

Five-flavour scheme predictions for $t\bar{t}b\bar{b}$ at next-to-leading order accuracy in MG5_aMC@NLO

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based on: [EPJC 84, 763 \(2024\)](#)

LHC Top WG meeting
11–13 November 2024

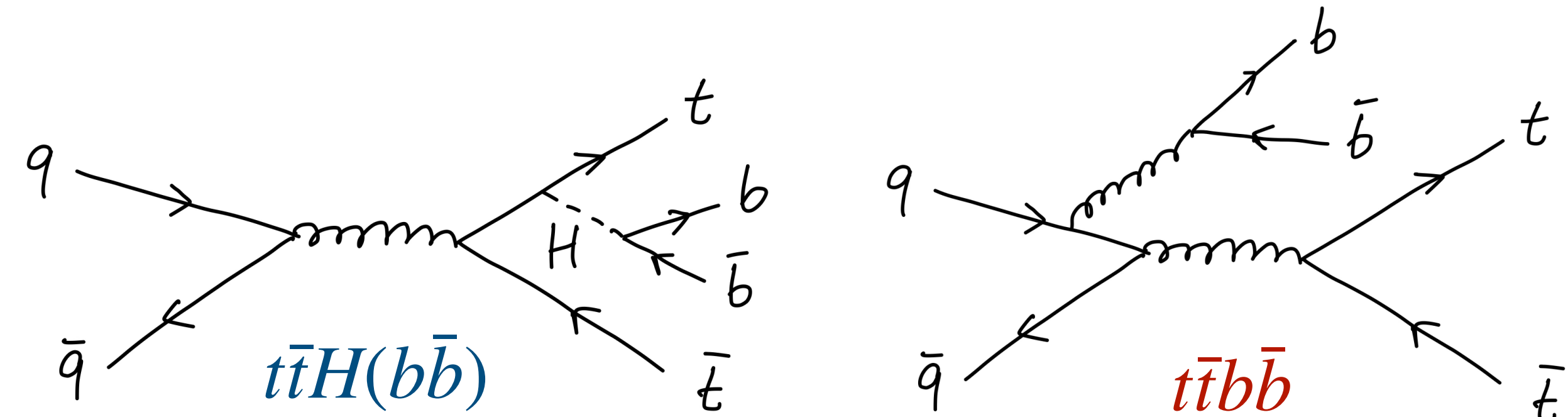
Role of the $t\bar{t}b\bar{b}$ process in the physics analyses

$t\bar{t}b\bar{b}$ represents a significant background in measurements probing the top Yukawa coupling

▶ $t\bar{t}H(\rightarrow b\bar{b})$ analyses

latest $t\bar{t}H(\rightarrow b\bar{b})$ from ATLAS [PLB 849 \(2024\)](#)
and CMS [arXiv:2407.10896](#)

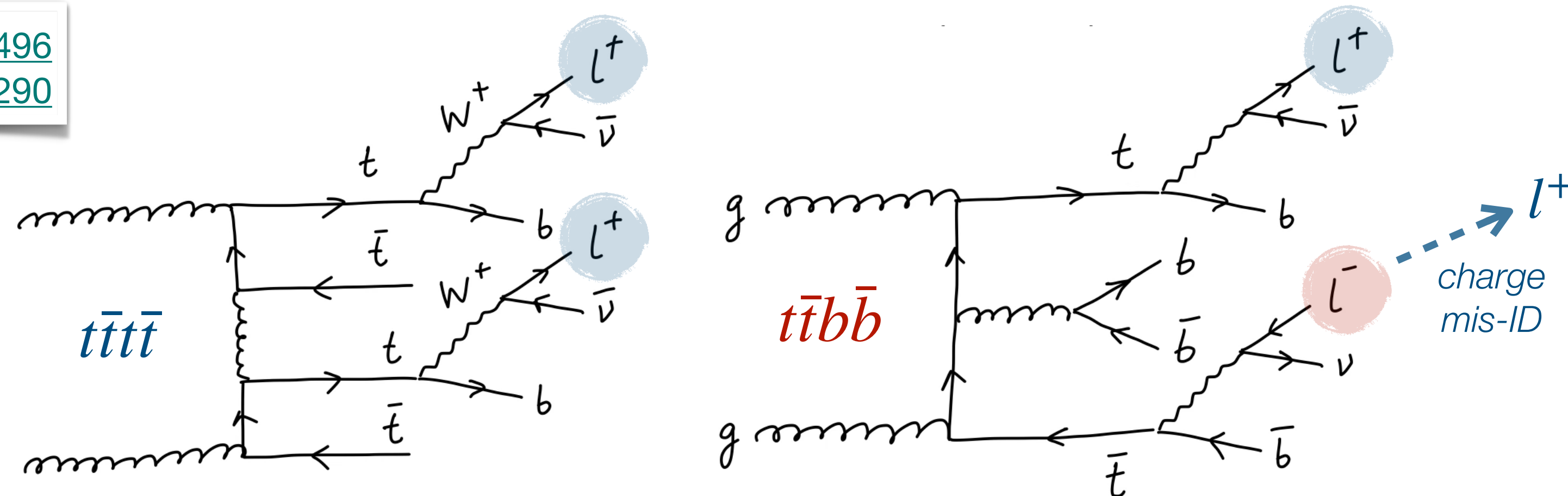
- $t\bar{t}b\bar{b}$ is the dominant background
- modelling uncertainty is currently a limitation



▶ 4-top analyses

latest $t\bar{t}t\bar{t}$ from ATLAS [EPJC 83 \(2023\) 6, 496](#)
and CMS [PLB 847 \(2023\) 138290](#)

- $t\bar{t}$ +jets (with additional b -jets) is the main source of fake/non-prompt and charge-mis-identification backgrounds



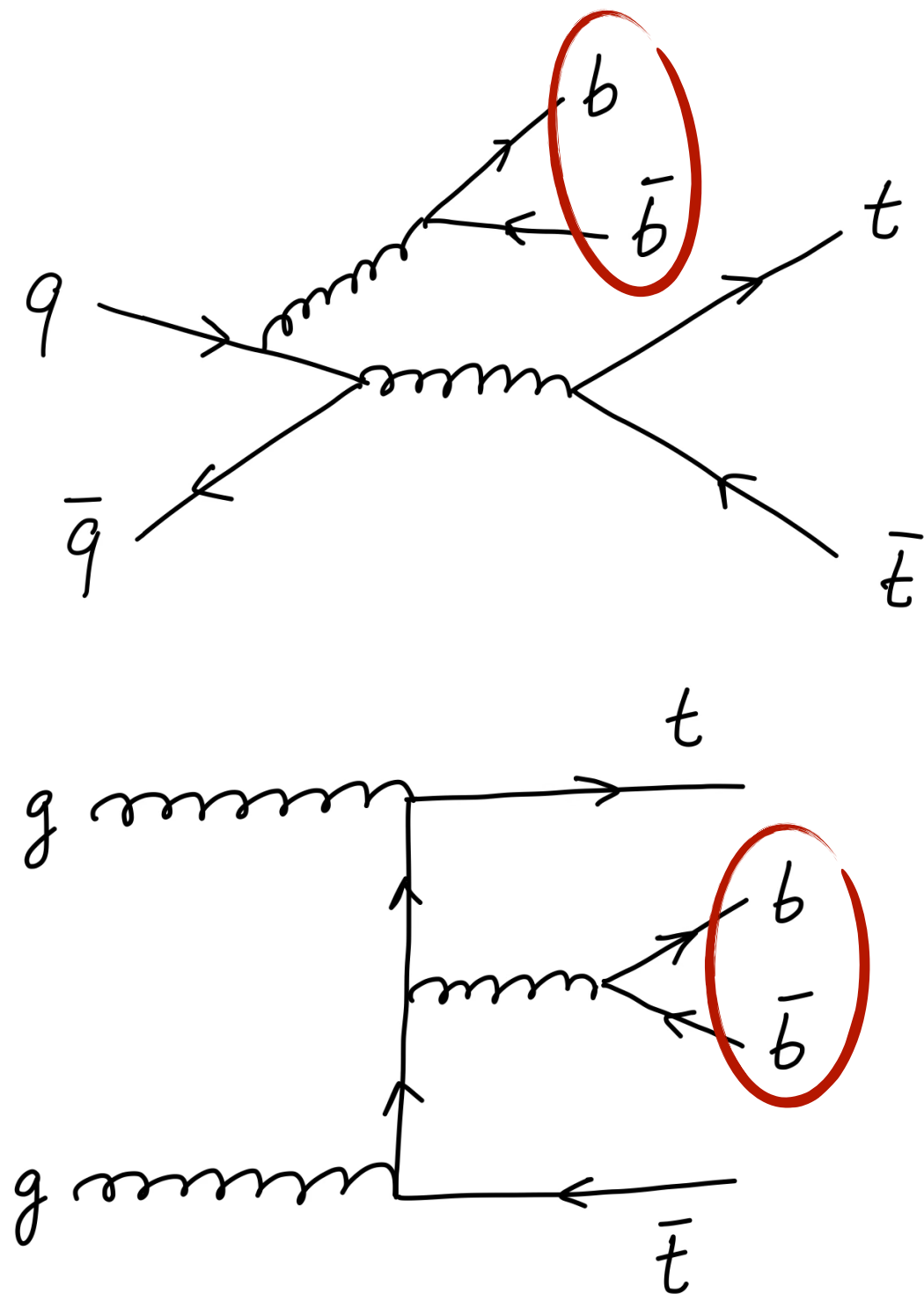
Simulation of the $t\bar{t}b\bar{b}$ process

▶ **Two primary theoretical frameworks:** four-flavour scheme (4FS) and five-flavour scheme (5FS)

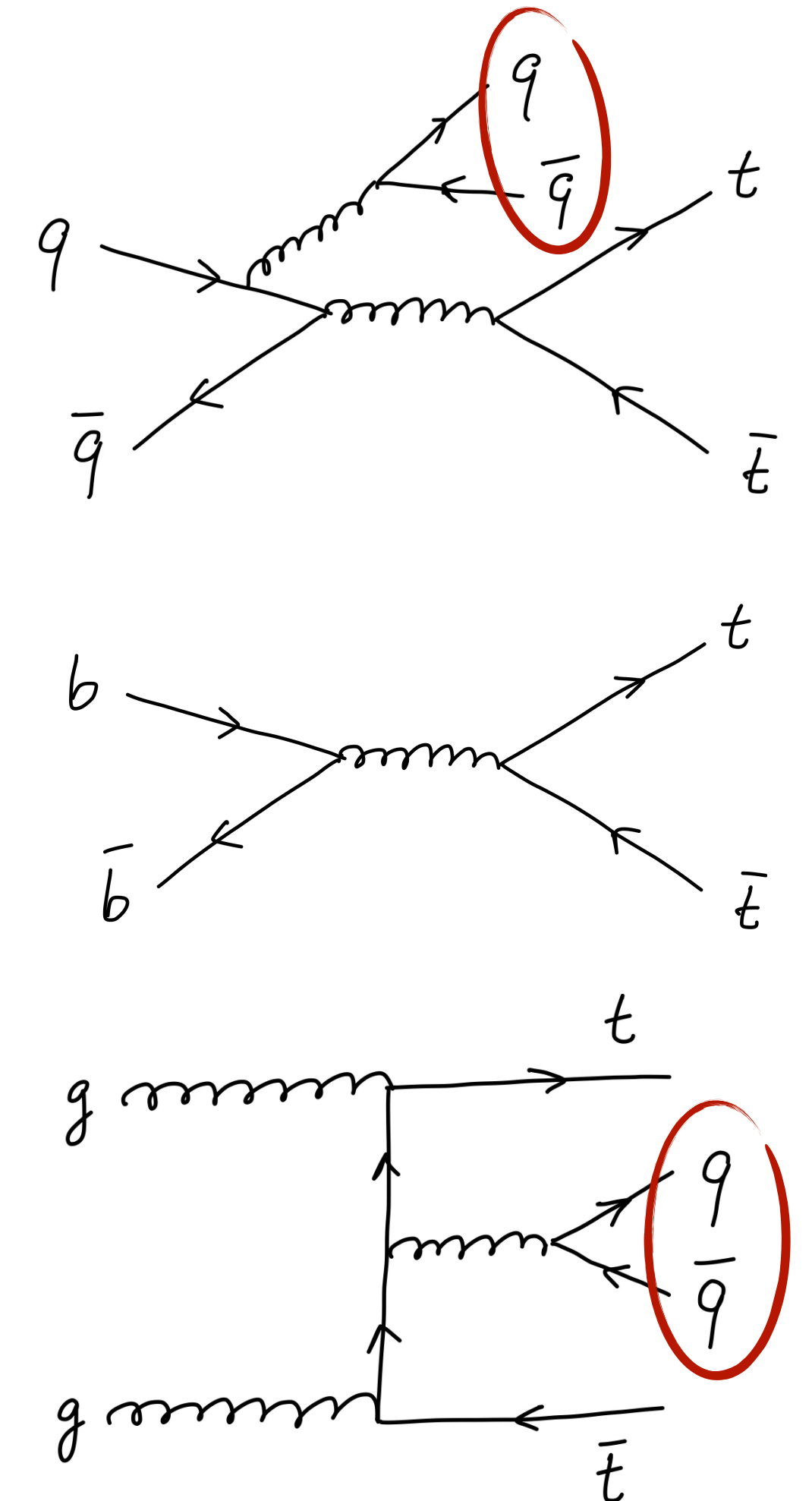
▶ **Alternative:** “fusion” method (or variable flavor number scheme)

[Höche, Krause, Siegert \(2019\)](#)
[Ferencz, Höche, Katzy, Siegert \(2024\)](#)

- Merges aspects of both the 4FS and 5FS calculations
- Currently, the additional jets in the 5FS component are only computed at LO



	4FS	5FS
b-quarks in the matrix element	massive	massless
b-quarks included in the PDF?	no	yes
renormalisation scheme	on-shell	$\overline{\text{MS}}$
final state	exclusively $t\bar{t}b\bar{b}$	inclusive $t\bar{t}$ + jets



Simulation of the $t\bar{t}b\bar{b}$ process in the 4FS

- ▶ 4FS calculations are usually the most precise at fixed order
 - b -quark mass effects taken into account
 - The process can be generated down to any energies

[Buccioni, Kallweit, Pozzorini, Zoller \(2019\)](#)

- ▶ Calculation with a certain number of jets at fixed order is reliable only if there are no scale hierarchies
 - $t\bar{t}b\bar{b}$ production is a multi-scale process

- Large mass difference between the top and bottom quarks → **large logarithms** $\left(\log(m_b/p_{T,b}) \text{ or } \log(p_{T,b}/\sqrt{\hat{s}})\right)$
- Difficult to choose optimal renormalisation and factorisation scales

see the discussion in the LHC Higgs Xsec WG report [arXiv:1610.07922](#)

- ▶ Challenges arise when matching to a parton shower:

- Parton shower radiation can produce additional b -quarks
- Jets generated by the shower can be harder than the matrix-element-level bottom quarks
- We need only the subleading b -quarks to come from the parton shower, but not the leading ones
 - Not fully understood how the parton shower radiation should be constrained

[Cascioli, Maierhöfer, Moretti, Pozzorini, Siegert \(2014\)](#)
[Ježo, Lindert, Moretti, Pozzorini \(2018\)](#)

Simulation of the $t\bar{t}b\bar{b}$ process in the 5FS

- ▶ Generate an inclusive $t\bar{t}$ + jets sample, select b -jets after parton showering

[Frixione, Nason, Webber \(2003\)](#)

[Frixione, Nason, Ridolfi \(2007\)](#)

[Hoeche, Krauss, Maierhoefer, Pozzorini, Schonherr, Siegert \(2015\)](#)

[Mazzitelli, Monni, Nason, Re, Wieseemann, Zanderighi \(2022\)](#)

- ▶ Massless b -quarks → **large logarithms do not arise in the matrix element**

- ▶ Large scale hierarchies between the top quarks and the jets can be resummed by a multi-jet merging procedure

- For example, FxFx merging in MadGraph5_aMC@NLO

[Frederix, Frixione \(2012\)](#)

- ▶ **Accurate parton-shower approximation for all softer jets**

- hardest parton shower jets are always softer than the softest matrix element jets, which is not always the case in the 4FS

$$p_T(\text{PS jets}) < \mu_Q < p_T(\text{ME jets})$$

merging scale (arrow pointing to μ_Q)

except for jets coming from the higher-multiplicity sample (arrow pointing to the inequality between $p_T(\text{PS jets})$ and μ_Q)

- ▶ b -quark mass effects:

- Important in the collinear/IR region ← incorporated into parton shower splitting functions
- Missing in the matrix element, but they are less relevant for the hard b -quarks

Simulation of the $t\bar{t}b\bar{b}$ process in the 5FS

- ▶ But generating $t\bar{t} + 0,1,2$ jets @ NLO accuracy requires substantial computing resources

[Hoeche, Krauss, Schonherr, Siegert \(2013\)](#)

[Frederix, Frixione \(2012\)](#)

[Plätzer \(2013\)](#)

- ▶ Selection efficiency of $t\bar{t}b\bar{b}$ is low (percent level)
 - $gg \rightarrow t\bar{t}gg$ dominates

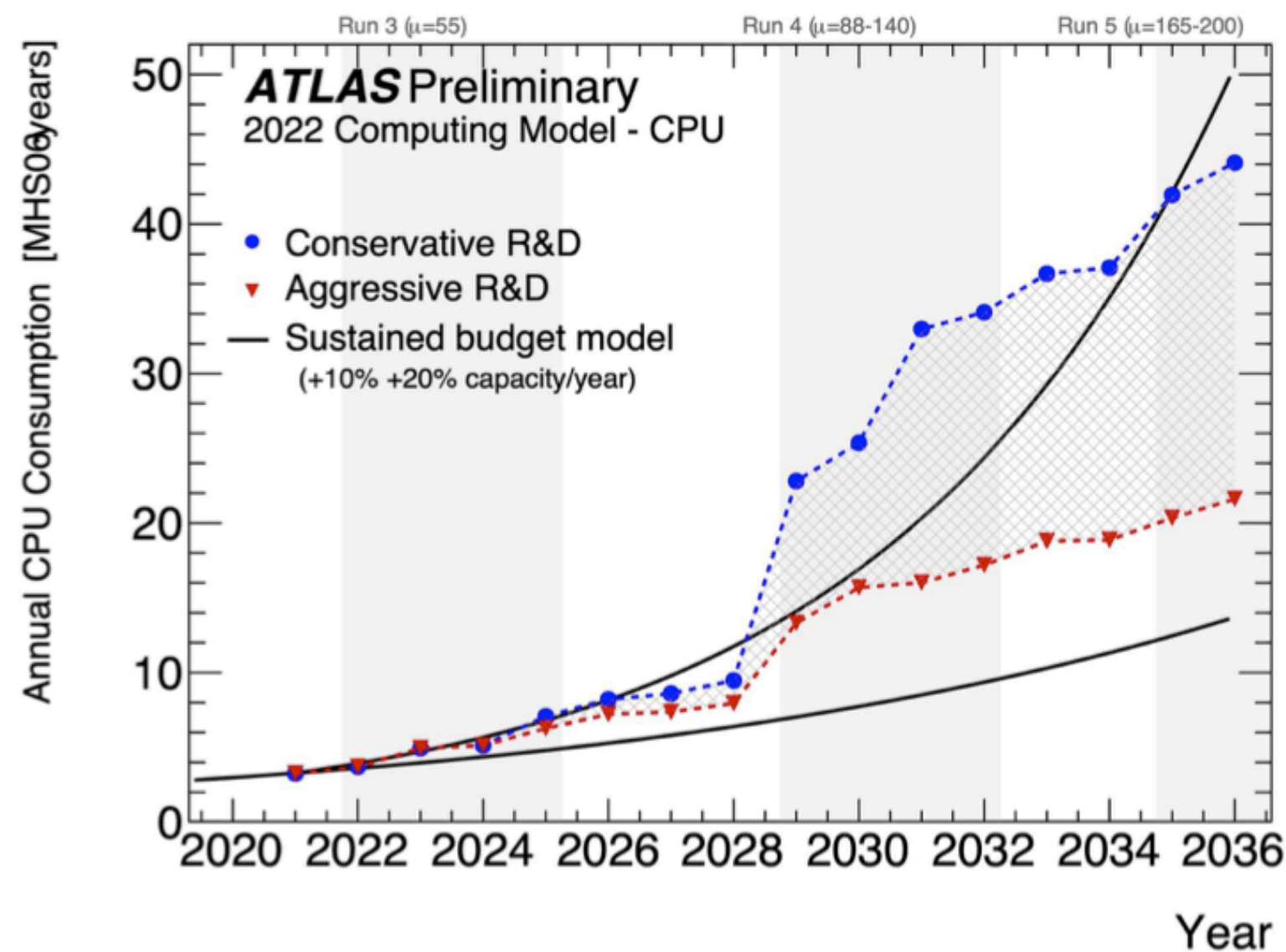
number of instructions to calculate a process in MadGraph5_aMC@NLO

	$gg \rightarrow t\bar{t}$	$gg \rightarrow t\bar{t}gg$	$gg \rightarrow t\bar{t}ggg$
madevent	13G	470G	11T
matrix1	3.1G (23%)	450G (96%)	11T (>99%)
└ ext	450M (3.4%)	3.3G (<1%)	7.3G (<1%)
└ int	1.9G (14%)	160G (35%)	2T (19%)
└ amp	530M (4.0%)	210G (44%)	5.5T (51%)

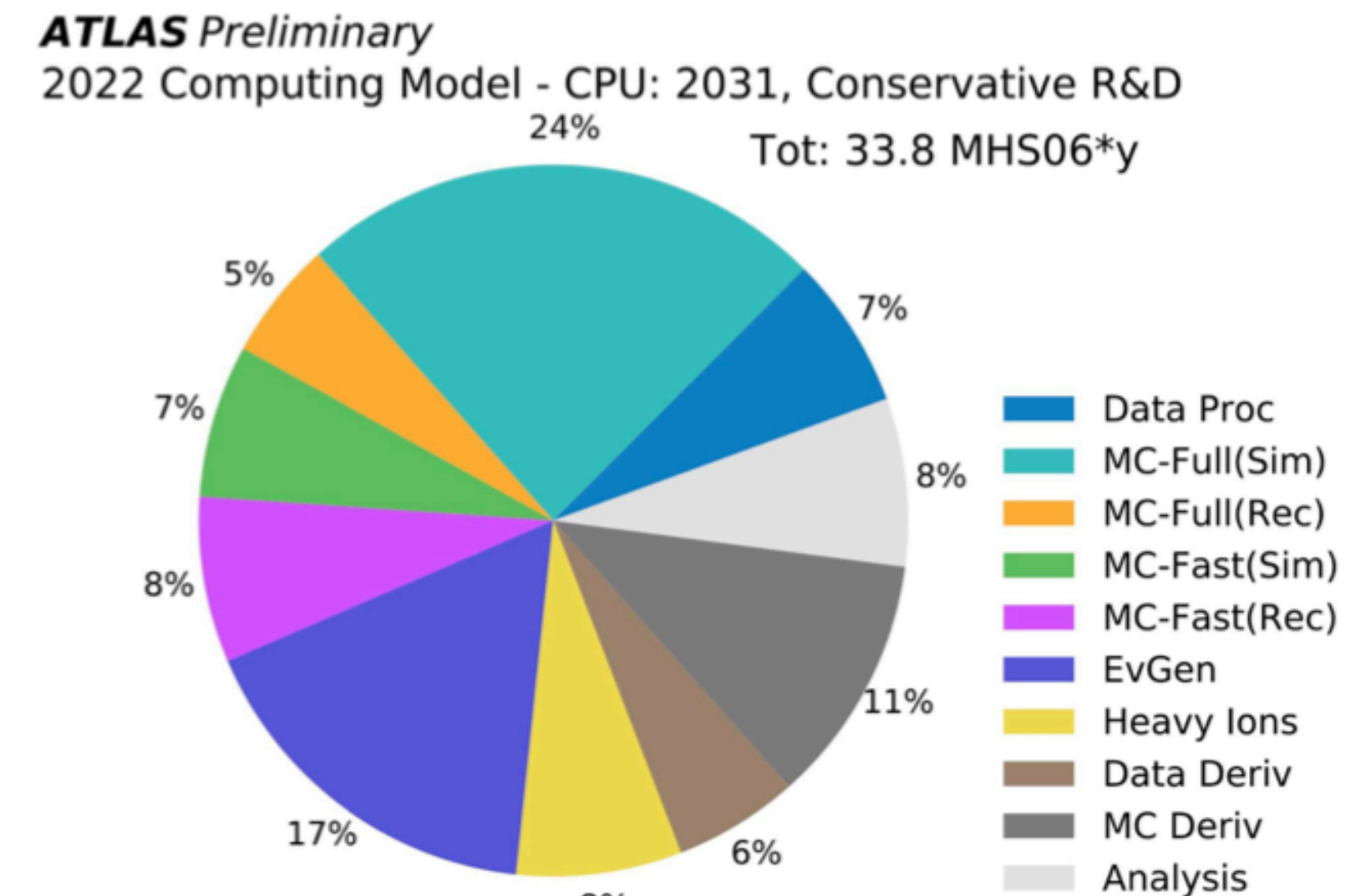
from [O. Mattelaer's talk](#)

5FS approach is computationally demanding!

- This will become even more relevant when producing Monte-Carlo for the HL-LHC era



[CERN-LHCC-2022-005](#)

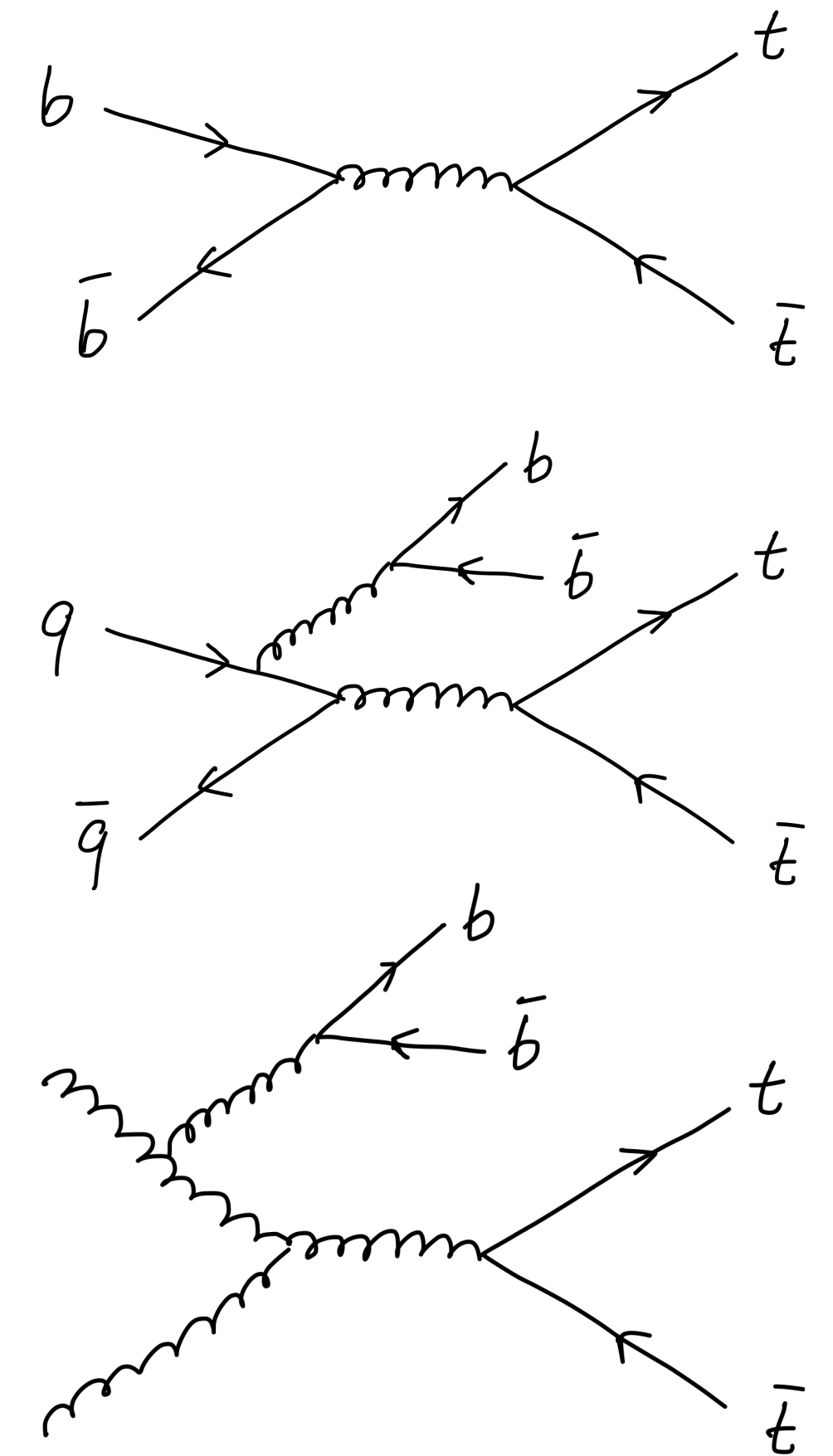


b -flavour enhancement in the matrix element

✓ We proposed a novel method to enhance the b -jet selection efficiency in the 5FS approach

R. Frederix, TM
[EPJC 84, 763 \(2024\)](#)

- ▶ Augment the generation probability of bottom quark flavour in the short-distance event generation
 - During phase-space integration and unweighting, multiply the weight of each contribution containing external b -quarks by w_{enh}
 - For bottom quarks can be generated in the initial or final state
 - $gg \rightarrow t\bar{t}b\bar{b}(g)$
 - $gb \rightarrow t\bar{t}bg(\rightarrow b\bar{b})$
 - $bb \rightarrow t\bar{t}q\bar{q}(g)$
 - ...
- ▶ To compensate for this and to preserve the cross-section, multiply the weight of events with external b -quarks by $1/w_{\text{enh}}$



examples of the enhanced subprocesses

b -flavour enhancement in the matrix element

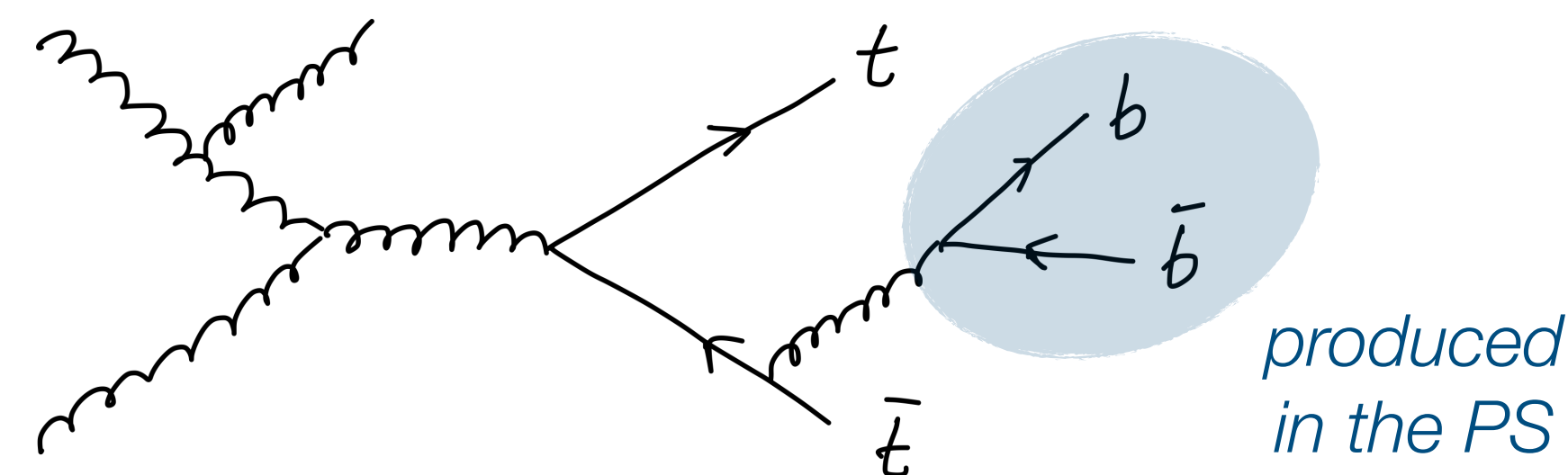
✓ We proposed a novel method to enhance the b -jet selection efficiency in the 5FS approach

R. Frederix, TM
[EPJC 84, 763 \(2024\)](#)

- ▶ This procedure is implemented in the MadGraph5_aMC@NLO
 - enhancement factor w_{enh} can be set by a new parameter, `bfLav_enhancement`, in the runcard file
 - **The new feature will become part of an upcoming release**

* **NB:** hard processes like $gg \rightarrow t\bar{t}gg$ which can yield a $t\bar{t}b\bar{b}$ event after a $g \rightarrow b\bar{b}$ splitting in the parton shower will not get enhanced \Rightarrow the fraction of $t\bar{t}b\bar{b}$ events is increased by a factor smaller than w_{enh}

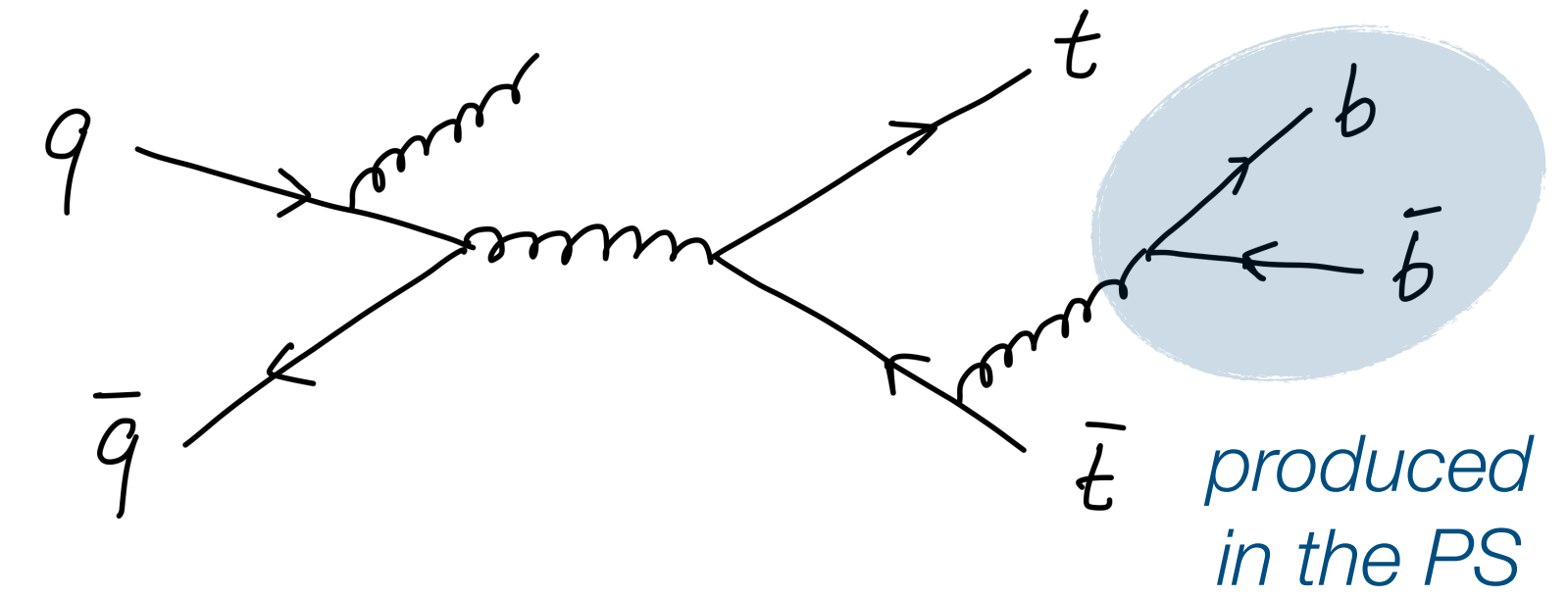
- also, too high enhancement factors (>100) cause instabilities which result in large statistical fluctuations



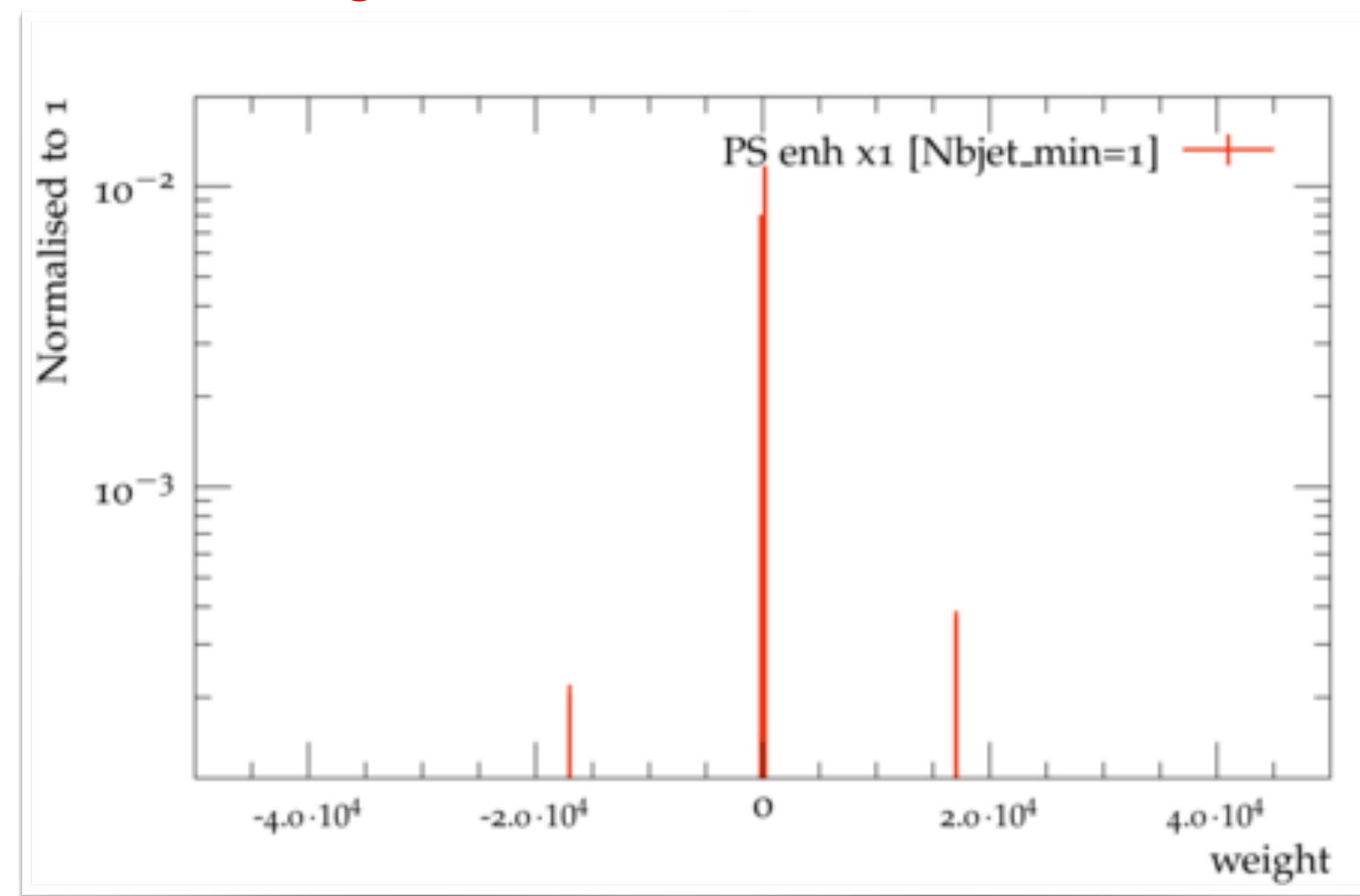
*this diagram is not enhanced
— can we enhance it in the PS?*

b -flavour enhancement in the parton shower?

- ▶ A similar biasing strategy can be potentially applied in the parton shower
- ▶ Pythia8 has a built-in mechanism for enhancing splitting probabilities, in particular $g \rightarrow b\bar{b}$ ones
 - In versions ≥ 8.311
- ▶ In practice:
 - Even moderate enhancement in the PS causes significant widening of the event weight distribution
 - **Large weights deteriorate the statistics** \rightarrow cancels the improvement from the b -enhancement completely

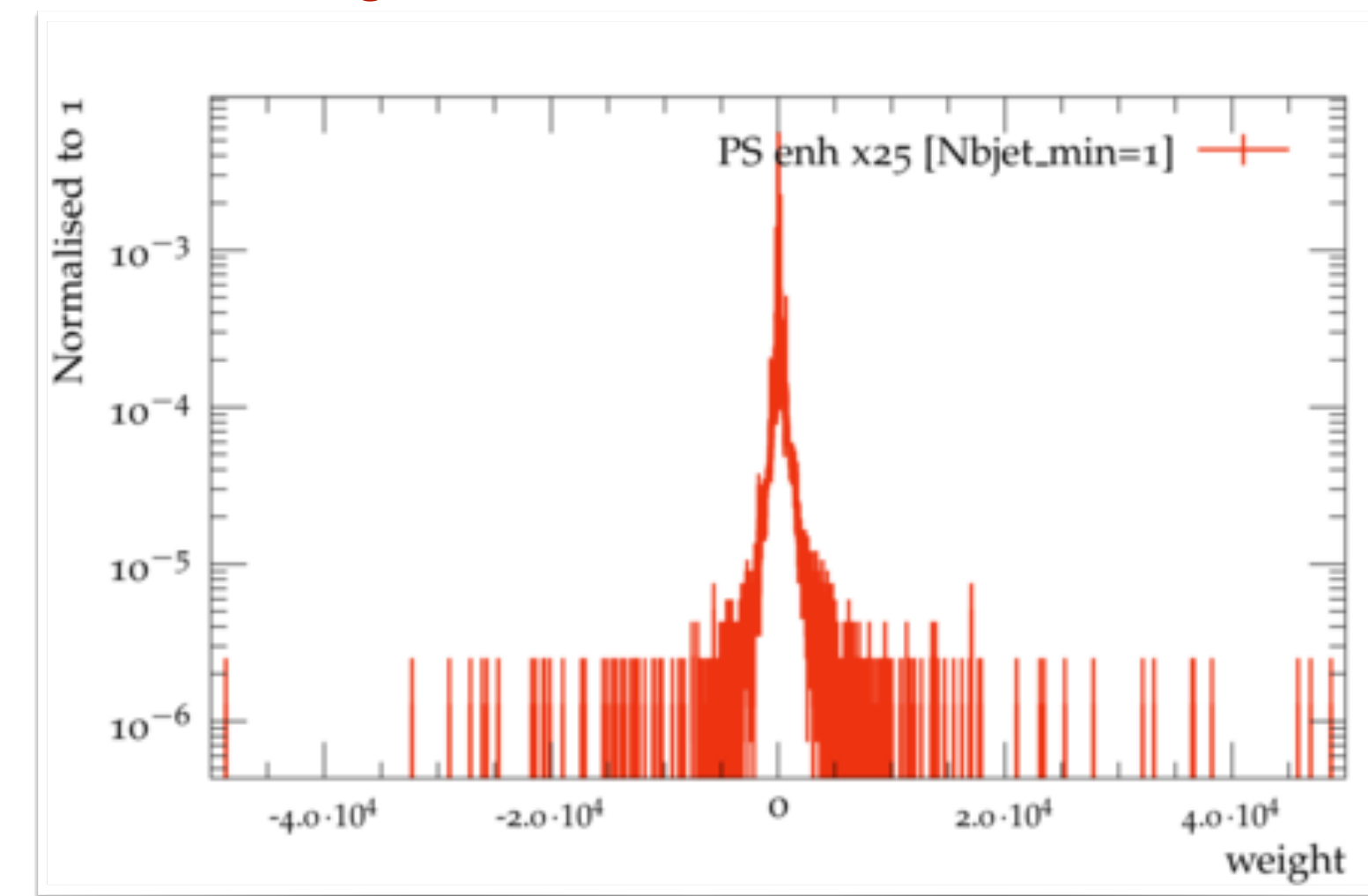


event weights w/o enhancement in the PS



$g \rightarrow b\bar{b}$
enhancement
 \rightarrow

event weights with enhancement in the PS



Generation setup for the 5FS sample and comparison to the 4FS

- ▶ **MadGraph5_aMC@NLO $t\bar{t}$ + 0,1,2 jets @NLO sample, FxFx merged**
- ▶ **Enhancement factor** $w_{\text{enh}} = 100$
- ▶ **Renormalisation/factorisation scales:** central values for are taken from the FxFx merging
 - 7-point variations
- ▶ **Merging scale:** 40 GeV
 - variations: 70 and 100 GeV
- ▶ **Shower starting scale:** $H_T/2$
 - variation: $H_T/4$
- ▶ Generation-level cut of 20 GeV on jet p_T
- ▶ **Matched to the Pythia8 parton shower**

*taking an envelope as
a total uncertainty*

- ▶ **Not including:**
 - hadronisation
 - underlying events
 - top quark decay

*to reduce the generation time
and to simplify the analysis,
and because we focus on the
differences in the ME*

Truth-level analysis in Rivet:

- ▶ anti- k_T jets ($R > 0.4$)
 - $p_T > 25$ GeV
 - $|\eta| < 2.5$
- ▶ jets containing at least one bottom quark are identified as b -jets
- ▶ consider two scenarios:
 - at least 1 b -jet
 - at least 2 b -jets

- ▶ **MadGraph5_aMC@NLO+Pythia8** NLO+PS $t\bar{t}b\bar{b}$ sample

- ▶ **Renormalisation/factorisation scales:**

- central values:

$$\mu_R = (E_{T,t}E_{T,\bar{t}}E_{T,b}E_{T,\bar{b}})^{1/4}$$

$$\mu_F = \frac{1}{2}(E_{T,t} + E_{T,\bar{t}} + E_{T,b} + E_{T,\bar{b}})$$

following the recommendations in
the LHC Higgs Xsec WG report
[arXiv:1610.07922](#)

- 7-point variations

- ▶ **Shower starting scale:** $H_T/2$

- ▶ Generation-level cut of 20 GeV on jet p_T

- ▶ **Matched to the Pythia8 parton shower**

- ▶ **Not including:**

- shower starting scale uncertainty
- matching scheme uncertainty
- hadronisation
- underlying events
- top quark decay

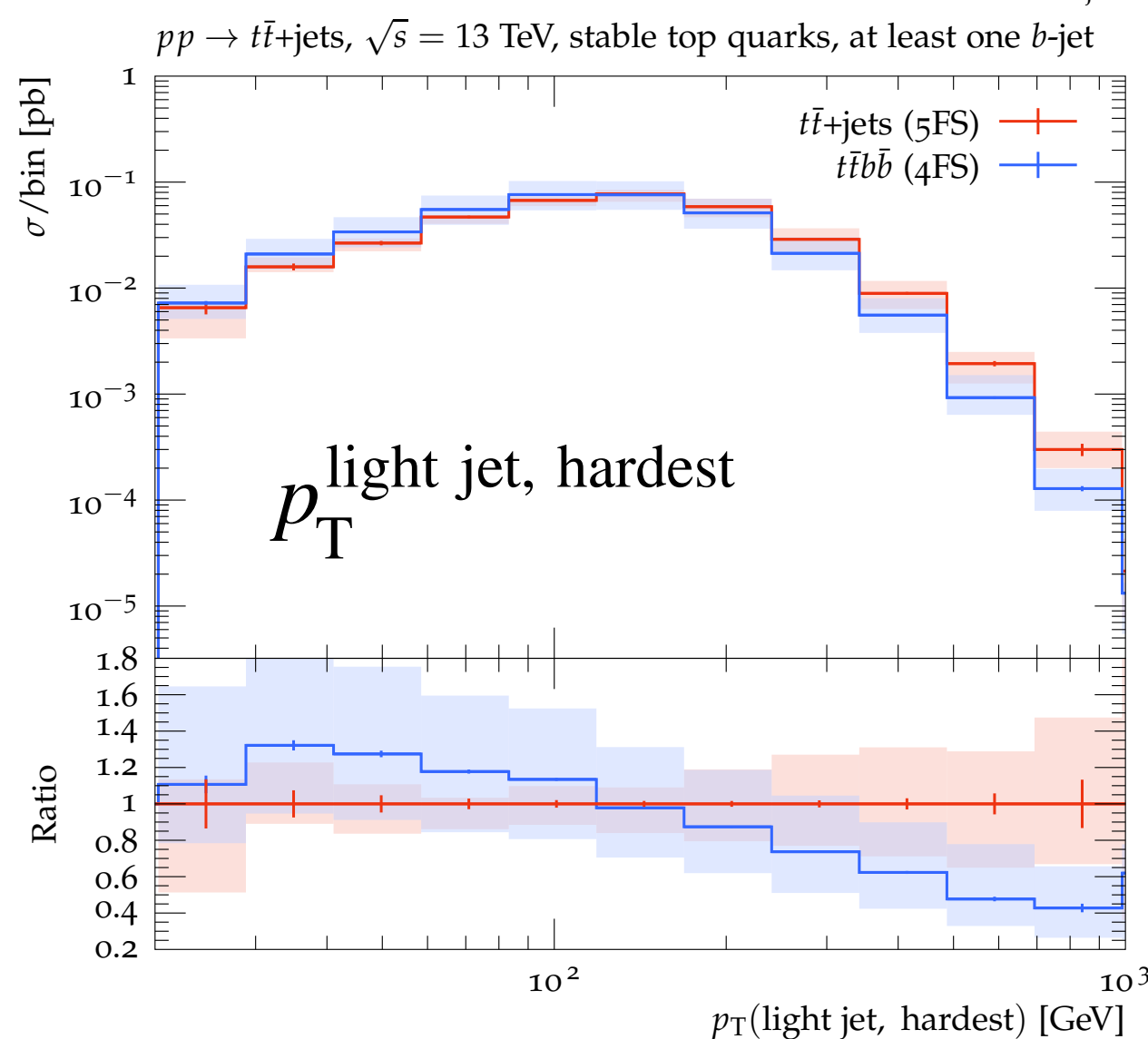
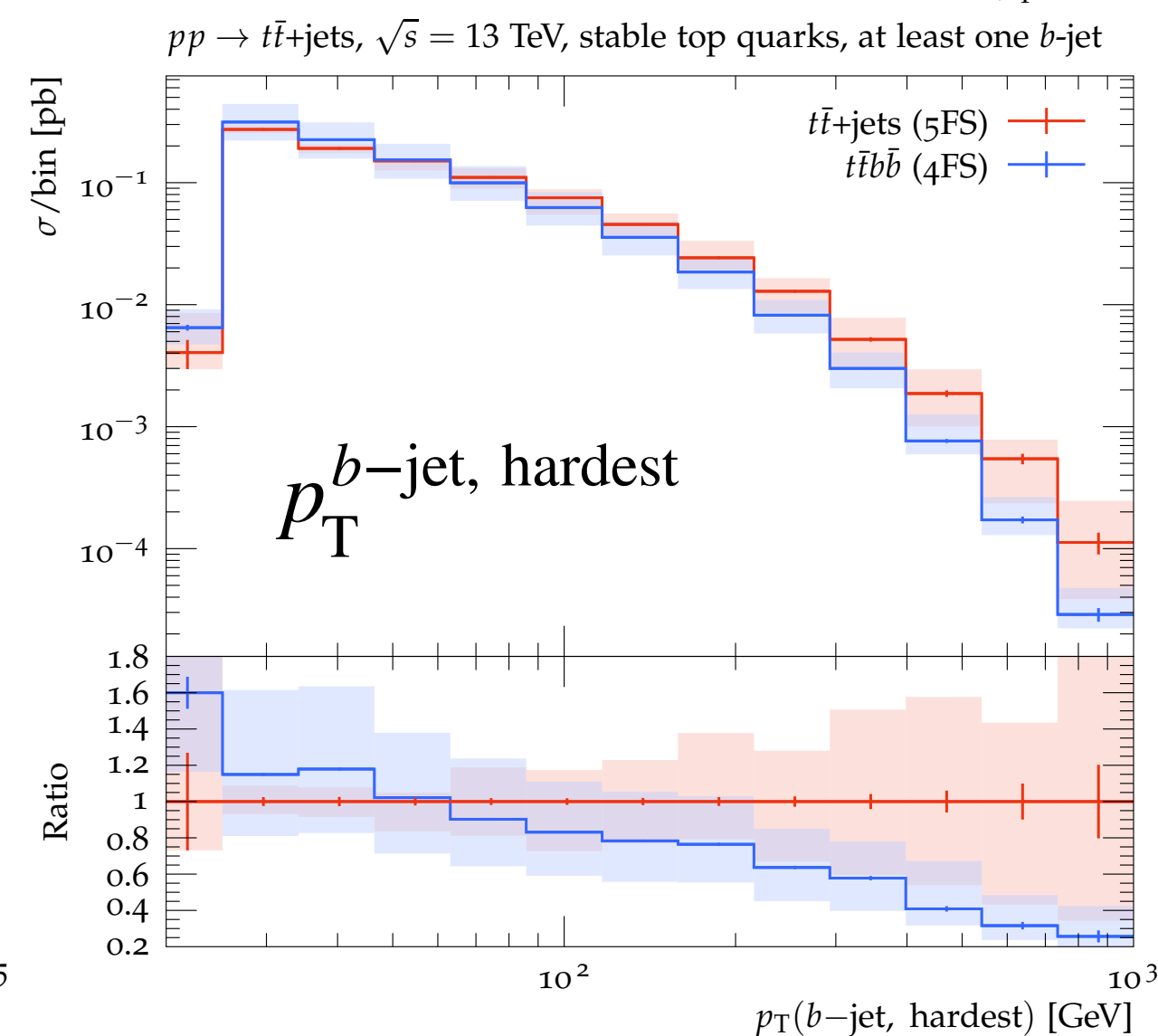
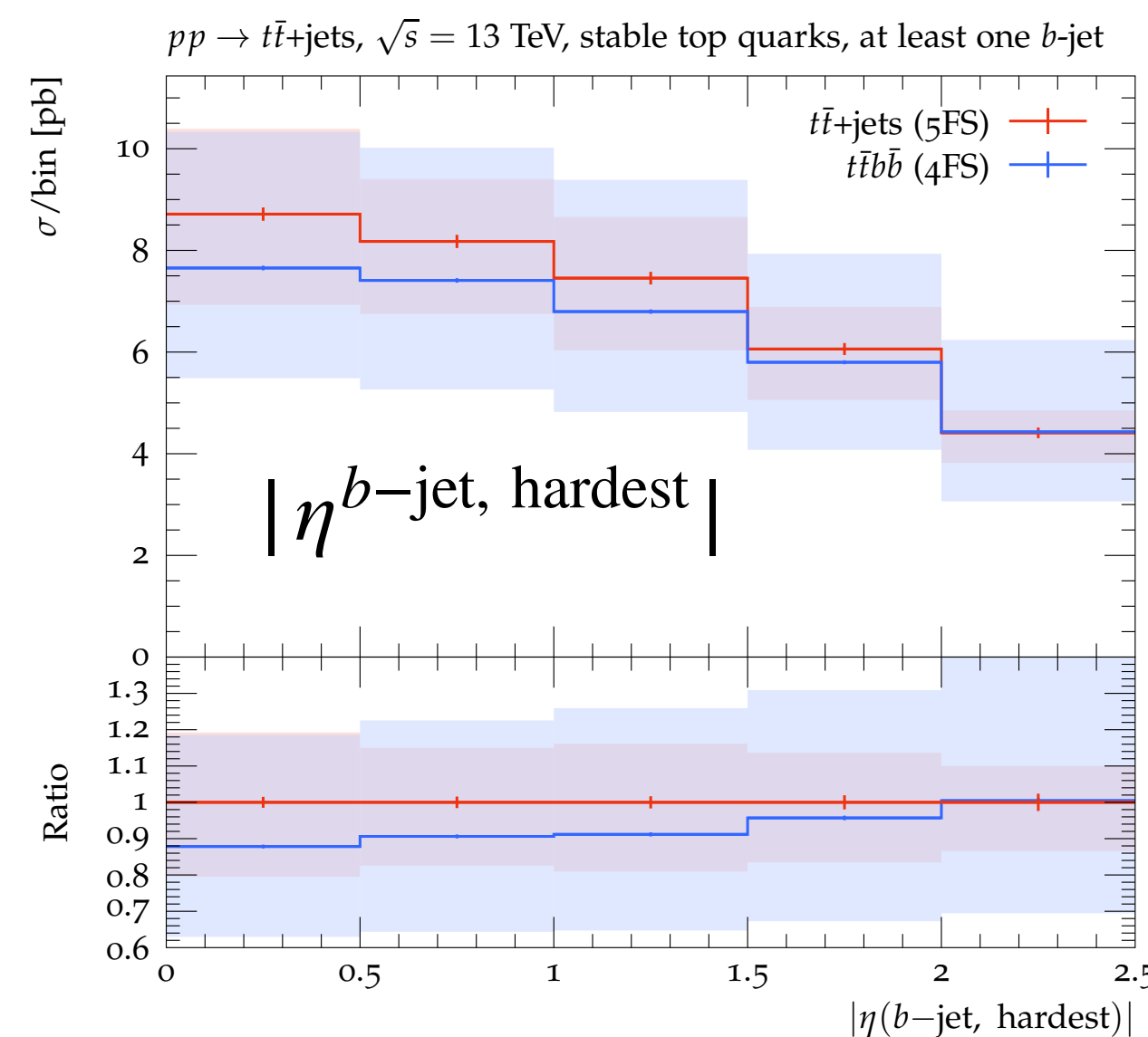
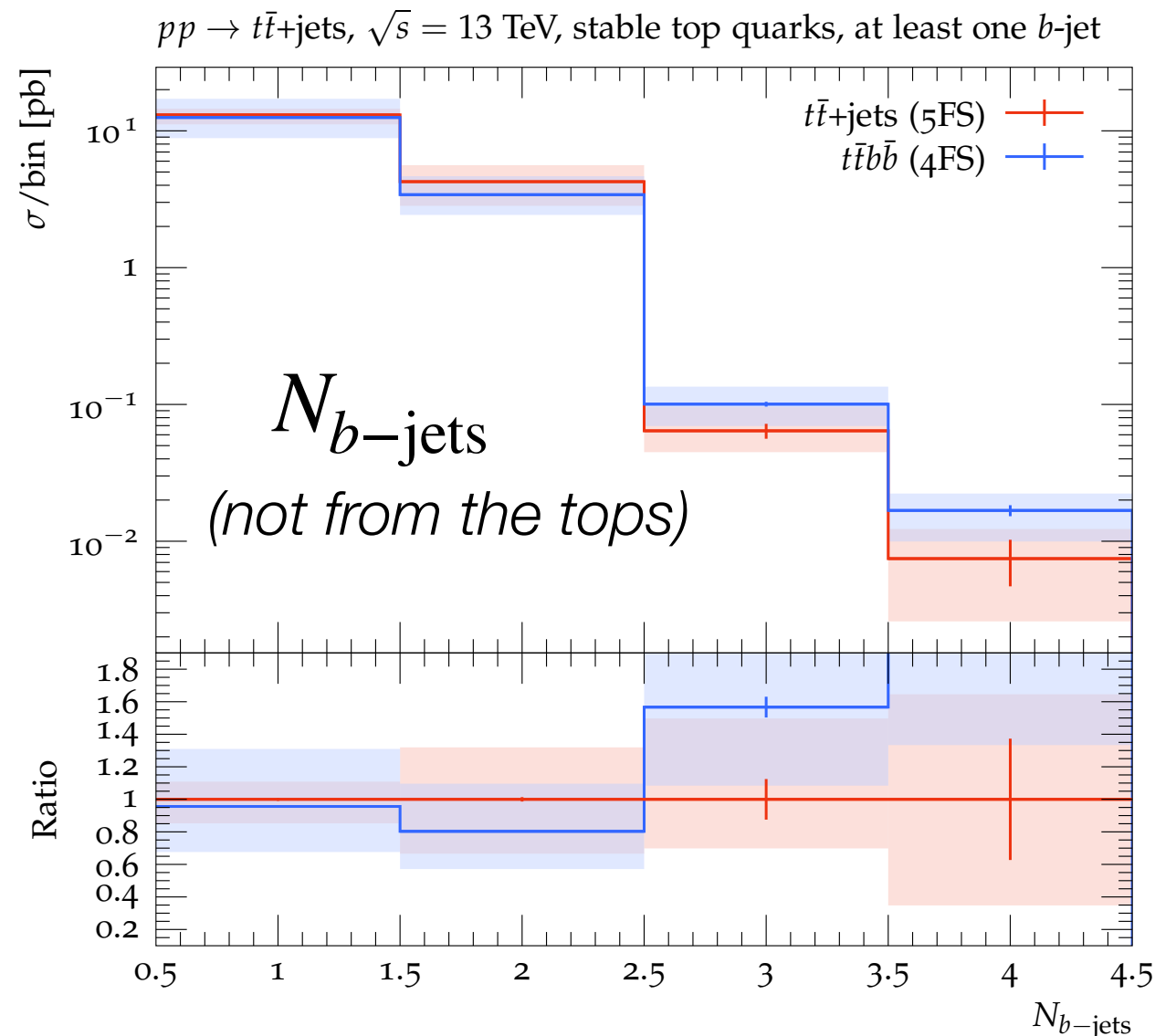
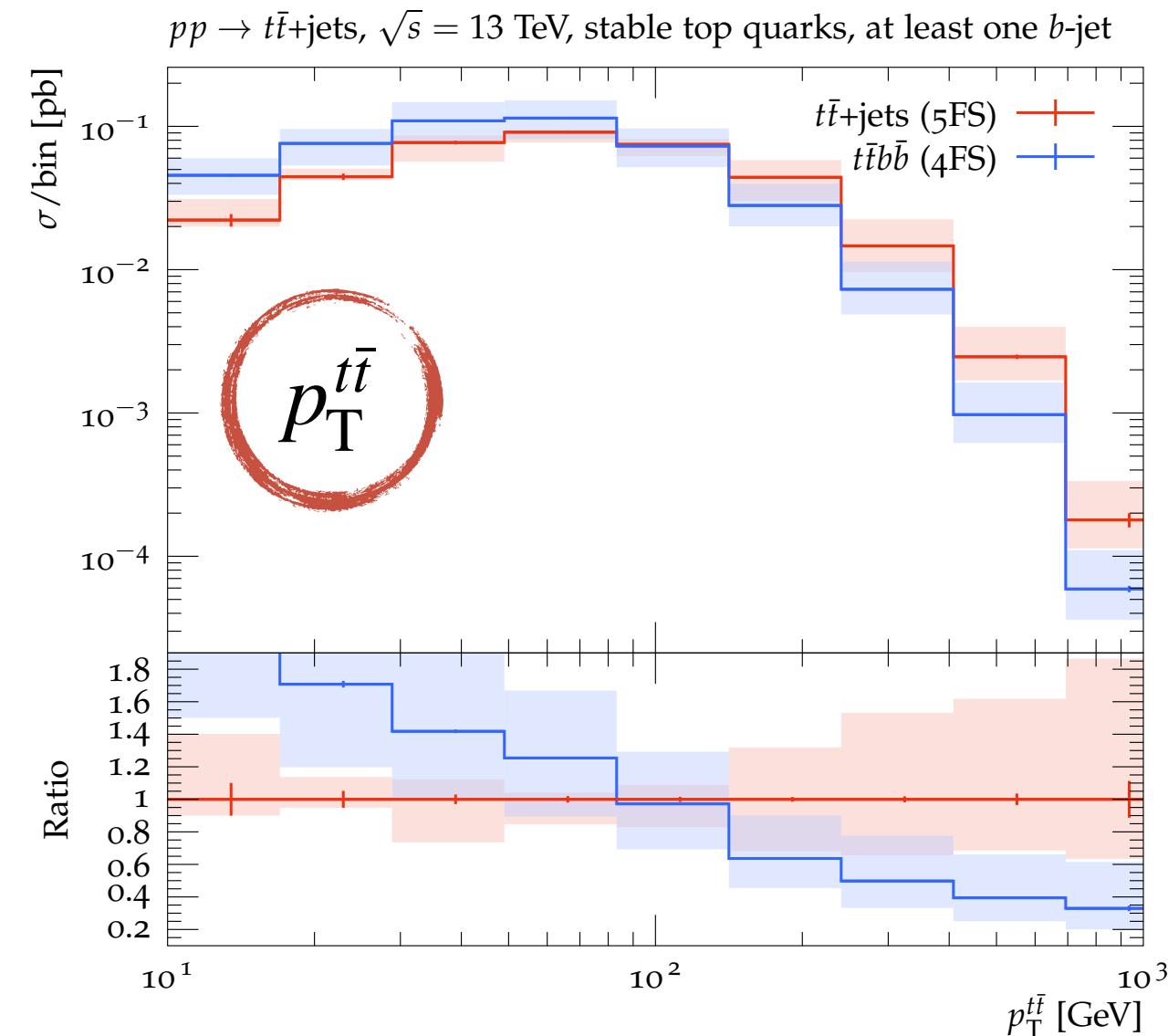
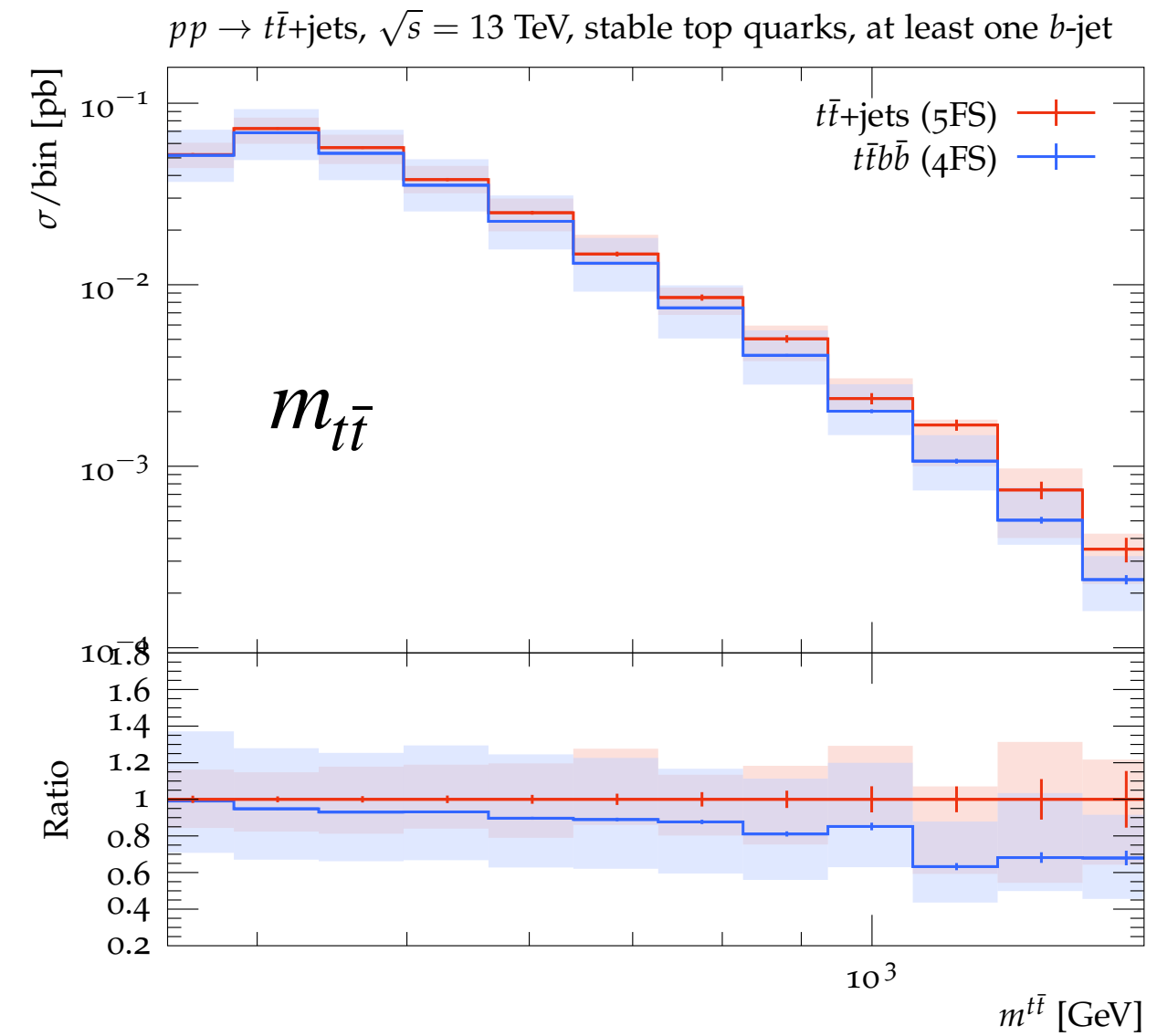
← *expected to be sizeable,
(see the LHC HXS WG report)
but is non-trivial to assess exactly*

Truth-level analysis in Rivet:

- ▶ anti- k_T jets ($R > 0.4$)
 - $p_T > 25$ GeV
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- ▶ jets containing at least one bottom quark are identified as b -jets
- ▶ consider two scenarios:
 - at least 1 b -jet
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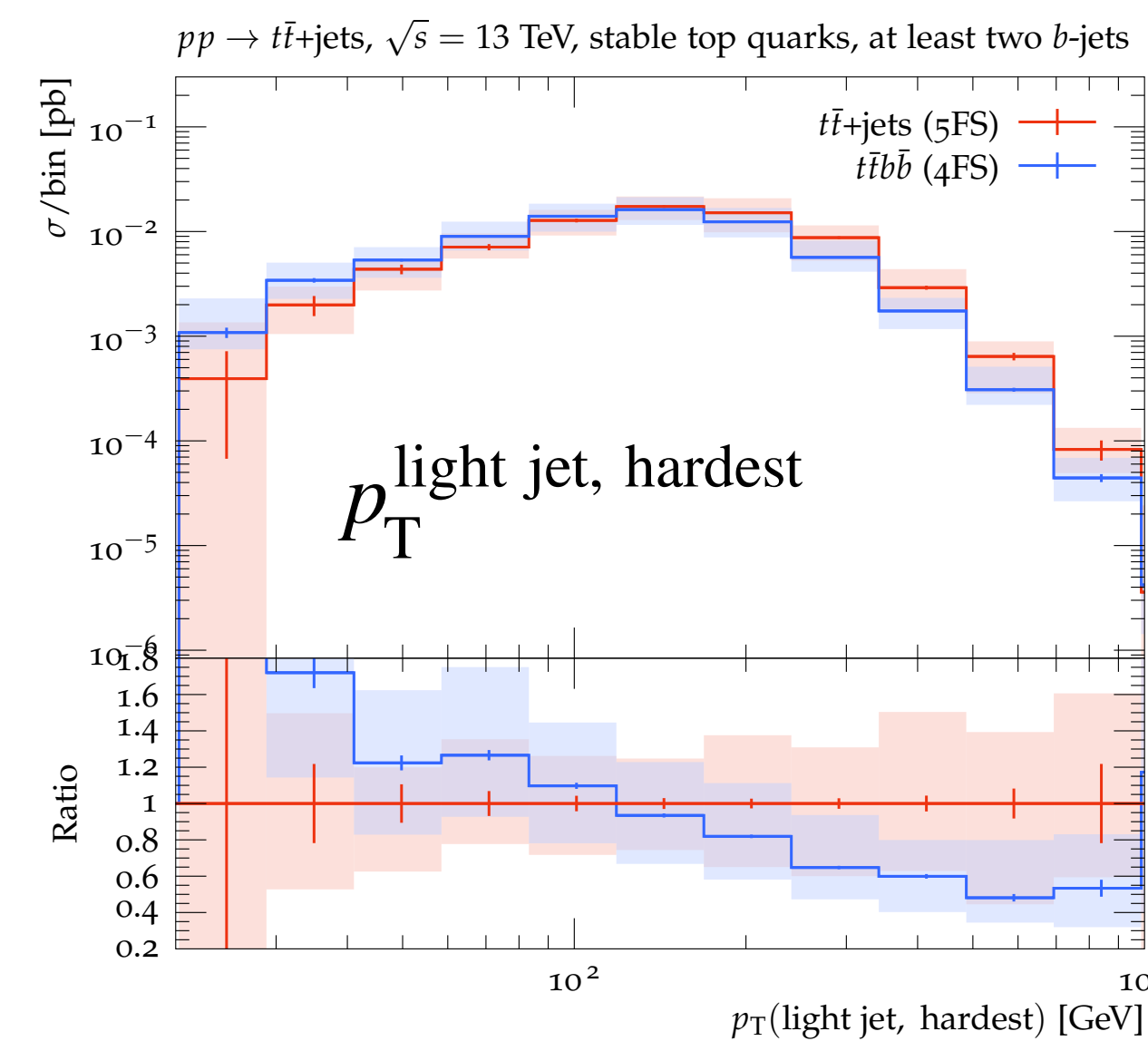
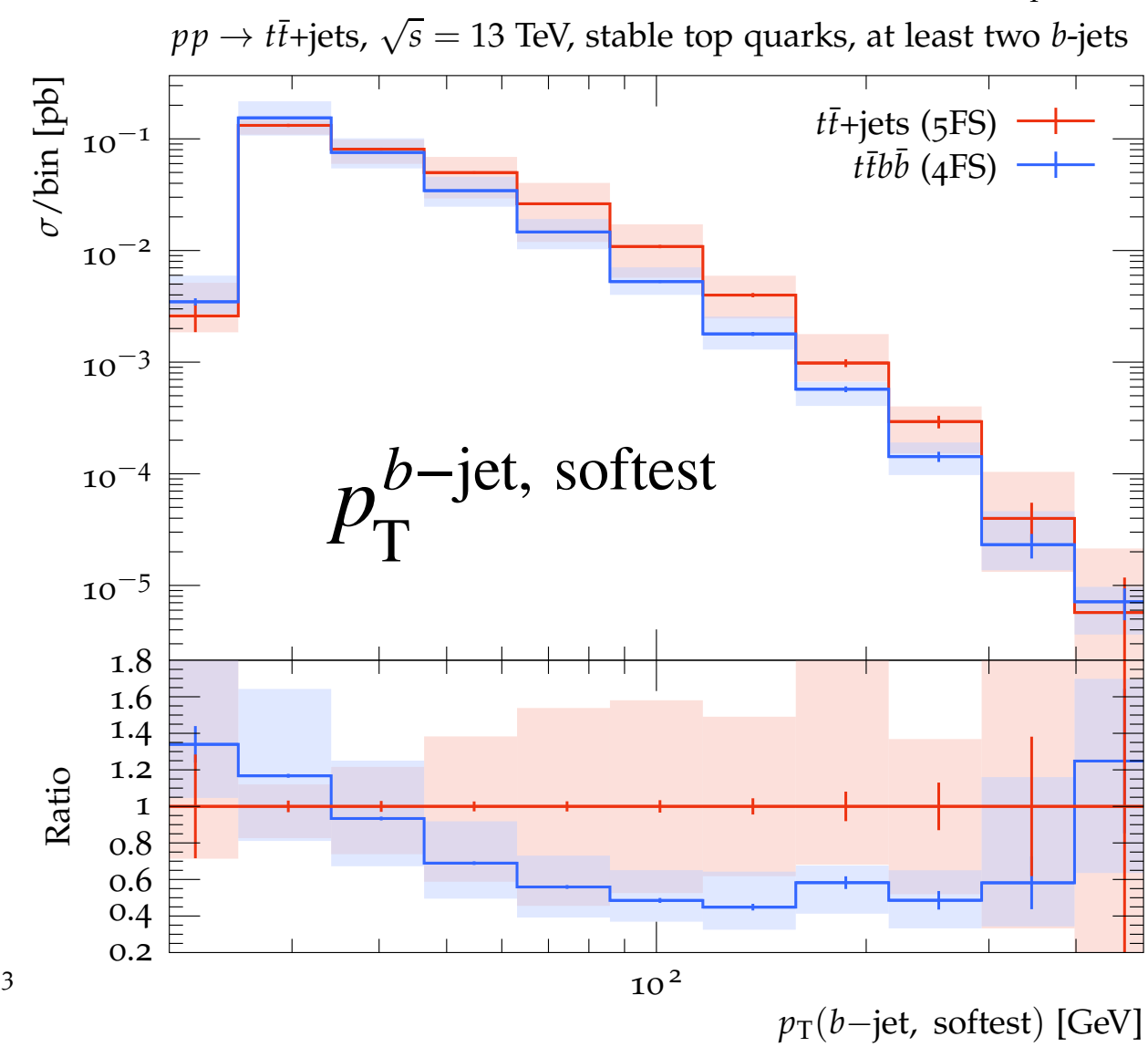
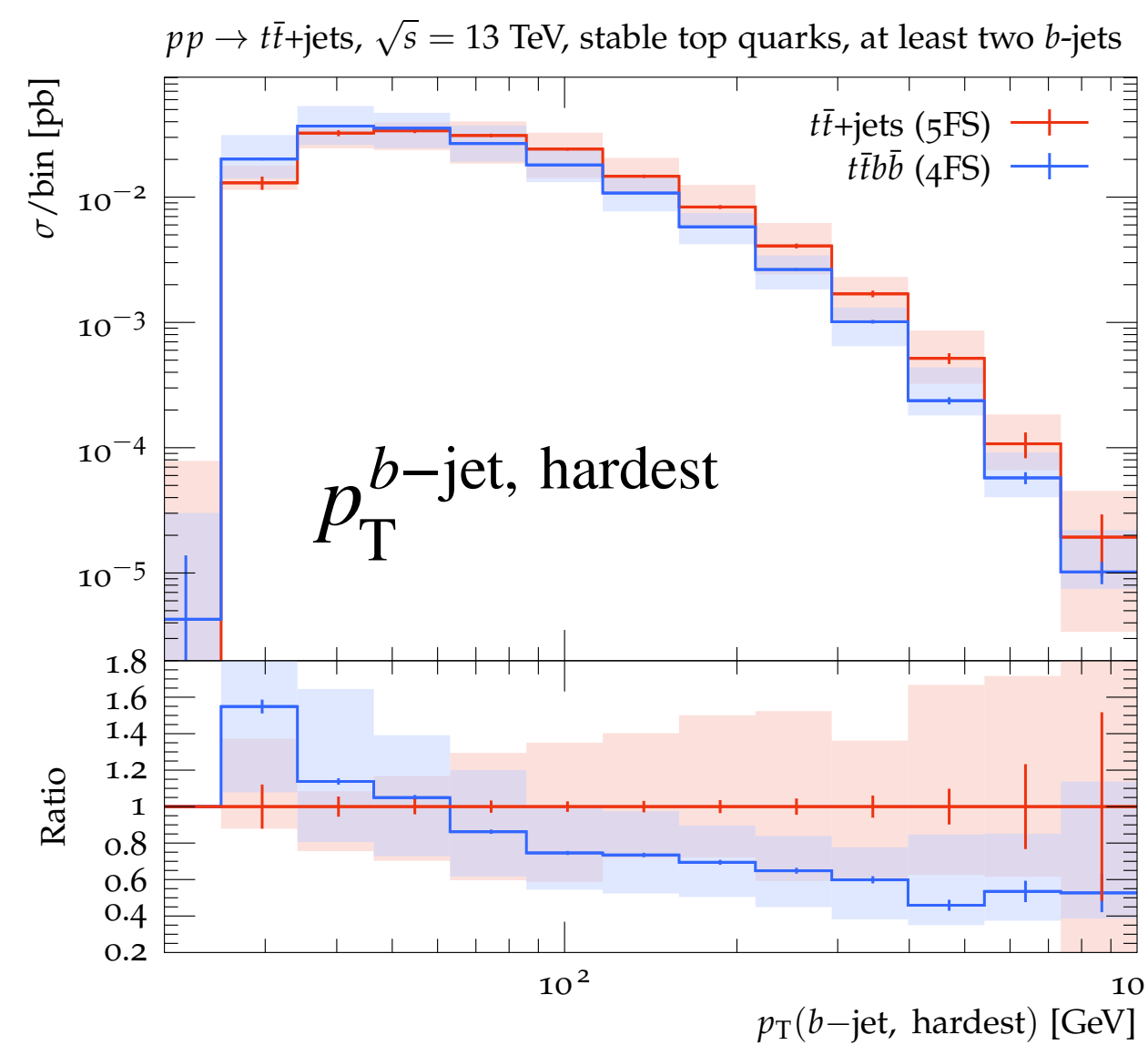
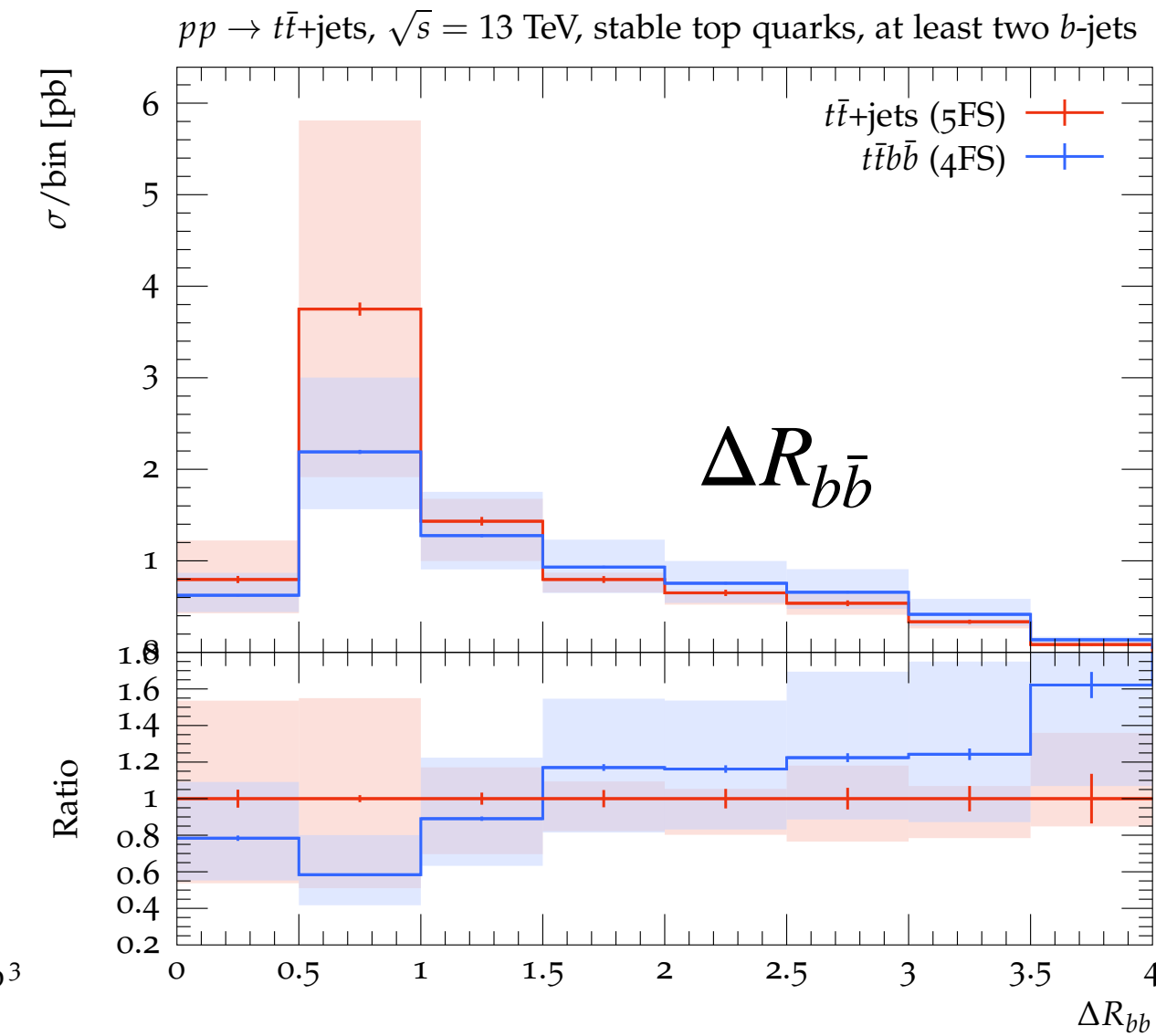
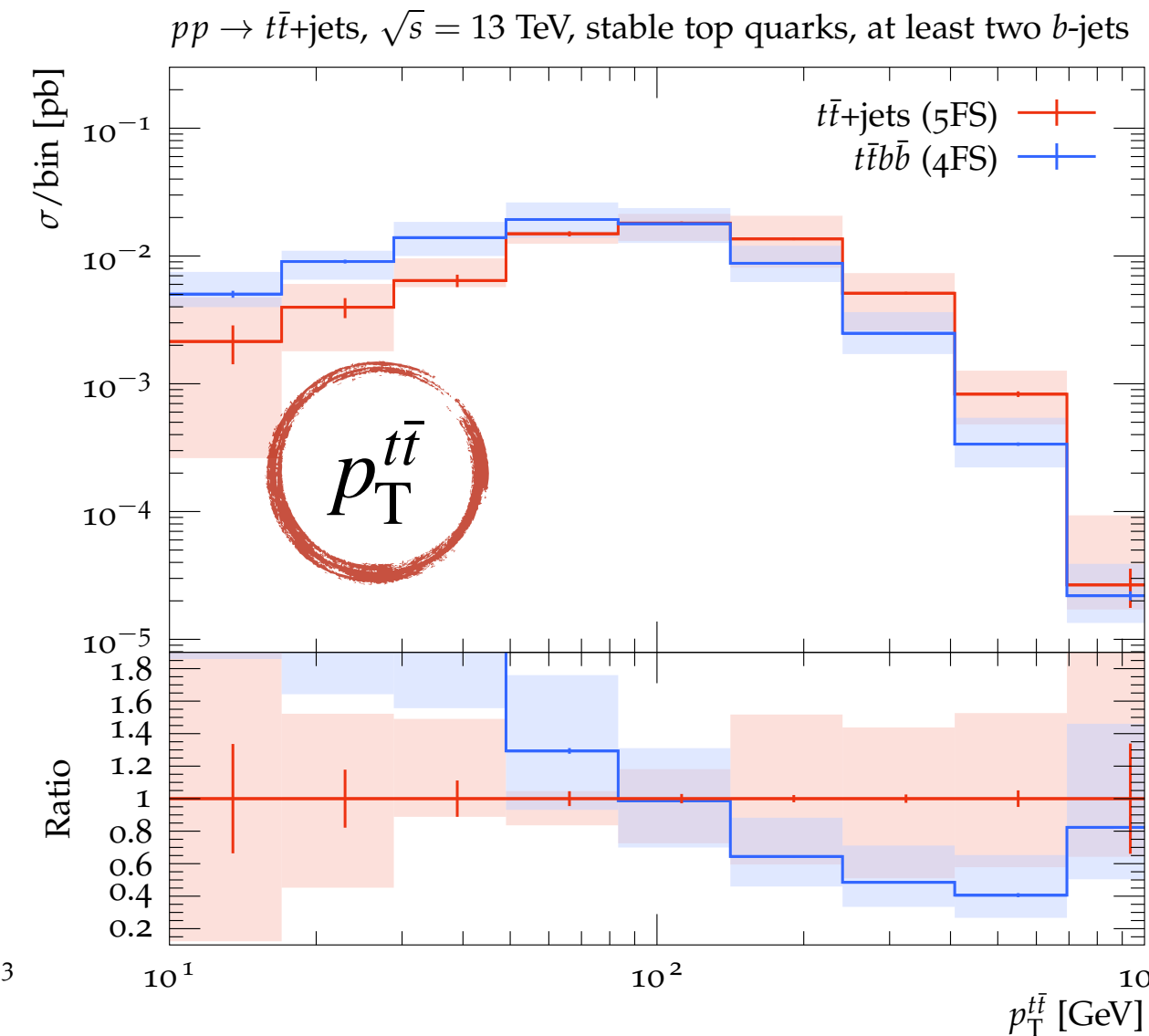
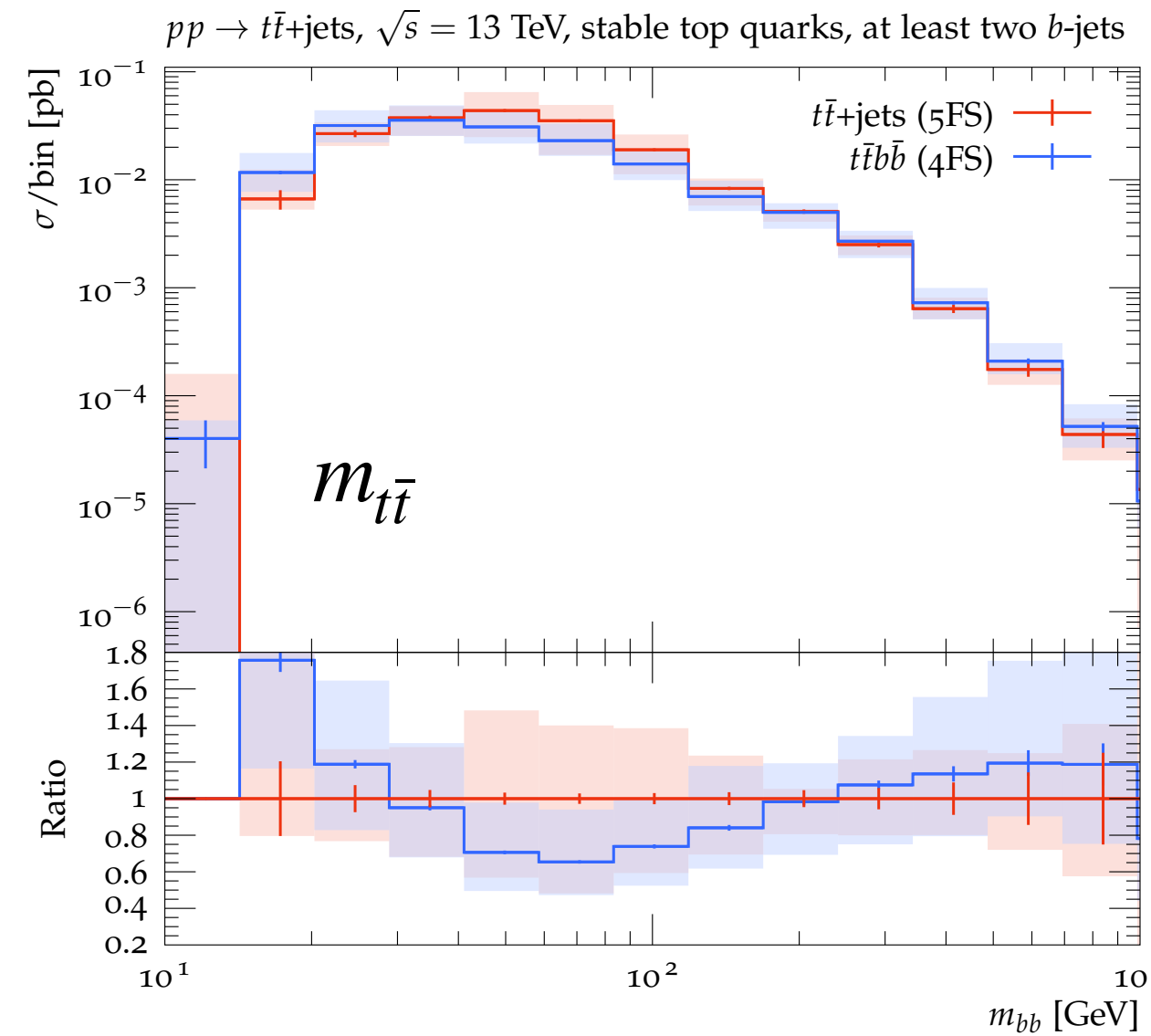
5FS vs 4FS: at least 1 b -jet scenario

- ▶ For most of the variables, 4FS and 5FS predictions are compatible within the uncertainty bands
- ▶ 5FS uncertainty is more reliable than the 4FS one, since the 4FS matching uncertainty is expected to be significant but is not included
- ▶ $p_T^{t\bar{t}}$: 5FS predicts a much harder spectrum than 4FS (but this difference is expected to be covered by the full 4FS uncertainty)
 - ➔ We investigated it further, see next slides

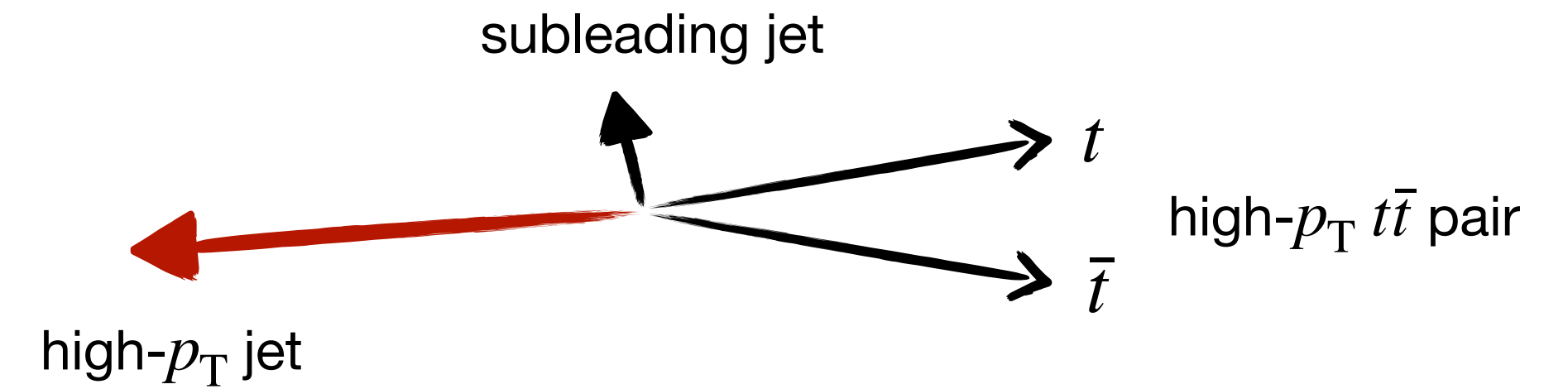


5FS vs 4FS: at least 2 b -jets scenario

- ▶ Similar picture as for the ≥ 1 b -jet selection
- ▶ Difference in the $p_T^{t\bar{t}}$ spectrum (expected to be covered by the full 4FS uncertainty)
- ▶ The rest of the variables are in agreement

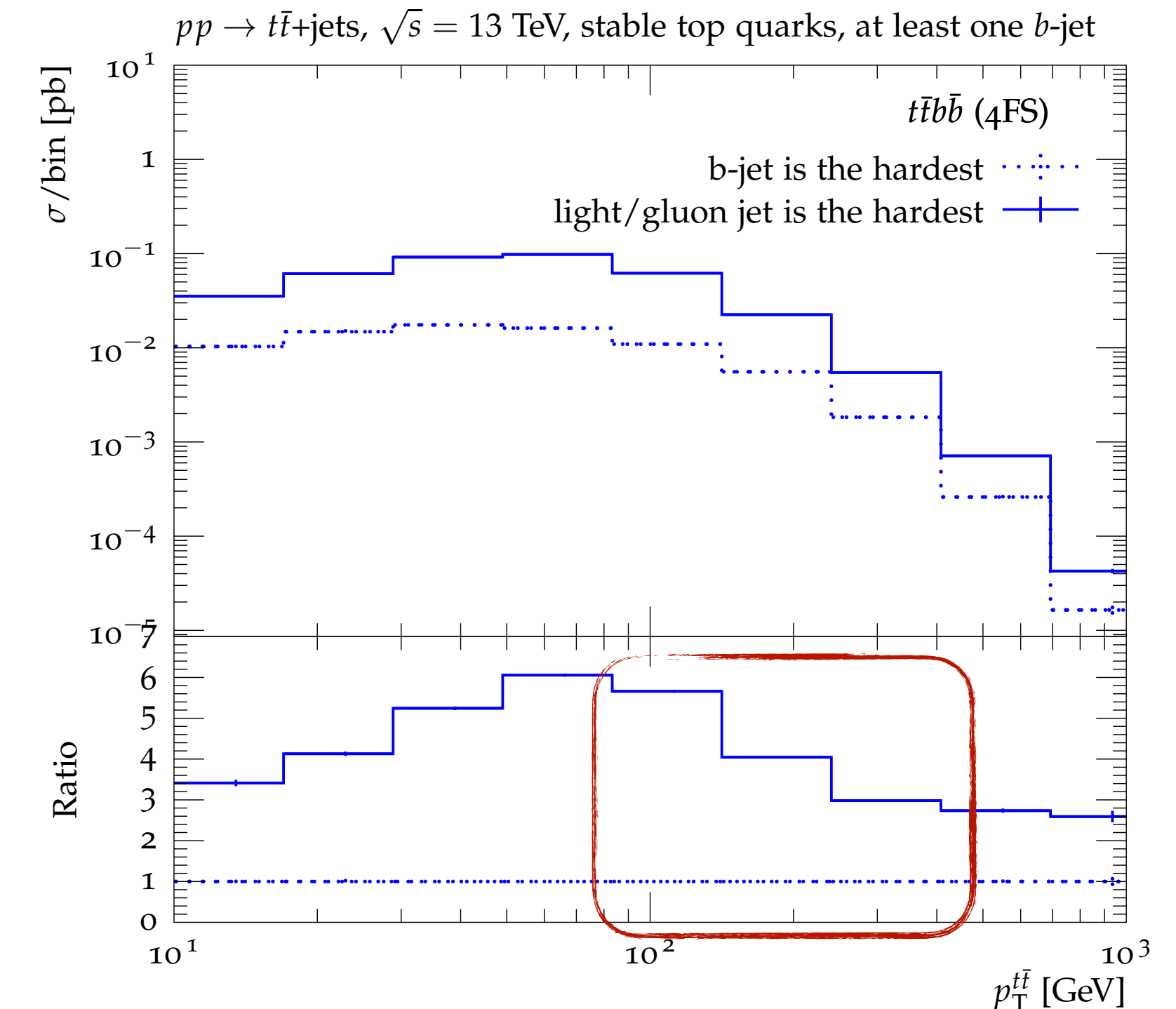
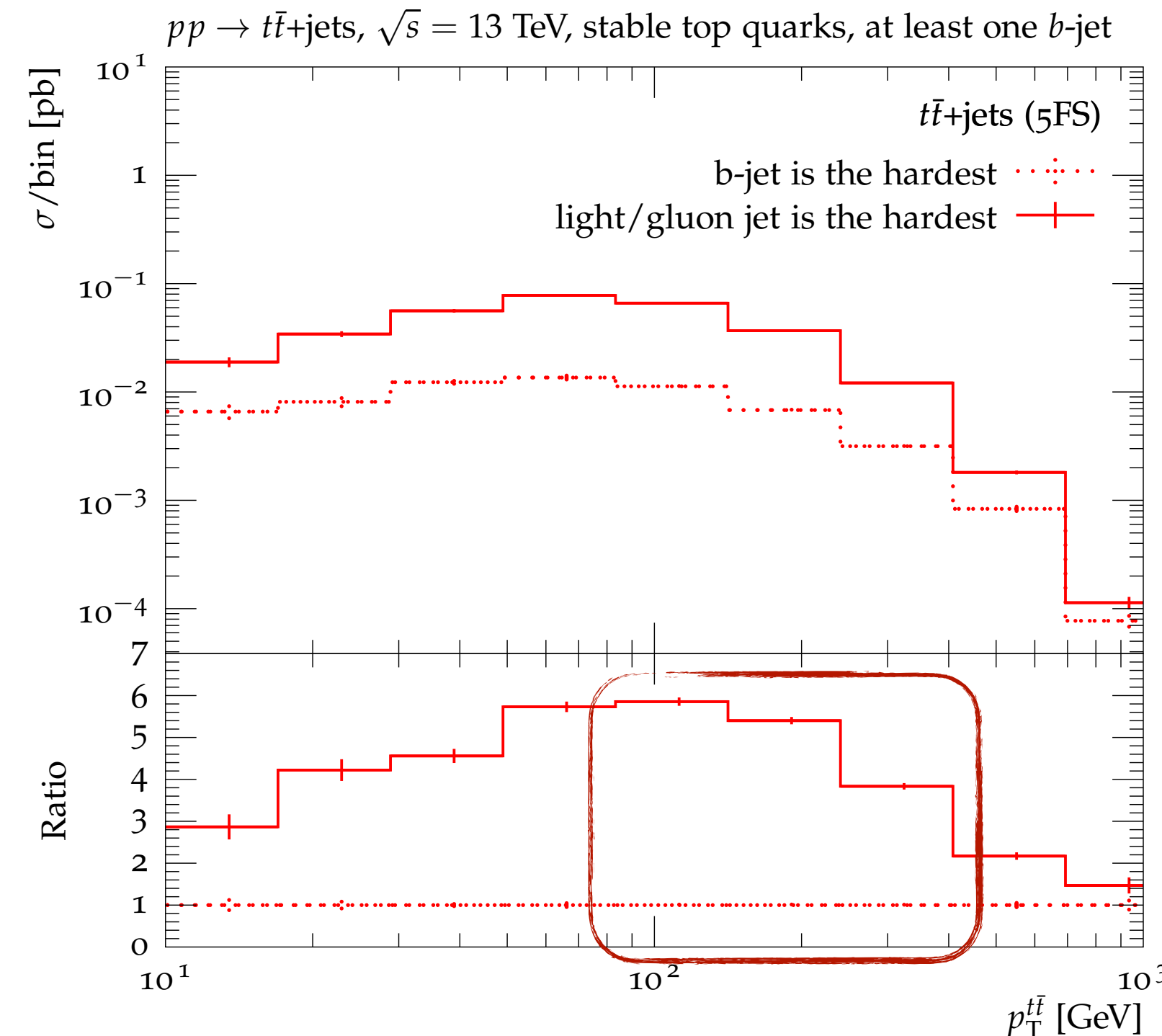


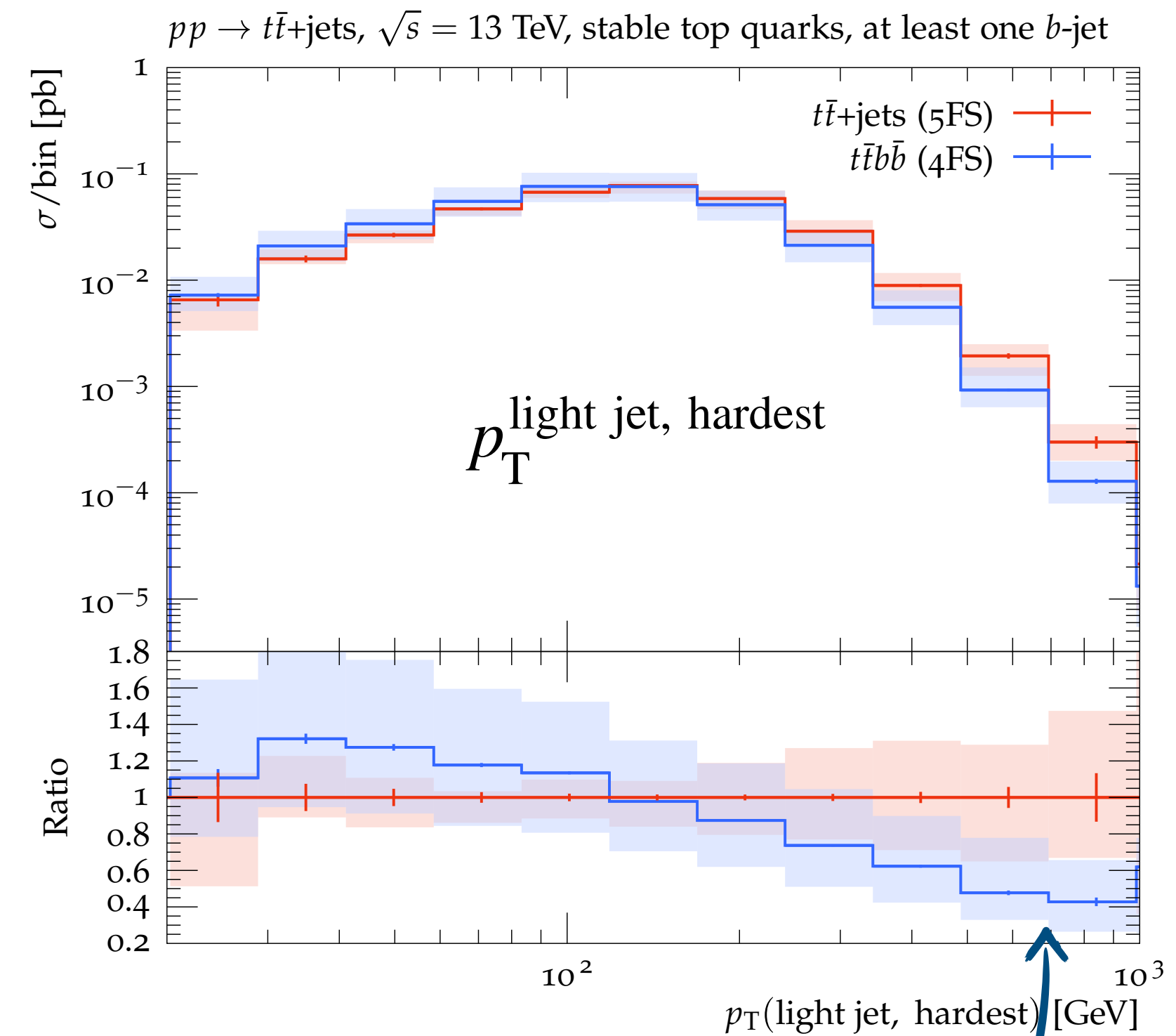
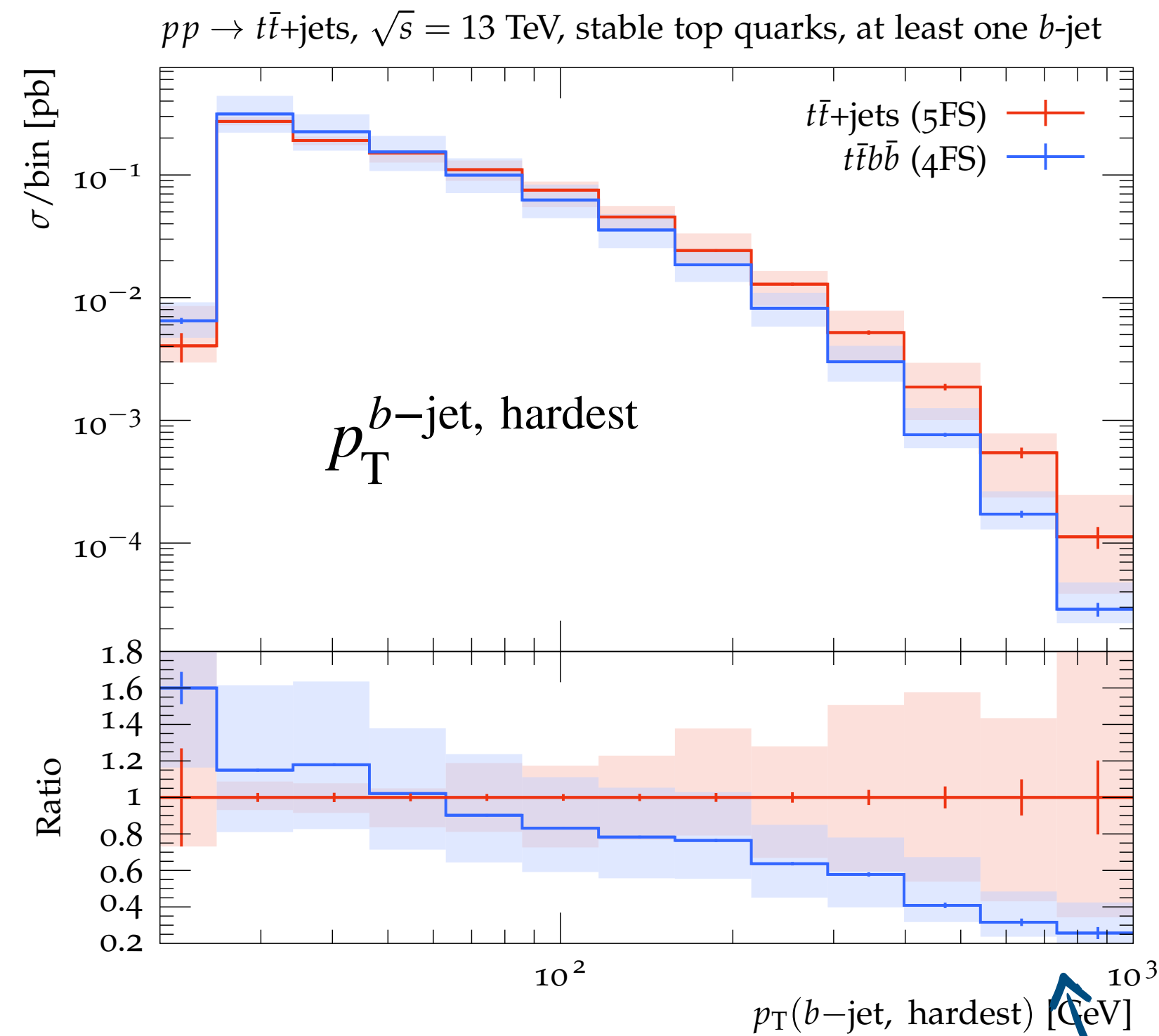
- ▶ At large $p_T^{t\bar{t}}$, it is kinematically most-likely that the $t\bar{t}$ pair recoils against a single hard jet
- ▶ If the hardest jet is a light jet:
 - 5FS: described at NLO (most likely it is a gluon jet)
 - 4FS: described at LO or by the PS
 - No $t\bar{t}gg$ events from the ME
 - There is no hard gluon to recoil from



at least 1 b -jet selection

- ▶ For high $p_T^{t\bar{t}}$, the fraction of events with the hardest jet being light-flavoured is indeed larger in the 5FS
- ▶ But after $p_T^{t\bar{t}} \sim 500$ GeV the situation is opposite — why?
 - Let's look again at the jet p_T distributions...

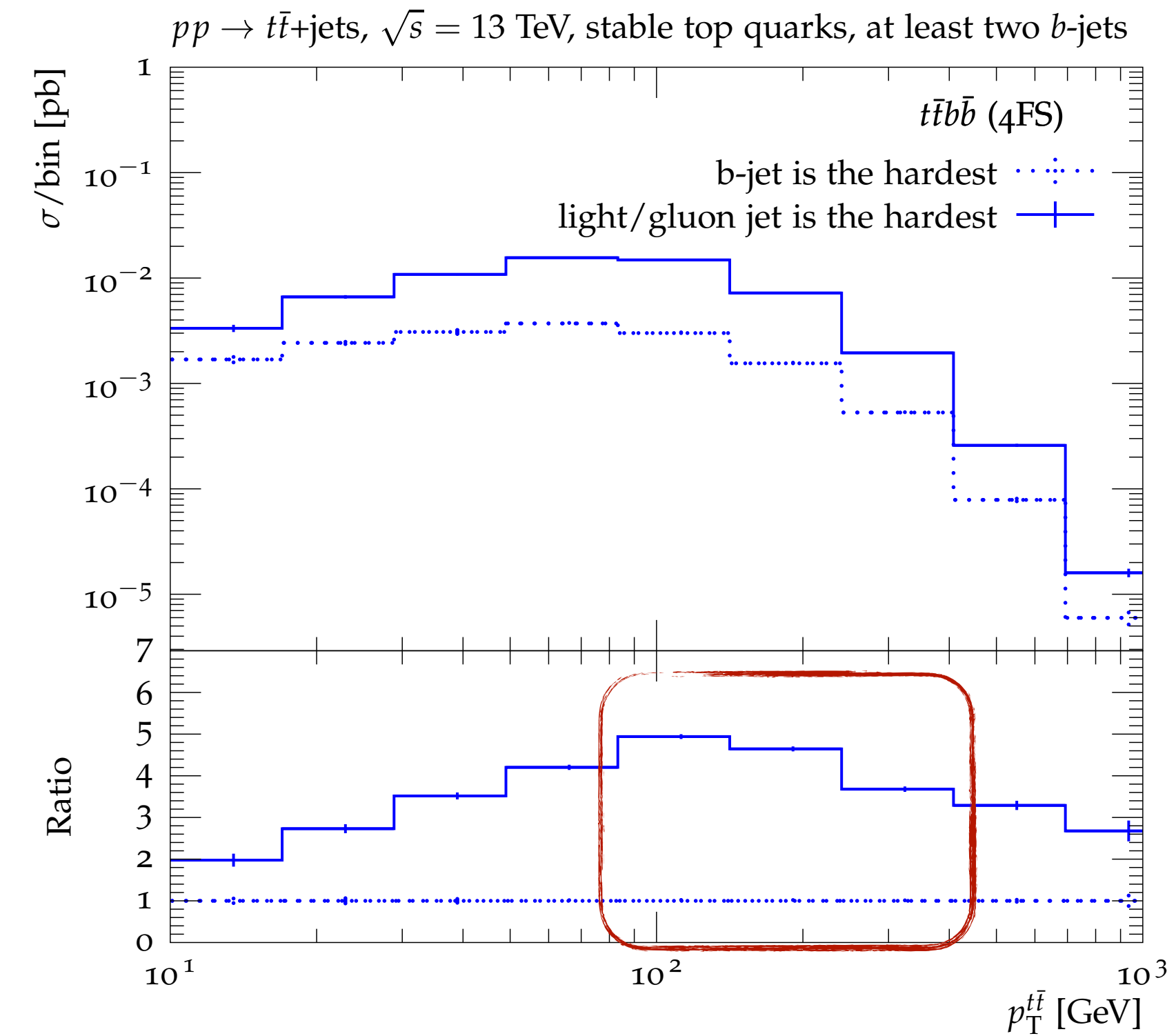
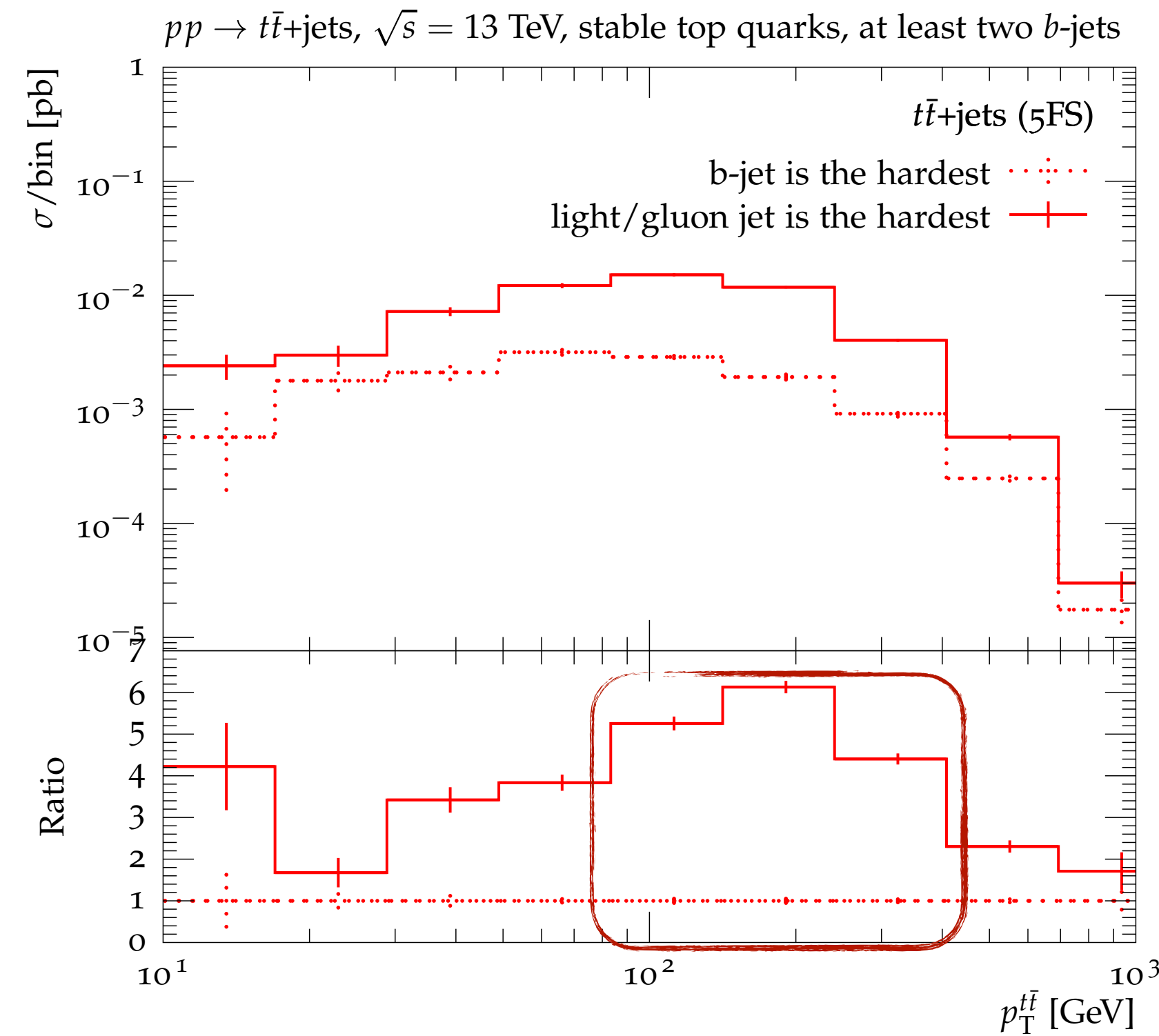




at very high p_T , the ratio of the 5FS over the 4FS predictions is larger for b -jets than for light jets

at least 2 b -jets selection

- ▶ The difference in the fraction of the hardest light jets in even more pronounced in the ≥ 2 b -jet selection



- ➔ The reason for the large 5FS–4FS difference in the $p_T^{t\bar{t}}$ spectrum at large momenta is
 - The correlation between $p_T^{t\bar{t}}$ and $p_T^{\text{light jet, hardest}}$
 - Expected 5FS–4FS difference between the fraction of events with the hardest jet being light-flavoured

What about a comparison to data?

- ▶ Not done yet
- ▶ A study is planned within the ATLAS Physics Modelling Group
- ▶ Once the new MadGraph5_aMC@NLO release is out, the idea is to produce an "official" sample (5FS, FxFx, with enhancement) and check:
 - data/MC comparison using the ATLAS $t\bar{t}b\bar{b}$ Rivet routine
 - performance of the enhancement feature

To summarise:

- ▶ $t\bar{t}b\bar{b}$ production serves as a significant background process across various high-energy physics phenomena
 - ▶ 5FS calculation of $t\bar{t}b\bar{b}$ at NLO yields the most accurate prediction for this process to date
 - no large logarithms appearing in the matrix element calculation
 - no complications when matching to a parton shower
 - ▶ We compute the $t\bar{t}$ + jets process with up to 2 jets at NLO using the FxFx merging prescription and match it to the Pythia8 shower
 - ▶ To improve the efficiency of selecting events with additional b -jets we enhance the probability of producing short-distance events with additional b -quarks using a newly implemented feature in the MadGraph5_aMC@NLO generator
 - This makes producing the $t\bar{t}b\bar{b}$ in the 5FS at NLO more viable, given the computational demands of the 5FS approach
- * *Similar heavy-flavour enhancement could also be applied to the “fusion” method in Sherpa*