# New Constraints on the Central Mass Contents of Omega Centauri from Combined Stellar Kinematics and Pulsar Timing

Andrés Bañares-Hernández

Co-authors: Francesca Calore, Jorge Martín Camalich, and Justin Read



#### Image credit: ESA/Gaia/DPAC (CC BY-SA 3.0 IGO)

#### Why is ω Cen interesting?

#### • Believed to be the nucleus of a disrupted dwarf galaxy

#### Ideal candidate for dark matter annihilation signals (only ~ 5 kpc away + high DM concentrations)

#### Argued to host a ~ 10<sup>4</sup> M<sub>O</sub> intermediate-mass black hole (IMBH)

#### ... and also a cluster of stellar-mass black holes

#### Has an abundant population of recently discovered pulsars

#### Mass contents from stellar kinematics

• Stellar velocity dispersions trace mass contents from the Jeans equations: 1  $\partial x \sigma^2 = 2R\sigma^2 = CM$ 

where: 
$$\frac{1}{\nu_{\star}} \frac{\partial \nu_{\star} \sigma_{r}^{2}}{\partial r} + \frac{2\beta \sigma_{r}^{2}}{r} = -\frac{GM}{r^{2}},$$
$$\sigma_{r/t}^{2} = \langle v_{r/t}^{2} \rangle - \langle v_{r/t} \rangle_{,}^{2} \quad \beta \equiv 1 - \frac{\sigma_{t}^{2}}{\sigma_{r}^{2}}$$

• Solve for the radial velocity dispersion for a given mass profile:

$$\sigma_r^2(r) = \frac{1}{\nu_\star(r)g(r)} \int_r^\infty \frac{GM(r')\nu_\star(r')}{r'^2} g(r')dr', \qquad g(r) \equiv \exp\left(2\int \frac{\beta(r)}{r} dr\right)$$

#### IMBH discovered in ω Cen?

- IMBH can produce a steep rise in central velocity dispersions
- However, studies claiming IMBHs have been susceptible to degeneracies and biases from limited data / modeling
- Concentrated cluster of remnants can produce degenerate effects

GEMINI AND HUBBLE SPACE TELESCOPE EVIDENCE FOR AN INTERMEDIATE-MASS BLACK HOLE IN  $\omega$  CENTAURI

EVA NOYOLA<sup>1</sup> AND KARL GEBHARDT Astronomy Department, University of Texas at Austin, Austin, TX 78712; noyola@mpe.mpg.de

AND

MARCEL BERGMANN Gemini Observatory, Tucson, AZ 85726 Received 2006 May 14; accepted 2007 December 22

# Black hole found in Omega Centauri

02/04/2008 9427 VIEWS 12 LIKES

# • Doppler-induced changes in apparent period derivatives trace LOS $a_{\text{LOS}}(r,l) = -G\left(\frac{l}{r}\right)\frac{M(r)}{r^2}$

Constraining mass models with pulsar timing

 $\dot{P}_{\rm obs} = \frac{(a_{\rm LOS} + a_{\rm S} + a_{\rm g})P}{c} + \dot{P}_{\rm int}$ 

acceleration due to GC potential

• Intrinsic spin-down can be modeled as magnetic dipole emission:

$$a_{\text{int}} \equiv c \left(\frac{\dot{P}}{P}\right)_{\text{int}} = 7.96 \times 10^{-10} \left(\frac{B}{2 \times 10^8 \text{ G}}\right)^2 \left(\frac{2\text{ms}}{P}\right)^2 \text{ m s}^{-2}$$

$$\dot{P}_{\mathrm{int}}$$

Image credit: Prager et al. 2017

# Generalized multi-component mass modeling

• Fully explore mass degeneracies by including both an IMBH and extended remnant and stellar distributions:

$$M(r) = M_{\star}(r) + M_{\rm cen}(r) + M_{\rm BH} + M_p(r)$$

- IMBH modeled as a point mass
- Central cluster of remnants as a Plummer component:

$$M_{\rm cen}(r) = M_{\rm cen} \frac{r^3}{r_{\rm cen}^3} \left(1 + \frac{r^2}{r_{\rm cen}^2}\right)^{-3/2}$$

- Stellar profile set by photometry
- Included a mass component traced by the pulsars

# Implementation

- Employed **GravSphere** (Read & Steger 2017) as a Jeans equations solver for stellar kinematics with full anisotropic modeling
- Extended GravSphere to include pulsar timing data and robust posterior sampling with **dynesty**
- Self-consistently binned, high-resolution data from HST + Gaia (PMs) and multiple ground-based observations (LOS)

## **Results: Velocity dispersions**

- Fit can reproduce the profile well
- Somewhat elevated, but flat at the centre (important)
- LOS ~ PM,t ~ PM,R (isotropy)



#### Mass model parameter results



#### Stellar-mass black holes over an IMBH

- Fit shows strong preference for two-component model with extended central mass
- IMBHs greater than  $6,000 \text{ M}_{\odot}$  excluded at  $3\sigma$ , limiting their kinematic relevance
- Pulsar profile has intermediate concentration and also has limited kinematic relevance



### Probing models of pulsar formation in $\omega$ Cen

- Milli-second pulsars (MSPs) are believed progeny of X-ray binaries
- Encounter models predict MSP abundances scale with stellar encounter rates:

$$\Gamma \propto \int dr r^2 \rho_{\star}(r)^2 / \sigma(r)$$

• We extend this for densities to describe the intra-cluster MSP distribution, showing excellent agreement (inset)



#### Conclusions

- Introduced promising methodology combining stellar kinematics and pulsar timing, providing additional constraints on mass models
- Simultaneously considered the presence of IMBH and central cluster of remnants, favoring remnants and setting a stringent  $3\sigma$  upper limit of 6,000 M<sub> $\odot$ </sub> on IMBH
- Analyzed MSP distribution, exploring stellar encounter formation scenario. Extended previous validation analyses to derive a profile for the intra-cluster distribution, showing excellent agreement

#### **Comments & future directions**

- Our methodology shows promise with the advent of rapidly growing and upcoming observations (e.g. SKA)
- Will address the presence of a DM halo and its comparison to simulations in a future publication (in prep.)
- Following realistic DM halo modeling, J-factors and escape velocities for  $\omega$  Cen should be revisited (in prep.)
- Followup analysis with new MUSE + HST data (~ 600,000 PMs)



#### IMBH discovered in $\omega$ Cen: a comeback?

- $3\sigma$  lower bound of  $8{,}200~{\rm M}_{\odot}$  from high-velocity stars, in apparent tension with ours
- Sensitive to the assumed escape velocity of the cluster
- Remnants already increase this by ~ 10 %, while kinematics allow an extended DM halo to ~ double it!

#### NASA's Hubble Finds Strong Evidence for Intermediate-Mass Black Hole in Omega Centauri

# $\begin{array}{l} & \mbox{ Article } \\ \hline Fast-moving stars around an intermediate- \\ mass black hole in $\omega$ Centauri \\ \end{array}$

https://doi.org/10.1038/s41586-024-07511-z Received: 15 December 2023

Accepted: 2 May 2024

Published online: 10 July 2024

Maximilian Häberle<sup>1™</sup>, Nadine Neumayer<sup>1</sup>, Anil Seth<sup>2</sup>, Andrea Bellini<sup>3</sup>, Mattia Libralato<sup>4,5</sup>, Holger Baumgardt<sup>6</sup>, Matthew Whitaker<sup>2</sup>, Antoine Dumont<sup>1</sup>, Mayte Alfaro-Cuello<sup>7</sup>, Jay Anderson<sup>3</sup>, Callie Clontz<sup>1,2</sup>, Nikolay Kacharov<sup>8</sup>, Sebastian Kamann<sup>9</sup>, Anja Feldmeier-Krause<sup>1,10</sup>, Antonino Milone<sup>11</sup>, Maria Selina Nitschai<sup>1</sup>, Renuka Pechetti<sup>9</sup> & Glenn van de Ven<sup>10</sup> View All News Releases >

July 10, 2024 11:00AM (EDT) | Release ID: 2024-015



#### Combining stellar kinematics + pulsar timing

• Included all data self-consistently into the likelihood:

$$\ln \mathcal{L}_{tot}(\boldsymbol{\theta}) = \ln \mathcal{L}_{LOS} + \ln \mathcal{L}_{PM,t} + \ln \mathcal{L}_{PM,R} + \ln \mathcal{L}_{VSP1} + \ln \mathcal{L}_{VSP2} + \ln \mathcal{L}_{p, LOS} + \ln \mathcal{L}_{pos} + \ln \mathcal{L}_{pos, R}$$

• For stellar kinematics and LOS accelerations:

$$\ln \mathcal{L}(\boldsymbol{\theta}) = -\frac{1}{2} \sum_{y} \chi_{y}^{2}, \qquad \chi_{y}^{2} \equiv \sum_{i} \frac{\left[y_{i, \text{ obs}} - y_{i}(\boldsymbol{\theta})\right]^{2}}{\delta y_{i}^{2}}$$

• For positional likelihoods of pulsars:  $n(r) = f(\alpha, r_0, n_0) \left[ 1 + \left(\frac{r}{r_0}\right)^2 \right]^{(\alpha-1)/2}$ 

#### **Pulsar accelerations**

- Favor ~ 20 % more massive and extended central mass than stellar kinematics alone
- Extremal central pulsar accelerations are IMBH 'smoking gun' signatures
- These are not favored by our analysis (except for a small IMBH very close to the center)



#### Stellar kinematics observables (LOS + PMs)

• Can't measure  $\sigma_r(r)$  and  $\beta(r)$  directly, but can constrain them with projected quantities (line-of-sight + proper motions):

$$\sigma_{\text{LOS}}^2(R) = \frac{2}{\Sigma_{\star}(R)} \int_R^{\infty} \left(1 - \frac{R^2}{r^2} \beta(r)\right) \frac{\nu_{\star}(r) \sigma_r^2(r) r}{\sqrt{r^2 - R^2}} dr,$$

$$\sigma_{\mathrm{PM,\,t}}^2(R) = \frac{2}{\Sigma_{\star}(R)} \int_R^{\infty} \left(1 - \beta(r)\right) \frac{\nu_{\star}(r)\sigma_r^2(r)r}{\sqrt{r^2 - R^2}} dr.$$

$$\sigma_{\rm PM,\,R}^2(R) = \frac{2}{\Sigma_\star(R)} \int_R^\infty \left(1 - \beta(r) + \frac{R^2}{r^2} \beta(r)\right) \frac{\nu_\star(r) \sigma_r^2(r) r}{\sqrt{r^2 - R^2}} dr$$

#### Stellar kinematics observables (VSPs)

• Can also measure higher velocity moments (virial shape paremeters) as an additional constraint:

$$VSP1 = \frac{2}{5} \int_0^\infty GMv_\star (5 - 2\beta) \sigma_r^2 r dr = \int_0^\infty \Sigma_\star \langle v_{LOS}^4 \rangle R dR$$

$$\text{VSP2} = \frac{4}{35} \int_0^\infty GM v_\star (7 - 6\beta) \sigma_r^2 r^3 dr = \int_0^\infty \Sigma_\star \langle v_{\text{LOS}}^4 \rangle R^3 dR$$

#### **Results: VSPs**



### The mass/density - anisotropy degeneracy

• Full consideration of these observables is important to avoid degeneracies with mass models

Monthly Notices of the ROYAL ASTRONOMICAL SOCIETY

MNRAS **471**, 4541–4558 (2017) Advance Access publication 2017 July 19



doi:10.1093/mnras/stx1798

### How to break the density-anisotropy degeneracy in spherical stellar systems

#### J. I. Read<sup>1</sup> $\star$ and P. Steger<sup>2</sup>

<sup>1</sup>Department of Physics, University of Surrey, Guildford GU2 7XH, UK <sup>2</sup>Institute for Astronomy, Department of Physics, ETH Zürich, Wolfgang-Pauli-Strasse 27, CH-8093 Zürich, Switzerland

Accepted 2017 July 14. Received 2017 June 23; in original form 2016 December 19

#### Dependence with enclosed encounter rates

• We define the enclosed encounter rate:

$$\Gamma(R) \propto \int_{-\infty}^{\infty} dl \int_{0}^{R} dR' R' \rho_{\star}(r)^{2} / \sigma(r),$$

- MSPs show a clear linear scaling, the other components don't
- Similar to the relation for total encounter rates of GCs (gray dots)
- X-ray sources are traced by stellar profile, not the encounter rate



