The unexpected uses of a bowling pin Exploiting ²⁰Ne isotopes for precision characterizations of collectivity in small systems

Govert Nijs

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November 11, 2024

Based on:

Giacalone, Bally, GN, Shen, Duguet, Ebran, Elhatisari, Frosini, Lähde, Lee, Lu, Ma, Meißner, Noronha-Hostler, Plumberg, Rodríguez, Roth, van der Schee, Somà, 2402.05995

One fluid to rule them all?

- Anisotropic flow is present in a great range of system sizes:
	- PbP_b.
	- **High multiplicity** pPb **.**
	- **High multiplicity pp.**
	- . . .
- \blacksquare Is this a sign of hydrodynamics?
	- Hydrodynamical simulations seem to work reasonably well.
	- But can a system that small really behave hydrodynamically?
	- Initial state geometry is poorly understood.

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■ We need a precision test of hydrodynamics in small systems.

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Recap: why do we believe PbPb is hydrodynamic?

- Not just the presence of $v_n\{k\}$.
- We understand where the $v_n\{k\}$ come from!
	- Hydrodynamics converts initial state anisotropic geometry into final state momentum anisotropy.
	- We understand very well what the initial geometry looks like!
- For pPb this is not the case.
	- **There is** $v_n\{k\}$ **measured.**
	- But we do not understand the initial geometry.

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■ No clear interpretation of experimental results.

Posing a precise question

- Can we describe PbPb and a small system in a hydrodynamical model with the same settings?
	- Hydro model used should describe a wide range of PbPb observables.
- \blacksquare Can we find a quantity to predict which does not suffer from huge theoretical uncertainties? Wishlist:
	- Initial geometry under control.
	- **Small sensitivity to proton substructure.**
	- No longitudinal structure issues.
	- Quantifiable and small theory uncertainty.

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Can ${}^{16}O{}^{16}O$ collisions help?

- 160^{16} O collisions are planned at the LHC for 2025.
- Shape of the proton and longitudinal structure are not an issue, but. . .

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- 160^{16} O collisions are planned at the LHC for 2025.
- Shape of the proton and longitudinal structure are not an issue, but. . .
- **Magnitude of fluctuations in the** initial state is poorly constrained.
- Different nuclear structure calculations give different answers!
- We have a handle on systematics, but errors are substantial.

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[Giacalone, Bally, GN, Shen et al., 2402.05995] 5/9

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The nuclear bowling pin: ^{20}Ne

- We use both the PGCM and NLEFT frameworks for our nuclear structure input.
	- **PGCM** computes the average deformed densities.
	- **NLEFT** simulates an effective theory on a lattice.
- 16 O is shaped like an irregular tetrahedron.
- \blacksquare ²⁰Ne is close in size, but has the most extreme shape in the Segrè chart.
- **Can we take a ratio between** systems to cancel the uncertainties?

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A careful look at uncertainties

- \blacksquare Trajectum systematic uncertainty contains contributions from:
	- **Uncertainties in parameters.**
	- Extrapolation to zero grid spacing.
- **PGCM** systematic uncertainty contains contributions from:
	- Sampling method: how to convert a density into a configuration.
	- Constraint application: order of operations in the PGCM computation.
- **NLEFT** systematic uncertainty contains contributions from:
	- Resolution of ambiguities from periodicity of the lattice.
	- Nuclear Hamiltonian parameters.

- **NLEFT and PGCM are consistent** within uncertainties.
- Ratio of v_2 {2} reaches percent level precision from 5% to 20% centrality!
- Difference of $\rho({\mathsf{v}_2}\{2\}^2, \langle {\mathsf{p}_\mathcal{T}} \rangle)$ has uncertainty reduced by up to a factor 5!
- **Larger PGCM uncertainty is mostly** due to ambiguity in how to generate configurations from densities.

 $NLEFT$ 1.174 $(8)_{stat.}$ (31) $_{syst.}$ (4) $_{syst.}$

 $\mathsf{PGCM} \parallel 1.139(6)_{\mathsf{stat.}} (27)_{\mathsf{syst.}}^{\mathsf{Traj.}} (28)_{\mathsf{sys.}}^{\mathsf{str.}}$

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$0-1\%$	$v_2\{2\}_{NeNe}/v_2\{2\}_{OO}$	$\rho_{2,NeNe} - \rho_{2,OO}$	
NLEFT	$1.174(8)_{stat.}(31)_{syst.}^{Traj.}(4)_{syst.}^{str.}$	$-0.124(14)_{stat.}(10)_{syst.}^{Traj.}(7)_{syst.}^{str.}$	CFAN
PGCM	$1.139(6)_{stat.}(27)_{syst.}^{Traj.}(28)_{syst.}^{str.}$	$-0.124(10)_{stat.}(10)_{syst.}^{Traj.}(29)_{syst.}^{str.}$	MeV-12.124(19)_{syst.}^{Taj.}(29)_{syst.}

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$$
\begin{array}{c|c|c|c} 0-1\% & v_2\{2\}_{N\in\mathbb{N}}\ell \vee_2\{2\}_{OO} & \rho_{2,N\in\mathbb{N}}\mathrm{e} - \rho_{2,OO}\\ \hline NLEFT & 1.174(8)_{stat.}(31)_{syst.}^{Traj.}(4)_{syst.}^{str.} & -0.124(14)_{stat.}(10)_{syst.}^{Traj.}(7)_{syst.}^{str.}\\ \hline PGCM & 1.139(6)_{stat.}(27)_{syst.}^{Traj.}(28)_{syst.}^{str.} & -0.124(10)_{stat.}(10)_{syst.}^{Traj.}(29)_{syst.}^{str.}\\ \hline & \end{array}
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- NLEFT and PGCM are consistent within uncertainties.
- Ratio of v_2 {2} reaches percent level precision from 5% to 20% centrality!
- Difference of $\rho({\mathsf{v}_2}\{2\}^2, \langle {\mathsf{p}_\mathcal{T}} \rangle)$ has uncertainty reduced by up to a factor 5!
- **Example 2** Larger PGCM uncertainty is mostly due to ambiguity in how to generate configurations from densities.

$$
\begin{array}{c|c|c|c} 0-1\% & v_2\{2\}_{N\in Ne}/v_2\{2\}_{OO} & \rho_{2,N\in Ne} - \rho_{2,OO} \\ \hline NLEFT & 1.174(8)_{stat.}(31)_{syst.}^{Traj.}(4)_{syst.}^{str.} & -0.124(14)_{stat.}(10)_{syst.}^{Traj.}(7)_{syst.}^{str.} \\ \hline PGCM & 1.139(6)_{stat.}(27)_{syst.}^{Traj.}(28)_{syst.}^{str.} & -0.124(10)_{stat.}(10)_{syst.}^{Traj.}(29)_{syst.}^{str.} \\ \hline & \end{array} \hspace{1.5cm} \begin{array}{c|c|c} \hline \end{array}
$$

Revisiting the wishlist

- Theory has a much better handle on $^{16}O^{16}O$ compared to pPb.
- Theory uncertainties can be substantially reduced by supplementing $16O^{16}O$ collisions with ²⁰Ne²⁰Ne collisions.
	- \bullet v_2 {2} ratio can be predicted to 1% precision between 5% and 20% centrality.
	- Different nuclear structure calculations give consistent results.

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