



Exploring Meson Structures from Lattice QCD

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Outline

- Electromagnetic Form Factor (EMFF) of Pion and Kaon
 - based on Phys. Rev. Lett. **133**, 181902
- Pion Light-cone Generalized Parton Distribution (GPD)
 - based on arXiv: 2407.03516

Motivation

EPJA 48 (2012) 187 JPG 48 (2021) 075106 arXiv: 2102.09222

Front. Phys. 16 (2021) 64701

- Experiment: JLab, EIC, EicC ...

Gao et al., PRD 96 (2017) 034024

- Effective theory: QCD sum rules, DSE ...

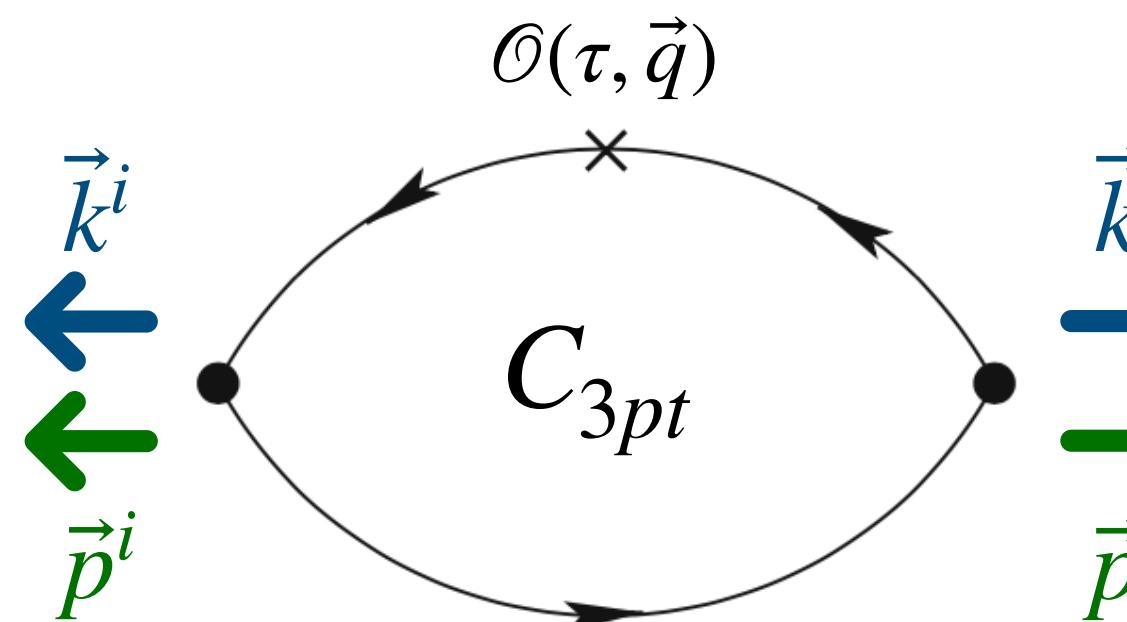
- Lattice QCD: first principle

PRD 96 (2017) 114509

ETMC, PRD 105 (2022) 054502

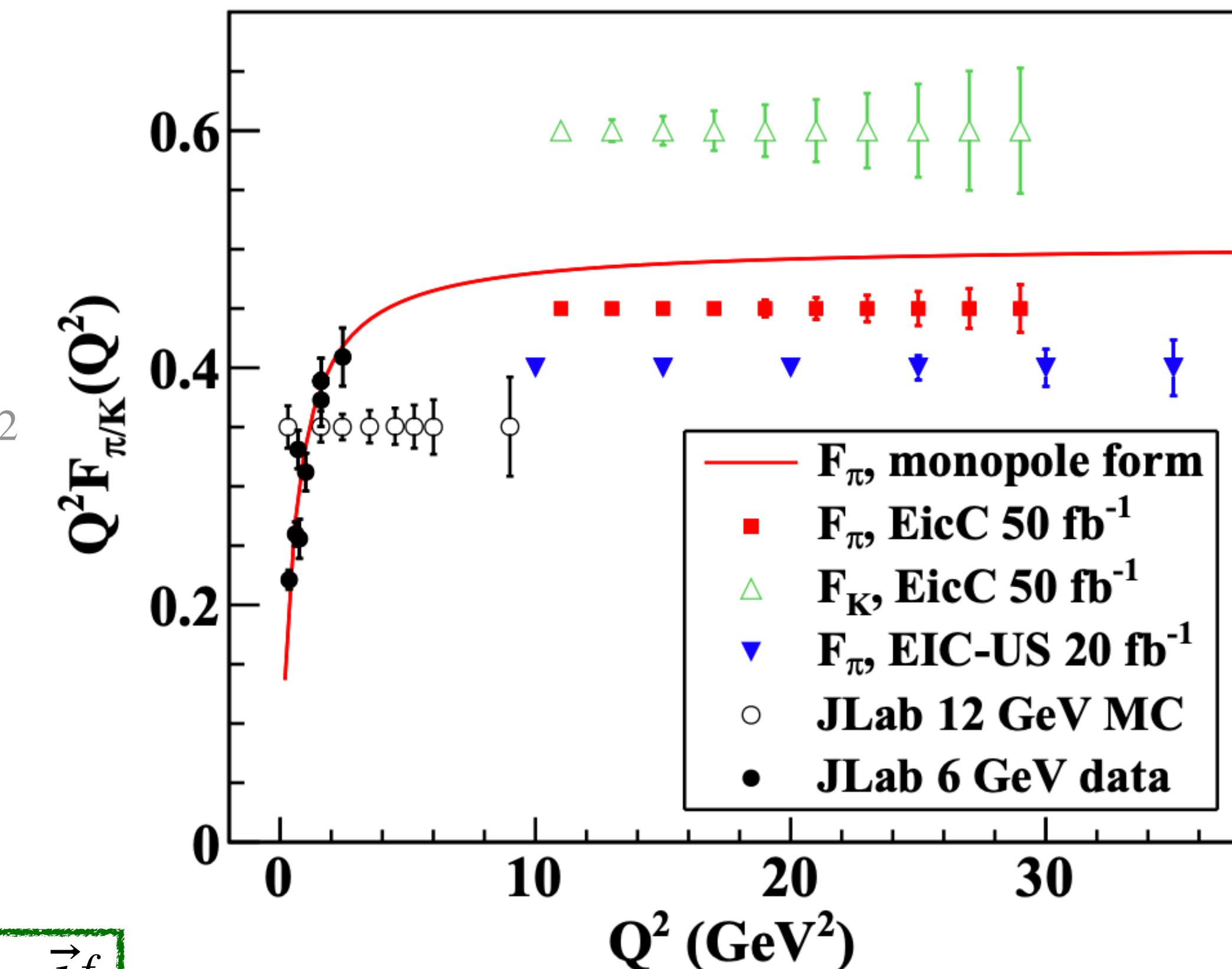
○ State-of-the-art: $Q^2 \leq 6$ (pion), 3 (kaon) GeV^2

○ This work: Q^2 up to 10, 28 GeV^2



$$\boxed{\text{Boost parameter } \zeta = \frac{\vec{k}^i}{\vec{p}^i} = \frac{\vec{k}^f}{\vec{p}^f}}$$

$$\boxed{\text{Momentum transfer } Q^2 = -t}$$

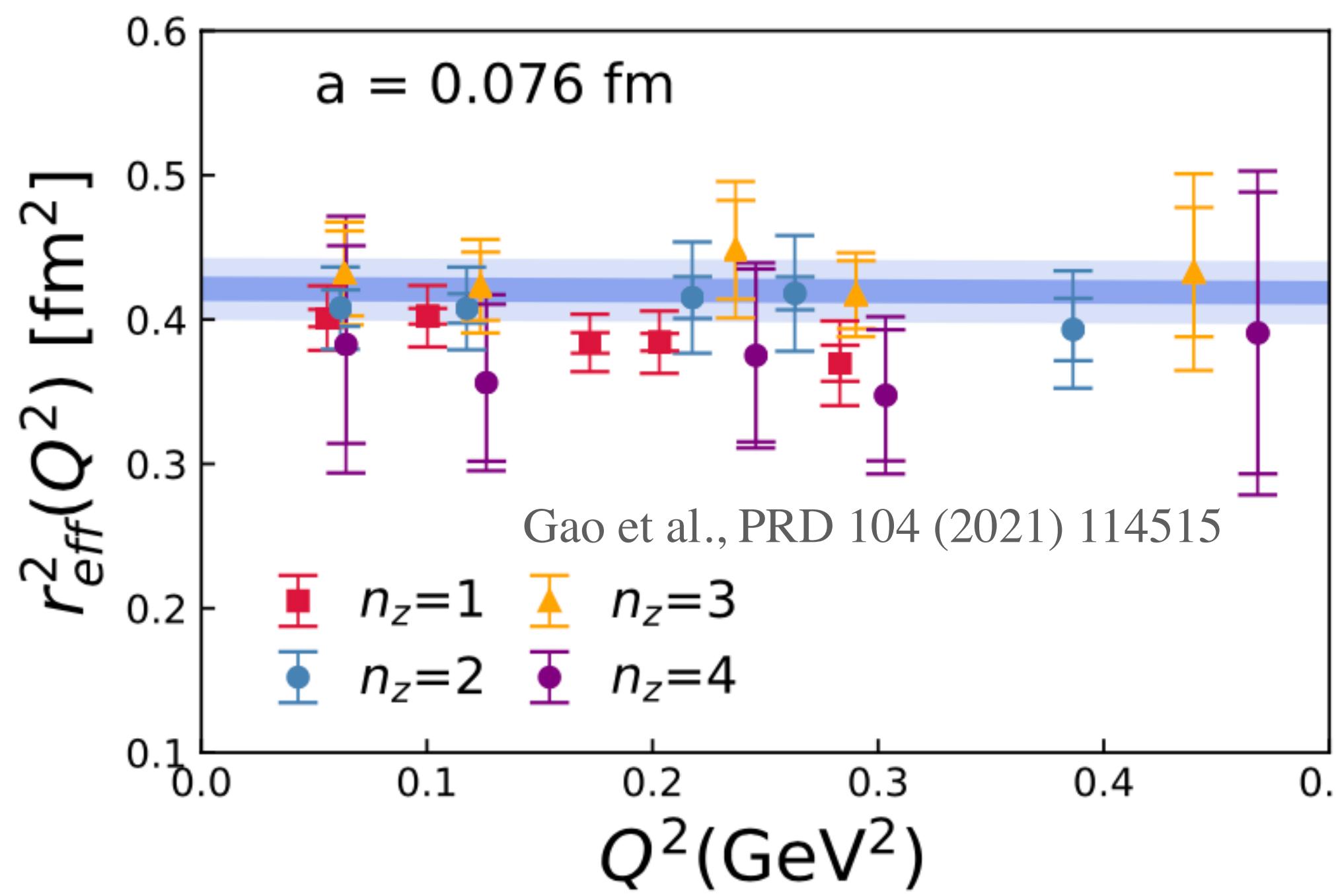


Motivation

Low Q^2 : Vector Meson Dominance

$$r_{\text{eff}}^2(Q^2) = 6[1/F_\pi(Q^2) - 1]/Q^2$$

$$\langle r_\pi^2 \rangle = 0.42(2) \text{ fm}^2, \quad \langle r_\pi^2 \rangle_{PDG} = 0.434(5) \text{ fm}^2$$

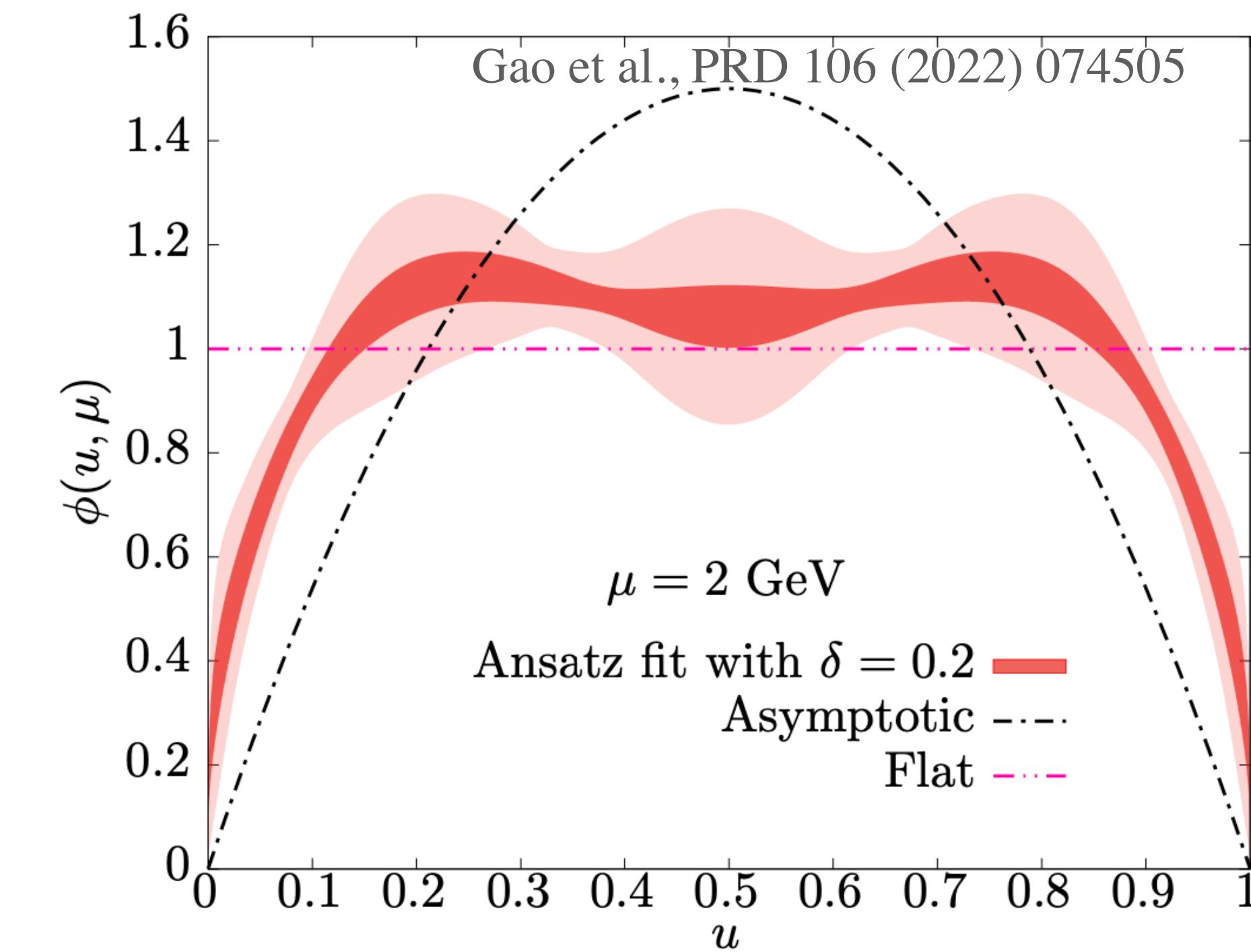


Larger Q^2

High Q^2 : Factorization framework

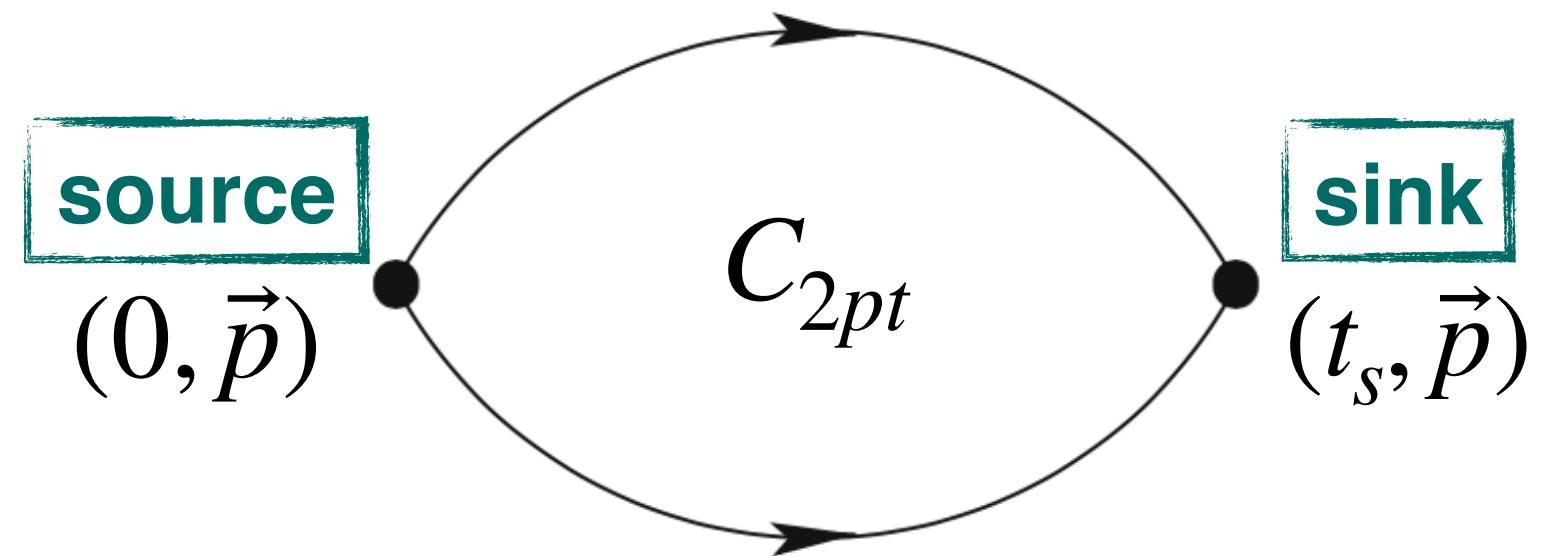
$$F(Q^2) = \int \int dx dy \Phi^*(y, \mu_F^2) T_H(x, y, Q^2, \mu_R^2, \mu_F^2) \Phi(x, \mu_F^2)$$

Hard-process kernel
Distribution amplitude



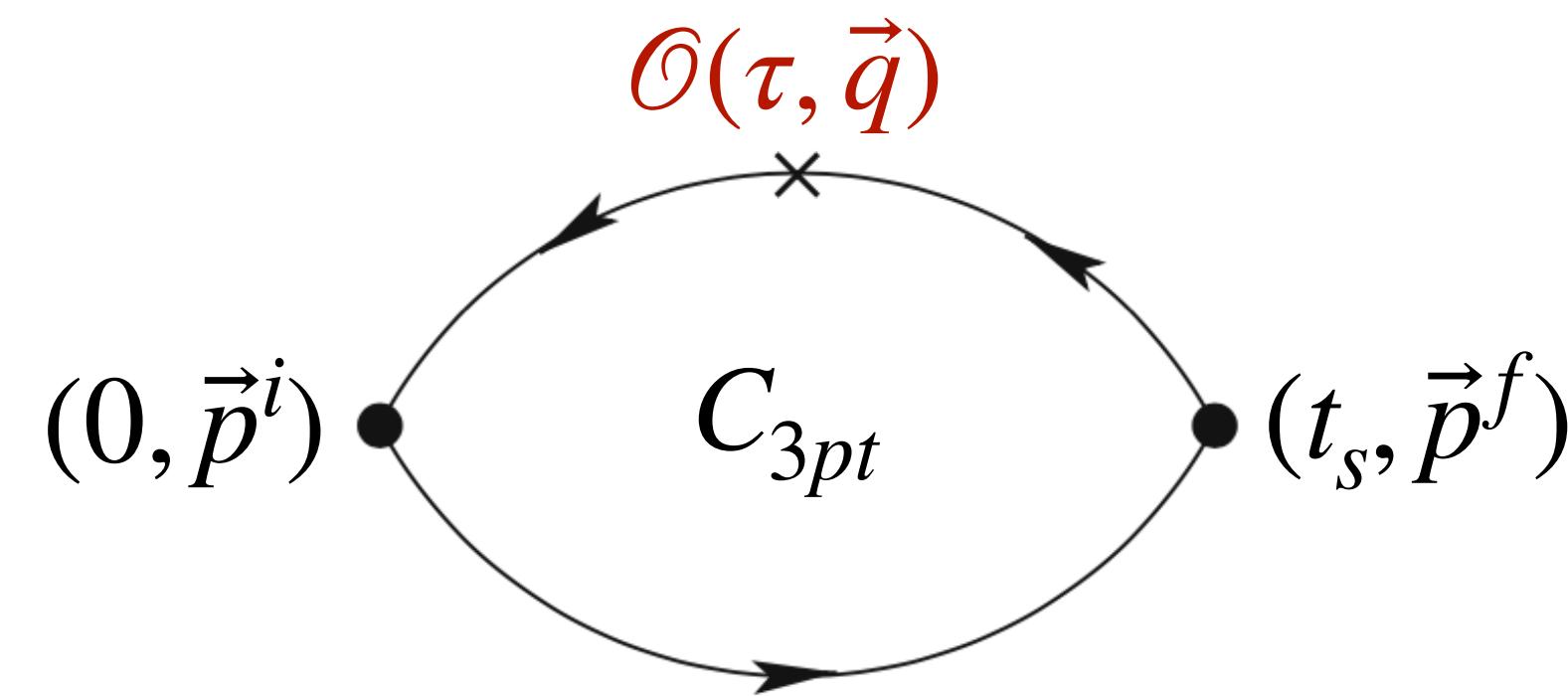
Test the factorization Q^2

How to get the form factor on the lattice



$$C_{2pt}(t, \vec{p}) = \langle H(t_s, \vec{p}) H^\dagger(0, \vec{p}) \rangle$$

+



$$C_{3pt}(\tau, t_s; \vec{p}^i, \vec{p}^f) = \langle H(t_s, \vec{p}^f) \hat{\mathcal{O}}_{\gamma_\mu}(\tau, \vec{q}) H^\dagger(0, \vec{p}^i) \rangle$$

$$R^{fi} \sim C_{3pt} / C_{2pt} \xrightarrow{t_s \rightarrow \infty}$$

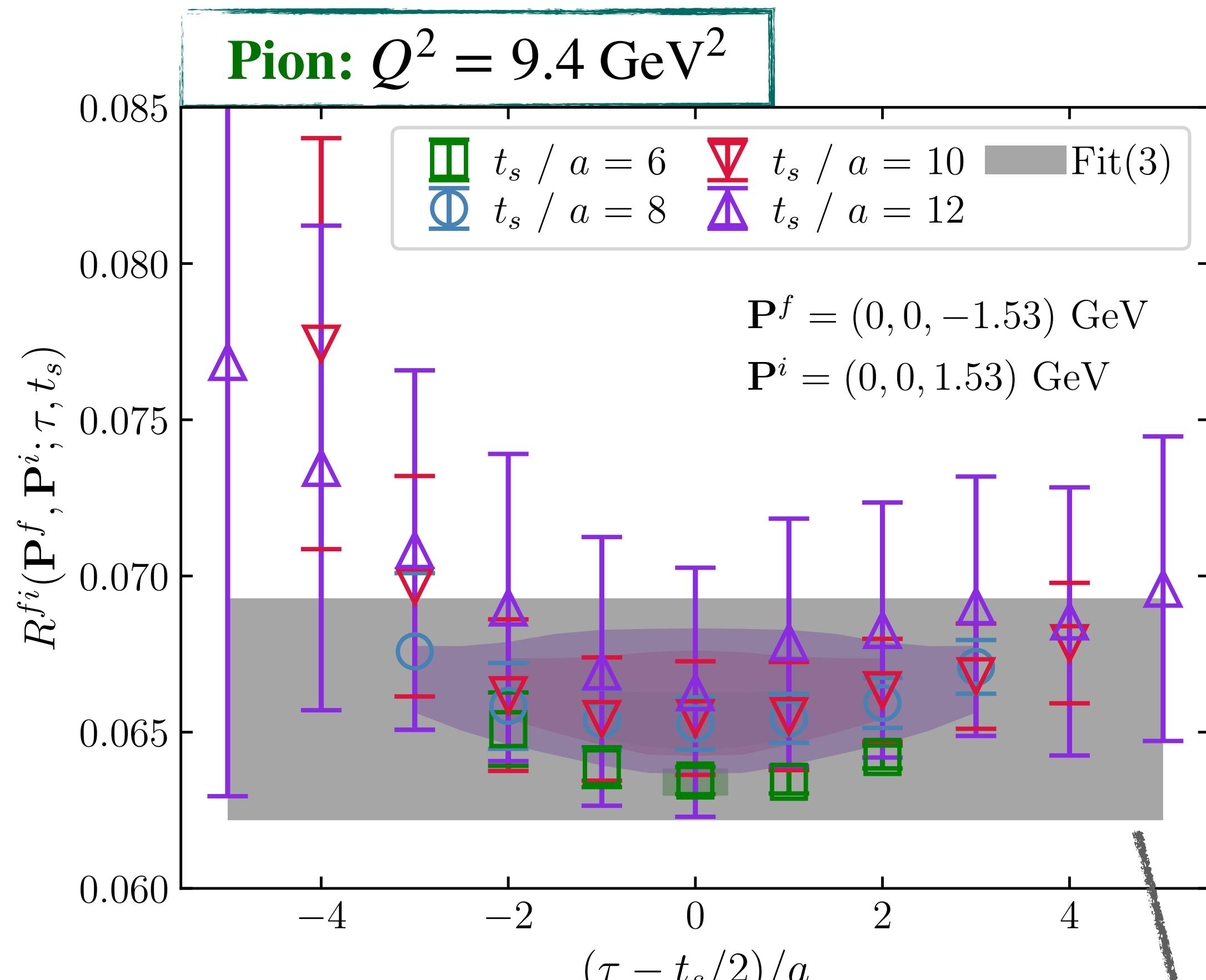
→ $F^B = \langle E_0, \vec{p}^f | \hat{\mathcal{O}}_{\gamma^\mu}(\tau, \vec{q}) | E_0, \vec{p}^i \rangle$ →

Bare Form factor

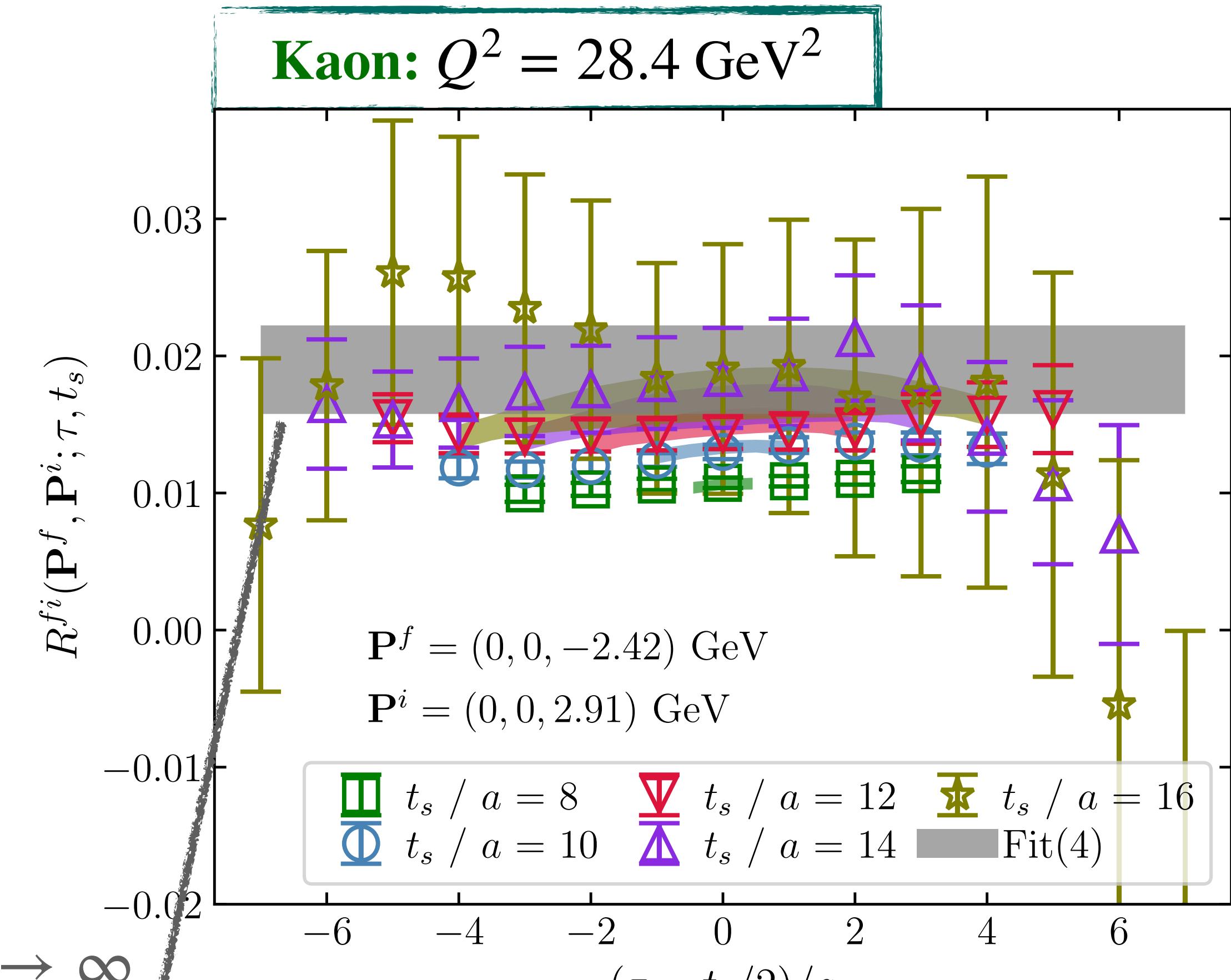
$$F(Q^2) = F^B \times Z_V^{-1}$$

$(Q^2 = -t)$

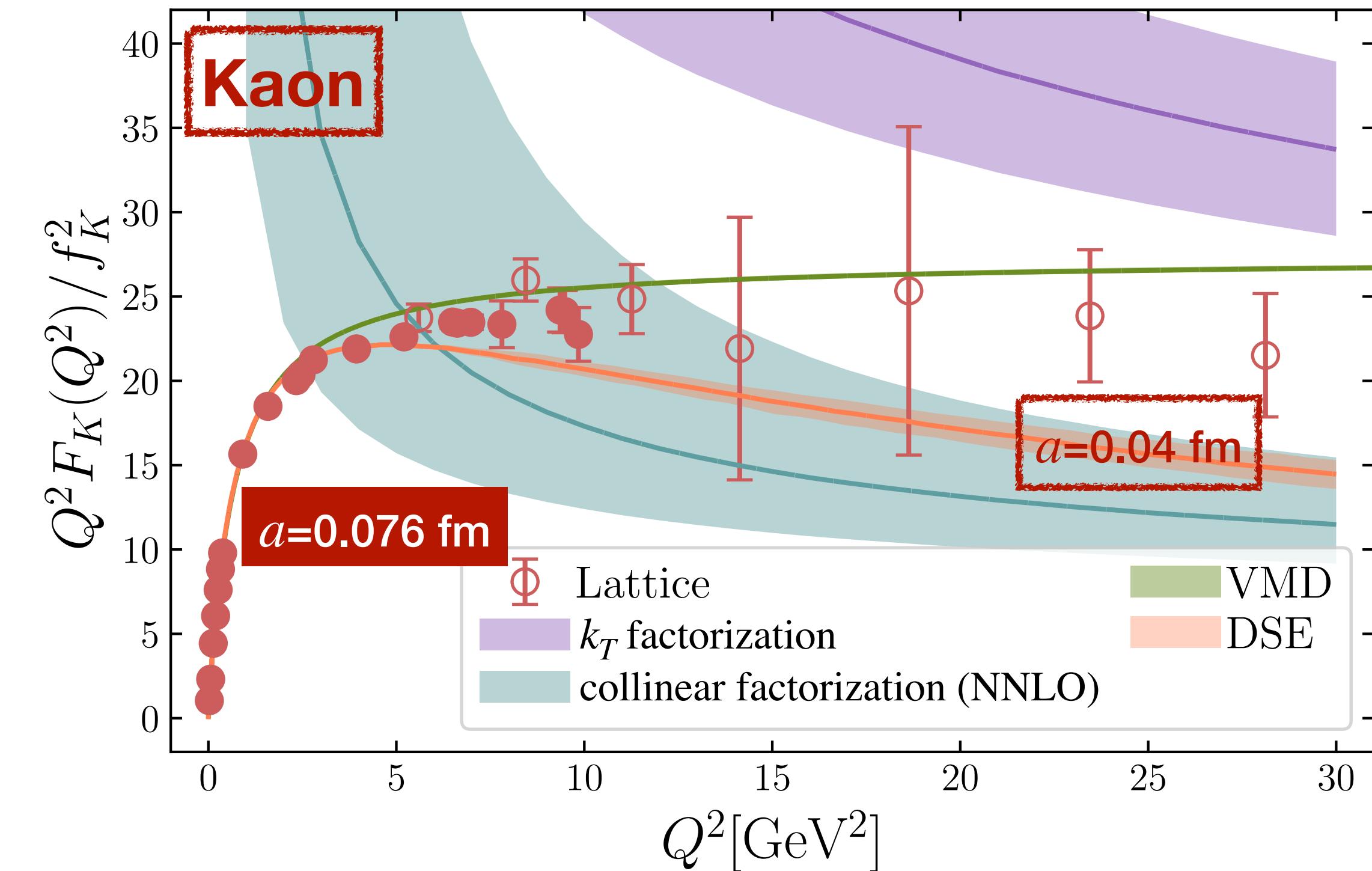
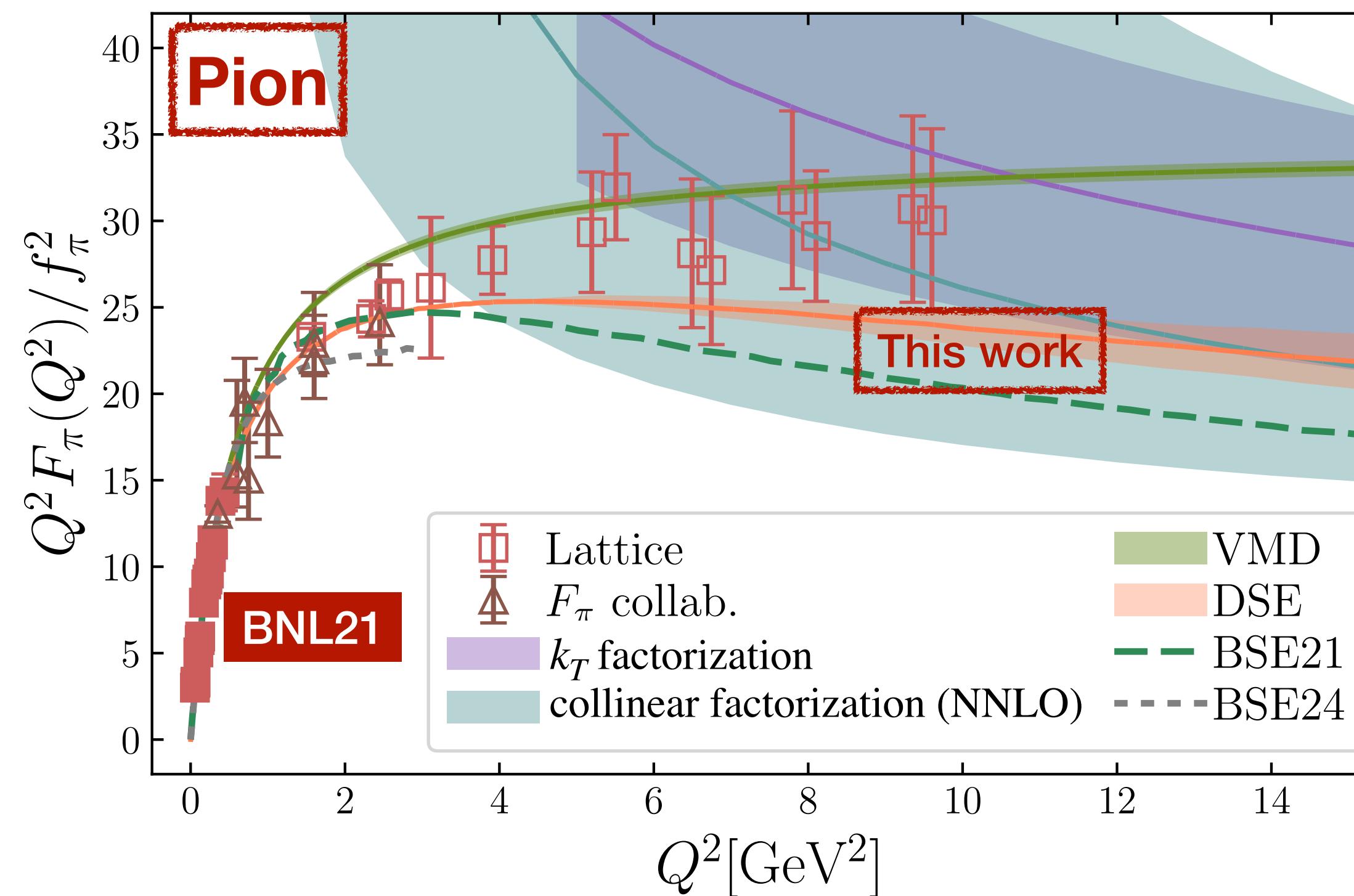
Bare Form Factor



Bare form factor $F^B \times Z_V^{-1} = F(Q^2)$



#F Electromagnetic Form Factor



Hard-process kernel

$$F(Q^2) = \int \int dx dy \Phi^*(y, \mu_F^2) T_H(x, y, Q^2, \mu_R^2, \mu_F^2) \Phi(x, \mu_F^2)$$

Distribution amplitude

$$F(Q^2 \rightarrow \infty) = 8\pi\alpha_s(Q^2)f^2/Q^2, \quad Q^2 F/f^2 \sim \text{Constant}$$

Lattice (BNL21): Gao et al., PRD 104 (2021) 114515

F^π collaboration: Huber et al., PRC 78 (2008) 045203

DSE (Dyson-Schwinger equation): Yao et al., PLB 855 (2024) 138823

BSE21 (Bethe-Salpeter equation): Ydrefors et al., PLB 820 (2021) 136494

BSE24: Jia and Cloët, arXiv:2402.00285

k_T factorization: Cheng, PRD 100 (2019) 013007

pion: Chai et al., EPJC 83 (2023) 556

kaon: in preparation

T_H in the collinear factorization: Chen et al., PRL 132 (2024) 201901

DA in the collinear factorization: Cloët et al., arXiv:2407.00206

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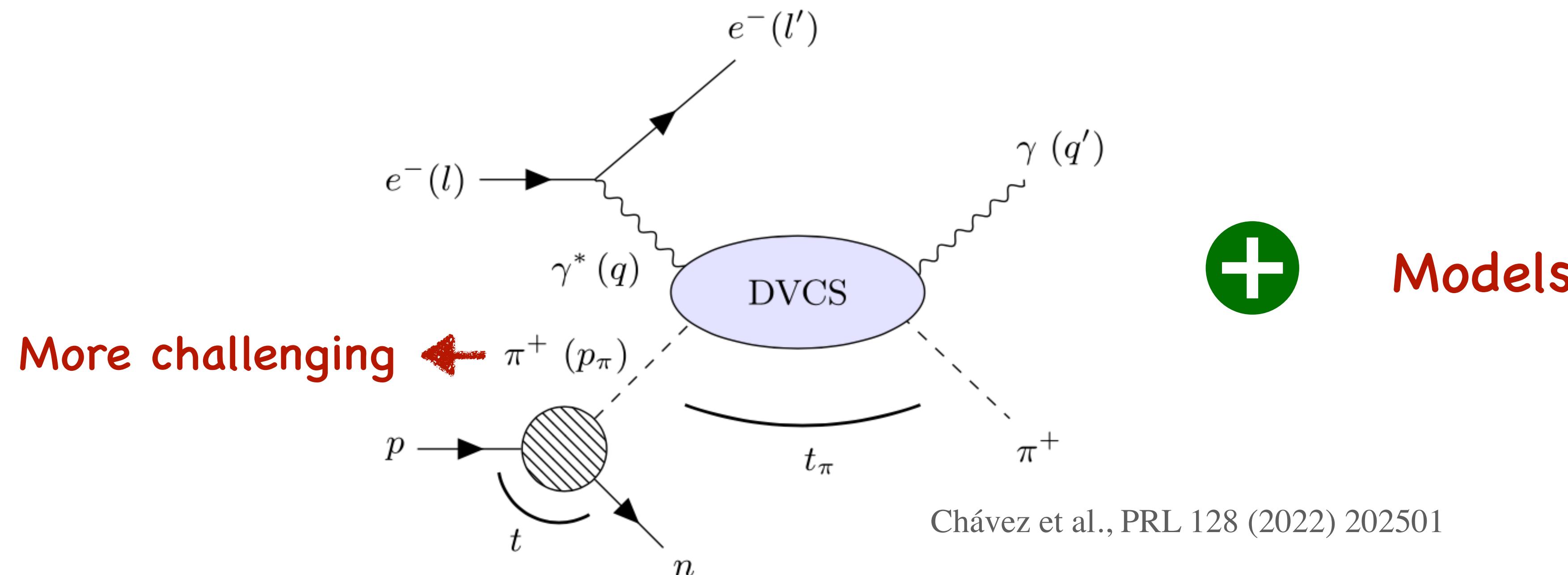
1D

Form Factor (pion-electron scattering)

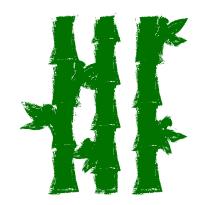
Parton Distribution Functions (Drell-Yan process)

3D

Generalized Parton Distributions



Lattice QCD: from first principle



Frame-independent approach

Bhattacharya et al., PRD 106 (2022) 114512

Calculate in the asymmetric frame
Save computational cost

- Lorentz-invariant amplitudes A_i 's (frame-independent)

$$M^\mu(P^m u, z^\mu, \Delta^\mu) = \bar{P}^\mu A_1 + m^2 z^\mu A_2 + \Delta^\mu A_3, \quad \bar{P}^\mu = (p_f^\mu + p_i^\mu)/2, \quad \Delta^\mu = p_f^\mu - p_i^\mu.$$

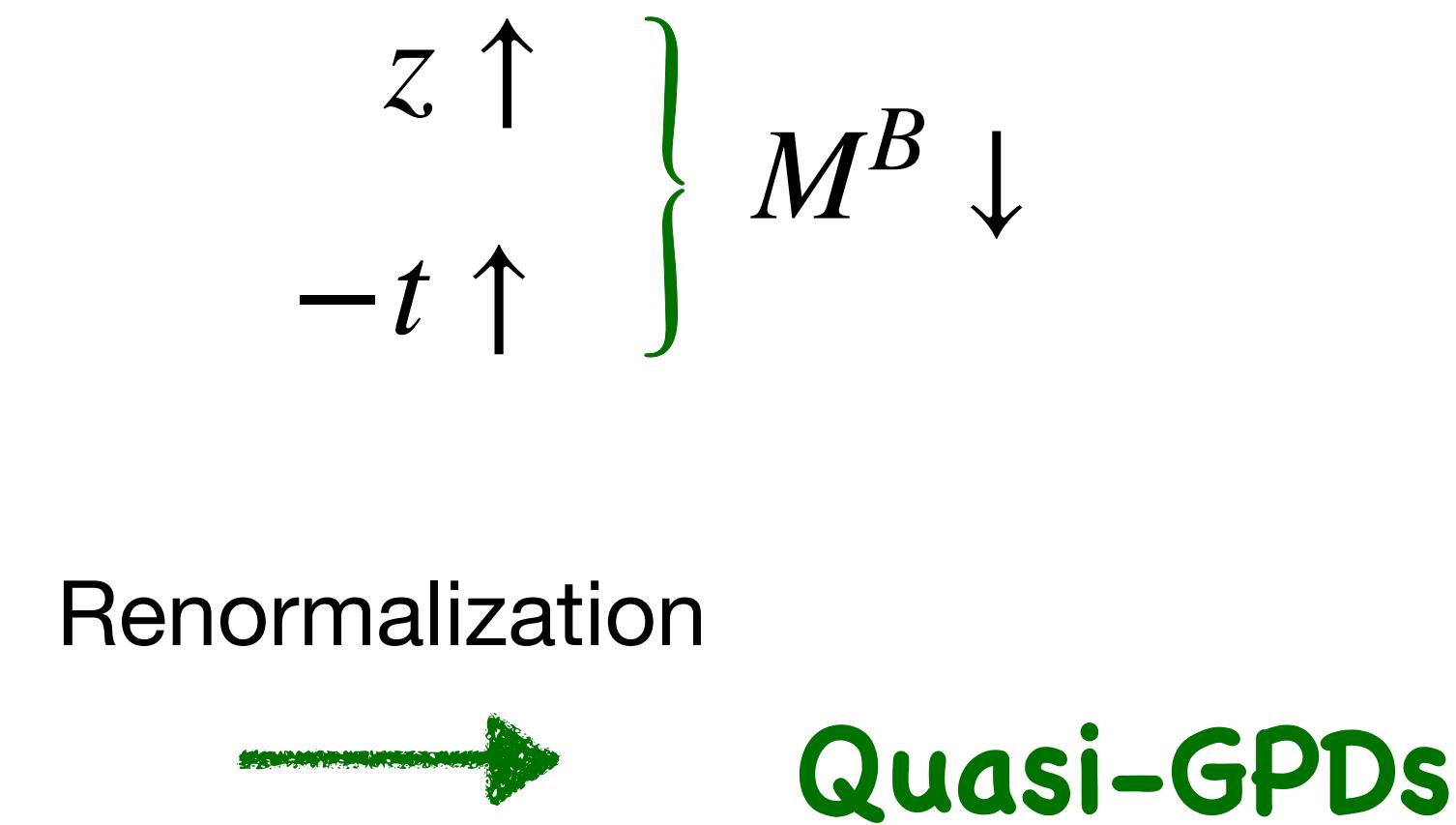
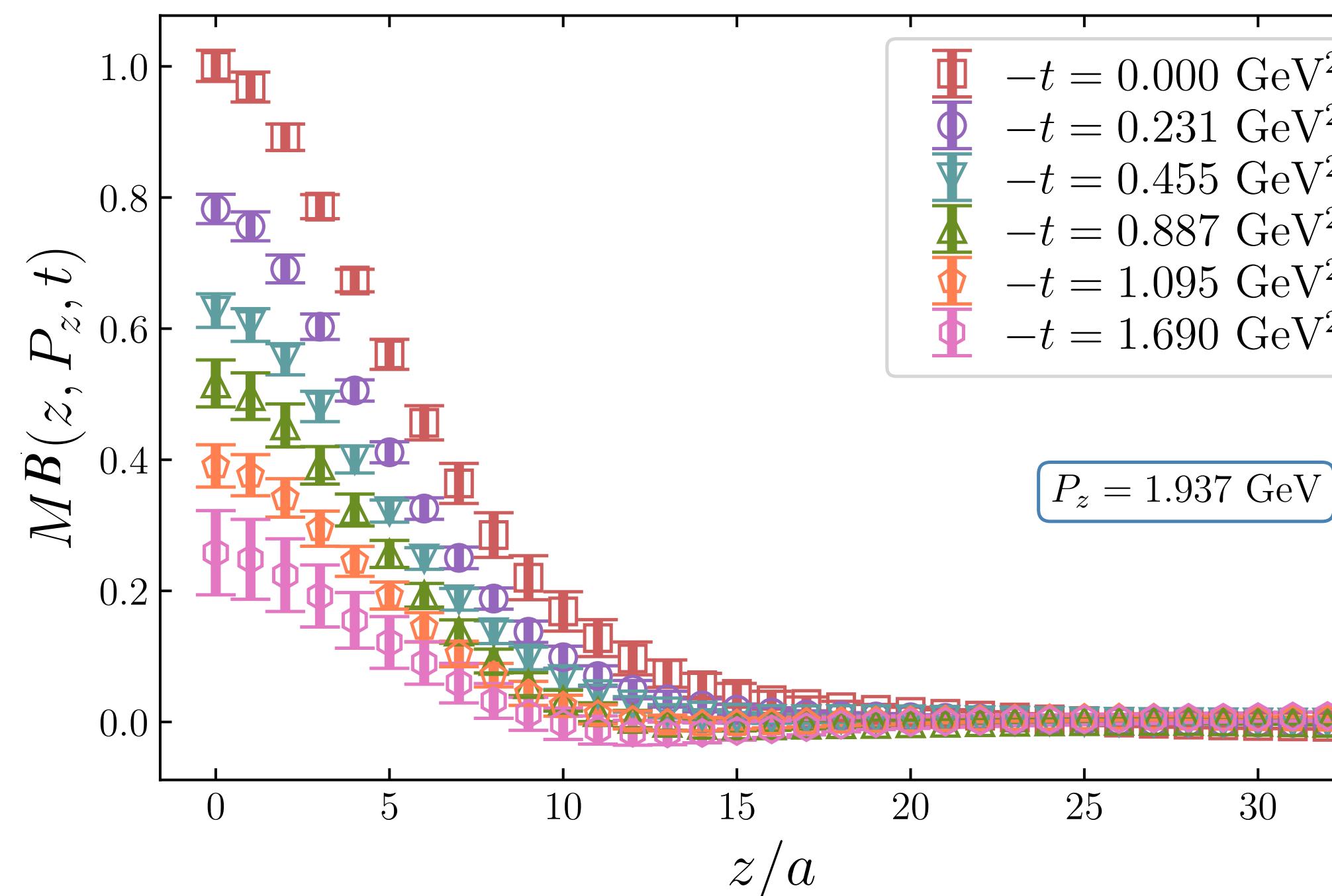
- Lorentz-invariant quasi-GPD \tilde{H}_{LI}

$$\tilde{H}_{\text{LI}}(z \cdot P, z \cdot \Delta, \Delta^2, z^2) = A_1 + \frac{z \cdot \Delta}{z \cdot \bar{P}} A_3 \quad A_3(-z \cdot \Delta) = -A_3(z \cdot \Delta) \quad A_3(z \cdot \Delta = 0) = 0$$

Bare matrix element

$$\tilde{H}_{\text{LI}} = A_1 = KM^t = M^B$$

K : normalization factor



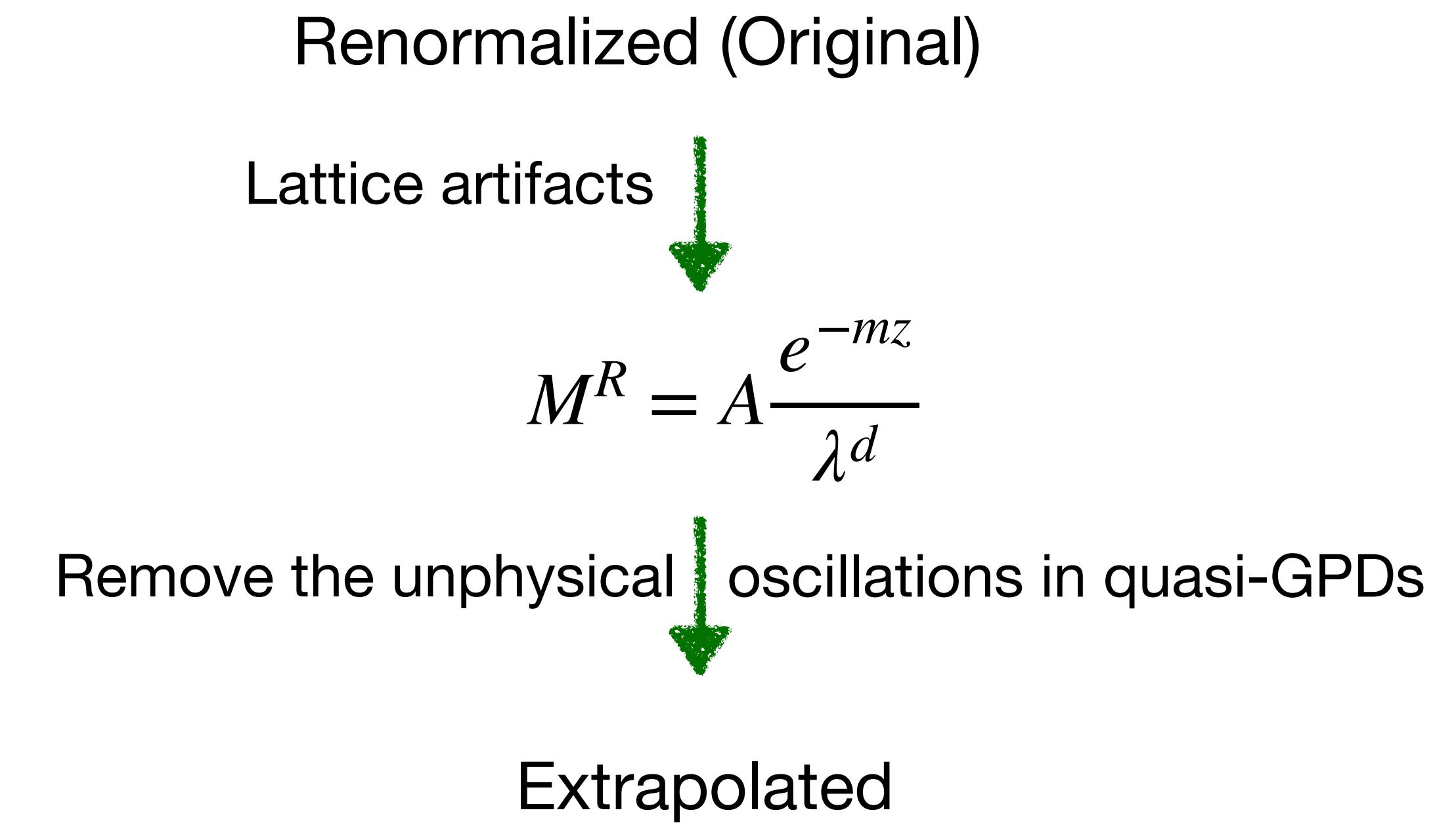
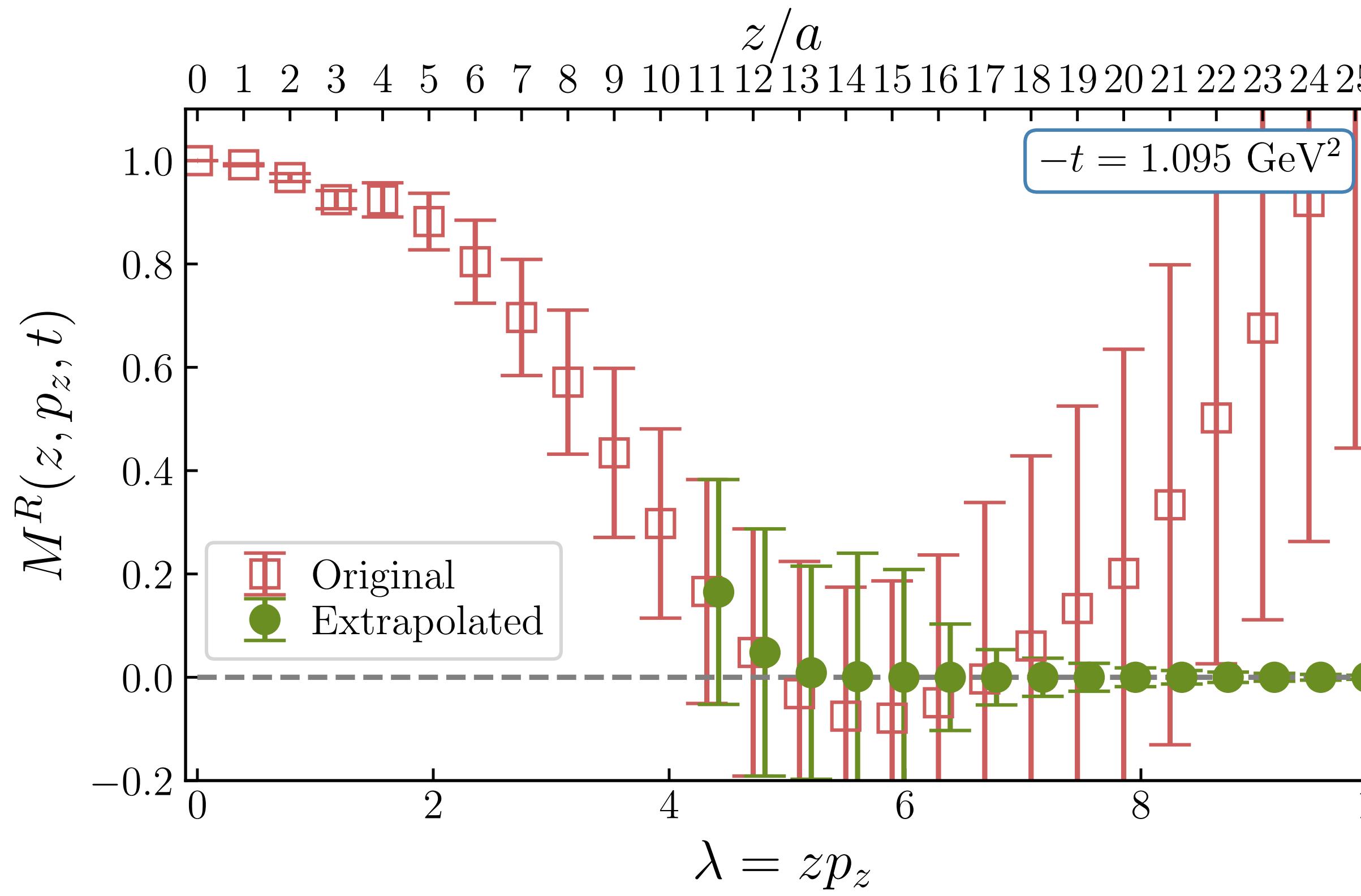


Renormalization & Extrapolation

Ji et al., NPB 964 (2021) 115311

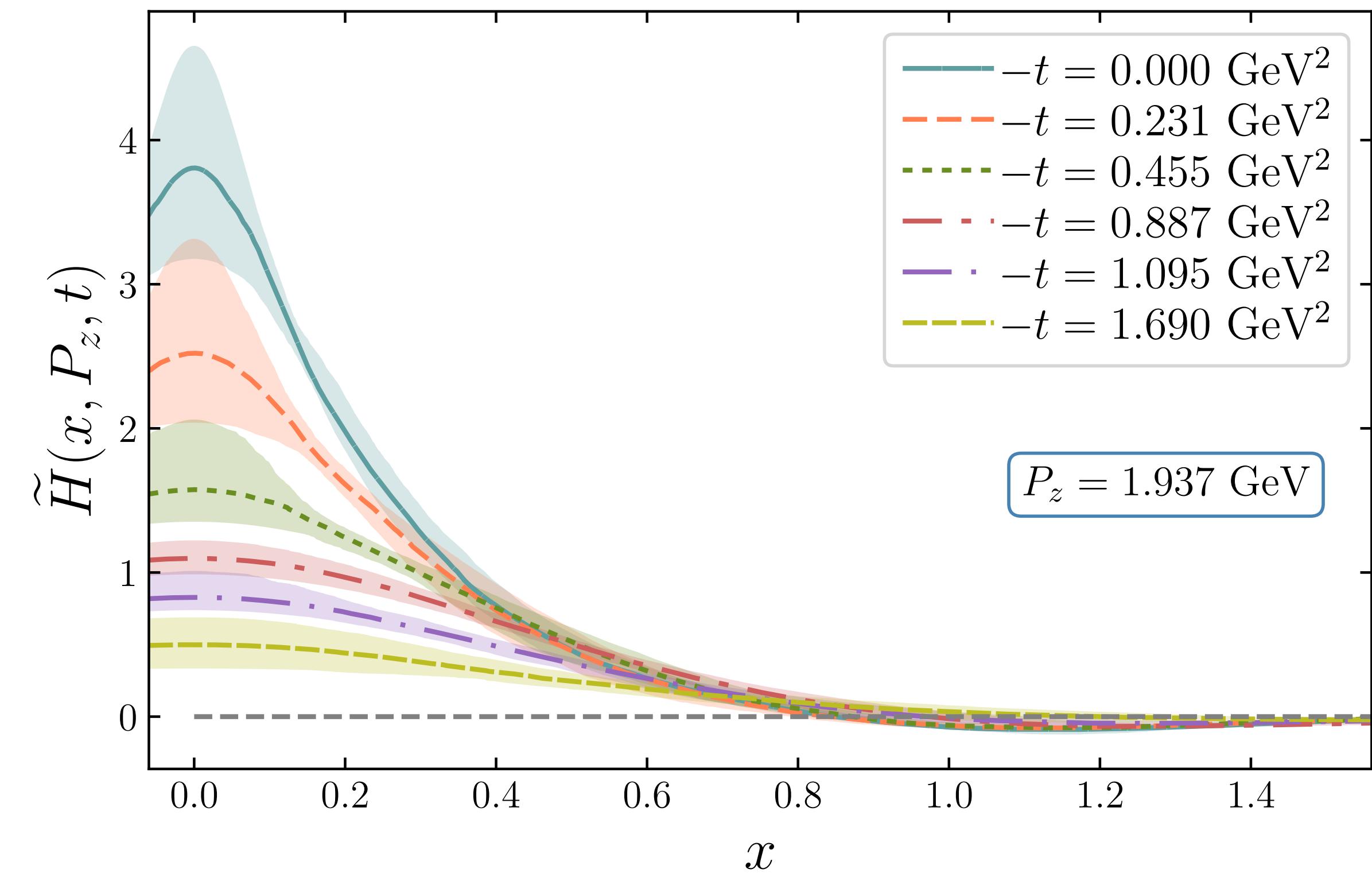
Hybrid scheme,

$$\left\{ \begin{array}{l} z \leq z_S : \frac{M^R(z, \vec{p}, \vec{q})}{M^R(z, 0, 0)} = \frac{M^B(z, \vec{p}, \vec{q})}{M^B(z, 0, 0)}, \text{ Ratio scheme} \\ z \geq z_S : \frac{M^R(z, \vec{p}, \vec{q})}{M^R(z_S, 0, 0)} = e^{(\delta m + \bar{m}_0)|z - z_S|} \frac{M^B(z, \vec{p}, \vec{q})}{M^B(z_S, 0, 0)}. \end{array} \right.$$



#F Quasi-GPDs

$$\tilde{H}(x, P_z, t) = \int \frac{d\lambda}{2\pi} e^{ix\lambda} M^R(z, P_z, t)$$



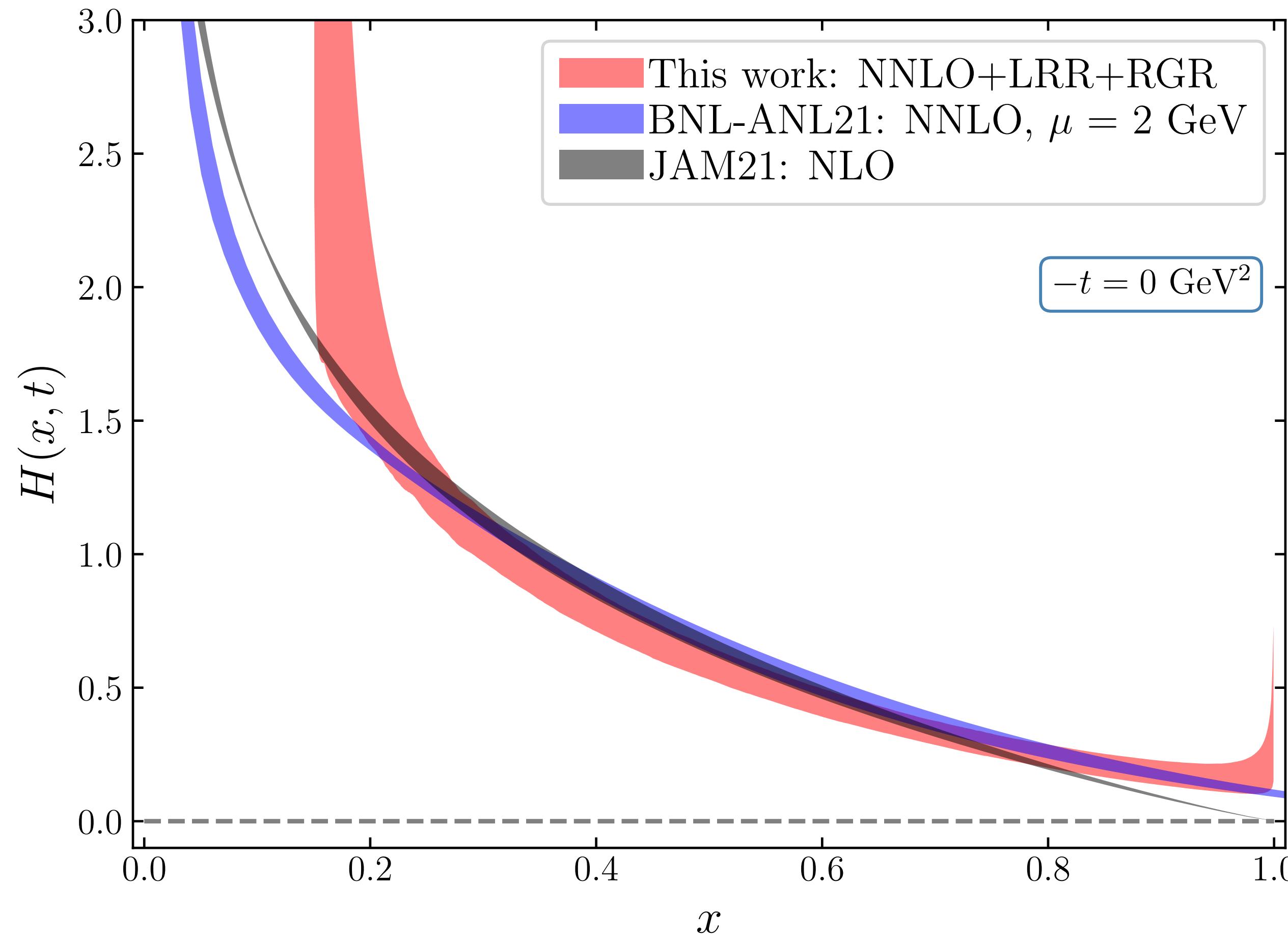
LC GPDs

$$H(x, t) = \int \frac{dk}{|k|} \int \frac{dy}{|y|} C_{\text{evo}}^{-1} \left(\frac{x}{k}, \frac{\mu}{\mu_0} \right) C^{-1} \left(\frac{k}{y}, \frac{\mu_0}{yP_z}, |y|\lambda_S \right) \tilde{H}(y, P_z, t, z_S, \mu_0) + \mathcal{O} \left(\frac{\Lambda_{\text{QCD}}^2}{(xP_z)^2}, \frac{\Lambda_{\text{QCD}}^2}{[(1-x)P_z]^2} \right)$$

NNLO+LRR+RGR ($\kappa = [1, \sqrt{2}, 2]$)

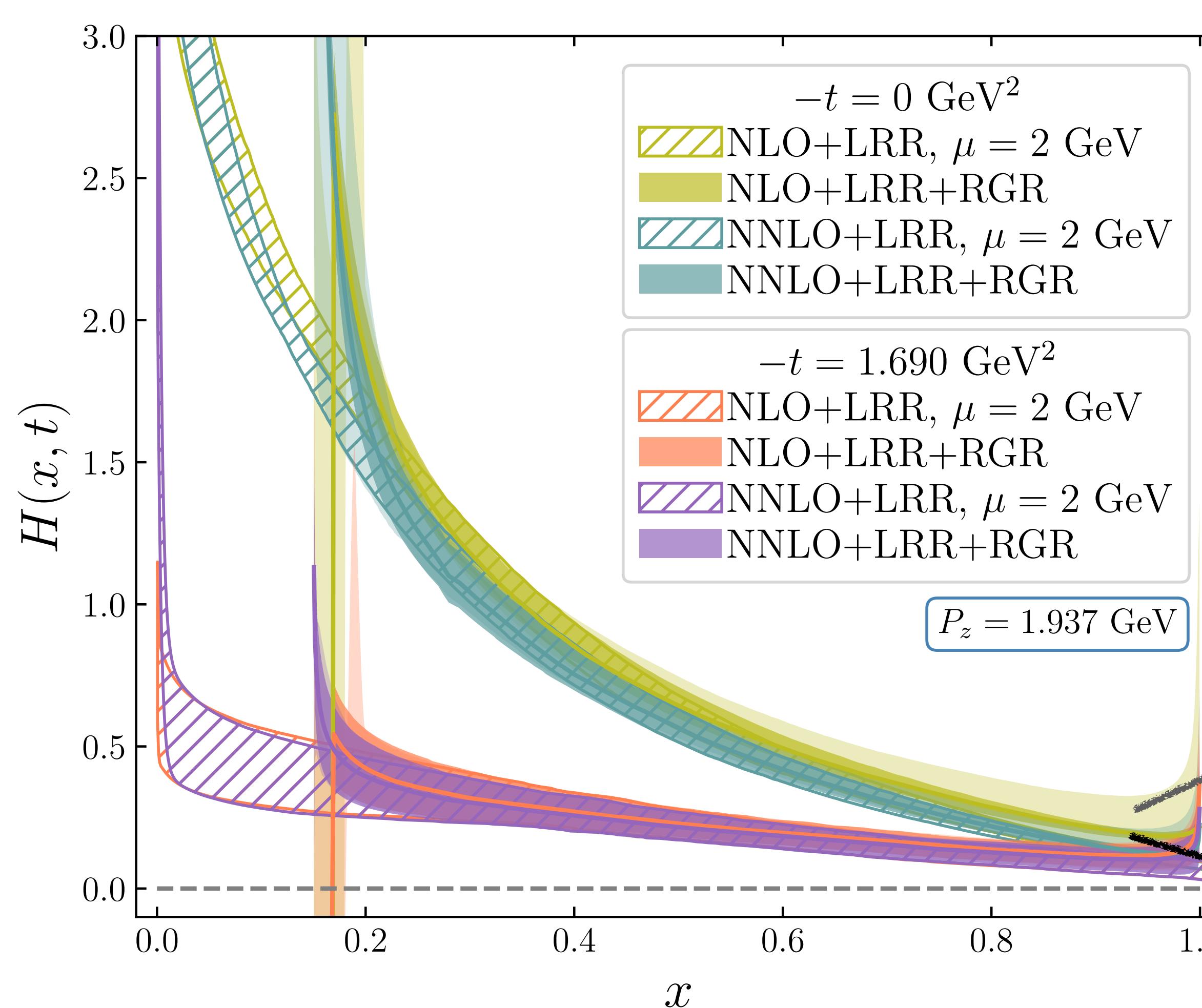
quasi-GPDs

Light-cone GPDs



- ▶ **GPDs at $-t = 0 \text{ GeV}^2$ —— PDFs**
agree well with BNL-ANL21 and JAM21

Light-cone GPDs



► Scale variation

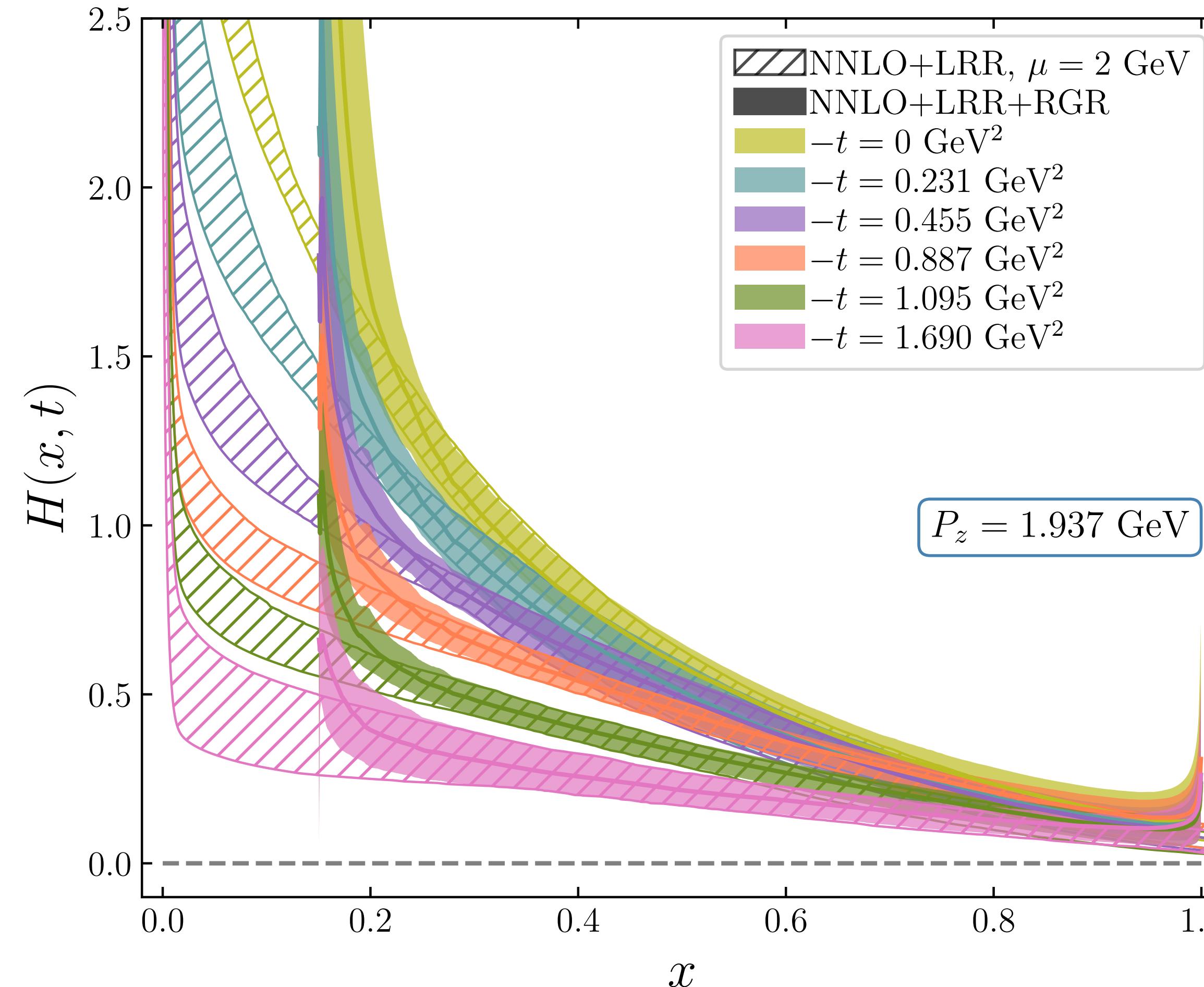
- More significant at small and large x
- Smaller at higher order
- Smaller at larger $-t$

► Good convergence: NLO & NNLO

Systematic error: $1 < \kappa < 2$

Statistical error: $\kappa = \sqrt{2}$

Light-cone GPDs: t -dependence

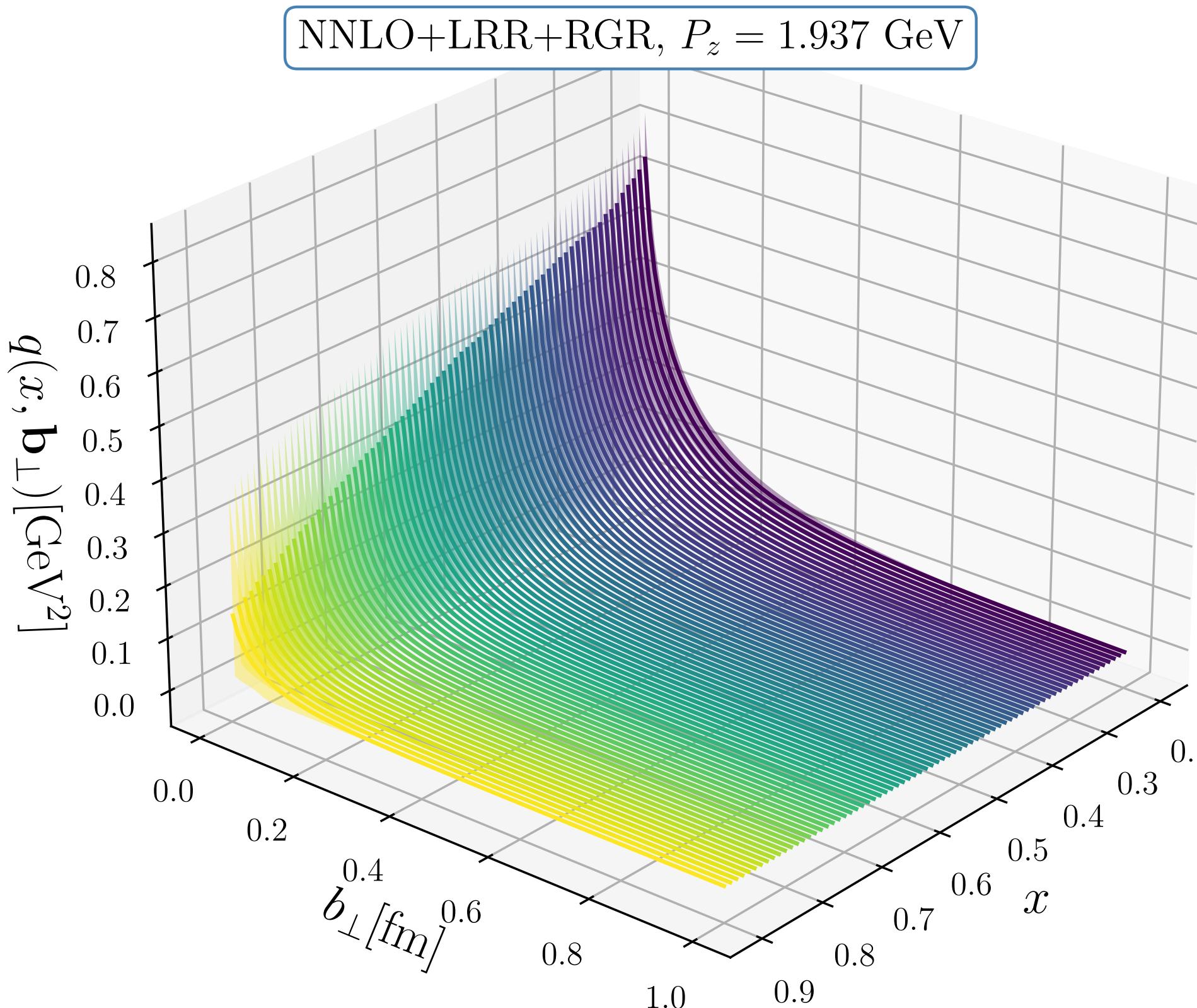


- ▶ At fixed x , H decreases as $-t$ increases
- ▶ The decrease of H along x is slower at larger $-t$
- ▶ RGR: logarithms diverge,
perturbation theory breaks down

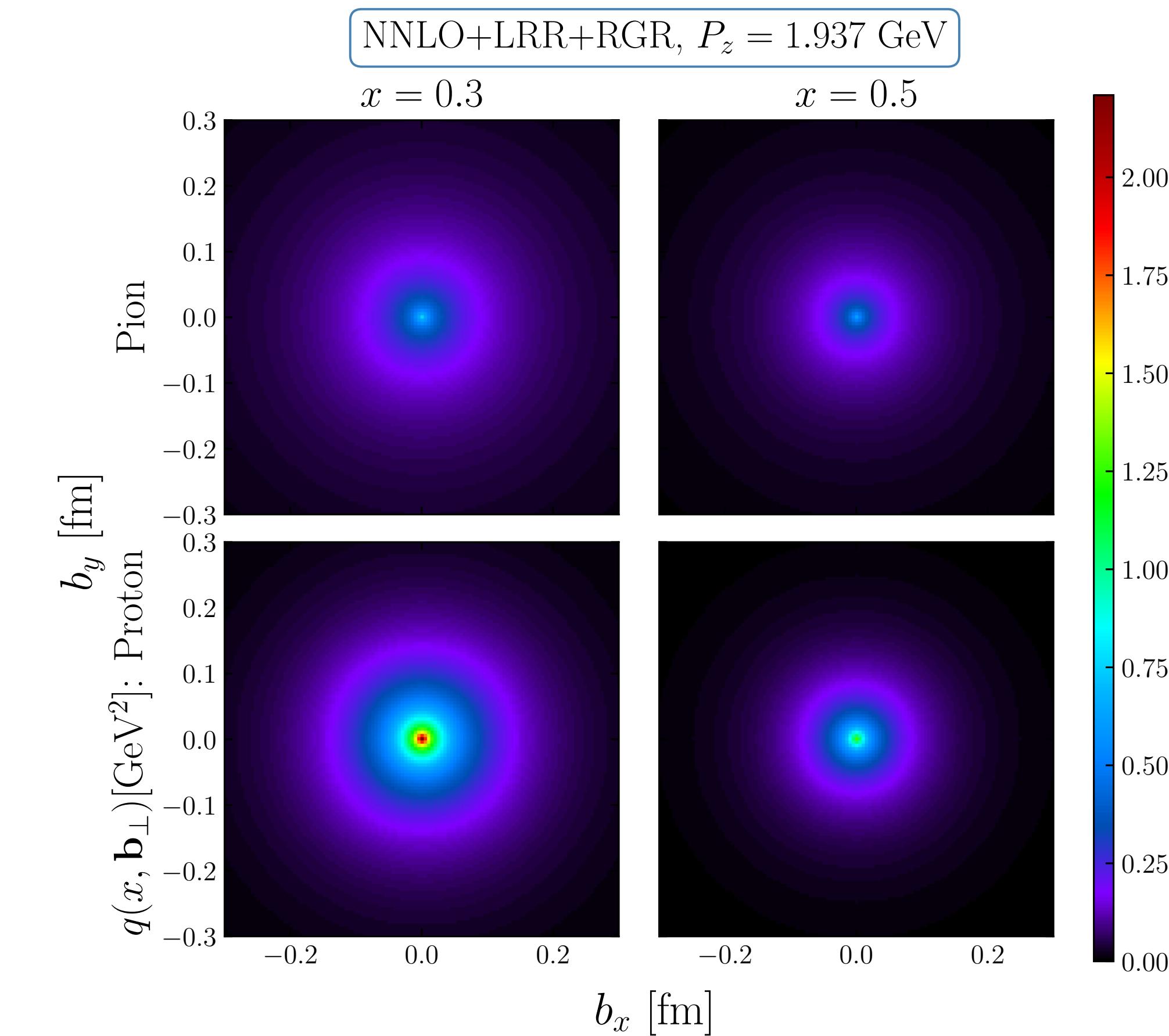


Impact-parameter-space parton distributions (IPDs)

$$q(x, \mathbf{b}_\perp) = \int \frac{d^2 \Delta_\perp}{(2\pi)^2} H(x, \xi = 0, \Delta_\perp^2) e^{i \mathbf{b}_\perp \cdot \Delta_\perp}$$



Quarks with higher x are more concentrated



- Distributions are more concentrated at larger x
- Proton is more broader than pion

Summary

- **Pion and kaon EMFF at the physical point**
 - $-t$ up to 10 and 28 GeV^2 for the pion and kaon
 - Consistent with the existing experimental and the collinear factorization results
 - Serve as benchmark QCD predictions for model-based studies and the future experimental measurements
- **Pion LC GPD in the asymmetric frame**
 - Hybrid-scheme renormalization
 - Matching with NNLO + LRR + RGR
 - t -, x -dependence of the LC GPDs, IPDs
 - Provide a comprehensive three-dimensional imaging of the pion structure

Thanks for your attention!

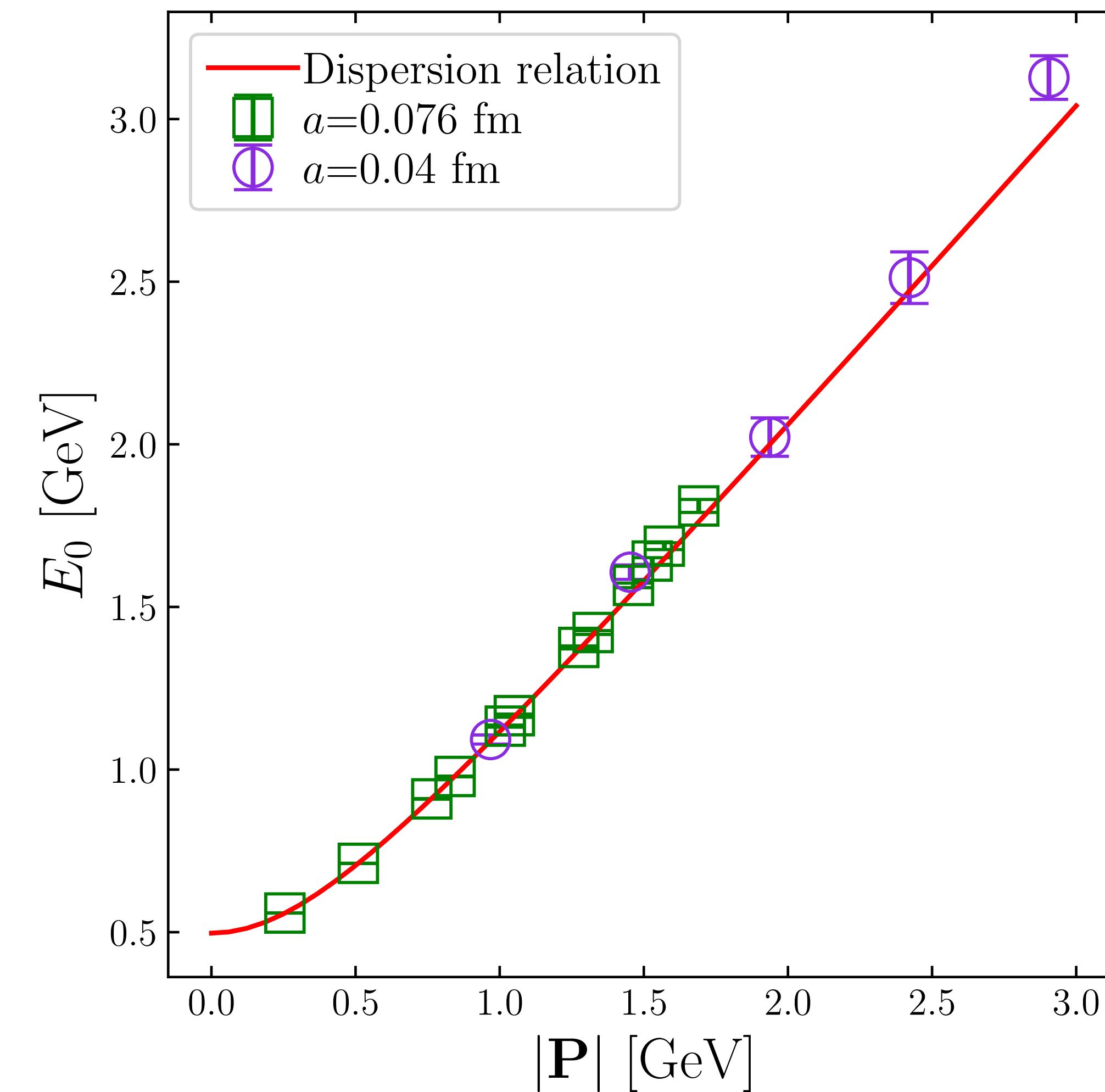
Backup

Lattice Setup

- $N_s^3 \times N_t = 64^3 \times 64$
- Pion / Kaon: $a = 0.076 \text{ fm} / a = 0.04, 0.076 \text{ fm}$
- HISQ action + Wilson-Clover action, **at the physical point**

Meson	$a[\text{fm}]$	$\mathbf{n}_P^f = (n_{P_1}^f, n_{P_2}^f, n_{P_3}^f)$	$n_{k_3}^f$	$\mathbf{n}_P^i = (n_{P_1}^i, n_{P_2}^i, n_{P_3}^i)$	$n_{k_3}^i$	$Q^2[\text{GeV}^2]$	(#ex, #sl)
Pion	0.076	(0, 0, -3)	-2	(0,0,2)		1.56	
				(0,0,3)		2.34	
				(0,0,4)		3.12	
				(2,0,3)		2.58	
		(0, 0, -5)	-4	(0,0,3)		3.90	
				(0,0,4)		5.20	
				(0,0,5)	4	6.50	(7, 224)
				(2,0,4)		5.50	
				(2,0,5)		6.75	
	0.04	(0, 0, -6)	-5	(0,0,5)		7.80	
				(0,0,6)		9.35	
				(2,0,5)		8.10	
				(2,0,6)		9.61	
							(18, 576)

Extract Energy and Amplitude



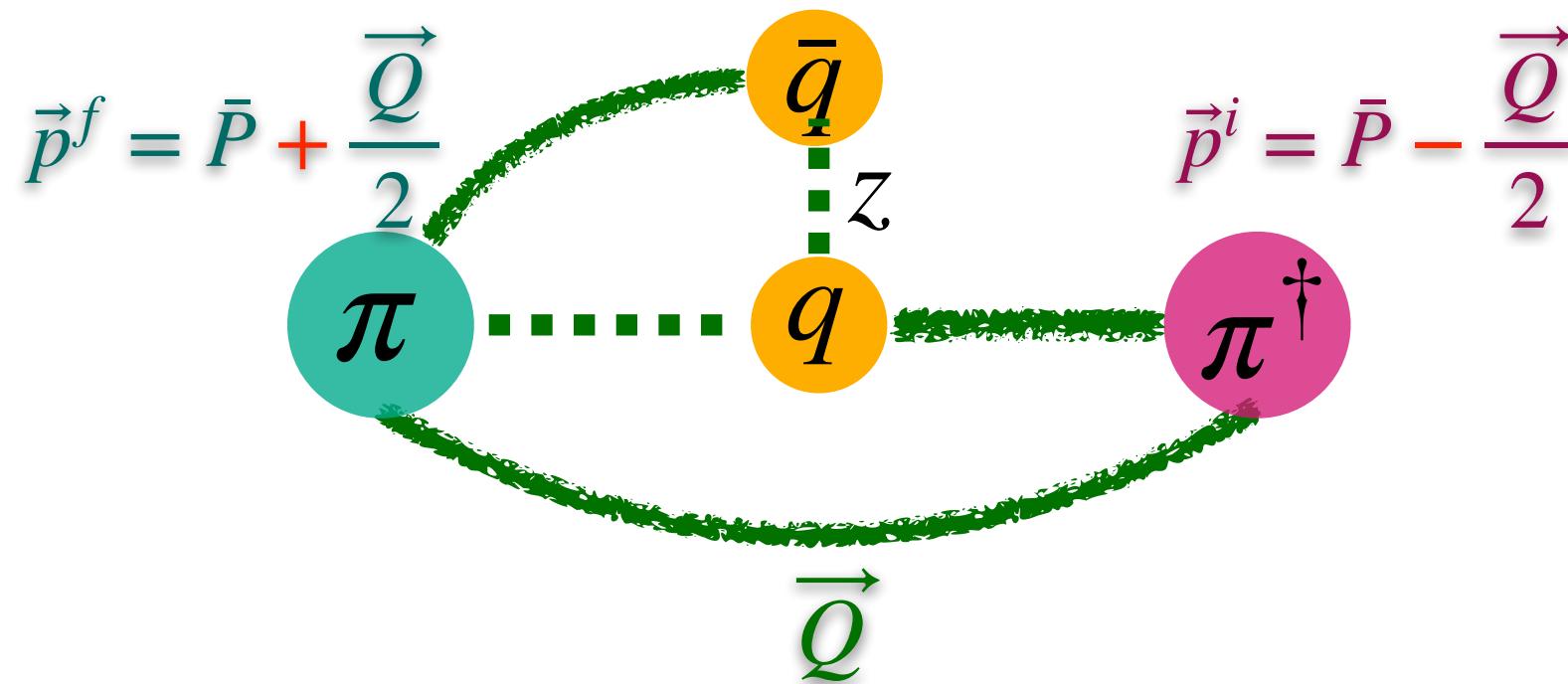


Frame-independent approach

$$N_{\text{color}} \otimes N_{\text{flavor}} \otimes N_{\text{spin}} \otimes N_{\text{space}}^3 \otimes N_{\text{time}} \gtrsim 10^9$$

Lattice: each \vec{p}^f requires a separate calculation -> Goal: reduce the number of \vec{p}^f

Traditional: Symmetric



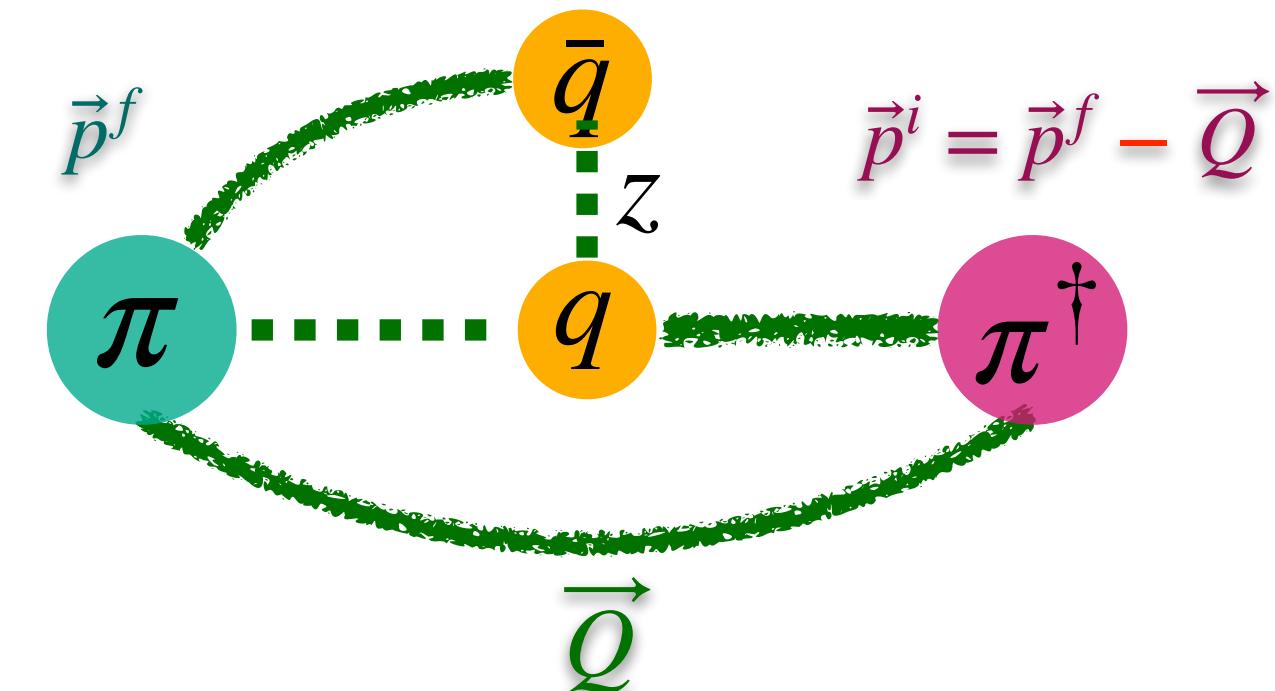
- One \vec{p}^f – only one $Q^2 = -t$ is useful
- Each Q^2 requires a separate calculation

Vary \vec{p}^f in several calcs



Newly proposed: Asymmetric

Bhattacharya, Constantinou et al., PRD 106 (2022)



- One \vec{p}^f – several Q^2 are useful

Fix \vec{p}^f , vary Q^2 in one calc

Computational cost



Lattice Setup

- $N_s^3 \times N_t = 64^3 \times 64$, $a = 0.04$ fm
- HISQ action + Wilson-Clover action $\Rightarrow m_\pi^{\text{val}} = 0.3$ GeV
- Using boost smearing to enhance the signal
- Momentum transfer Q^2 : $0 \sim 1.7$ GeV 2

Frame	t_s/a	$\mathbf{n}^f = (n_x^f, n_y^f, n_z^f)$	m_z	$P_z[\text{GeV}]$	$\mathbf{n}^\Delta = (n_x^\Delta, n_y^\Delta, n_z^\Delta)$	$-t[\text{GeV}^2]$	#cfgs	(#ex, #sl)
Breit	9,12,15,18	(1, 0, 2)	2	0.968	(2, 0, 0)	0.938	115	(1, 32)
non-Breit	9,12,15,18	(0,0,0)	0	0	(0,0,0)	0	314	(3, 96)
	9,12,15,18	(0,0,2)	2	0.968	(1,2,0)	0.952	314	(4, 128)
	9,12,15	(0,0,3)	2	1.453	$[(0,0,0), (1,0,0)$ $(1,1,0), (2,0,0)$ $(2,1,0), (2,2,0)]$	$[0, 0.229, 0.446,$ $0.855, 1.048, 1.589]$	314	(4, 128)
	9,12,15	(0,0,4)	3	1.937		$[0, 0.231, 0.455,$ $0.887, 1.095, 1.690]$	564	(4, 128)