



Modeling of coasting beams in Xsuite

G. Iadarola

Acknowledgements:

H. Bartosik, R. De Maria, S. Lopaciuk, K. Paraschou



- **Introduction**
 - Particle slippage
 - Implications for collective effects
- **Simulation method description**
 - Reference speed β_{sim}
 - From turns to frames
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 - Justification of the method
 - Generalization of ζ coordinate
 - ζ update at start turn
 - ζ update on frame jump
- **Numerical tests**

In a ring the particles **revolution frequency depends on the on the particle energy.**

To first order, this dependence is described by the **slip factor**:

$$\Delta f_{\text{rev}} = -\eta f_0 \frac{\Delta P}{P_0}$$

In **bunched beams**, the momentum of the particles **oscillates around the equilibrium momentum**:

- The **average revolution period** (over a synchrotron period) is the **same for all particles**
- The **differences in time of arrival among particles never exceeds one revolution period**
- Over a given time interval, **all particles make the same number of turns**
 - The turn index can be used as independent variable in the simulation

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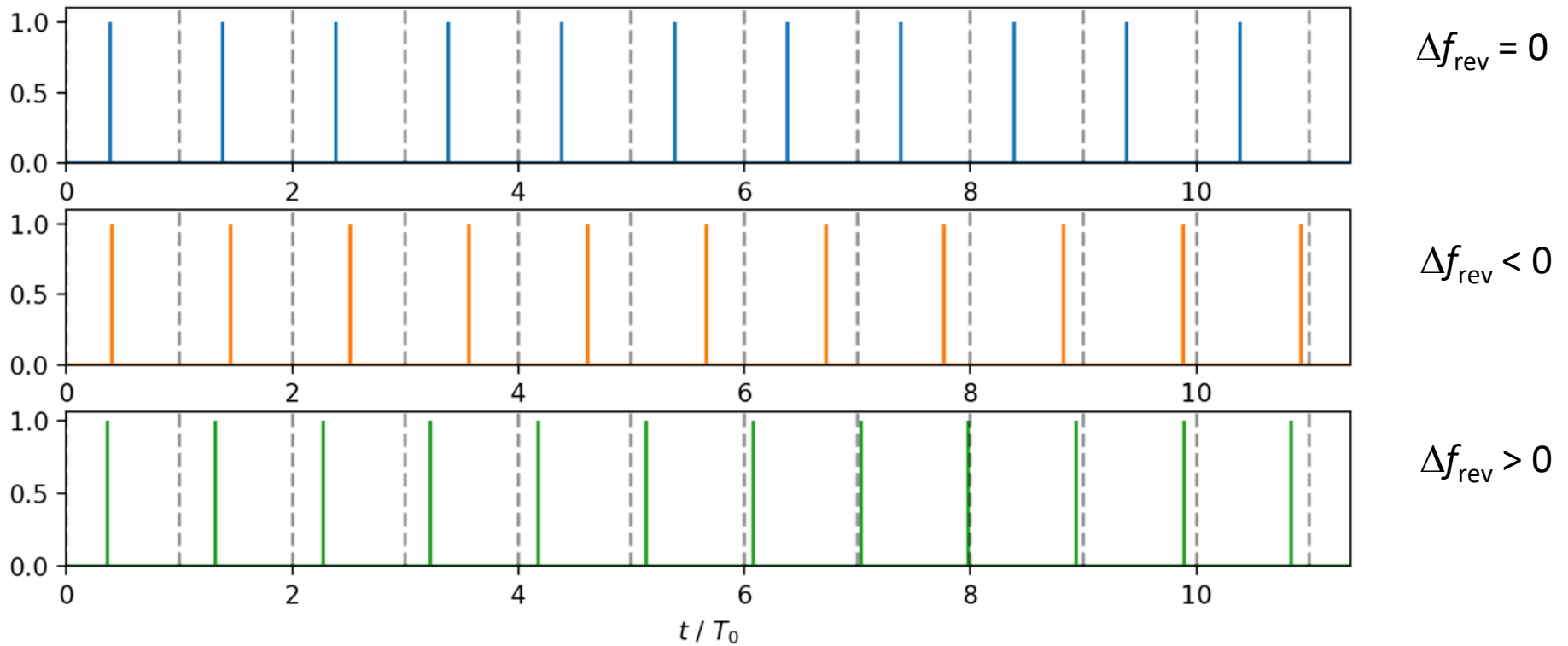
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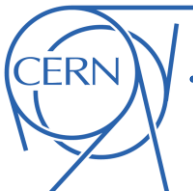
In **coasting beams**, there is no RF focusing:

- Particles **keep their momentum deviation indefinitely**
- Hence, **differences in revolution frequency are also kept indefinitely**
- Over a given time interval, **different particles perform a different number of turns**

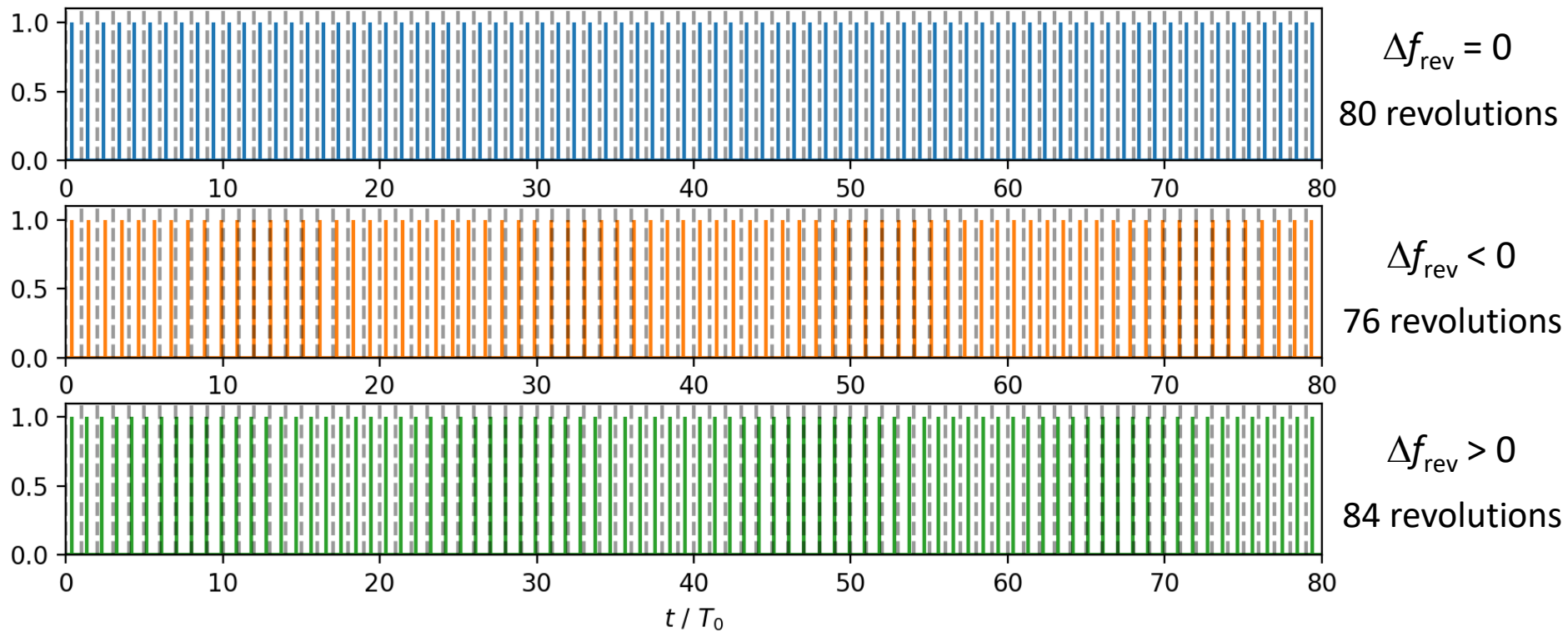


- To visualize this effect, in the case of a coasting beam, we **compare three particles having three different revolution frequencies** by marking the times at which they are detected at a given location of the ring

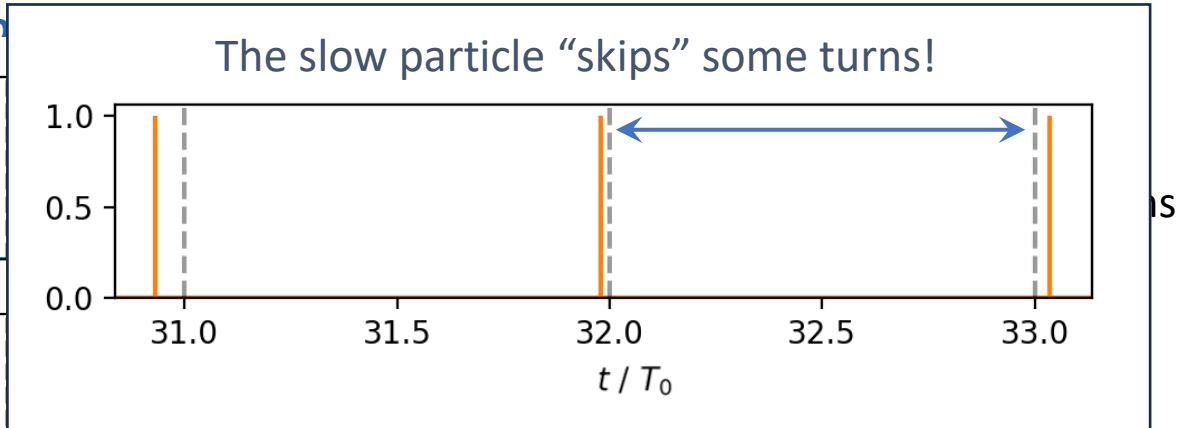
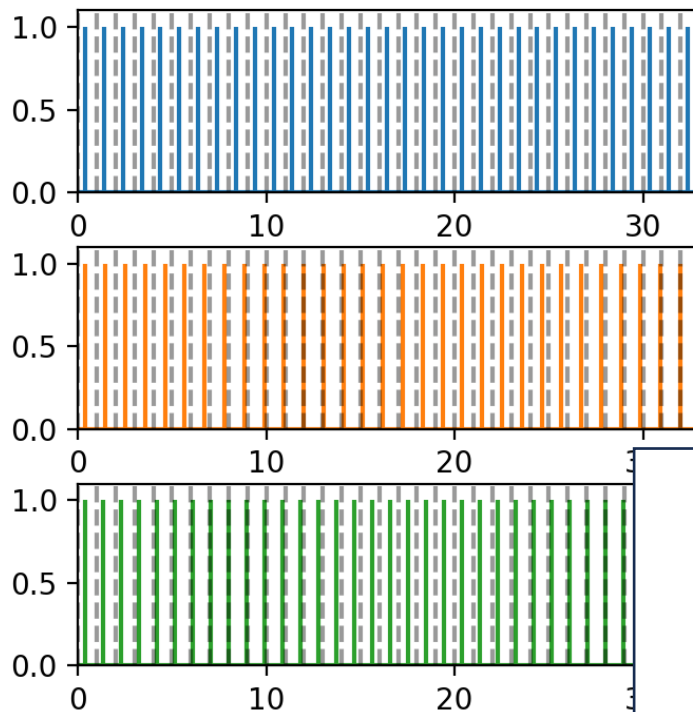




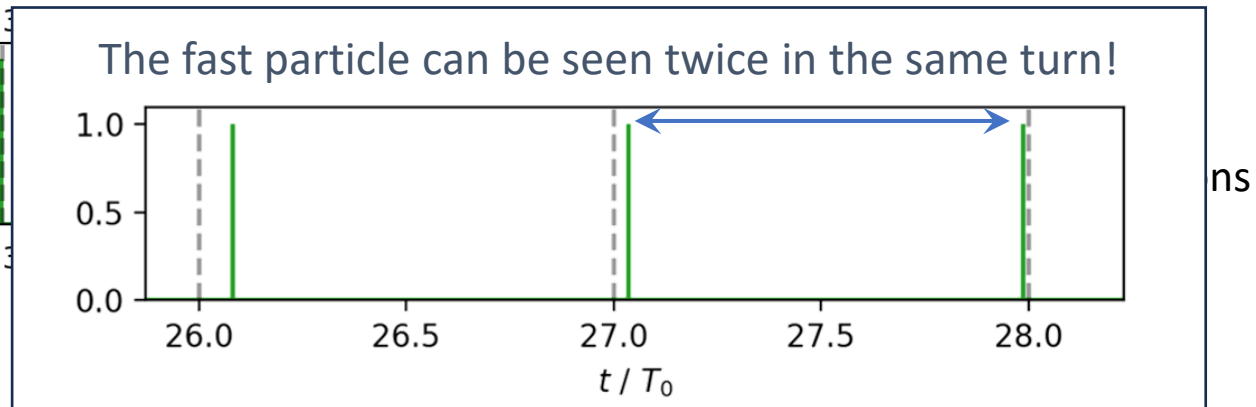
- To visualize this effect, in the case of a coasting beam, we **compare three particles having three different revolution frequencies** by **marking the times at which they are detected at a given location of the ring**
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- To visualize this effect, in the case of a coasting beam, we **compare three particles having three different revolution frequencies** by marking the times at which they are detected at a given location of the ring
 - Over a given time interval, the three particles **perform a different number of revolutions**



76 revolutions



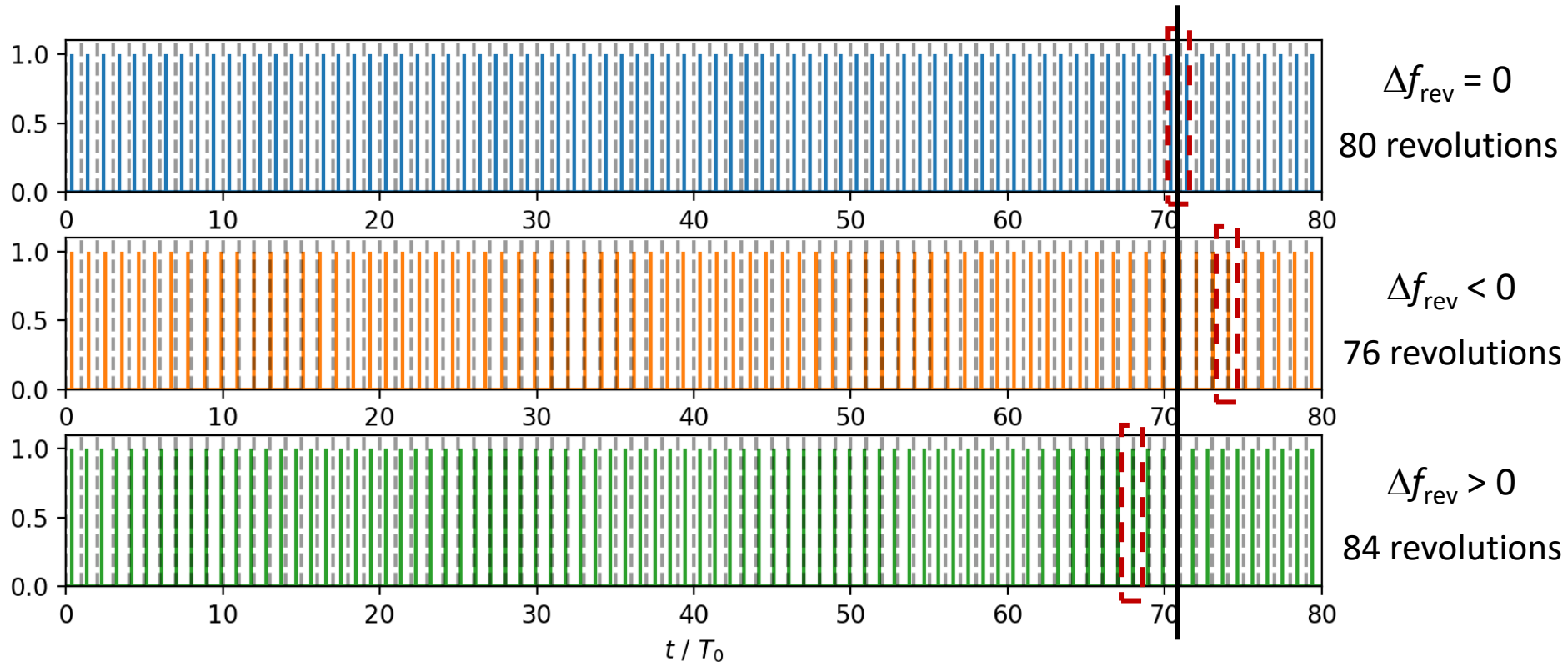


In **conventional tracking simulations** we use the **longitudinal coordinate s** and the **turn number as independent variables** of our simulations

- At a given simulation stage **all particles have travelled the same distance**, but in the case of coasting beams **their arrival time can be several turns apart**

This creates an **issue for the simulation of collective effects**.

- For example, to compute **space charge** forces on a particle we **need to know the position of all the other particles at the same time**. Same for getting distribution moments in **wakefield simulations**.





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One **radical solution to this problem** would be to change simulation approach and **use time as independent variable**:

- While conceptually simple, it would **extremely heavy in terms of development effort and very expensive in terms of follow-up and maintenance** (it's basically another code)
- Instead, we have **devised a method** allowing us to **reuse without changes** our conventional simulation modules (tracking elements, space-charge, wakefields, etc.) while **achieving the required synchronization for collective effects**
 - The **implementation** turns out to be quite **simple** and **confined in a dedicated module** (great advantage to preserve Xsuite extensibility and maintainability)
 - The **derivation is a bit cumbersome**, so bear with me...



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We define for each particles:

$$\hat{\beta}(s_1, s_2) = \frac{1}{c} \frac{s_2 - s_1}{t(s_2) - t(s_1)} \quad \text{Speed measured on the reference trajectory}$$

Then we choose a **"reference speed"** β_{sim} such that for all particles

$$\hat{\beta}(s_1, s_2) < \beta_{\text{sim}} \quad \text{i. e. No particle moves faster than } \beta_{\text{sim}}$$

$$\hat{\beta}(s_1, s_2) > \frac{\beta_{\text{sim}}}{2} \quad \text{i. e. No particle moves slower } \beta_{\text{sim}} / 2$$

Easy to find: $\beta_{\text{sim}} = 1.1 \beta_0$ works practically for any storage ring

In a nutshell, what we will do in the following is to **use β_{sim} as reference velocity**

→ **Particles can be too slow** (skip a turn) **but not too fast** (arrive twice in a turn)

Note that in the following of this presentation, if we replace β_{sim} with β_0 we get back the definitions and relations used for bunched beams.

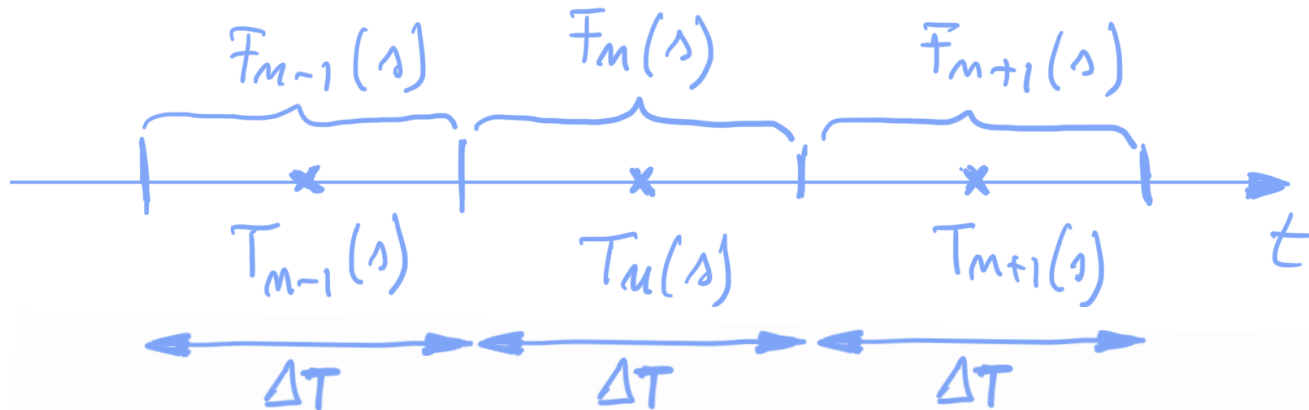


From turns to time frames

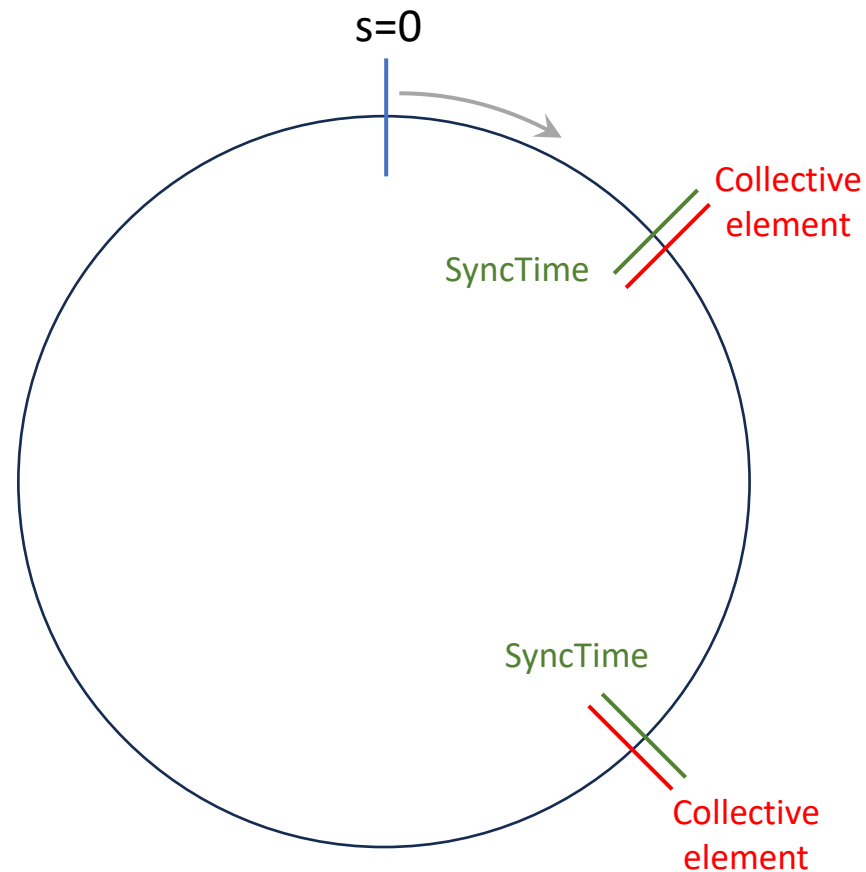
We use a “virtual particle” moving at the “reference speed” β_{sim} on the reference trajectory to **generalize the concept of turn**:

- We call **frame** a time interval of length: $\Delta T = \frac{L}{\beta_{\text{sim}} c}$
- The **mid point of each frame** depends on the s position and is defined by the **passage of the virtual particle at the given s** :

$$T_n(s) = n\Delta T + \frac{s}{\beta_{\text{sim}} c}$$

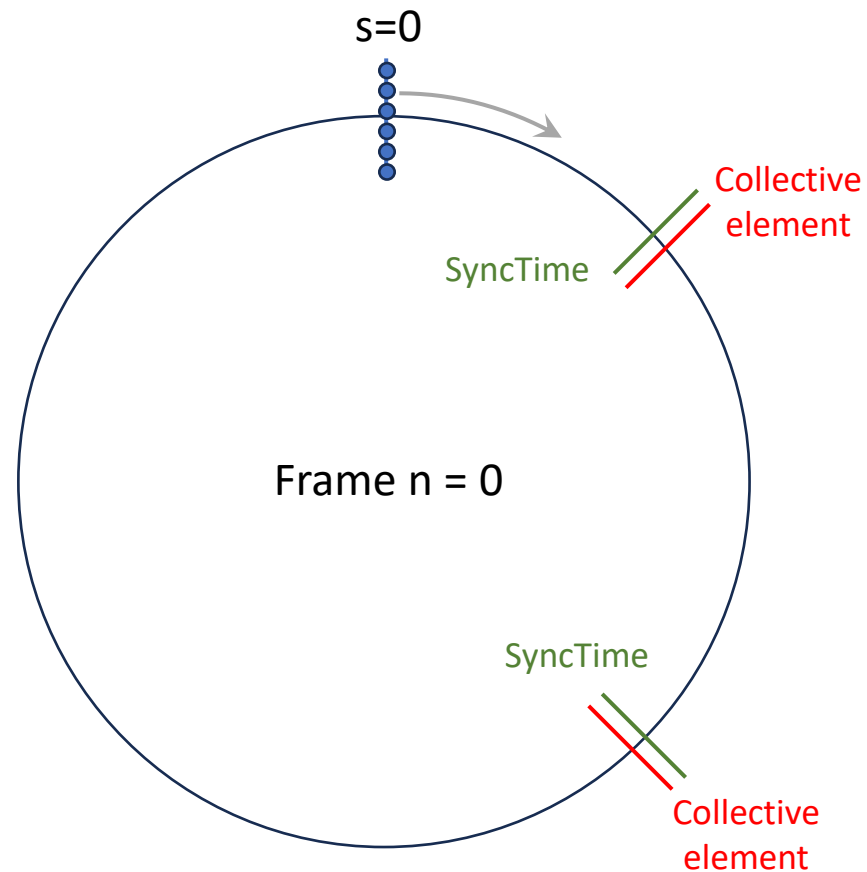


- The **time frames act as the turns** of the bunched beam case:
 - For **each frame we track particles around the ring**
 - We **compute forces due to collective effects only for the present frame**
- We need to ensure that **at each collective interaction we pass to the collective elements all particles arriving during the present time frame** (independently of their number of performed revolutions)
- For this purpose, **we install in front of each collective element, a “SyncTime” element**
 - The SyncTime takes care of **disabling particles that fall out of the current time frame** (arrive too late) and re-enabling them at the following frame



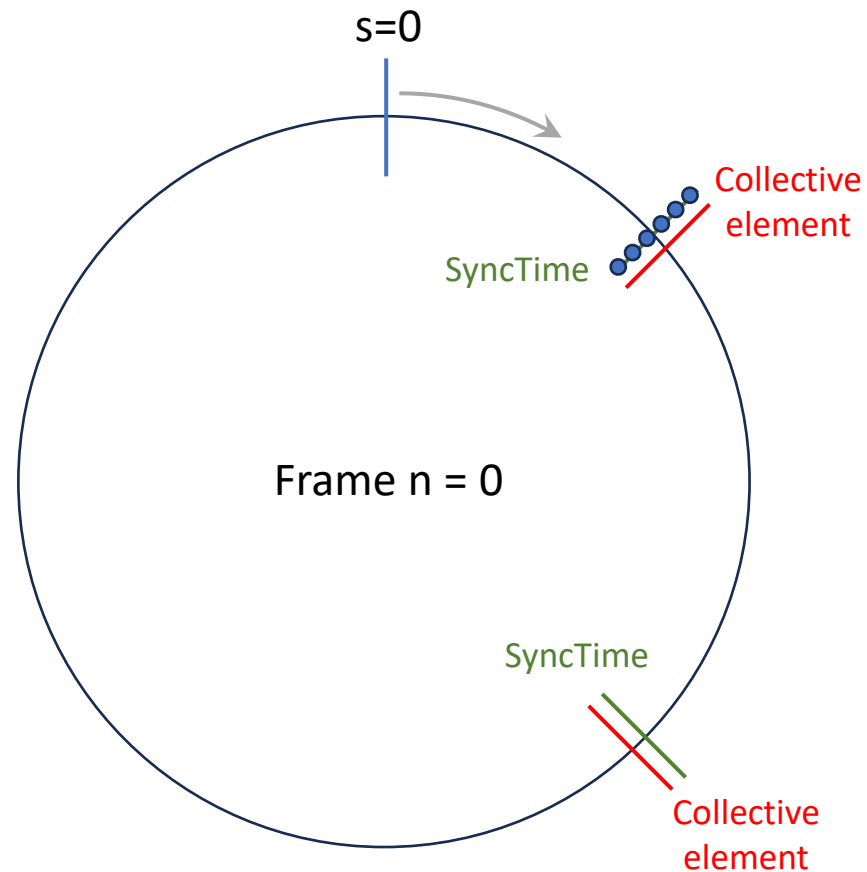
We start by simulating frame $F_0(s)$:

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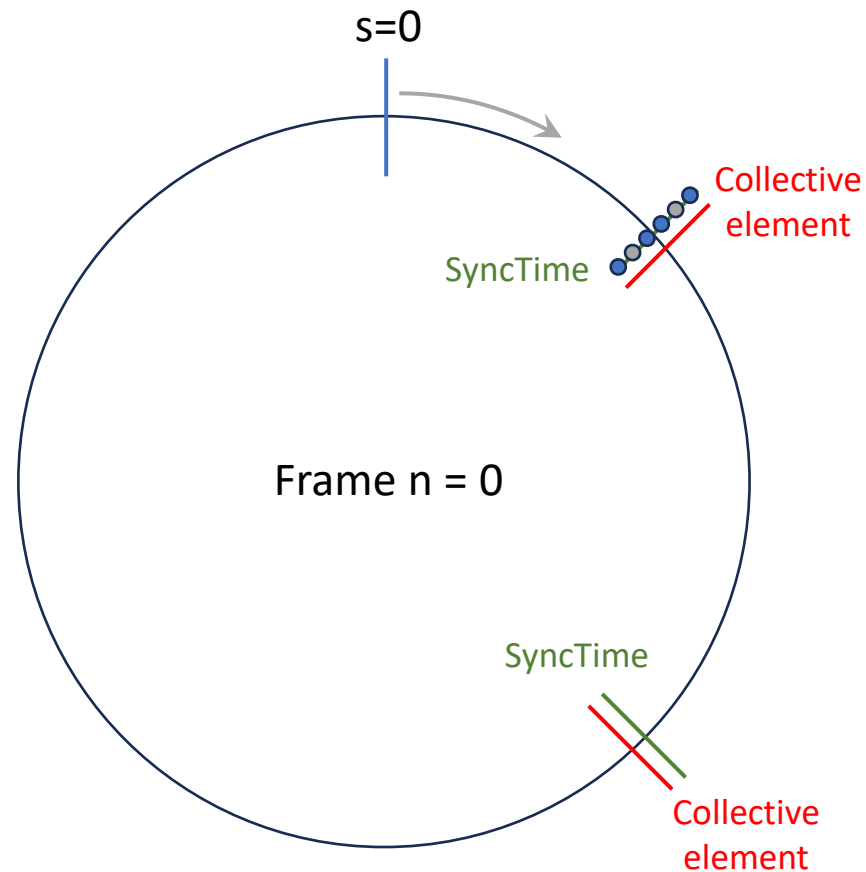
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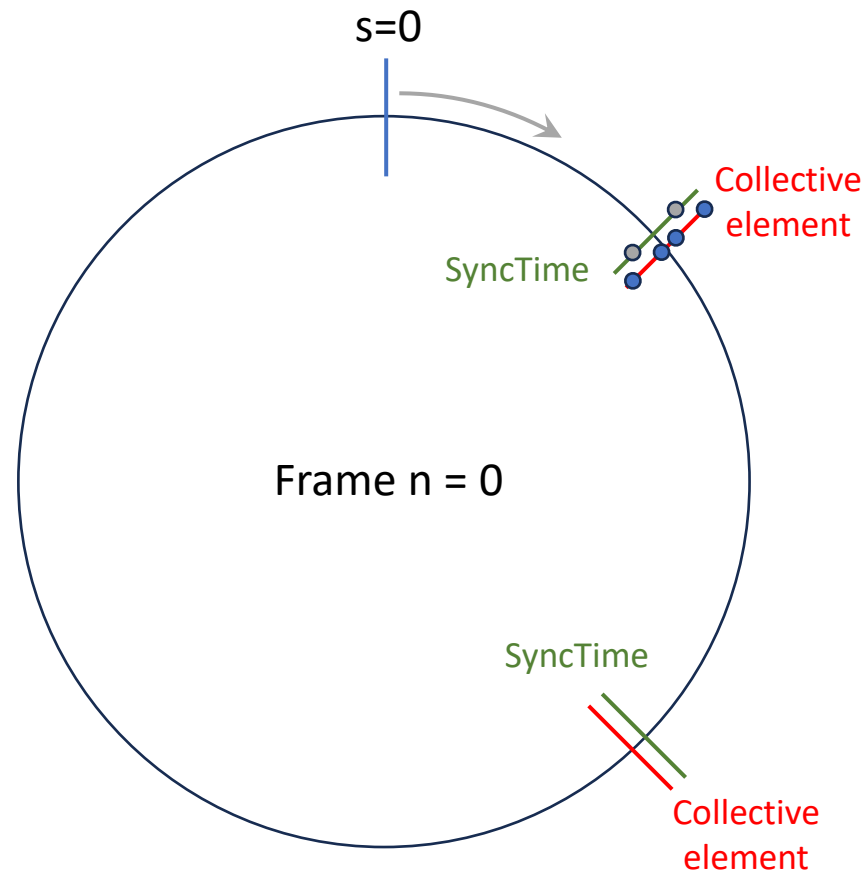
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 → these particles are **deactivated**



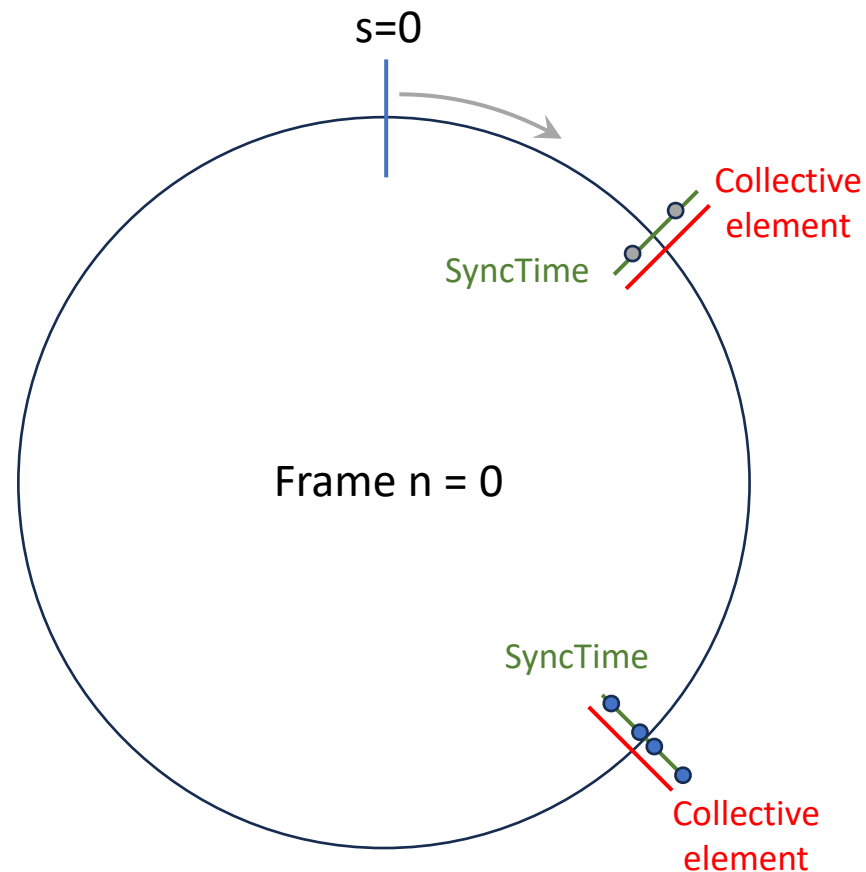
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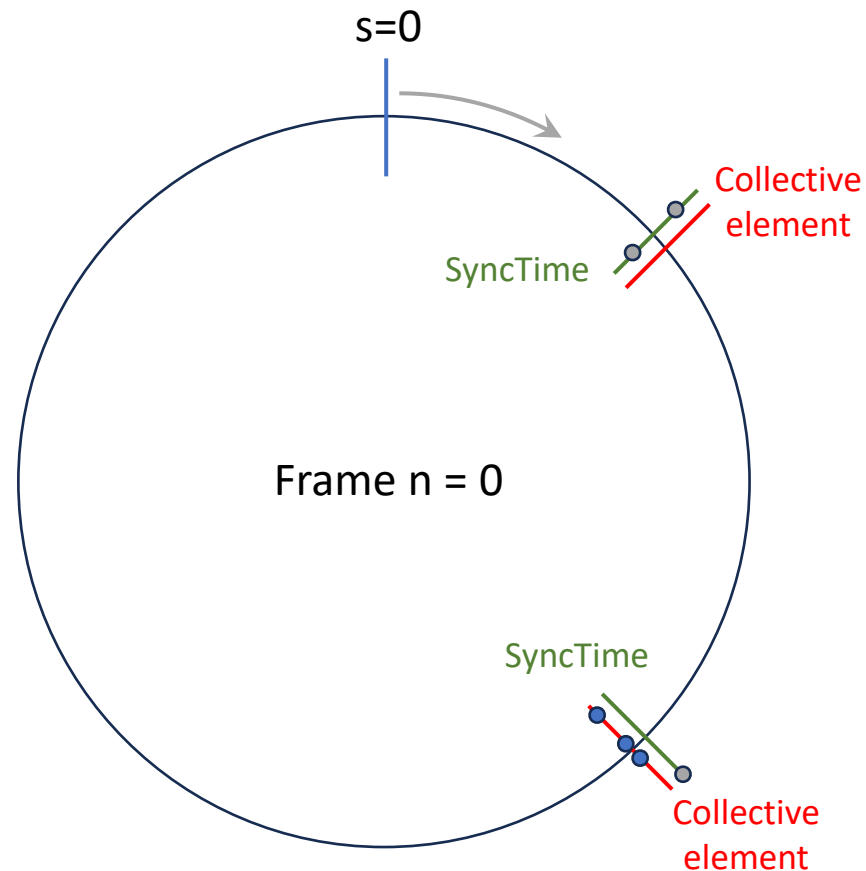
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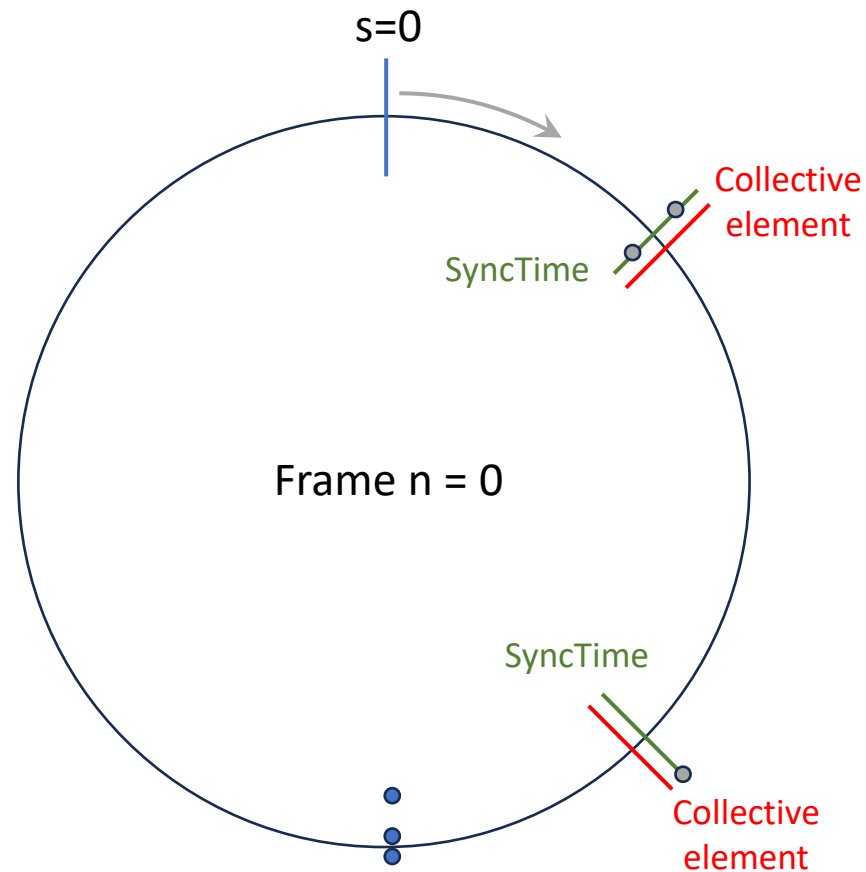
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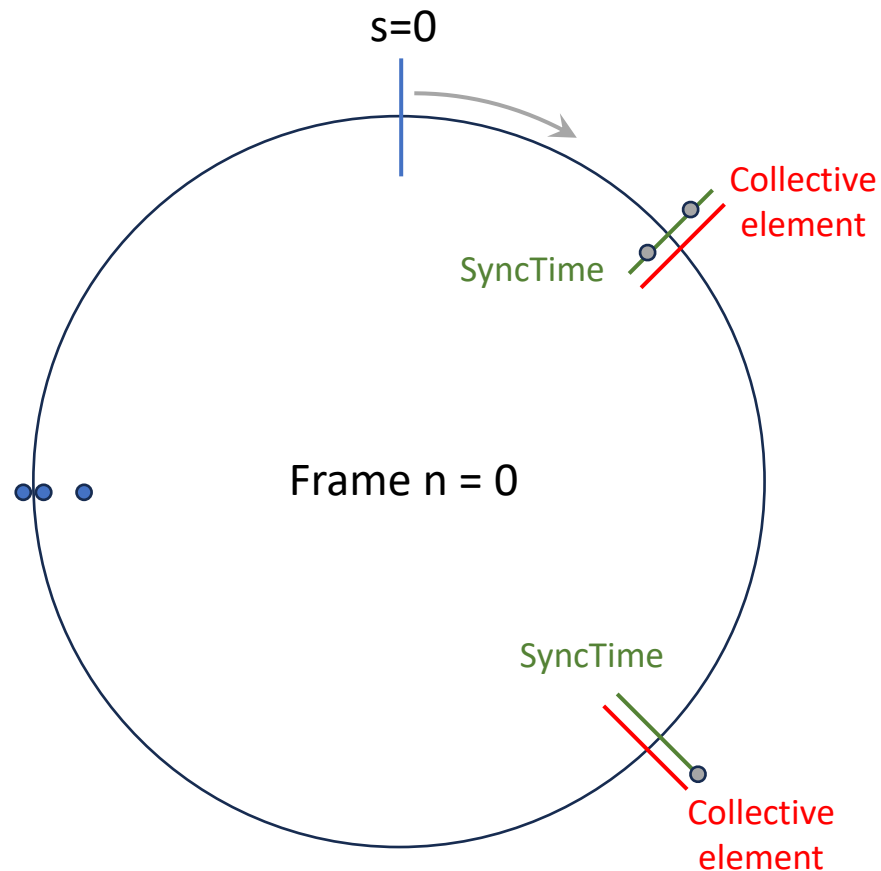
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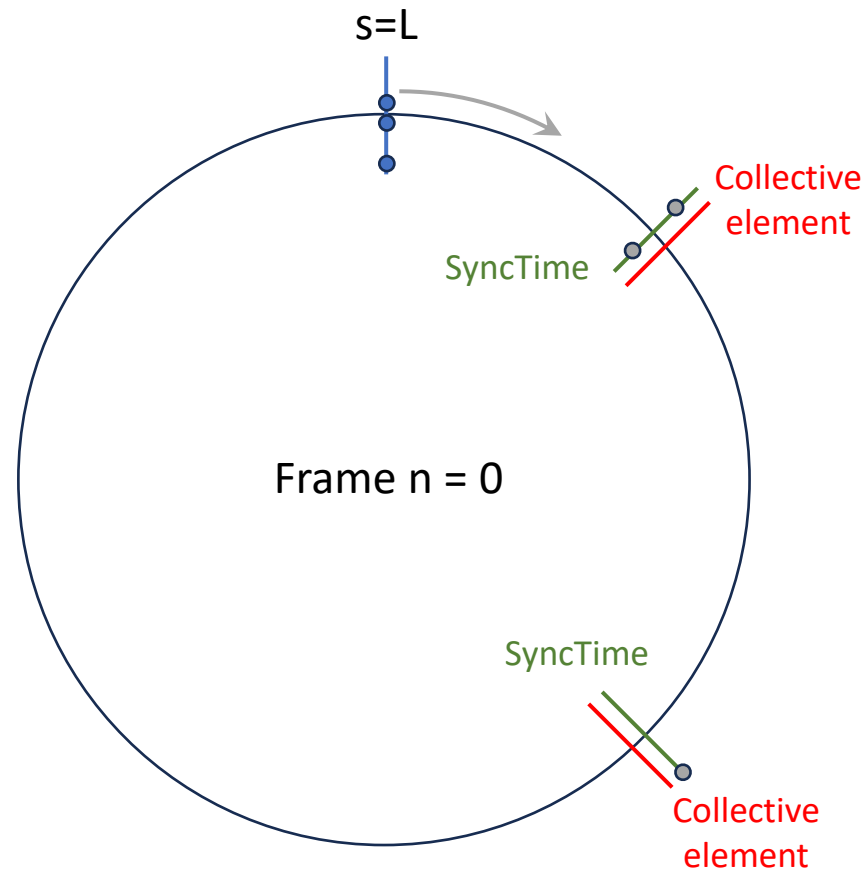
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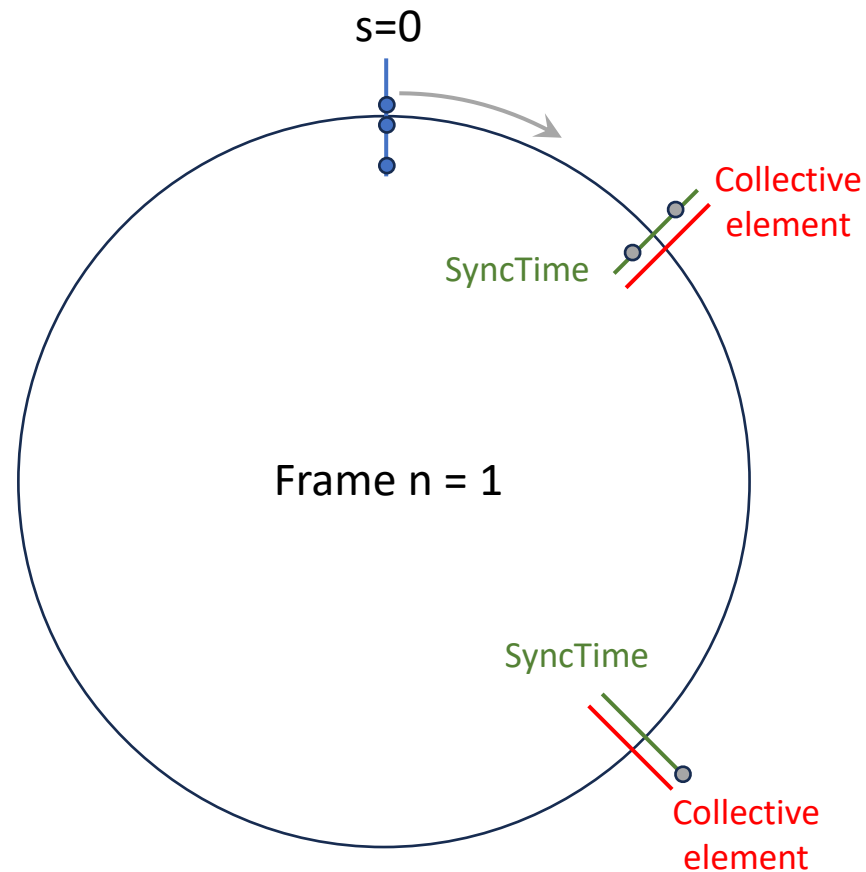
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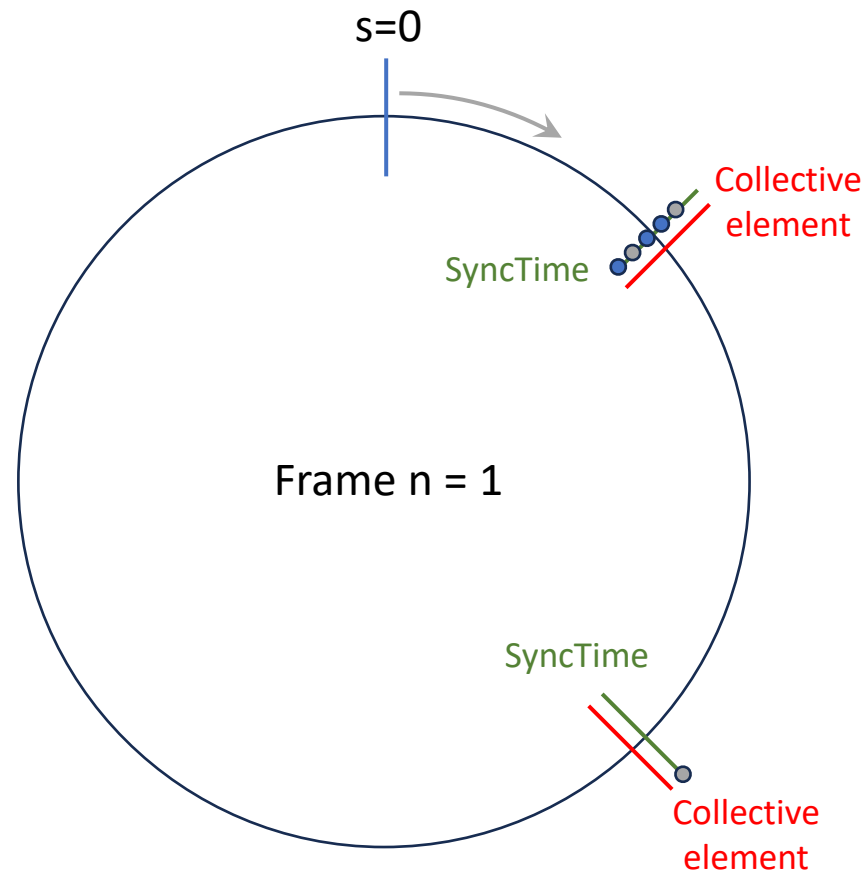
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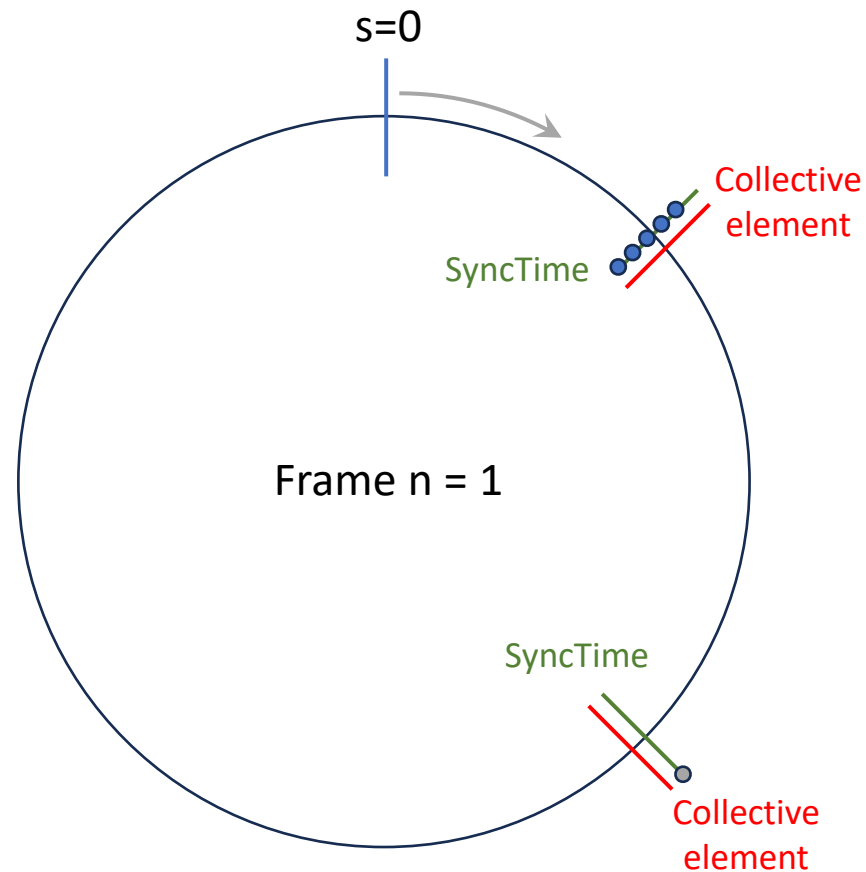
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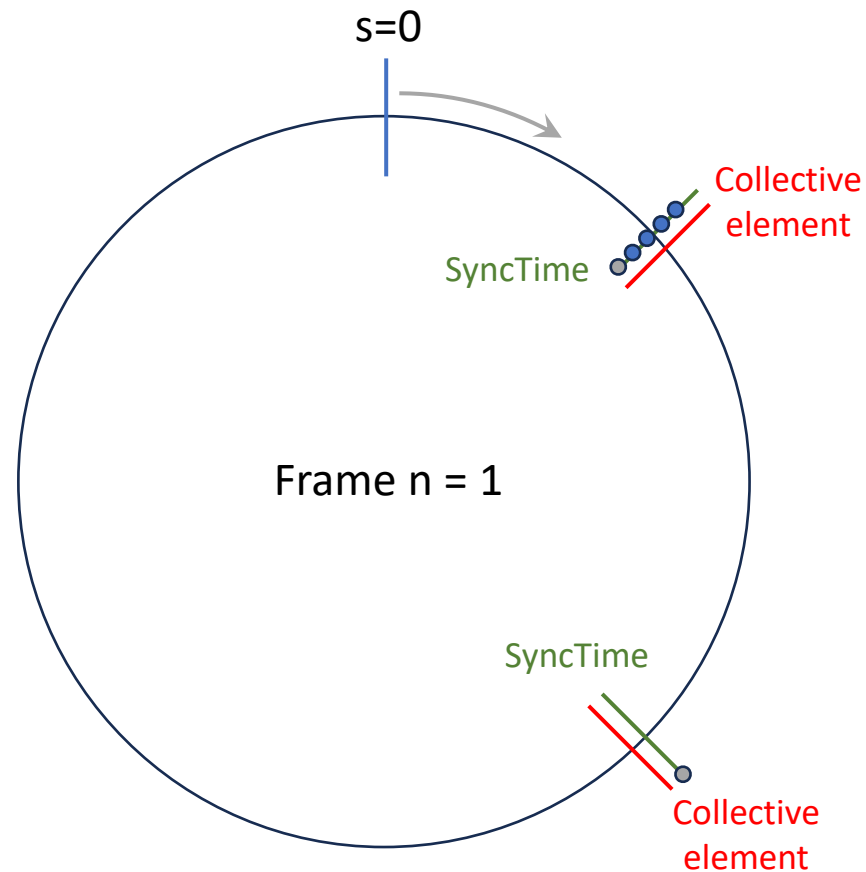
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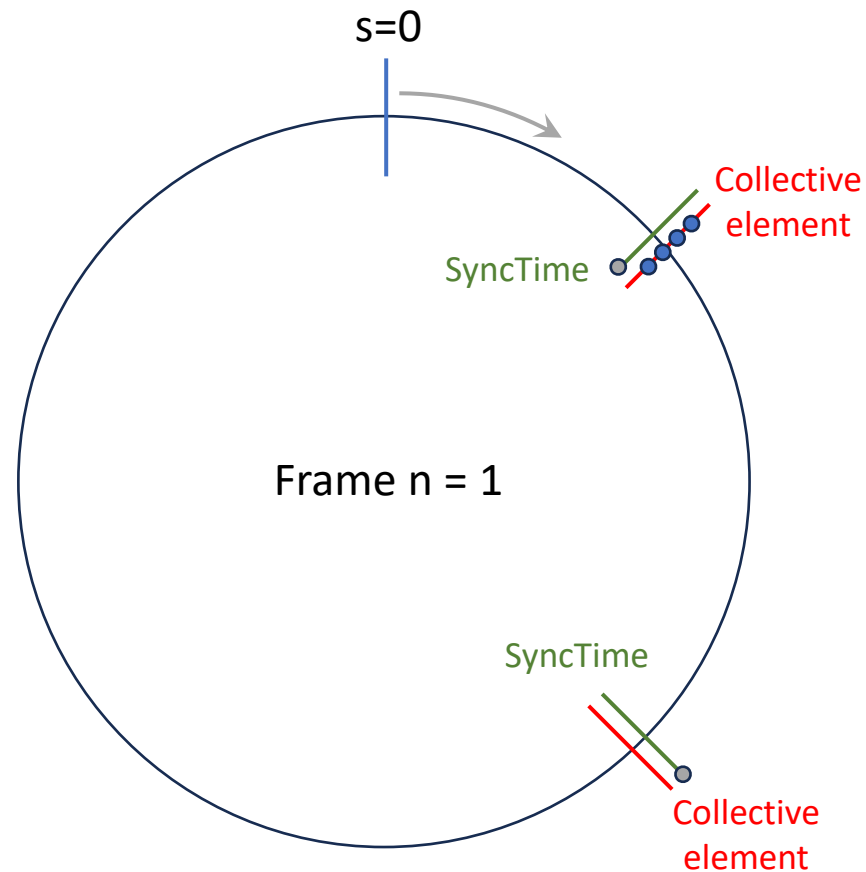
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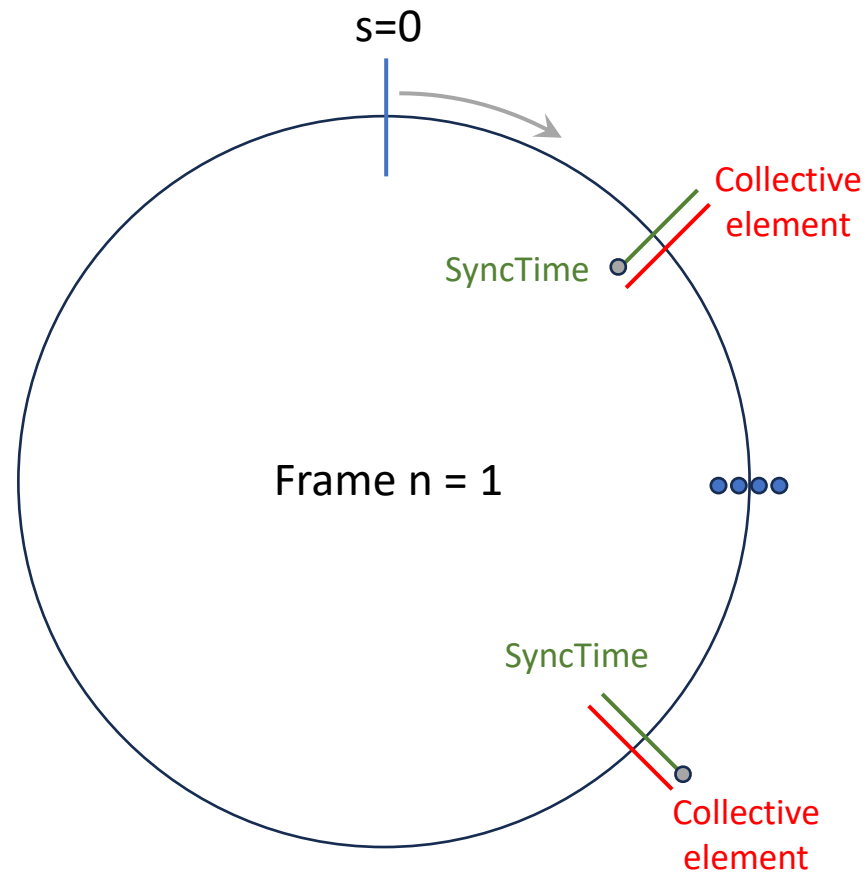
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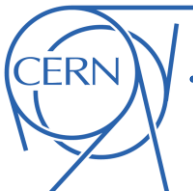
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Propositions on particles arrival times

The procedure illustrated so far can be **rigorously justified mathematically**. Full derivation available in the [Xsuite Physics Guide](#).

A central role is played by the following **two propositions**:

Proposition 1: If the time $t_k(s_1)$ defining the k -th arrival of a particle at location s_1 falls in the frame $F_n(s_1)$, then the particle arrives at location $s_2 > s_1$ either in the frame $F_n(s_2)$ or in the following frame $F_{n+1}(s_2)$. In symbols:

$$t_k(s_1) \in F_n(s_1) \Rightarrow t_k(s_2) \in F_n(s_2) \cup F_{n+1}(s_2) \quad \text{for any } s_1 < s_2 \quad (5.10)$$

Proposition 2: If the time $t_k(s_2)$ defining the k -th arrival of a particle at location s_2 falls in the frame $F_n(s_2)$, then the time of $(k+1)$ -th arrival at any location $s_1 < s_2$ falls in the frame $F_{n+1}(s_1)$ or in the following frame $F_{n+2}(s_1)$. In symbols:

$$t_k(s_2) \in F_n(s_2) \Rightarrow t_{k+1}(s_1) \in F_{n+1}(s_1) \cup F_{n+2}(s_1) \quad \text{for any } s_1 < s_2 \quad (5.11)$$

In the following I give you a glimpse of the proof...



Reminders:

$$\hat{\beta}(s_1, s_2) = \frac{1}{c} \frac{s_2 - s_1}{t(s_2) - t(s_1)} \quad \hat{\beta}(s_1, s_2) < \beta_{\text{sim}} \quad T_n(s) = n\Delta T + \frac{s}{\beta_{\text{sim}}c}$$
$$\hat{\beta}(s_1, s_2) > \frac{\beta_{\text{sim}}}{2}$$

We prove an auxiliary result:

$$\hat{\beta} < \beta_{\text{sim}} \Rightarrow \frac{1}{\hat{\beta}} > \frac{1}{\beta_{\text{sim}}} \Rightarrow \left(\frac{1}{\hat{\beta}} - \frac{1}{\beta_{\text{sim}}} \right) > 0$$

$$\hat{\beta} > \frac{\beta_{\text{sim}}}{2} \Rightarrow \frac{1}{\hat{\beta}} < \frac{2}{\beta_{\text{sim}}} \Rightarrow \left(\frac{1}{\hat{\beta}} - \frac{1}{\beta_{\text{sim}}} \right) < \frac{1}{\beta_{\text{sim}}}$$

Combining the two: $0 < \left(\frac{1}{\hat{\beta}} - \frac{1}{\beta_{\text{sim}}} \right) < \frac{1}{\beta_{\text{sim}}}$

Another relation that we will use:

$$T_n(s_2) - T_n(s_1) = \frac{s_2 - s_1}{\beta_{\text{sim}}c}$$



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We want to prove: $t_k(s_1) \in F_n(s_1) \Rightarrow t_k(s_2) \in F_n(s_2) \cup F_{n+1}(s_2)$ for any $s_1 < s_2$

or equivalently: $T_n(s_2) - \frac{\Delta T}{2} < t_k(s_2) < T_{n+1}(s_2) + \frac{\Delta T}{2}$

$$t_k(s_1) \in F_n(s_1) \Rightarrow t_k(s_1) < T_n(s_1) + \frac{\Delta T}{2}$$

$$t_k(s_2) = t_k(s_1) + \frac{s_2 - s_1}{\hat{\beta} c}$$

$$t_k(s_2) < T_n(s_1) + \frac{\Delta T}{2} + \frac{s_2 - s_1}{\hat{\beta} c}$$

$$t_k(s_2) < T_n(s_2) - \frac{s_2 - s_1}{\beta_{\text{sim}} c} - \frac{\Delta T}{2} + \frac{s_2 - s_1}{\hat{\beta} c} \Rightarrow t_k(s_2) < T_n(s_2) + \frac{\Delta T}{2} + \frac{s_2 - s_1}{c} \left(\frac{1}{\hat{\beta}(s)} - \frac{1}{\beta_{\text{sim}}} \right)$$

$$t_k(s_2) < T_n(s_2) + \frac{\Delta T}{2} + \frac{s_2 - s_1}{\beta_{\text{sim}} c} \Rightarrow t_k(s_2) < T_n(s_2) + \frac{\Delta T}{2} + \Delta T \Rightarrow t_k(s_2) < T_{n+1}(s_2) + \frac{\Delta T}{2}$$

$$\frac{s_2 - s_1}{\beta_{\text{sim}} c} < \frac{L}{\beta_{\text{sim}} c} = \Delta T$$

This proves the upper bound.



Reminders:

$$\hat{\beta}(s_1, s_2) = \frac{1}{c} \frac{s_2 - s_1}{t(s_2) - t(s_1)} \quad 0 < \left(\frac{1}{\hat{\beta}} - \frac{1}{\beta_{\text{sim}}} \right) < \frac{1}{\beta_{\text{sim}}} \quad T_n(s_2) - T_n(s_1) = \frac{s_2 - s_1}{\beta_{\text{sim}} c}$$

We want to prove: $t_k(s_1) \in F_n(s_1) \Rightarrow t_k(s_2) \in F_n(s_2) \cup F_{n+1}(s_2)$ for any $s_1 < s_2$

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$$t_k(s_1) \in F_n(s_1) \Rightarrow t_k(s_1) < T_n(s_1) + \frac{\Delta T}{2}$$

$$t_k(s_2) = t_k(s_1) + \frac{s_2 - s_1}{\hat{\beta} c}$$

$$t_k(s_2) < T_n(s_1) + \frac{\Delta T}{2} + \frac{s_2 - s_1}{\hat{\beta} c}$$

$$t_k(s_2) < T_n(s_2) - \frac{s_2 - s_1}{\beta_{\text{sim}} c} - \frac{\Delta T}{2} + \frac{s_2 - s_1}{\hat{\beta} c} \Rightarrow t_k(s_2) < T_n(s_2) + \frac{\Delta T}{2} + \frac{s_2 - s_1}{c} \left(\frac{1}{\hat{\beta}(s)} - \frac{1}{\beta_{\text{sim}}} \right)$$

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The proof for the lower bound is very similar and can be found on the [Xsuite Physics Manual](#)

This proves the upper bound.



How these proofs justify the simulation procedure

Prop. 1: $t_k(s_1) \in F_n(s_1) \Rightarrow t_k(s_2) \in F_n(s_2) \cup F_{n+1}(s_2)$

Prop. 2: $t_k(s_2) \in F_n(s_2) \Rightarrow t_{k+1}(s_1) \in F_{n+1}(s_1) \cup F_{n+2}(s_1)$

for any $s_1 < s_2$

- We detect that the particle is **not in the current frame $F_n(s)$ at $s=s_a$**

$$t_k(s_a) \notin F_n(s_a)$$

- We know from **Prop. 1** that the particles is **already in the next frame**

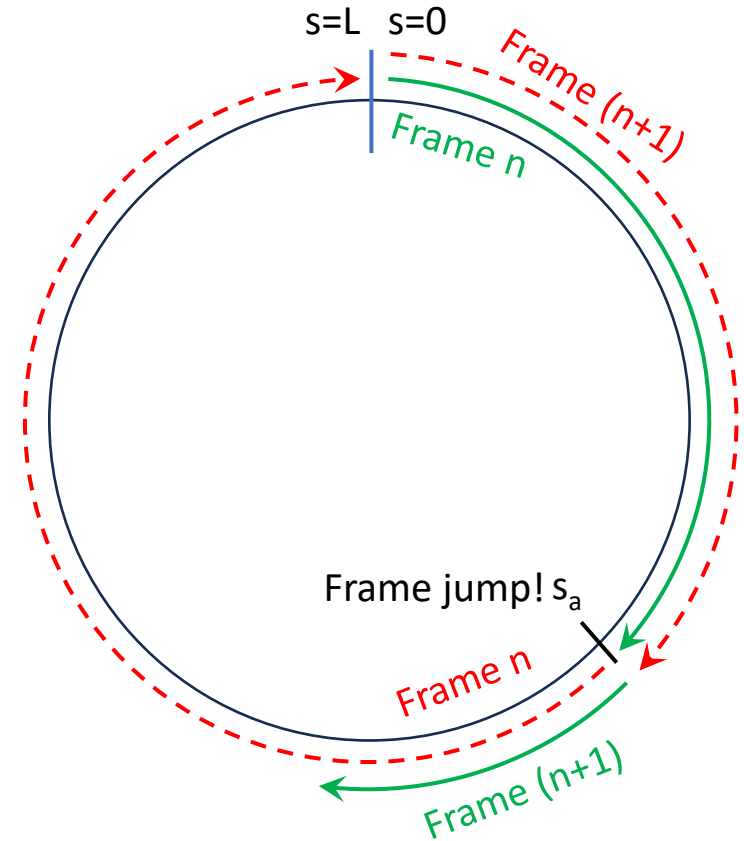
$$t_k(s_a) \in F_{n+1}(s_a)$$

- We know from **Prop. 1** that **for $s > s_a$:**

$$\begin{aligned} t_k(s_a) \in F_{n+1}(s_a) &\Rightarrow t_k(s) \in F_{n+1}(s) \cup F_{n+2}(s) \\ &\Rightarrow t_k(s) \notin F_n(s) \end{aligned}$$

- We know from Prop. 2 for $s < s_a$:

$$\begin{aligned} t_k(s_a) \in F_{n+1}(s_a) &\Rightarrow t_{k+1}(s) \in F_{n+2}(s) \cup F_{n+3}(s) \\ &\Rightarrow t_{k+1}(s) \notin F_{n+1}(s) \end{aligned}$$



This is exactly what we do in the simulator!



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Generalization of the zeta coordinate

Reminder on usual definitions (bunched beams)

In Xsuite (as in all beam dynamics codes) we don't use t to track the particle arrival time. We instead use ζ **defined as**:

$$\zeta = S - \beta_0 c t$$

where **S is the total traveled length** on the reference trajectory:

$$S = s + nL = s - n\beta_0 T_0$$

Combining the relations above we obtain:

$$\zeta = s - \beta_0 c \underbrace{(t - nT_0)}$$

Time within the turn

Generalization to coasting beams:

For coasting beam we define ζ as:

$$\zeta = s - \beta_0 c \underbrace{(t - n\Delta T)}$$

Time within the frame

Reminder

$$\Delta T = \frac{L}{\beta_{\text{sim}} c}$$



Reminder:

$$\zeta = s - \beta_0 c(t - n\Delta T) \quad \Delta T = \frac{L}{\beta_{\text{sim}} c}$$

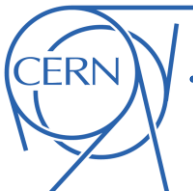
One of the implications of our definition of ζ is that we **need to update ζ at the start of each turn:**

$$\text{End of turn } n \text{ (s=L):} \quad \zeta^- = L - \beta_0 c(t - n\Delta T)$$

$$\text{Start of turn } n+1 \text{ (s=0):} \quad \zeta^+ = 0 - \beta_0 c(t - (n+1)\Delta T)$$

Subtracting one from the other we obtain the **update equation:**

$$\zeta^+ = \zeta^- - L \left(1 - \frac{\beta_0}{\beta_{\text{sim}}} \right)$$



Conditions on zeta for arrival time in the current frame

We can prove⁽¹⁾ the following condition for the **arrival time of a particle to be in the current frame $F_n(s)$** :

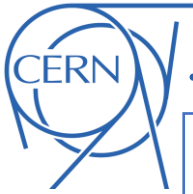
$$t \in F_n(s) \quad \longleftrightarrow \quad -\frac{\Delta\zeta}{2} + \zeta_{\text{corr}}(s) < \zeta < \frac{\Delta\zeta}{2} + \zeta_{\text{corr}}(s)$$

where: $\Delta\zeta = \beta_0 c \Delta T = \frac{\beta_0}{\beta_{\text{sim}}} L$ $\zeta_{\text{corr}}(s) = s \left(1 - \frac{\beta_0}{\beta_{\text{sim}}} \right)$

(1) Proof:

Reminder: $\zeta = s - \beta_0 c(t - n\Delta T)$

$t \in F_n(s)$ $-\frac{\Delta T}{2} < t - T_n(s) < \frac{\Delta T}{2}$ $-\frac{\Delta T}{2} < t - n\Delta T - \frac{s}{\beta_{\text{sim}}c} < \frac{\Delta T}{2}$ $-\frac{\Delta T}{2} < \left(\frac{s}{\beta_0 c} - \frac{\zeta}{\beta_0 c} + n\Delta T \right) - n\Delta T - \frac{s}{\beta_{\text{sim}}c} < \frac{\Delta T}{2}$	$\frac{\Delta T}{2} > -\frac{s}{\beta_0 c} + \frac{\zeta}{\beta_0 c} + \frac{s}{\beta_{\text{sim}}c} > -\frac{\Delta T}{2}$ $-\frac{\Delta T}{2} < -\frac{s}{\beta_0 c} + \frac{\zeta}{\beta_0 c} + \frac{s}{\beta_{\text{sim}}c} < \frac{\Delta T}{2}$ $-\beta_0 c \frac{\Delta T}{2} < \zeta - s \left(1 - \frac{\beta_0}{\beta_{\text{sim}}} \right) < \beta_0 c \frac{\Delta T}{2}$
---	--



$$t \in F_n(s) \iff -\frac{\Delta\zeta}{2} + \zeta_{\text{corr}}(s) < \zeta < \frac{\Delta\zeta}{2} + \zeta_{\text{corr}}(s)$$

Reminder:

$$\text{where: } \Delta\zeta = \beta_0 c \Delta T = \frac{\beta_0}{\beta_{\text{sim}}} L \quad \zeta_{\text{corr}}(s) = s \left(1 - \frac{\beta_0}{\beta_{\text{sim}}} \right)$$

Particles **arrive out of the current frame** (arrive too late) and jump to the next when:

$$t > T_n(s) + \frac{\Delta T}{2} \iff \zeta < -\frac{\Delta\zeta}{2} + s \left(1 - \frac{\beta_0}{\beta_{\text{sim}}} \right)$$

As discussed, the particle is stopped, and its tracking is resumed when handling frame (n+1)

- When this happens the **ζ coordinates needs to be updated:**

$$\zeta_{\text{before jump}} = s - \beta_0 c (t - n\Delta T)$$

$$\zeta_{\text{after jump}} = s - \beta_0 c (t - (n + 1)\Delta T)$$

Subtracting one from the other we obtain the **update equation:**

$$\zeta_{\text{after jump}} = \zeta_{\text{before jump}} + \beta_0 c \Delta T = \zeta_{\text{before jump}} + \Delta\zeta$$

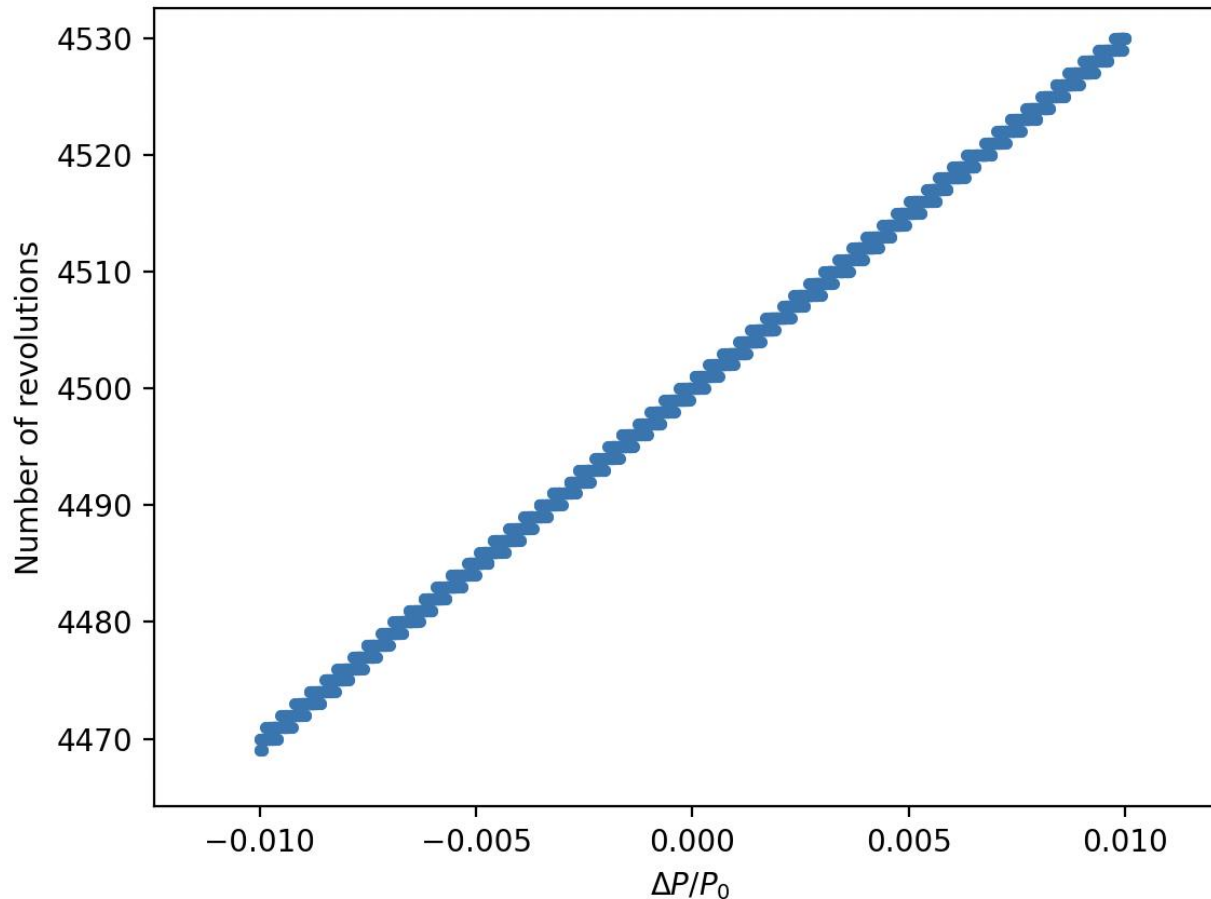


- **Introduction**
 - Particle slippage
 - Implications for collective effects
- **Simulation method description**
 - Reference speed β_{sim}
 - From turns to frames
 - Algorithm step by step
- **Mathematical description**
 - Propositions on time of arrivals
 - Justification of the method
 - Generalization of ζ coordinate
 - ζ update at start turn
 - ζ update on frame jump
- **Numerical tests**



Test case: CERN PS Booster $E_{\text{kin}} = 160 \text{ MeV}$
Full lattice RF off

Tracking a beam with an artificially **large momentum spread** we can clearly see that different particles perform a **different number of revolution over the simulated time**.



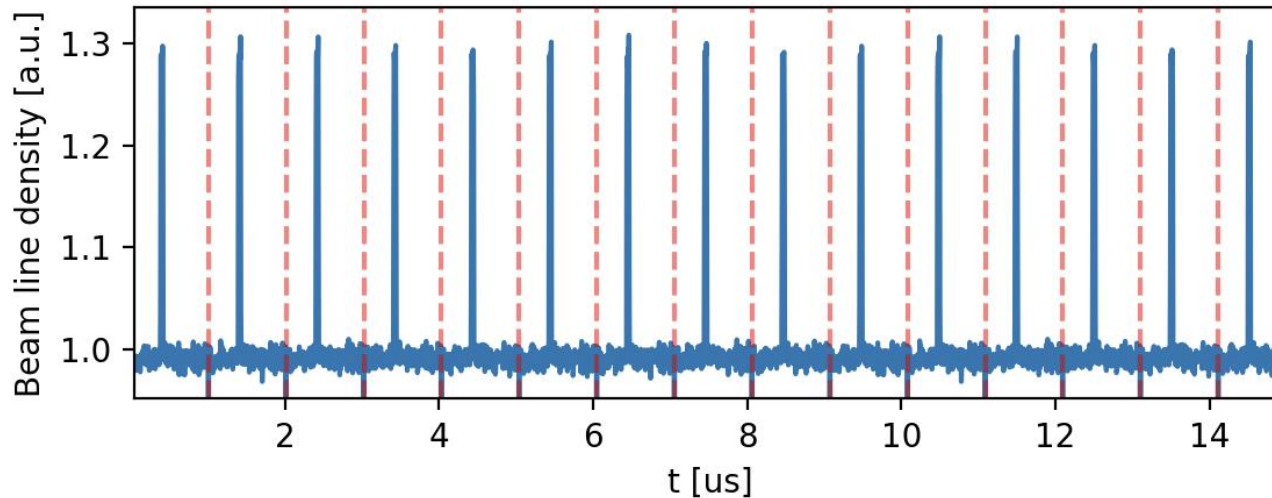


Test simulations – measuring the revolution frequency

We track a **beam with no momentum spread**, and we introduce a **perturbation on the beam line density**.

→ As there is no slippage, the perturbation is observed at each turn

→ It is possible to **measure the revolution frequency** by extracting the **main Fourier component** of the longitudinal profile (using nafflib)



The beam is generated with $\delta = 0$

`f_nominal (on momentum):` 991.96599 kHz

`f_measured (off momentum):` 991.96599 Hz

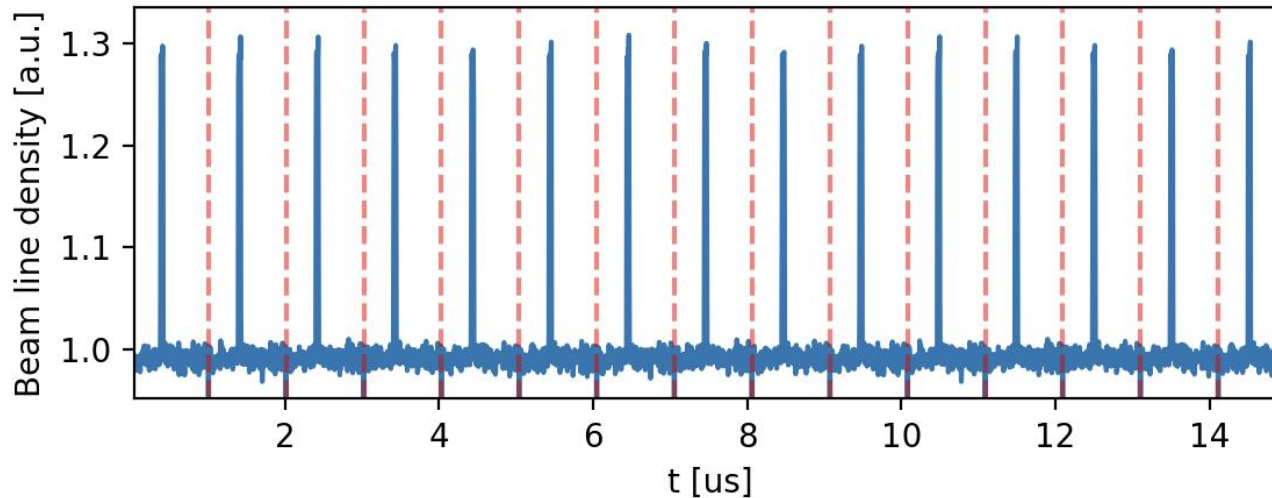


Test simulations – measuring the revolution frequency

We track a **beam with no momentum spread**, and we introduce a **perturbation on the beam line density**.

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→ It is possible to **measure the revolution frequency** by extracting the **main Fourier component** of the longitudinal profile (using nafflib)



The beam is generated with $\delta = -0.01$

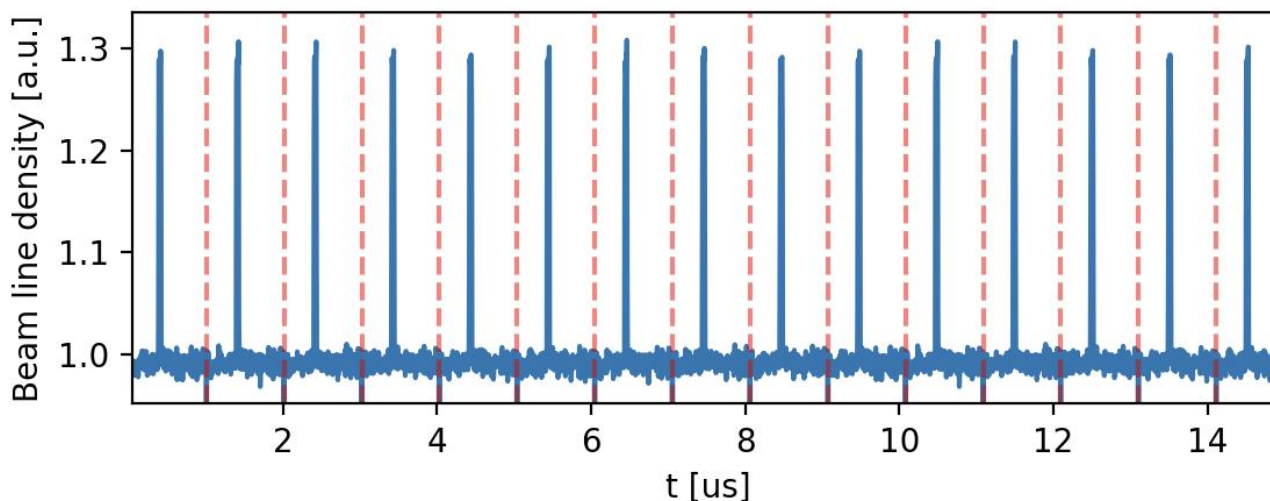
<code>f_nominal</code> (on momentum):	991.966 kHz
<code>f_expected</code> (off momentum):	985.271 kHz
<code>f_measured</code> (off momentum):	985.271 kHz



Test simulations – measuring the revolution frequency

We track a **beam with no momentum spread**, and we introduce a **perturbation on the beam line density**.

- As there is no slippage, the perturbation is observed at each turn
- It is possible to **measure the revolution frequency** by extracting the **main Fourier component** of the longitudinal profile (using nafflib)



The beam is generated with $\delta = +0.01$

<code>f_nominal</code> (on momentum):	991.966 kHz
<code>f_expected</code> (off momentum):	998.580 kHz
<code>f_measured</code> (off momentum):	998.581 kHz

Frequency obtained from the line density agrees very well with expected one 😊

A method has been devised **to simulate coasting beams with tracking codes that use the reference path length s as independent variable**

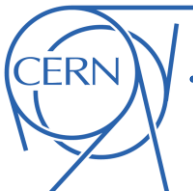
→ The different revolution frequency among particles is accurately modeled

The method has been **implemented in Xsuite:**

- **No modifications in the conventional tracking elements**
- **Time synchronization** of particles is achieved by installing **dedicated SyncTime elements** in front of the collective elements (e.g. space charge, impedances)

Planned applications include:

- PSB with space charge (benchmark of coasting beam experiments)
- Fermilab IOTA ring with space charge
- ISIS-2 stability studies



Thanks for your attention!

```

6 class SyncTime:
7
8     def __init__(self, circumference, id, frame_relative_length=None,
9                 at_start=False, at_end=False):
10         if frame_relative_length is None:
11             frame_relative_length = DEFAULT_FRAME_RELATIVE_LENGTH
12         assert id > COAST_STATE_RANGE_START
13         self.id = id
14         self.frame_relative_length = frame_relative_length
15         self.circumference = circumference
16         self.at_start = at_start
17         self.at_end = at_end
18
19     def track(self, particles):
20
21         assert isinstance(particles._context, xo.ContextCpu), (
22             'SyncTime only available for CPU for now')
23
24         beta0 = particles._xobject.beta0[0]
25         beta1 = beta0 / self.frame_relative_length
26         beta0_beta1 = beta0 / beta1
27
28         mask_alive = particles.state > 0
29
30         zeta_min = -self.circumference / 2 * beta0_beta1 + particles.s * (
31             1 - beta0_beta1)
32
33         if (self.at_start and particles.at_turn[0] == 0
34             and not (particles.state == -COAST_STATE_RANGE_START).any()): # done by the user
35             mask_stop = mask_alive * (particles.zeta < zeta_min)
36             particles.state[mask_stop] = -COAST_STATE_RANGE_START
37             particles.zeta[mask_stop] += self.circumference * beta0 / beta1
38
39         # Resume particles previously stopped
40         particles.state[particles.state == -self.id] = 1
41         particles.reorganize()
42
43         # Identify particles that need to be stopped
44         zeta_min = -self.circumference / 2 * beta0_beta1 + particles.s * (1 - beta0_beta1)
45         mask_stop = mask_alive & (particles.zeta < zeta_min)
46
47         # Check if some particles are too fast
48         mask_too_fast = mask_alive & (
49             particles.zeta > zeta_min + self.circumference * beta0_beta1)
50         if mask_too_fast.any():
51             raise ValueError('Some particles move faster than the time window')
52
53         # Update zeta for particles that are stopped
54         particles.zeta[mask_stop] += beta0_beta1 * self.circumference
55
56         # Stop particles
57         particles.state[mask_stop] = -self.id
58
59         if self.at_end:
60             mask_alive = particles.state > 0
61             particles.zeta[mask_alive] -= (
62                 self.circumference * (1 - beta0_beta1))
63
64         if self.at_end and particles.at_turn[0] == 0:
65             particles.state[particles.state == -COAST_STATE_RANGE_START] = 1
66

```

```

67 def install_sync_time_at_collective_elements(line, frame_relative_length=None):
68
69     circumference = line.get_length()
70
71     ltab = line.get_table()
72     tab_collective = ltab.rows[ltab.iscollective]
73     for ii, nn in enumerate(tab_collective.name):
74         cc = x=SyncTime(circumference=circumference,
75                        frame_relative_length=frame_relative_length,
76                        id=COAST_STATE_RANGE_START + ii + 1)
77         line.insert_element(element=cc, name=f'syncntime_{ii}', at=nn)
78
79     syncntime_start = SyncTime(circumference=circumference,
80                               frame_relative_length=frame_relative_length,
81                               id=COAST_STATE_RANGE_START + len(tab_collective)+1,
82                               at_start=True)
83     syncntime_end = SyncTime(circumference=circumference,
84                              frame_relative_length=frame_relative_length,
85                              id=COAST_STATE_RANGE_START + len(tab_collective)+2,
86                              at_end=True)
87
88     line.insert_element(element=syncntime_start, name='syncntime_start', at_s=0)
89     line.append_element(syncntime_end, name='syncntime_end')
90
91 def prepare_particles_for_sync_time(particles, line):
92     syncntime_start = line['syncntime_start']
93     beta0 = particles._xobject.beta0[0]
94     beta1 = beta0 / syncntime_start.frame_relative_length
95     beta0_beta1 = beta0 / beta1
96     zeta_min = -syncntime_start.circumference / 2 * beta0_beta1 + particles.s * (
97         1 - beta0_beta1)
98     mask_alive = particles.state > 0
99     mask_stop = mask_alive * (particles.zeta < zeta_min)
100     particles.state[mask_stop] = -COAST_STATE_RANGE_START
101     particles.zeta[mask_stop] += syncntime_start.circumference * beta0 / beta1

```