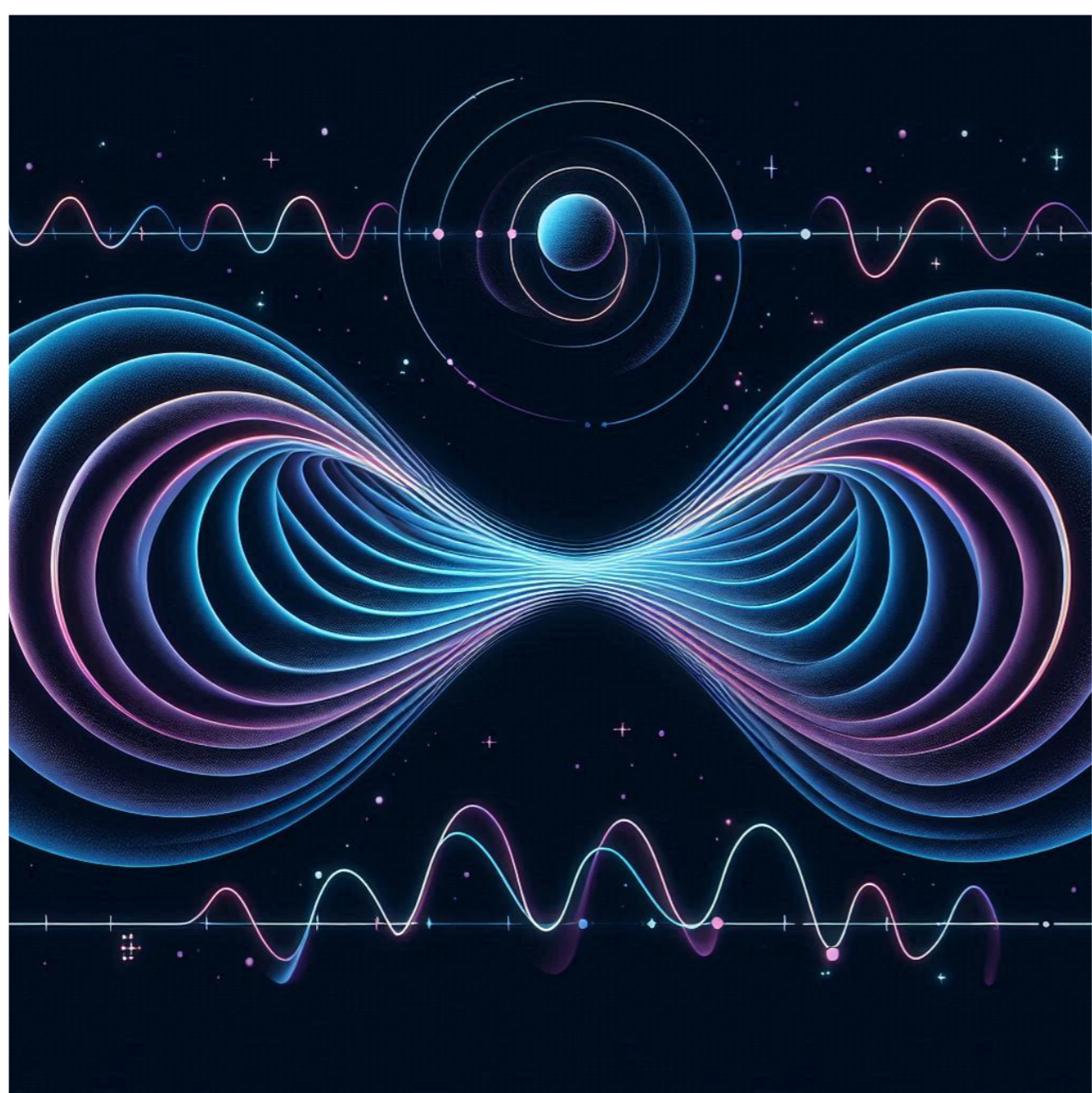


# FINITE PARTS OF INFLATIONARY LOOPS



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WORKSHOP  
“LOOPING IN  
THE PRIMORDIAL  
UNIVERSE”





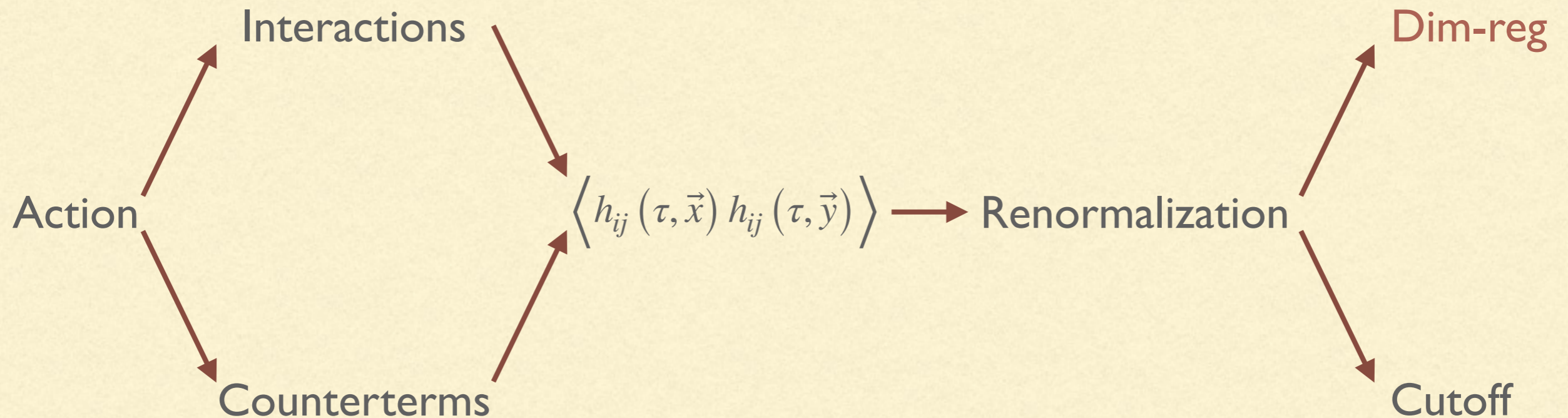
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# GOAL & MAP OF THE TALK

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📖 Goal: present our tools and procedure to compute loop corrections in inflationary cosmologies through **dimensional regularization**

📱 Application: one-loop power spectrum of tensor modes in interaction with scalar modes on dS background





# ACTION & VERTICES

Degrees of freedom ( $\delta\phi$ -gauge):  $\delta\phi, h_{ij}$

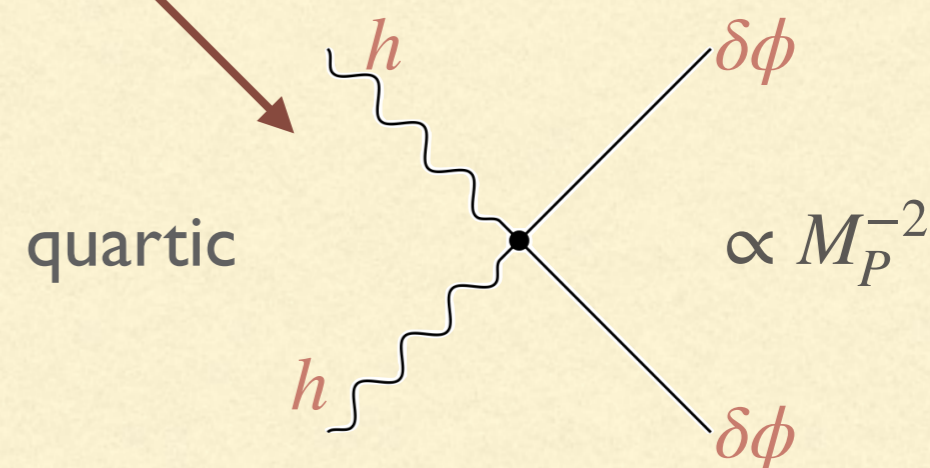
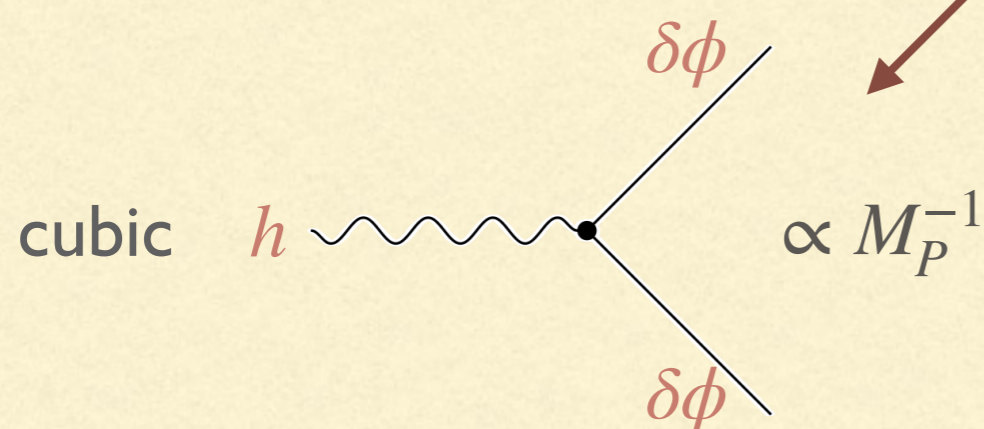
$$S = S_h + S_{\delta\phi} + S_{\text{int}}$$

dS background  
 $a(\tau) = -\frac{1}{H\tau}$

$$S_h = \frac{M_P^2}{8} \int d\tau d^3x a^2(\tau) \left( h'_{ij} h'_{ij} - \partial_k h_{ij} \partial_k h_{ij} \right)$$

$$S_{\delta\phi} = \frac{1}{2} \int d\tau d^3x a^2(\tau) \left[ (\delta\phi')^2 - \partial_i \delta\phi \partial_i \delta\phi \right]$$

$$S_{\text{int}} = \int d\tau d^3x a^2(\tau) \left( \frac{1}{2} h_{ij} \partial_i \delta\phi \partial_j \delta\phi - \frac{1}{4} h_{ik} h_{kj} \partial_i \delta\phi \partial_j \delta\phi \right)$$





# CORRELATION FUNCTION & POWER SPECTRUM

$$\left\langle h_{ij}(\tau, \vec{x}) h_{kl}(\tau, \vec{y}) \right\rangle \stackrel{\text{in-in}}{\downarrow} = \langle 0 | F^{-1}(\tau, -\infty_+) h_{ij}^I(\tau, \vec{x}) h_{kl}^I(\tau, \vec{y}) F(\tau, -\infty_-) | 0 \rangle \Big|_{\text{no bubbles}}$$

$$F(\tau, \tau_0) = T \exp \left[ -i \int_{\tau_0}^{\tau} d\tau' H_I(\tau') \right]$$

fields in the interaction picture

$$\delta\phi^I(\tau, \vec{x}) = \int \frac{d^3k}{(2\pi)^{3/2}} e^{i\vec{k}\cdot\vec{x}} \left[ \delta\phi_k(\tau) a_{\vec{k}} + \delta\phi_k^*(\tau) a_{-\vec{k}}^\dagger \right]$$

$$h_{ij}^I(\tau, \vec{x}) = \int \frac{d^3k}{(2\pi)^{3/2}} e^{i\vec{k}\cdot\vec{x}} \sum_{\lambda=+, \times} \epsilon_{ij}^\lambda(\vec{k}) \left[ h_k(\tau) a_{\lambda, \vec{k}} + h_k^*(\tau) a_{\lambda, -\vec{k}}^\dagger \right]$$

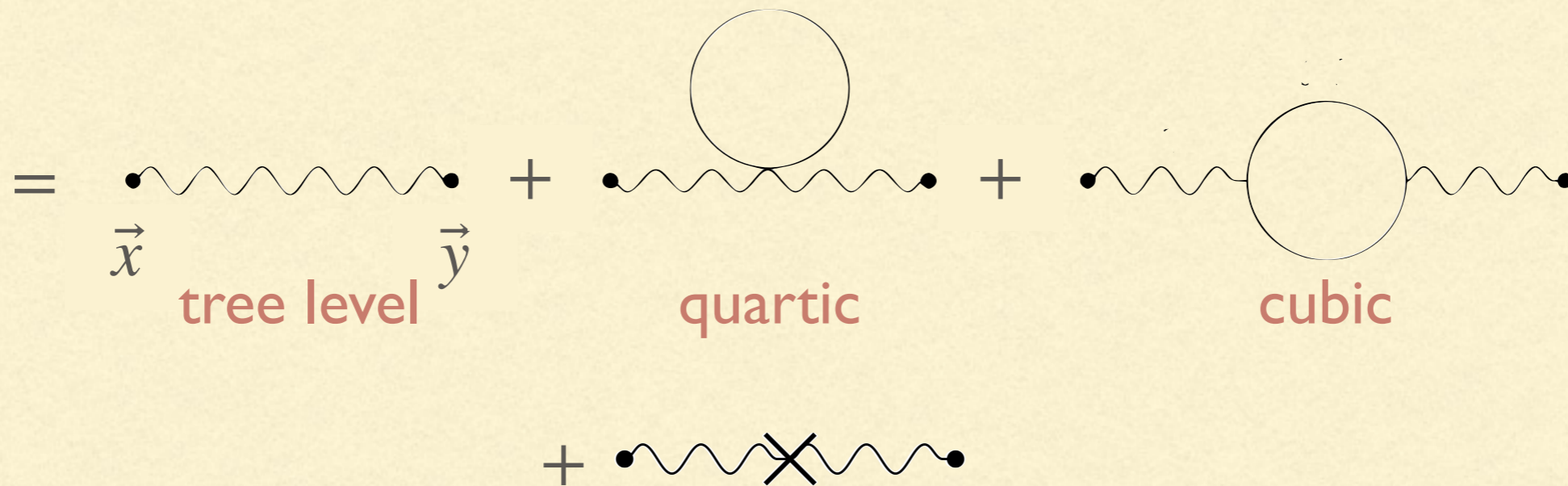


# CORRELATION FUNCTION & POWER SPECTRUM

$$\langle h_{ij}(\tau, \vec{x}) h_{ij}(\tau, \vec{y}) \rangle = \int \frac{d^3k}{(2\pi)^3} e^{i\vec{k}\cdot(\vec{x}-\vec{y})} \frac{2\pi^2}{k^3} 2\mathcal{P}_h(\tau, k)$$

power spectrum

polarizations



how do counterterms look like?



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# COUNTERTERMS

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$$S_{\text{HDO}} = \int d^4x \sqrt{-g} \left( A_1 R^2 + A_2 R_{\mu\nu} R^{\mu\nu} + A_3 R_{\mu\nu\rho\sigma} R^{\mu\nu\rho\sigma} + A_4 \square R \right)$$

$$S_{\text{cts}} = - \int d\tau d^3x a^2 \left[ C_1 H^2 h'_{ij} h'_{ij} + C_2 H^2 \partial_k h_{ij} \partial_k h_{ij} + \frac{C_3}{a^2} \left( \partial_k h'_{ij} \partial_k h'_{ij} - \partial^2 h_{ij} \partial^2 h_{ij} \right) \right]$$



Other operators break gauge invariance  
(discuss them with cutoff)

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# RENORMALIZATION

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Regularization



Redefinition



Subtraction



Regularization  $\rightarrow 0$

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# DIMENSIONAL REGULARIZATION: MODE FUNCTIONS

$$d = 3 + \delta$$

Not present on flat space!



Equations of motion  
(scalar)

$$\begin{cases} \delta\phi_k'' + (2 + \delta)aH \delta\phi_k' + k^2\delta\phi_k = 0 \\ \delta\phi_k(\tau) \xrightarrow{k\tau \rightarrow -\infty} \frac{1}{a^{1+\delta/2}(\tau)\mu^{\delta/2}} \frac{e^{-ik\tau}}{\sqrt{2k}} \end{cases}$$

$\exists$  analytical solution, but we solve it perturbatively:

$$\delta\phi_k(\tau) = \delta\phi_k^{(0)}(\tau) + \delta \delta\phi_k^{(1)}(\tau) + \mathcal{O}(\delta^2)$$

$$\delta\phi_k(\tau) = i \frac{H}{\sqrt{2k^3}} e^{-ik\tau} (1 + ik\tau) \left\{ 1 + \frac{\delta}{2} \left[ \frac{1 - ik\tau}{1 + ik\tau} e^{2ik\tau} (i\pi - \text{Ei}(-2ik\tau)) + \frac{2}{1 + ik\tau} + \log\left(-\frac{H\tau}{\mu}\right) \right] + \mathcal{O}(\delta^2) \right\}$$



# DIMENSIONAL REGULARIZATION: INTEGRALS

$$I(\delta) = \int_0^\infty dp p^\delta f(\delta, p) = \underbrace{\int_0^L dp p^\delta f(\delta, p)}_{\int_0^L dp f(0, p) + \mathcal{O}(\delta)} + \int_L^\infty dp p^\delta f(\delta, p)$$

from the d-dimensional measure

$\underbrace{\sum_{n=-\infty}^N c_n(\delta) p^n}$

Logarithmic divergence  
as single pole to subtract

$$I(\delta) = -\frac{c_{-1}(0)}{\delta} - \left. \frac{dc_{-1}}{d\delta} \right|_{\delta=0} + \lim_{L \rightarrow \infty} \left[ \int_0^L dp f(0, p) - c_{-1}(0) \log L - \sum_{n=0}^N \frac{L^{n+1}}{n+1} c_n(0) \right] + \mathcal{O}(\delta)$$

Only here we need  $\mathcal{O}(\delta)$  corrections to the mode functions



# RENORMALIZED TENSOR POWER SPECTRUM

$$\mathcal{P}_h(\tau, k) \Big|_{k\tau=y} = \alpha(1 + y^2) + \frac{3}{80}\alpha^2 \left\{ c_1 + c_2 y^2 + c_3 y^4 + (1 + y^2) \log \left( -\frac{H}{\mu} y \right) + \operatorname{Re} \left[ e^{2iy} (i + y)^2 (\operatorname{Ei}(-2iy) - i\pi) \right] \right\} + \mathcal{O}(\alpha^3)$$

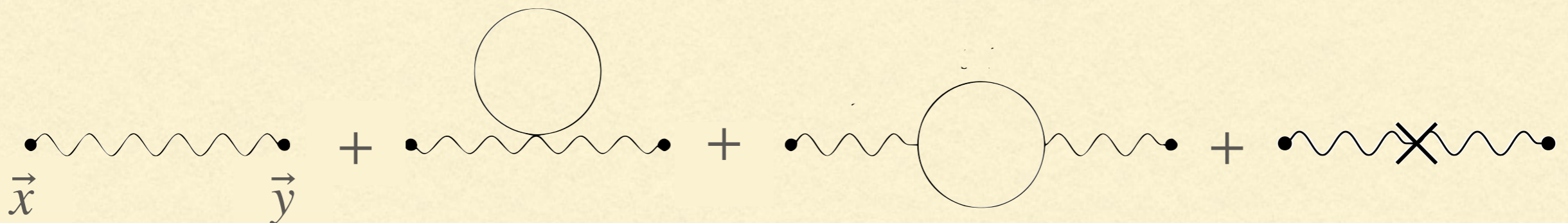
expansion  
parameter

tree level

$$\alpha = \left( \frac{H}{\pi M_P} \right)^2$$

$y \rightarrow 0$

$$\mathcal{P}_h(0, k) = \alpha + \frac{3}{80}\alpha^2 \left( c_1 - \gamma_E + \log \frac{H}{2\mu} \right) + \mathcal{O}(\alpha^3)$$





# CUTOFF: QED

$$\begin{aligned}
 \Pi^{\mu\nu}(q) &= \text{Diagram} \stackrel{\text{WT}}{=} \Pi_2(q^2) (q^\mu q^\nu - q^2 \eta^{\mu\nu}) \\
 &= \Pi_2(q^2) (q^\mu q^\nu - q^2 \eta^{\mu\nu}) + A(q^2) \eta^{\mu\nu}
 \end{aligned}$$

↑
cuttoff

**dim-reg**  $\mathcal{L}_{cts} = -\frac{1}{4} (Z_3 - 1) F_{\mu\nu} F^{\mu\nu}$

**cuttoff**  $\mathcal{L}_{cts} = \frac{1}{2} \left( M^2 A_\mu A^\mu + c \partial_\mu A_\nu \partial^\mu A^\nu + c' \partial_\mu A_\nu \partial^\nu A^\mu \right)$

~~WT~~  
~~U(1)~~



# CUTOFF: TENSOR POWER SPECTRUM

More freedom in the counterterms

$$S_{\text{cts}} = \int d\tau d^3x a^4 \left[ C_1 H^4 h_{ij} h_{ij} + \frac{H^2}{a^2} \left( C_2 h'_{ij} h'_{ij} + C_3 \partial_k h_{ij} \partial_k h_{ij} \right) + \frac{1}{a^4} \left( C_4 \partial^2 h_{ij} \partial^2 h_{ij} + C_5 \partial_k h'_{ij} \partial_k h'_{ij} \right) \right]$$



$$\mathcal{P}_h(\tau, k) \Big|_{k\tau=y} = \alpha(1 + y^2) + \frac{3}{80} \alpha^2 \left\{ c_1 + c_2 y^2 + c_3 y^4 + (1 + y^2) \log \left( -\frac{H}{\mu} y \right) + c_4 \text{Re} \left[ e^{2iy} (i + y)^2 (\text{Ei}(-2iy) - i\pi) \right] \right\} + \mathcal{O}(\alpha^3)$$

You must impose by hand the cancellation of the late-time divergence choosing  $c_4 = 1$

$$\alpha = \left( \frac{H}{\pi M_P} \right)^2$$



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# CONCLUSIONS

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- Importance of loops in inflation  $\Rightarrow$  Renormalization
- Procedure for dimensional regularization  $I(\delta)$
- Cutoff: Analogy with QED (HIC SUNT LEONES)
- GWs in inflation: computation of the one-loop PS, familiar  
 $\propto \log \left( \frac{H}{\mu} \right)$ , computation up to finite parts