

NLO QCD and other recent results in POWHEG-BOX-ew

Clara Del Pio
with Mauro Chiesa and Fulvio Piccinini
INFN Pavia
EWWG 4-09-2024



This talk

- Recent results from Eur. Phys. J. C 84 (2024) 5, 539
most recent report on Z_ew-BMNNPV with comparison of different input schemes
- Results including NLO QCD

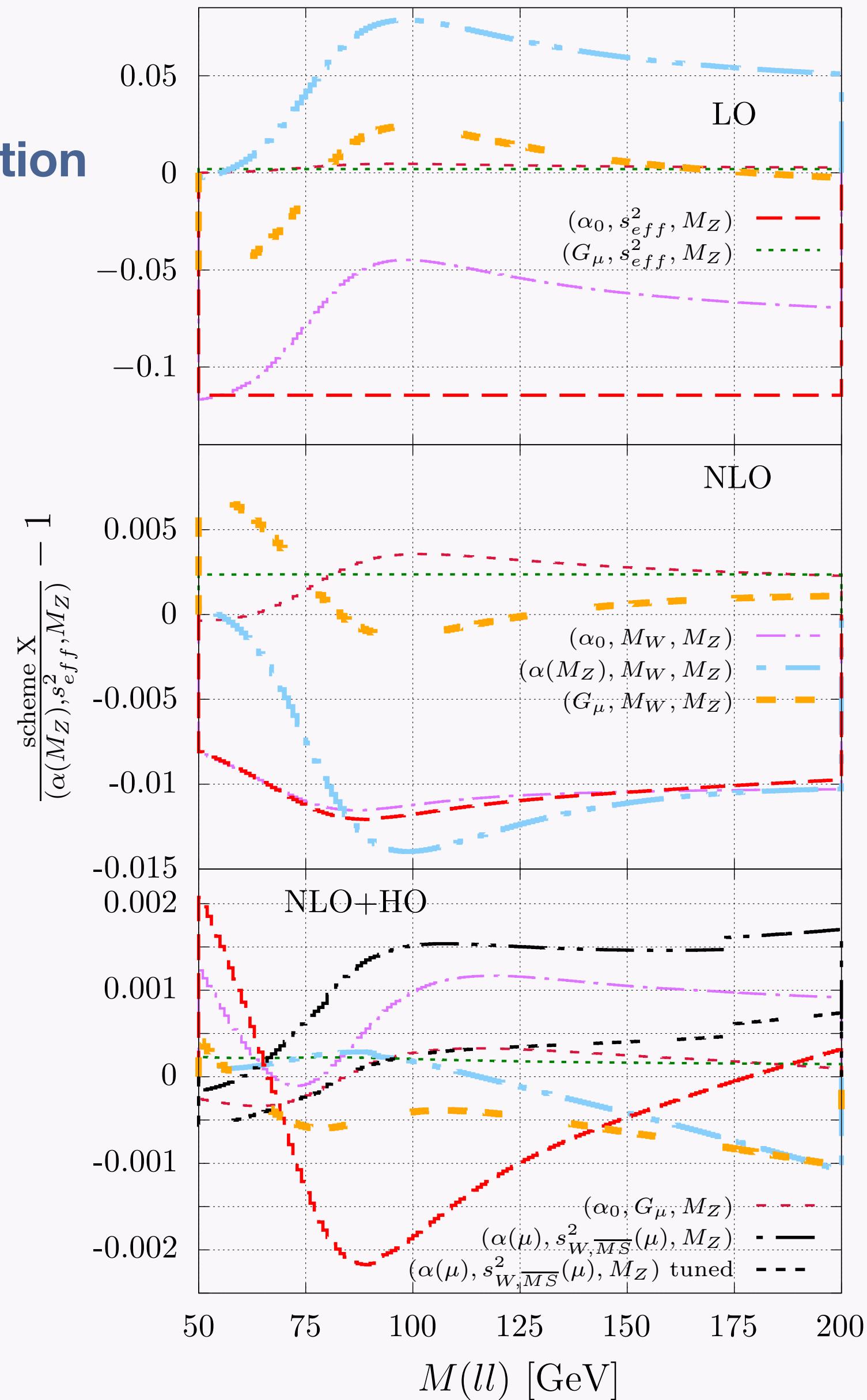
Main features of POWHEG Z_ew-BMNNPV

- Possibility to select among different EW input and renormalization schemes:

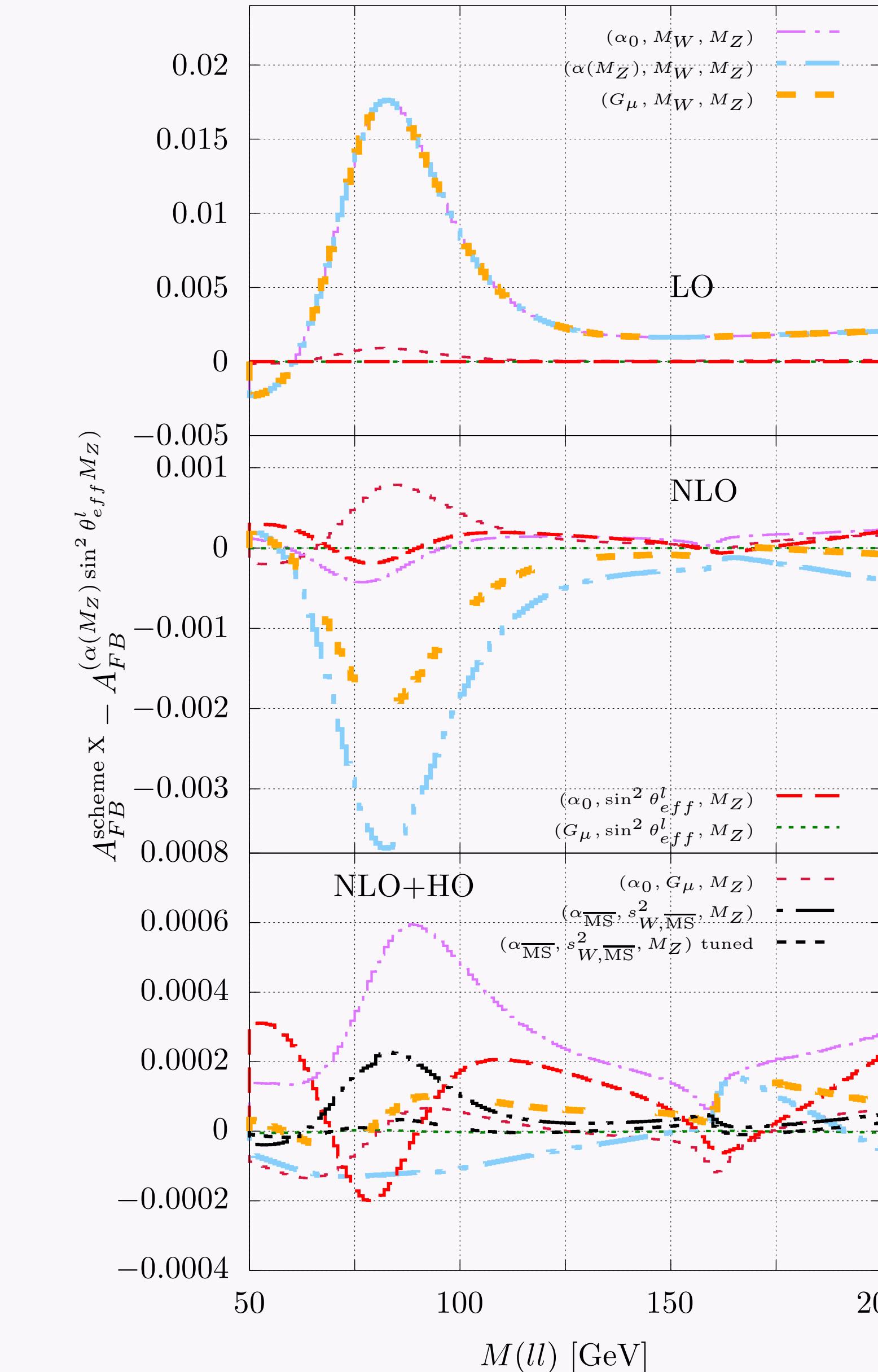
| | |
|--|--------------------------------------|
| $(\alpha_0 / \alpha(M_Z^2) / G_\mu, M_W, M_Z)$ | (α, G_μ, M_Z) |
| $(\alpha_0 / \alpha(M_Z^2) / G_\mu, s_{eff}^2, M_Z)$ | $(\alpha(\mu^2), s_w^2(\mu^2), M_Z)$ |
- Different options for resonance treatment: complex-mass scheme, pole scheme, factorization scheme
- Results for weak corrections are presented in the Z-peak region and high energy region (Sudakov regime)

Comparison among input schemes

Cross section



Asymmetry



Tuning (LEP inspired) I

Tuning from reference scheme (α_0, G_μ, M_Z)

$$(\alpha_0, s_{eff}^2, M_Z)$$

$$s_{eff}^2|_{G_\mu} = \frac{1}{2} - \sqrt{\frac{1}{4} - \frac{\pi}{\sqrt{2}G_\mu M_Z^2} \alpha(M_Z^2) \left(1 + \Delta\tilde{r}_{HO}\right)}$$

$$(\alpha_0, M_W, M_Z)$$

$$M_W|_{G_\mu} = \frac{M_Z^2}{2} \left(1 + \sqrt{1 - \frac{\sqrt{8}\pi}{G_\mu M_Z^2} \alpha(M_Z^2) \frac{1 + \Delta r^{(1)} - \Delta\alpha + \frac{c_W^2}{s_W^2} \Delta\rho^{(1,X)}}{1 + \frac{c_W^2}{s_W^2} \Delta\rho^{(X)}}} \right)$$

$$\alpha(M_Z^2) = \frac{\alpha_0}{1 - \Delta\alpha}$$

$$\Delta\tilde{r}_{rem} = \Delta\tilde{r}^{(1)} - \Delta\alpha + \Delta\rho^{(1)}$$

$$\Delta\tilde{r}_{HO} = \Delta\tilde{r}_{rem} - \Delta\rho$$

$$\Delta\rho = \Delta\rho^{(1)} + \Delta\rho^{(2)}$$

$$\Delta\rho^{(1,X)} = \frac{\Sigma_T^{ZZ}(M_Z)}{M_Z^2} - \frac{\Sigma_T^W(M_W)}{M_W^2} \Big|_{fin, \mu_{dim}=M_Z}$$

$$\Delta\rho^{(X)} = \Delta\rho^{(1,X)} + \Delta\rho^{(2,X)}$$

Tuning (LEP inspired) II

Higher orders in Born Improved Approximation with $\alpha(M_Z^2)$ and s_{eff}^2

(α_0, G_μ, M_Z)

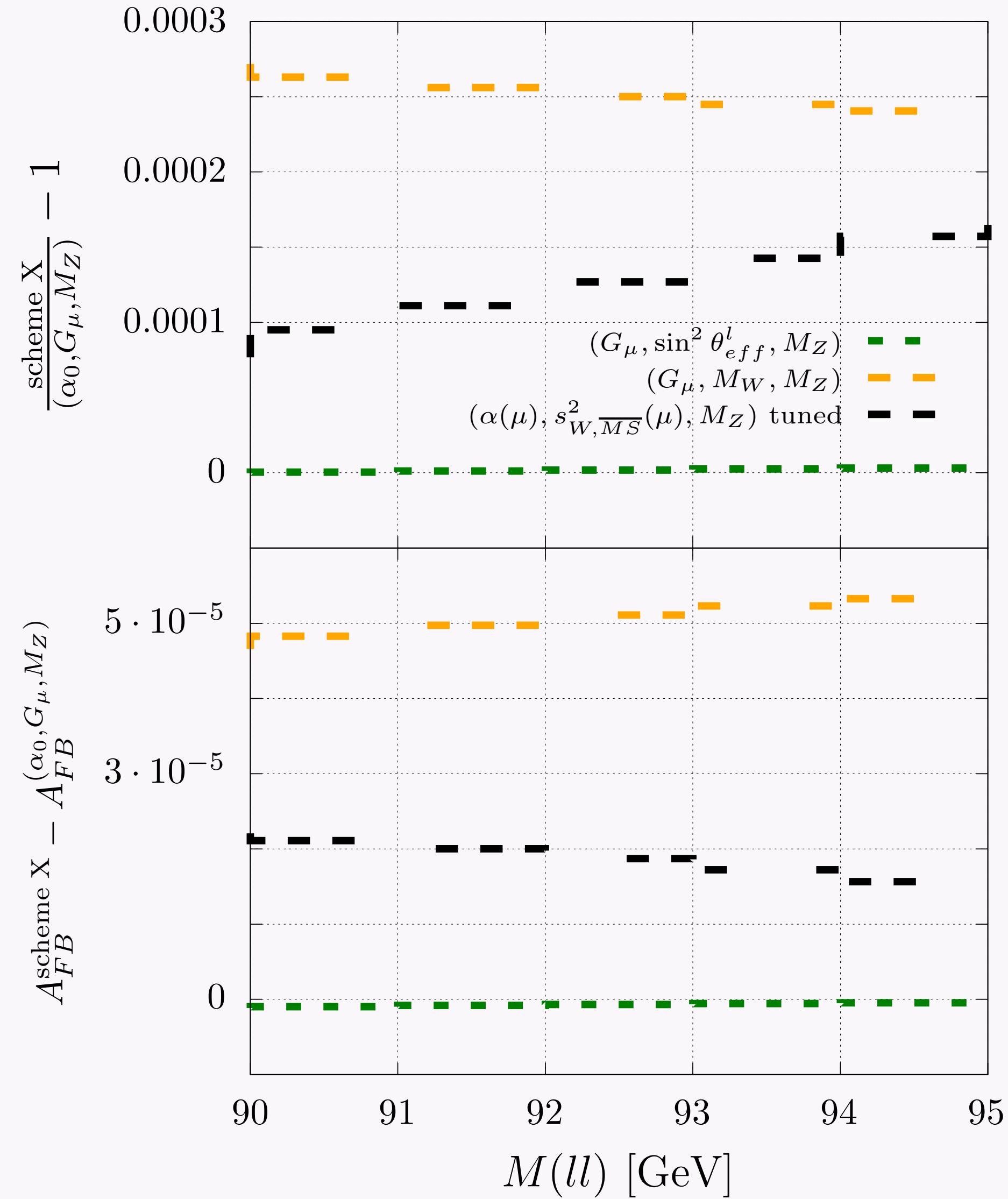
$$\tilde{s}_w^2, \text{NLO+HO} = \frac{1}{2} - \sqrt{\frac{1}{4} - \frac{\pi\alpha}{\sqrt{2}G_\mu M_Z^2}(1 + \Delta\tilde{r}|_{s_w^2})}$$

(α_0, M_W, M_Z)

$$\tilde{s}_w^2, \text{NLO+HO} = s_w^2 \left(1 + \frac{c_w^2}{s_w^2} \Delta\rho^{(X)} \right) \left[1 - \frac{c_w^2}{s_w^2} \Delta\rho^{(1,X)} + \frac{1}{s_w^2} \frac{1}{2} \frac{g_L g_R}{(g_L - g_R)^2} \text{Re} \left(\frac{\delta g_L}{g_L} - \frac{\delta g_R}{g_R} \right) \right]$$

$(\alpha_0, s_{eff}^2, M_Z)$ already ok

Results of tuning



How can this impact
on our uncertainty estimation?

Resonance treatment

Complex-mass scheme

Denner, Dittmaier, Roth, Wackerlo, Nucl. Phys. B 560 no. 1-3, 33–65, 1999
Denner, Dittmaier, Roth, Wieders, Nucl. Phys. B 724 no. 1-2, 247–294, 2005
Denner, Dittmaier, Nucl. Phys. B - Proceedings Supplements 160, 22–26, 2006

$$\mu_Z = M_Z - i\Gamma_Z M_Z \quad \mu_W = M_W - i\Gamma_W M_W$$

Pole scheme

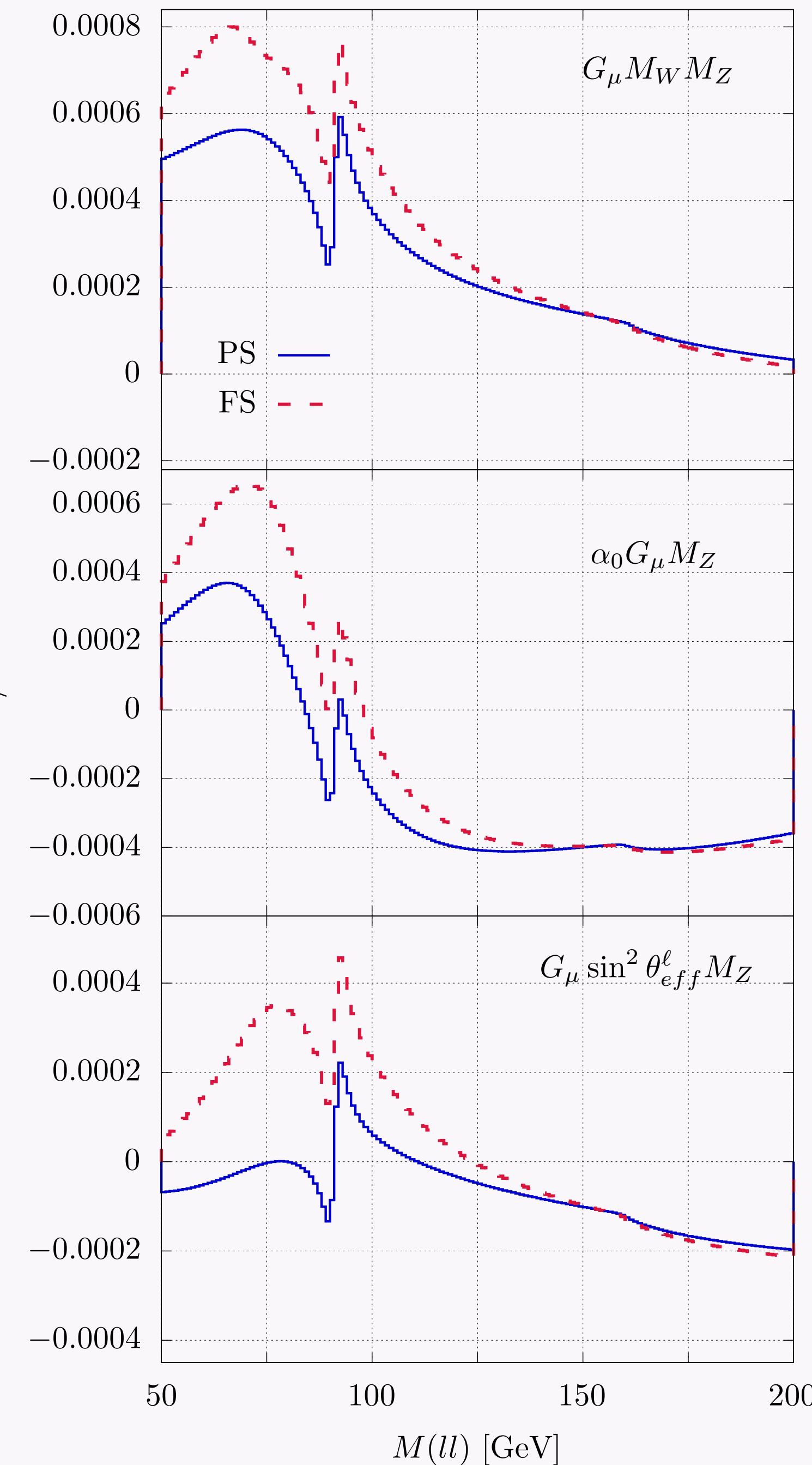
Stuart, Phys. Lett. B 262 no. 1, 113–119, 1991 - Sirlin, Phys. Lett. B 267 no. 2, 240–242, 1991
Gambino, Grassi, Phys. Rev. D 62 no. 7, 2000 - Grassi, Kniehl, Sirlin, Phys. Rev. D 65 no. 8, 2002
Stuart, Phys. Rev. Lett. 70, 3193–3196, 1993 - Dittmaier, Huber, JHEP 2010 no. 1, 2010)

$$\mathcal{M} = \frac{\tilde{R}(\mu_P^2)}{p^2 - \mu_P^2} + \frac{R(p^2) - R(M_P^2)}{p^2 - M_P^2} + \tilde{N}(p^2)$$

Factorization scheme

Argyres et al., Phys. Lett. B 358 no. 3-4, 339–346, 1995
Kurihara, Perret-Gallix, Shimizu, Phys. Lett. B 349 no. 3, 367–374, 1995
S. Dittmaier and M. Krämer, Phys. Rev. D 65 no. 7, 2002

$$f_P(p^2) = \frac{p^2 - M_P^2}{p^2 - \mu_P^2}$$



Sudakov regime

Channel with d-quarks only: no PDFs dependence
 → no large unphysical distortions at high energies

Sudakov logs

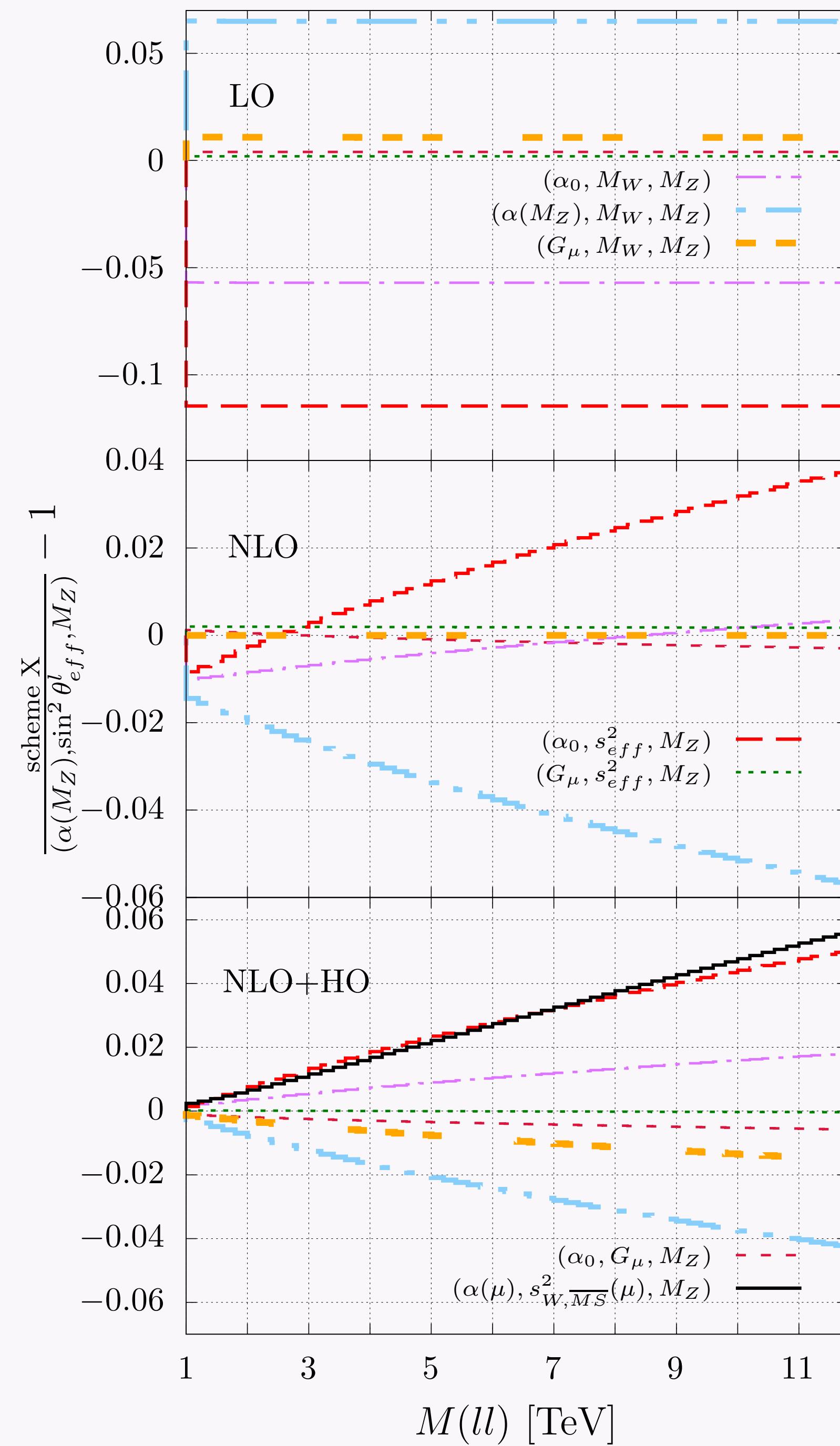
$$A(\alpha, s_w^2) \ln^2 \frac{s}{M_Z^2} + B(\alpha, s_w^2) \ln \frac{s}{M_Z^2}$$

Parameter renormalization logs

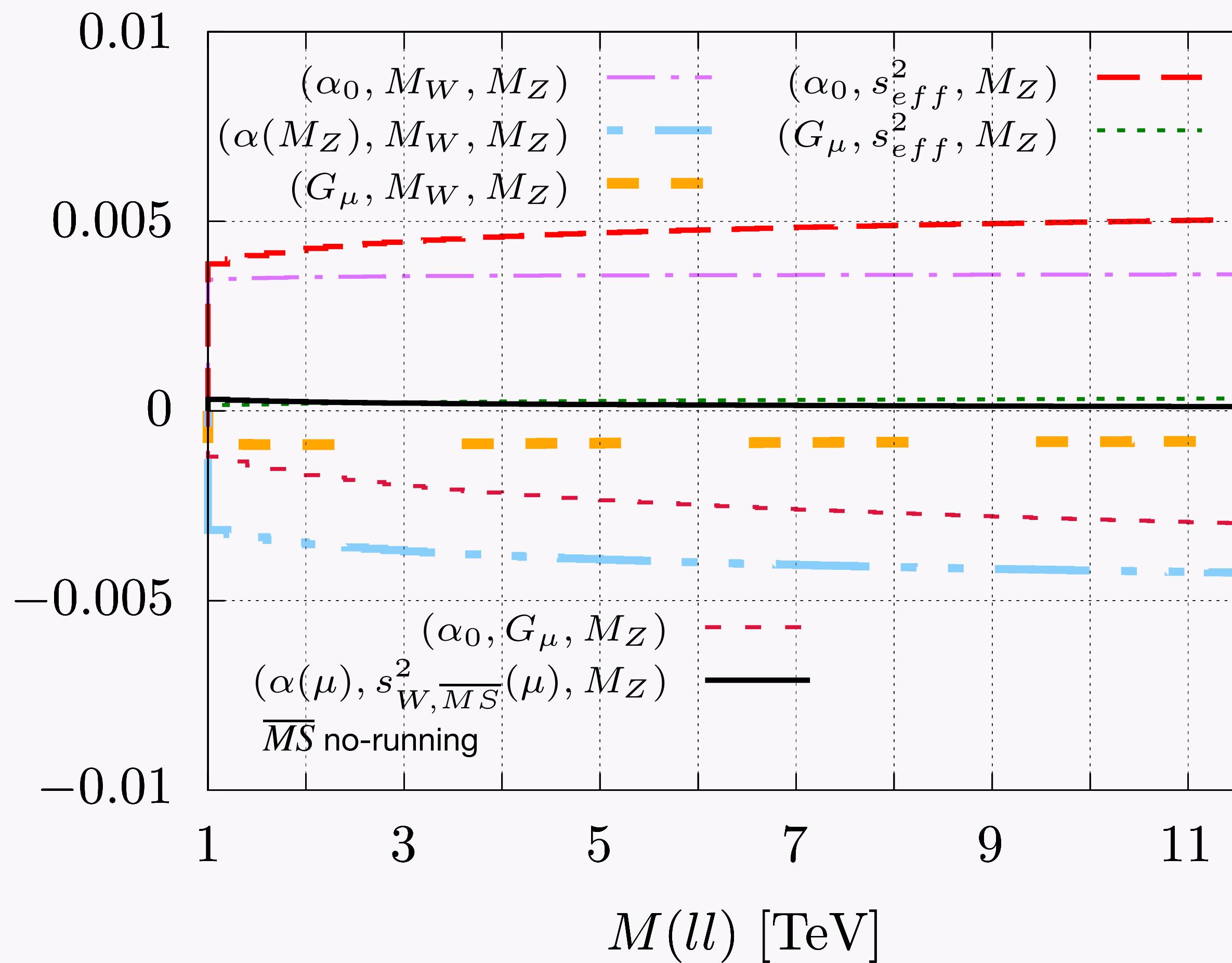
$$\frac{1}{\epsilon} - \ln \frac{r_{ct}^2}{\mu_{dim}^2} - \frac{1}{\epsilon} + \ln \frac{r_{bare}^2}{\mu_{dim}^2} = \ln \frac{r_{bare}^2}{r_{ct}^2} \sim \ln \frac{M_{ll}^2}{m^2}$$

Counterterms from parameter renorm.

Bare diagrams



Sudakov regime

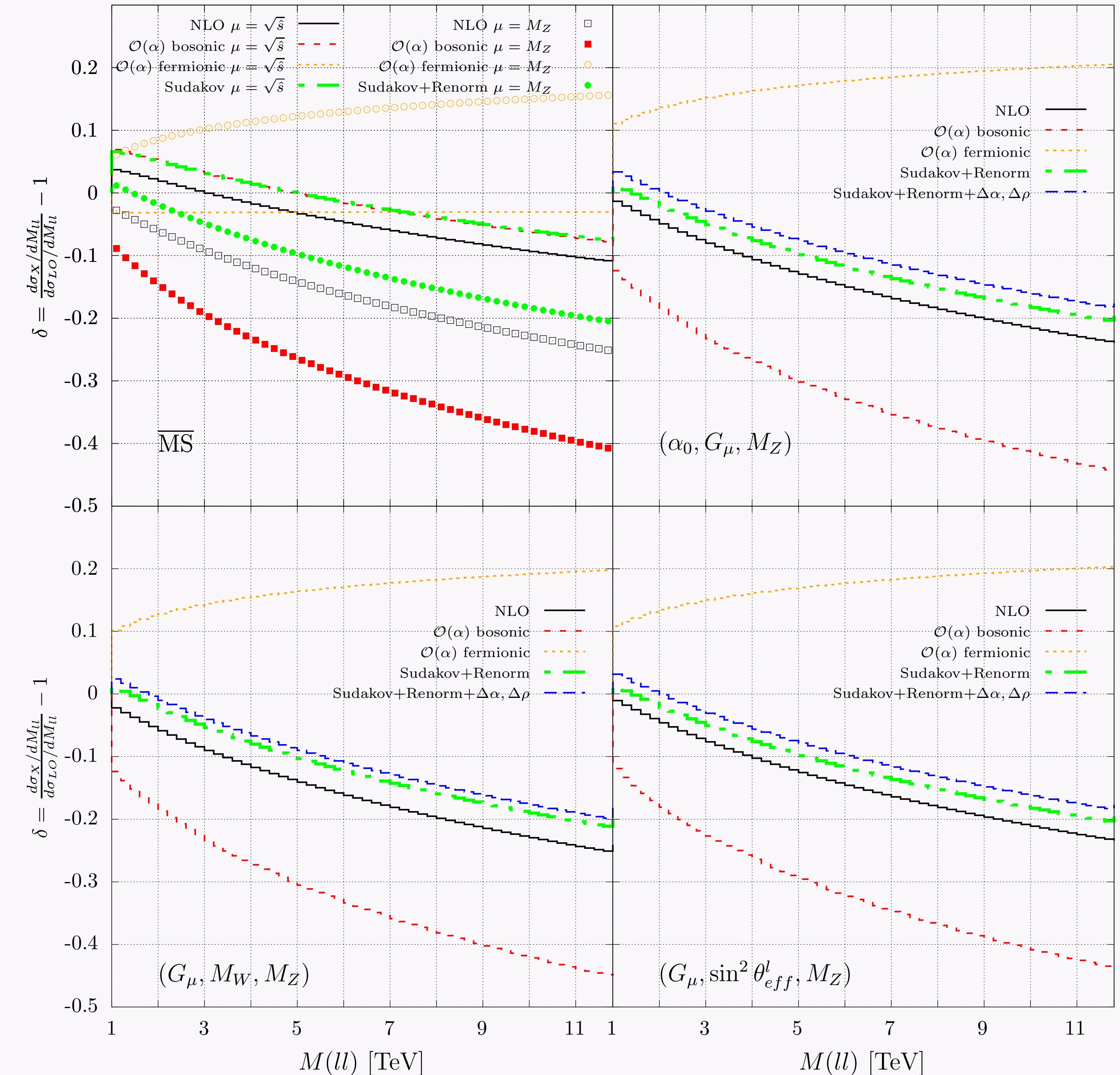


Cross section
at NLO+ho - (Sudakov + param.
renorm. logs)

$$\frac{\text{scheme } X}{(\alpha(M_Z^2), s_{eff}^2, M_Z)} - 1$$

Sudakov regime

True NLO - approx. $\sim 5\%$



NLO QCD effects

Code inputs and definitions

- EW scheme: (α, M_Z, M_W)
- CMS scheme
- $\sqrt{s} = 8 \text{ TeV}$
- PDF: MSTW2008nlo68cl (LHA 21100)
- factorisation scheme: \overline{MS}
- $\mu_F = \mu_R = M_{\ell\bar{\ell}}$

| Parameter | Value |
|--------------|----------------------------|
| α | $7.297353 \cdot 10^{-3}$ |
| G_μ | $1.166389 \cdot 10^{-5}$ |
| M_Z^* | 91.1876 |
| Γ_Z^* | 2.4952 |
| M_W^* | 80.385 |
| Γ_W^* | 2.085 |
| M_H | 125 |
| m_e | $0.51099907 \cdot 10^{-3}$ |
| m_μ | 0.1056583 |
| m_τ | 1.77705 |
| m_b | 4.7 |
| m_t | 173 |

| Parameter | Value |
|-----------|---------|
| m_u | 0.06983 |
| m_d | 0.06984 |
| m_s | 0.15 |
| m_c | 1.2 |

* Automatically translated to pole values

QED corrections

$$A_4 = 8/3 A_{FB}$$

| Code: | $89 < M_{\ell\bar{\ell}}[\text{GeV}] < 93$ | $60 < M_{\ell\bar{\ell}}[\text{GeV}] < 81$ | $81 < M_{\ell\bar{\ell}}[\text{GeV}] < 101$ | $101 < M_{\ell\bar{\ell}}[\text{GeV}] < 150$ |
|--|--|--|---|--|
| $8/3 \cdot [A_{FB}(\text{NLO QED ISR}) - A_{FB}(\text{LO})]/10^{-4}$ | | | | |
| MCSANC | 0.2(3) | -5(2) | 0.2(3) | 5(2) |
| WZGRAD2 | 0.2(5) | -5(3) | 0.3(5) | 6(4) |
| KKMC-hh | -1.0(6) | 0(1) | -0.5(5) | -8(2) |
| KKMC-hh (NISR) | -1(2) | 0(4) | 0(1) | 6(8) |
| RADY (CMS) | 0.16(4) | -4.05(3) | 0.12(3) | 4.90(3) |
| A. Huss | 0.17(1) | -4.07(1) | 0.11(1) | 4.94(4) |
| POWHEG _{ew} | 0.1(1) | -4.0(4) | 0.1(1) | 4.5(7) |
| $8/3 \cdot [A_{FB}(\text{NLO QED IFI}) - A_{FB}(\text{LO})]/10^{-4}$ | | | | |
| MCSANC | -2.8(5) | -34(2) | -4.0(4) | -60(3) |
| WZGRAD2 | -1.1(5) | -37(3) | -2.3(5) | -51(4) |
| KKMC-hh | -3.8(6) | -25(1) | -2.1(1) | -53(1) |
| KKMC-hh (NISR) | -3.1(6) | -17(1) | -3.2(5) | -60(3) |
| RADY (CMS) | -1.5(1) | -33.6(4) | -2.49(7) | -59.5(1) |
| A. Huss | -1.42(6) | -33.9(6) | -2.57(7) | -58.7(3) |
| POWHEG _{ew} | $\mu_F = M_{\ell\bar{\ell}\gamma}$ | -1.2(3) | -2.5(4) | -59(2) |
| | $\mu_F = M_{\ell\bar{\ell}}$ | -1.3(6) | -34(2) | -59(3) |

- POWHEG_{ew} $\mu_F(1) \implies M_{II}$ for real rad calculated with underlying Born momenta
- POWHEG_{ew} $\mu_F(2) \implies M_{II}$ for real rad calculated with radiative event momenta
- differences between $\mu_F(1)$ and (2) expected to decrease when including also QCD corrections

From last General Meeting
July 2024

Overall agreement of the codes

QED corrections

$$A_4 = 4 \langle \cos \vartheta \rangle$$

| Code: | $89 < M_{\ell\bar{\ell}}[\text{GeV}] < 93$ | $60 < M_{\ell\bar{\ell}}[\text{GeV}] < 81$ | $81 < M_{\ell\bar{\ell}}[\text{GeV}] < 101$ | $101 < M_{\ell\bar{\ell}}[\text{GeV}] < 150$ |
|--|--|--|---|--|
| $[A_4(\text{NLO QED ISR}) - A_4(\text{LO})]/10^{-4}$ | | | | |
| RADY (CMS) | 0.15(3) | -4.05(3) | 0.10(2) | 4.89(2) |
| A. Huss | 0.16(1) | -4.07(1) | 0.11(1) | 4.87(2) |
| POWHEG _{ew} | 0.07(9) | -4.0(3) | 0.10(7) | 4.8(4) |
| $[A_4(\text{NLO QED IFI}) - A_4(\text{LO})]/10^{-4}$ | | | | |
| RADY (CMS) | -1.7(1) | -42.3(4) | -2.97(6) | -71.6(2) |
| A. Huss | -1.68(6) | -42.4(6) | -3.05(8) | -71.2(3) |
| POWHEG _{ew} | $\mu_F = M_{\ell\bar{\ell}\gamma}$ | -1.5(5) | -70(1) | -3.0(4) |
| | $\mu_F = M_{\ell\bar{\ell}}$ | -1.5(5) | -43(1) | -3.0(4) |

- POWHEG_{ew} $\mu_F(1) \implies M_{II}$ for real rad calculated with underlying Born momenta
- POWHEG_{ew} $\mu_F(2) \implies M_{II}$ for real rad calculated with radiative event momenta
- at low and high M_{II} virtual QED boxes and I-F real radiation interference break factorization assumption for angular coefficients and the LO equality between the two A_4 def's
- differences between $\mu_F(1)$ and (2) expected to decrease when including also QCD corrections

From last General Meeting
July 2024

Overall agreement of the codes

NLO QED and QCD corrections - how to combine them?

1: additive method

$$\sigma_{\text{QED+QCD}} = \sigma_{\text{NLOQCD}} + \sigma_{\text{NLOQED}} - \sigma_{\text{LO}} = \sigma_{\text{LO}}(1 + \delta_{\text{QCD}} + \delta_{\text{QED}})$$

$$A_{\text{FB}}^{\text{QED+QCD}} = \frac{(\sigma_F - \sigma_B)^{\text{NLOQED}} + (\sigma_F - \sigma_B)^{\text{NLOQCD}} - (\sigma_F - \sigma_B)^{\text{LO}}}{(\sigma_F + \sigma_B)^{\text{NLOQED}} + (\sigma_F + \sigma_B)^{\text{NLOQCD}} - (\sigma_F + \sigma_B)^{\text{LO}}}$$

Equivalent at NLO

2: factorized method

difference due to $\mathcal{O}(\alpha\alpha_S)$ terms

$$\sigma_{\text{QED+QCD}} = \sigma_{\text{LO}}(1 + \delta_{\text{QCD}})(1 + \delta_{\text{QED}})$$

$A_{\text{FB}}^{\text{QED+QCD}}$ to be computed bin-by-bin from $d\sigma/d\cos\theta$ - has more complicated expression

starting point:
$$\frac{d\sigma_{\text{QED+QCD}}}{d\cos\theta} = \frac{d\sigma_{\text{NLOQCD}}}{d\cos\theta} \left(\frac{d\sigma_{\text{NLOQED}}}{d\cos\theta} / \frac{d\sigma_{\text{LO}}}{d\cos\theta} \right)$$

QED and QCD corrections

Numbers obtained with additive method

$$\mu_R = \mu_F = M_{\ell\bar{\ell}}$$

$$A_4 = 8/3 \cdot A_{FB}$$

| 89 – 93 GeV | 60 – 81 GeV | 81 – 101 GeV | 101 – 150 GeV |
|--|-------------|--------------|---------------|
| $8/3 \cdot A_{FB}(\text{LO})$ | | | |
| 0.12414(1) | -0.54107(2) | 0.119512(8) | 0.60292(4) |
| $8/3 \cdot [A_{FB}(\text{NLO QCD}) - A_{FB}(\text{LO})]/10^{-2}$ | | | |
| -0.38(4) | 2.10(2) | -0.36(3) | -1.86(3) |
| $8/3 \cdot [A_{FB}(\text{NLO QCD + QED}) - A_{FB}(\text{LO})]/10^{-2}$ | | | |
| -0.067(9) | 27.24(1) | -0.152(6) | -1.79(4) |

QED : $[A_{FB}(\text{NLO QED}) - A_{FB}(\text{LO})]/10^{-2}$

NLO QCD + QED : $[A_{FB}(\text{NLO QCD + QED}) - A_{FB}(\text{NLO QCD})]/10^{-2}$

| | $89 < M_{\ell\bar{\ell}}[\text{GeV}] < 93$ | $60 < M_{\ell\bar{\ell}}[\text{GeV}] < 81$ | $81 < M_{\ell\bar{\ell}}[\text{GeV}] < 101$ | $101 < M_{\ell\bar{\ell}}[\text{GeV}] < 150$ |
|-------------------|--|--|---|--|
| QED | 0.423(2) | 28.296(3) | 0.262(2) | 0.27(1) |
| NLO QCD + QED | 0.32(1) | 25.14(3) | 0.209(9) | 0.07(5) |
| QED FSR | 0.439(2) | 28.604(7) | 0.289(1) | 0.87(1) |
| NLO QCD + QED FSR | 0.33(1) | 25.42(3) | 0.231(9) | 0.58(5) |
| QED ISR | 0.001(1) | -0.040(4) | 0.001(1) | 0.045(6) |
| NLO QCD + QED ISR | 0.002(8) | -0.04(2) | 0.002(7) | 0.05(3) |
| QED IFI | -0.013(6) | -0.34(2) | -0.026(7) | -0.59(3) |
| NLO QCD + QED IFI | -0.01(1) | -0.30(4) | -0.02(1) | -0.52(5) |

QED and QCD corrections

Numbers obtained with additive method

$$\mu_R = \mu_F = M_{\ell\bar{\ell}}$$

$$A_4 = 4 < \cos \theta >$$

| 89 – 93 GeV | 60 – 81 GeV | 81 – 101 GeV | 101 – 150 GeV |
|--|--------------|--------------|---------------|
| $A_4 = 4 < \cos \theta > (\text{LO})$ | | | |
| 0.031034(1) | -0.135266(3) | 0.029878(1) | 0.150727(2) |
| $[A_4(\text{NLO QCD}) - A_4(\text{LO})]/10^{-2}$ | | | |
| -0.384(4) | 2.1(1) | -0.363(3) | -1.89(1) |
| $[A_4(\text{NLO QCD + QED}) - A_4(\text{LO})]/10^{-2}$ | | | |
| -0.071(7) | 27.154(6) | -0.160(5) | -1.94(3) |

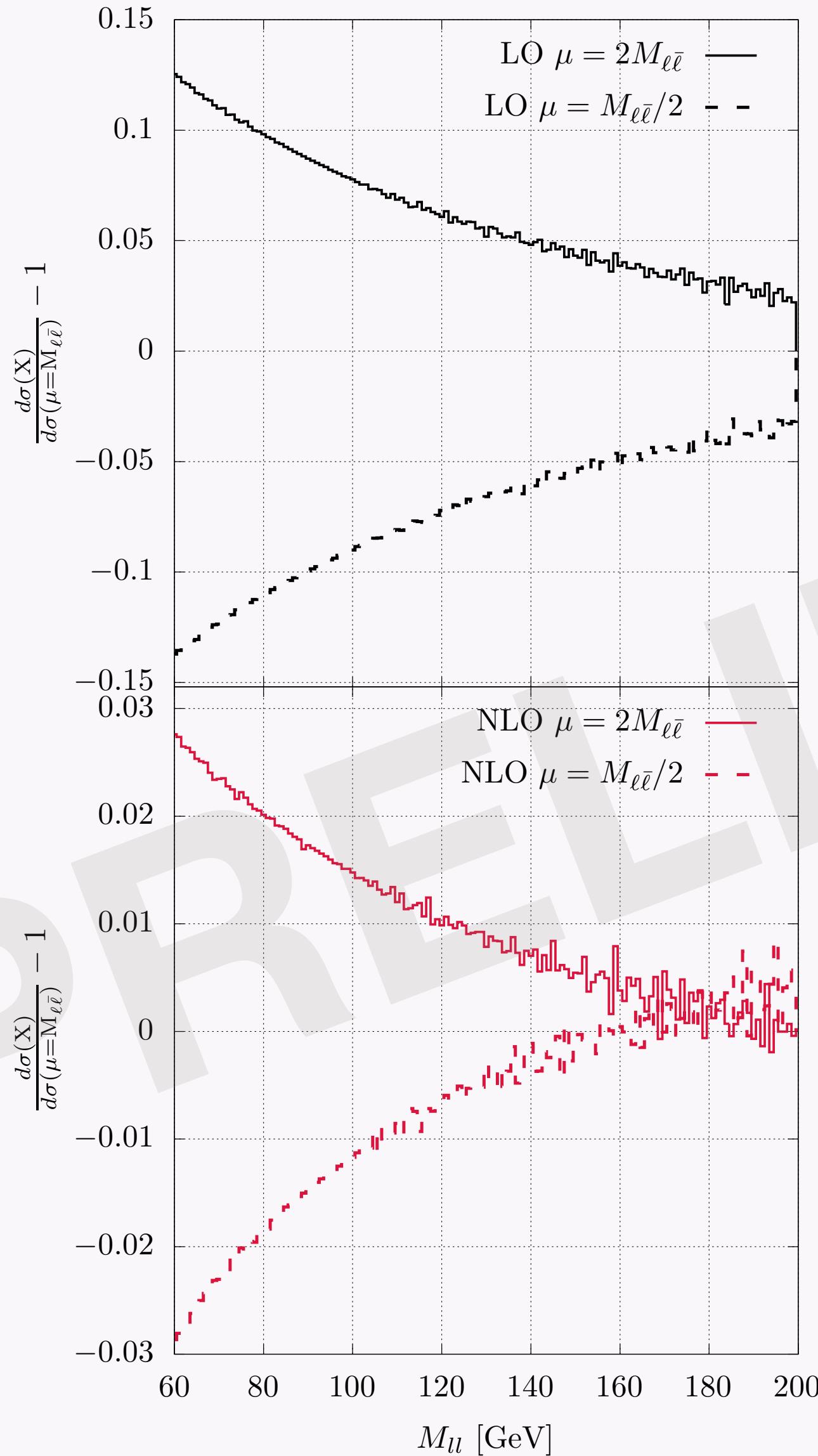
QED : $[A_4(\text{NLO QED}) - A_4(\text{LO})]/10^{-2}$

NLO QCD + QED : $[A_4(\text{NLO QCD + QED}) - A_4(\text{NLO QCD})]/10^{-2}$

| | 89 < $M_{\ell\bar{\ell}}[\text{GeV}] < 93$ | 60 < $M_{\ell\bar{\ell}}[\text{GeV}] < 81$ | 81 < $M_{\ell\bar{\ell}}[\text{GeV}] < 101$ | 101 < $M_{\ell\bar{\ell}}[\text{GeV}] < 150$ |
|-------------------|--|--|---|--|
| QED | 0.420(2) | 28.196(3) | 0.256(1) | 0.131(8) |
| NLO QCD + QED | 0.313(9) | 25.05(1) | 0.203(7) | -0.05(4) |
| QED FSR | 0.438(1) | 28.553(6) | 0.288(1) | 0.862(9) |
| NLO QCD + QED FSR | 0.327(9) | 25.37(2) | 0.229(7) | 0.57(4) |
| QED ISR | 0.0007(9) | -0.040(3) | 0.0010(7) | 0.048(4) |
| NLO QCD + QED ISR | 0.002(6) | -0.04(1) | 0.002(4) | 0.05(3) |
| QED IFI | -0.015(5) | -0.43(1) | -0.030(4) | -0.72(3) |
| NLO QCD + QED IFI | -0.013(9) | -0.37(2) | -0.026(7) | -0.62(5) |

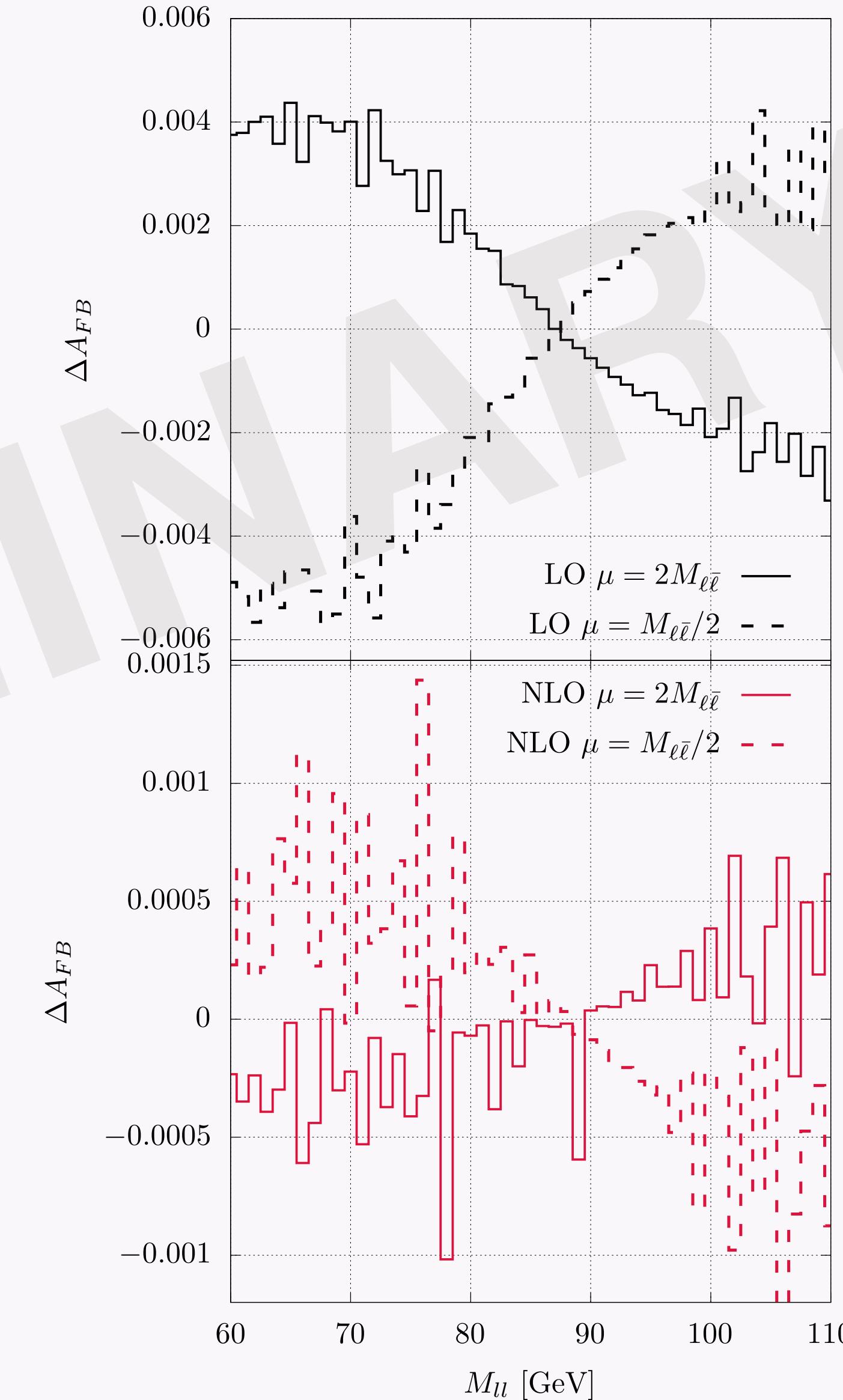
Scale variations at NLO QCD

Cross section



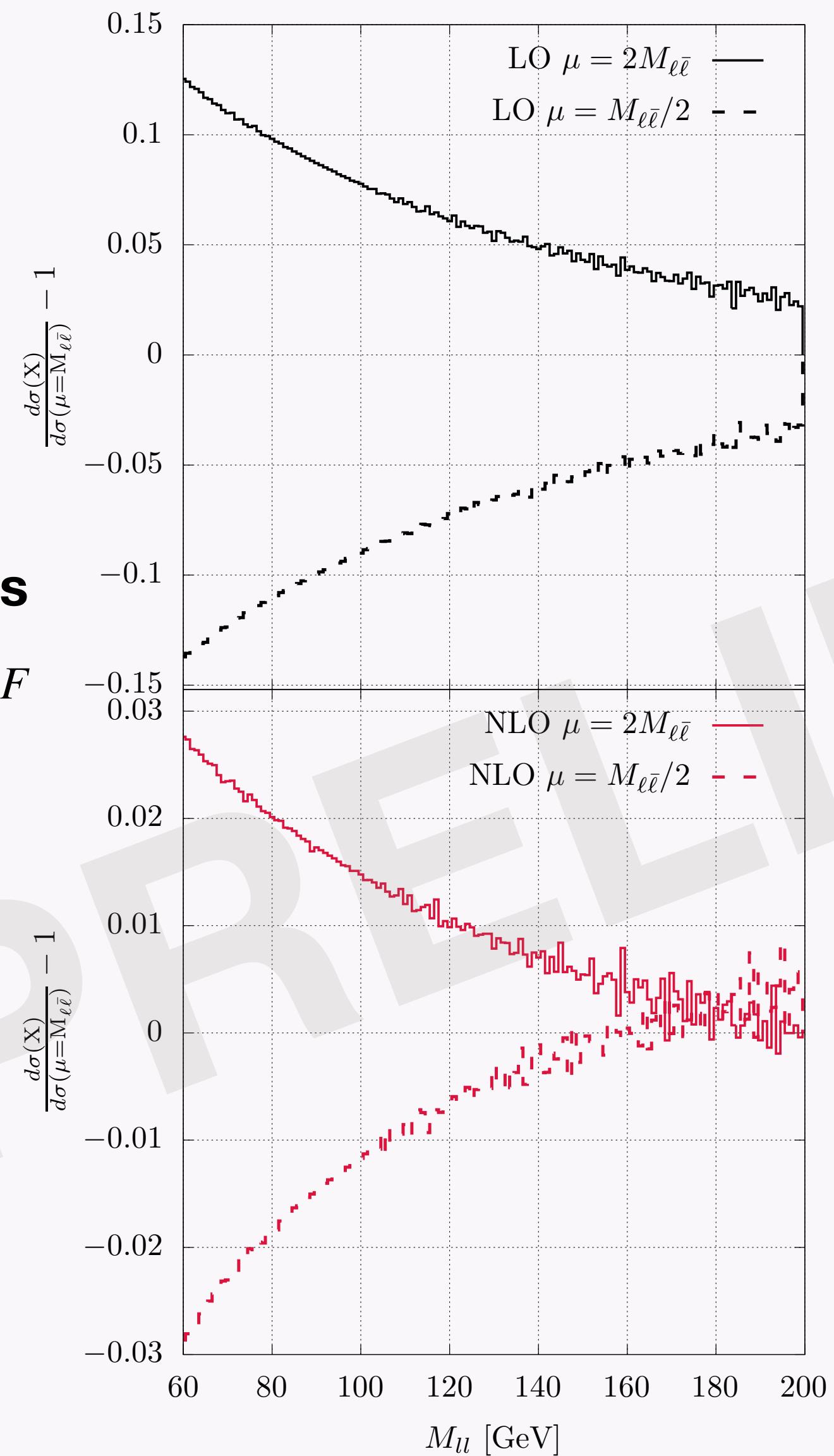
Asymmetry

Still low statistics

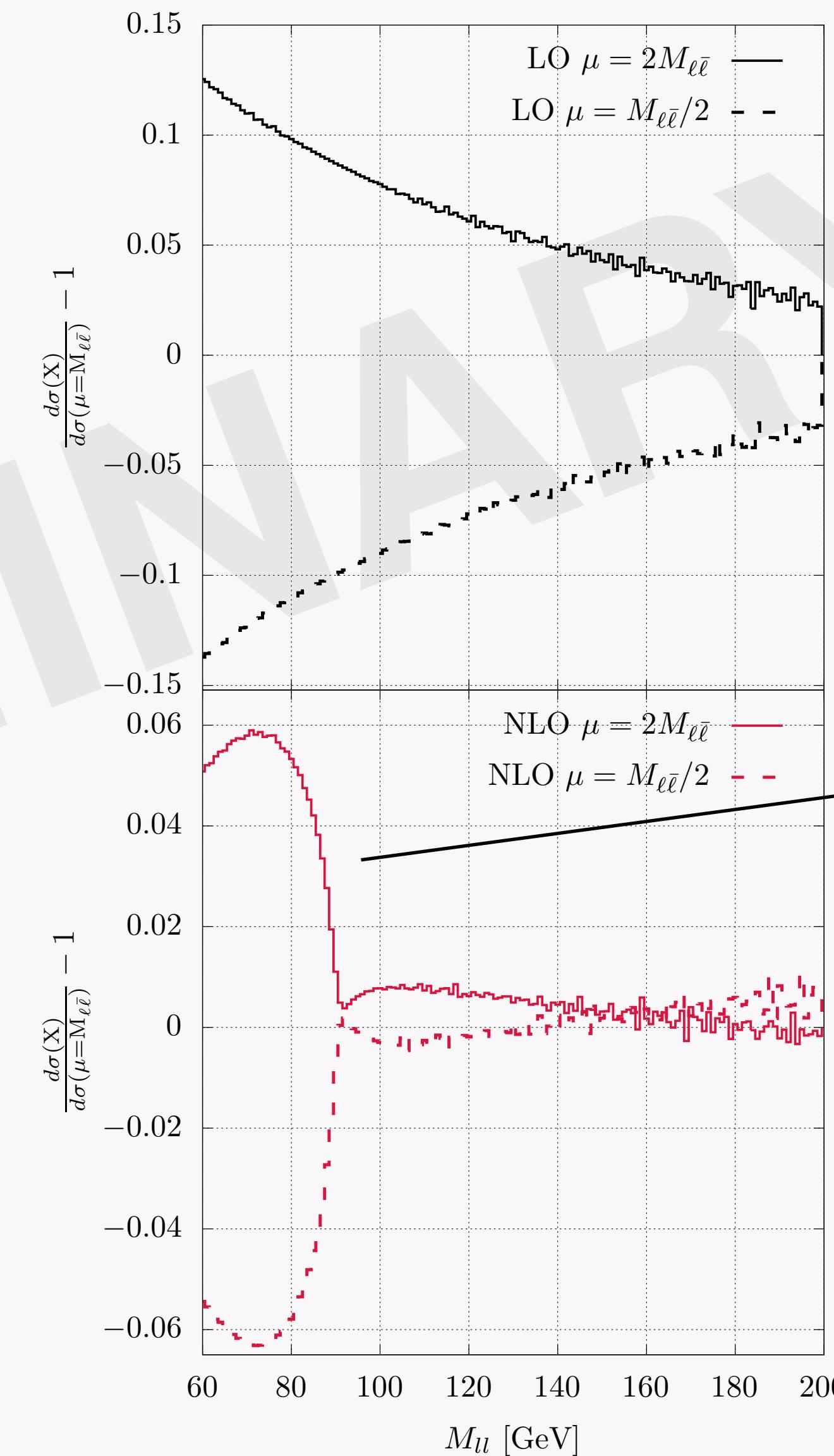


Scale variations at NLO QCD+QED

Cross section
factorised method



Hp. QED corrections
do not depend on μ_F



Cross section
additive method

This bulk is due to
NLO QED corrections:
multi-photon emission
would reduce the
effect

Summary and roadmap

- Systematic study of different renormalization and input parameter schemes, as well as tuned realisations of some schemes, can be useful in our discussion on theoretical uncertainties
- Comparison of different schemes at high energy and effects of the Sudakov approximation for NC DY
- NLO QCD corrections with POWHEG and their combination with QED effects studied for A_4 and A_{FB} together with preliminary exploration of scale variations

To be done:

- plots for asymmetry with QCD+QED combination (higher stat will probably be needed)
- quantification of the impact of using $\mu_F = M_{\ell\bar{\ell}}$ or $\mu_F = M_{\ell\bar{\ell}\gamma}$ at NLO QCD+QED to follow-up previous discussions