Beyond the Standard Model

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Plan of the Lectures

- The Standard Model and issues beyond it
- Origin of particle masses: Higgs boson or?
- Supersymmetry
- Searches for supersymmetry: LHC & dark matter
- Extra dimensions and string theory

Beyond the Standard Model for Montañeros

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http://arXiv.org/pdf/09

9.pdf

Summary of the Standard Model

Particles and SU(3) \times SU(2) \times U(1) quantum numbers:

L_L E_R	$\left(\begin{array}{c}\nu_{e}\\e^{-}\end{array}\right)_{L}, \left(\begin{array}{c}\nu_{\mu}\\\mu^{-}\end{array}\right)_{L}, \left(\begin{array}{c}\nu_{\tau}\\\tau^{-}\end{array}\right)_{L}\\e_{R}^{-}, \mu_{R}^{-}, \tau_{R}^{-}\end{array}\right)_{L}$	(1,2 ,-1) (1,1 ,-2)
Q_L U_R D_R	$ \begin{pmatrix} u \\ d \end{pmatrix}_{L}, \begin{pmatrix} c \\ s \end{pmatrix}_{L}, \begin{pmatrix} t \\ b \end{pmatrix}_{L} \\ u_{R}, c_{R}, t_{R} \\ d_{R}, s_{R}, b_{R} $	(3,2 ,+1/3) (3,1 ,+4/3) $(\mathbf{3,1,-2/3})$

Lagrangian: $\mathcal{L} = -\frac{1}{\Lambda} F^a_{\mu\nu} F^{a\ \mu\nu}$ + $i\bar{\psi}$ /D ψ + h.c. + $\psi_i y_{ij} \psi_j \phi + h.c.$ + $|D_{\mu}\phi|^2 - V(\phi)$

Yukawa interactions Higgs potential

matter fermions

gauge interactions

Gauge Interactions of the Standard Model

- Three separate gauge group factors:
 - SU(3) × SU(2) × U(1)
 - Strong × electroweak
- Three different gauge couplings:
 - $-g_3, g_2, g$
- Mixing between the SU(2) and U(1) factors:

 $\begin{pmatrix} Z^{\mu} \\ A^{\mu} \end{pmatrix} = \begin{pmatrix} \cos(\theta_W) & \sin(\theta_W) \\ -\sin(\theta_W) & \cos(\theta_W) \end{pmatrix} \begin{pmatrix} W_3^{\mu} \\ B^{\mu} \end{pmatrix} \quad \sin^2(\theta_W) = \frac{g'^2}{g'^2 + g^2}$

• Experimental value: $\sin^2 \theta_{\rm W} = 0.23120 \pm 0.00015$

Important clue for Grand Unification

Weak Interactions

• Interactions of lepton doublets: *I*

$$L = \left(\begin{array}{c} \nu_e \\ e^- \end{array}\right)_L$$

• Charged-current interactions:

$$\mathcal{L}_{cc} = \frac{-g}{\sqrt{2}} \sum_{\alpha=e,\mu,\tau} \nu_{L_{\alpha}} \gamma_{\mu} l_{L_{\alpha}} W^{\mu} + h.c_{\alpha}$$

• Neutral-current interactions:

$$\mathcal{L}_{nc} = \frac{-g}{2\cos\theta_W} \sum_{\alpha=e,\mu,\tau} \nu_{L_\alpha} \gamma_\mu l_{L_\alpha} Z^\mu + h.c$$

• Effective four-fermion interaction:

$$G_{\rm F}/\sqrt{2} = g^2/8m_{\rm W}^2 \quad \mathcal{L} = \frac{G_F}{\sqrt{2}}\bar{\nu}_{\mu}\gamma^{\alpha}(1-\gamma^5)\mu\bar{e}\gamma_{\alpha}(1-\gamma^5)\nu_e$$

Status of the Standard Model

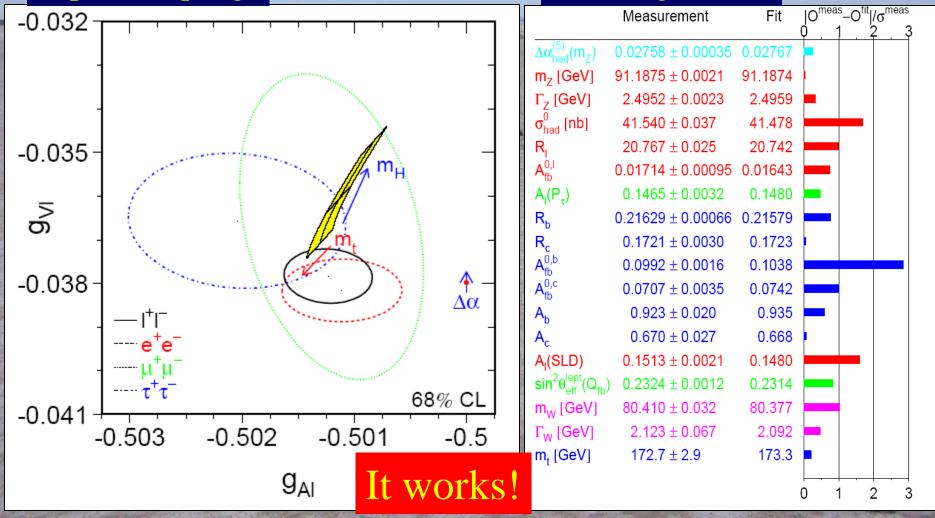
- Perfect agreement with all *confirmed* accelerator data
- Consistency with precision electroweak data (LEP et al) *only if there is a Higgs boson*
- Agreement seems to require a relatively light Higgs boson weighing < ~ 180 GeV
 - Raises many unanswered questions:

mass? flavour? unification?

Precision Tests of the Standard Model

Lepton couplings

Pulls in global fit



Parameters of the Standard Model

• Gauge sector:

- -3 gauge couplings: g_3 , g_2 , g
- 1 strong CP-violating phase
- Yukawa interactions:
 - 3 charge-lepton masses
 - 6 quark masses
 - 4 CKM angles and phase
- Higgs sector:
 - -2 parameters: μ , λ
- Total: 19 parameters

Unification?

Flavour?



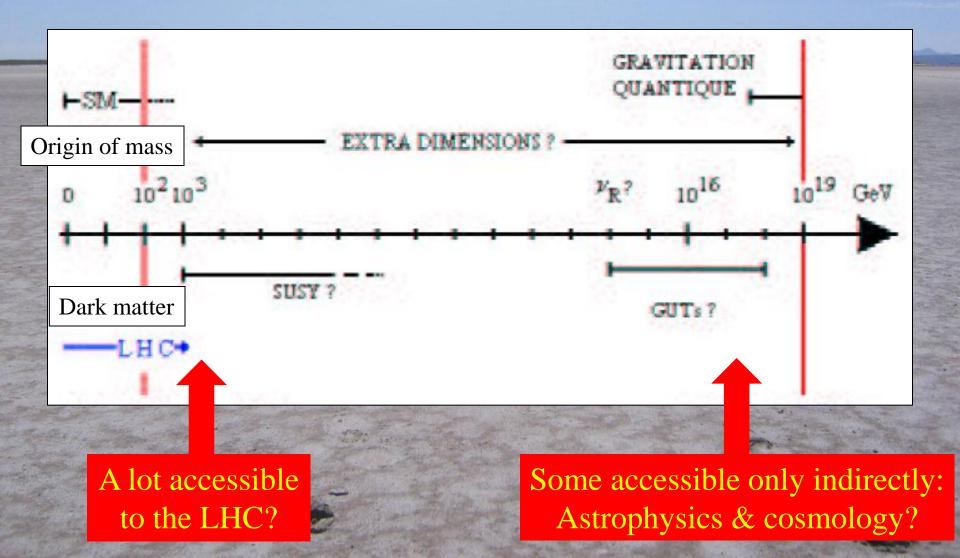
Open Questions beyond the **Standard Model**

- What is the origin of particle masses? due to a Higgs boson? + other physics? solution at energy < 1 TeV (1000 GeV)
- Why so many types of matter particles? matter-antimatter difference?
- Unification of the fundamental forces? at very high energy ~ 10^{16} GeV? probe directly via neutrino physics, indirectly via masses, couplings
 - Quantum theory of gravity? Supersymmetry (super)string theory: extra space-time dimensions?

Supersymmetry

Supersymmetry

At what Energy is the New Physics?



Why do Things Weigh?

Newton:

Weight proportional to Mass Einstein:

Energy related to Mass Neither explained origin of Mass

Where do the masses come from?

Are masses due to Higgs boson? (the physicists' Holy Grail)



The Seminal Papers

BROKEN SYMMETRY AND THE MASS OF GAUGE VECTOR MESONS*

F. Englert and R. Brout

Faculté des Sciences, Université Libre de Bruxelles, Bruxelles, Belgium (Received 26 June 1964)

BROKEN SYMMETRIES, MASSLESS PARTICLES AND GAUGE FIELDS

P.W. HIGGS

Tail Institute of Mathematical Physics, University of Edunburgh, Scotland

Received 27 July 1964

VOLUME 13, NUMBER 16

PHYSICAL REVIEW LETTERS

19 October 1964

BROKEN SYMMETRIES AND THE MASSES OF GAUGE BOSONS

Peter W. Higgs

Tait Institute of Mathematical Physics, University of Edinburgh, Edinburgh, Scotland (Received 31 August 1964)

GLOBAL CONSERVATION LAWS AND MASSLESS PARTICLES*

G. S. Guralnik,[†] C. R. Hagen,[‡] and T. W. B. Kibble Department of Physics, Imperial College, London, England (Received 12 October 1964)

The Englert-Brout-Higgs Mechanism

- Vacuum expectation value of scalar field
- Englert & Brout: June 26th 1964
- First Higgs paper: July 27th 1964
- Pointed out loophole in argument of Gilbert if gauge theory described in Coulomb gauge
- Accepted by Physics Letters
- Second Higgs paper with explicit example sent on July 31st 1964 to Physics Letters, rejected!
- Revised version (Aug. 31st 1964) accepted by PRL

The Englert-Brout-Higgs Mechanism

Englert & Brout

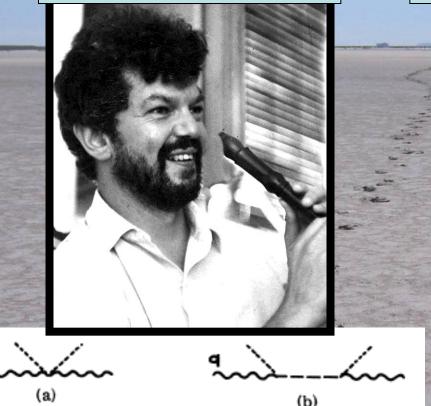


FIG. 1. Broken-symmetry diagram leading to a mass for the gauge field. Short-dashed line, $\langle \varphi_1 \rangle$; long-dashed line, φ_2 propagator; wavy line, A_{μ} propagator. (a) $\rightarrow (2\pi)^4 i e^2 g_{\mu\nu} \langle \varphi_1 \rangle^2$, (b) $\rightarrow -(2\pi)^4 i e^2 (q_{\mu}q_{\nu}/q^2) \times \langle \varphi_1 \rangle^2$.

Guralnik, Hagen & Kibble

We consider, as our example, a theory which was partially solved by Englert and Brout,⁵ and bears some resemblance to the classical theory of Higgs.⁶ Our starting point is the ordinary electrodynamics of massless spin-zero particles, characterized by the Lagrangian

$$\mathcal{L} = -\frac{1}{2} F^{\mu\nu} (\partial_{\mu} A_{\nu} - \partial_{\nu} A_{\mu}) + \frac{1}{4} F^{\mu\nu} F_{\mu\nu}$$
$$+ \varphi^{\mu} \partial_{\mu} \varphi + \frac{1}{2} \varphi^{\mu} \varphi_{\mu} + i e_{0} \varphi^{\mu} q \varphi A_{\mu}$$

With no loss of generality, we can take $\eta_2 = 0$, and find

$$(-\partial^{2} + \eta_{1}^{2})\varphi_{1} = 0,$$

$$-\partial^{2}\varphi_{2} = 0,$$

$$(-\partial^{2} + \eta_{1}^{2})A_{k}^{T} = 0,$$

where the superscript T denotes the transverse part. The two degrees of freedom of A_k^T combine with φ_1 to form the three components of a

The Higgs Boson

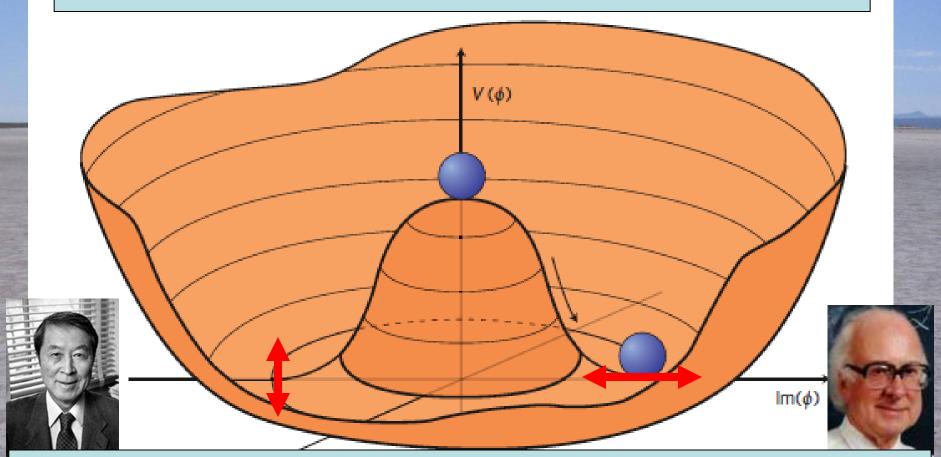
Higgs pointed out a massive scalar boson

 $\{\partial^2 - 4\varphi_0^2 V''(\varphi_0^2)\}(\Delta \varphi_2) = 0, \qquad (2b)$

Equation (2b) describes waves whose quanta have (bare) mass $2\varphi_0 \{V''(\varphi_0^2)\}^{1/2}$

- "... an essential feature of [this] type of theory ... is the prediction of incomplete multiplets of vector and scalar bosons"
- Englert, Brout, Guralnik, Hagen & Kibble did not comment on its existence
- Discussed in detail by Higgs in 1966 paper

Nambu EB, GHK and Higgs



Spontaneous symmetry breaking: massless Nambu-Goldstone boson **'eaten' by gauge boson**

Accompanied by massive particle

Without Higgs ...

- ... there would be no atoms
 - Electrons would escape at the speed of light
- ... weak interactions would not be weak
 - Life would be impossible: there would be no nuclei, everything would be radioactive

How does the Higgs trick work?

Masses for Gauge Bosons

• Kinetic terms for SU(2) and U(1) gauge bosons:

$$\mathcal{L} = -\frac{1}{4} G^{i}_{\mu\nu} G^{i\mu\nu} - \frac{1}{4} F_{\mu\nu} F^{\mu\nu}$$

where $G^i_{\mu\nu} \equiv \partial_\mu W^i_\nu - \partial_\nu W^i_\mu + ig\epsilon_{ijk}W^j_\mu W^k_\nu$ $F_{\mu\nu} \equiv \partial_\mu W^i_\nu - \partial_\nu W^i_\mu$

• Kinetic term for Higgs field:

SAC.

$$\mathcal{L}_{\phi} = -|D_{\mu}\phi|^2 D_{\mu} \equiv \partial_{\mu} - i \ g \ \sigma_i \ W^i_{\mu} - i \ g' \ Y \ B_{\mu}$$

Expanding around vacuum: $\phi = < 0|\phi|0 > + \hat{\phi}$

$$\mathcal{L}_{\phi} \ni -\frac{g^2 v^2}{2} \ W^+_{\mu} \ W^{\mu-} - g'^2 \ \frac{v^2}{2} \ B_{\mu} \ B^{\mu} + g \ g' v^2 \ B_{\mu} \ W^{\mu3} - g^2 \ \frac{v^2}{2} \ W^3_{\mu} \ W^{\mu3}$$

$$m_{W^{\pm}} = \frac{gv}{2} \qquad Z_{\mu} = \frac{gW_{\mu}^3 - g'B_{\mu}}{\sqrt{g^2 + g'^2}} : \quad m_Z = \frac{1}{2}\sqrt{g^2 + g'^2}v \; ; \quad A_{\mu} = \frac{g'W_{\mu}^3 + gB_{\mu}}{\sqrt{g^2 + g'^2}} : \quad m_A = 0$$

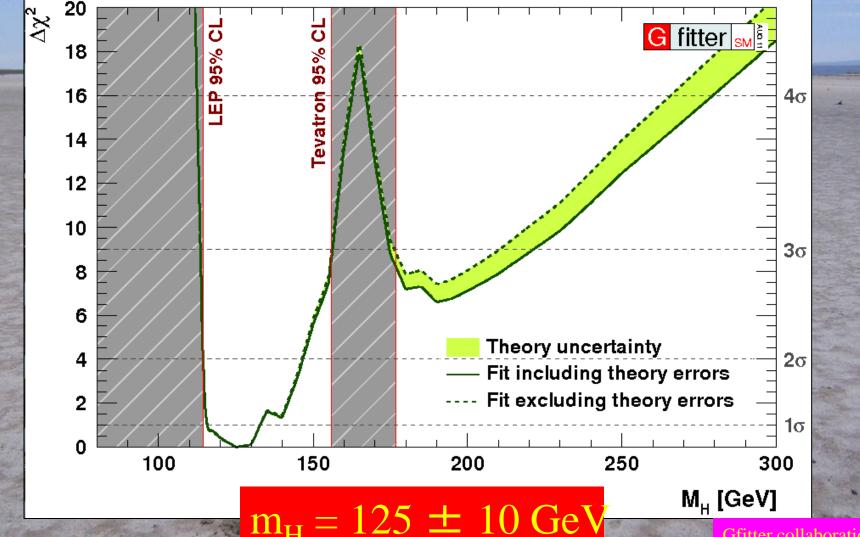
Constraints on Higgs Mass

- Electroweak observables sensitive via quantum loop corrections: $m_W^2 \sin^2 \theta_W = m_Z^2 \cos^2 \theta_W \sin^2 \theta_W = \frac{\pi \alpha}{\sqrt{2} G_F} (1 + \Delta r)$
- Sensitivity to top, Higgs masses:

$$\frac{3\mathbf{G}_F}{8\pi^2\sqrt{2}}m_t^2 \qquad \frac{\sqrt{2}\mathbf{G}_F}{16\pi^2}m_W^2(\frac{11}{3}\ln\frac{M_H^2}{m_Z^2}+\ldots), \, M_H >> m_W$$

- Preferred Higgs mass: m_H ~ 100 ± 30 GeV
 Compare with lower limit from direct searches: m_H > 114 GeV
- No conflict!

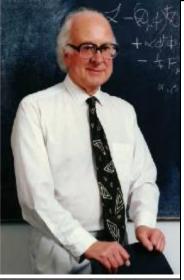
2011: Combining Information from Previous Direct Searches and Indirect Data

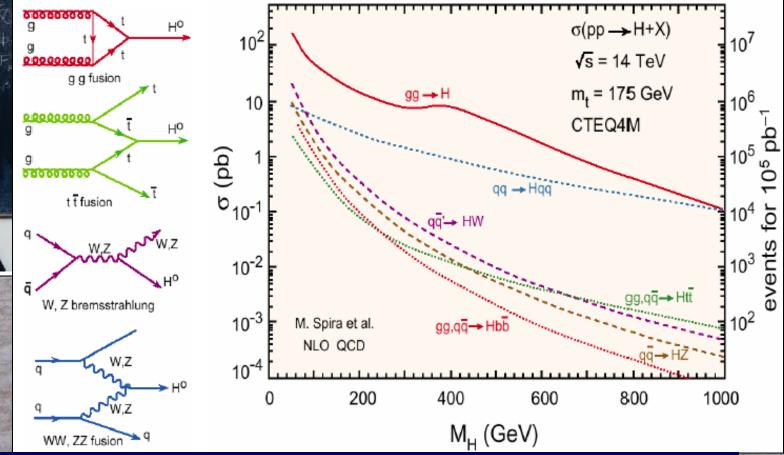


Gfitter collaboration

A la recherche du Higgs perdu

Higgs Production at the LHC

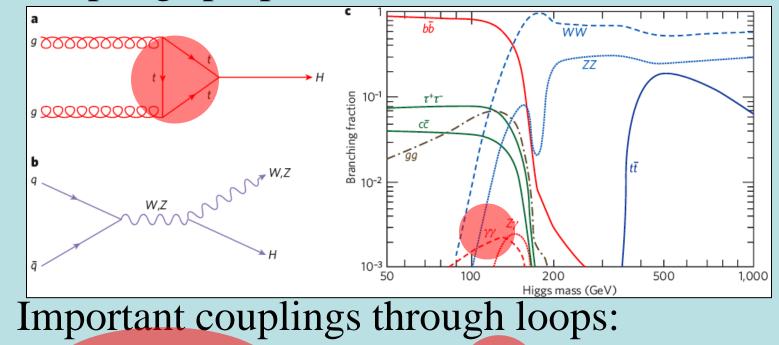




Many production modes measurable if $M_h \sim 125$ GeV

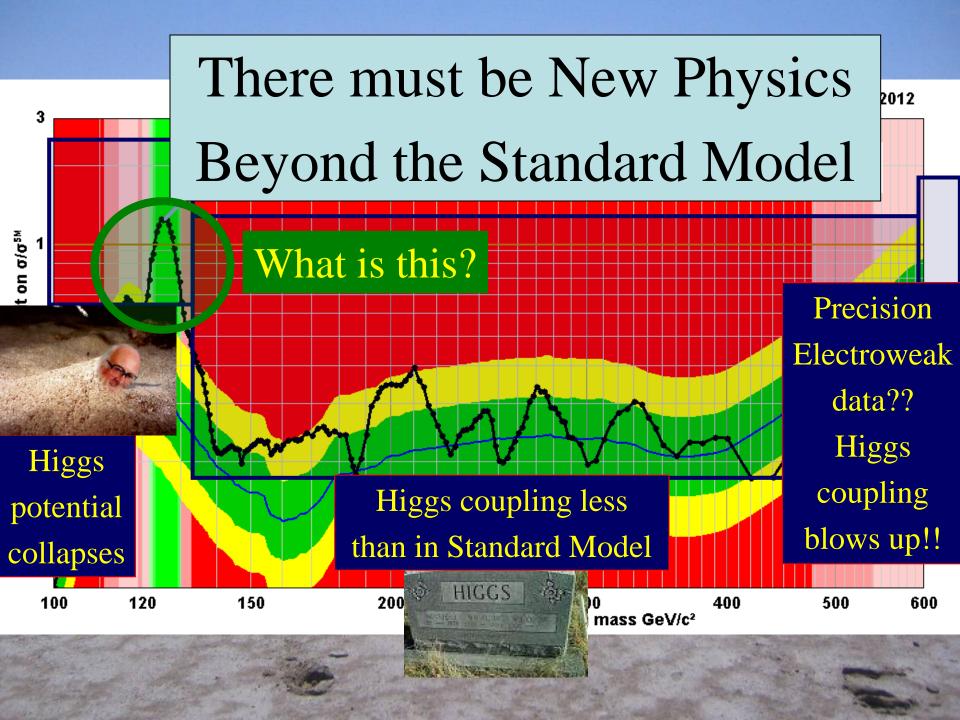
Higgs Decay Branching Ratios

• Couplings proportional to masses (?)

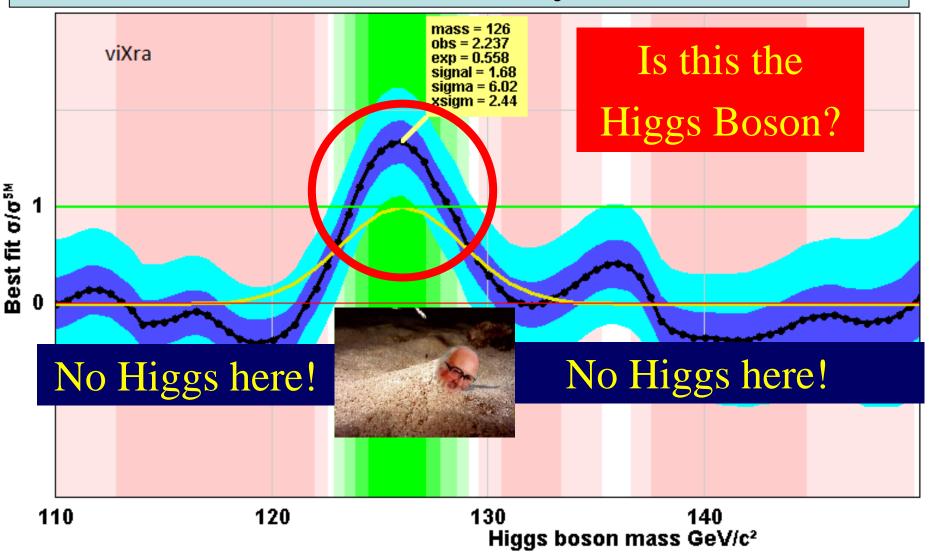


 $-gluon + gluon \rightarrow Higgs \rightarrow \gamma\gamma$

Many decay modes measurable if $M_h \sim 125 \text{ GeV}$



Unofficial Combination of Higgs Search Data from July 4th

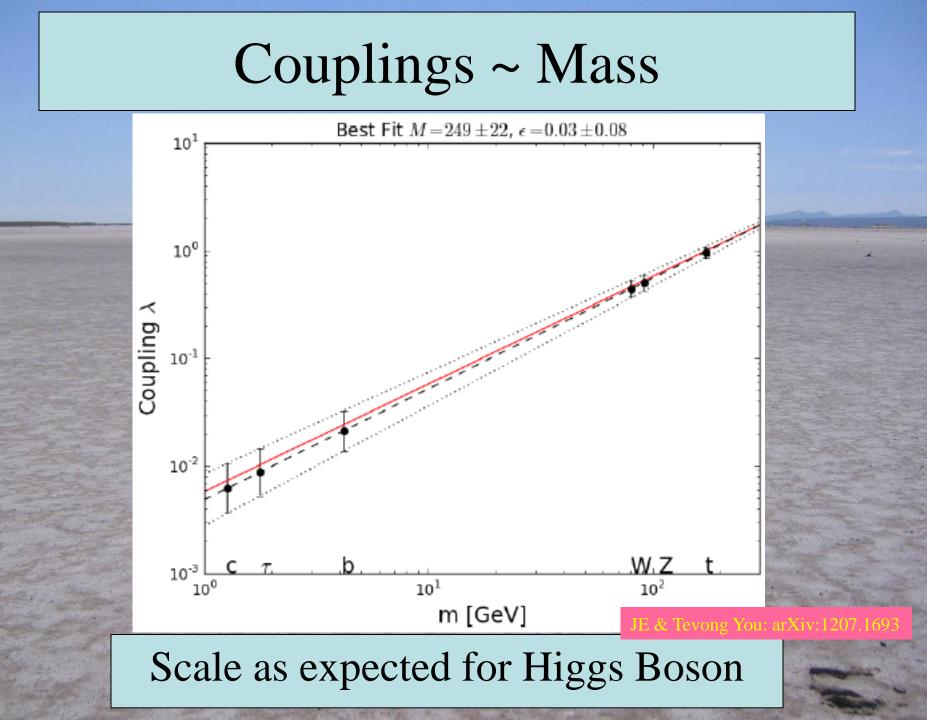


The Particle Higgsaw Puzzle

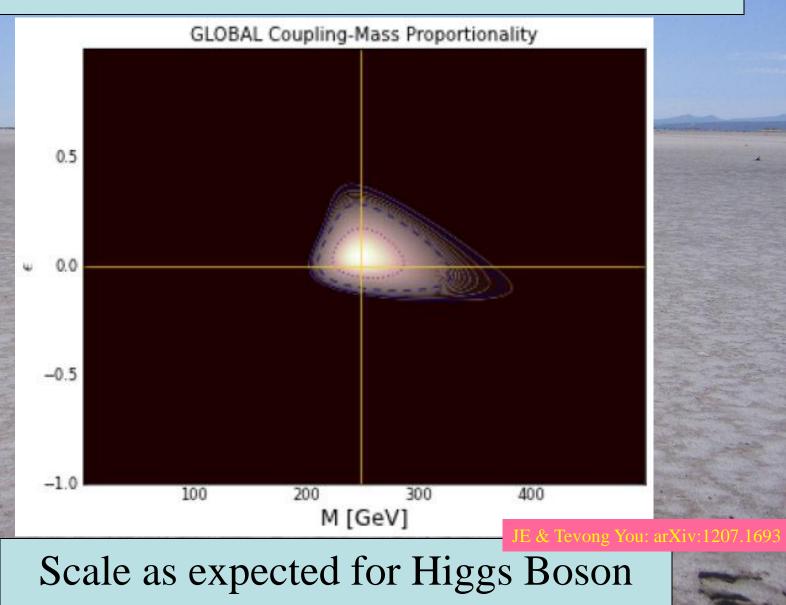
Is LHC finding the missing piece? Is it the right shape? Is it the right size?

Do we already know the 'Higgs' has Spin Zero?

- Decays into $\gamma\gamma$, so cannot have spin 1
- Spin 0 or 2?
- If it decays into ττ or b-bar: spin 0 or 1 or orbital angular momentum
- Can diagnose spin via
 - angular distribution of $\gamma\gamma$
 - angular correlations of leptons in WW, ZZ decays
- Does selection of WW events mean spin 0?



Couplings ~ Mass

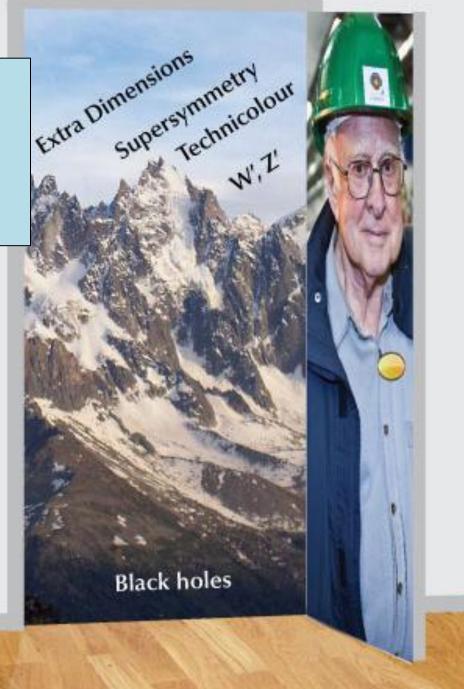


Imagine a Room ...



... Open The Door

What lies Beyond?



Black holes

Supersymmetry

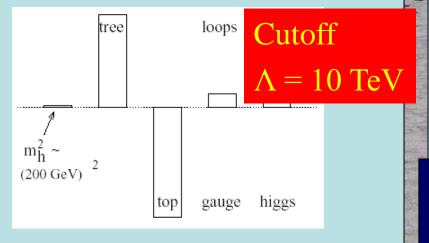
Technicolour

W',Z'

Extra Dimensions

Elementary Higgs or Composite?

- Higgs field: $<0|H|0> \neq 0$
- Quantum loop problems



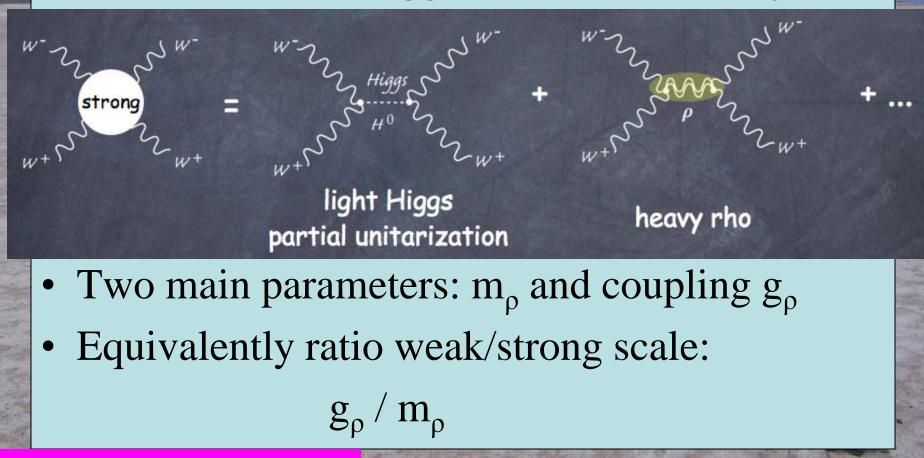
Cut-off $\Lambda \sim 1$ TeV with Supersymmetry?

- Fermion-antifermion condensate
- Just like QCD, BCS superconductivity
- Top-antitop condensate? needed m_t > 200 GeV
- New technicolour force?
- Heavy scalar resonance?
- Inconsistent with

precision electroweak data?

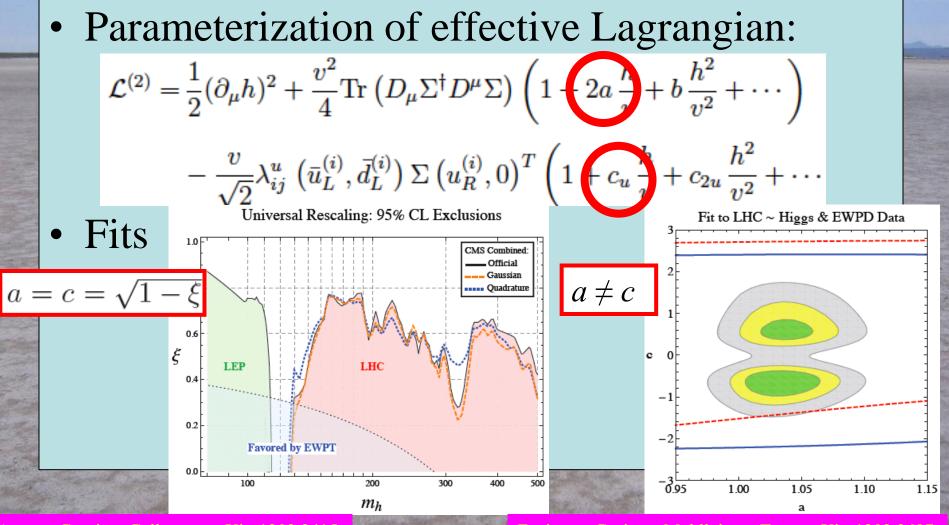
Interpolating Models

• Combination of Higgs boson and vector ρ



Grojean, Giudice, Pomarol, Rattazzi

General Analysis of Generalized Higgs Models

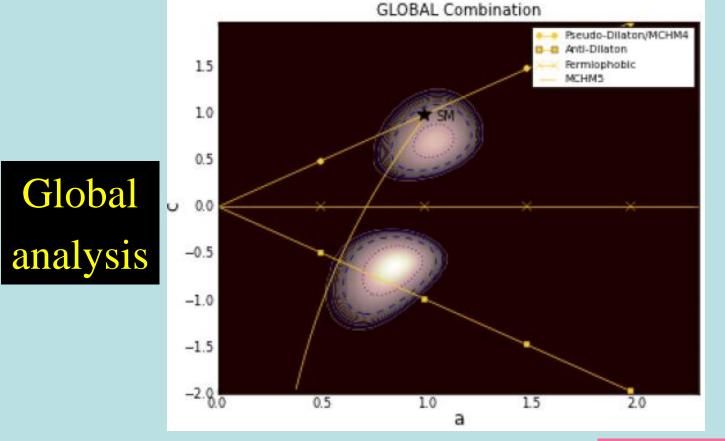


Azatov, Contino, Galloway: arXiv:1202.3415

Espinosa, Grojean, Muhlleitner, Trott: arXiv:1202.3697

Combination of LHC & Tevatron

• Rescale couplings: to bosons by a, to fermions by c



• Standard Model: a = c = 1

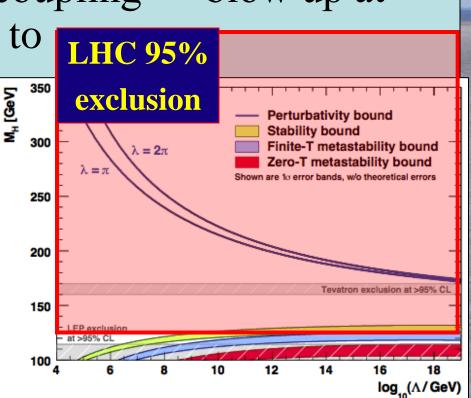
JE & Tevong You: arXiv:1207.1693

The Stakes in the Higgs Search

- How is gauge symmetry broken?
- Is there any elementary scalar field?
- Likely portal to new physics
- Would have caused phase transition in the Universe when it was about 10⁻¹² seconds old
- May have generated then the matter in the Universe: electroweak baryogenesis
- A related **inflaton** might have expanded the Universe when it was about 10⁻³⁵ seconds old
- Contributes to today's dark energy: 10⁶⁰ too much!

Theoretical Constraints on Higgs Mass

- Large $M_h \rightarrow$ large self-coupling \rightarrow blow up at low-energy scale Λ due to renormalization $\begin{bmatrix} 1 & 1 & 1 \\ 1 & 2 & 3 \end{bmatrix} = \begin{bmatrix} 1 & 1 & 1 \\ 1 & 2 & 3 \end{bmatrix} = \begin{bmatrix} 1 & 1 & 2 & 3 \\ 1 & 2 & 3 \end{bmatrix} = \begin{bmatrix} 1 & 1 & 2 & 3 \\ 1 & 2 & 3 \end{bmatrix} = \begin{bmatrix} 1 & 1 & 2 & 3 \\ 2 & 3 & 3 & 3 \\ 1 & 2 & 3 & 3 \end{bmatrix} = \begin{bmatrix} 1 & 1 & 2 & 3 \\ 2 & 3 & 3 & 3 \\ 1 & 2 & 3 & 3 \\ 1 & 2 & 3 & 3 \\ 1 & 2 & 3 & 3 \\ 1 & 2 & 3 & 3 \\ 1 & 2 & 3 & 3 \\ 1 & 2 & 3 & 3 \\ 1 & 2 & 3 & 3 \\ 1 & 2 & 3 & 3 \\ 1 & 2 & 3 & 3 \\ 1 & 2 & 3 & 3 \\ 1 & 2 & 3 & 3 \\ 1 & 2 & 3 & 3 \\ 1 & 2 & 3 & 3 \\ 1 & 2 & 3 & 3 \\ 1 & 2 & 3 & 3 \\ 1 & 2 & 3 & 3 \\ 1 & 3 & 3 & 3$
- Small: renormalization due to t quark drives quartic coupling < 0 at some scale Λ
 - \rightarrow vacuum unstable



Espinosa, JE, Giudice, Hoecker, Riotto, arXiv0906.0954

• Vacuum could be stabilized by **Supersymmetry**