

# A new high-precision measurement of the $B_d^0$ meson lifetime at the ATLAS experiment

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- Introduction, Physics motivation
- ATLAS detector and performance in Run2
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  - Data and event selection
- Maximum-likelihood fit
- Efficiencies and corrections
- Systematic uncertainties
- Results
- Comparison to others
- Summary and Conclusions

- Precise measurements on B-lifetimes and their Ratios test our understanding of weak interactions.
- In Heavy Quark Expansion (HQE) theory the total decay rate  $\Gamma = 1/\tau$  of a weakly decaying Heavy Hadron  $B_q$  is calculated by formula consisting of two terms:

$$\Gamma(\mathcal{B}_q) = \Gamma_3 + \delta\Gamma(\mathcal{B}_q)$$

leading

subleading

## Free $b$ -quark decay:

- + free of non-perturbative uncertainties
- 0 Looks like the muon decay

$$\Gamma_3 \propto \frac{G_F^2 m_b^5}{192\pi^3} V_{cb}^2$$

- Quark masses are difficult to define, huge dependence on definition can be reduced by higher order **perturbative corrections**

## Power-suppressed terms on the HQE:

- + suppressed with at least 2 powers of  $1/m_b \Rightarrow$  small
- 0 Individual contributions are products of **perturbative** Wilson coefficients and **non-perturbative matrix elements** (determined with lattice-QCD, sum rules and/or from fits of experimental data of inclusive semi-leptonic decays -  $V_{cb}$ )

- Overview of HQE calculations: historic development [1405.3601](#), and status quo [2402.04224](#), Feb 2024.

- The HQE prediction of decay rates suffers from huge theory uncertainties due to  $m_b^5$  in  $\Gamma_3$  term  
 $\Gamma_d = 0.63_{-0.07}^{+0.11} \text{ ps}^{-1}$  [Lenz, Piscopo, Rusov 2023](#) 
- In the lifetimes ratios - the free quark decay rate  $\Gamma_3$  exactly cancels leading to smaller uncertainties:  
 $\Gamma_d/\Gamma_s = 1.003 \pm 0.006$  [Lenz, Piscopo, Rusov 2023](#) 

- Lifetimes measurements can also serve to test Theory models of New Physics (BSM):

$$\Gamma(\mathcal{B}_q) = \Gamma_3^{\text{SM}} + \Gamma_3^{\text{BSM}} + \delta\Gamma(\mathcal{B}_q)^{\text{SM}} + \delta\Gamma(\mathcal{B}_q)^{\text{BSM}}$$

- From theory considerations it would be very unlikely to have e.g. a 5% BSM contribution to the total decay rate, but it is also hard to exclude such a possibility at present.
- On the other hand excluding BSM effects to be larger than 5%, will already constrain many BSM scenarios Lenz2021 
- In all scenarios: SM or BSM, the improvement in experimental measurements of lifetimes and their ratios, will serve to constrain theory models.

The lifetime we measured in  $B^0 \rightarrow J/\psi K^{*0}$  is the effective lifetime  $\tau_{B_d^0}$  related to decay widths:  $\Gamma_L, \Gamma_H$  of the light and heavy mass eigenstates of  $B_d^0-\bar{B}_d^0$  system via: [Fleischer et al, 2011](#)

$$\tau_{B^0} = \frac{1}{\Gamma_d} \frac{1}{1-y^2} \left( \frac{1+2Ay+y^2}{1+Ay} \right), \quad (1)$$

$\Gamma_d = (\Gamma_L + \Gamma_H)/2$ ;  $y = \Delta\Gamma_d/(2\Gamma_d) = (\Gamma_L - \Gamma_H)/(2\Gamma_d)$ . The asymmetry  $A$ :

$$A = \frac{R_H^f - R_L^f}{R_H^f + R_L^f}.$$

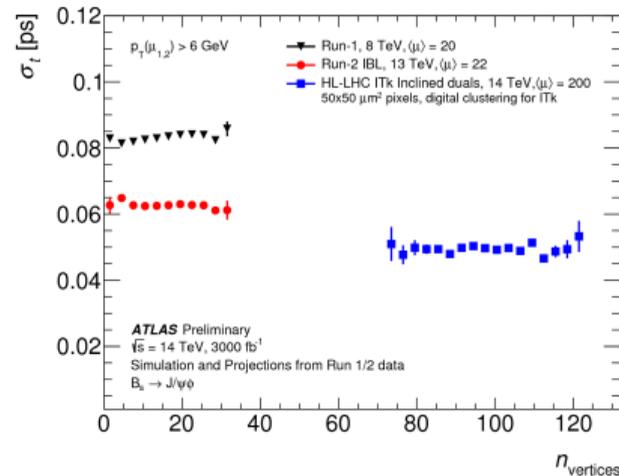
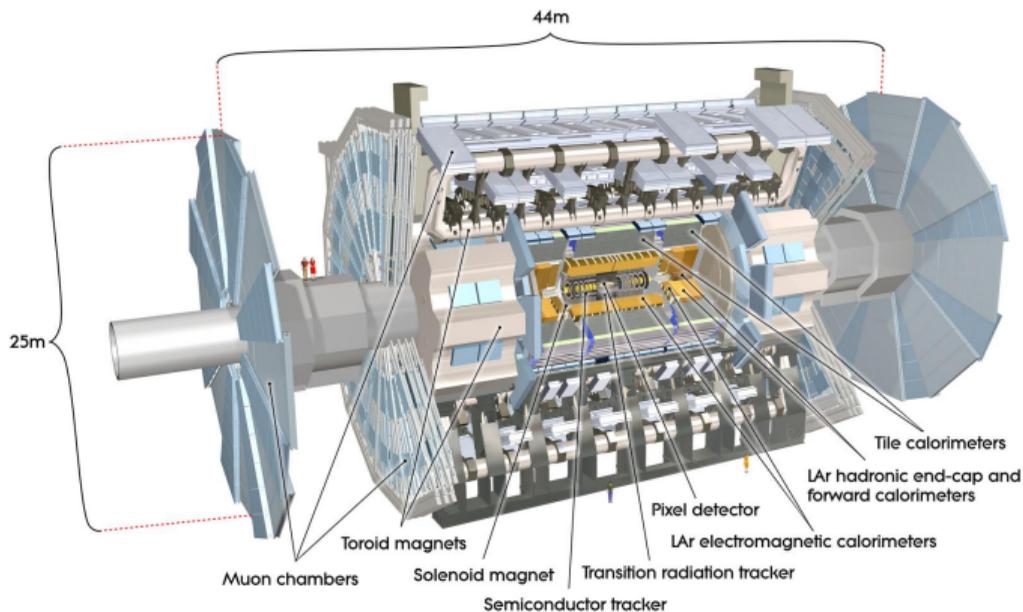
amplitudes  $R_L^f$  and  $R_H^f$  are defined via the summed decay rate of the members of the  $B_d^0-\bar{B}_d^0$  system: [Fleischer et al, 2011](#)

$$\langle \Gamma(B^0(t)) \rangle = \Gamma(B^0(t)) + \Gamma(\bar{B}^0(t)) = R_H^f \exp(-\Gamma_H t) + R_L^f \exp(-\Gamma_L t).$$

Using the values of  $y$  and  $A$  from Heavy Flavour Averaging group (HFLAV) Ref. [HFLAV:2023](#) , our  $\tau_{B_d^0}$  value and Eg(1) allows  $\Gamma_d$  to be extracted.

ATLAS measured  $\Gamma_s = 0.6703 \pm 0.0014(\text{stat.}) \pm 0.0018(\text{syst.}) \text{ ps}^{-1}$  from  $B_s^0 \rightarrow J/\psi \phi$  [Eur. Phys. J. C 81 \(2021\) 342](#) . This result combined with  $B_d^0 \rightarrow J/\psi K^*$  allowed us to determine the ratio  $\Gamma_d/\Gamma_s$ .

# ATLAS detector and feature important for this measurement

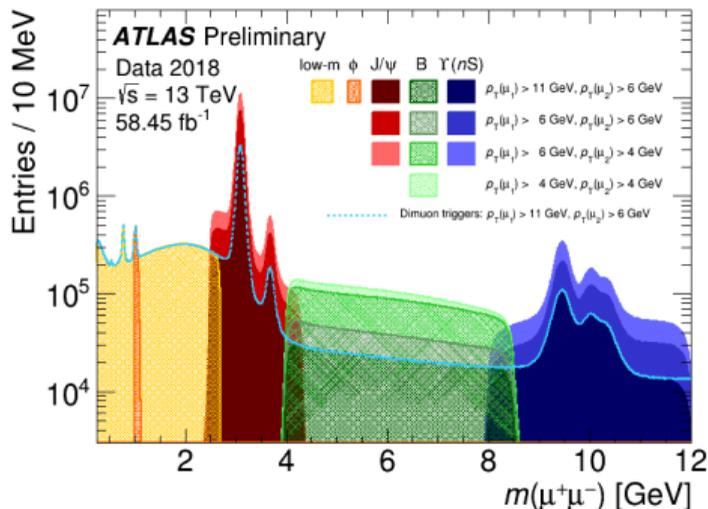


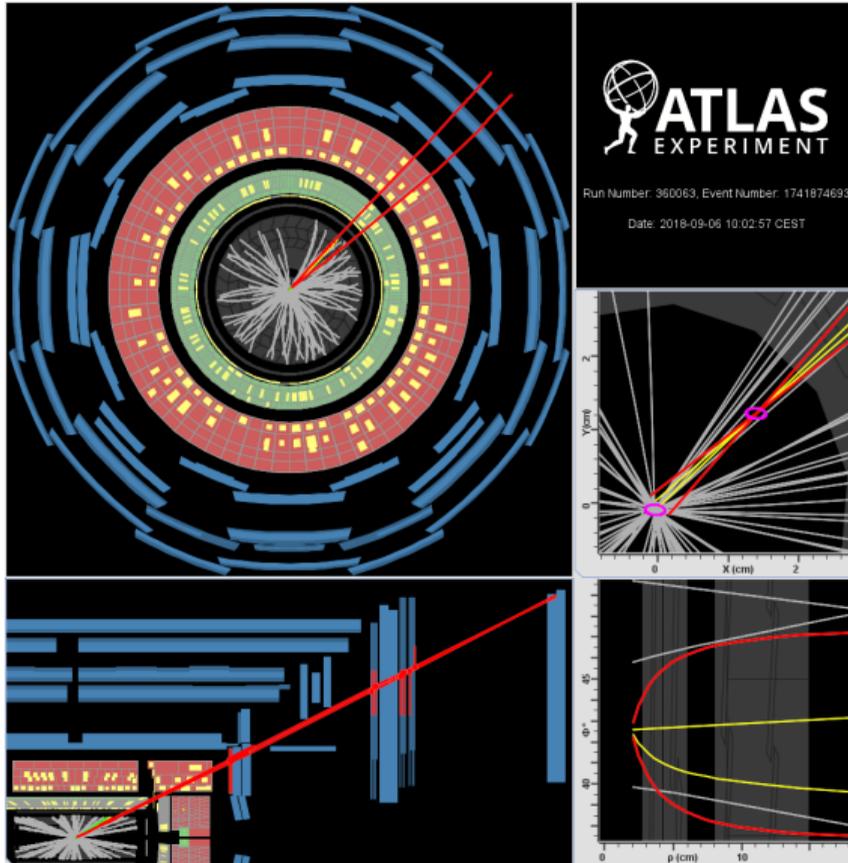
Time resolution of  $B_s^0 \rightarrow J/\psi\phi$  for different numbers of reconstructed PV in the same bunch crossing [ATL-PHYS-PUB-2013-010](#)

- Inner Detector: PIX, SCT and TRT,  $p_T > 0.5 \text{ GeV}$ ,  $|\eta| < 2.5$ 
  - Run2: new IBL 25% improvement of time resolution with respect to Run1.
  - Time, mass resolutions remain stable within increasing pileup in Run 2
- Muon Spectrometer: triggering ( $|\eta| < 2.4$ ), precision tracking ( $|\eta| < 2.7$ )

# Data-taking conditions for this analysis

- The measurement uses 2015-2018  $pp$  collision data at  $\sqrt{s} = 13$  TeV.
- Triggers:  $J/\psi \rightarrow \mu^+ \mu^-$ , the muon  $p_T$  thresholds varying:
  - 4 GeV, 6 GeV and 11 GeV.
  - Low  $p_T$  thresholds were activated in the end of fills when the instantaneous luminosity decreases.
- For events accepted for this analysis, an average number of  $pp$  interactions per bunch crossing (pile-up) was 31.
- No displaced  $J/\psi$  vertex cuts applied. Trigger tracking: transverse parameter  $d_0 < 10$ mm on all tracks.





## $B^0 \rightarrow J/\psi K^{*0}$ candidates

- At least one  $J/\psi \rightarrow \mu^+ \mu^-$  with  $\chi^2/\text{ndof} < 10$ ; within mass window retaining 99.7%  $J/\psi$  candidates.
- $K^*$ : Out of two hypothesis  $K^+ \pi^- / K^- \pi^+$  the one closer to  $K^*$  PDG mass selected.
- $J/\psi$  and  $K^*$  - fitted to a common vertex, constrained by fixing di-muon mass to  $J/\psi$  PDG. Only  $\chi^2/\text{ndof} < 3$  retained.
- 10% events have multiple  $J/\psi K^*$  candidates (in average 2.1); the one with smallest  $\chi^2/\text{ndof}$  selected.

## Primary vertex (PV) selection

- For selected events the average pileup is 31: need to choose the best PV candidate where  $B_d^0$  is produced.
- PV positions are recalculated after removing any tracks used in the  $B_d^0$ .
- The PV candidate with the smallest 3D impact parameter,  $a_0$  (min. distance between PV and the line extrapolated from the  $B_d^0$  vertex in  $B_d^0$  momentum direction), is used.

For each  $B_d^0$  candidate, the proper decay time  $t$  is determined:

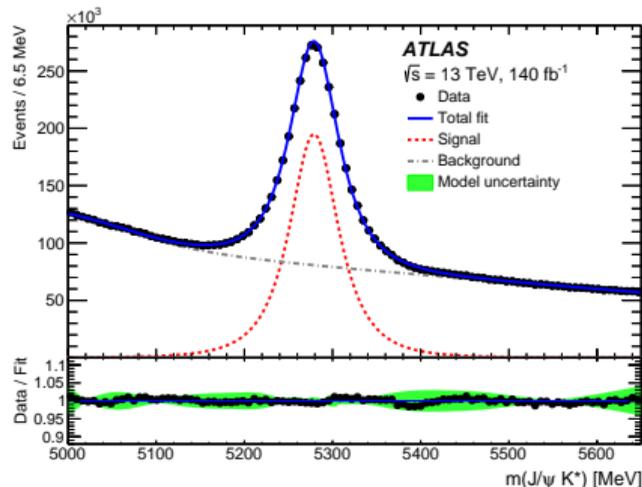
$$t = \frac{L_{xy} m_B}{p_{\Gamma B}}, \quad (2)$$

$p_{\Gamma B}$  is transverse momentum of  $m_B = \text{PDG } B_d^0$  mass. The transverse decay length,  $L_{xy}$ , is the distance in the transverse plane from PV to the  $B_d^0$  decay vertex, projected onto  $p_{\Gamma B}$ .

2D unbinned maximum-likelihood fit applied simultaneously to mass and proper decay time of  $B_d^0$  candidates. Likelihood Function formula:

$$\ln L = \sum_{i=1}^N w(t_i) \ln [f_{\text{sig}} \mathcal{M}_{\text{sig}}(m_i) \mathcal{T}_{\text{sig}}(t_i, \sigma_{t_i}, p_{T_i}) + (1 - f_{\text{sig}}) \mathcal{M}_{\text{bkg}}(m_i) \mathcal{T}_{\text{bkg}}(t_i, \sigma_{t_i}, p_{T_i})].$$

- $f_{\text{sig}}$  - fraction of signal events in the total number of events,  $N$ .
- $\mathcal{M}_{\text{sig}}, \mathcal{T}_{\text{sig}}$  mass and time signal PDFs
- $\mathcal{M}_{\text{bkg}}, \mathcal{T}_{\text{bkg}}$  mass and time Background PDFs
- The mass  $m_i$ , the proper decay time  $t_i$ , its uncertainty  $\sigma_i$  and the  $B_d^0$  candidate transverse momentum  $p_{T_i}$  are the values measured from the data for each event  $i$ .
- weight  $w_i$  accounts for event selection efficiency, that will be explained later.



Signal mass is modelled with a Johnson  $S_U$ -distribution [Johnson](#)

$$\mathcal{M}_{\text{sig}}(m_i) = \frac{\delta}{\lambda \sqrt{2\pi} \sqrt{1 + \left(\frac{m_i - \mu}{\lambda}\right)^2}} \exp \left[ -\frac{1}{2} \left( \gamma + \delta \sinh^{-1} \left( \frac{m_i - \mu}{\lambda} \right) \right)^2 \right]$$

where  $\mu$ ,  $\gamma$ ,  $\delta$  and  $\lambda$  are free parameters of fit.

**Background** has two components: 1. **prompt** -  $J/\psi$  produced in  $pp \rightarrow J/\psi X$ , combined with random  $K^*$ . 2. **combinatorial** -  $J/\psi$  from any  $b$ -hadron decay combined with random  $K^*$ . They are modelled by sum of linear and sigmoid functions:

$$\mathcal{M}_{\text{bkg}}(m_i) = f_{\text{poly}}(1 + p_0 \cdot m_i) + (1 - f_{\text{poly}}) \left( 1 - \frac{s(m_i - m_0)}{\sqrt{1 + (s(m_i - m_0))^2}} \right)$$

where  $f_{\text{poly}}$  relative size of the two components and  $m_0$ ,  $s$  and  $p_0$  are parameters of the fit.

**Signal PDF:** exponential function convolved by Resolution function  $R$ .

$$P_{\text{sig}}(t_i | \sigma_{t_i}, p_{T_i}) = E(t', \tau_{B^0}) \otimes R(t' - t_i, \sigma_{t_i})$$

$E(t', \tau_{B^0}) = (1/\tau_{B^0}) \exp(-t'/\tau_{B^0})$  for  $t' \geq 0$ , with  $\tau_{B^0}$  - the fitted  $B^0$  lifetime.

**Background PDF:**

$$P_{\text{bkg}}(t_i | \sigma_{t_i}, p_{T_i}) = \left( f_{\text{prompt}} \cdot \delta_{\text{Dirac}}(t') + (1 - f_{\text{prompt}}) \sum_{k=1}^3 b_k \prod_{l=1}^{k-1} (1 - b_l) E(t', \tau_{\text{bkg}_k}) \right) \otimes R(t' - t_i, \sigma_{t_i})$$

Dirac function  $\delta_{\text{Dirac}}$ : direct background; 3 exponentials  $E(t', \tau_{\text{bkg}_k})$ : components of combinatorial background;  $\tau_{\text{bkg}_k}$ , and fractions:  $f_{\text{prompt}}$  and  $b_k$  are free parameters of fit.

**Resolution function  $R$ :** modelled as a sum of three Gaussian distributions with widths:  $S^{(k)} \sigma_{\tau_i}$

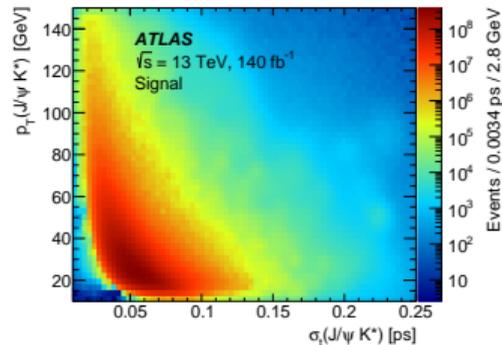
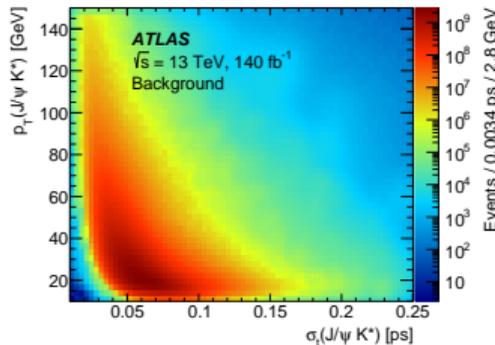
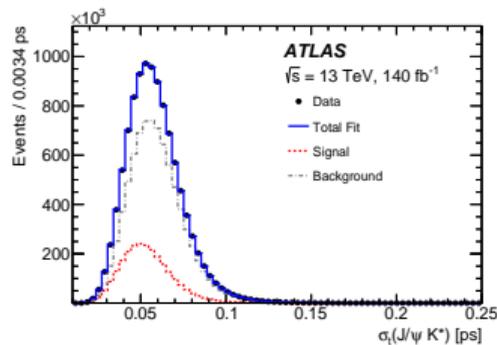
$$R(t' - t_i, \sigma_{t_i}) = \sum_{k=1}^3 f_{\text{res}}^{(k)} \frac{1}{\sqrt{2\pi} S^{(k)} \sigma_{t_i}} \exp\left(\frac{-(t' - t_i)^2}{2(S^{(k)} \sigma_{t_i})^2}\right)$$

three  $f_{\text{res}}^{(k)}$  fractions and scale factors  $S^{(k)}$  are free parameters of fit,  $\sigma_{\tau_i}$  is the per-candidate time error, extracted from data in the vertex fit of each  $B^0 \rightarrow J/\psi K^{*0}$ .

# Conditional probability of the time uncertainty $\sigma_{\tau_i}$ and $p_{\Gamma_i}$ of $B_d^0$

- Per-candidate time errors  $\sigma_{\tau_i}$ , extracted from data in the vertex fit of each  $B_d^0$  are used in the Resolution function  $R$  for deconvolution of proper decay times.
- $\sigma_{\tau_i}$  are different for signal and background, see fig Left. Same it true for  $p_{\Gamma}$  of  $B_d^0$ .
- To account for these differences in the likelihood fit, the 2D probability terms  $C_j(\sigma_{\tau_i}, p_{\Gamma_i})$  are introduced into Time PDF terms  $\mathcal{T}_j$  for Signal and Bg, j - stands for Signal or Bg. The method first used in [G.Punzi](#)

$$\mathcal{T}_j(\tau_i, \sigma_{\tau_i}, p_{\Gamma_i}) = P_j(\tau_i, \sigma_{\tau_i}, p_{\Gamma_i}) \cdot C_j(\sigma_{\tau_i}, p_{\Gamma_i}), \quad (3)$$



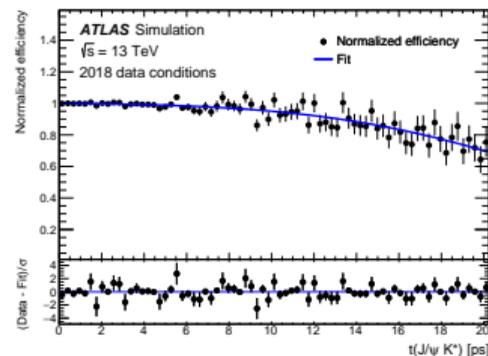
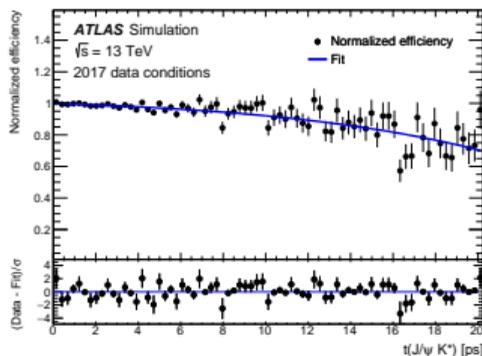
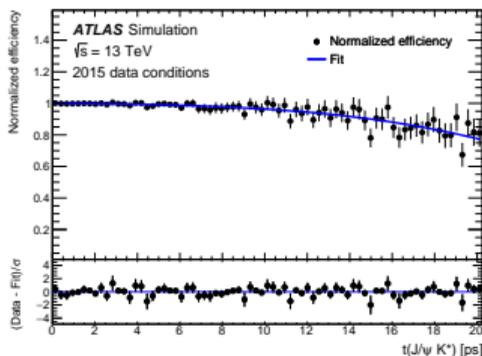
Left:  $\sigma_{\tau_i}$  for signal and Background. Middle and Right: 2D conditional probability distribution of the per-candidate  $\sigma_{\tau_i}$  and  $p_{\Gamma_i}$  of  $B_d^0$  candidates, illustrating the strong dependence between these observables for both (a) signal and (b) background. The sPlot technique [Pivk:2004ty](#) is used to separate the signal and background distributions in data.

- Trigger, offline reconstruction, event selections bias reconstructed proper-decay time distribution.
- Trigger and offline tracking impose  $|d_0| < 10$  mm, for all four final-state tracks of  $B^0 \rightarrow J/\psi K^{*0}$ , resulting in inefficiency at large times.
- Inefficiency effects are determined by signal MC, passed through simulation of detector response & triggers, followed by offline tracking, vertexing and event selections - as data.

Ratio of proper decay time distributions before and after the whole chain, were fitted by:

$$1/w(t_i) = p_0 \cdot [1 - p_1 \cdot (\text{Erf}((t_i - p_3)/p_2) + 1)] \quad (4)$$

$w(t_i)$  are used to re-weight each event in Likelihood fit. The fit is applied individually for each year, due to their different trigger and data-taking conditions.



Source of uncertainty	Systematic uncertainty [ps]
ID alignment	0.00108
Choice of mass window	0.00104
Time efficiency	0.00130
Best-candidate selection	0.00041
Mass fit model	0.00152
Mass-time correlation	0.00229
Proper decay time fit model	0.00010
Conditional probability model	0.00070
Fit model test with pseudo-experiments	0.00002
<b>Total</b>	<b>0.0035</b>

# Systematic uncertainties: Inner detector alignment

- ID misalignment effects: dominated by global length scale biases originating from ID geometry radial and longitudinal distortions along track trajectory.
- These manifest themselves as a shift in the reconstructed masses of known resonances, e.g.  $J/\psi \rightarrow \mu^+ \mu^-$  [Eur. Phys. J. C 80 \(2020\) 1194](#) ↗.

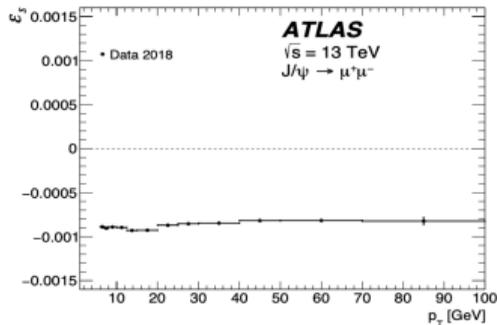
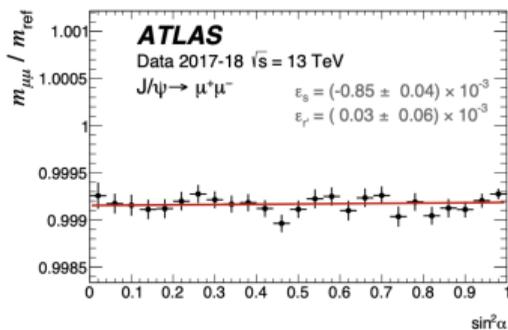


Fig left: Kinematic variable  $\sin^2_{\alpha} = E^+ \cdot E^- \cdot (p_T^+ / E^+ - p_T^- / E^-) / m(\mu^+ \mu^-)$  introduced to allow differentiation between  $\epsilon_r$  and  $\epsilon_z$ .

Scale biases  $\epsilon_S = \epsilon_z$  and  $\epsilon' = \epsilon_r - \epsilon_z$  measured from fit. Difference between the radial and longitudinal biases negligible.

Fig right:  $\epsilon_S$  as a function of track  $p_T$

- In our analysis  $B_d^0$  tracks are re-fitted with the  $J/\psi$  mass constrained to PDG, effectively removing the misalignment effect from the data. The impact of misalignment - estimated by alternative fit:  $B_d^0$  vertex re-fitted without PDG constrain.
- Additionally, to account for the momentum bias affecting hadrons from  $K^{*0}$  - their  $p_T$  are altered by  $-0.085\%$  [Eur. Phys. J. C 80 \(2020\) 1194](#) ↗.
- The two effects summed in quadrature, give the systematics:  $0.9 \sigma$

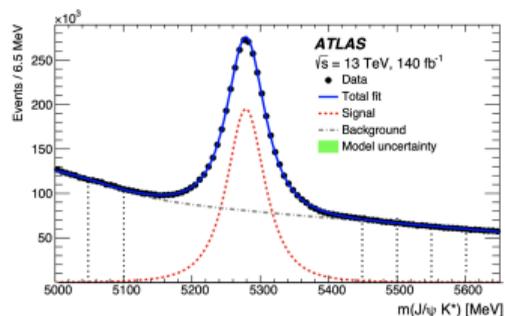
## Kinematic reflections

- MC models of not negligible kinematic reflections  $B^+ \rightarrow J/\psi K^+$ ,  $B_s \rightarrow J/\psi K^+ K^-$  and  $B^0 \rightarrow J/\psi \rho(\pi^+ \pi^-)$ , are included in the alternative fit to data with freely floating fractions.

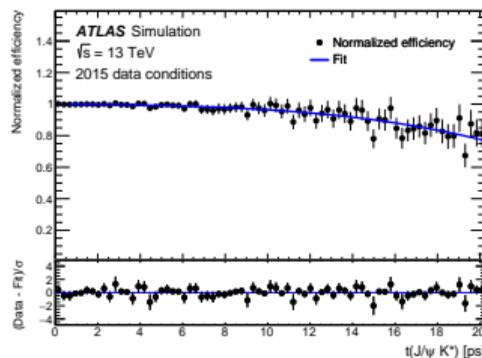
## Varying Sig, Bg mass PDF

- Alternative signal mass models, the Johnson  $S_U$ -distribution is replaced by a double-sided Crystal Ball distribution or Student's  $t$ -distribution, giving systematics
- Alternative background mass models: the sigmoid component is replaced by arctan function
- Both items summed give systematics  $1.3 \sigma$

- Correlations between the invariant mass and the pseudo-proper decay time, and their potential impact on the fit results, are studied.
- No mass-time correlation in signal PDF proven by MC.
- For background the correlations are studied in data. First, applying the default time Background PDF  $P_{\text{bkg}}(t_i, \sigma_i, \rho_{T_i})$  fit in each of the 6 mass sideband bins in data. Fractions  $f_{\text{prompt}}$  and  $b_1, b_2$  determined in each bin were then fitted to extract their dependence on mass. The best description - achieved by linear functions.



- Based on this information, an alternative fit model, in which the background PDF term  $P_{\text{bkg}}(t_i, \sigma_i, \rho_{T_i})$  is constructed with the parameters accounting for mass-dependence:  
 $f_{\text{prompt}}(m_i) = a + b(m_i - 5.279 \text{ GeV})$ ,  $b_1(m_i) = c_1 + d_1(m_i - 5.279 \text{ GeV})$  and  
 $b_2(m_i) = c_2 + d_2(m_i - 5.279 \text{ GeV})$ .  $a, b, c_1, d_1, c_2$  and  $d_2$  are free parameters of the fit. A difference of lifetime w.r.t default fit is 0.00228 ps ( $1.9 \sigma$ )

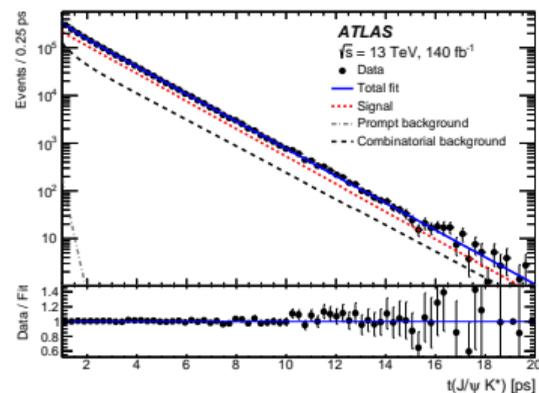
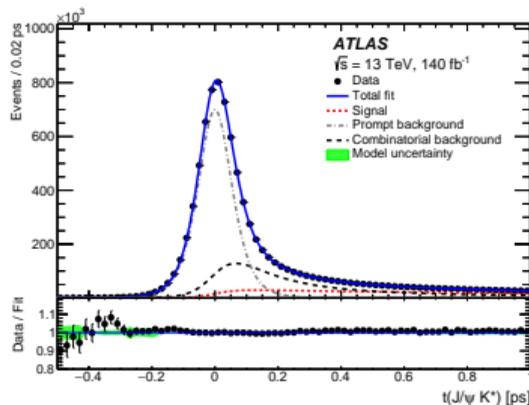
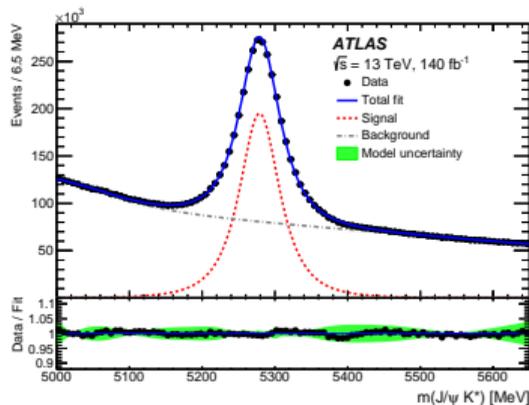


- Two alternative functions fitting the time efficiency histogram, replacing the default error function, were: hyperbolic tangent function or  $(x^2 + 1)^{-1/2}$ . The larger of the changes in the fitted lifetime, of size  $0.6\sigma_{\text{stat}}$  - taken as a systematic uncertainty.
- Systematic effects due to limited MC statistic, used to build the time efficiency histogram, was estimated by repeating the fit with a large number of alternative time efficiency functions, obtained by smearing the number of MC events in the time bins, leading to  $0.8\sigma_{\text{stat}}$ .
- An alternative fit using data events with  $t < 8$  ps, is performed to validate the modelling of the efficiency for high lifetimes, where the efficiency decrease is large, lead to small systematics  $0.5\sigma_{\text{stat}}$ .

# Results: The $B^0$ effective lifetime and the mass and time projections of the likelihood fit

The  $B^0$  effective lifetime value measured with a total of  $2\,450\,500 \pm 2400$   $B^0 \rightarrow J/\psi K^{*0}$  signal events. The measured effective lifetime is

$$\tau = 1.5053 \pm 0.0012 \text{ (stat.)} \pm 0.0035 \text{ (syst.) ps.}$$



- Mass fit projection (left). Proper decay time fit projections in two different ranges:  $(-0.5; 1.0)$  ps (Middle) and  $(1; 20)$  ps (Right).
- Solid blue line - total fit, dashed red line - signal.
- The lower panels: ratio of data point to the fit value. The green band - the envelope of model variations included in the systematic uncertainty, the bars on the data points indicate statistical uncertainties. Plot -right - the model variation band too small to be visible.

# Parameters fitted in mass-time Likelihood fit

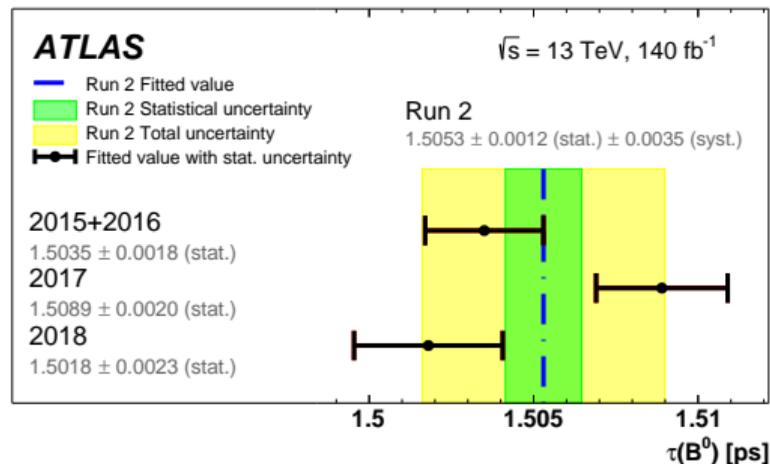
Parameter	Starting Value	Fitted Value	Error
$p_0$	-0.00015	-0.00015	1.8018e-07
$m_0$	4755.77315	5067.68398	2.2599
$s$	0.00940	0.00990	0.00028
$\delta_{sig}$	1.26524	1.10875	0.00388
$\gamma_{sig}$	-0.05542	-0.00000	2e-05
$\lambda_{sig}$	40.57214	36.27316	0.13244
$m_0$	5277.87937	5278.96001	0.0285
$f_{prompt}$	0.60000	0.93427	0.00138
$f_{prompt}$	0.60000	0.63547	0.00048
$f_{res;1}$	0.82000	0.83456	0.00099
$f_{res;2}$	0.16000	0.15896	0.00096
$f_{sig}$	0.22000	0.23175	0.00022
$f_{bkg;1}$	0.19891	0.29218	0.00387
$f_{bkg;2}$	0.41310	0.41696	0.00376
$\tau_{bkg;1}$	0.50000	0.41574	0.00589
$\tau_{bkg;2}$	1.54400	1.57023	0.00483
$\tau_{bkg;3}$	0.13970	0.12932	0.00128
$SF_1R$	1.00000	0.99267	0.00068
$SF_2R$	1.70000	1.74252	0.00438
$SF_3R$	5.40000	5.35241	0.04473
$\tau(B^0)$	1.51600	1.50529	0.00116

There are no high correlations of the fitted  $B_d^0$  lifetime with other parameters. The highest ones are:

- 18% with  $f_{sig}$  fraction of signal events
- 12% with  $\tau_{bkg_2}$  - the lifetime of one of three components of Combinatorial Background, that has the fitted value  $\tau_{bkg_2} = 1.5702$  ps rather close to the signal one.

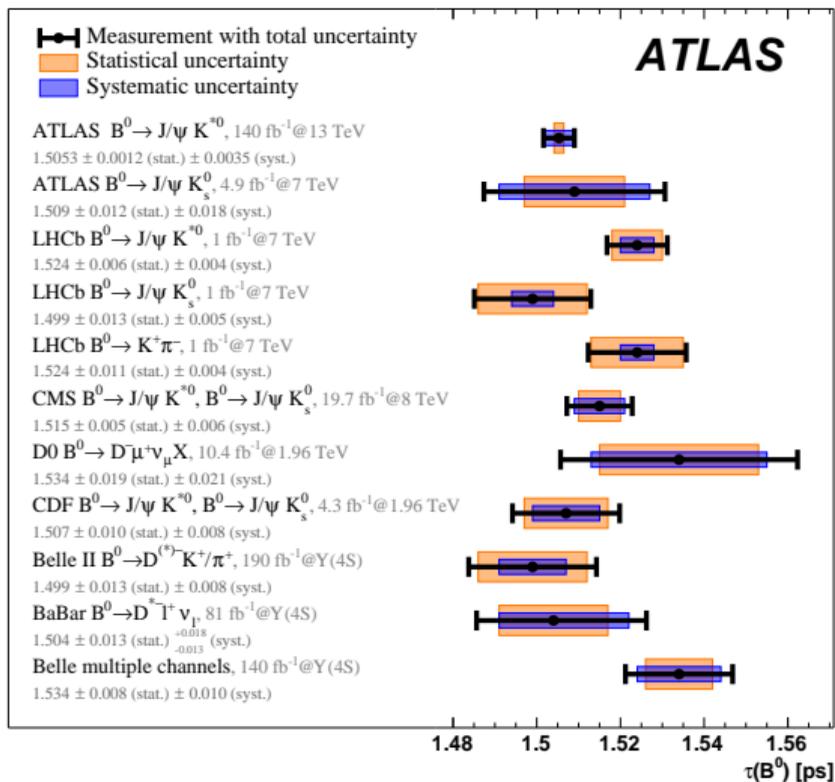
## Results: A consistency and stability test over data taking periods

As a consistency and stability test, the  $B^0$  lifetime value was fitted separately for each data-taking period (2015+2016, 2017 and 2018). Figure shows the degree of stability over time. The  $p$ -value for consistency of the three individual results, accounting for just statistical uncertainties is 0.038.



The fitted values of the  $B^0$  lifetime, measured with  $B^0 \rightarrow J/\psi K^{*0}$  decays, for the 2015+2016, 2017 and 2018 subsamples compared to the value for the whole sample. The  $B^0$  lifetime value for each subsample is shown by a black point, with the error bar indicating the statistical uncertainty.

# A comparison ATLAS $B^0$ lifetime with the latest results of other experiments.



- The current ATLAS result in  $B^0 \rightarrow J/\psi K^{*0}$  channel.
- The previous ATLAS result [Phys. Rev. D 87 \(2013\) 032002](#) in the  $B^0 \rightarrow J/\psi K_S^0$  channel.
- Latest LHCb results [JHEP 04 \(2014\) 114](#) in  $B^0 \rightarrow J/\psi K^{*0}$  and  $B^0 \rightarrow J/\psi K_S^0$  decays and [Phys. Lett. B 736 \(2014\) 446](#) in the  $B^0 \rightarrow K^+ \pi^-$ .
- Latest CMS [Eur. Phys. J. C 78 \(2018\) 457](#) combined result for  $B^0 \rightarrow J/\psi K^{*0}$  and  $B^0 \rightarrow J/\psi K_S^0$  decays.
- Tevatron experiments: D0 [Phys. Rev. Lett. 114 \(2015\) 062001](#) in the  $B^0 \rightarrow D^- \mu^+ \nu_\mu$  channel, and CDF [Phys. Rev. Lett. 106 \(2011\) 121804](#) with a combined result for  $B^0 \rightarrow J/\psi K^{*0}$  and  $B^0 \rightarrow J/\psi K_S^0$ .
- $e^+ e^-$  colliders: Belle II [Phys. Rev. D 107 \(2023\) L091102](#) in the  $B^0 \rightarrow D^{(*)-} K^+ / \pi^+$  channel and the last result from BaBar [Phys. Rev. D 73 \(2006\) 012004](#) in the  $B^0 \rightarrow D^{*+} \ell^- \bar{\nu}_\ell$ . Belle [PhysRevD.71.07990, 2005](#) this combination includes  $B_d^0$  decays to  $D^{*-} \ell^+ \nu$ ,  $D^{*-} \pi^+$ ,  $D^- \pi^+$ ,  $D^{*-} \rho^+$ ,  $J/\psi K^{*0}$ ,  $J/\psi K_S^0$ .

# Results: Determination of the $B_d^0$ average decay width $\Gamma_d$ and the ratio $\Gamma_d/\Gamma_s$

## $\Gamma_d$

- We determine  $\Gamma_d$  from our measured effective lifetime  $\tau_{B^0}$ , using Eq (1) and input values  $2y = \Delta\Gamma_d/\Gamma_d = 0.001 \pm 0.010$  and asymmetry  $A = -0.578 \pm 0.136$  from [HFLAV:2023](#) 

$$\Gamma_d = 0.6639 \pm 0.0005 \text{ (stat.)} \pm 0.0016 \text{ (syst.)} \pm 0.0038 \text{ (ext.) ps}^{-1}$$

The uncertainty denoted 'ext.' originates from the HFLAV.

- The value  $\Gamma_d$  is in agreement with HQE theory of  $0.63_{-0.07}^{+0.11}$  ps<sup>-1</sup> [Lenz et al. 2023](#) 

## $\Gamma_d/\Gamma_s$

- Using  $\Gamma_s = 0.6703 \pm 0.0014$  (stat.)  $\pm 0.0018$  (syst.) ps<sup>-1</sup> measured by the ATLAS [Eur. Phys. J. C 81 \(2021\) 342](#) 

$$\Gamma_d/\Gamma_s = 0.9905 \pm 0.0022 \text{ (stat.)} \pm 0.0036 \text{ (syst.)} \pm 0.0057 \text{ (ext.)}$$

- the statistical, systematic and external uncertainties are propagated from the quantities above. In  $\Gamma_d/\Gamma_s$  - systematic uncertainties of the ATLAS measurements of  $\tau_{B^0}$  and  $\Gamma_s$  primarily come from different sources. They are therefore treated as uncorrelated.
- $\Gamma_d/\Gamma_s$  agrees with theory HQE and lattice QCD models prediction.

- ATLAS performed a measurement of  $B_d^0$  effective lifetime and the average decay width  $\Gamma_d$  using  $B^0 \rightarrow J/\psi K^{*0}$  events reconstructed from a  $140 \text{ fb}^{-1}$  data sample of  $pp$  collisions collected with the ATLAS detector during the  $\sqrt{s}=13 \text{ TeV}$  LHC run.
- The  $B^0$  effective lifetime is measured to be  $1.5053 \pm 0.0012(\text{stat}) \pm 0.0035(\text{syst}) \text{ ps}$ . This result is compatible with other experimental measurements and is the most precise measurement to date.
- The measured average decay width of the heavy and light  $B^0$  mass eigenstates is  $\Gamma_d = 0.6639 \pm 0.0005 \text{ (stat.)} \pm 0.0016 \text{ (syst.)} \pm 0.0038 \text{ (ext.) ps}^{-1}$ . This value is in good agreement with the theory prediction.
- The measured average decay width  $\Gamma_d$  is combined with the average decay width  $\Gamma_s$  measured previously by ATLAS to obtain the ratio  $\Gamma_d/\Gamma_s = 0.9905 \pm 0.0022 \text{ (stat.)} \pm 0.0036 \text{ (syst.)} \pm 0.0057 \text{ (ext.)}$ . This result is compatible with the theory predictions from HQE and lattice QCD calculations as well as with the experimental average.

# Backup Slides

- ATLAS  $B^0$  lifetime  $1.5053 \pm 0.0012(\text{stat}) \pm 0.0035(\text{syst})$  ps, is the most precise measurement to date.
  - Comparing to latest LHC measurements
    - LHCb  $B^0 \rightarrow J/\psi K_S^0$  lifetime is smaller than ATLAS and well compatible within  $0.4 \sigma$ . The other two LHCb results  $B^0 \rightarrow J/\psi K^{*0}$  and  $B^0 \rightarrow K^+ \pi^-$  have lifetimes larger than ATLAS by  $2.3 \sigma$  and  $1.5 \sigma$ . The CMS combining  $B^0 \rightarrow J/\psi K_S^0$   $B^0 \rightarrow J/\psi K^{*0}$  is within  $1.1 \sigma$  from ATLAS.
  - Comparing to non LHC experiments
    - CDF and Belle II results combining the  $B^0 \rightarrow J/\psi K_S^0$   $B^0 \rightarrow J/\psi K^{*0}$  are compatible with ATLAS within  $0.1 \sigma$  and  $0.4 \sigma$  respectively, while the older Belle result, combining six channels, has lifetime larger than ATLAS by  $2.1 \sigma$ .
  - The latest world average
    - $1.517 \pm 0.004$  ps [PDG2024](#) differs from ATLAS by  $2.1 \sigma$ .
- ATLAS  $\Gamma_d = 0.6639 \pm 0.0005$  (stat.)  $\pm 0.0016$  (syst.)  $\pm 0.0038$  (ext.)  $\text{ps}^{-1}$  value
  - is compatible with the HQE theory of  $0.63_{-0.07}^{+0.11} \text{ps}^{-1}$  [Lenz et al. 2023](#) within  $0.3 \sigma$ ,  $\sigma$  combining both sources.
- ATLAS  $\Gamma_d/\Gamma_s = 0.9905 \pm 0.0022$  (stat.)  $\pm 0.0036$  (syst.)  $\pm 0.0057$  (ext.)
  - is compatible with the theory HQE and lattice QCD models [Lenz et al. 2023](#) within  $1.4 \sigma$  and  $0.4 \sigma$ , respectively, and with the experimental average  $(1.001 \pm 0.004)$  [HFLAV:2023](#) within  $1.3 \sigma$

- Measurement of effective lifetime in  $B_{(s)}^0 \rightarrow \mu^+ \mu^-$   
[JHEP 09 \(2023\) 199 W](#) 
- Measurement of  $\Lambda_b^0$  lifetime in  $\Lambda_b^0 \rightarrow J/\psi \Lambda^0$  and  $B_d^0$  lifetime in  $B_d^0 \rightarrow J/\psi K_S^0$  decays  
[PhysRevD.87.032002, 2013](#) 
- Measurement of the relative width difference  $\Delta\Gamma_d/\Gamma_d$  of the  $B_d^0$ - $\overline{B}_d^0$  system with the ATLAS detector  
[JHEP06\(2016\)081](#) 
- Measurement of the CP-violating phase, the width difference  $\Delta\Gamma_s$  between the meson mass eigenstates and the average decay width  $\Gamma_s$  in the  $B_s \rightarrow J/\psi \phi$  decay.  
[Eur. Phys. J. C 81 \(2021\) 342](#) 