

# Quantum network technology

The second life of rare-earth crystals

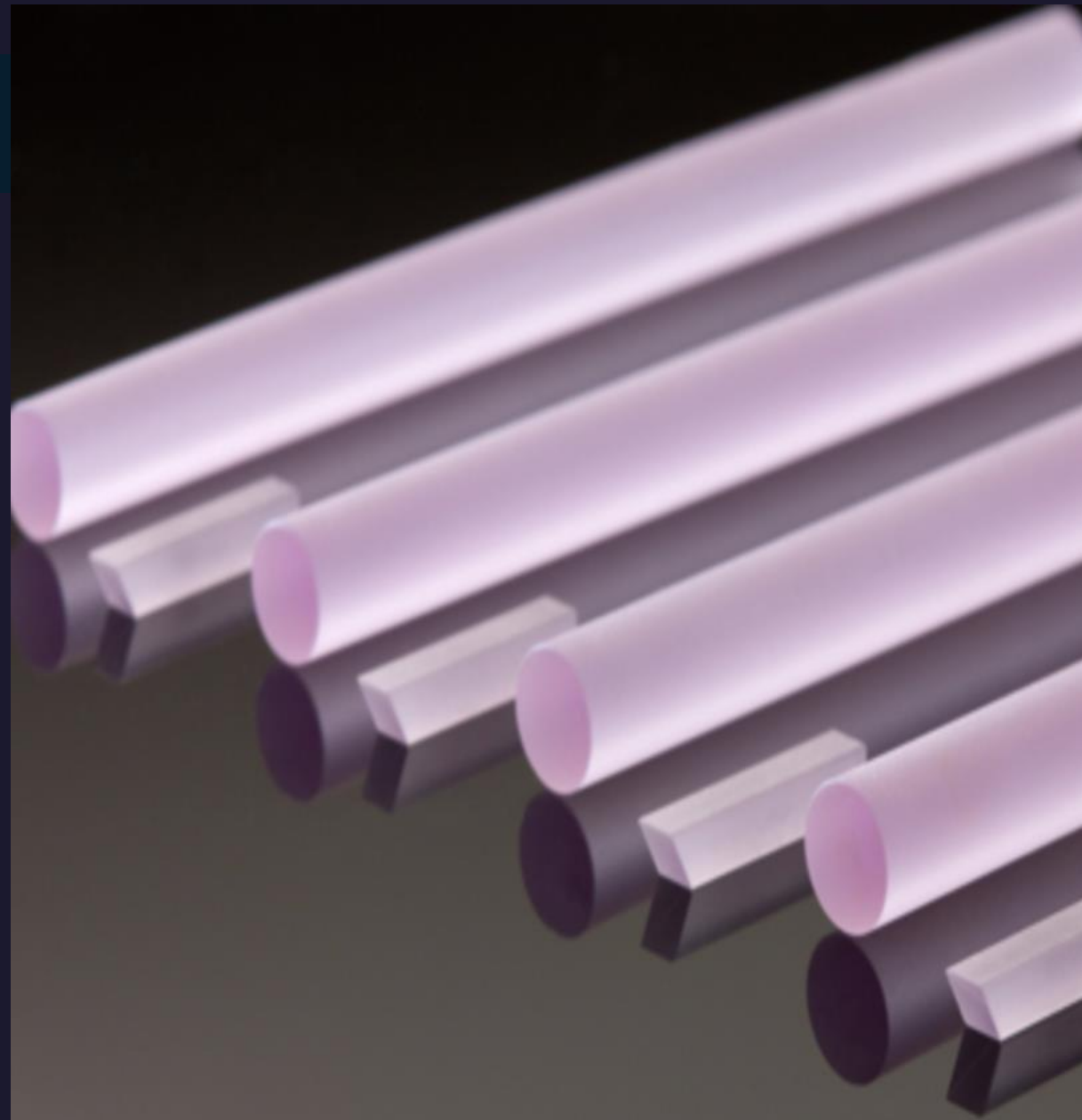
Wolfgang Tittel

*University of Geneva*

*Constructor Institute of Technology & Constructor University*



**C>ONSTRUCTOR**



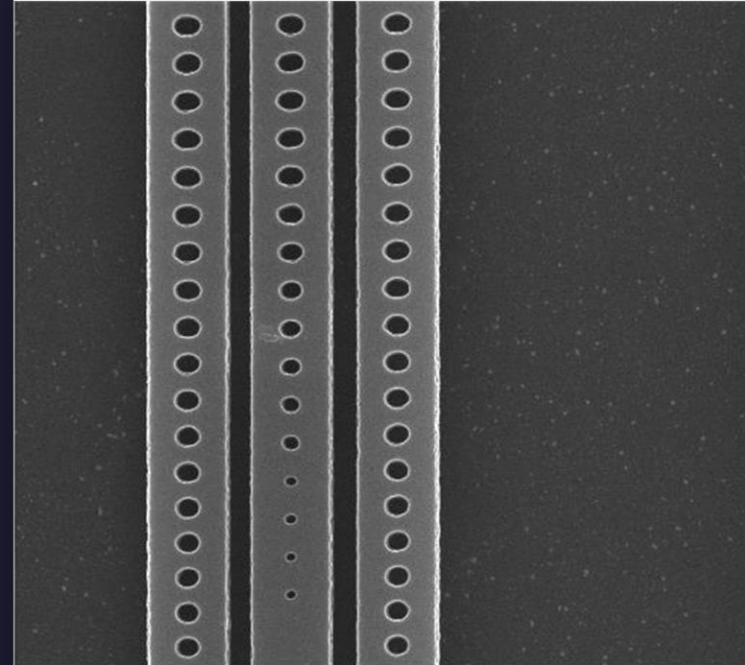
# Elements of a quantum network

- Interacting centers for matter-based qubits
- Individual optical centers for single and entangled photons
- Large ensembles for high-capacity, feed-forward-controlled optical quantum memories
- For compatibility and scalability, these components should be based on the same type of defect and be on-chip integrable using standard photonics technology.



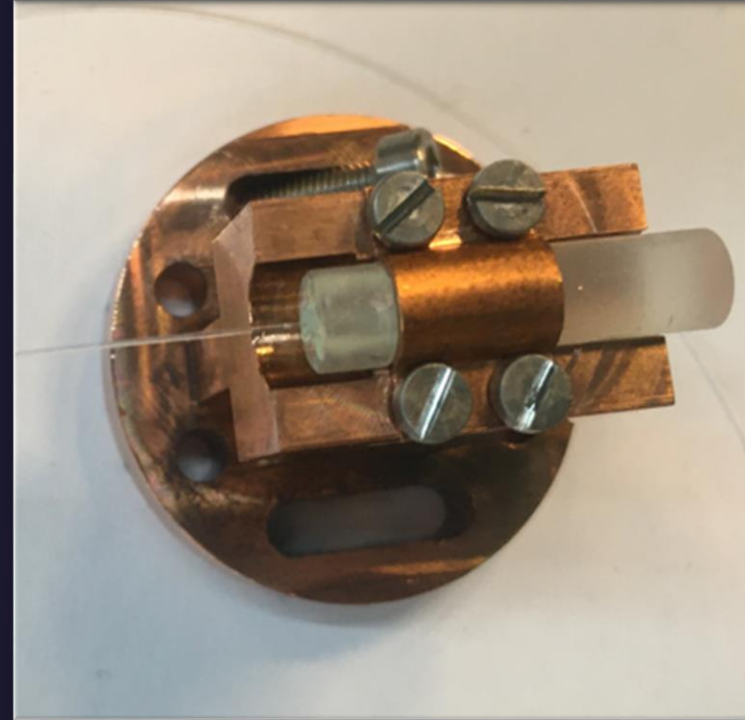
# Outline

- Rare-earth ion-doped crystals
- Single photons based on individual rare-earth ions
- Optical quantum memory using ensembles of rare-earth ions
- Outlook and conclusion



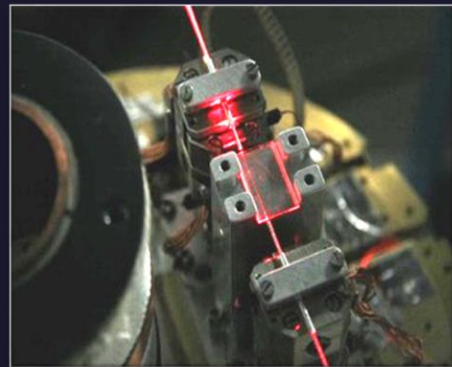
HPW	WD	mag	det	mode
8.29 $\mu\text{m}$	5.5 mm	25 000 x	TLD	SE

2  $\mu\text{m}$

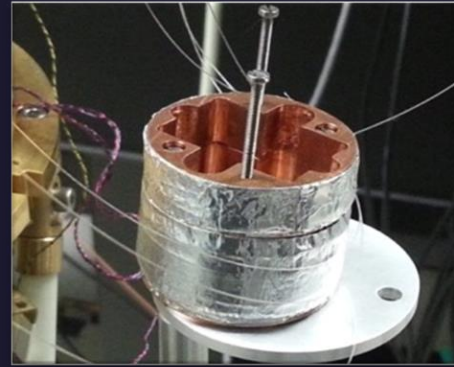


# Rare-earth ion-doped crystals

Periodic table showing the lanthanide and actinide series highlighted in red.

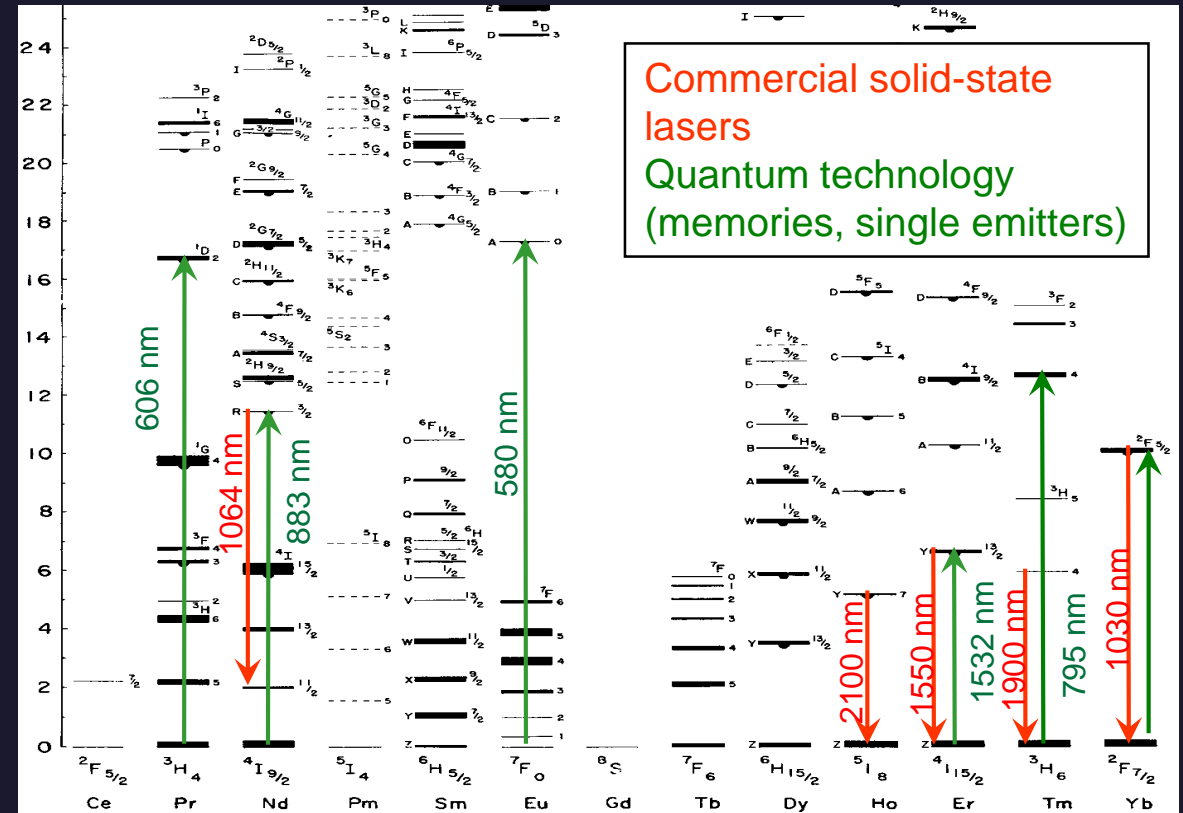


Saglamyurek, WT *et al.*, Nature 469, 512-515 (2011).



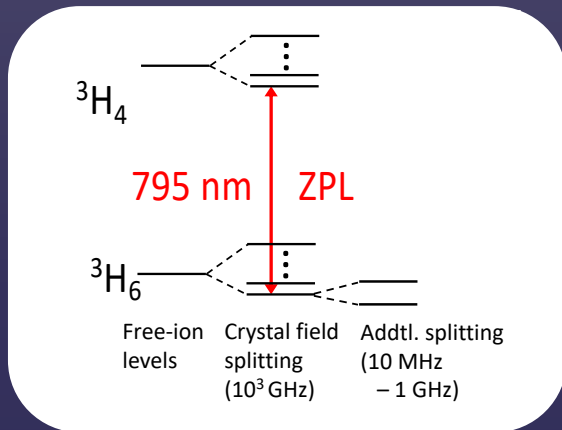
Saglamyurek, WT *et al.*, Nature Phot. 9, 83 (2015).

Energy of electronic levels [ $\times 10^3 \text{ cm}^{-1}$ ]

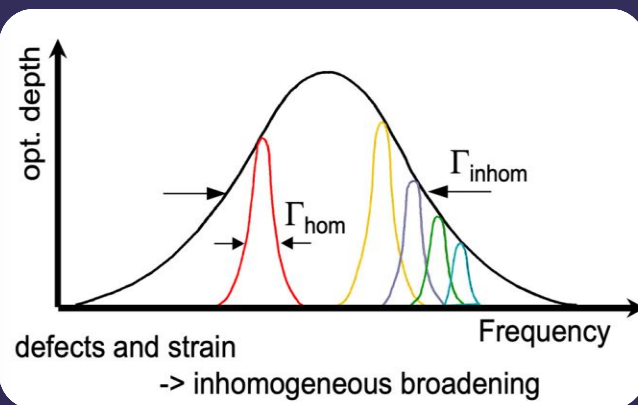


Commercial solid-state lasers  
Quantum technology (memories, single emitters)

# Rare-earth crystals: a brief introduction



Simplified level structure of  $Tm^{3+}$ -doped crystals

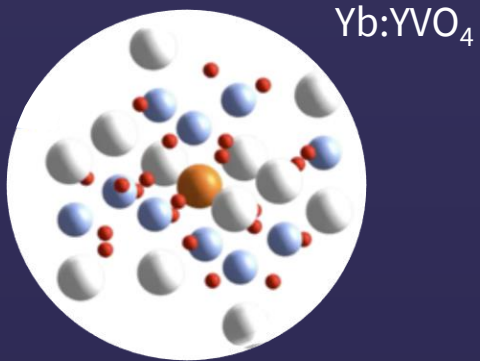


- 1) Transitions (zero-phonon lines) in the visible and near infrared  
-> quantum communication
- 2)  $\Gamma_{\text{inhom}} \approx 100 \text{ MHz} - 500 \text{ GHz}$   
-> *broadband quantum memory*
- 3) Excited states with very long lifetimes (ms)  
-> *difficult to observe single-photon emissions*
- 4) At  $T < 2 \text{ K}$ :  $\Gamma_{\text{hom}}^{\text{opt}} \approx 50 \text{ Hz} - 100 \text{ kHz}$  ->  $T_2 = 4 \text{ ms}$   
-> *high-capacity and long-lived quantum memory*
- 5) At  $T < 2 \text{ K}$ : ground states with long  $T_1$  (d) and long  $T_2$  (h)  
-> *long-lived quantum memory and qubits*
- 6) Electric dipole-dipole interaction between neighboring ions  
-> *quantum gates*

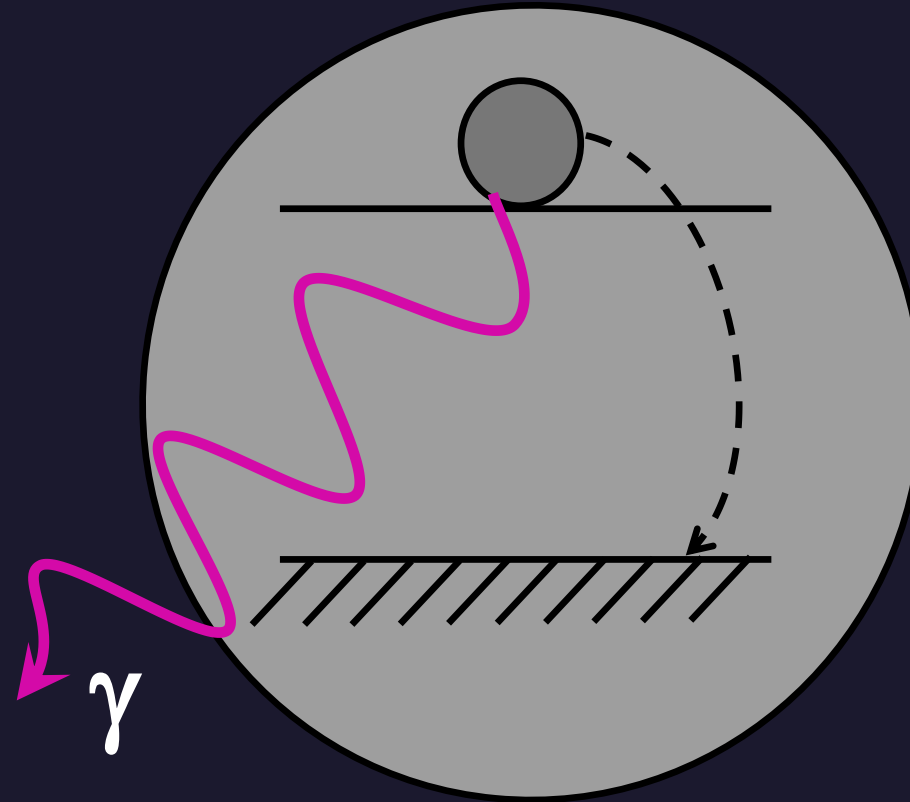
Promising for optical quantum memory and QIP. But not for single-photon emitters.

# How to create a single photon? Spontaneous emission from a single emitter

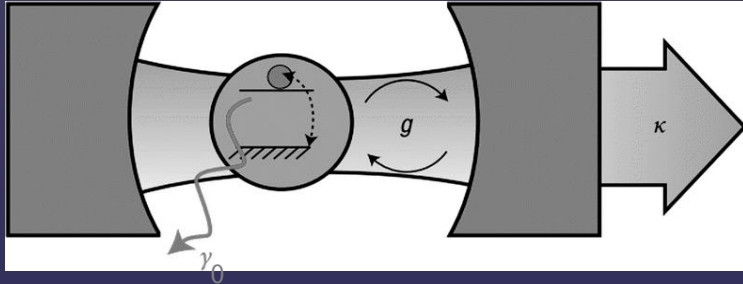
- The long optical lifetimes in rare-earth ions result in very small decay rates
- Photons will be emitted into random directions



Single rare-earth ion



# Creating (and observing) single photons



$$g = \frac{\mu_{eg}}{\hbar} \sqrt{\frac{\hbar\omega}{2\epsilon_M V_{mode}}}$$

For  $\kappa \gg g \gg \gamma_0$  (weak coupling regime):

$$\gamma' = F_P \gamma_0$$

$\gamma_0$ : vacuum emission rate

$$F_P = \frac{3}{4\pi^2} \left(\frac{\lambda}{n}\right)^3 \frac{Q}{V} \left| \frac{E(\mathbf{r})}{E_{\max}} \right|^2 \quad \text{Purcell factor}$$

# for atom in max field region

Quality factor

$$Q = \frac{\nu}{\Delta\nu}$$

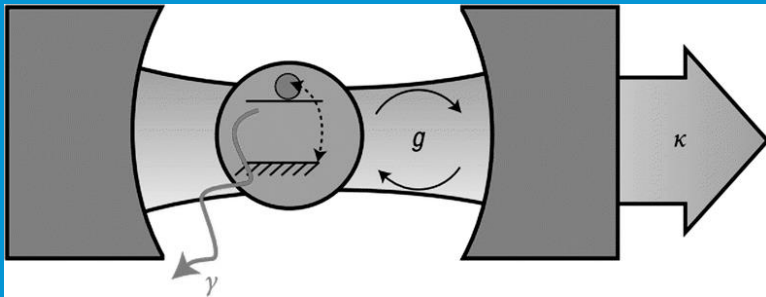
Mode volume

$$V = \int d^3\mathbf{r} \frac{\epsilon(\mathbf{r}) |E(\mathbf{r})|^2}{\max\{\epsilon(\mathbf{r}) |E(\mathbf{r})|^2\}}$$

Needed: cavity with small V and large Q

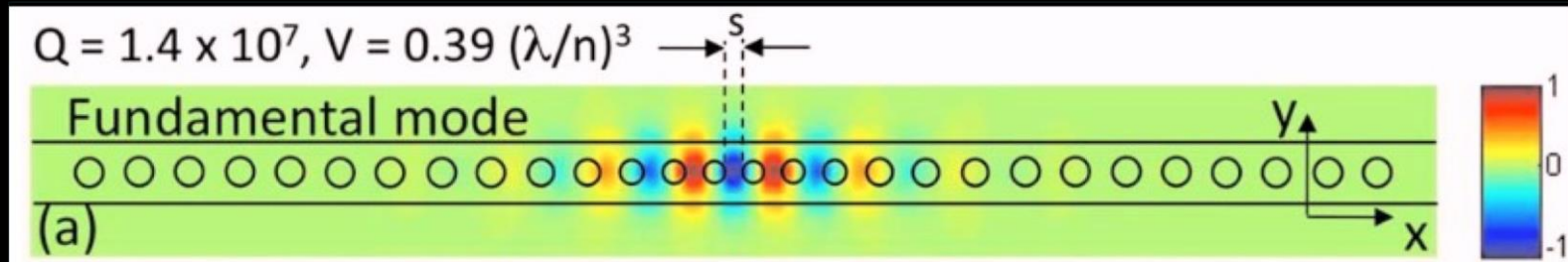
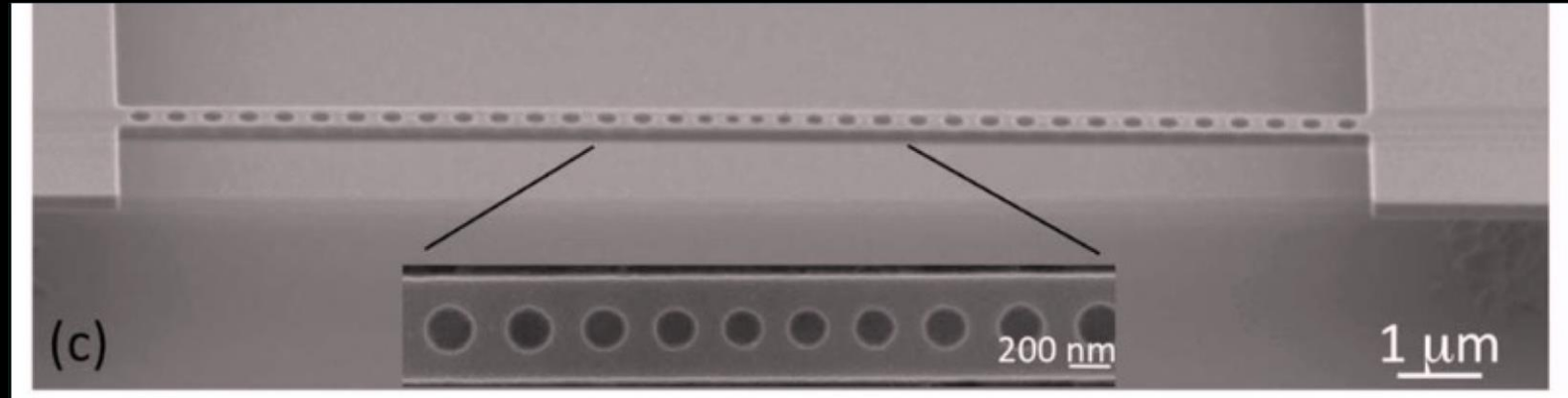
## Nano-photonic crystal cavities

The Purcell effect:  
atom-light interaction  
in the weak coupling  
regime



$$g = \frac{\mu_{eg}}{\hbar} \sqrt{\frac{\hbar\omega}{2\varepsilon_M V_{\text{mode}}}}$$

$$F_P = \frac{3}{4\pi^2} \left(\frac{\lambda}{n}\right)^3 \frac{Q}{V} \quad \text{Purcell factor}$$



Deotare et al. *Appl. Phys. Lett.* 94, 121106 (2009)



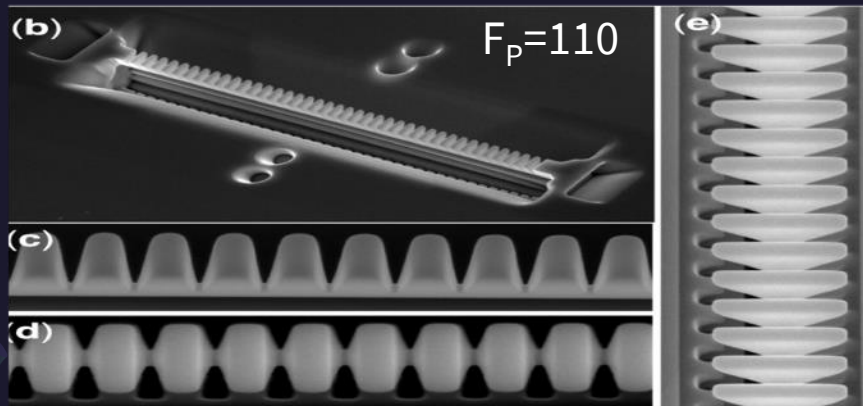
# Single photons from individual rare-earth ions

- **Purcell-entangled light-matter interaction** has enabled the observation of true single photons from individual rare-earth ions coupled to nanocavities.

*Heterogeneous approach*

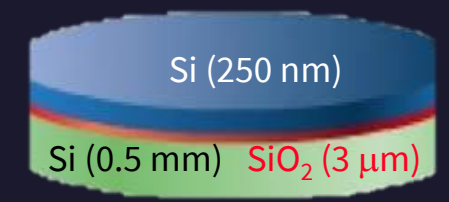
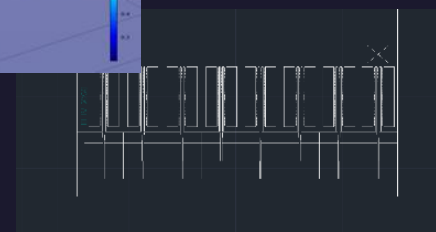
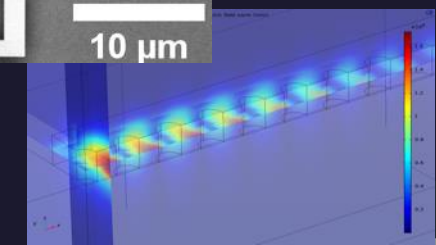
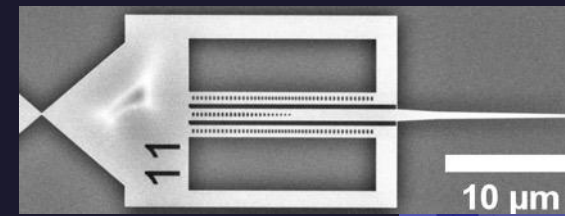


*Homogeneous approach*



Faraon group: photonic crystal nano-cavity milled out of Nd:YVO<sub>4</sub> (PRL 2018)

# A single-photon source based on Er:LiNbO<sub>3</sub>



Design and Simulations

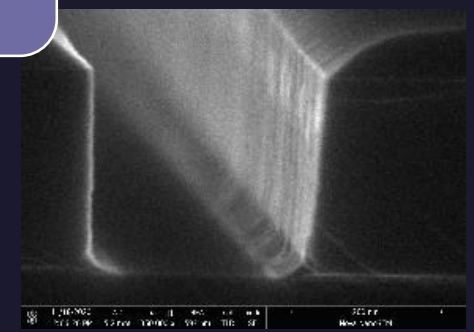
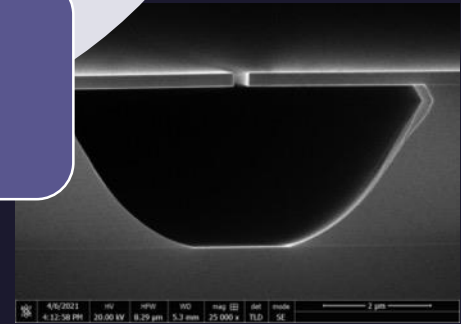
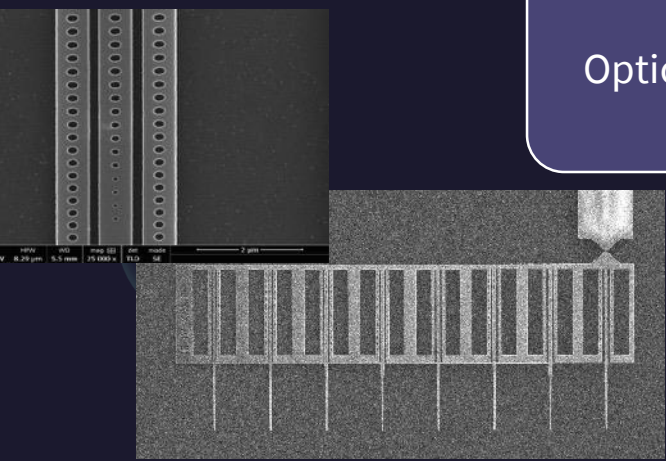
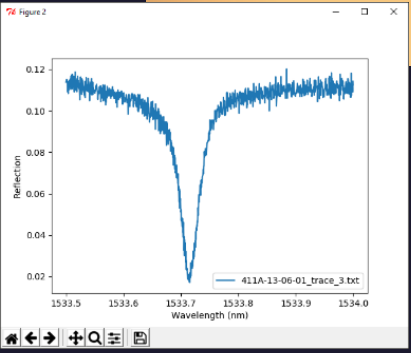
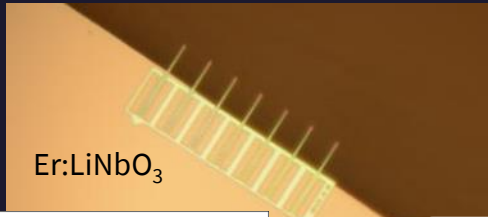
Patterning:  
E-beam lithography

Patter transfer:  
Reactive ion etching

Under-cutting:  
Hydrofluoric acid

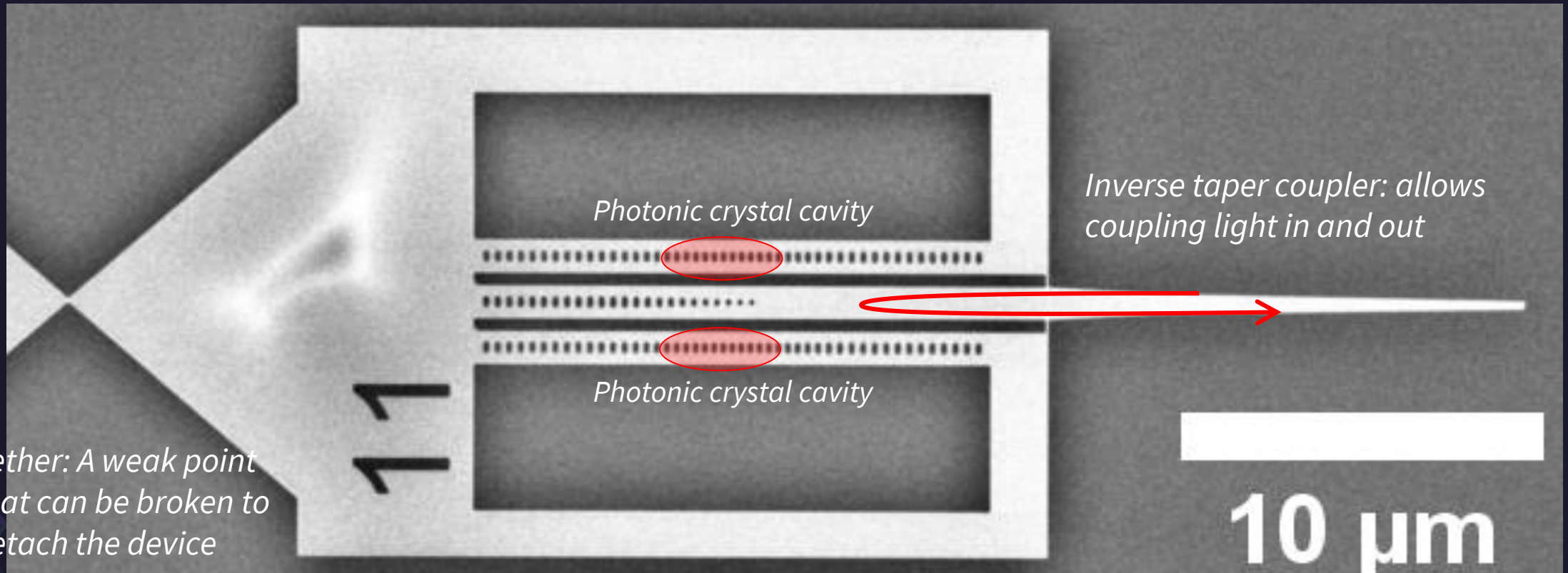
Optical characterization in air,  
and after transfer on  
Er:LiNbO<sub>3</sub>, at T=293K and T=4K  
Spectra and quality factor

Inspection:  
Optical microscope &  
SEM



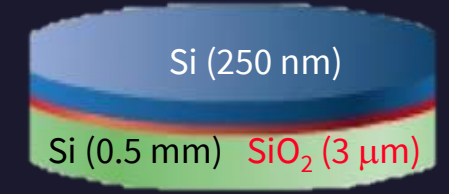
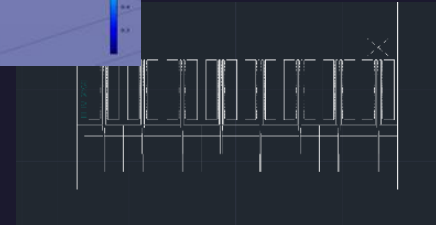
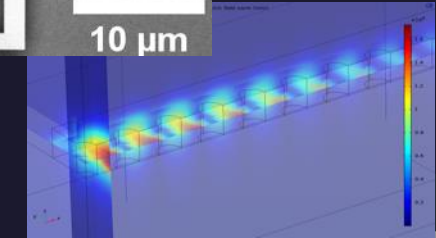
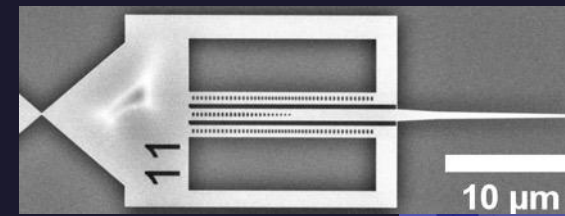
# Design

*A “bus” waveguide (with Bragg reflector):  
couples to 2 nanobeam cavities*



*Tether: A weak point  
that can be broken to  
detach the device*

# A single-photon sources based on Er:LiNbO<sub>3</sub>



Design and Simulations

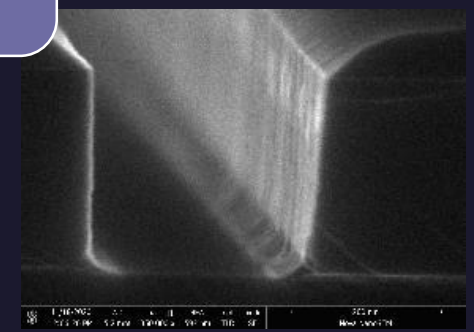
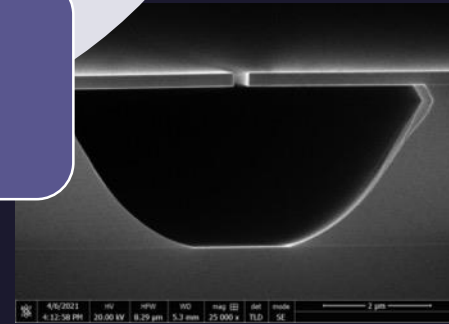
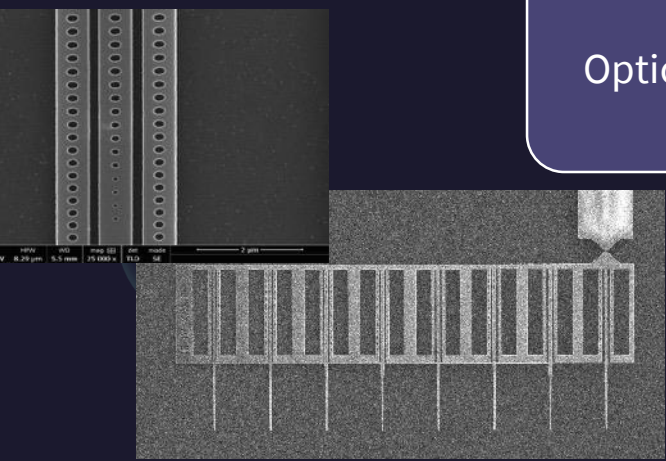
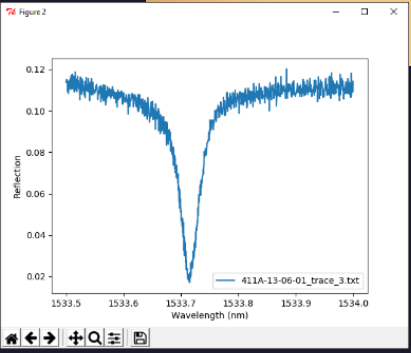
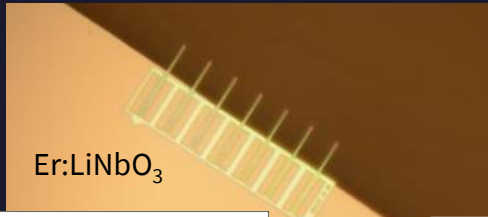
Patterning:  
E-beam lithography

Patter transfer:  
Reactive ion etching

Under-cutting:  
Hydrofluoric acid

Inspection:  
Optical microscope &  
SEM

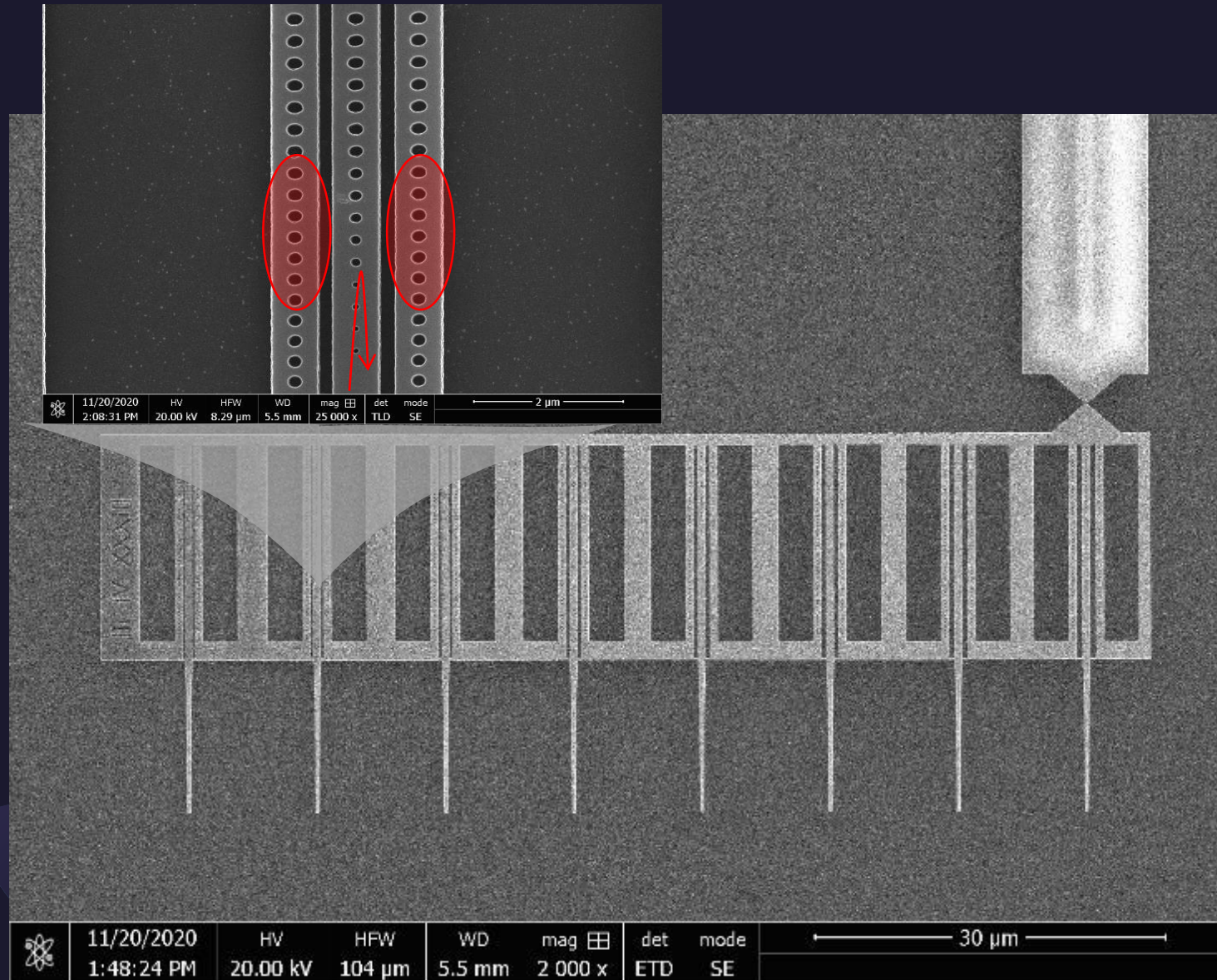
Optical characterization in air,  
and after transfer on  
Er:LiNbO<sub>3</sub>, at T=293K and T=4K  
Spectra and quality factor



# Cavity characterization

Mode volume	calculated	$0.1 \mu\text{m}^3$
Quality factor (on Er: LiNbO <sub>3</sub> )	measured	50 000
Purcell factor ( $E=E_{\text{max}}$ , $\beta=1$ )	predicted	1 000

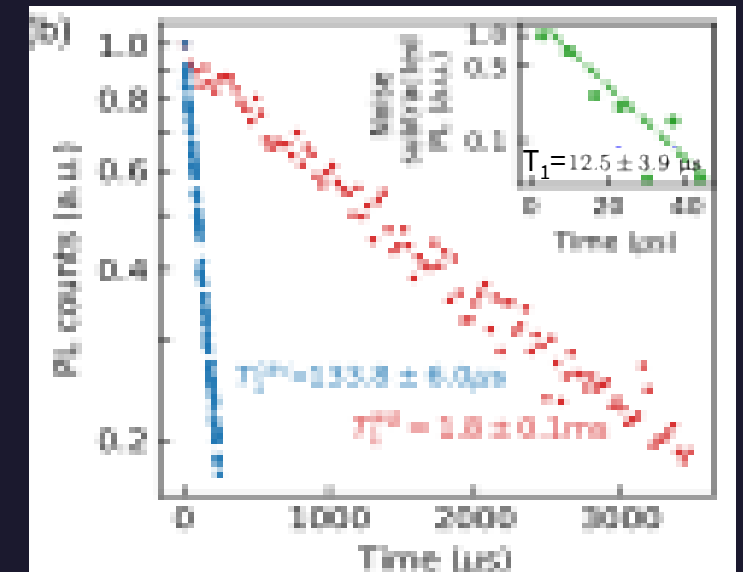
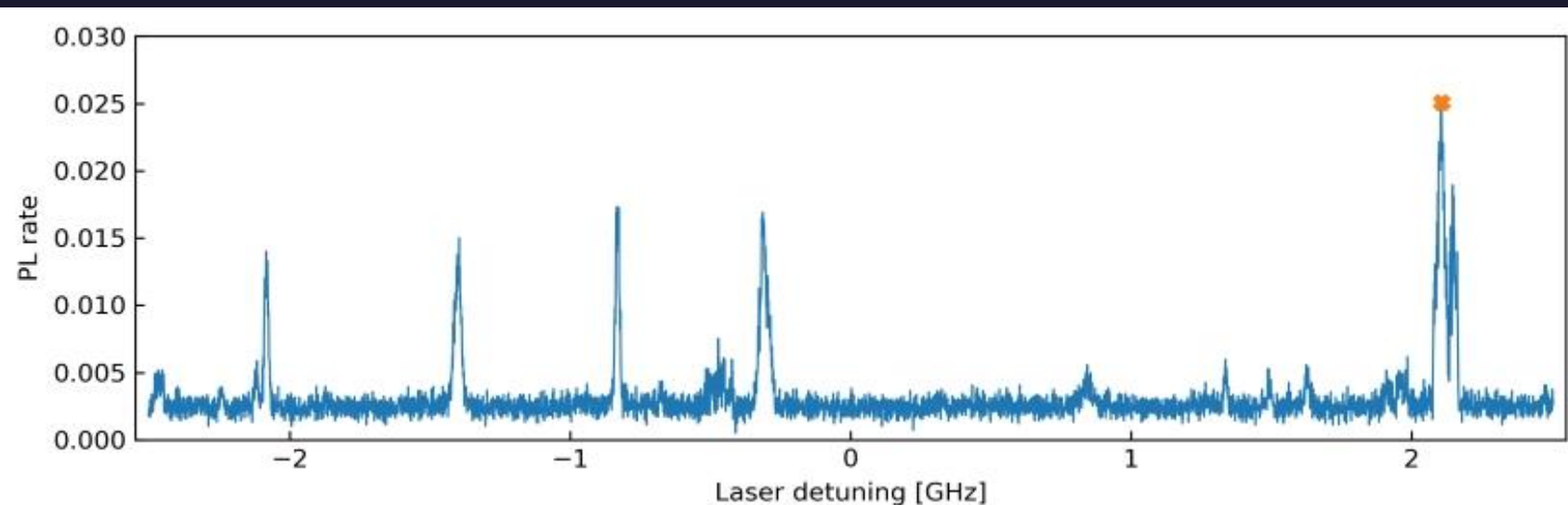
$$F_P = \frac{3}{4\pi^2} \left(\frac{\lambda}{n}\right)^3 \frac{Q}{V} \quad \text{Purcell factor}$$



# Purcell-enhanced emission

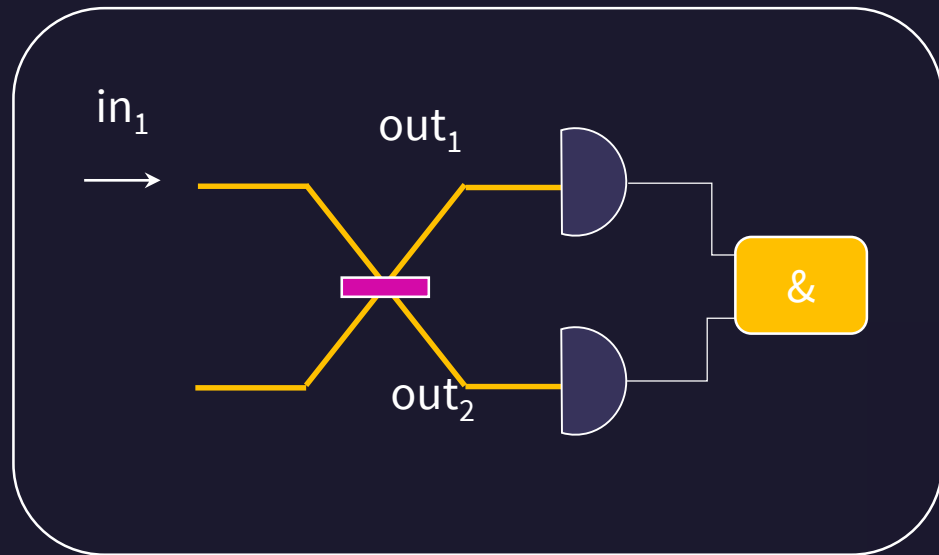
- Si cavity on 0.005% Er:LiNbO<sub>3</sub>
- T approx. 50 mK
- Observation of isolated photoluminescence lines off line-center
- 13-fold(144-fold) reduction of decay constant from 1.8 ms to 134  $\mu$ s (12.5  $\mu$ s)
- $T_1 < 10 \mu$ sec and radiatively limited emission ( $T_1 = T_2/2$ ) seem possible  
-> Fourier-limited photons

*Joint work with Gröblacher group, TU Delft*



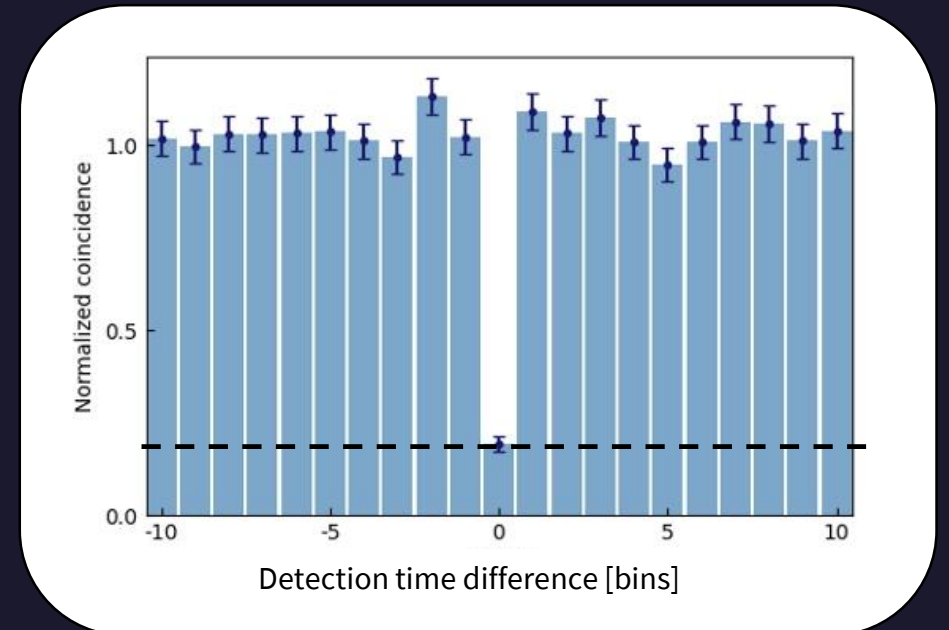
# Purcell-enhanced emission

- Measurement of auto-correlation coefficient shows non-classical (single-photon) nature of emissions and confirms interaction with individual erbium ions
  - -> single-photon source and possibility for qubit read-out



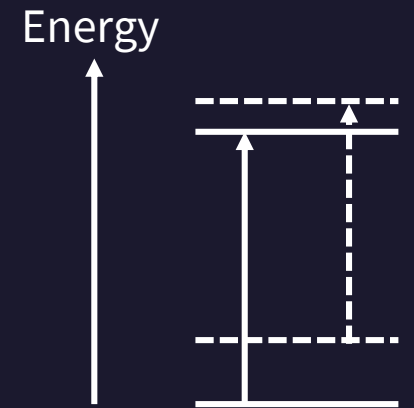
Experimental setup

$$g^2(0) = 0.190 \pm 0.02$$

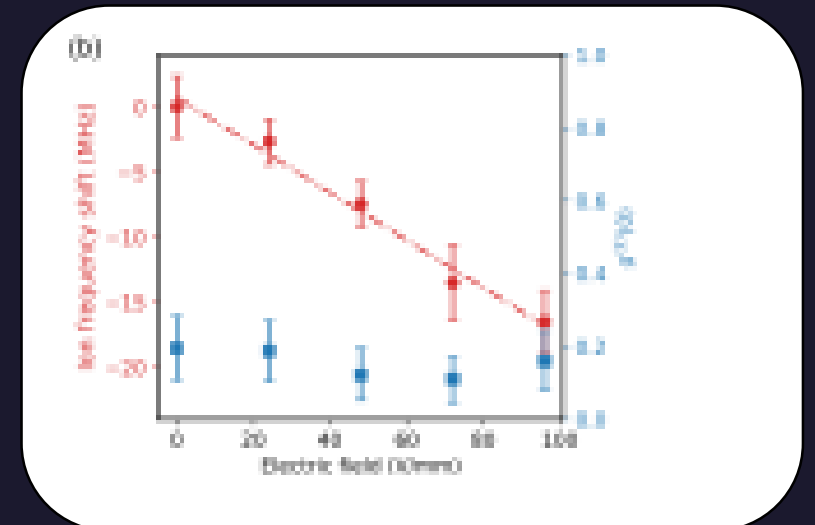
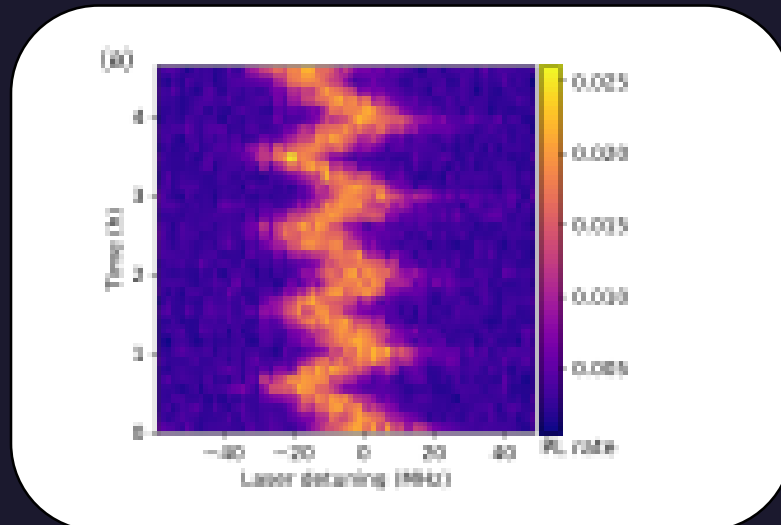


# Purcell-enhanced emission

- Demonstration of Stark tunability of single ion
- Single-photon character not affected
- Feedback mechanism to counter spectral diffusion
  - > indistinguishable single photons
  - > distant spin-spin entanglement
  - > heralded entangled photon pairs

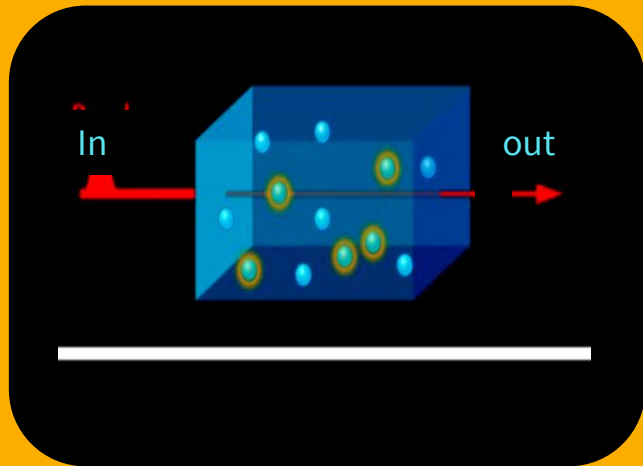


$$\Delta\nu = (\vec{d}_e - \vec{d}_g) \vec{E} / h$$

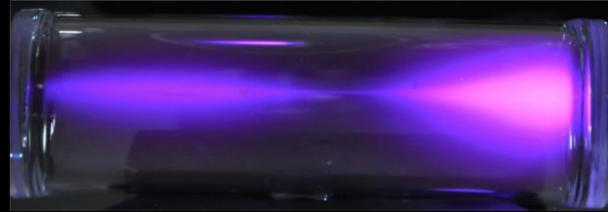




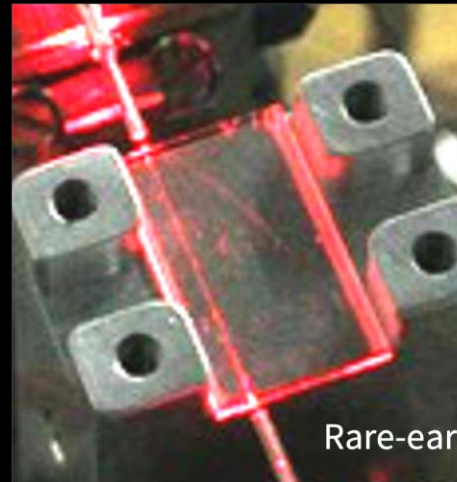
How to store photonic quantum states in a multiplexed manner?  
Use large ensembles of atoms



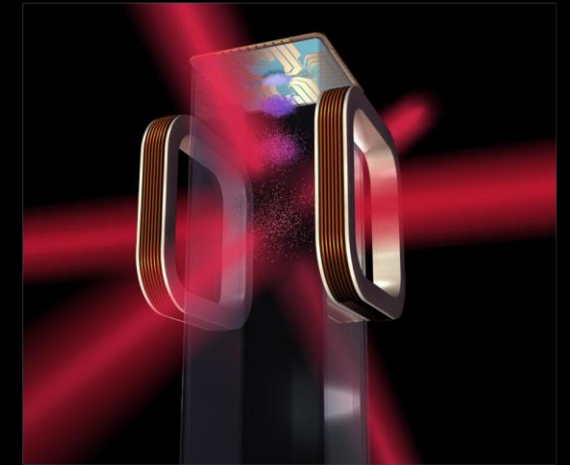
$$|\psi'\rangle = \mathbb{1}|\psi\rangle$$



Room-temperature vapor

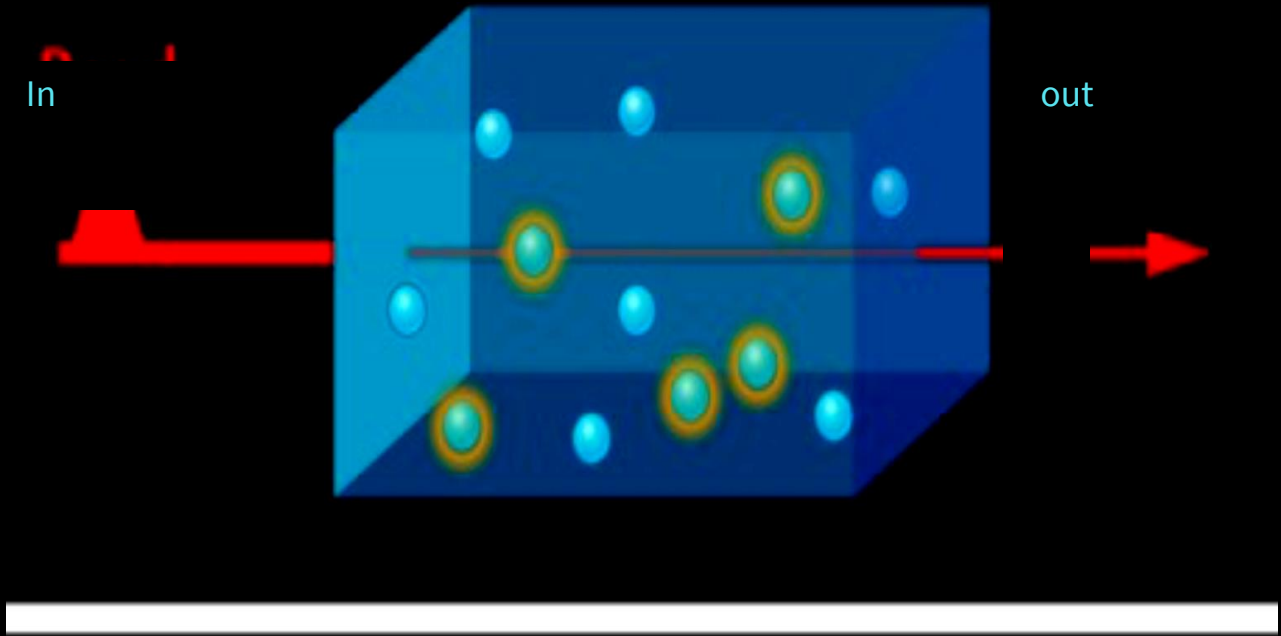


Rare-earth crystals



Laser-cooled atoms

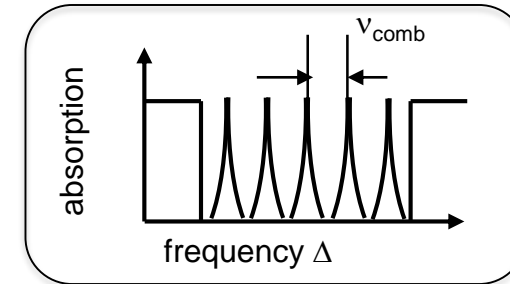
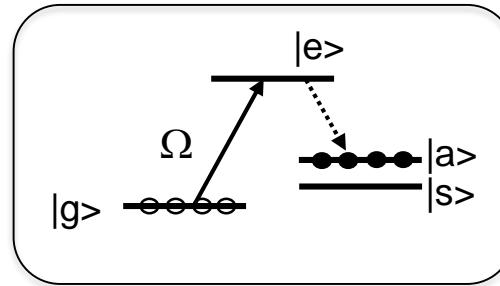
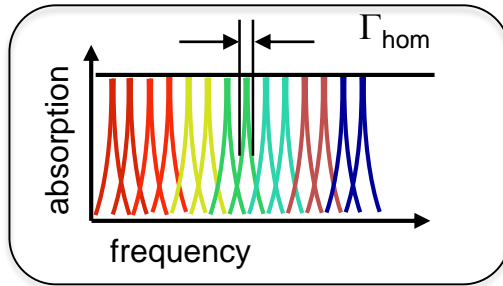
How to store photonic quantum states in a multiplexed manner?  
Use large ensembles of atoms



$$|\psi'\rangle = \mathbb{1}|\psi\rangle$$

# Photon echo quantum memory (AFC)

## 1. Preparation of an atomic frequency comb



## 2. Absorption of a photon -> fast dephasing

$$|y\rangle = \frac{1}{\sqrt{N}} \sum_{j=1}^N c_j e^{-i2\rho D_j t} e^{ikz_j} |g_1 \dots e_j \dots g_N\rangle$$

*Experiments: Geneva, Lund, Paris, Calgary, Delft, Barcelona, Hefei, Caltech*

## 3. Rephasing at $t=1/\nu_{\text{comb}}$ : $2\pi\Delta_j t = 2\pi(n\nu_{\text{comb}})/\nu_{\text{comb}} = n 2\pi$

- > Re-emission of photonic qubits with unity efficiency\* and fidelity
- > Possibility to use control pulses for on-demand read-out (if needed)

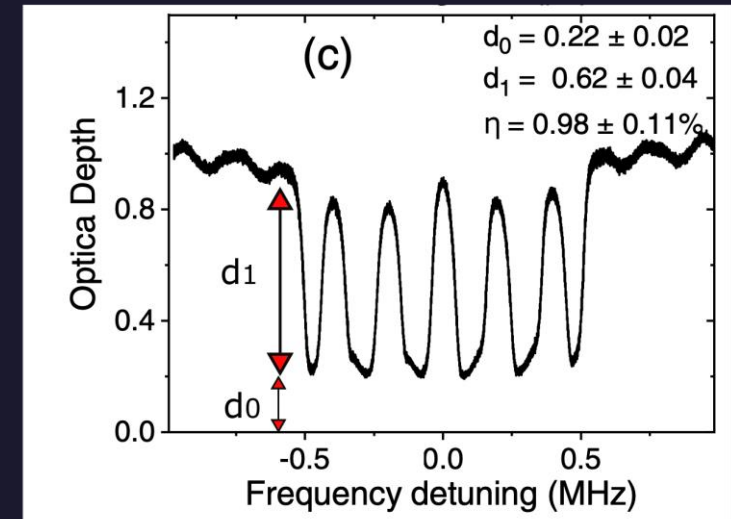
\* requires phase matching or cavity

**Needed:** inhomogeneously broadened transition, long-lived auxiliary state, narrow homogeneous linewidth

# Towards efficient quantum memory

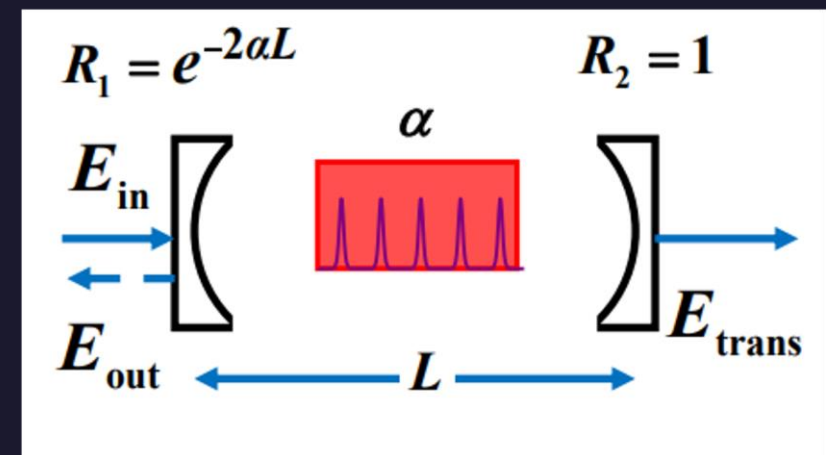
The efficiency of the AFC quantum memory is limited by its optical depth

$$\eta = \left(\frac{d_1}{F}\right)^2 e^{-d_1/F} e^{-d_0} e^{-7/F^2}$$



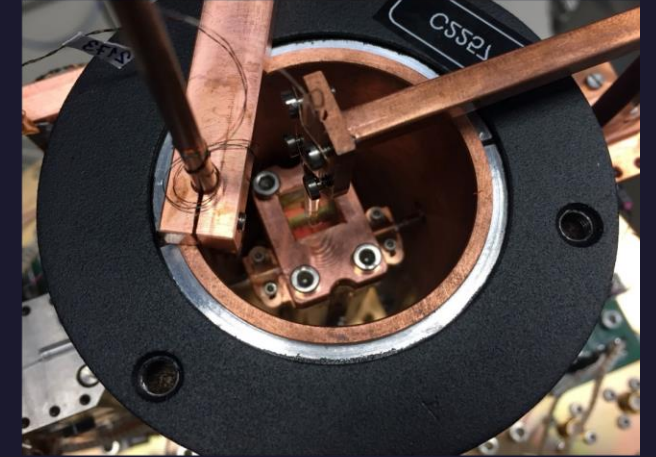
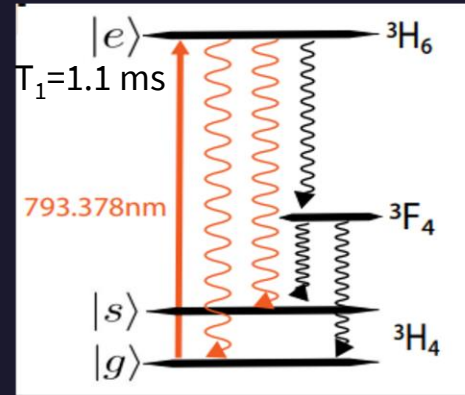
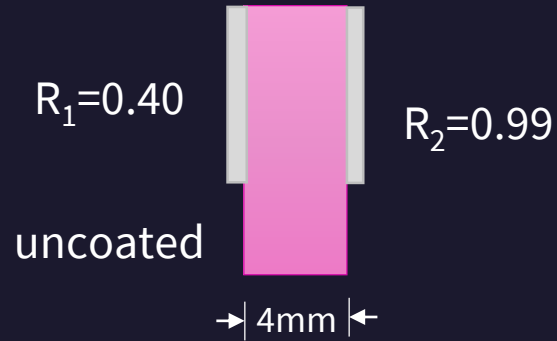
Using an impedance-matched cavity allows in principle to increase the efficiency to 1 despite small single-pass absorption  $e^{-\alpha l}$

Condition:  $R_1 = e^{-2\alpha l} R_2$  with  $R_2 = 1$

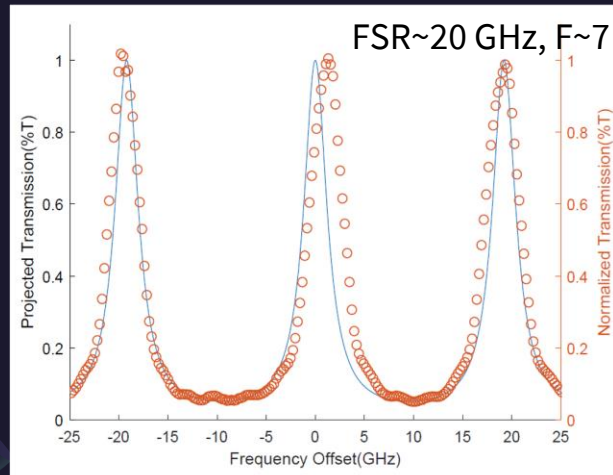


# Towards efficient quantum memory

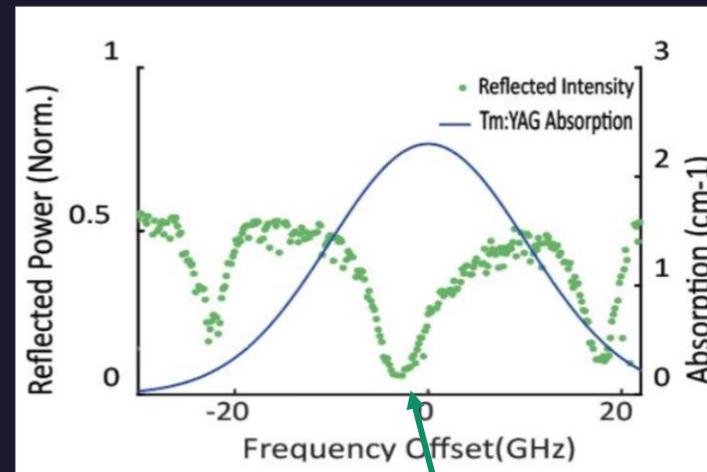
Reflection-coated Tm:Y<sub>3</sub>Al<sub>5</sub>O<sub>12</sub> (YAG) crystal



Cavity transmission spectrum  
(outside Tm resonance)

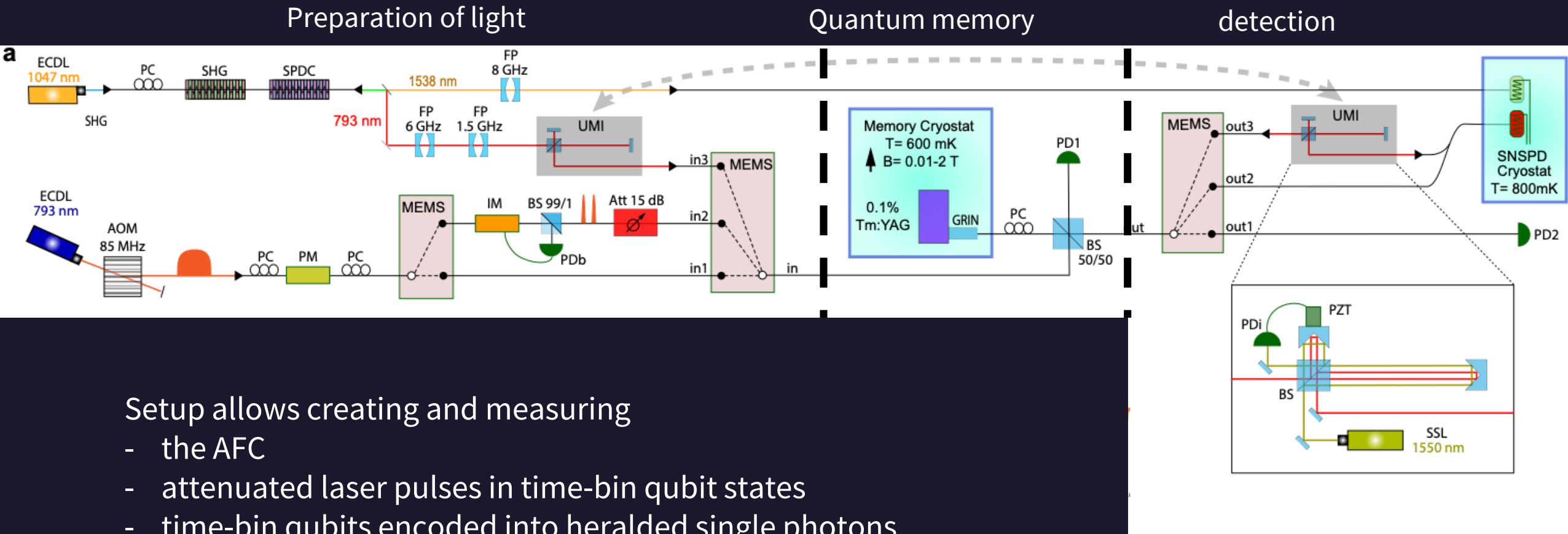


Cavity reflection spectrum  
(within Tm resonance)



Almost impedance-matched!

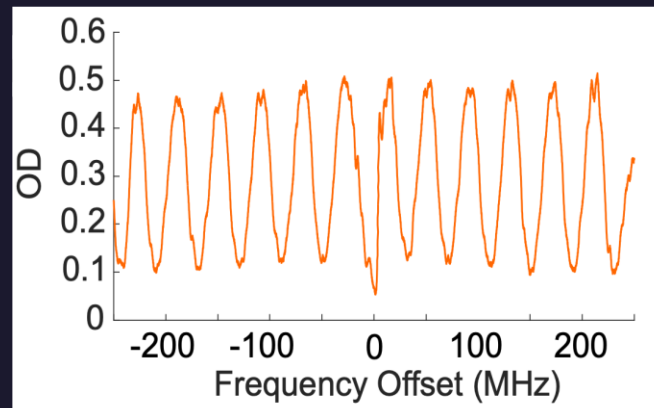
# Towards efficient quantum memory



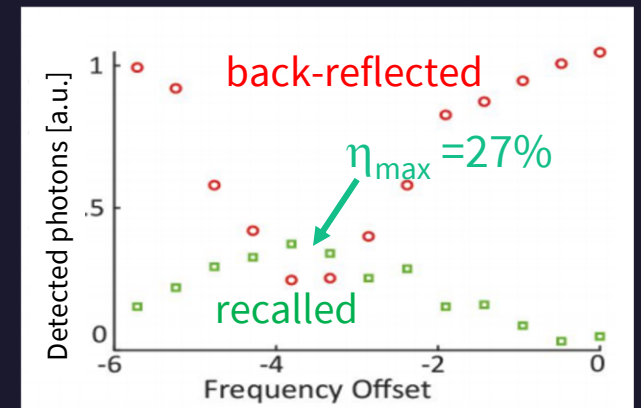
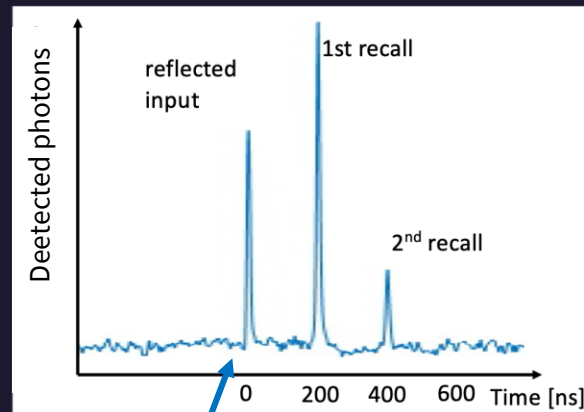
$$|\psi\rangle = \alpha|e\rangle + \beta e^{i\phi}|l\rangle$$

# Results

AFC created using the non-coated part of the crystal.  $\eta \sim 1\%$



AFC-based storage of attenuated laser pulses ( $\mu=0.7$ ) using coated part of crystal



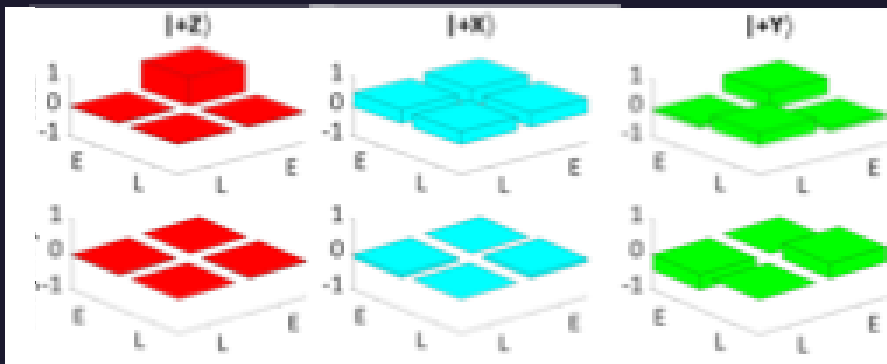
Imperfect impedance matching

# More Results

Measurement of **non-classical cross-correlations** with photon pairs before and after storage:

$$g^{(2)}_{\text{before}} = 61.8 \pm 3.8, g^{(2)}_{\text{after}} = 9.1 \pm 1.2$$

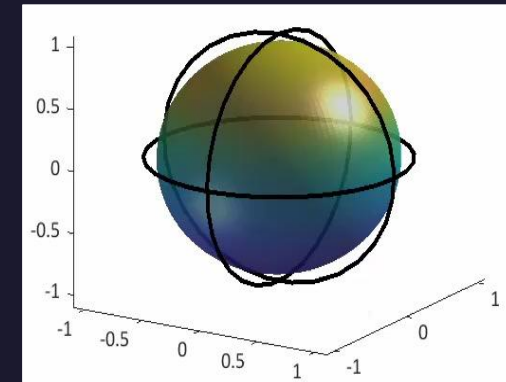
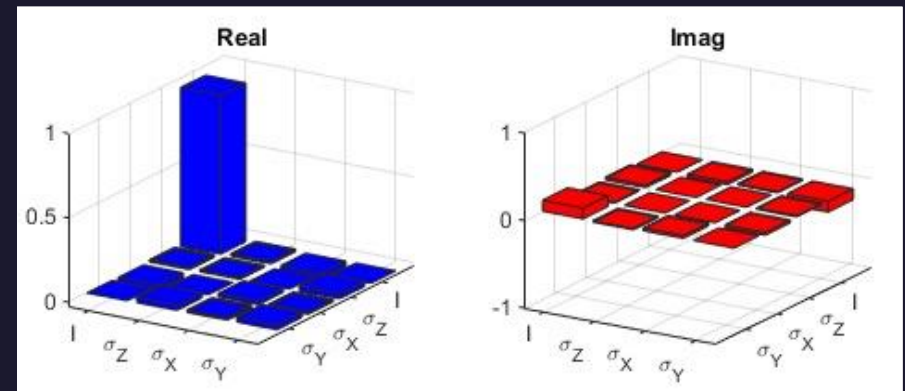
**Quantum state tomography** of time-bin qubits encoded into heralded single photons



->  $F = (85 \pm 1)\% \gg 0.66$

**Quantum process tomography** of time-bin qubits encoded into attenuate laser pulses ( $\mu=0.7$ )

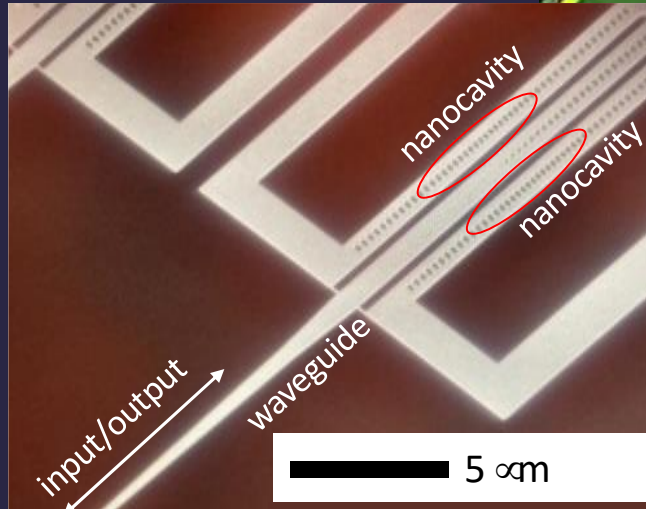
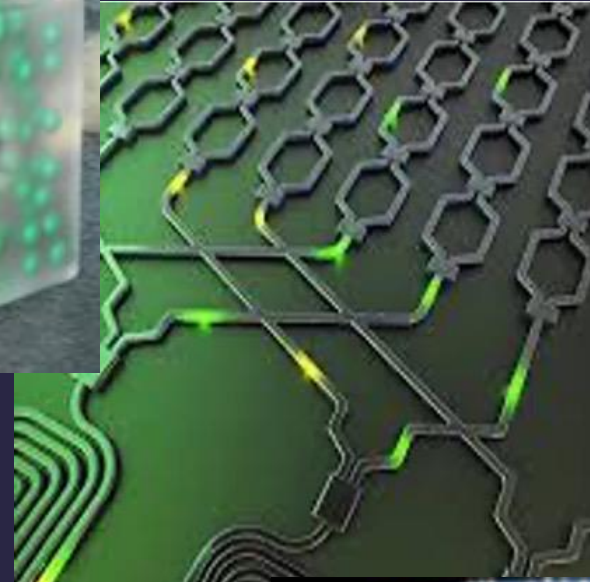
$$\rho_{out} = \sum_{m,n=0}^3 \chi_{mn} \sigma_m \rho_{in} \sigma_n^\dagger$$





# Towards fully quantum-enabled networks based on an integrated platform

- **Single (and soon? entangled) photons** based on Purcell-enhanced emission from single rare-earth ions
- Compatible **quantum memories** based on large ensembles of rare-earth ions
- **(Quantum computing nodes** using interacting rare-earth ions coupled to nano-cavities for readout)
- Exploit maturity of Si/SiN/LiNbO<sub>3</sub> photonics (foundries) to create **compatible, scalable** and **integrated quantum network technology**
- More **fundamental research into materials, protocols, applications,...**



# Thank you!

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# Quantum repeater - how to mitigate loss

**Goal:** Overcome the exponential scaling of photon transmission over a long (lossy) quantum channel

Note: multiplexing does not lead to better scaling

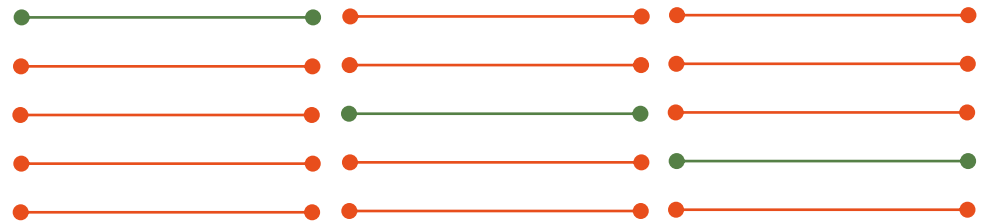
## Solution

- 1) Break long link into shorter *elementary links*.
- 2) Distribute *heralded and long-lived entanglement* across each elementary link.
- 3) Multiplex distribution (any degree of freedom) to make it efficient.
- 4) Mode mapping based on feed-forward info allows connecting “good” links using Bell-state measurements.

*Exponential scaling*



*Same scaling*

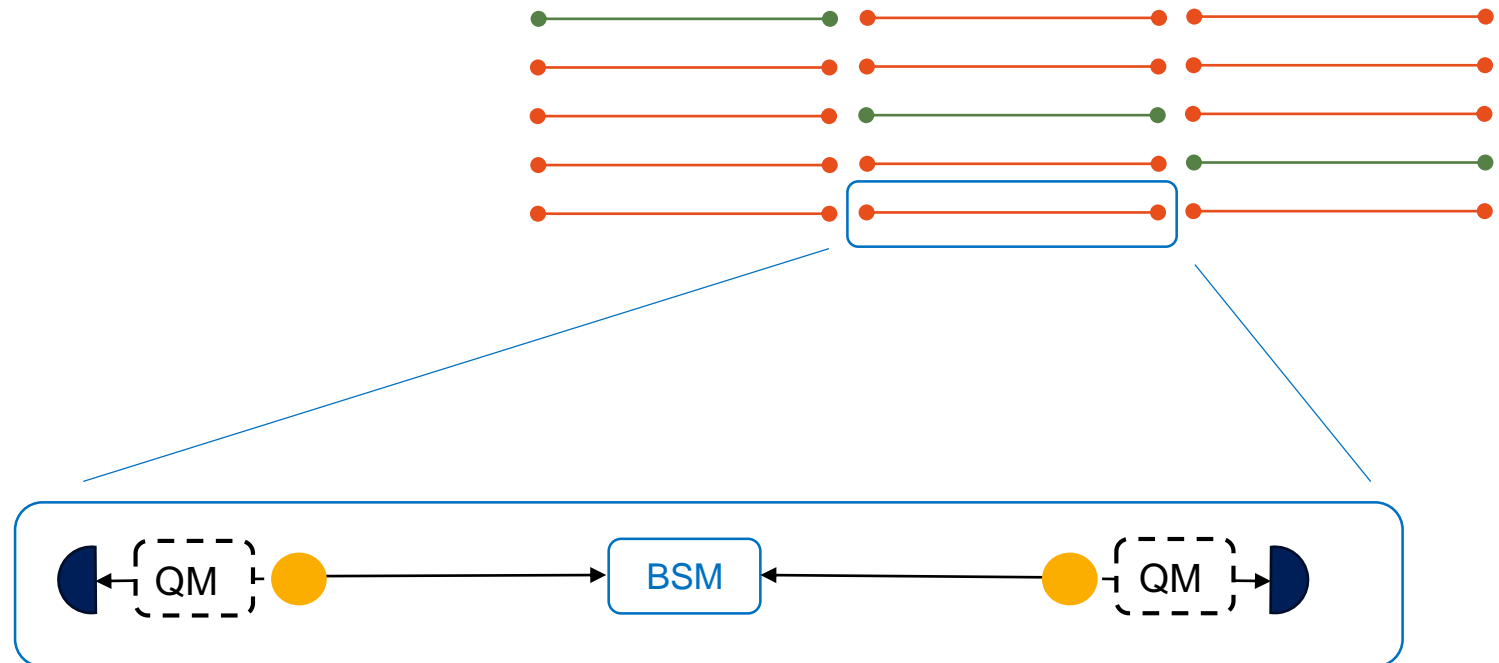


*Better scaling*

No need for photons to travel *in one go* over the entire link.

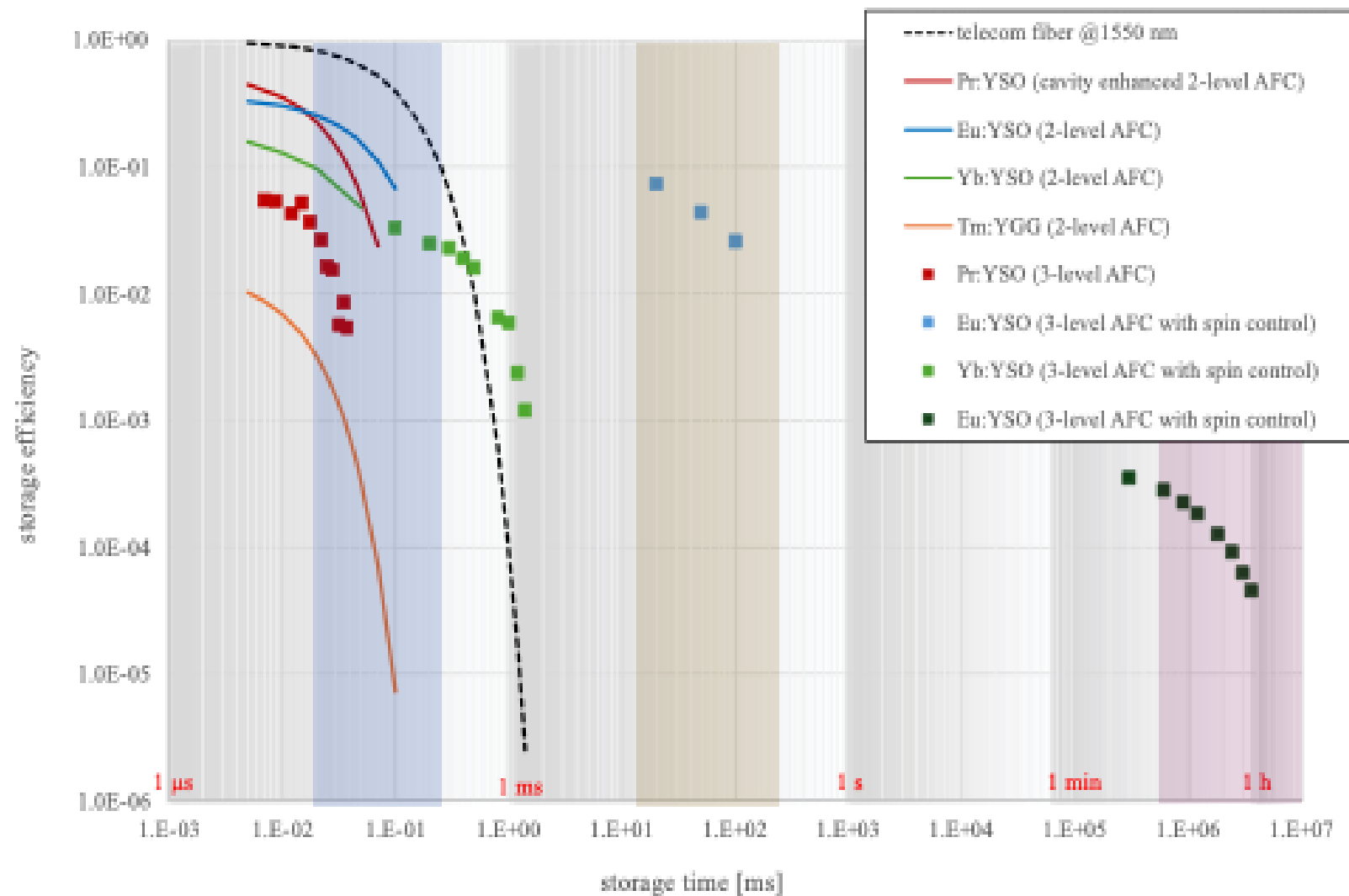
# Quantum memory requirements

- 1) Large storage efficiency
- 2) Sufficient storage time
- 3) Fidelity  $\rightarrow 1$
- 4) Feed-forward mode mapping
- 5) High multiplexing capacity
- 6) Wavelength of operation
- 7) Bandwidth per qubit
- 8) Integrability



# State-of-the-art

- **Comparing results taken under different conditions...** But while not all experiments demonstrate quantum nature, all used a quantum protocol
- **2-level AFC:** not yet better than fiber, but close. Materials with sufficient  $T_2^{\text{opt}}$  for  $\tau = 1\text{ms}$  exist. Need to reduce technical noise and add cavities.
- **AFC spin-storage:** scaling already better than fiber, but efficiencies still small. Materials with sufficient  $T_2^{\text{spin}}$  exist. Need to improve efficiency of  $\pi$ -pulses and noise and to add cavities.
- **Three use cases** (assuming sufficient multiplexing)
  - Quantum repeaters for fiber networks
  - Quantum repeaters for satellite networks
  - Physical qubit transport



Fiber-based networks

Satellite-based networks

Physical qubit transport-based networks