

# Accessing the strong interaction in three-hadron systems with ALICE

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#### **Bhawani Singh**

Technical University of Munich (TUM), Germany On behalf of the ALICE Collaboration (based on Phys. Rev. X 14, 031051 (2024)) LHC Seminar, 26.11.2024

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Non-perturbative QCD  $\rightarrow Q \sim 1 \text{ GeV}$ 

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- Non-perturbative QCD  $\rightarrow Q \sim 1 \text{ GeV}$
- Use effective field theories (residual strong interaction)
  - Hadrons as degrees of freedom (baryons, mesons)





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- Use effective field theories (residual strong interaction)
  - Hadrons as degrees of freedom (baryons, mesons)
  - Need for experimental data of hadronic interactions
  - Constrain low-energy constants in the EFTs



Two-body interaction

Many-body interaction



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#### Nuclei/hypernuclei





 Explanation for nucleon-deuteron scattering observables: requires the presence of three-body interaction<sup>[1]</sup>

[1] K. Sekiguchi, Few-Body Syst 60, 56 (2019)







#### Nuclei/hypernuclei



 $\rho_0$ 

- Explanation for nucleon-deuteron scattering observables: requires the presence of three-body interaction<sup>[1]</sup>
- 3-body interaction contributes ~10% to the binding energies of light nuclei<sup>[2]</sup>

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#### Dense nuclear matter: neutron star













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- NNN and NNA interactions used in the modeling of the equation of the state of neutron stars<sup>[3-4]</sup>
- We need new tools to study three-body hadronic interaction

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#### Dense nuclear matter: neutron star









# Access interaction three-hadron system with hadron-deuteron correlation











$$C(k^*) = \mathcal{N}\frac{N_{\text{same}}(k^*)}{N_{\text{mixed}}(k^*)} = \int S(\vec{r}^*) \left| \psi(\vec{k}^*, \vec{r}^*) \right|^2$$
  
experimental definition theoretical definition

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S.E. Koonin et al, Phys. Lett. B 70 43 (1977) L. Fabbietti et al, Ann. Rev. Nucl. Part.Sci. 71 (2021) 377-402





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$$d^3 r^* \xrightarrow{k^* \to \infty} 1$$

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### The emission source in pp collisions

- Source modeling is based on
  - Emission of all primordial particles (Gaussian)

$$S(r^{*}) = \frac{1}{(4\pi r_{\rm core}^{2})^{3/2}} exp\left(-\frac{r^{*2}}{4r_{\rm core}^{2}}\right)$$



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Gaussian core source



### The emission source in pp collisions

- Source modeling is based on
  - Emission of all primordial particles (Gaussian)
  - Short-lived resonances ( $c\tau \sim 1 \text{ fm} : \Delta, N^*, \Sigma^*$ )

$$S(r^*) = \frac{1}{(4\pi r_{\rm core}^2)^{3/2}} exp\left(-\frac{r^{*2}}{4r_{\rm core}^2}\right)$$

 $\oplus$  resonance tail

- **Resonance contributions** 
  - Dependent on the particle species
  - Fixed from the statistical hadronization model<sup>[1]</sup> and EPOS<sup>[2]</sup>
  - r<sub>core</sub> : particle-emitting source can be studied using particle pairs with known interaction



Gaussian core  $\oplus$  resonance contributions

[1] V. Vovchenko et al, Comput. Phys. Comm. 244 (2019)

[2] T. Pierog et al, Phy. Rev. C 92, 034906 (2015)







### A common source for all hadrons in pp collisions

- Particle-emitting sources studied with
  - well-known hadronic interaction
  - p-p ALICE, Phys. Lett. B 811 135849 (2020)
  - p-K<sup>+</sup> <u>ALICE</u>, arXiv:2311.14527
  - $\pi^{\pm} \pi^{\pm} ALICE$ , arXiv:2311.14527
  - $p-\pi^{\pm}$  (paper in preparation)





**ALI-PREL-576328** 

$$m_{\rm T} = \sqrt{\bar{m}^2 + k_{\rm T}^2} \text{ and } \vec{k}_{\rm T} = \frac{1}{2} \left( \vec{p}_{\rm T,1} + \vec{p}_{\rm T,2} \right)$$



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  - $\pi^{\pm} \pi^{\pm}$  ALICE, arXiv:2311.14527
  - $p-\pi^{\pm}$  (paper in preparation)
- A common primordial source for all hadrons in high-multiplicity pp collisions!
- Use the source size for particle pairs with unknown interaction
- Possibility to study interaction for exotic pairs



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### A common source for all hadrons in pp collisions

#### **Recent ALICE femtoscopy measurements**

PRC 99 (2019) 2, 024001 PLB 797 (2019), 134822 PRL 123 (2019), 112002 PRL 124 (2020) 092301 PLB 805 (2020), 135419 PLB 811 (2020), 135849 Nature 588 (2020) 232-238 PRL 127 (2021), 172301 PLB 822 (2021) 136708 PRC 103 (2021) 5, 055201 PLB 833 (2022), 137272 PLB 829 (2022), 137060 PRD 106 (2022) 5, 052010 PLB 833 (2022) 137335 PLB 844 (2022), 137223 EPJA 59 (2023) 7, 145 EPJA 59 (2023) 12, 298 EPJC 83 (2023) 4, 340 PLB 845 (2023), 138145 PRD 110, 032004 PRX 14 (2024) 3, 031051

 $p-p, p-\Lambda, \Lambda-\Lambda$  (methods)  $\Lambda - \Lambda$ p-Ξ p–K **p**–Σ<sup>0</sup> р–р, р–∧  $p-\Xi, p-\Omega$ p-¢ p–K  $\Lambda - K$ p–∧ baryon-(anti)baryon p–D K<sup>0</sup>–K<sup>0</sup>, K<sup>ch</sup>–K<sup>0</sup> ∧–Ξ р–р–р, р–р–∧ p-p-K p–K  $\Lambda - K$  $K/\pi -D$ p-d, K-d

Possibility to study interaction for exotic pairs





ALI-PREL-576328

$$m_{\rm T} = \sqrt{\bar{m}^2 + k_{\rm T}^2} \text{ and } \vec{k}_{\rm T} = \frac{1}{2} \left( \vec{p}_{\rm T,1} + \vec{p}_{\rm T,2} \right)$$



### **Today: three-hadron systems**



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Hadron-deuteron correlations provide an indirect way to study the strong interaction in system of three hadrons



### K<sup>+</sup>-d and p-d systems



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### So far... hadron-deuteron correlations

At very low energy (~ GeV beam energy), fixed target experiments<sup>[1-4]</sup>



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[1] C. B. Chitwood et al, Phys. Rev. Lett. 54, 302 (1985)

[2] J. Pochodzalla et al, Phys. Rev. C 35, 1695 (1986)

[3] J. Pochodzalla et al, Phys. Lett. B 175 (1986)

[4] K. Wosinska et al, Eur. Phys. J. A 32, 55–59 (2007)



### So far... hadron-deuteron correlations

At very low energy (~ GeV beam energy), fixed target experiments<sup>[1-4]</sup>



- No **full-fledged calculations** and unconstrained source distributions

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### A Large Ion Collider Experiment

Excellent tracking and particle identification (PID) capabilities

Inner Tracking System (ITS) Tracking, vertex

**Time Projection Chamber (TPC)** Tracking, PID (dE/dx)

Time Of Flight detector (TOF) PID (TOF measurement)

- Run 2 data-set
  - 10<sup>9</sup> high-multiplicity pp collisions at  $\sqrt{s} = 13$  TeV





#### ALICE : <u>ITS</u> and <u>TPC</u> upgrades

Int.J.Mod.Phys.A 29 (2014) 1430044 JINST 3 (2008) S08002





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ALI-PERF-131248







# A Large Ion Collider Experiment

- Excellent tracking and particle identification (PID) capabilities Inner Tracking System (ITS) Tracking, vertex Time Projection Chamber (TPC) Tracking, PID (dE/dx) Transition Radiation Detector (TRD) Time Of Flight detector (TOF) PID (TOF measurement) Run 2 data-set
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### Hadron-deuteron correlations in pp collisions at LHC







The femtoscopic correlation consists of various contributions<sup>[1-2]</sup> 



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### $C_{\text{femto}}(k^*) = \lambda_{\text{gen}} C_{\text{gen}}(k^*) \bigoplus \lambda_{\text{feed}} C_{\text{feed}}(k^*) \bigoplus \lambda_{\text{misid}} C_{\text{misid}}(k^*) \bigoplus \dots$

[1] D. Mihaylov et al. Eur. Phys. J. C78 (2018) 394 [2] R. Lednicky, Phys. Part. Nuclei 40, 307–352 (2009)





- The femtoscopic correlation consists of various contributions<sup>[1-2]</sup>
  - Genuine interaction from primordial particle





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Particles from material knock-outs and misidentifications



- - Purity of the individual particles ( $\mathscr{P}_i$ ) and feed-down fractions ( $f_i$ )

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A data-driven approach to quantify contributions (lambda parameters  $\lambda_{ij} = \mathcal{P}_i \cdot f_i \times \mathcal{P}_j \cdot f_j$  with  $\sum \lambda_{ij} = 1$ )

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### Source for kaon-deuteron and proton-deuteron pairs

Primordial source size for K<sup>+</sup>–d and p–d systems 



Source size	mean value: p–d	mean value: K+–d
r <sub>core</sub>	0.99±0.05 fm	1.04±0.04 fm

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### Source for kaon-deuteron and proton-deuteron pairs

- Primordial source size for K<sup>+</sup>–d and p–d systems
- Source radius is effectively increased by short-lived strongly decaying resonance



Source size	mean value: p–d	mean value: K+–d
r <sub>core</sub>	0.99±0.05 fm	1.04±0.04 fm
r <sub>eff</sub>	1.08±0.06 fm	1.35±0.05 fm

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### **Theoretical approach correlation functions**

Potential approach: solve Schrödinger equation for the two-hadron system<sup>[1]</sup> 



[1] D. Mihaylov et al. Eur. Phys. J. C78 (2018) 394 [2] R. Lednicky, Phys. Part. Nuclei 40, 307–352 (2009)





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- Potential approach: solve Schrödinger equation for the two-hadron system<sup>[1]</sup>
- Lednický-Lyuboshits approach: for pointlike distinguishable particles and asymptotic wavefunction<sup>[2]</sup>
  - Only **s-wave** two-particle relative wave function
  - **Considers Coulomb effects**

 $\psi_{-k^*}(r^*) = e^{i\delta_c} \sqrt{A_c(\eta)} e^{-ik}$ 



$$\tilde{K}^{*r*}F(-i\eta,1,i\zeta) + f_c\left(k^*\right)\frac{\tilde{G}(\rho,\eta)}{r^*}$$

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- $f_c$  Coulomb normalized scattering amplitude for strong interaction, F and  $\tilde{G}$  are Coulomb functions
  - *a*<sub>0</sub>: scattering length
  - $d_0$ : effective range

[1] D. Mihaylov et al. Eur. Phys. J. C78 (2018) 394 [2] R. Lednicky, Phys. Part. Nuclei 40, 307–352 (2009)





### Kaon-deuteron correlation function

- Assuming  $m_T$ -scaling holds for d,  $r_{eff} = 1.35 \pm 0.05$  fm
- Coulomb potential: disagree

[1] R. Lednicky', Phys. Part. Nuc. 40, (2009)

[2] provided by Prof. Johann Haidenbaur

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### Kaon-deuteron correlation function

- Assuming  $m_{\rm T}$ -scaling holds for d,  $r_{\rm eff} = 1.35 \pm 0.05$  fm
- Coulomb potential: disagree
- Strong interaction in  $K^+$ –d as an **effective two-body** system: Lednický-Lyuboshits approach<sup>[1]</sup>
  - Effective-Range approx. (ER):  $a_0 = -0.47$  fm,  $d_0 = -1.75$  fm<sup>[2]</sup>
  - Fixed-center approx. (FCA):  $a_0 = -0.54$  fm,  $d_0 = 0$  fm<sup>[3]</sup>

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- **Deuterons follow the m<sub>T</sub>-scaling, and an effective two-body** approach can describe the K<sup>+</sup>-d system



[1] R. Lednicky', Phys. Part. Nuc. 40, (2009)

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# **Proton-deuteron correlation function**

- Assuming p–d as an effective two-body: Lednický-Lyuboshits approach<sup>[1]</sup>
- Source size  $r_{\text{eff}} = 1.08 \pm 0.06$  fm
- Strong interaction: constrained from the scattering measurements<sup>[2]</sup>



[1] R. Lednicky', Phys. Part. Nuc. 40, (2009)

[2] Scattering parameters from N-d scattering





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- Source size  $r_{\text{eff}} = 1.08 \pm 0.06$  fm
- Strong interaction: constrained from the scattering measurements<sup>[2]</sup>
- A point-like particle within the LL approach does not work
  - Pauli-blocking for p–(pn) system
  - Asymptotic strong interaction insufficient



- [1] R. Lednicky', Phys. Part. Nuc. 40, (2009)
- [2] Scattering parameters from N-d scattering





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[2] Scattering parameters from N-d scattering

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Start from p–(pn) system that forms p–d state asymptotically: 

$$C_{pd}(k^{*}) = \frac{1}{6A_{d}} \sum_{m_{1},m_{2}} \int d^{3}r_{1}d^{3}r_{2}d^{3}r_{3}S_{1}(r_{1})S_{1}(r_{2})S_{1}(r_{3}) \left|\Psi_{m_{1},m_{2},k^{*}}\right|^{2}$$
$$= \frac{1}{16A_{d}} \int S(\rho, R_{M}) \left|\Psi(k^{*},\rho)\right|^{2} \rho^{5}d\rho d\Omega$$

M. Viviani, **B. Singh** et al. Phys. Rev. C 108, 064002 (2023) Michele Viviani, Alejandro Kievsky, and Laura Marcucci from Pisa group Sebastian König from NC state University

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$$= \frac{1}{16A_{d}} \int S(\rho, R_{M}) \left|\Psi(k^{*},\rho)\right|^{2} \rho^{5}d\rho d\Omega$$

A<sub>d</sub>: deuteron formation probability<sup>[1]</sup>

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[1] S. Mrowczynski, Acta Physica Polonica B 51, 1739 (2020)





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$$= \frac{1}{16A_d} \int S(\rho, R_M) \left| \Psi(k^*, \rho) \right|^2 \rho^5 d\rho d\Omega$$
$$R_M = 1.43 \pm 0.16 \text{ fm}$$

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$$= \frac{1}{16A_d} \int S(\rho, R_M) \left| \Psi(k^*, \rho) \right|^2 \rho^5 d\rho d\Omega$$

$$R_M = 1.43 \pm 0.16 \text{ fm}$$

$$\Psi(k^*, \rho) \text{ the three-nucleon}$$

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 $\Psi(\kappa^{\star}, \rho)$  the three-nucleon wave function

[1] S. Mrowczynski, Acta Physica Polonica B 51, 1739 (2020)







### Asymptotic form of strong interaction in p-d system

Coulomb only: does not describe the data

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## Asymptotic form of strong interaction in p-d system

- Coulomb only: does not describe the data
- Born approximated wavefunction NN<sup>[1-2]</sup> NNN potentials<sup>[3]</sup>
  - Perform antisymmetrization
  - Approximate the wavefunction by ignoring centrifugal core interaction
  - Asymptotic form of strong interaction is insufficient to capture the dynamics of nucleons ~ 1 fm

[1] M. Viviani, **B. Singh** et al. Phys. Rev. C108,064002 (2023)

[2] AV 18 NN potential: R. B. Wiringa et al. Phys. Rev. C 51, 38 (1995)

[3] UIX NNN potential: B. S. Pudliner et al. Phys. Rev. Lett. 74, 4396 (1995)

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### Two- and three-body interaction at short distance

- Full three-body dynamics at short distances using AV18+UIX potentials<sup>[1-3]</sup>
  - s-wave: undershoots due to repulsion in s-wave

[1] M. Viviani, **B. Singh** et al. Phys. Rev. C108,064002 (2023)

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B. Singh





### Two- and three-body interaction at short distance

- Full three-body dynamics at short distances using AV18+UIX potentials<sup>[1-3]</sup>
  - s-wave: undershoots due to repulsion in s-wave
  - All partial waves up to *d*-waves: excellent description ( $n\sigma \sim 1$  for  $k^*$  up to 400 MeV/c)

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- Pionless EFT NLO (s+p+d waves):
  - Agree with data within  $n_{\sigma} \sim 2$  for  $k^* < 120$  MeV/c

- Dynamics of the three-body p–(pn) system at short distances!
- Inclusion of the higher partial waves

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  - Ratio of CF with and without UIX potential<sup>[1]</sup>
  - Upto ~5% effects due to genuine three-body strong interaction
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Avenue for the study of hadron-deuteron systems, including charm and strange hadrons!





### Study correlations among three unbound hadrons (3 to 3 scattering process)

### Three-body femtoscopy

Study interaction in hadron-triplets via three-particle correlations 



B. Singh

$$C(Q_{3}) = N \frac{N_{\text{same}}(Q_{3})}{N_{\text{mixed}}(Q_{3})} \qquad C(Q_{3}) = \int S(\rho) \left| \Psi(Q_{3}, \rho) \right|^{2} \rho^{5} \, \mathrm{d}\rho$$
  
experimental definition <sup>[1-2]</sup> theoretical definition <sup>[3]</sup>



$$q_{23}^2 - q_{13}^2$$

 $\vec{p}_3$ 

 $\dot{p_2}$ 

 $\vec{p}$ 

$$\rho = 2\sqrt{r_{12}^2 + r_{23}^2 + r_{31}^2}$$

[1] ALICE Coll, Eur. Phys. J. A 59, 145 (2023) [2] R. Del Grande et al, Eur. Phys. J. C 82 (2022) 244 [3] A. Kievsky et al, Phys. Rev. C 109 (2024) 3, 034006





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Exp:

ALICE Coll., EPJ A 59, 145 (2023) ALICE Coll., EPJ A 59, 298 (2023)

Theory (Munich and PISA group)

- R. Del Grande et al. EPJC 82 (2022) 244
- M. Viviani et al, PRC 108 (2023) 6, 064002
- A. Kievsky, et al., PRC 109 (2024) 3, 034006
- B. E. Garrido et al., arXiv: 2408.01750 (2024)

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## Three-body femtoscopy with ALICE

Hadron-triplets via three-particle correlations: p-p-p and  $p-p-\Lambda$ 



Direct access to two- and three-body forces in p-p-p and  $p-p-\Lambda$  systems

B. Singh

LHC Seminar | Strong interaction in three-body systems

Projector: Del Grande et al, Eur. Phys. J. C 82, 244 (2022)



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## p-p-p correlation using AV18 potential

Three-body correlation function AV18 and UIX potential<sup>[1]</sup>

$$C(Q_3) = \int S(\rho) \left| \Psi(Q_3, \rho) \right|^2 \rho^5 d\rho$$

- $\Psi(Q_3, \rho)$  computed using **only pp AV18** strong interaction, Coulomb corrections, and quantum statistics
- Negligible contribution from NNN (via UIX) found < 1%
- Attractive AV18 interaction: results peak
- Pauli-blocking: depletion in C(Q<sub>3</sub>)







[1] A. Kievsky et al, Phys. Rev. C 109 (2024) 3, 034006



### p-p-p correlation in Run 3



By the end of Run 3, 100 times more triplets w.r.t Run 2 statistics, estimated with dedicated software triggers!

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[1] A. Kievsky et al, Phys. Rev. C 109 (2024) 3, 034006



### Theoretical $p-p-\Lambda$ correlation

- Three-particle emission source modeled as three single-particle emitters constrained to data<sup>[1]</sup>
- Modeling includes experimental corrections (e.g. feed-down)









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- Three-particle emission source modeled as three single-particle emitters constrained to data<sup>[1]</sup>
- Modeling includes experimental corrections (e.g. feed-down)
- The most interesting region  $Q_3 < 100 \text{ MeV/c}$ not yet accessed by data







![](_page_68_Picture_9.jpeg)

### Theoretical $p-p-\Lambda$ correlation

![](_page_69_Figure_1.jpeg)

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#### B. Singh

![](_page_69_Figure_6.jpeg)

![](_page_69_Picture_7.jpeg)

### **Conclusions and Outlook**

- **K<sup>+</sup>-d:** deuterons follow source size scaling for all hadrons in pp collisions
- p-d
  - Access to three-body strong interaction
  - Sensitive to the inclusion of higher partial waves

![](_page_70_Picture_5.jpeg)

![](_page_70_Figure_7.jpeg)

![](_page_70_Figure_8.jpeg)

![](_page_70_Picture_10.jpeg)

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- **p–p–** $\Lambda$ : 3-body force with strangeness up to 40%

![](_page_71_Picture_7.jpeg)

![](_page_71_Figure_11.jpeg)

![](_page_71_Picture_13.jpeg)
### **Conclusions and Outlook**

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- **p**–**p**–**p**: insignificant three-body force due to Pauli-blocking effects
- **p–p–** $\Lambda$ : 3-body force with strangeness up to 40%
- Large statistics of LHC Run 3 and Run 4
  - p-p correlation in LHC Run 3: source constrained for all interaction studies
  - Ongoing studies for p–d,  $\Lambda$ –d, p–p–p, and p–p– $\Lambda$  from LHC Run 3









### Thank you for your attention!

Credit: D. Chinellato, DOI 10.5281/zenodo.13284731



# Additional slides

## Asymptotic form of strong interaction in p-d system

- Coulomb only: does not describe the data
- Born approximated wavefunction AV18(2N) <sup>[1-2]</sup> UIX (NNN) potentials <sup>[3]</sup>
- Asymptotic form of strong interaction is insufficient to capture the dynamics of nucleons ~ 1 fm

$$\begin{split} \Psi_{LSJJ_z} &= \sum_{n,\alpha} \frac{u_{n,\alpha}(\rho)}{\rho^{5/2}} \mathcal{Y}_{n,\alpha}(\Omega) \\ &+ \frac{1}{\sqrt{3}} \sum_{\ell}^{\text{even perm.}} \left\{ Y_L(\hat{\boldsymbol{y}}_\ell) \left[ \varphi^d(i,j) \chi(\ell) \right]_S \right\}_{JJ_z} \frac{F_L(\eta, ky_\ell)}{ky_\ell} \\ &+ \sum_{L'S'} T_{LS,L'S'}^J \frac{1}{\sqrt{3}} \sum_{\ell}^{\text{even perm.}} \left\{ Y_{L'}(\hat{\boldsymbol{y}}_\ell) \left[ \varphi^d(i,j) \chi(\ell) \right]_{S'} \right\} \\ &\times \frac{\overline{G}_{L'}(\eta, ky_\ell) + iF_{L'}(\eta, ky_\ell)}{ky_\ell} \, . \end{split}$$

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## **Coulomb interaction in p-d system**

- Complete p-pn dynamics, but the strong interaction is absent at very short-range!
  - r<sup>NN</sup>eff =1.43±0.16 fm (nucleon-nucleon distance)
- In the case of the two-body picture Coulomb-only interaction differs from the one using the p-(pn) dynamics
- Two-body source 1.08±0.06 fm (proton-deuteron distance)
  - More repulsion due to the Pauli-blocking







## AV18+UIX vs NVIa3 3N Chiral potentials

Precise calcualtion using AV18+UIX as well NVIa3/3N chiral potentials

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M. Viviani, **B. Singh** et al. Phys. Rev. C 108, 064002 (2023)

## Theoretical $p-p-\Lambda$ correlation

- Three-particle emission source modeled as three single-particle emitters constrained to data <sup>[1]</sup>
- Modeling includes experimental corrections (e.g. feed-down)







## AV18+UIX vs NVIa3 3N Chiral potentials

- Comparisition with Chiral potentials (Full three-body dynamics)<sup>[1]</sup>
- Argonne v18+Urbana IX interaction<sup>[2,3]</sup>
- All partial waves upto d-waves: describes data within  $n_{\sigma} \sim 1$  for  $k^*$  up to 400 MeV/c
- Calculations using chiral potential from NVIa+3N
  - Very good agreement with AV18+UIX
- AV18 alone: just two-body NN interaction
- Current data cannot resolve the effect of three-body force





### The pA interaction in the femtoscopy

**Improvement:** combined analysis of femtoscopic and scattering data 



### B. Singh