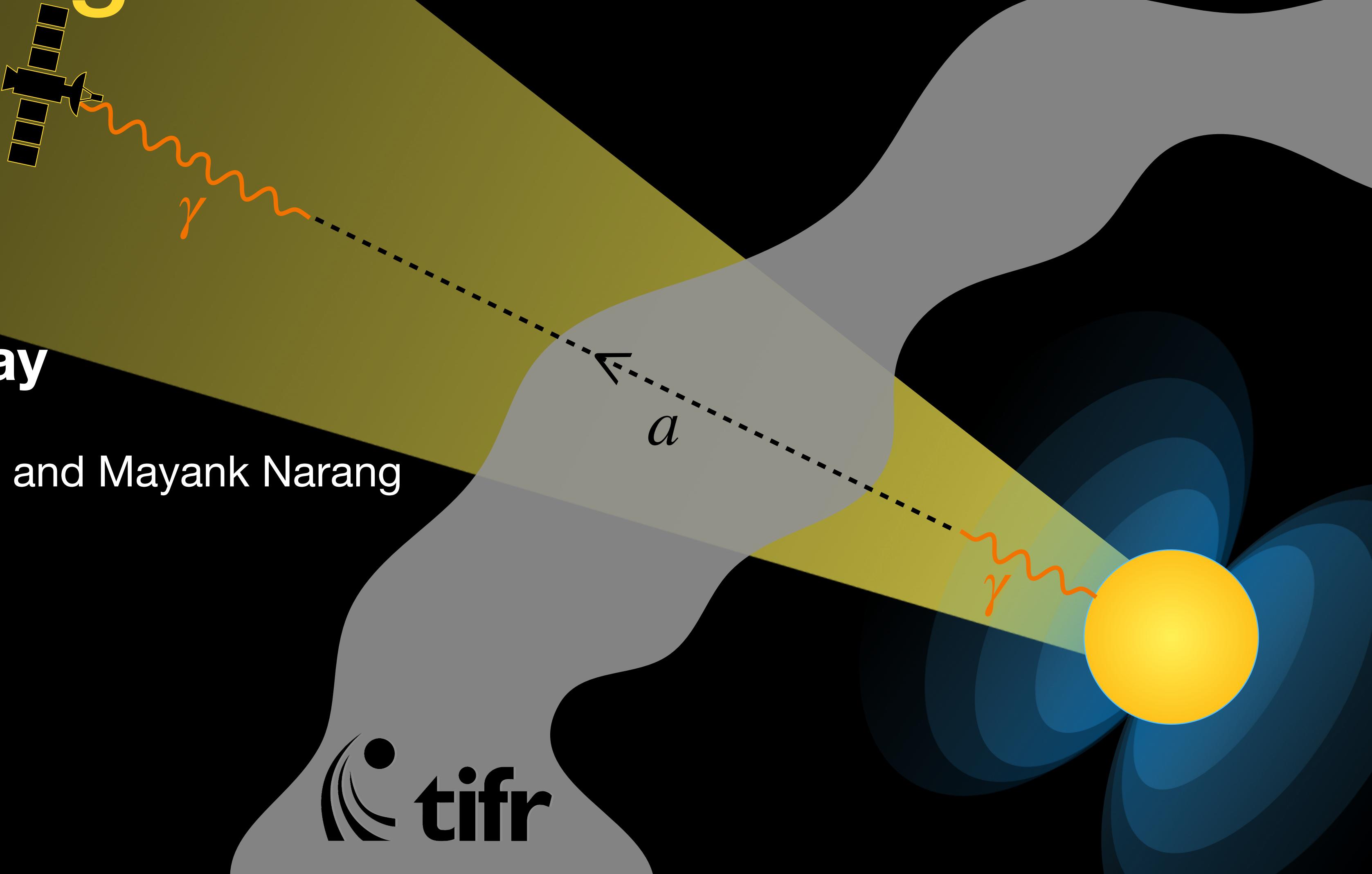
# Shedding Light on Axions Through Obscured Magnetars

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FPP 2024, CHEP, IISc  
9 Aug, 2024

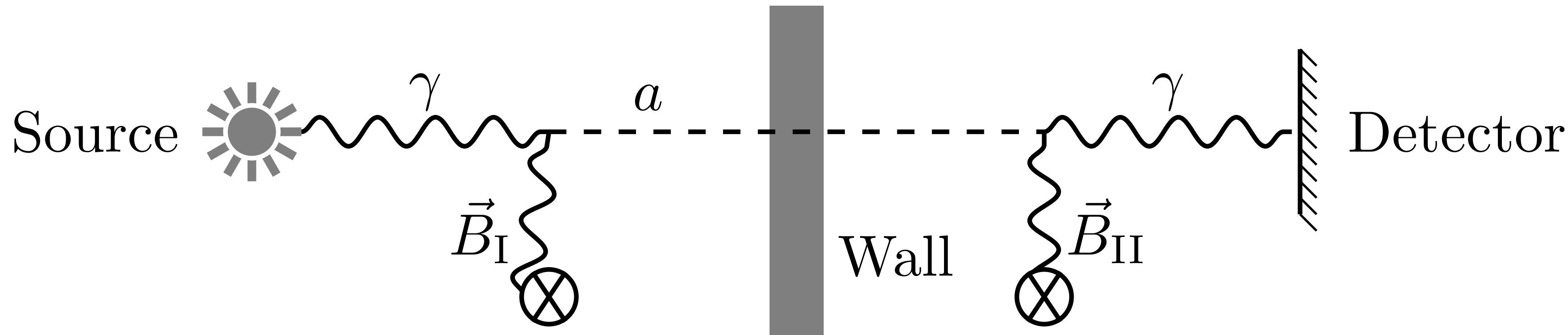


# Recap...

- The **Light Shining through the Wall (LSW)** technique of looking for ALPs

$$\mathcal{L}_{\text{ALP}} = \frac{1}{2} \partial^\mu a \partial_\mu a - \frac{1}{2} m_a^2 a^2 - \frac{1}{4} g_{a\gamma} a F_{\mu\nu} \tilde{F}^{\mu\nu}$$

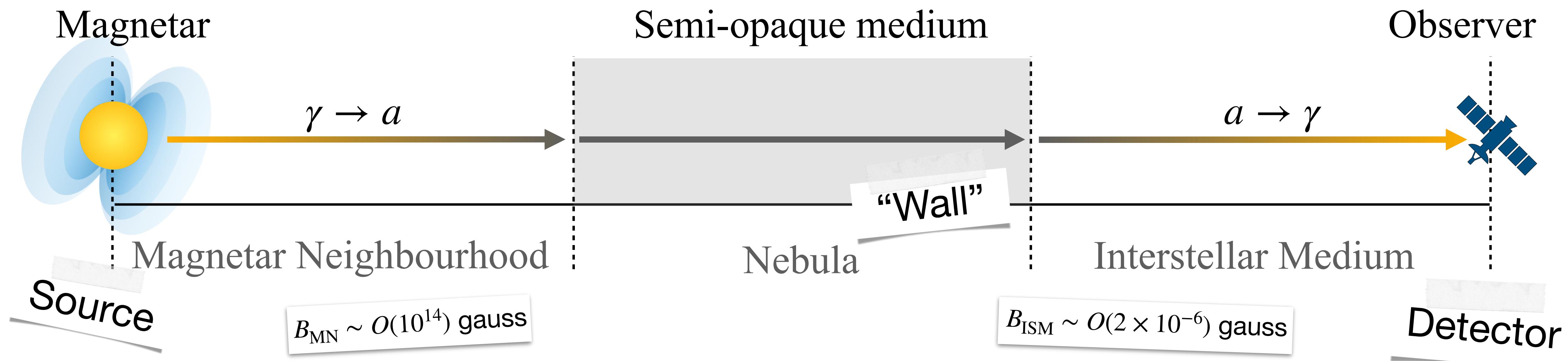
$$\mathcal{L}_{\text{ALP}} \subset g_{a\gamma} \vec{E} \cdot \vec{B} a$$



- **Lab-based experiments:** OSQAR, CROWS, ALPS, ALPS - II (upcoming...)

# Idea

- Applying the **LSW technique in astrophysics** by finding a suitable “laboratory”.
- **Obscured Magnetars** are an excellent candidate.



# Obtaining the constraint on $g_{a\gamma}$

- The **fraction** of photons that are finally **observed** must always be **larger than** the **fraction** of photons that **may escape** through the  $\gamma \rightarrow a \rightarrow \gamma$  process.
- Calculating  $P(\gamma \rightarrow a \rightarrow \gamma)$  dependence on  $g_{a\gamma}$  allows us to **constrain** the **ALP-photon coupling**.
- The “**escape probability**” is given by

$$P(\gamma \rightarrow a \rightarrow \gamma) = P_{\text{MN}}(\gamma \rightarrow a) \times P_{\text{ISM}}(a \rightarrow \gamma)$$

$$P_{\text{sur}} \gtrsim P_{\text{MN}} P_{\text{ISM}}$$

## I. $P_{\text{sur}}$

# The fraction of photons that survive...

- A magnetar candidate for the **LSW technique**: PSR J1622-4950

$$P_{\text{sur}}(E, E + \delta E) = \frac{F_{\text{obs}}(E, E + \delta E)}{F_0(E, E + \delta E)}$$

$$F_{\text{obs}}(\text{bin}) \approx (0.68 - 2.01) \times 10^{-18} \text{ erg cm}^{-2} \text{ s}^{-1}$$

- The ratio between the observed flux vs. the “expected” flux.

$$P_{\text{sur}} \approx (0.25 - 4.58) \times 10^{-4}$$

$$F_0(\text{bin}) = (0.44 - 2.72) \times 10^{-14} \text{ erg cm}^{-2} \text{ s}^{-1}$$

(0.71 – 0.98) keV

$F_{\text{obs}}(\text{total})$ (erg cm $^{-2}$ s $^{-1}$ )	$F_0(\text{total})$ (erg cm $^{-2}$ s $^{-1}$ )	$kT$ (keV)	$N_H$ (10 $^{22}$ cm $^{-2}$ )
$3.0^{+0.8}_{-0.6} \times 10^{-14}$	$11^{+9}_{-4} \times 10^{-14}$	$0.5 \pm 0.1$	$5.4^{+1.6}_{-1.4}$

Anderson et al., MULTI-WAVELENGTH OBSERVATIONS OF THE RADIO MAGNETAR PSR J1622–4950  
AND DISCOVERY OF ITS POSSIBLY ASSOCIATED SUPERNOVA REMNANT

# Conversion near the magnetar

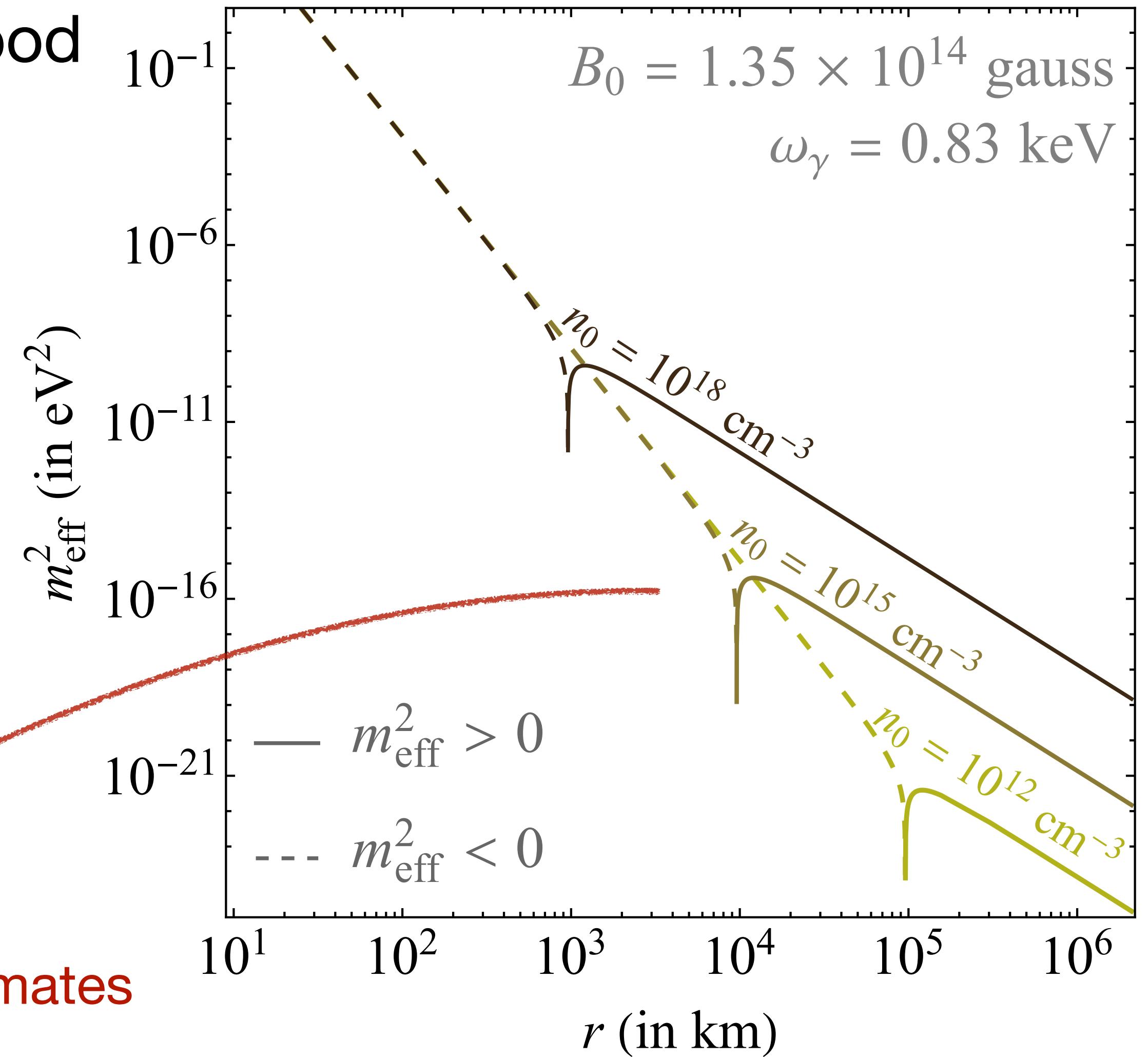
- **Resonance** in the Magnetar Neighbourhood

$$\underline{m_{\text{eff}}^2 = \frac{4\pi\alpha n_e}{m_e} - \frac{88\alpha^2\omega_\gamma^2}{135m_e^4} \frac{B^2}{2}}$$

$$n_e \approx n_0 \left( \frac{r}{r_0} \right)^{-3} \quad B \approx B_0 \left( \frac{r}{r_0} \right)^{-3}$$

- Resonance at  $m_{\text{eff}}^2 = m_a^2$
- Fluctuations may lead to multiple pairs of resonances.

Large uncertainties in charge density estimates



## II. $P_{\text{MN}}$

# Conversion near the magnetar...

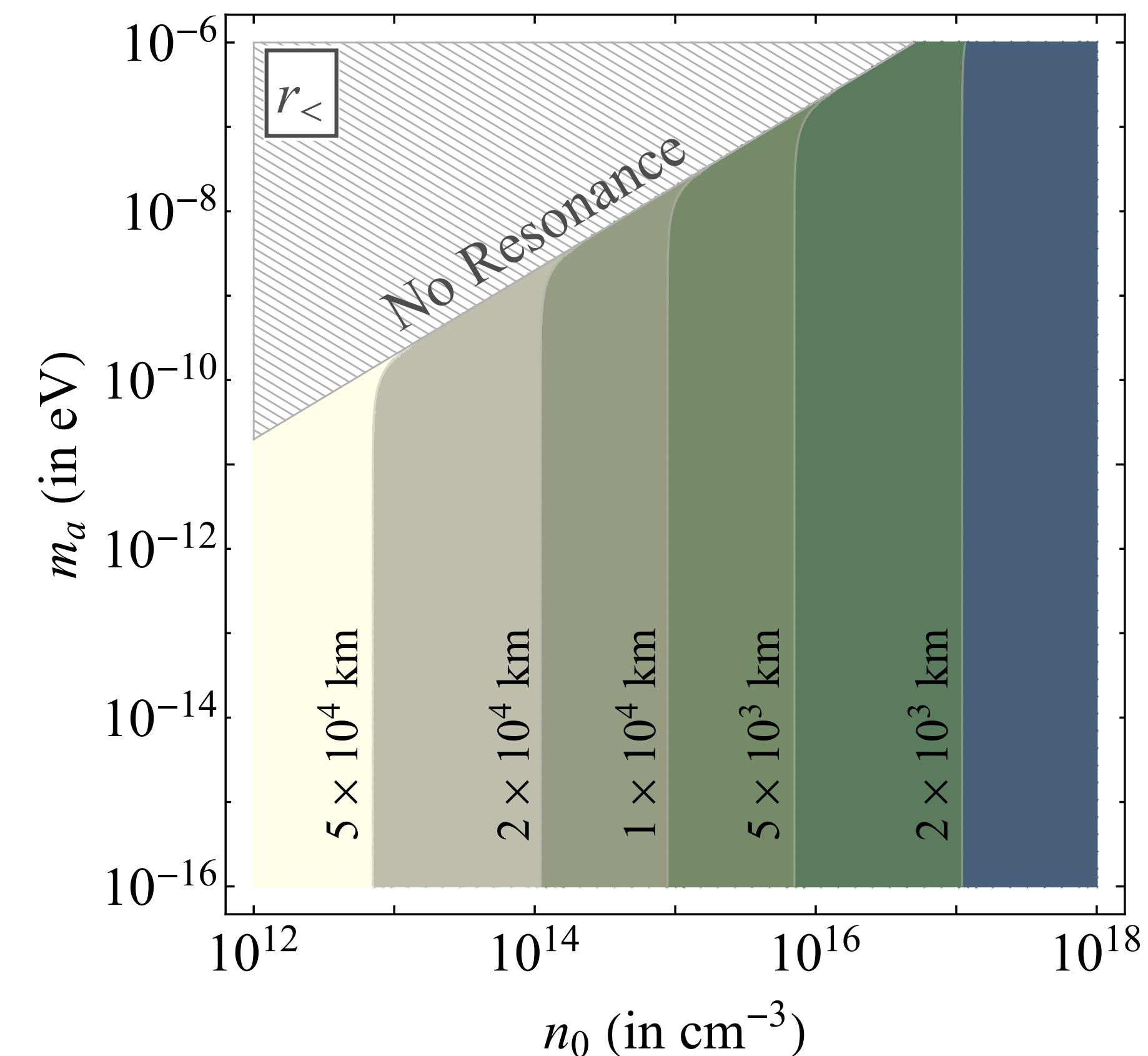
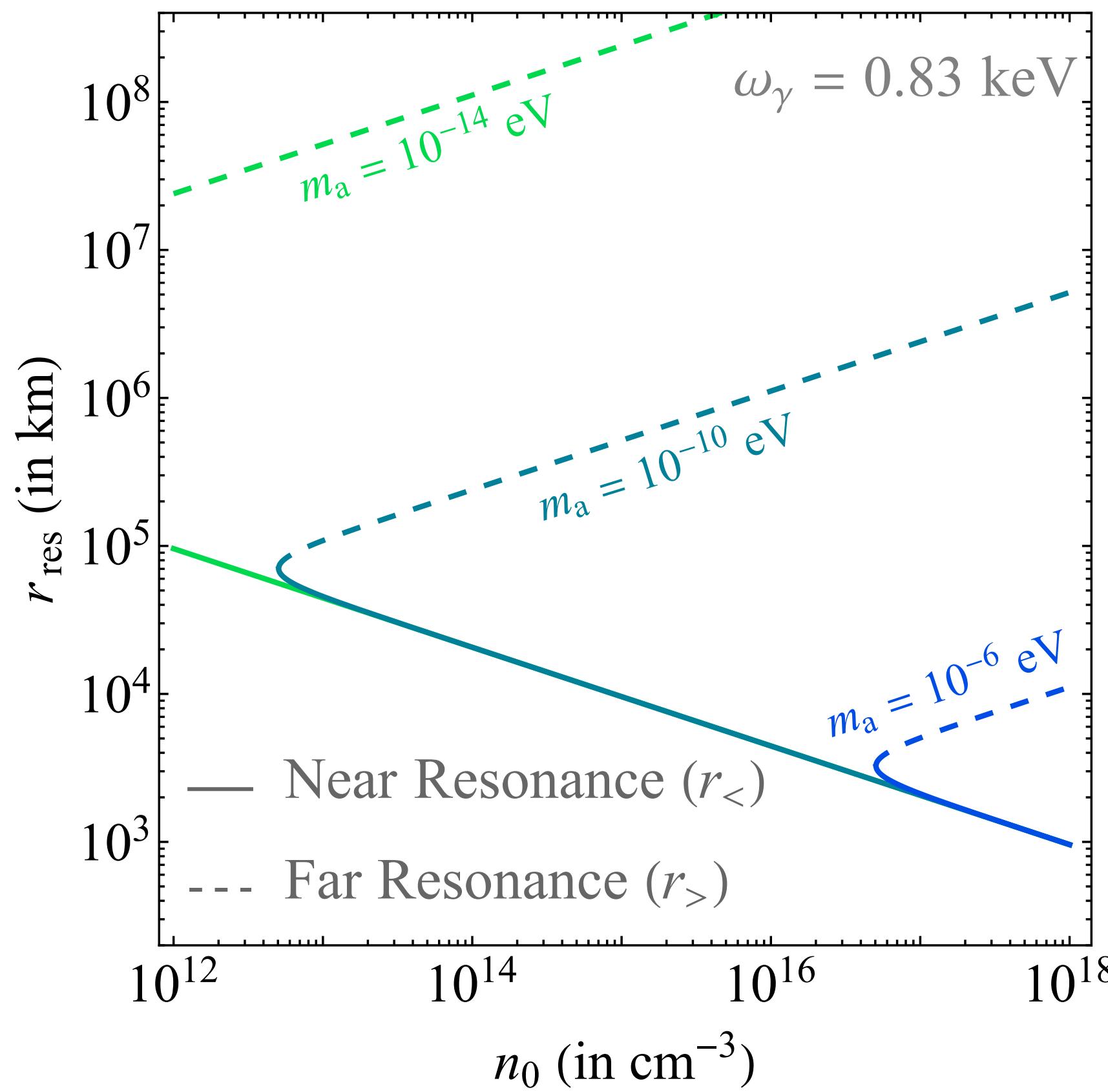
- Even in the simplified approximation: at least two resonances...

$$r_{\lessgtr} = r_0 \left[ \frac{2 c_B \omega_\gamma^2 B_0^2}{c_n n_0 \pm \sqrt{c_n^2 n_0^2 - 4 c_B B_0^2 \omega_\gamma^2 m_a^2}} \right]^{1/3}$$

$c_n \equiv \frac{4\pi\alpha}{m_e}, \quad c_B \equiv \frac{88\alpha^2}{270 m_e^4}$

$$r_< \approx r_0 \left[ \frac{c_B \omega_\gamma^2 B_0^2}{c_n n_0} \right]^{1/3}$$

$$r_> \approx r_0 \left[ \frac{c_n n_0}{m_a^2} \right]^{1/3}$$



## II. $P_{\text{MN}}$

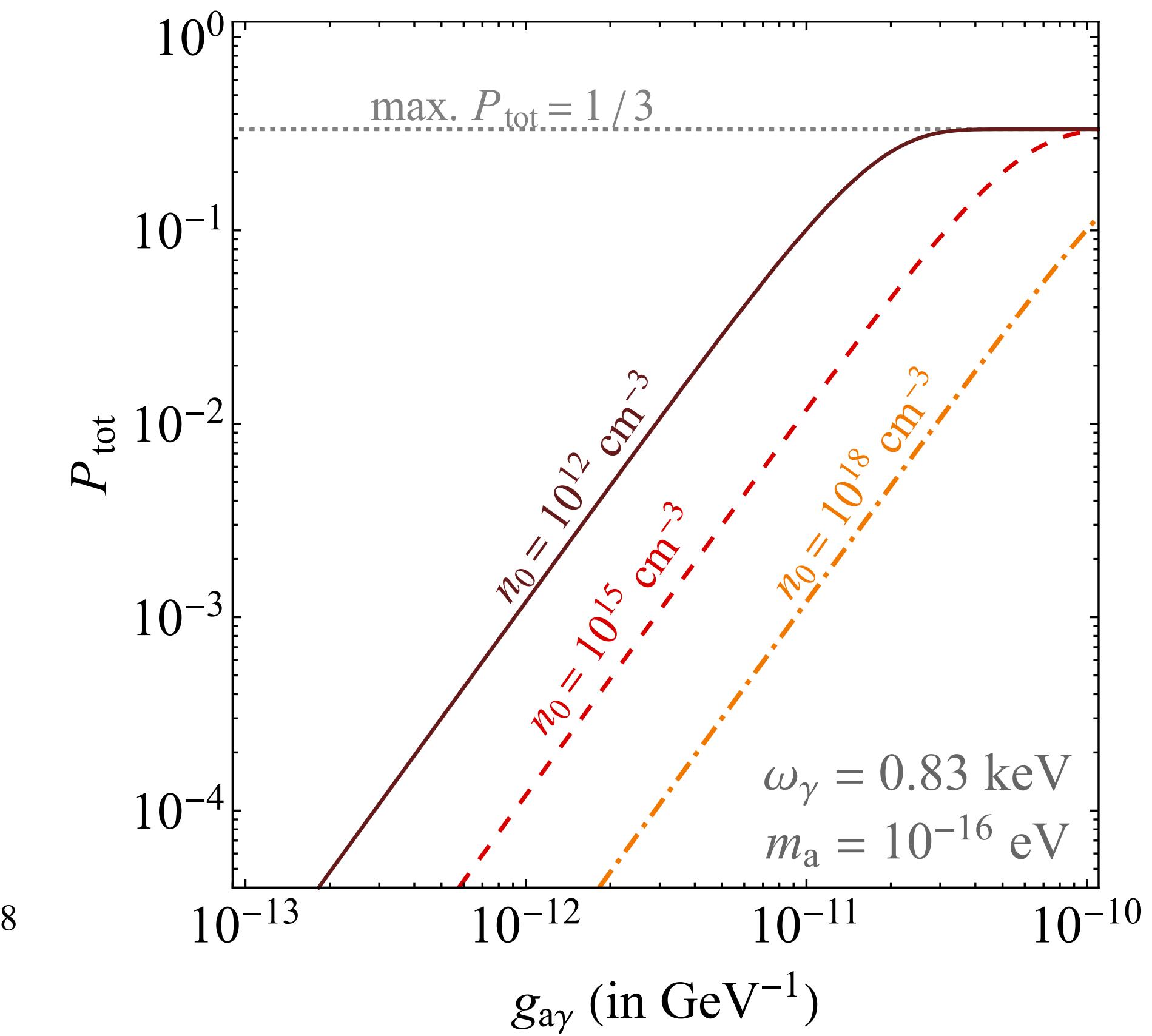
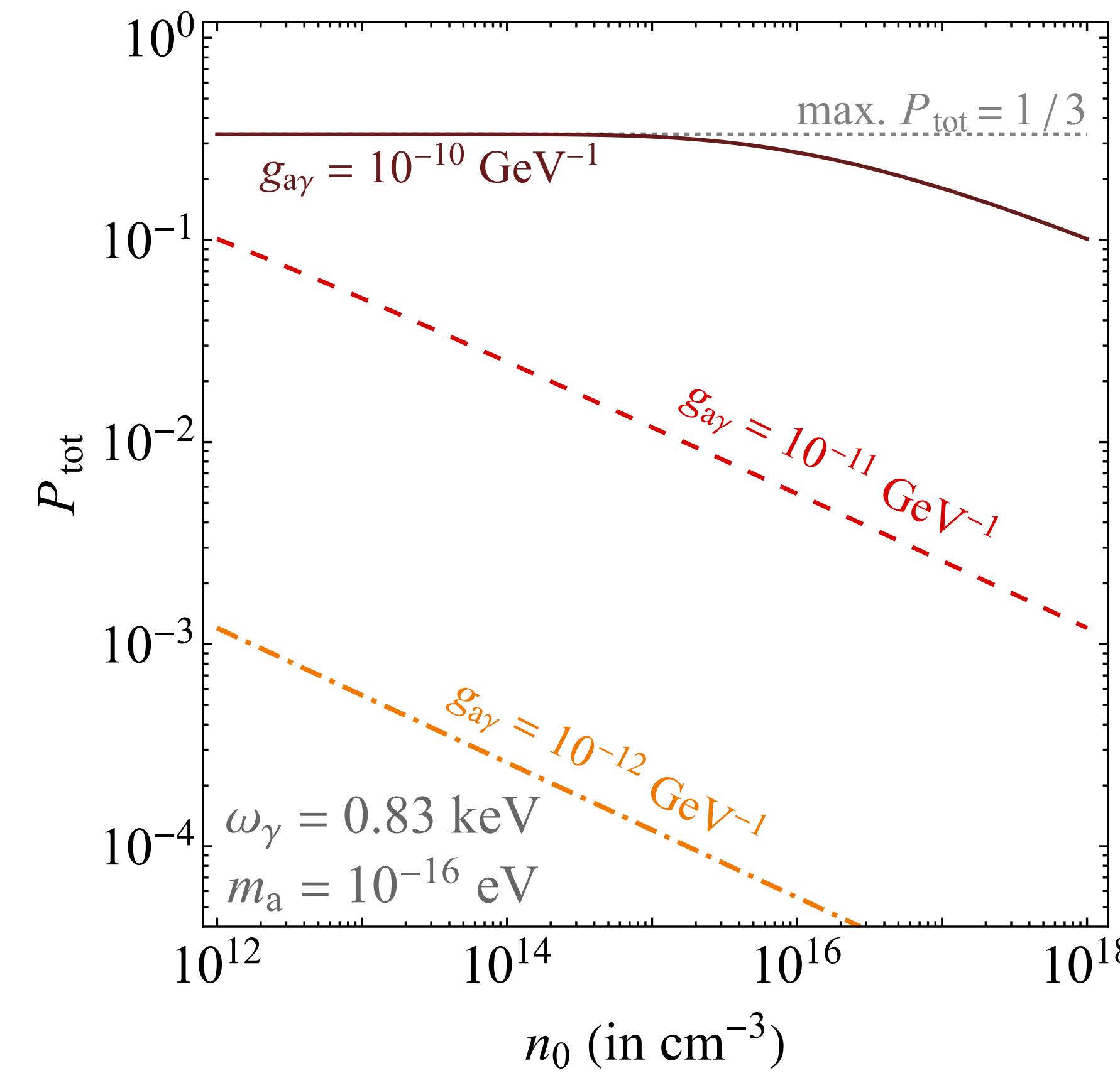
# Conversion near the magnetar...

- We use conservative estimates for  $P_{\text{MN}} \equiv P_{\text{tot}}$

$$P_{\text{tot}} \approx \frac{1}{3} \left( 1 - e^{-\frac{3\pi}{4} \Gamma_{\text{tot}}} \right)$$

$$\Gamma_{\text{tot}} = \frac{2g_{a\gamma}^2 \omega}{m_a^2} \sum_i^{n_r} B_{T,i}^2 \mathcal{R}_i$$

$$\mathcal{R}_i \equiv \left| \frac{d \ln m_{\text{eff}}^2}{dl} \right|_{l=l_i}^{-1}$$



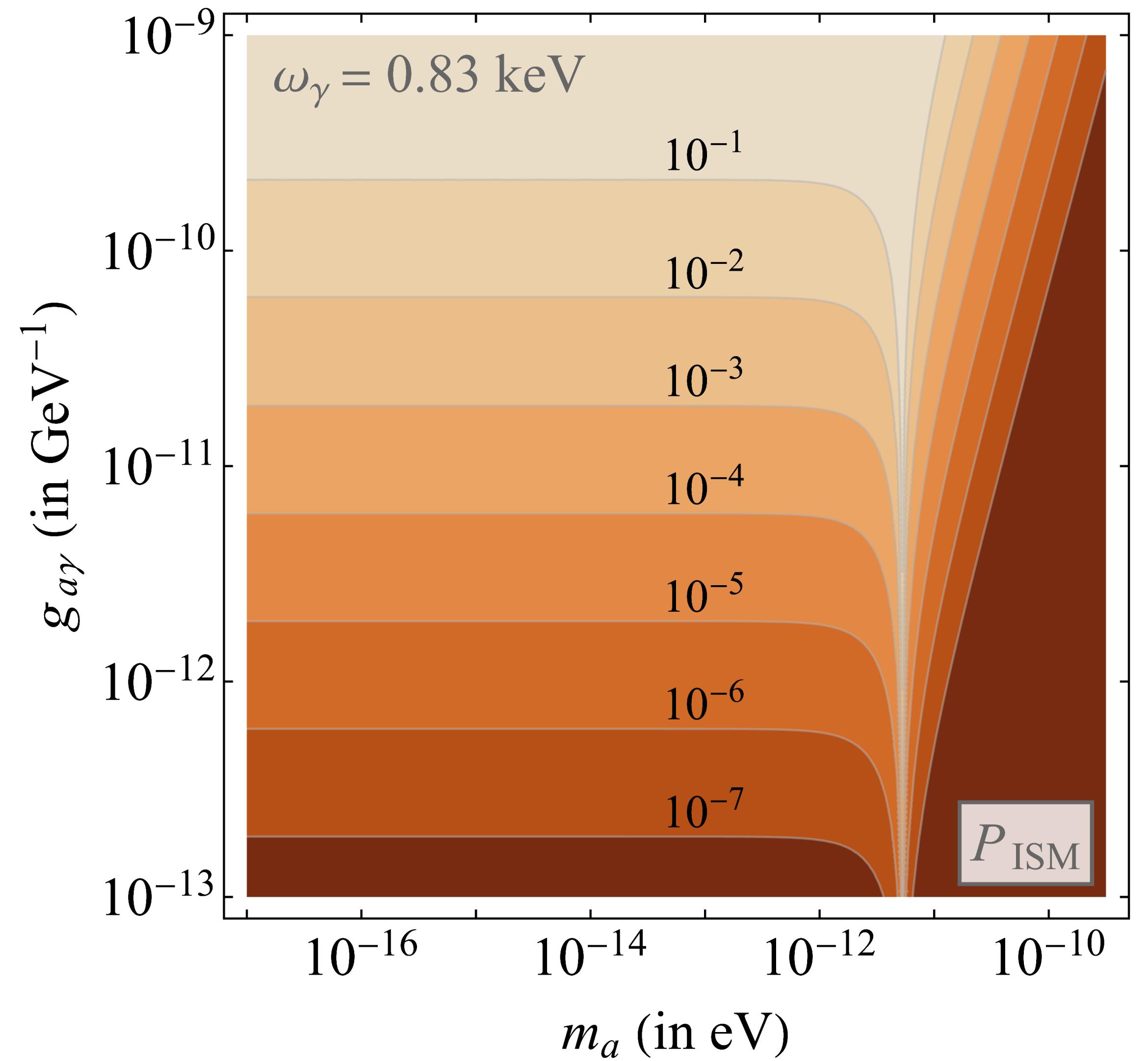
### III. $P_{\text{ISM}}$

# Conversion probability in ISM

- The conversion probability in the ISM will be averaged out.
- We take the electron density and the magnetic field value to be:

$$n_{\text{ISM}} \approx 2 \times 10^{-2} \text{ cm}^{-3}$$

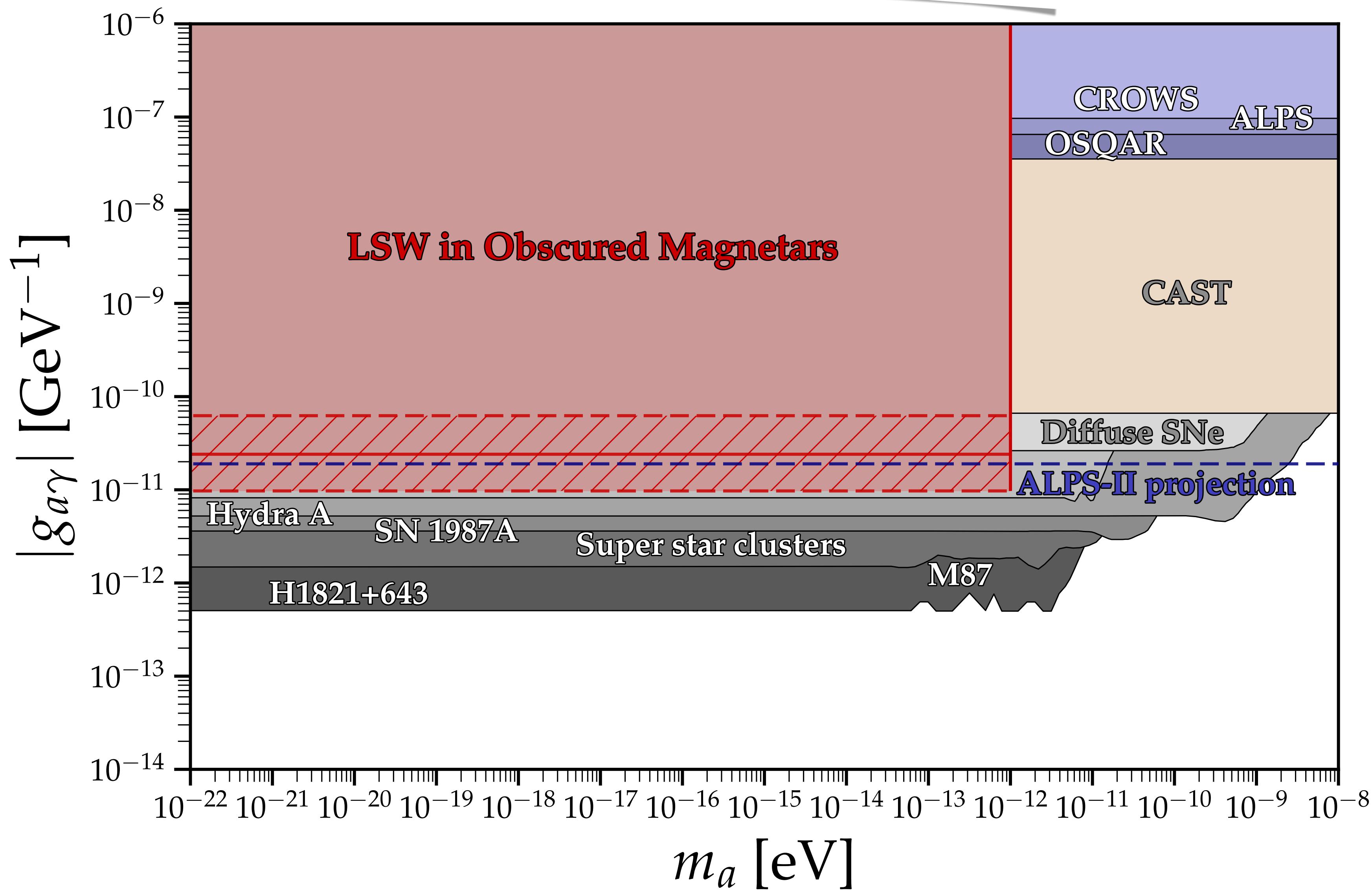
$$B_{\text{ISM}} \approx 2 \times 10^{-6} \text{ gauss}$$



# Results

- Complementary to existing astrophysical bounds.
- Better than all current lab-based LSW bounds.
- Competitive even with ALPS-II projections for  $m_a \lesssim 10^{-12}$  eV.

$$g_{a\gamma} \lesssim (10^{-11} - 10^{-10}) \text{ GeV}^{-1}$$



# Take home message

- The idea: LSW + astrophysical systems
- The candidate: obscured magnetars
- The result:  $g_{a\gamma} \lesssim (10^{-11} - 10^{-10}) \text{ GeV}^{-1}$  for low mass ALPs ( $m_a \lesssim 10^{-12} \text{ eV}$ ).

arXiv:2311.14298 [hep-ph]

