# Gravitational Wave Signals in Modified Cosmologies

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# Probing Non-Standard Cosmologies with the Stochastic GW Background

- What contributes to the stochastic GW background?
  - Astrophysical processes (e.g., BH & neutron star mergers)
  - Cosmology: inflation, cosmic strings, etc.
  - GW are induced by curvature perturbations at second order!

K. Ananda, C. Clarkson, D. Wands, Phys.Rev.D 75 (2007) 123518

- Key Questions:
  - What does this signal look like?
  - How does the cosmological history of the universe affect these?

#### **Curvature-Induced Gravitational Waves**

#### Why interesting? Resonance-like enhancement!

First pointed out in K. Ananda, C. Clarkson, D. Wands, Phys Rev D 75 (2007) 123518 Modified curvature power spectrum:  $\mathcal{P}_{\zeta}(k) \propto \delta(k - k_{\mathrm{in}})$ 



Why? Oscillating gravitational potential ↔ oscillating sound wave resonantly produces GW

Inomata et. al. Phys.Rev.D 101 (2020) 12, 123533

#### **Curvature-Induced Gravitational Waves**

- Option 1: Modify  $\mathcal{P}_{\zeta}(k)$  like in Ananda et. al. (Will return to this later)
- Option 2: Use inflationary curvature spectrum-How does cosmological evolution affect it?
  - Consider an early matter dominated epoch



# Induced GW with Early Matter Dominated Epoch

#### Depends sensitively on how the eMD epoch ENDS



Left: Slow end- signal is suppressed because of a cancellation between modes that enter in MD & RD

K. Inomata, K. Kohri, T. Nakama, T. Terada, JCAP 10 (2019) 071

# Induced GW with Early Matter Dominated Epoch

Depends sensitively on how the eMD epoch ENDS



Right: Instantaneous end-Resonance-like enhancement!

K. Inomata, K. Kohri, T. Nakama, T. Terada, Phys.Rev.D 100 (2019) 043532

#### This Resonance...

• My question:

Instantaneous  $\rightarrow$  infinite slope  $\rightarrow$  signal depends on  $\Phi'$ How sure are we that this is real???



Used tanh profile for matter decay rate M. Pearce, L. Pearce, G. White, C. Balazs, JCAP 06 (2024) 021

#### This Resonance...

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Cause: Sudden decay leads to oscillating sound waves in radiation bath Resonantly produce GW *K. Inomata et. al. Phy.Rev.D 101 (2020) 12, 123533* 

#### How Can We Get a "Fast" Phase Transition?

- Need heavy objects whose decays speed up:
  - PBHs: K. Inomata et. al. Phy.Rev.D 101 (2020) 12, 123533
  - Q-balls: G. White., L. Pearce, D. Vagie, A. Kusenko PRL 127 (2021) 18, 181601



Temp scales: eMD starts around  $10^4$  GeV, ends around  $10^2$  GeV.

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  - Q-balls: G. White., L. Pearce, D. Vagie, A. Kusenko PRL 127 (2021) 18, 181601
    - Can be embedded in e.g., high scale SUSY
       M. Flores et. al. Phys.Rev.D 108 (2023) 12, 123002



#### What About Modifying Curvature Power Spectrum?

- Need to modify curvature power spectrum on superhorizon scales
- Phase transitions give  $\mathcal{P}_{\zeta} \propto k^3$ 
  - M. Lewicki, P. Toczek, and V. Vaskonen, arXiv:2402.04158





#### What Else?

- Work has also been done with kination instead of early matter domination
  - G. Domenech, S. Pi, M. Sasaki, JCAP 08 (2020) 017
- Connects to pulsar timing arrays
  - E.g., K. Harigaya, K. Inomata, T. Terada Phys Rev D 108 (2023) 12, 123538



#### Landscape of GW Observations



# Stochastic GW Background (low frequency)

Sources:

- Population of unresolved supermassive black hole binaries
- Early Universe physics (e.g., inflation, phase transitions, topological defects, enhanced scalar perturbations)

Observations:

- Current: evidence from pulsar timing array (PTA) experiments
- Upcoming: LISA will probe higher frequencies



#### PTAs



Figure credit: Andreas Freise

#### **PTA Experiments**



# **Evidence for SGWB**

NANOGrav (2306.16213)



#### **New-Physics Interpretation**



NANOGrav (2306.16219)

#### Astrophysics interpretation





#### Back-Up Slides (Lauren)



# Formalism (1)

Conformal Newtonian Gauge (using conformal time  $\eta$ ):

$$ds^{2} = -a^{2} \left(1 + 2\Phi\right) d\eta^{2} + a^{2} \left(\left(1 - 2\Psi\right)\delta_{ij} + \frac{1}{2}h_{ij}\right) dx^{i} dx^{j}$$

Fourier components of tensor modes satisfy equation of motion:

$$h''_{\vec{k}}(\eta) + 2\mathcal{H}h'_{\vec{k}}(\eta) + k^2h_{\vec{k}}(\eta) = 4S_{\vec{k}}(\eta)$$

with the source expressed in terms of the gravitational potential:

$$S_{\vec{k}} = \int \frac{d^3q}{(2\pi)^{3/2}} e_{ij}(\vec{k}) q_i q_j \left( 2\Phi_{\vec{q}} \Phi_{\vec{k}-\vec{q}} + \frac{4}{3(1+w)} \left( \mathcal{H}^{-1} \Phi_{\vec{q}}' + \Phi_{\vec{q}} \right) \left( \mathcal{H}^{-1} \Phi_{\vec{k}-\vec{q}}' + \Phi_{\vec{k}-\vec{q}} \right) \right)$$

GWs are a second order effect

Following 1804.08577; earlier references mentioned later

# Formalism (2)

Solve via Green's function approach:  $a(\eta)h_{\vec{k}}(\eta) = 4 \int^{\eta} d\bar{\eta} G_{\vec{k}}(\eta, \bar{\eta})a(\bar{\eta})S_{\vec{k}}(\bar{\eta})$ 

$$\mbox{Green's function solves:} \ \ G_{\vec{k}}''(\eta,\bar{\eta}) + \left(k^2 - \frac{a''(\eta)}{a(\eta)}\right) G_{\vec{k}}(\eta,\bar{\eta}) = \delta(\eta - \bar{\eta})$$

Gravitational potential obeys:  $\Phi_{\vec{k}}''(\eta) + \frac{6(1+w)}{(1+3w)\eta} \Phi_{\vec{k}}'(\eta) + wk^2 \Phi_{\vec{k}}(\eta) = 0$ 

Express in terms of transfer function  $\Phi_{\vec{k}} = \Phi(k\eta)\phi_{\vec{k}}$ 

which is normalized to 1 at early times and

$$\left\langle \phi_{\vec{k}}\phi_{\vec{k}'} \right\rangle = \delta(\vec{k} + \vec{k}') \cdot \frac{2\pi^2}{k^3} \left(\frac{3+3w}{5+3w}\right)^2 \mathcal{P}_{\zeta}(k)$$

Following 1804.08577; earlier references mentioned later

# Formalism (3)

Putting it all together, the GW power spectrum is:

$$\mathcal{P}_{h}(\eta, k) = 4 \int_{0}^{\infty} dv \int_{|1-v|}^{1+v} du \left( \frac{4v^{2} - (1+v^{2} - u^{2})^{2}}{4vu} \right)^{2} I^{2}(v, u, x) \mathcal{P}_{\zeta}(kv) \mathcal{P}_{\zeta}(ku)$$
where  $I(v, u, x) = \int_{0}^{x} d\bar{x} \frac{a(\bar{\eta})}{a(\eta)} kG_{k}(\eta, \bar{\eta}) f(v, u, \bar{x})$ 

and the source is

$$\begin{split} f(v,u,\bar{x}) &= \frac{6(w+1)}{3w+5} \Phi(v\bar{x}) \Phi(u\bar{x}) + \frac{6(1+3w)(w+1)}{(3w+5)^2} \left( \bar{x} \partial_{\bar{\eta}} \Phi(v\bar{x}) \Phi(u\bar{x}) + \bar{x} \partial_{\bar{\eta}} \Phi(v\bar{x}) \Phi(v\bar{x}) \right) \\ &+ \frac{3(1+3w)^2(1+w)}{(3w+5)^2} \bar{x}^2 \partial_{\bar{\eta}} \Phi(v\bar{x}) \partial_{\bar{\eta}} \Phi(u\bar{x}) \end{split}$$

Following 1804.08577; earlier references mentioned later