

Neutron irradiation experiments on HTS

Raphael Unterrainer¹, Davide Gambino², Florian Semper¹, Alexander Bodenseher¹, Danielle Torsello^{3,4}, Francesco Laviano^{3,4}, David X. Fischer⁵ <u>Michael Eisterer</u>¹

¹Atominstitut, TU Wien, Vienna, Austria
 ²IFM, Linköping University, Linköping, Sweden
 ³DiSAT, Politecnico di Torino, Torino, Piemonte, Italy
 ⁴INFN, Sezione di Torino, Torino, Piemonte, Italy
 ⁵PSFC, MIT, Cambridge, Massachusetts, United States

RADSUM, CERN, January 16th 2025

Collaborations and Funding



This work has been carried out within the framework of the EUROfusion Consortium, funded by the European Union via the Euratom Research and Training Programme (Grant Agreement No 101052200 — EUROfusion). Views and opinions expressed are however those of the author(s) only and do not necessarily reflect those of the European Union or the European Commission. Neither the European Union nor the European Commission can be held responsible for them.



Funded by the European Union















Irradiation Damage



- Compact fusion devices high fields for confinement
 - Currently REBCO coated conductors most promising
 - High quality, long lengths (800+ m)
 - Change of properties under irradiation conditions "well" known
 - Monotonic decrease of T_c

Irradiation Damage



15 T 30 K

- Non-monotonic behavior of critical current as function of the neutron fluence
- Increase at low fluence
- Decrease at higher fluences



Irradiation Damage



Compact devices

- Much higher neutron flux at the magnets
- Maximum lifte-time fluence of approx. $3-3.3 \cdot 10^{22} \, m^{-2}$



Irradiation Methods







[1] SuperPower®, superpower-inc.com

Thorough pre-characterization!							
ì	 HTS thicknes 	s: ~1 µm					
	 substrate: 	Hastelloy					
	• el. stabilized:	Cu					
	chem. stabiliz	zed: 1 µm Ag					

Supplier	Туре	REBCO	APCs	method	nomenclature
SuperPower	SCS4050 2009	GdBCO	None	MOCVD	SP SCS09
SuperPower	SCS4050 2013	(Y,Gd)BCO	BaZrO ₃	MOCVD	SP SCS13
SuNAM	HCN04150	GdBCO	None	RCE-DR	SuNAM HCN











Irradiation Environments









Neutron Irradiation – Shielded

TRIGA MARK II at TU Wien

- Irradiation in the central irradiation facility
- Fast / thermal neutron flux $3.2 / 4 \times 10^{16} \, m^{-2} \, s^{-1}$
- Irradiation with and without thermal (< 0.55 eV) neutrons
- Sample identifiers denoted with "S"



< 70 °C at sample



Neutron Irradiation – Unshielded

TRIGA MARK II at TU Wien

- Irradiation in the central irradiation facility
- Fast / thermal neutron flux $3.2 / 4 \times 10^{16} \, m^{-2} \, s^{-1}$
- Irradiation with and without thermal (< 0.55 eV) neutrons
- Sample identifiers denoted with "U"







p+ Irradiation - Bridged

General Ionix 1.7 MV tandem accelerator

- Irradiation with 1.2 MeV protons
- Room temperature irradiation
- Bridged samples 0.2 mm width
- Samples pre-characterized in Vienna
- On-sample temperature control to monitor beam heating





Defect Formation



Fast Neutron Irradiation



left – TEM picture of neutron induced defects **right** – FFT of selected regions ¹

[1] with friendly permission by Yatir Linden, *Analysing neutron radiation damage in YBa2Cu3O7–x high-temperature superconductor tapes*, <u>https://doi.org/10.1111/jmi.13078</u> Department of Materials, University of Oxford, Oxford, UK

Fast Neutron Irradiation



left – TEM picture of neutron induced defects **right** – FFT of selected regions ¹

- 1. Undisturbed GdBCO
- 2. Crystalline BZO rod
- 3. Amorphous cascade



Cascades Enhance Pinning

[1] with friendly permission by Yatir Linden, *Analysing neutron radiation damage in YBa2Cu3O7–x high-temperature superconductor tapes*, <u>https://doi.org/10.1111/jmi.13078</u> Department of Materials, University of Oxford, Oxford, UK





Fast Neutron Irradiation



left – TEM picture of neutron induced defects **right** – FFT of selected regions ¹

- 1. Undisturbed GdBCO
- 2. Crystalline BZO rod

۲

3. Amorphous cascade



- $3.3 \times 10^{19} 5 \times 10^{22}$ cascades per 10^{22}
 - ~ 0.01 % reduction of superconducting cross-section

What drives the degradation? → Must be small (invisible) defects

[1] with friendly permission by Yatir Linden, *Analysing neutron radiation damage in YBa2Cu3O7–x high-temperature superconductor tapes*, <u>https://doi.org/10.1111/jmi.13078</u> Department of Materials, University of Oxford, Oxford, UK

Thermal Neutron Irradiation





Thermal neutrons excite Gd — Recoil of 29 – 32 e

Recoil of 29 – 32 eV gamma emission displaces the nucleus

- Very high defect densities achievable
- Add to fast neutron induced defects

Thermal Neutron Irradiation



- Position enables introduction of many defects close to the planes
- Defects are small in comparison to coll. cascades
- Defects may be modelled with MDS
- 3 energies close to experimental value simulated (30, 35, 40 eV)



430 simulation runs per energy

- Most defects are oxygen vacancies
- Gd returns / stays in lattice position
- Different defects originating from Gd PKA (primary knock on atom)
- Defect distribution changes with energy
- 1-2 defects per incident particle





- Most defects are oxygen vacancies
- Gd returns / stays in lattice position
- Different defects originating from Gd PKA (primary knock on atom)
- Defect distribution changes with energy
- 1-2 defects per incident particle





- Most defects are oxygen vacancies
- Gd returns / stays in lattice position
- Different defects originating from Gd PKA (primary knock on atom)
- Defect distribution changes with energy
- 1-2 defects per incident particle





- Most defects are oxygen vacancies
- Gd returns / stays in lattice position
- Different defects originating from Gd PKA (primary knock on atom)
- Defect distribution changes with energy
- 1-2 defects per incident particle



- Mainly point-like defects
- Small defects form clusters
- Up to 1 nm in size
- Slightly improves pinning behavior



1.2 MeV p⁺ Irradiation



- SRIM/TRIM
- Most defects are oxygen displacements (low binding energy)
- Little is known about actual defects and recombination
- Large defects are possible but rare
- Most defects are point-like or small clusters like with thermal neutrons



Results





 $T_{\rm c}$ degrades ~14 x faster due to Gd-point defects

(MARS calculations: dpa only higher by about 50%)



- Maximum occurs at much lower neutron fluences
- *J*_c at maximum is smaller
- Degradation is much faster

Fluence is not a good measure for the disorder!

\square Influence of thermal neutrons - J_c



- $J_{\rm c}$ maximum shifted to lower $T_{\rm c}$
- Degradation with similar slope
 - Accumulation of similar defects?
- $T_{\rm c}$ is efficient disorder parameter (decrease of superfluid density)

•
$$D = T_{\rm c} - T_{\rm c}^{\rm p}$$



- Different defect densities
- Parallel degrading branch
- Specific defects origin of degradation?
- Accumulation in all irradiation techniques?

Focus on degrading branch





Sample irradiated with 1.2 MeV proton at room temperature



- Sample irradiated with 1.2 MeV proton at room temperature
- --- Shielded sample with APCs



- Different defect densities
- Parallel degrading branch
- Specific defects origin of degradation?
- Accumulation in all irradiation techniques?



Fluence dependence of critical current



Sketch of resistive transition



Degrading scattering

- T_c decreases
- J_d decreases
- *n*-value decreases

Sketch of resistive transition



Degrading scattering

- T_c decreases
- J_d decreases
- *n*-value decreases

Enhancing pinning

- $J_c = \eta J_d$
- η increases
 (pinning efficiency)







Degradation function

$$\frac{I_{\rm c}}{I_{\rm c,p}} = \sqrt{\frac{(\alpha_{\rm p}+1)t_{\rm c}^3}{\alpha_{\rm p}\left(1-K_{\rm p}(1-t_{\rm c})\right)+t_{\rm c}}} \left(\frac{E_{\rm crit}}{\sqrt{\phi_0 B}\nu_0}\right)^{\frac{1}{n}-\frac{1}{n_{\rm p}}} \frac{\eta_{\rm max}}{\eta^0} \tanh\left(\frac{D}{D_{\eta_{\rm max}}}(\pi-A)+A\right) \qquad t_{\rm c} = \frac{T_{\rm c}}{T_{\rm c,p}}$$



Attempt frequency $v_0 = 2.5 \times 10^7 \text{ Hz}$

Electric. field criterion $E_c = 1 \ \mu V \ cm$ $\alpha = \frac{\xi_0}{l} = 3$ $K_\rho = \frac{T_{c,p}}{\rho_{n,p}} \frac{\partial \rho_n}{\partial T_c} = -16.5$ Experimetal $n = n_0 - \Delta nD$

- Universal
- All parameters fairly easily accessible
- Curve hard to measure directly due to pinning

Degradation function

$$\frac{I_{\rm c}}{I_{\rm c,p}} = \sqrt{\frac{(\alpha_{\rm p}+1)t_{\rm c}^3}{\alpha_{\rm p}\left(1-K_{\rm p}(1-t_{\rm c})\right)+t_{\rm c}}} \left(\frac{E_{\rm crit}}{\sqrt{\phi_0 B}\nu_0}\right)^{\frac{1}{n}-\frac{1}{n_{\rm p}}} \frac{\eta_{\rm max}}{\eta^0} \tanh\left(\frac{D}{D_{\eta_{\rm max}}}(\pi-A)+A\right) \qquad t_{\rm c} = \frac{T_{\rm c}}{T_{\rm c,p}}$$



- $J \propto E^n$
- Linear fit to experimental data $n = n_0 \Delta nD$
- Smoothening of scattered data
- Envelope of expectations for large and small defects

Pinning enhancement

$$\frac{I_{\rm c}}{I_{\rm c,p}} = \sqrt{\frac{(\alpha_{\rm p}+1)t_{\rm c}^3}{\alpha_{\rm p}\left(1-K_{\rm p}(1-t_{\rm c})\right)+t_{\rm c}}} \left(\frac{E_{\rm crit}}{\sqrt{\phi_0 B}\nu_0}\right)^{\frac{1}{n}-\frac{1}{n_{\rm p}}} \frac{\eta_{\rm max}}{\eta^0} \tanh\left(\frac{D}{D_{\eta_{\rm max}}}(\pi-A)+A\right) \qquad t_{\rm c} = \frac{T_{\rm c}}{T_{\rm c,p}}$$



$$A = \tanh^{-1} \frac{\eta^0}{\eta_{\max}}$$

- tanh was chosen for the good • correspondence to the data
- Pinning efficiency can not increase ٠ above ~ 0.3







$$\frac{f_{c}}{f_{c,p}} = \sqrt{\frac{(\alpha_{p}+1)t_{c}^{3}}{\alpha_{p}\left(1-K_{\rho}(1-t_{c})\right)+t_{c}}} \left(\frac{E_{crit}}{\sqrt{\phi_{0}B}\nu_{0}}\right)^{\frac{1}{n}-\frac{1}{n_{p}}} \frac{\eta_{max}}{\eta^{0}} \tanh\left(\frac{D}{D_{\eta_{max}}}(\pi-A)+A\right)$$

Degradation is nearly universal and driven by the loss of superfluid density.

- Physical origin: Pair breaking scattering on small defects
- Relevant parameter: Scattering rate τ^{-1} (numerical modelling?)
- Experimental manifestation: Decrease of transition temperature (D), normal state resistivity (ρ_n)

Increase in critical current results from improved pinning

- Non-universal, depending on initial defect structure and type of radiation.
- Physical origin: Flux pinning by large defects
- Relevant parameter: pinning efficiency η (TDGL modelling?)
- Experimental benchmarking: Saturation value of η (suitable proxies for neutrons?)



Annealing



Thermal stability of small vs large defects



- $T_{\rm c}$ regenerates "linearly" with $T_{\rm a}$
- All neutron and proton irradiated samples anneal to same point
- Annealing defects have same/similar distribution and activation barrier
- n_{therm}, n_{fast} & p⁺ irradiated samples

Samples annealed in pure O_2 atmosphere at 1 atm

Thermal stability of small vs large defects



- $T_{\rm c}$ regenerates "linearly" with $T_{\rm a}$
- All neutron and proton irradiated samples anneal to same point
- Annealing defects have same/similar distribution and activation barrier
- n_{therm}, n_{fast} & p⁺ irradiated samples

Normalizing $\Delta T_{c}(T_{a})$ to $\Delta T_{c}(T_{a} = 25 \text{ °C})$

Thermal stability of small vs large defects



- $T_{\rm c}$ regenerates "linearly" with $T_{\rm a}$
- All neutron and proton irradiated samples anneal to same point
- Annealing defects have same/similar distribution and activation barrier
- n_{therm}, n_{fast} & p⁺irradiated samples

Normalizing $\Delta T_{c}(T_{a})$ to $\Delta T_{c}(T_{a} = 25 \text{ °C})$

 \square Annealing: I_c



- Recovery is non-monotonous, although with a linear trend.
- Degraded samples was recovered above its initial value.
- Optimum annealing protocol to be derived.

Contribution to the workflow

- Fast neutron irradiation is a benchmarking experiment
 - Among the closest available neutron sprectrum
- Checks the reability of Tasks 2) Defect Structure and 3) Flux Pinning and Superfluid Density.
 - > Task 2) provides the input parameters (τ^{-1} , defect density and size) for Task 3)
 - > Task 3) predicts changes of I_c and T_c .
- Comparison with experimental values.