

# Gas-based detectors

Most gas-based detectors rely on the same principles:

- a charged particle ionizes gas atoms/molecules along its track,
- an electric field transports electrons (and ions) towards electrodes,
- electrons are multiplied in a strong electric field,
- the moving electrons/ions induce a current on the readout electrodes.

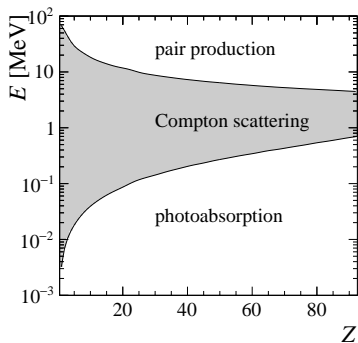
- Most examples on the following slides focus on argon, which is a commonly used detection gas.
- Its ionization potential is 15.7 eV.
- Typically, it is used with an admixture of a “quenching gas” (e. g. CH<sub>4</sub> or CO<sub>2</sub>).



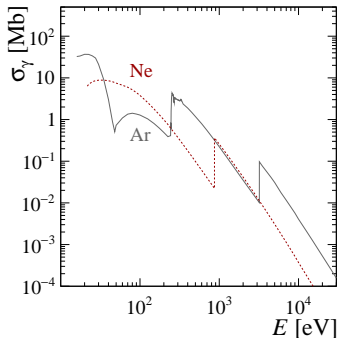
## Why argon?

- Third most common gas in the atmosphere (0.93%).
- Cheap (produced commercially by the distillation of liquid air).
- Inert (noble gas).
- Safe (non-toxic, non-flammable, ...).

## Reminder: photon interactions



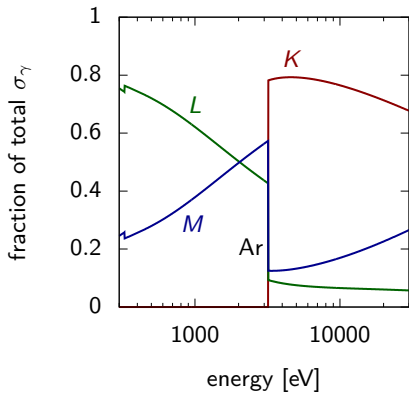
# Photoabsorption



Photoabsorption cross-section  $\sigma_\gamma$  of neon and argon.

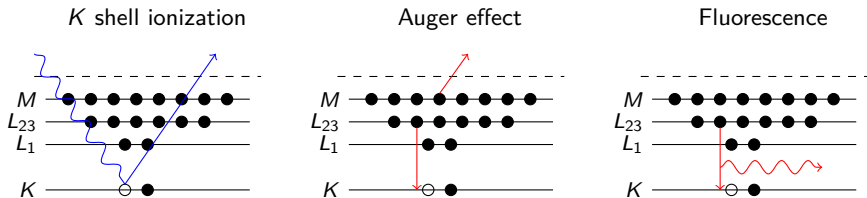
## Example: 5.9 keV $\gamma$ in Ar

- $\sigma_\gamma \sim 0.019$  Mbarn
- Absorption length at room temperature, atm. pressure:  $\lambda \sim 2$  cm

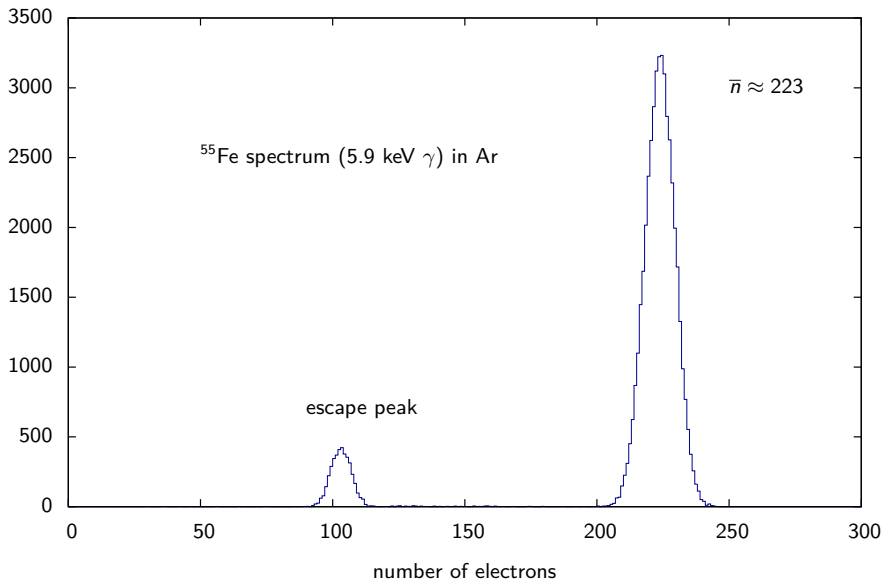


D. A. Verner *et al.*, *At. Data Nucl. Data Tables* **55** (1993), 233-280

- The ionization of inner shells is followed by atomic relaxation processes, resulting in the emission of Auger electrons and fluorescence photons.



- The fluorescence yield (probability to fill a vacancy radiatively) increases with the atomic number  $Z$ .
- Depending on their energy, the emerging photoelectrons and Auger electrons can ionize and excite additional gas atoms along their path and produce secondary electrons.





## W value and Fano factor

- Mean number of electrons:

$$\bar{n} = \frac{E_{\gamma}}{W}$$

- Variance:

$$\sigma^2 = F\bar{n}$$

Gas	W [eV]	F
Ne	35.4	0.13 - 0.17
Ar	26.4	0.16
Xe	22.1	0.17
CO <sub>2</sub>	33.0	0.32
CH <sub>4</sub>	27.3	0.26 - 0.29

ICRU Technical Report 31 (1979)  
IAEA TECDOC 799 (1995)

The interaction of **charged particles** with a gas is described by the differential cross-section  $d\sigma/dE$ , where  $E$  is the energy transfer in a collision.

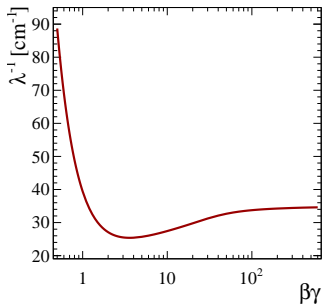
- Average number of collisions per cm (cluster density):

$$\lambda^{-1} = N \int_0^{E_{\max}} dE \frac{d\sigma}{dE}$$

- Average energy loss per cm (stopping power):

$$\frac{dE}{dx} = N \int_0^{E_{\max}} dE \frac{d\sigma}{dE} E$$

## Cluster density



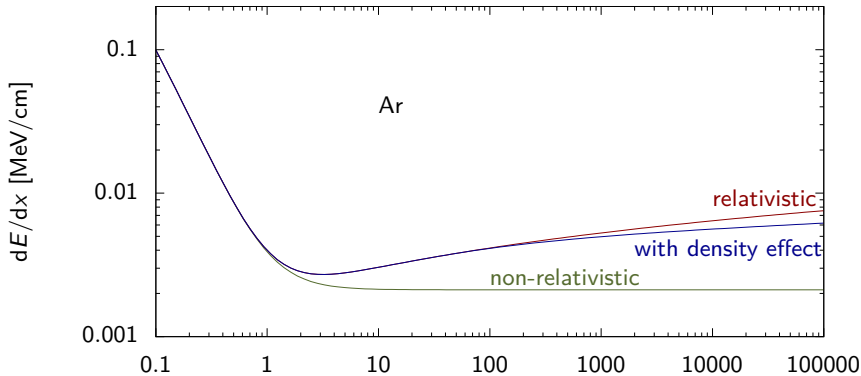
Cluster density in Ar at 20° C, 1 atm.

Measurements with “minimum ionizing” electrons:

Gas	$\lambda^{-1}$ [ $\text{cm}^{-1}$ ]
Ne	11
Ar	23
Xe	43
CO <sub>2</sub>	34
CH <sub>4</sub>	25

F. Rieke and W. Prepejchal,  
Phys. Rev. A **6** (1972), 1507-1519

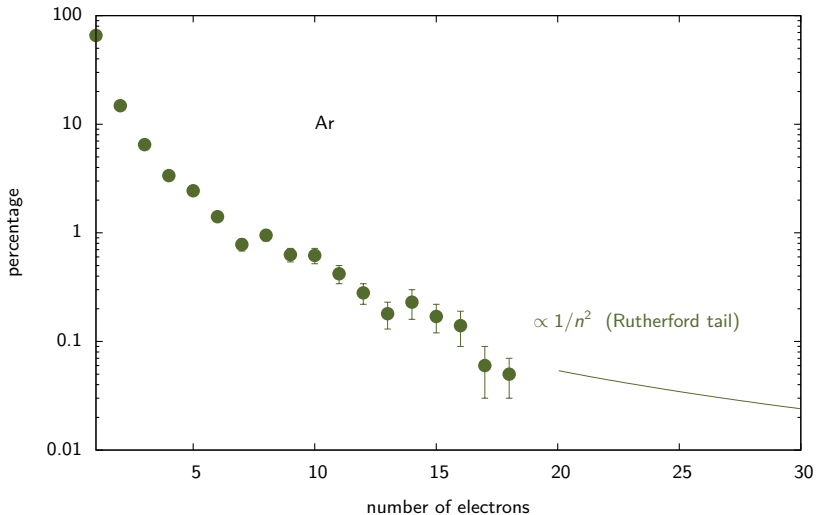
## Stopping power: Bethe-Bloch formula



$$\frac{dE}{dx} = -\frac{4\pi z^2 e^4}{mc^2} NZ \frac{1}{\beta^2} \left[ \ln \frac{2mc^2 \beta^2 \gamma^2}{I} - \beta^2 - \frac{\delta(\beta\gamma)}{2} \right].$$

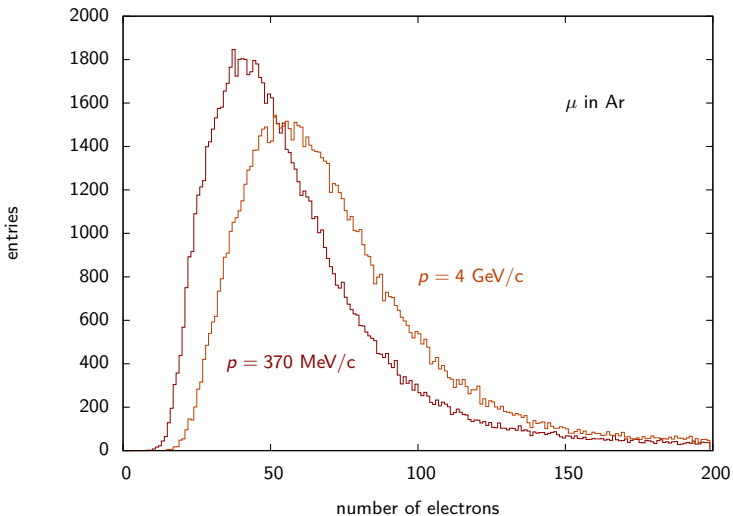
The number of electron/ion pairs deposited in the gas will fluctuate from track to track because of fluctuations

- in the number of collisions a particle suffers over a distance  $x$  (follows a Poisson distribution around  $x/\lambda$ ), and
- in the number of electron/ion pairs produced in a each collision (“cluster size”).



Cluster size distribution in Ar. H. Fischle *et al.*, NIM A **301** (1991), 202-214





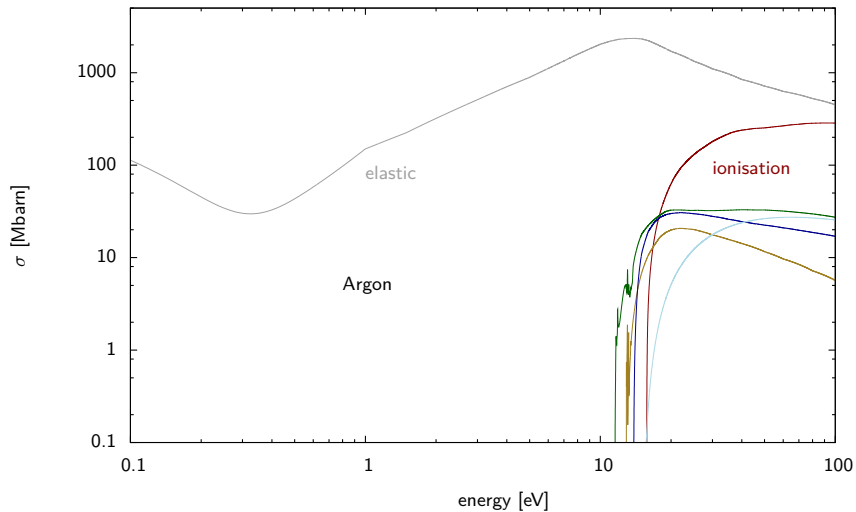
Number of ionization electrons in a 1 cm layer of Ar.



- To detect the electrons and ions produced in the gas by the passage of a charged particle, we need to move them by applying an electric field.

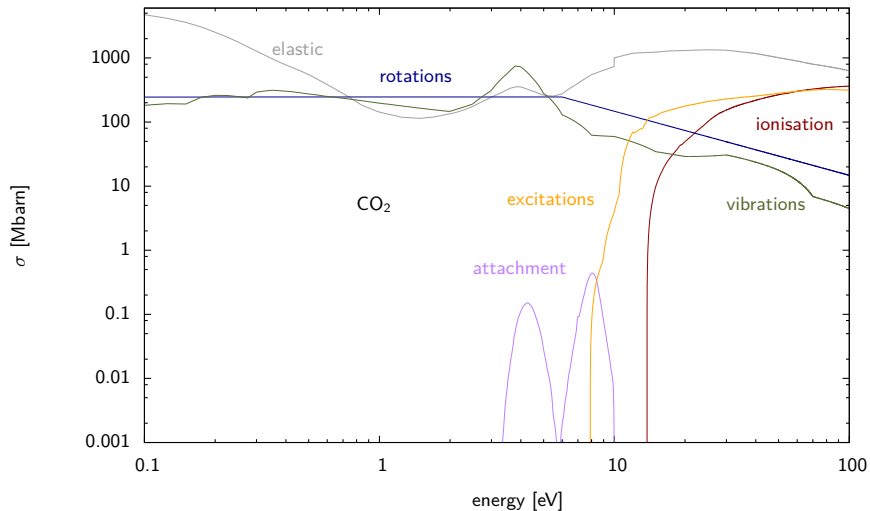
## Electron transport

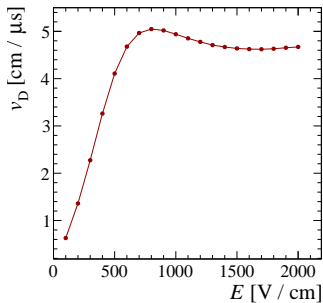
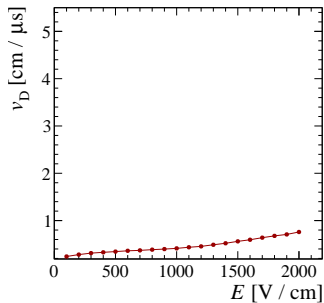
- The microscopic motion of an electron in a gas is a stochastic process.
- At atmospheric pressure, the mean free path between collisions of an electron with a gas atom/molecule is of order  $0.1 \dots 1 \mu\text{m}$ .
- Between collisions, the electron is accelerated by the electric field.
- In a collision, an electron may transfer part of its energy to the atom/molecule and/or change direction.
- To design a detector, we need to know (among other things) the average macroscopic drift velocity of an electron at a given electric field.
- We are also interested in how much individual electrons deviate from the average (diffusion).



- Pure argon,  $E = 1 \text{ kV} / \text{cm}$

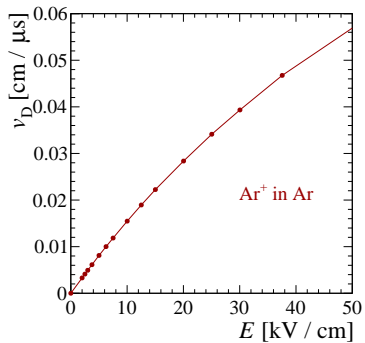
- Ar/CO<sub>2</sub> (90:10),  $E = 1 \text{ kV} / \text{cm}$





Electron drift velocity in pure argon (left) and Ar/CO<sub>2</sub> 90:10 (right).

## Ion transport



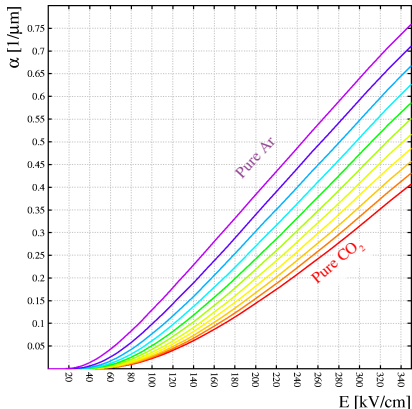
## Gain

- The amount of charge produced by a charged particle crossing a gas detector is usually too small to be measured directly ( $40 e^- \sim 6.4 \times 10^{-3} \text{ fC}$ ).
- We need to amplify the signal by creating secondary avalanche electrons in a strong electric field (few tens of  $\text{kV / cm}$ ).
- The avalanche growth is characterised by the Townsend coefficient  $\alpha$ ,

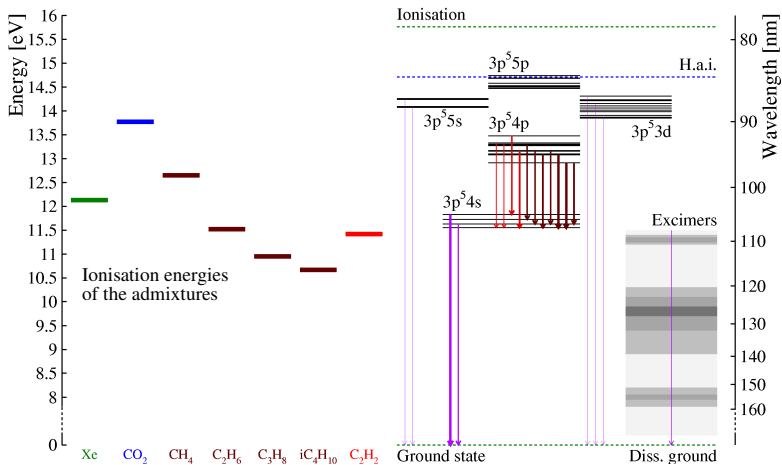
$$dn = n\alpha dx$$

- In a constant electric field, the gas gain is given by  $\exp(\alpha x)$ .



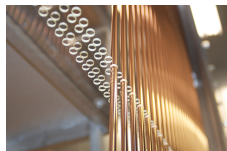
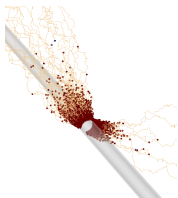
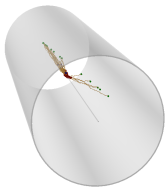


- The above plot (stolen from Rob Veenhof) shows the Townsend coefficient for Ar/CO<sub>2</sub> mixtures, calculated using the Magboltz program.
- Does the calculated gain of the gas mixture match measurements?



## Penning transfer

- Argon has excited states above the ionization energy of typical admixtures, including CO<sub>2</sub> (13.78 eV ionization energy).
- With a certain probability, the energy of such excited states can be transformed into ionization of the admixture.
- This process (Penning effect), can contribute significantly to the effective gain of gas mixtures.

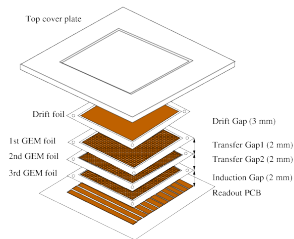
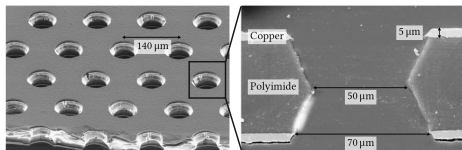


## How to create a strong electric field?

- One technique (going back to the Geiger counter) is to use a wire ( $1/r$  dependence of the electric field).
- Wire-based chambers (MWPC, drift chambers, straw tubes, ...) are still widely used today, but they have limitations in terms of particle rate and track density.
- This motivated the development of detectors with finer granularity (cell size of order  $100\ \mu\text{m}$ ), profiting from advances in photolithography: Micropattern Gas Detectors (MPGD).

# Gaseous Electron Multiplier (GEM)

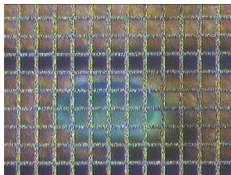
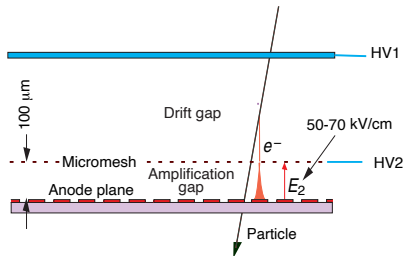
- Thin metal-clad polymer foil, with a regular, high-density pattern of chemically etched holes.
- Applying a potential difference between the metal layers creates a high-field amplification region in the holes.
- To achieve gains of order  $10^4 - 10^5$ , one often puts multiple (typically three) GEMs in series.
- GEMs are used e. g. in the CMS muon system and in the ALICE TPC.



Movie

# Micromegas

- Consists of an amplification gap (few tens of  $\text{kV/cm}$ ,  $50 - 150 \mu\text{m}$ ) and a drift gap ( $\sim \text{kV/cm}$ ).
- The two regions are separated by a micromesh ( $\sim 15 \mu\text{m}$  wires) supported by pillars.
- Ionization electrons produced in the drift gap are funneled into the amplification gap by the high ratio of electron fields between the two gaps.
- Micromegas are used, for instance, in the ATLAS muon system.



## Induced signals

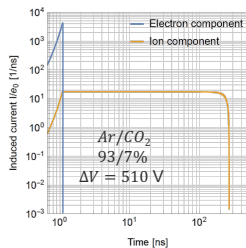
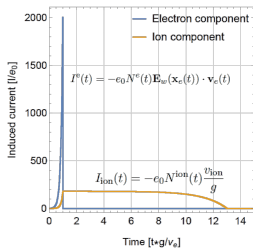
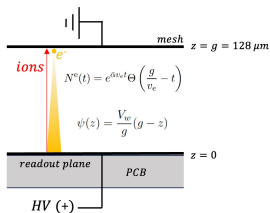
- Recap: motion of charges (electrons/ions) in the gas induces a current on the readout electrodes.
- To calculate the induced current, we use the Ramo-Shockley theorem

$$i(t) = q\mathbf{v} \cdot \mathbf{E}_w$$

where  $E_w$  is the so-called weighting field.

- To calculate the weighting field, we apply 1 V to the readout electrode and ground all other electrodes.
- In a parallel-plate chamber (gap  $d$ ), the weighting field is constant ( $E_w = 1/d$ ).

## Signal in a Micromegas



Slide stolen from Djunes Janssens.