#### **Pre-recombination resolutions to the Hubble tension**

• Angular structure of the CMB must remain  $\simeq$  constant

$$r_s(z_{\rm rec}) = \int_{z_{\rm rec}}^{\infty} \frac{c_s(z)}{H(z)} dz$$

$$\theta_s = \frac{r_s(z_{\text{rec}})}{D_A(z_{\text{rec}})} \sim \frac{c_s(z_{\text{rec}})/H(z_{\text{rec}})}{F(\Omega_m)/H_0} = \frac{H_0}{H(z_{\text{rec}})} \frac{c_s(z_{\text{rec}})}{F(\Omega_m)}$$



• If  $H(z_{rec})$  increases then Silk damping angular scale <u>must</u> increase

See Poulin, TLS, and Karwal 2302.09032

# Stop calling it the Hubble tension!

Aylor++ 1811.00537 Knox and Millea 1908.03663 Bernal, Verde++ 2102.05066 Poulin, TLS++ 2407.18292 Pedrotti++ 2408.04530

In this paper we called it the "cosmic calibration tension"



## Stop calling it the Hubble tension!

*See Poulin, TLS++* 2407.18292



#### Two case studies: axion-like EDE and 'new' EDE

- Axion-like EDE is a cosmological scalar field initially fixed by Hubble friction which then oscillates
- 'New' EDE is a field in a false vacuum which undergoes a phase transition





## **Extensions:**

• Coupling DM & EDE to address  $S_8$   $\mathcal{L} = \frac{1}{2} (\partial \phi)^2 + i \bar{\psi} D \psi - V(\phi) - m_{\text{DM}}(\phi) \bar{\psi} \psi$   $M_{cDonough ++ 1811.04083}$   $\mathcal{L} = \frac{1}{2} (\partial \phi)^2 + i \bar{\psi} D \psi - V(\phi) - m_{\text{DM}}(\phi) \bar{\psi} \psi$  $m_{\text{DM}}(\phi) = m_0 e^{c\phi/M_{\text{pl}}}$ 

Leads to enhanced DM growth:  $G_{eff} = G_N \left( 1 + \frac{2c^2k^2}{k^2 + a^2d^2V/d\phi^2} \right)_{Bean ++ 0808.1105}$ 

• Modified gravity  
Adi and Kovetz 2011.13853  
Abellan, Braglia++ 2308.12345  

$$S = \int d^4x \sqrt{-g} \left[ \frac{F(\sigma)}{2} R - \frac{g^{\mu\nu}}{2} \partial_{\mu}\sigma\partial_{\nu}\sigma - \Lambda - V(\sigma) + \mathcal{L}_m \right]$$

$$V(\sigma) = \lambda \sigma^4/4$$
• Non-minimal coupling to address  $S = \int d^4x \sqrt{-g} \left[ \frac{M_{\rm pl}^2 R(g)}{2} - \frac{1}{2} \nabla_{\mu}\phi \nabla^{\mu}\phi - V(\phi) \right]$ 
fine tuning Sakstein and Trodden 1911.11760  
Gonzalez, Liang, Sakstein and Trodden 2011.09895  $+ S_{\nu}[\tilde{g}_{\mu\nu}], \qquad \ddot{\phi} + 3H\dot{\phi} + V'(\phi) = \frac{\beta}{M_{\rm pl}}\Theta(\nu).$ 
Lin, McDonough, Hill, and Hu 2212.08098  
 $m_{\rm DM}(\phi) = m_0 \left( 1 + g \frac{\phi^2}{M_{\rm pl}^2} \right)$ 

$$V_{\rm eff} \approx V(\phi) + g \frac{\phi^2}{M_{\rm pl}^2}\rho_{\rm DM}$$

### 'Model independent' approaches?

 $\text{\textit{Hojjati et al. 1304.3724}} \quad H^2(a) = \frac{8\pi G}{3} [\rho_m(a) + \rho_r(a) + \rho_\Lambda] \left[1 + \delta(a)\right] \qquad \delta = \sum_i \delta_i \left[\frac{1}{1 + e^{(\ln a - \ln a_{i+1})/\tau}} - \frac{1}{1 + e^{(\ln a - \ln a_i)/\tau}}\right]$ 



## Very Early Dark Energy



## **Projecting into the future**



Tristan L. Smith, ICC, 6 June 2024

## Stop calling it the Hubble tension!

Lynch, Knox, and Chluba, 2406.10202



## The preferred shape of the potential

