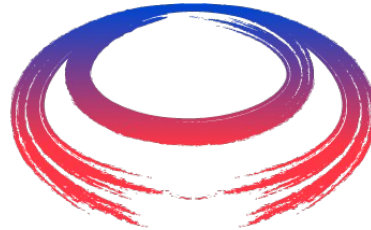




# Rectilinear Cooling DA

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 International  
UON Collider  
Collaboration

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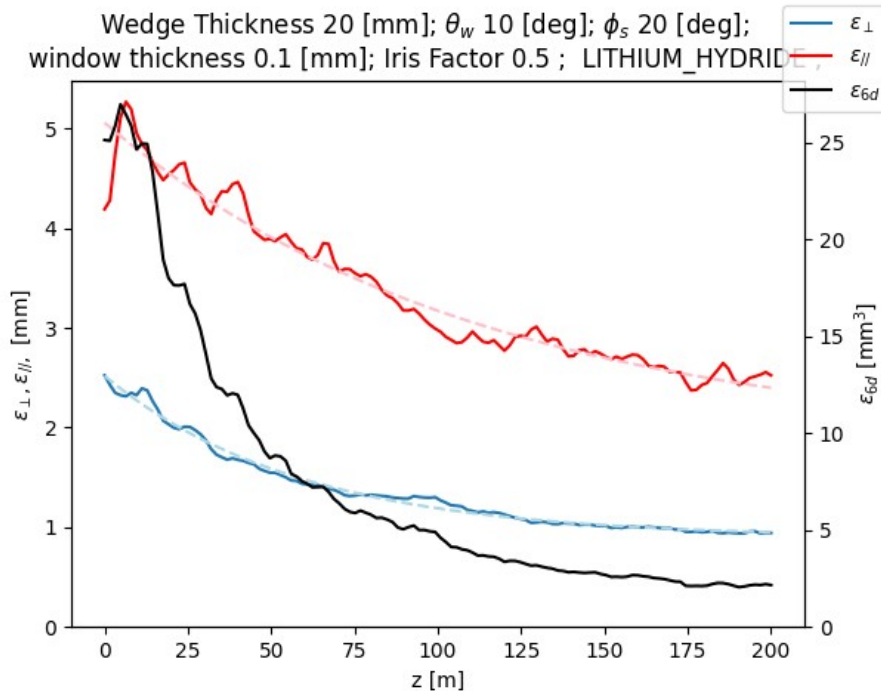
Science & Technology Facilities Council

**ISIS**

# DA and cooling optimisation

- Rectilinear cooling design is approaching a release version
  - Thanks to Ruihu!
  - Ruihu's design now frozen → publication
- Performance is improved over MAP
  - Ongoing discussion about cost/etc
- **But** still want higher efficiency
  - Would like ~ factor 2 more muons from the production system
  - Needed to make baseline parameters
  - Luminosity goes with  $N^2$
- Transmission of all rectilinear cooling system designs is rather low
  - Why?
  - Can we do better?

# Ionisation cooling



- Emittance change goes as

$$\epsilon_n(z) = \left[ \epsilon_n(0) - \epsilon_n^{eqm} \right] \exp\left(\frac{-z}{L_{cool}}\right) + \epsilon_n^{eqm}$$

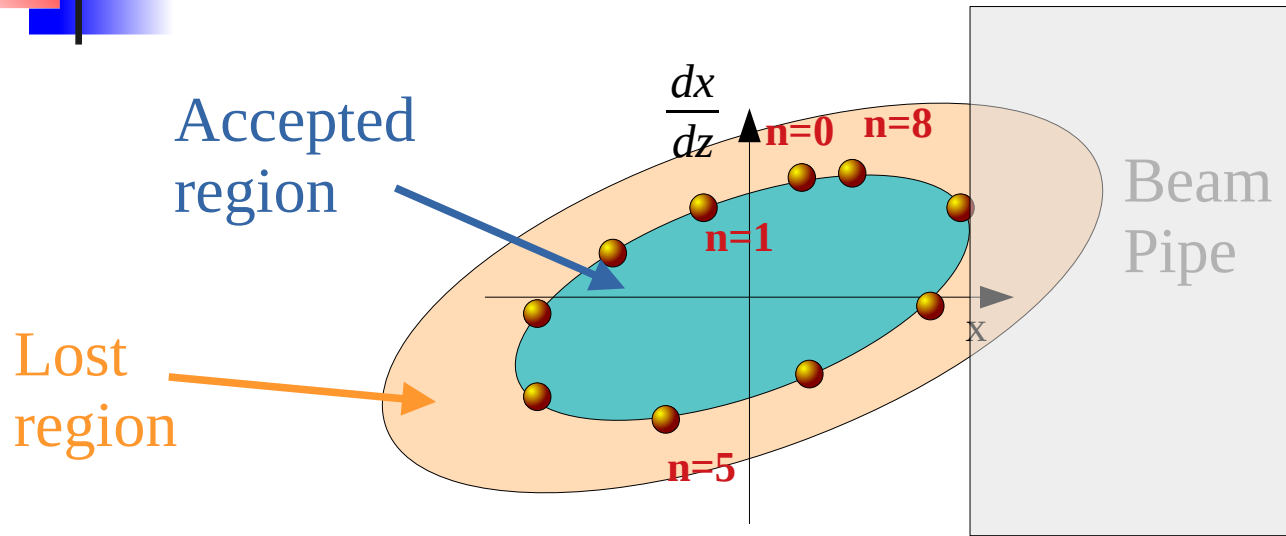
- $L_{cool}$  is cooling length, characteristic of the channel
- $\epsilon_n^{eqm}$  is equilibrium emittance, characteristic of the channel
- Quick (cheap) cooling  $\rightarrow \epsilon_n(0) \gg \epsilon_n^{eqm}$

# Acceptance

- We want  $\varepsilon_n(0)$  as large as possible
- The maximum emittance is determined by the *acceptance* of the cooling system
  - Maximum emittance particle that makes it through the cooling cells
- Two classes of acceptance (aperture)
  - Physical acceptance - equipment intercepts the beam
  - Dynamical acceptance - aberrations in the focusing at large emittance cause particles to get lost



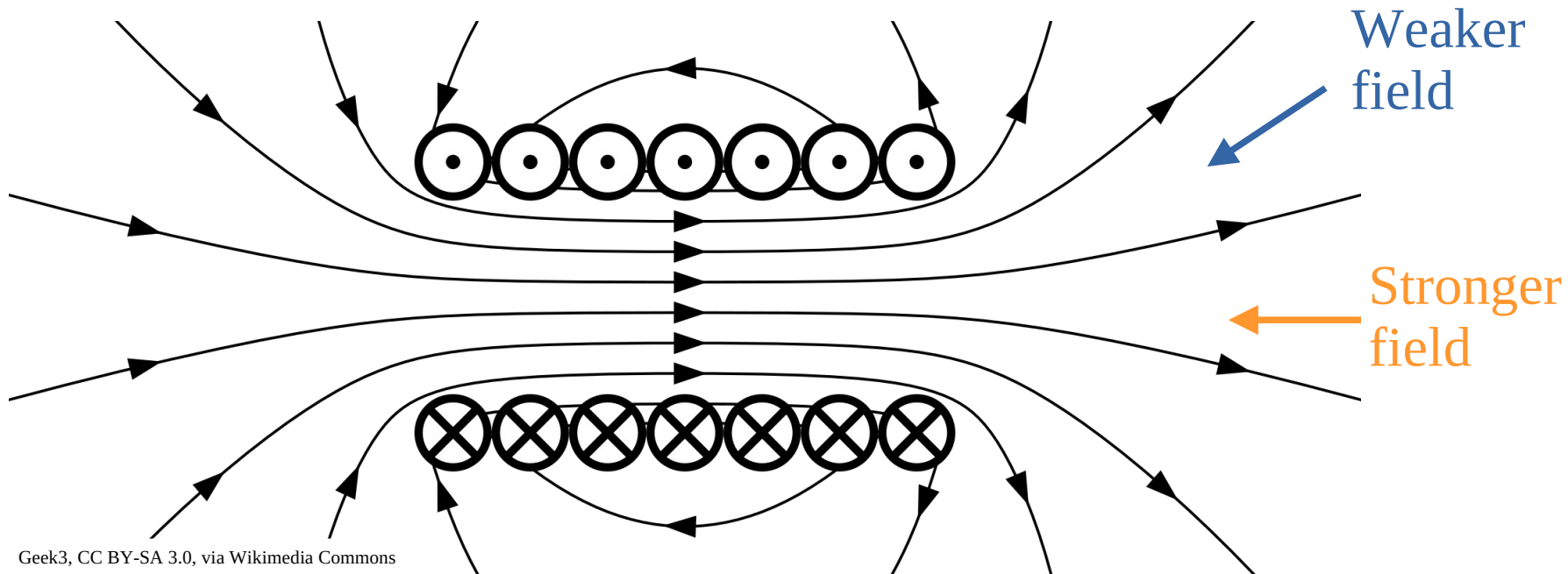
# Physical Acceptance



- Particles undergoing focusing follow elliptical trajectories in (position, momentum) space
- Beam pipe/aperture → maximum position transmitted
- Trajectories that touch the beam pipe are lost
- Maximum emittance trajectories that do not touch the beam pipe → “Acceptance”

# Dynamical Acceptance

- Focusing of solenoids varies across the magnet face



- Variation is only in the fringe of the solenoid
- Variation → determined purely in terms of the field on axis

# Dynamical Acceptance

- Consider Maxwell's equations

$$\nabla \cdot B = 0$$

$$\nabla \times B = 0$$

- For a generalised solenoid symmetry

$$\nabla \cdot B = \frac{1}{r} \partial_r (r B_r) + \partial_z B_z$$

$$= \partial_r B_r + \frac{B_r}{r} + \partial_z B_z$$

$$\nabla \times B = (\partial_z B_r - \partial_r B_z) \hat{\phi} = 0$$

- We can solve the PDE to get a generalised expression for B



# Dynamical Acceptance

- Consider Maxwell's equations (magnetostatic)

$$\nabla \cdot B = 0$$

$$\nabla \times B = 0$$

- For a generalised solenoid symmetry

$$\nabla \cdot B = \frac{1}{r} \partial_r (r B_r) + \partial_z B_z$$

$$= \partial_r B_r + \frac{B_r}{r} + \partial_z B_z$$

$$\nabla \times B = (\partial_z B_r - \partial_r B_z) \hat{\phi} = 0$$

- We can solve the PDE to get a generalised expression for B





# Dynamical Acceptance

- Assume generalised Taylor series for the field

$$B_r = \sum_{i=0} a_i(z) r^i$$

$$B_z = \sum_{i=0} b_i(z) r^i$$

- Substitute into Maxwell and solve

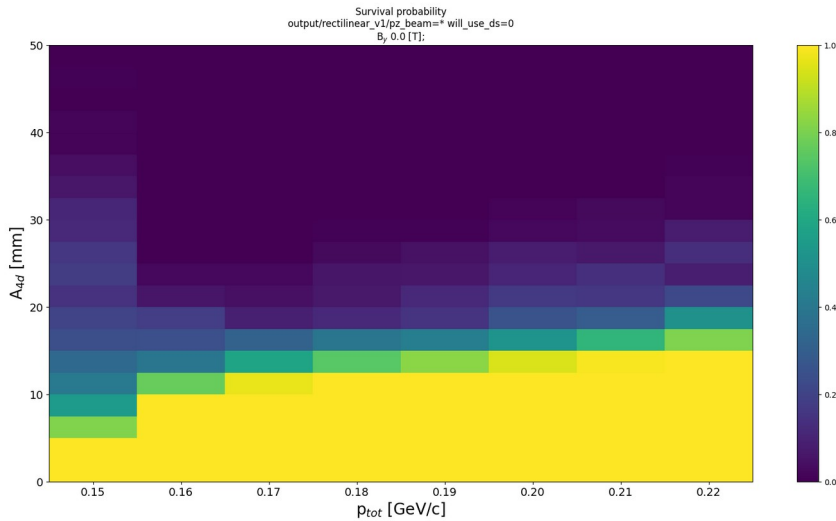
$$b_{i+2}(z) = -\frac{1}{(i+2)^2} \partial_z^2 b_i(z)$$

$$a_i = -\frac{1}{i+1} \partial_z b_{i-1}(z)$$

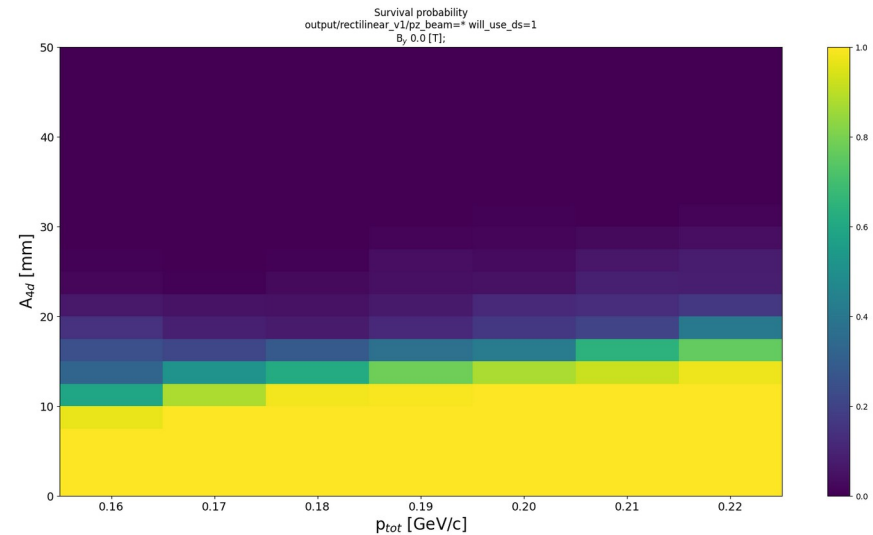
- Off-axis behaviour of solenoid is unique function of the field on-axis**
  - Particular coil arrangement is **irrelevant** for beam dynamics
  - Of course it is important for many other reasons!

# Example - Demo Lattice

- Consider Dynamic Aperture of demo lattice
  - Implementing “derivatives solenoid” into G4BL
  - Use  $B = h_0 \sin(kz) + h_1 \sin(2kz) + h_2 \sin(3kz)$



2024-05-24-release  
Solenoids only



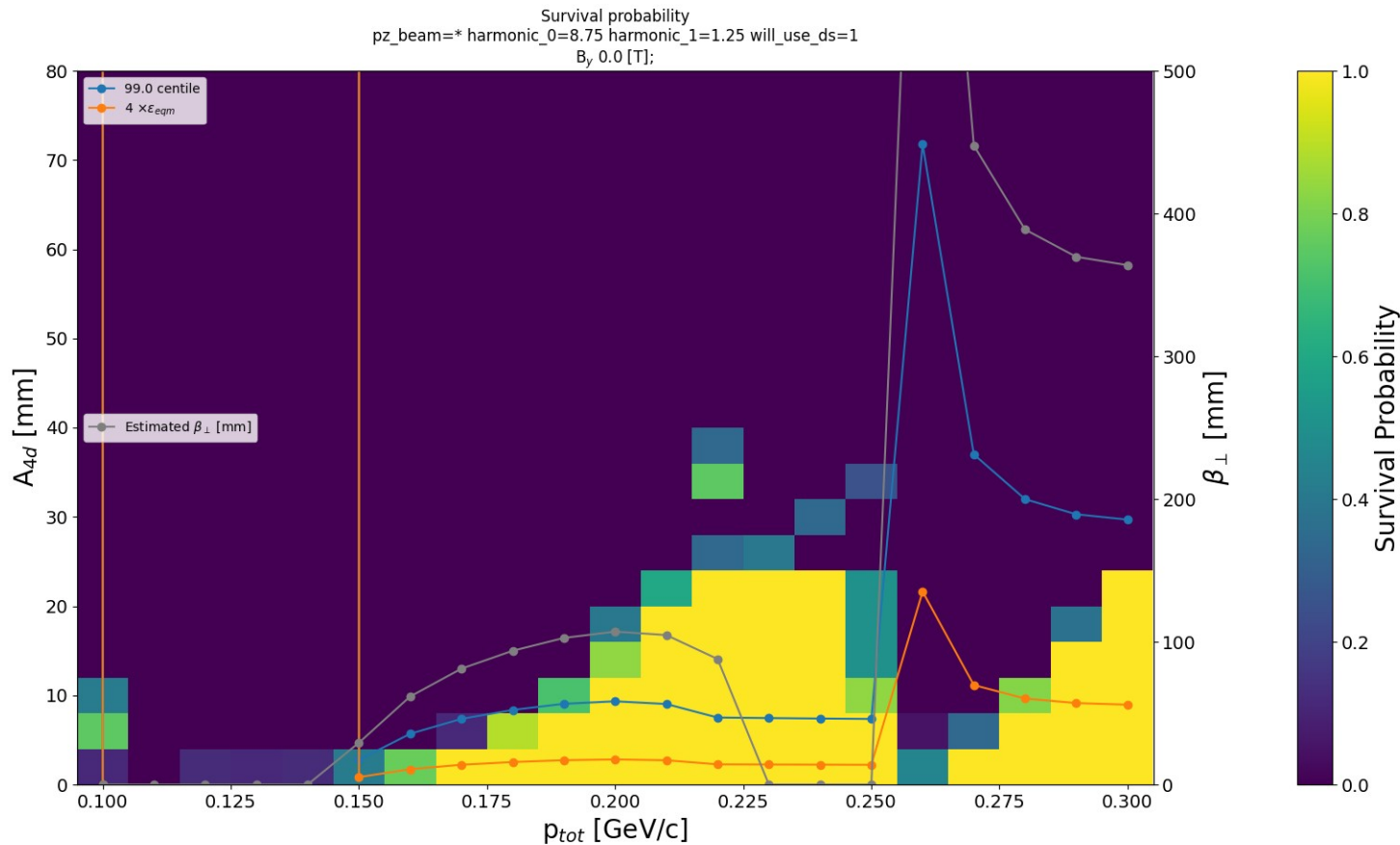
Field expansion  
 $h_0 = 8.75$  T  
 $h_1 = 1.25$  T  
Truncate at  $r^9$

# Nb Amplitude distribution

- For a beam that follows a gaussian distribution...
- 4D Amplitude follows a  $\chi^2$  distribution with 4 degrees of freedom
- By comparing DA with emittance we can estimate number of muons outside acceptance
  - Assuming gaussian beam
  - (For ionisation cooling, expect a bit more tail than in a Gaussian)
- By comparing DA with equilibrium emittance, we can develop “figure of merit” for cooling lattice
  - “Dynamic range” of a given cooling lattice

# Demo - zoom out

- Interesting features if we zoom out
  - Acceptance even in the stop band
  - Here I assume beta for 200 MeV/c (reference) trajectory to calculate amplitude



# Back to cooling...

- So what about cooling?

Stratakis et al, PRAB 18 031003, 2015

Stage	$\epsilon_T^{\text{sim}}$ [mm]	$\epsilon_L^{\text{sim}}$ [mm]	$P_z^{\text{sim}}$ [MeV/c]	$T$ [%]
Begin	17.00	46.00	255	
A1	6.28	14.48	238	70.6
A2	3.40	4.64	229	87.5
A3	2.07	2.60	220	88.8
A4	1.48	2.35	215	94.6

51.9 %

Zhu et al, in progress

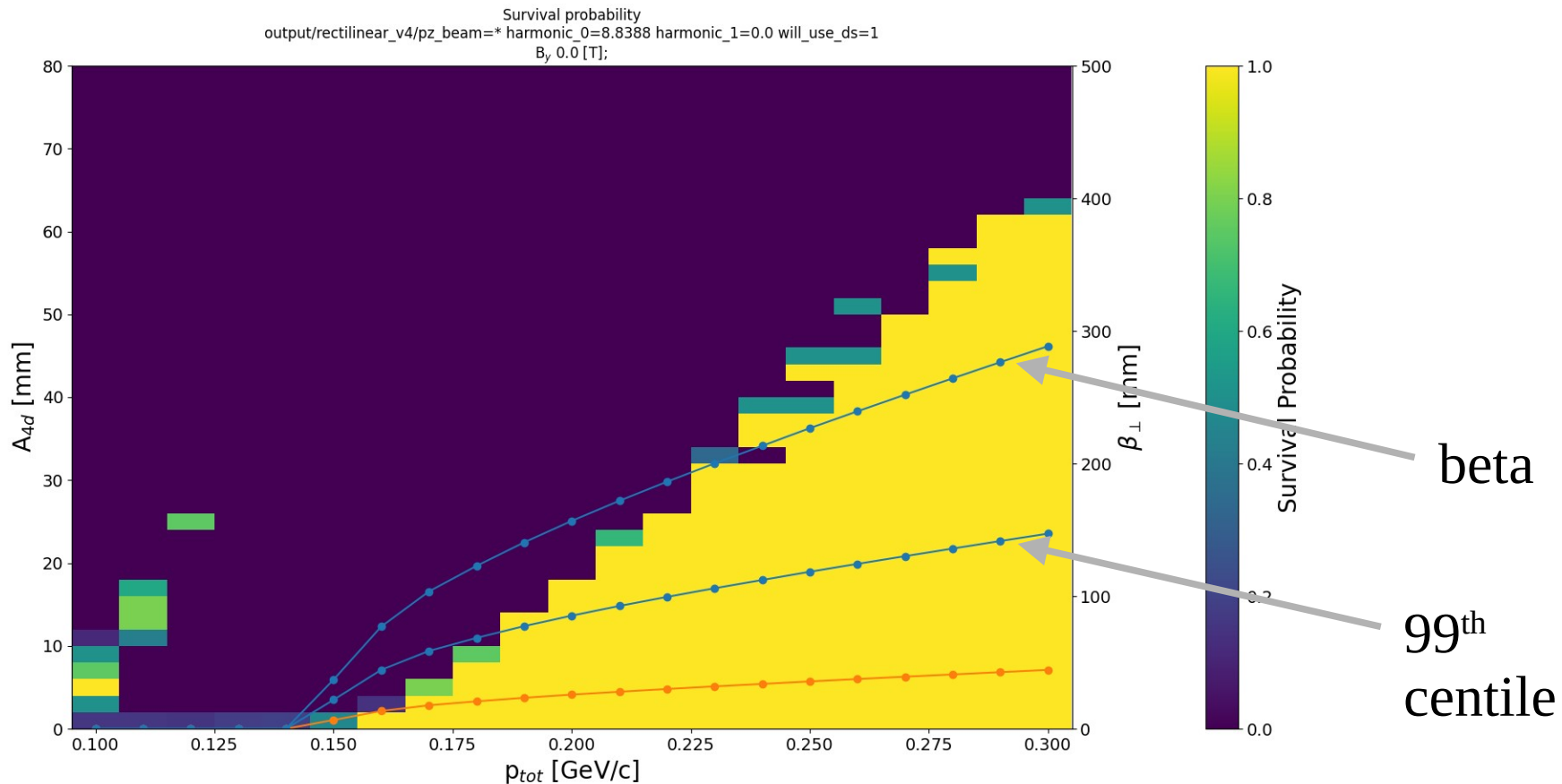
	$\epsilon_{T,\text{sim}}$ (mm)	$\epsilon_{L,\text{sim}}$ (mm)	$\epsilon_{6D,\text{sim}}$ (mm <sup>3</sup> )	Transmission
Start	16.96	45.53	13500	
A-Stage 1	5.165	18.31	492.6	75.2%
A-Stage 2	2.473	7.113	44.03	84.4%
A-Stage 3	1.556	3.880	9.594	85.6%
A-Stage 4	1.239	1.741	2.861	91.3%

Just  
decays:  
80 %

49.6 %

# A-type Acceptance

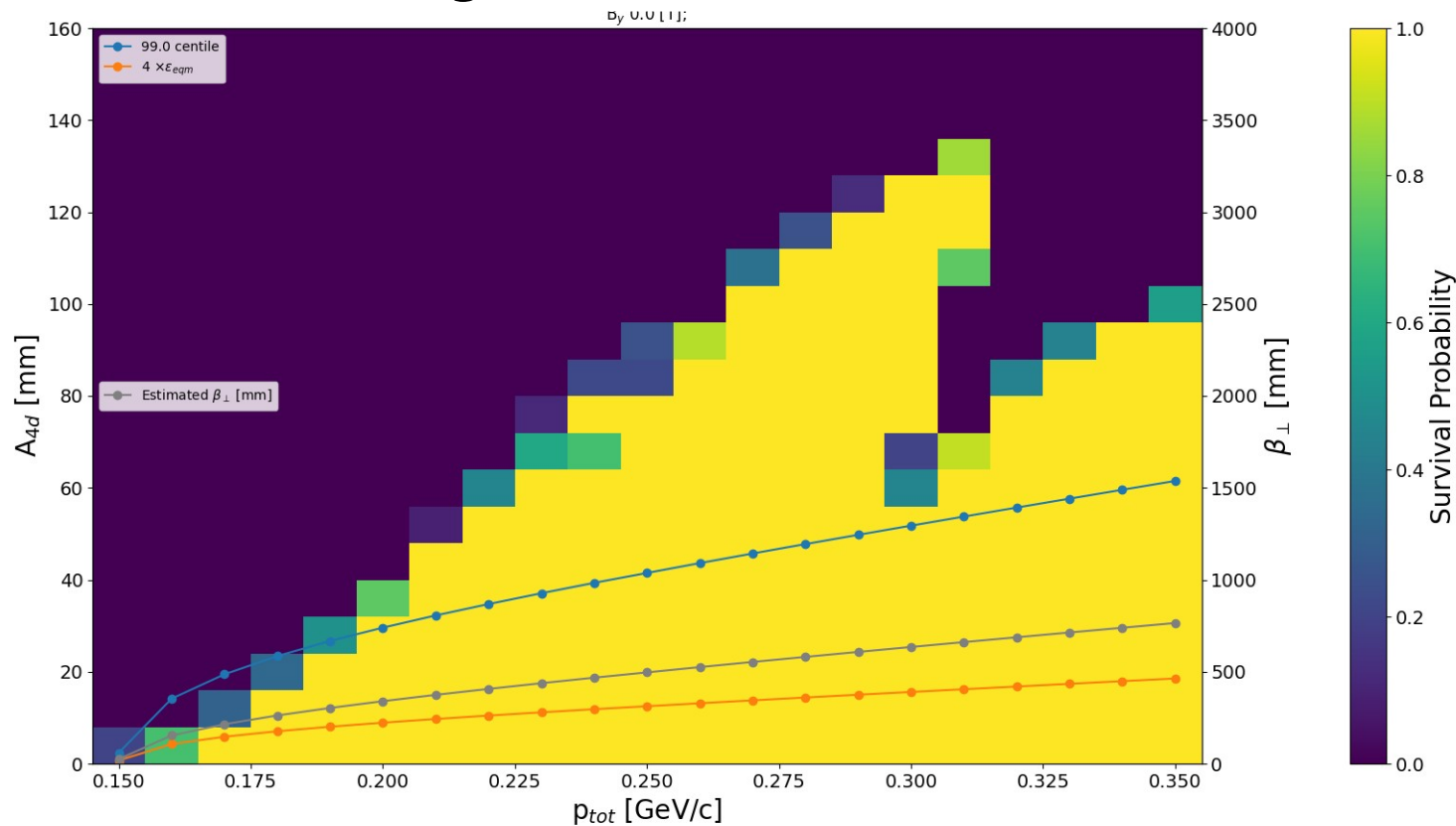
- A-type lattices
  - Set  $H_1$  to 0 and adjust  $H_0$  so that integral  $B_z^2 dz$  is constant
    - Average focusing strength goes with integral  $B_z^2 dz$
  - Field is approximately sine wave – stop band disappears



# Momentum range

- Can scale the momentum by scaling  $B_z$

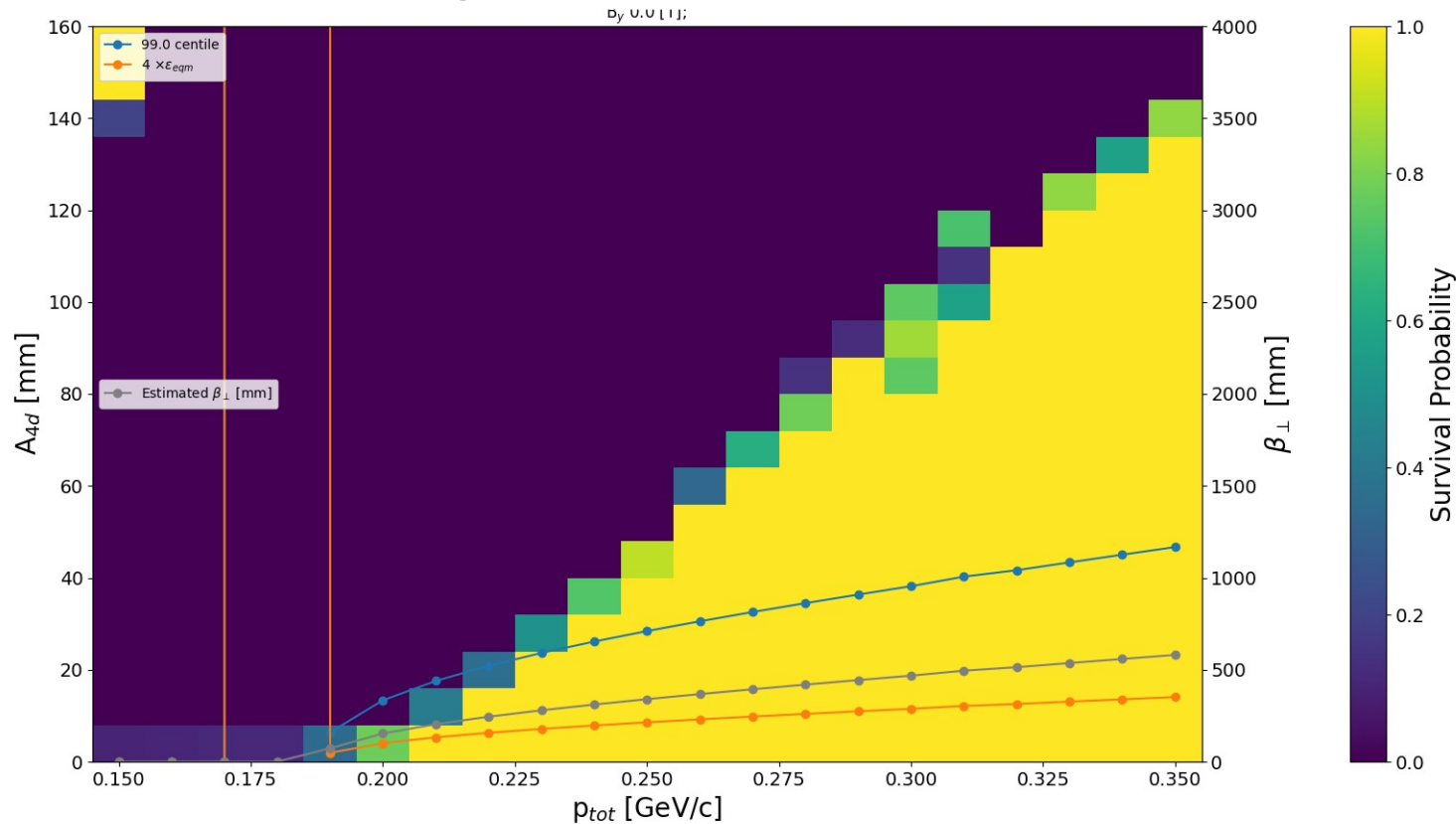
Cell length = 1.8 m;  $h_0 = 4.0$



# Momentum range

- Can scale the momentum by scaling  $B_z$

Cell length = 1.8 m;  $h_0 = 5.0$

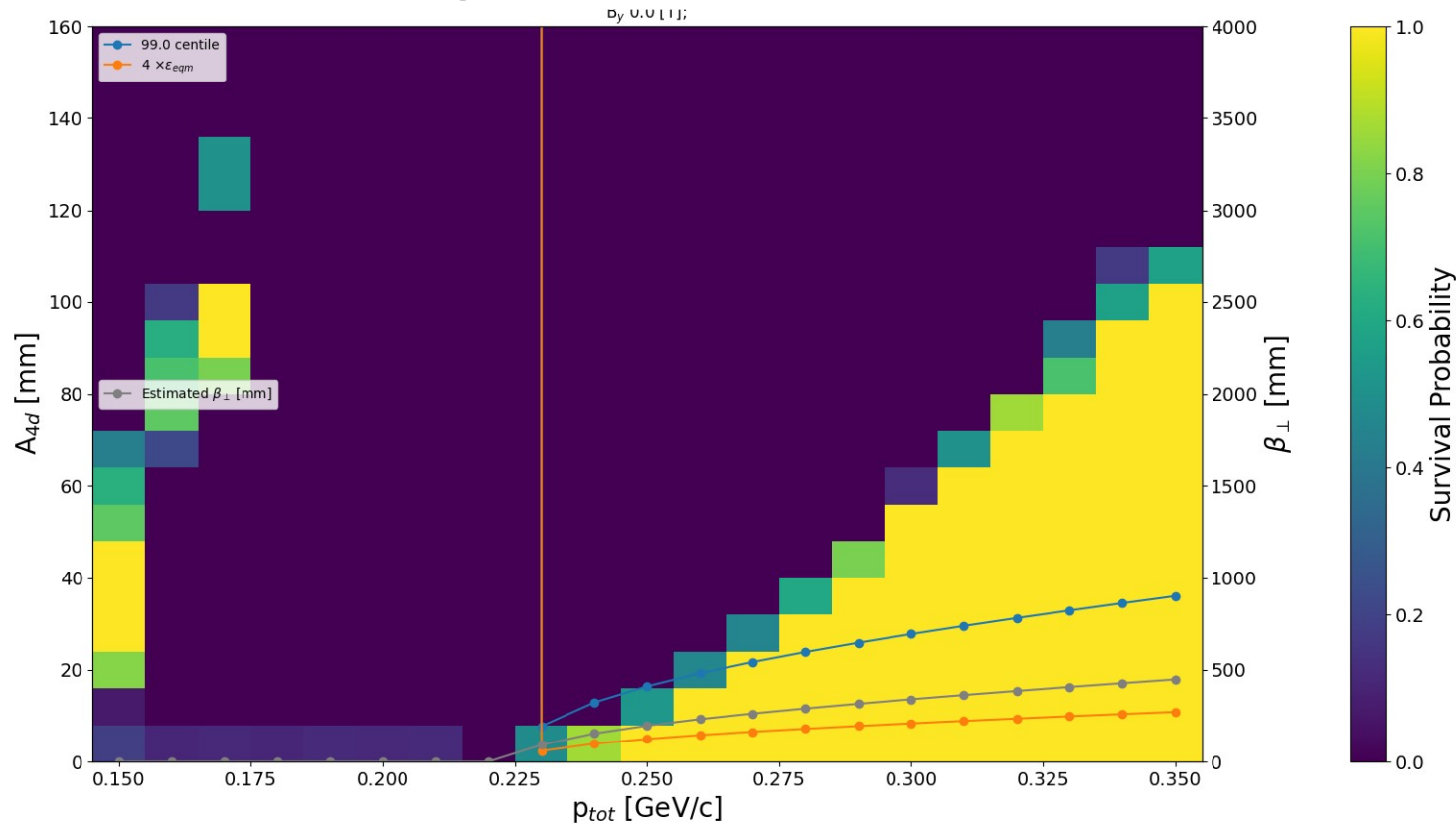




# Momentum range

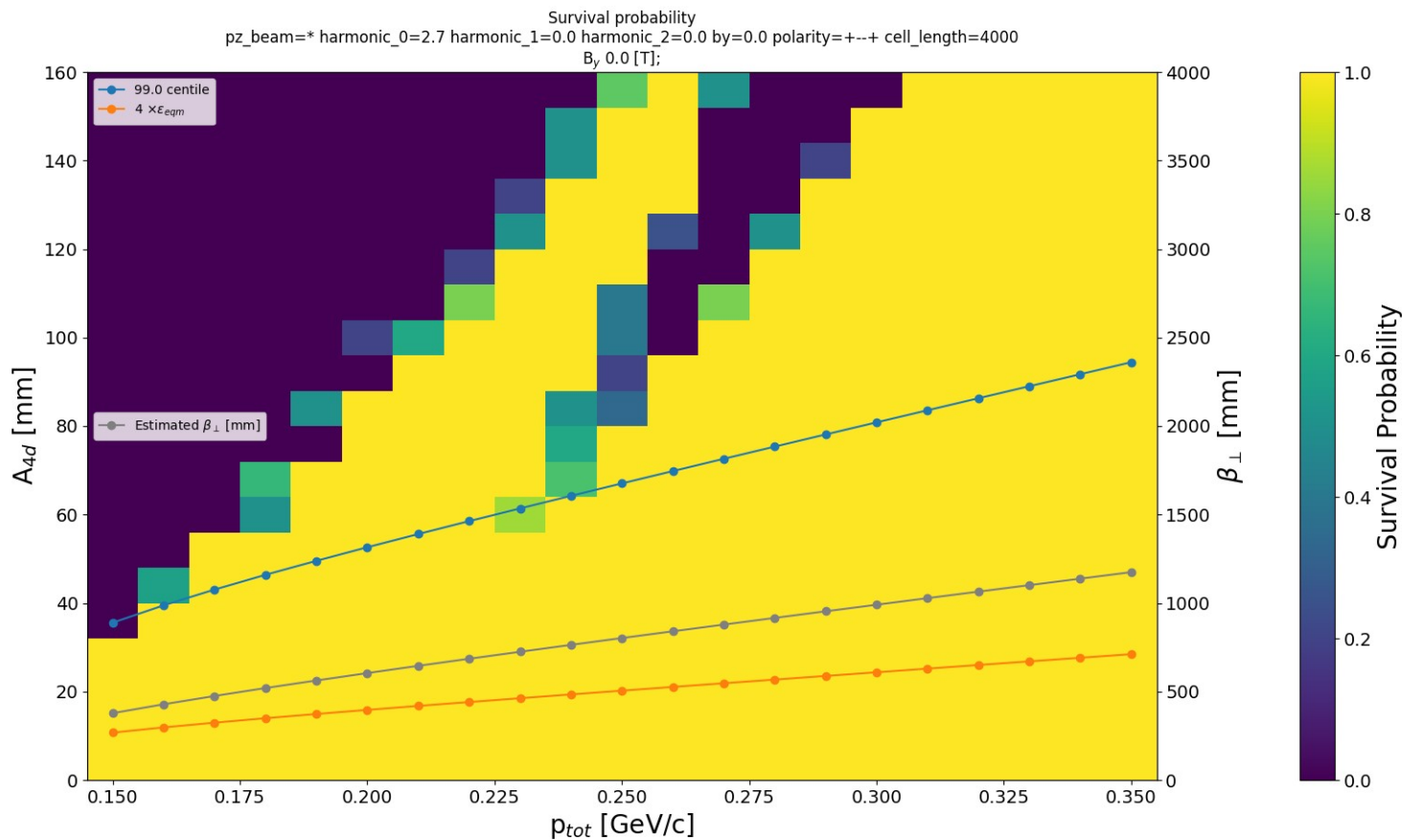
- Can scale the momentum by scaling  $B_z$

Cell length = 1.8 m;  $h_0 = 6.0$



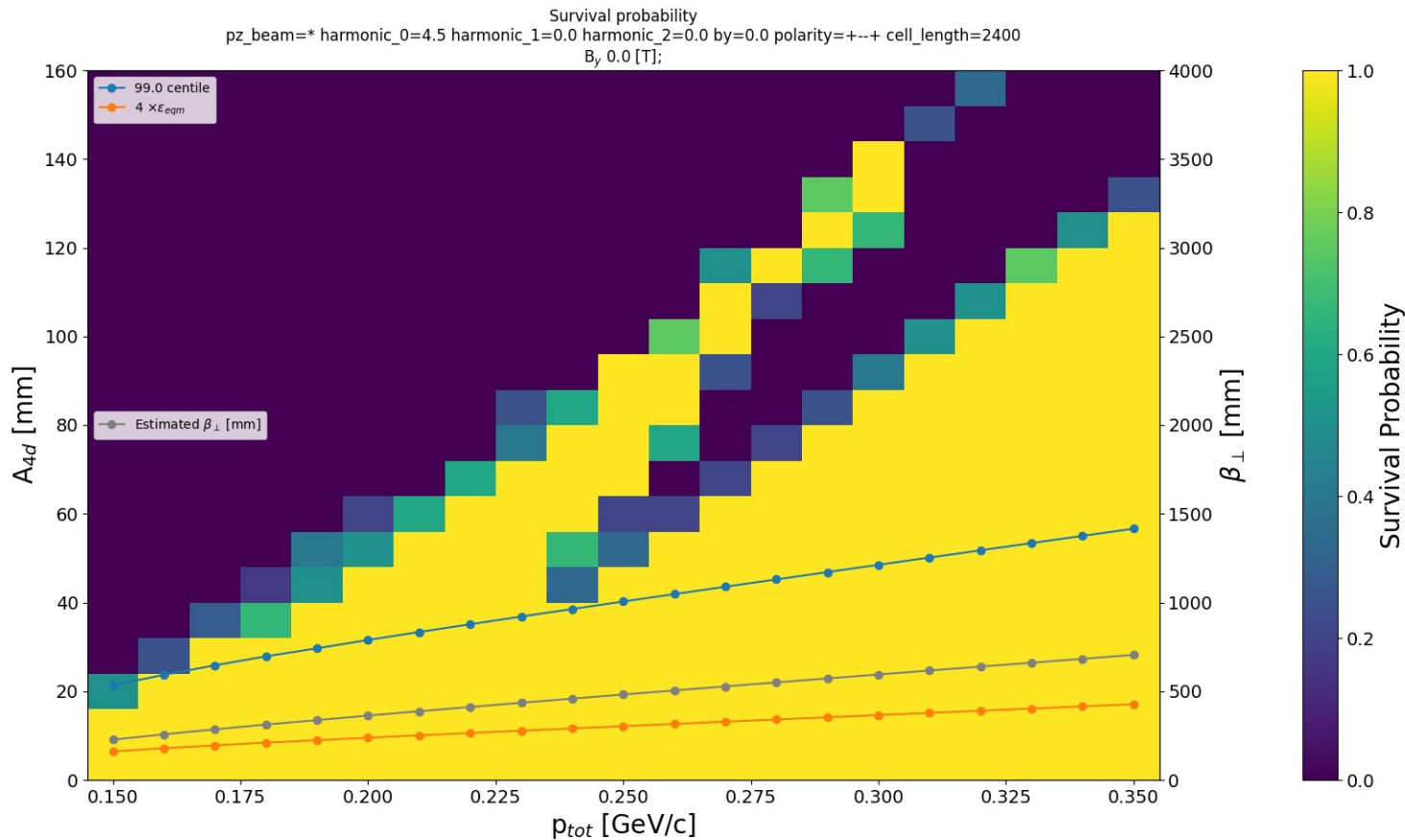
# Beta scaling & acceptance

- Can scale the beta function by scaling  $B_z$  and cell length
  - Keep (cell length) \*  $B_z$  constant



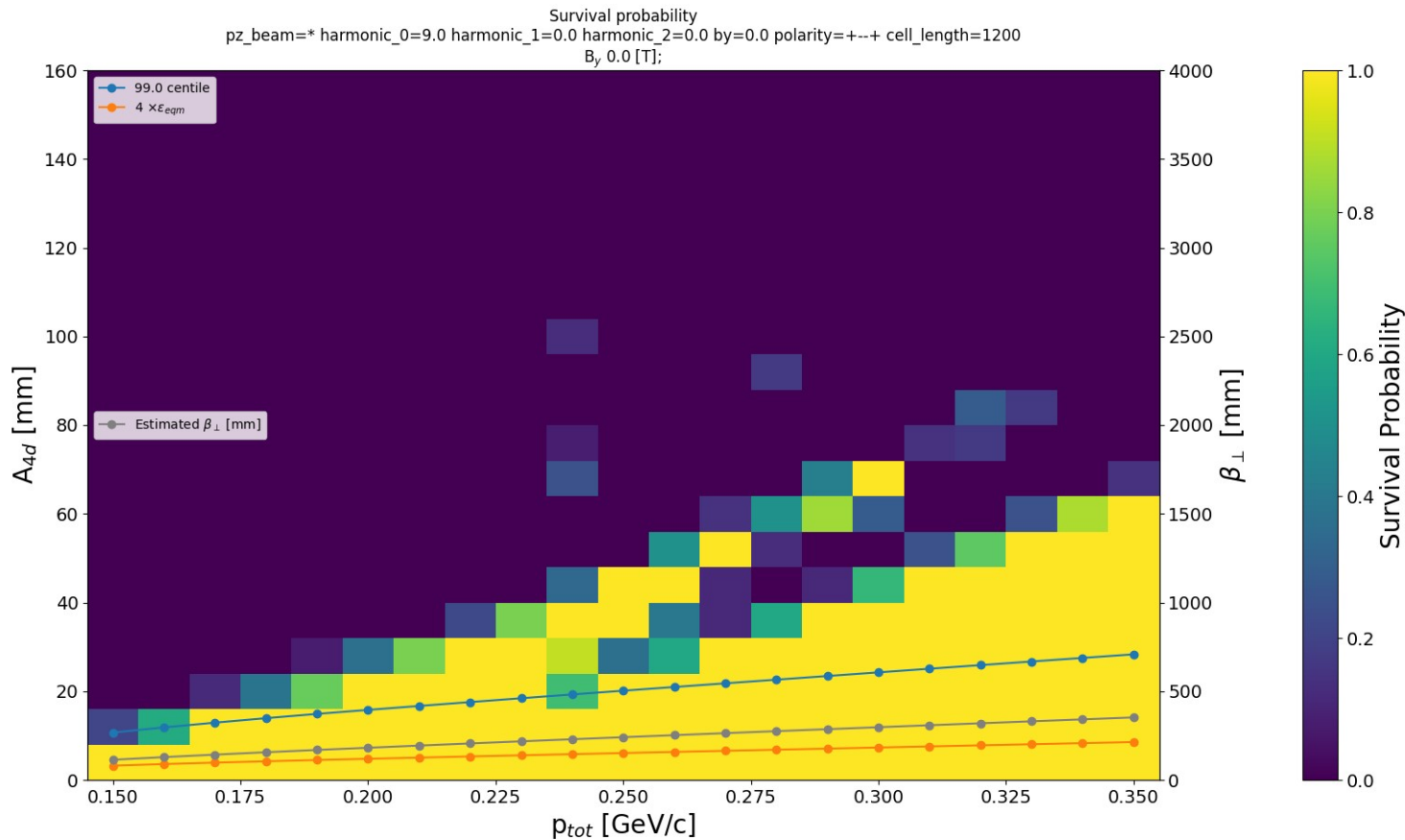
# Beta scaling & acceptance

- Can scale the beta function by scaling  $B_z$  and cell length
  - Keep (cell length) \*  $B_z$  constant



# Beta scaling & acceptance

- Can scale the beta function by scaling  $B_z$  and cell length
  - Keep (cell length) \*  $B_z$  constant



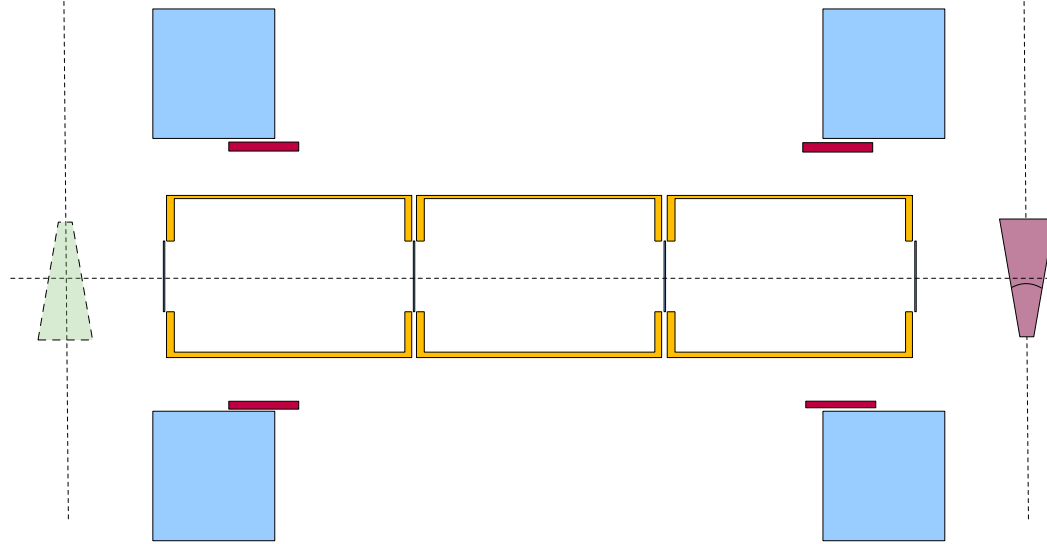


# Acceptance

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- We can generate *any* acceptance we like
- Dynamic range is conserved
- High acceptance is *easier* to generate than low acceptance
- Note also that dynamic range is better at high momentum
  
- So why do we have a bad transmission for A1+???

# Physical Acceptance

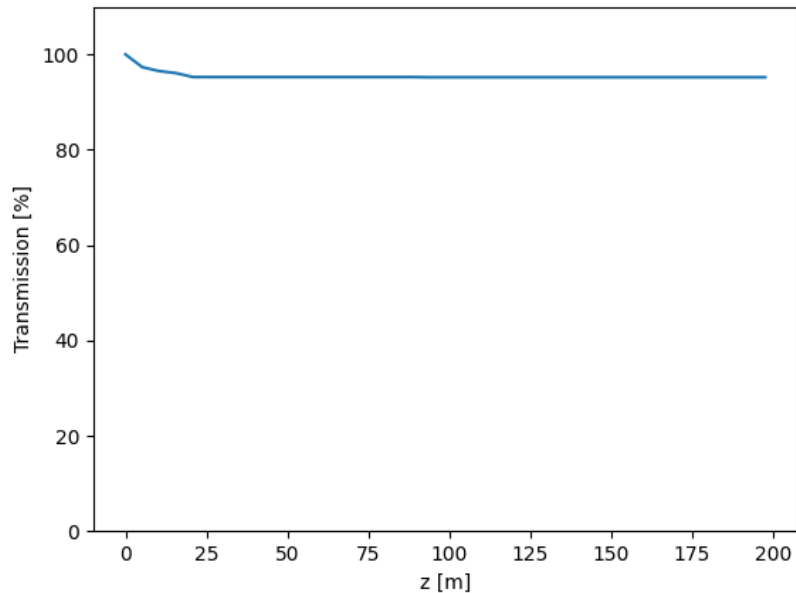


- Physical acceptance driven by RF cavity iris radius
- RF team advice:  $(\text{iris radius}) = 0.5 * (\text{ideal cavity radius})$
- Define “iris factor” =  $(\text{iris radius}) / (\text{ideal cavity radius})$

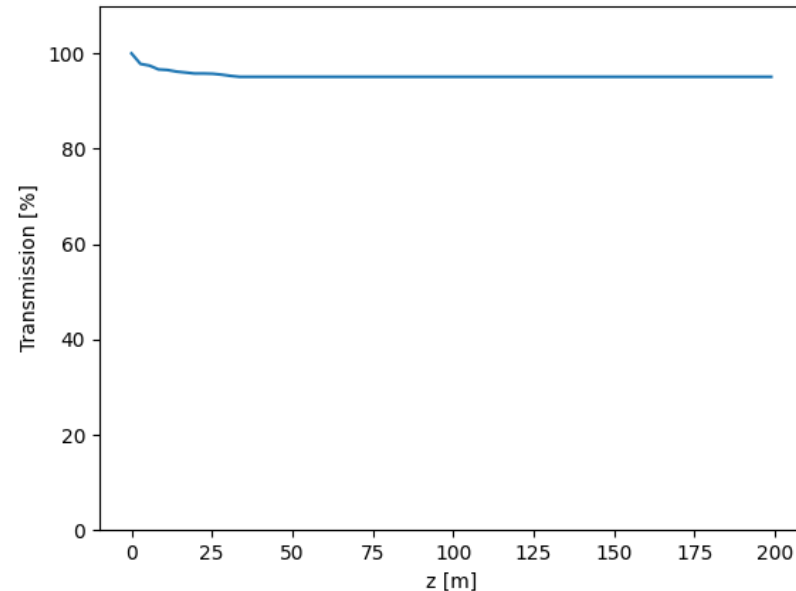
# Dynamic Acceptance

- Consider the beam used by Ruihu at A1
- Scan DA
- Transmission is high
- No dependence on DA

$B_0$  1.7308 [T]; N(rf) 0 ; Cell Length 5200 [mm]



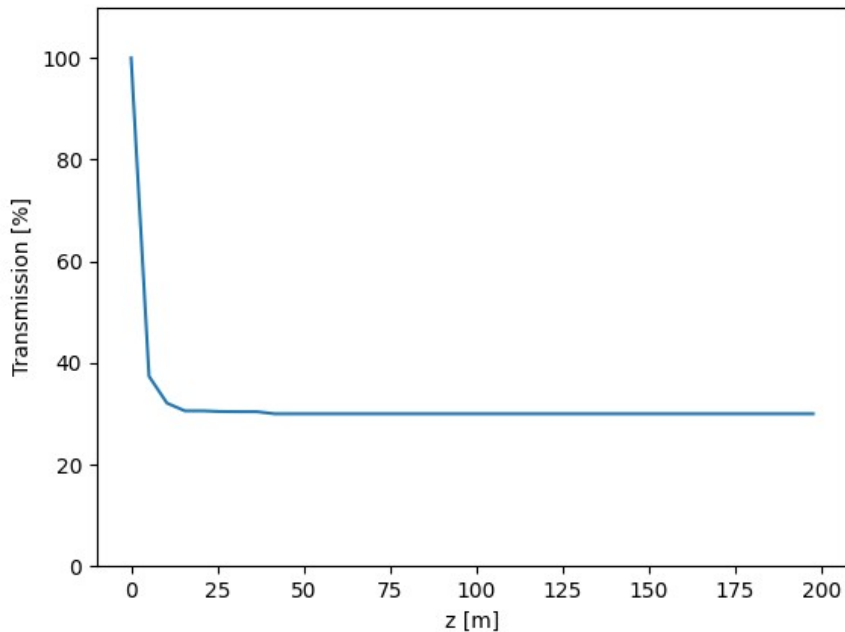
$B_0$  3.2143 [T]; N(rf) 0 ; Cell Length 2800 [mm]



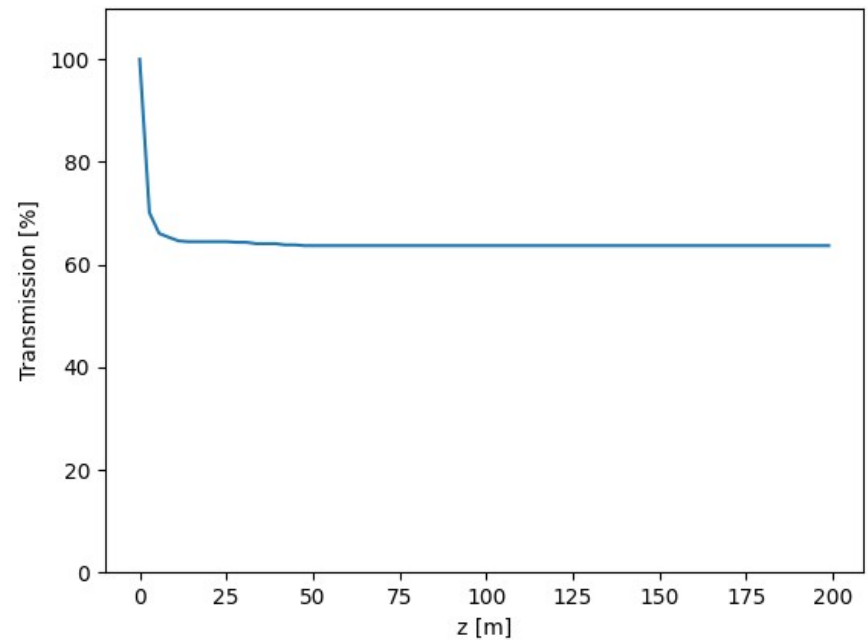
# Physical Acceptance

- Consider the beam used by Ruihu at A1
- Add a single RF cavity @ 352 MHz (iris factor 0.5)
- Transmission is terrible

$B_0$  1.7308 [T]; N(rf) 1 ; Cell Length 5200 [mm]



$B_0$  3.2143 [T]; N(rf) 1 ; Cell Length 2800 [mm]

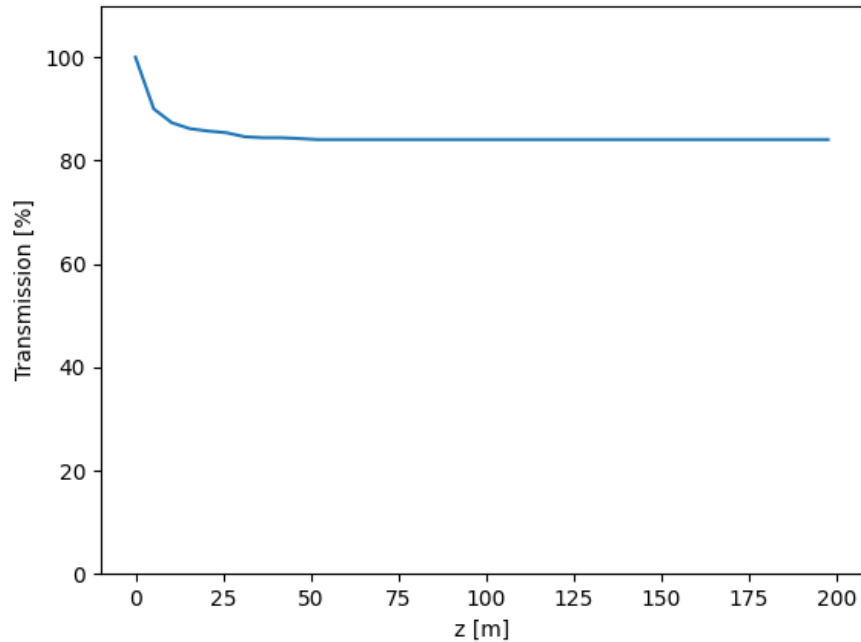




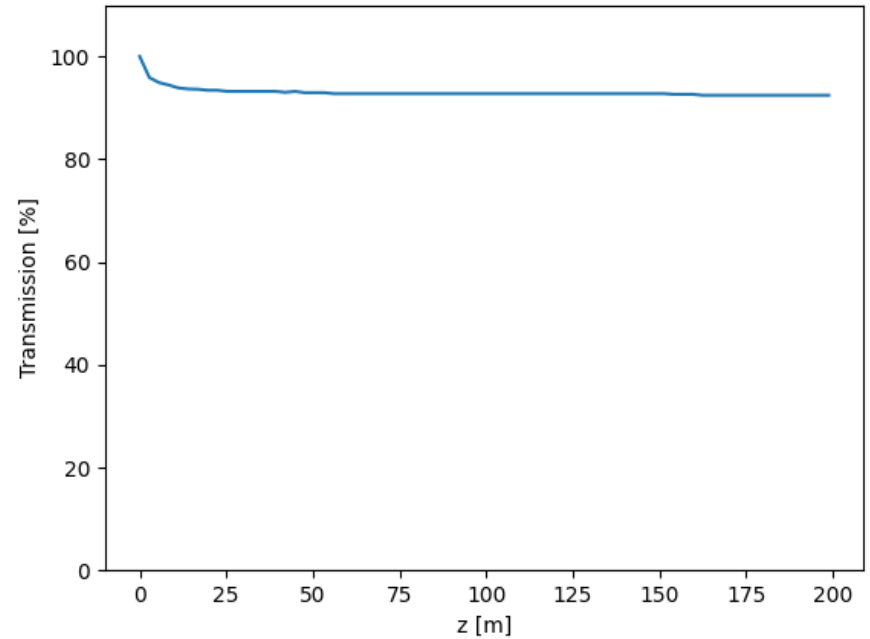
# RF frequency

- Higher radius  $\rightarrow$  lower frequency (176 MHz)

$B_0$  1.7308 [T]; N(rf) 1 ; Cell Length 5200 [mm]



$B_0$  3.2143 [T]; N(rf) 1 ; Cell Length 2800 [mm]



# Conclusions

- Physical acceptance is limiting performance of A-type lattices
- Improved physical acceptance → lower frequency
- Need to consider 176 MHz RF for A-type lattice
  
- We can generate a good dynamic range for the A-type lattices
  - Should be possible to make a good cooling performance with decay-dominated transmission
  
- Addendum... can we capture into 176 MHz RF?