Hadronic Top-quark Polarimetry with ParticleNet

LHCTopWG meeting, Nov 12, 2024

Alberto Navarro

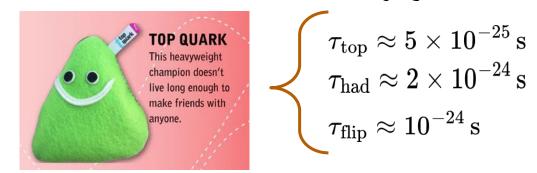


Based on arXiv:2407.01663 and arXiv:2407.07147

with Z. Dong, D. Gonçalves, K. Kong, and A. J. Larkoski

Motivation

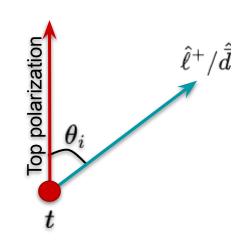
- > Spin correlation studies in top quark physics plays an important role in new physics searches (e.g., resonant or non-resonant searches), observation of entanglement, etc.
- > Polarization is a valuable tool when studying top-quark physics.



> The top polarization can be probed from the kinematics of its decay products.

> The direction of the top decay products correlates with the top polarization axis

$$rac{1}{\Gamma}rac{d\Gamma}{d\cos heta_i}=rac{1}{2}(1+peta_i\cos heta_i)$$

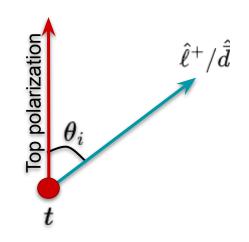


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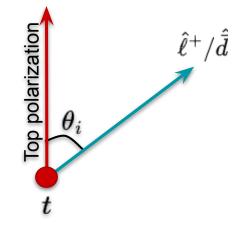
Spin analyzing power

$\ell^+/ar{d}$	u/u	$oldsymbol{b}$
1	-0.32	-0.41



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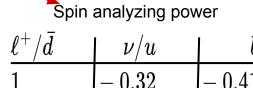
$$rac{\ell^+/ar{d}}{1} \hspace{0.5cm} |\hspace{0.2cm}
u/u \hspace{0.5cm} |\hspace{0.2cm}
ldot \ 1 \hspace{0.5cm} |\hspace{0.2cm} -0.41$$

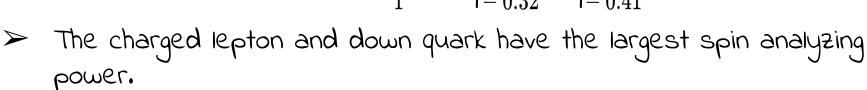
Spin analyzing power

> The charged lepton and down quark have the largest spin analyzing power.

> The direction of the top decay products correlates with the top polarization axis

$$rac{1}{\Gamma}rac{d\Gamma}{d\cos heta_i} = rac{1}{2}(1+poldsymbol{eta_i}\cos heta_i)$$



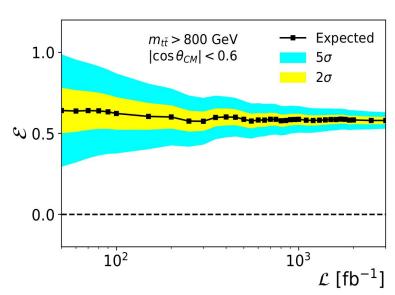


Hadronic decaying top offers higher statistics so fully recovering its polarization is crucial.

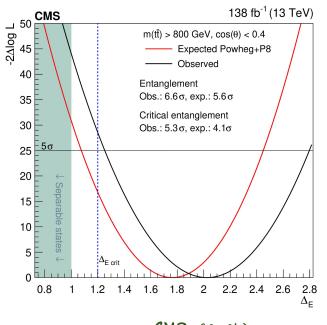
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Entanglement with semileptonic tops

> very recent application of hadronic top polarization in measuring entanglement with semileptonic top quarks



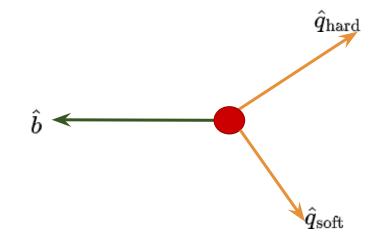
Dong, Gonçalves, Kong and **AN** (2023) See also Han, Low, WU (2023) Alberto Navarro



CMS (2024)

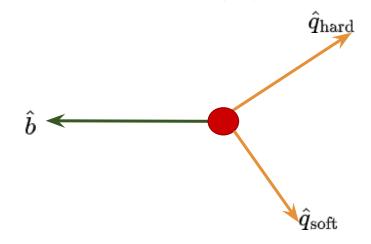
Hadronic top polarimetry

> Identifying a down-quark in a collider environment is challenging.



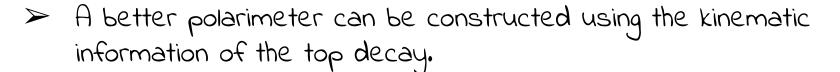
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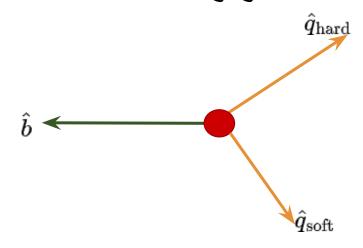
- > Identifying a down-quark in a collider environment is challenging.
- > An alternative approach is to use the light jet that is softest in the top rest frame.
- This choice gives a spin analyzing power equals to 0.5 Jeżabek (1994)



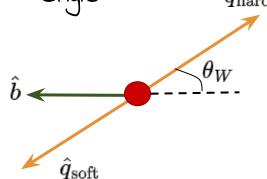
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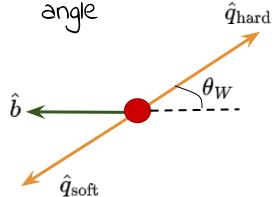




 \succ Look at the top decay from the W-rest frame and define the helicity angle $\hat{q}_{
m hard}$



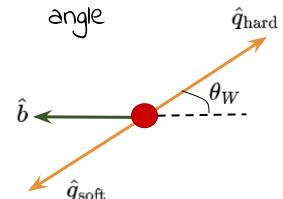
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$$p(d
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$$p(d
ightarrow q_{ ext{soft}}|c_W) = rac{p(-|c_W|)}{p(|c_W|) + p(-|c_W|)}$$

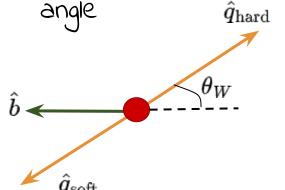
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$$p(d o q_{
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 down-quark emitted in forward hemisphere

$$p(d o q_{
m soft}|c_W) = rac{p(-|c_W|)}{p(|c_W|) + p(-|c_W|)} iggr\}$$
 down-quark emitted in backward hemisphere

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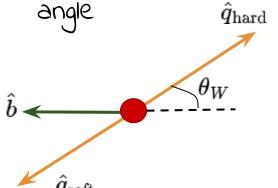


$$p(d o q_{
m hard}|c_W) = rac{p(|c_W|)}{p(|c_W|) + p(-|c_W|)} egin{displays l} ext{down-quark emitted in} \ ext{forward hemisphere} \end{cases}$$

$$p(d o q_{ ext{soft}}|c_W) = rac{p(-|c_W|)}{p(|c_W|) + p(-|c_W|)} egin{cases} ext{down-quark emitted in} \ ext{backward hemisphere} \end{cases}$$

$$>$$
 Define the optimal hadronic direction through kinematics as $ec{q}_{
m opt}^{
m \, kin}=p(d o q_{
m hard}|c_W)\hat{q}_{
m hard}+p(d o q_{
m soft}|c_W)\hat{q}_{
m soft}$

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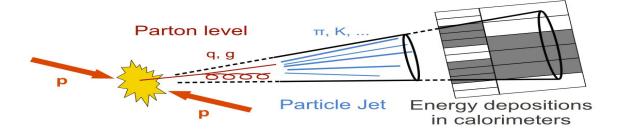
$$ec{q}_{ ext{opt}}^{ ext{ kin}} = p(d
ightarrow q_{ ext{hard}} | c_W) \hat{q}_{ ext{hard}} + p(d
ightarrow q_{ ext{soft}} | c_W) \hat{q}_{ ext{soft}}$$

> The length of this vector is the spin analyzing power and equals 0.64.

Tweedie (2014)

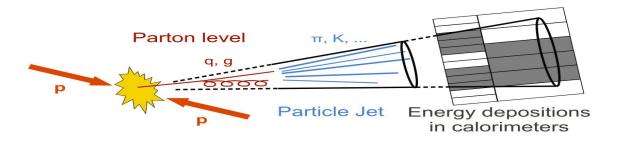
Going beyond kinematics

> At particle level, there's more information available from the jets.



Going beyond kinematics

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Assume some set of observables $\{\mathcal{O}\}$ is measured and generalized the optimal hadronic direction to

$$ec{q}_{ ext{opt}} = p(d
ightarrow q_{ ext{hard}} | \{\mathcal{O}\}) \hat{q}_{ ext{hard}} + p(d
ightarrow q_{ ext{soft}} | \{\mathcal{O}\}) \hat{q}_{ ext{soft}}$$

Dong, Gonçalves, Kong, Larkoski and AN (2024) Dong, Gonçalves, Kong, Larkoski and AN (2024)

The analytical approach

Start by considering "global jet observables" such as kinematics, jet charge and particle multiplicity.

$$p(ar{d}
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ightarrow q_{ ext{hard}} | c_W) + rac{p(\mathcal{Q}_{\kappa,h} | u
ightarrow q_{ ext{hard}}, N_h)}{p(\mathcal{Q}_{\kappa,h} | ar{d}
ightarrow q_{ ext{hard}}, N_h)} \, p(u
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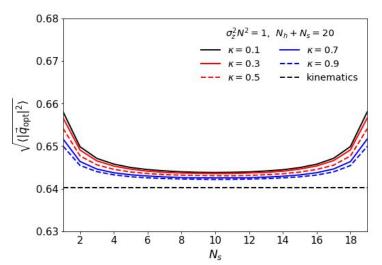
$$p(\bar{d} \rightarrow q_{\mathrm{hard}}|c_{W}, \mathcal{Q}_{\kappa,h}, N_{h}, \mathcal{Q}_{\kappa,s}, N_{s}) = \frac{1}{p(\bar{d} \rightarrow q_{\mathrm{hard}}|c_{W}) + \underbrace{\frac{p(\mathcal{Q}_{\kappa,h}|u \rightarrow q_{\mathrm{hard}}, N_{h})}{p(\mathcal{Q}_{\kappa,h}|\bar{d} \rightarrow q_{\mathrm{hard}}, N_{h})}} p(u \rightarrow q_{\mathrm{hard}}|c_{W})}_{\text{Likelihood ratio of the down of the down and up being the soft jet}} \times \underbrace{\frac{1}{p(\mathcal{Q}_{\kappa,s}|\bar{d} \rightarrow q_{\mathrm{soft}}, N_{s})}} p(\bar{d} \rightarrow q_{\mathrm{soft}}|c_{W}) + p(u \rightarrow q_{\mathrm{soft}}|c_{W})}_{\text{the hard jet}}}_{\text{the hard jet}}$$

Dong, Gonçalves, Kong, Larkoski and AN (2024)

- > Jet charge distributions conditioned on multiplicity are Gaussians.

 Kang, Larkoski and Yang (2023)
- Compute the spin analyzing power as the square root of the mean squared-vector length

$$\langle |ec{q}_{ ext{opt}}|^2
angle = \int dc_W \, p(c_W) \int d\mathcal{Q}_{\kappa,h} \, d\mathcal{Q}_{\kappa,s} \, p(\mathcal{Q}_{\kappa,h},\mathcal{Q}_{\kappa,s}|c_W,N_h,N_s) \, |ec{q}_{ ext{opt}}|^2$$

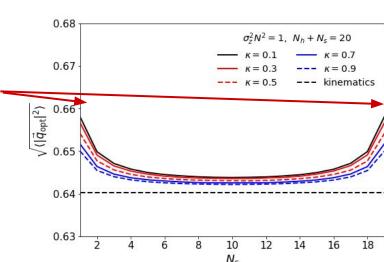


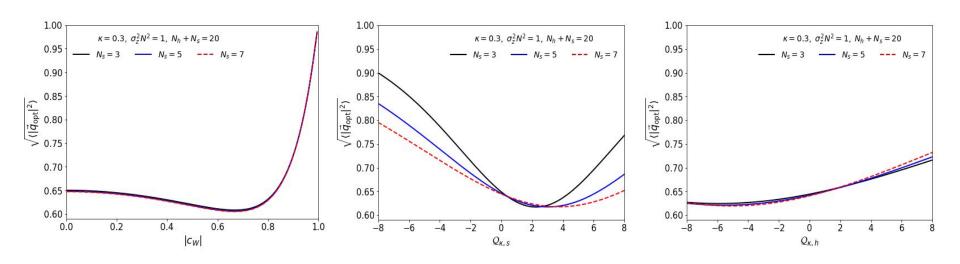
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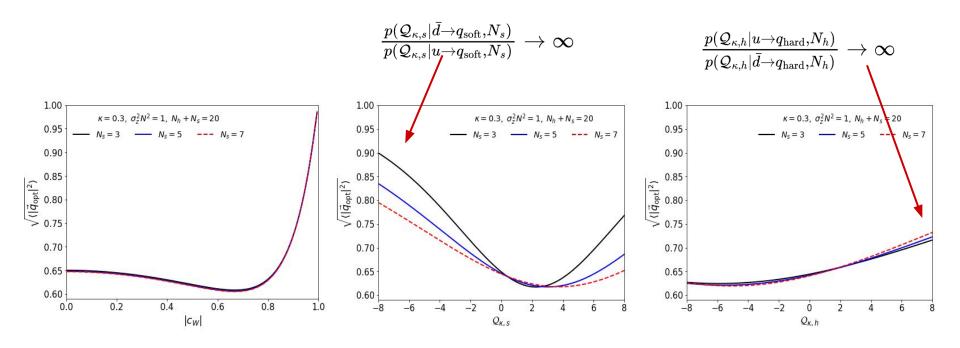
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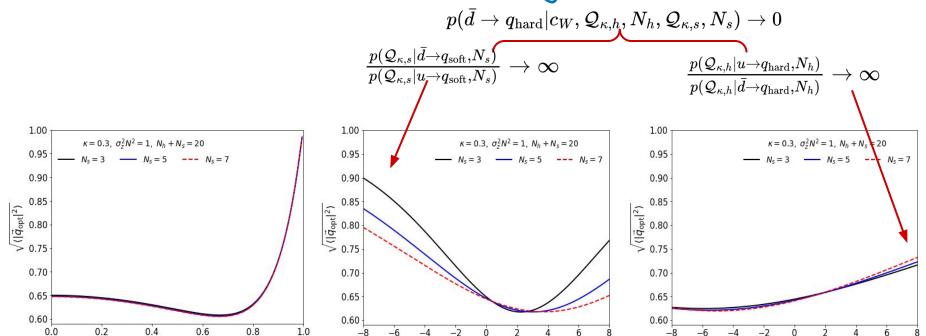
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Better discrimination for low $N_{\rm s}$ and $N_{\rm h}$





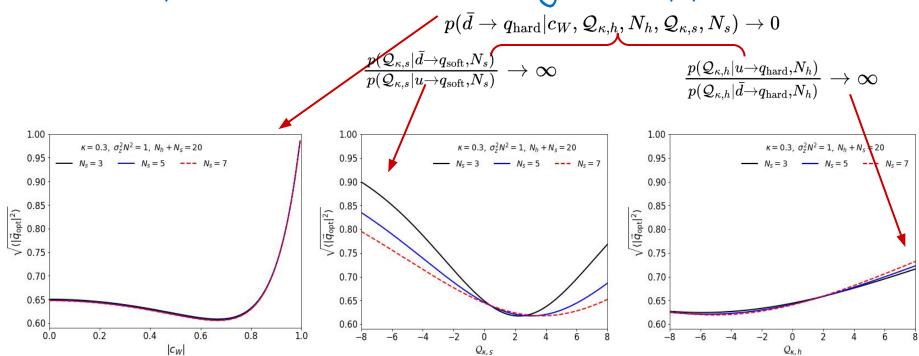




 $Q_{K,S}$

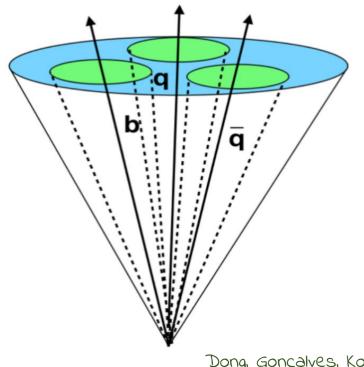
 $Q_{\kappa,h}$

 $|c_W|$



The ML approach

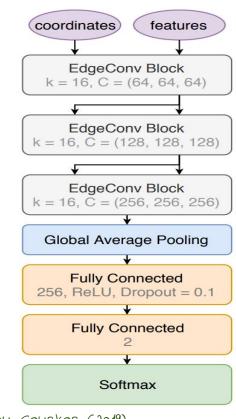
- Could improve result further by incorporating information of jet constituents such as charge, momenta and ID.
- > Train network to identify the down jet within the top jet.
- > Using a simple DNN may not be ideal due to complexity of data and number of features.



Dong, Gonçalves, Kong, Larkoski and AN (2024)

ParticleNet

- Idea is to represent a jet as an unordered, permutation invariant set of particles (a particle cloud).
- > Analogous to the point cloud representation of 3D shapes used in computer vision.
- > Has been successfully used for top and quark-gluon tagging.
- > need to modify for our purpose.



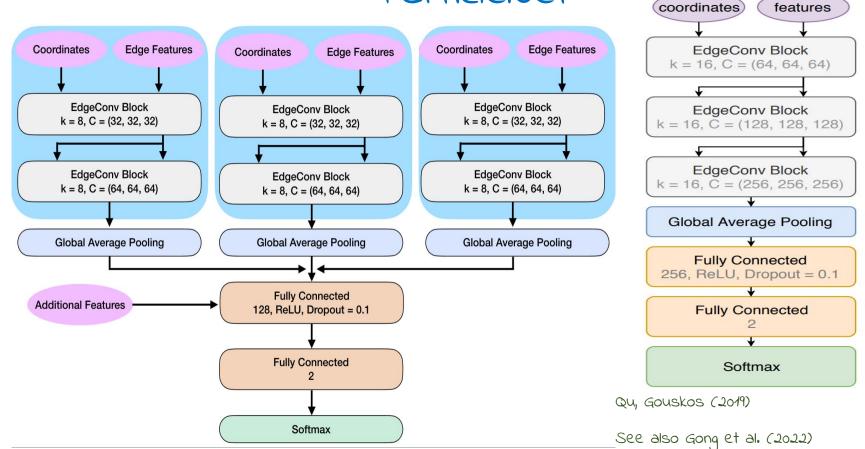
Qu, Gouskos (2019)

See also Gong et al. (2022)

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Bogatskiy, Hoffman, Miller, Offermann (2022)

ParticleNet

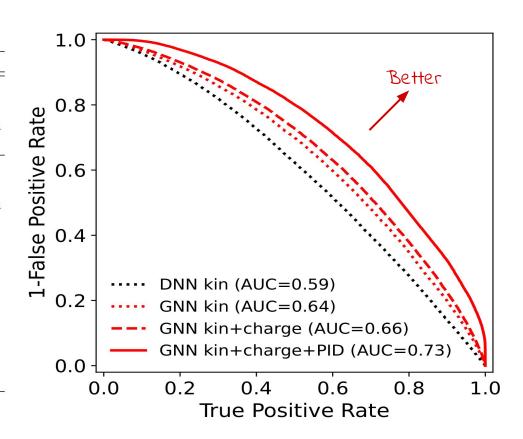


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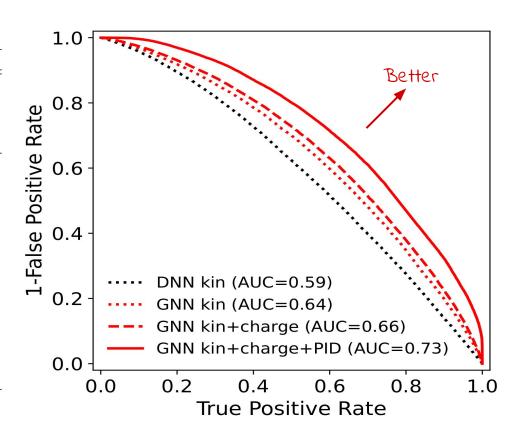
Bogatskiy, Hoffman, Miller, Offermann (2022)

Variable	Definition	
$\Delta \eta_t$	difference in pseudorapidity between	
	the particle and the top jet axis	
$\Delta\phi_t$	difference in azimuthal angle between	
	the particle and the top jet axis	
$\Delta \eta_j$	difference in pseudorapidity between	
	the particle and the subjet axis	
$\Delta\phi_j$	difference in azimuthal angle between	
	the particle and the subjet axis	
$\log p_T$	logarithm of the particle's p_T	
$\log E$	logarithm of the particle's Energy	
q	electric charge of the particle	
isElectron	if the particle is an electron	
isMuon	if the particle is a muon	
isPhoton	if the particle is a photon	
is Charged Hadron	if the particle is a charged hadron	
isNeutralHadron	if the particle is a neutral hadron	

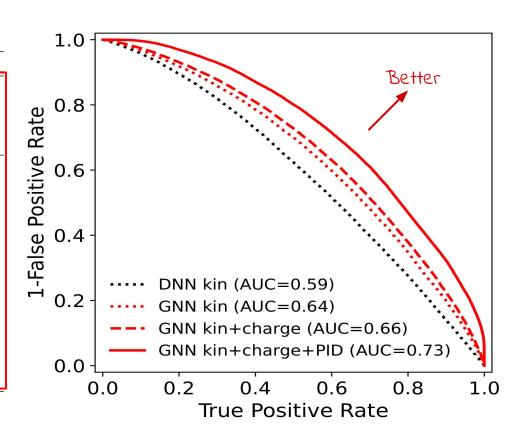
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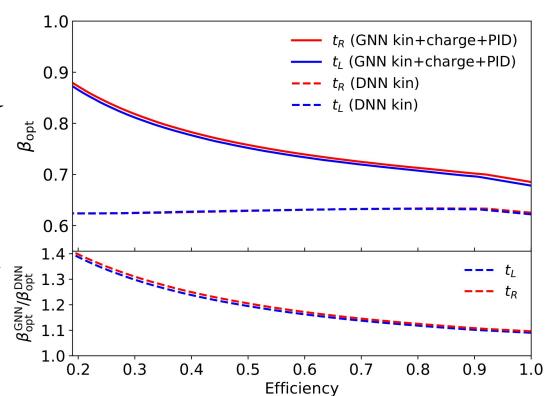


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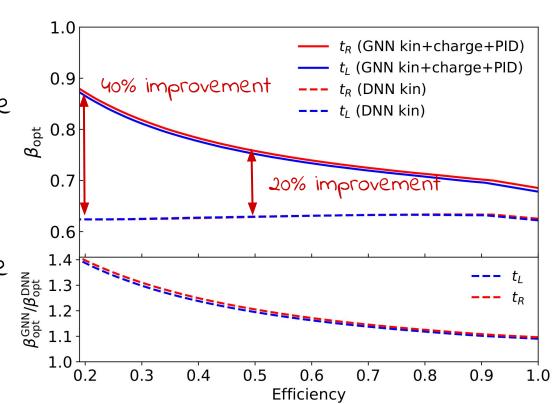
Improvements with ML approach

- > Use network score as the probability of the down-type jet being the soft jet.
- Apply cuts on the network score to enhanced the results as long as total rate is still larger than in di-leptonic case.



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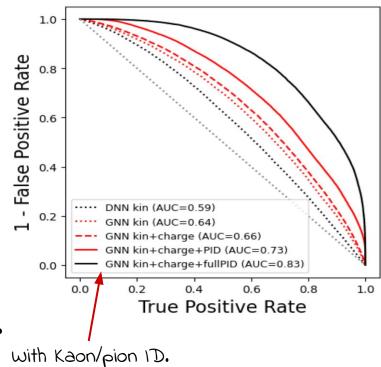
Summary

- > Hadronic top-quark polarization can be improved by incorporating information of global jet dynamics and constituents.
- > Owing to the larger rate (factor of for semileptonic w.r.t dileptonic case), the spin analyzing power in hadronic decays can improve by approximately 20% (40%) compared to the kinematic approach, assuming an efficiency of 0.5 (0.2) for the GNN.
- This results could be used to further improve studies were the top polarization is relevant such as new physics searches, spin correlations, etc.

Backup

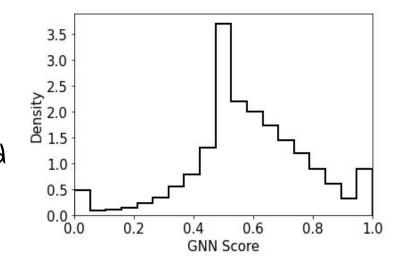
Path for further improvements

- Performance can be improved with pion/kaon discrimination.
- They are not well distinguished at the LHC. This may change with future detectors, though.
- > Charm tagging can also be incorporated.
- > Any extra information should further improve our current results.



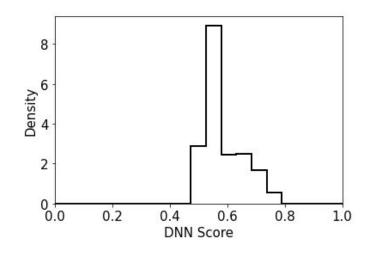
GNN score

- > Total of 77k trainable parameters.
- > Trained on 200k unpolarized sample.
- > Tested on 50k right-hand and left-hand polarized samples.



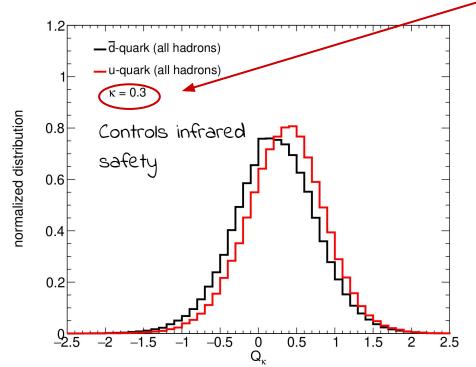
DNN score

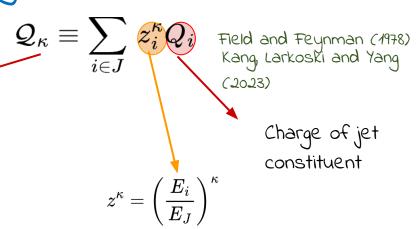
- Using jets momenta as input.
- > Total of 20k trainable parameters.
- > Trained on 200k unpolarized sample.
- > Tested on 50k right-hand and left-hand polarized samples.



Jet charge







Energy fraction of jet constituent

Parametric form of jet charge distributions Kang, Larkoski and Yang (2023)

> Assume hadrons in the jet are produced through identical independent processes and the multiplicity is large

$$p(\mathcal{Q}_{\kappa}|u,N) = rac{1}{\sqrt{2\pi\sigma^2}}\,e^{-rac{(\mathcal{Q}_{\kappa}-\mu_u)^2}{2\sigma^2}} \qquad p(\mathcal{Q}_{\kappa}|ar{d},N) = rac{1}{\sqrt{2\pi\sigma^2}}\,e^{-rac{(\mathcal{Q}_{\kappa}-\mu_{ar{d}})^2}{2\sigma^2}}$$

> Assume particles in the jets are pions and the SU(2) isospin of the pions is exact

$$\mu_u = rac{2}{3}N^{-\kappa}\left(1+rac{\kappa}{2}(\kappa-1)\sigma_z^2N^2+\cdots
ight) \quad \mu_{ar{d}} = rac{1}{3}N^{-\kappa}\left(1+rac{\kappa}{2}(\kappa-1)\sigma_z^2N^2+\cdots
ight) \quad \sigma^2 = rac{2}{3}N^{1-2\kappa}\left(1+\kappa(2\kappa-1)\sigma_z^2N^2+\cdots
ight)$$

The discrimination power

$$\eta = rac{(\langle Q_{\kappa}
angle_u - \langle Q_{\kappa}
angle_{ar{d}})^2}{\sigma^2} \sim rac{1}{N}ig(1-\kappa^2\sigma_z^2N^2+\ldotsig)$$

Data preparation

- > we generate 14 TeV $pp \to t\bar{t} \to \ell^\pm \nu 2b2j$ events using MG5, with no cuts except for pT(t) > 200 GeV.
- > Three sets of samples where the top quark is unpolarized, left hand polarized and right hand polarized in the $t\overline{t}$ rest frame.
- > Parton shower and hadronization are done with PYTHIA8 without MPI.
- > Identify the top jet using CA algorithm with R=1.5, and pT>250 GeV. And decluster to find the subjets with m < 30 GeV.
- > Keep the three subjets that satisfy 165 GeV < m(top) < 195 GeV.
- \succ Match the hadron level jets with true parton level momenta, by using the smallest ΔR between the two.

EdgeConv block

- Represents a point cloud as a graph, whose vertices are the points themselves, and the edges are constructed as connections between each point to its k nearest neighboring points
- A local patch needed for convolution is defined for each point as the k nearest neighboring points connected to it.
- > Can be viewed as a mapping from a point cloud to another point cloud with the same number of points.

