



EUROPEAN
SPALLATION
SOURCE



Neutron Sources

PART 1

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European Spallation Source





Acknowledgment

This lecture relies on the precious input and materials from:

ESS: Axel Steuwer, Mats Lindroos, Colin Carlite, Ferenc Mezei,
Konstantín Batkov, Luca Zanini

EPFL: Giorgio Margaritondo

INFN/ESS: Santo Gammíno

CNRS/ESS: Sébastien Bousson and Alex Mueller

A special thanks to Mats Lindroos, who donated the money received for his 50th birthday to invite few students to this school.



PART 1

- Background for neutron course
- Neutrons properties and their interactions
- Applications using Neutrons

PART 2

- How to generate intense neutron beams
- High power proton linear accelerator

PART 3

- Examples of world-wide neutron sources



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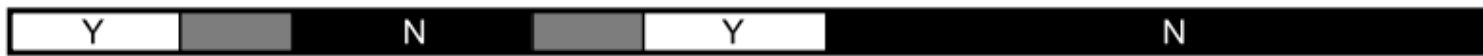


Context – Reference to ASP2012 Classes:

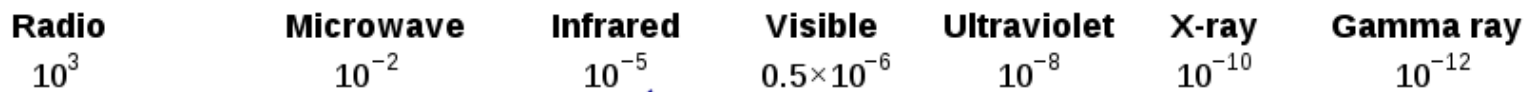
- Keteví Assamagan: Nuclear Physics (July 16 and 18)
- Lenny Rivkin: Particle beam dynamics, beam optics, instrumentations and light sources
- Lyn Evans: Accelerator sciences and the LHC
- Herman Winick: Light Sources

Electro-magnetic Spectrum

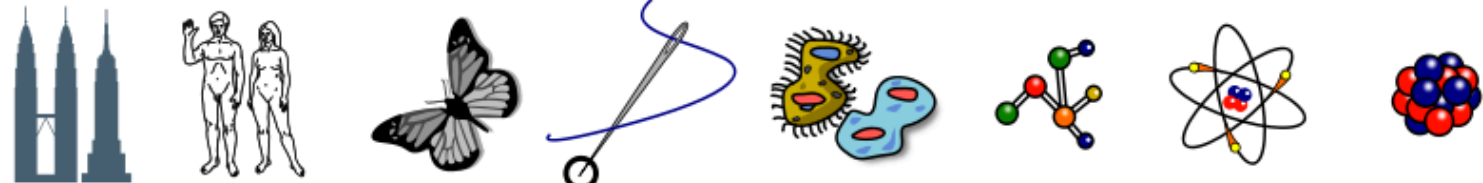
Penetrates Earth's Atmosphere?



Radiation Type
Wavelength (m)

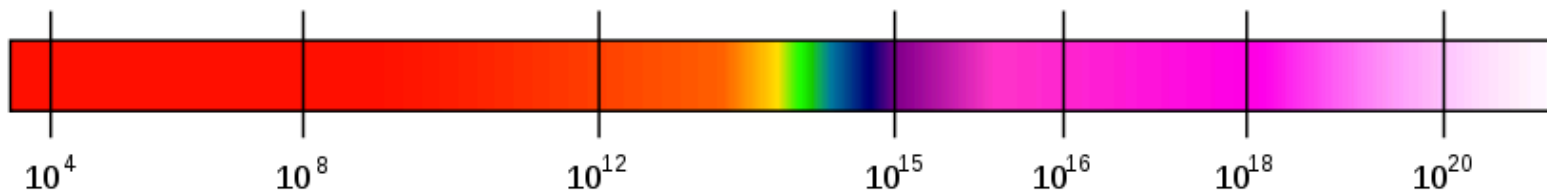


Approximate Scale of Wavelength

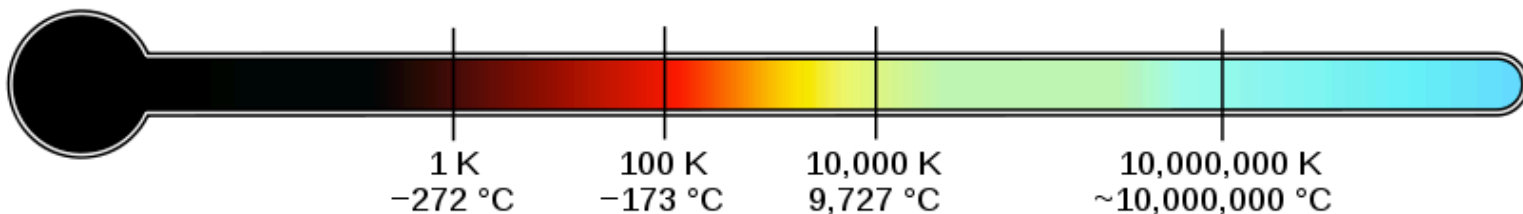


Buildings Humans Butterflies Needle Point Protozoans Molecules Atoms Atomic Nuclei

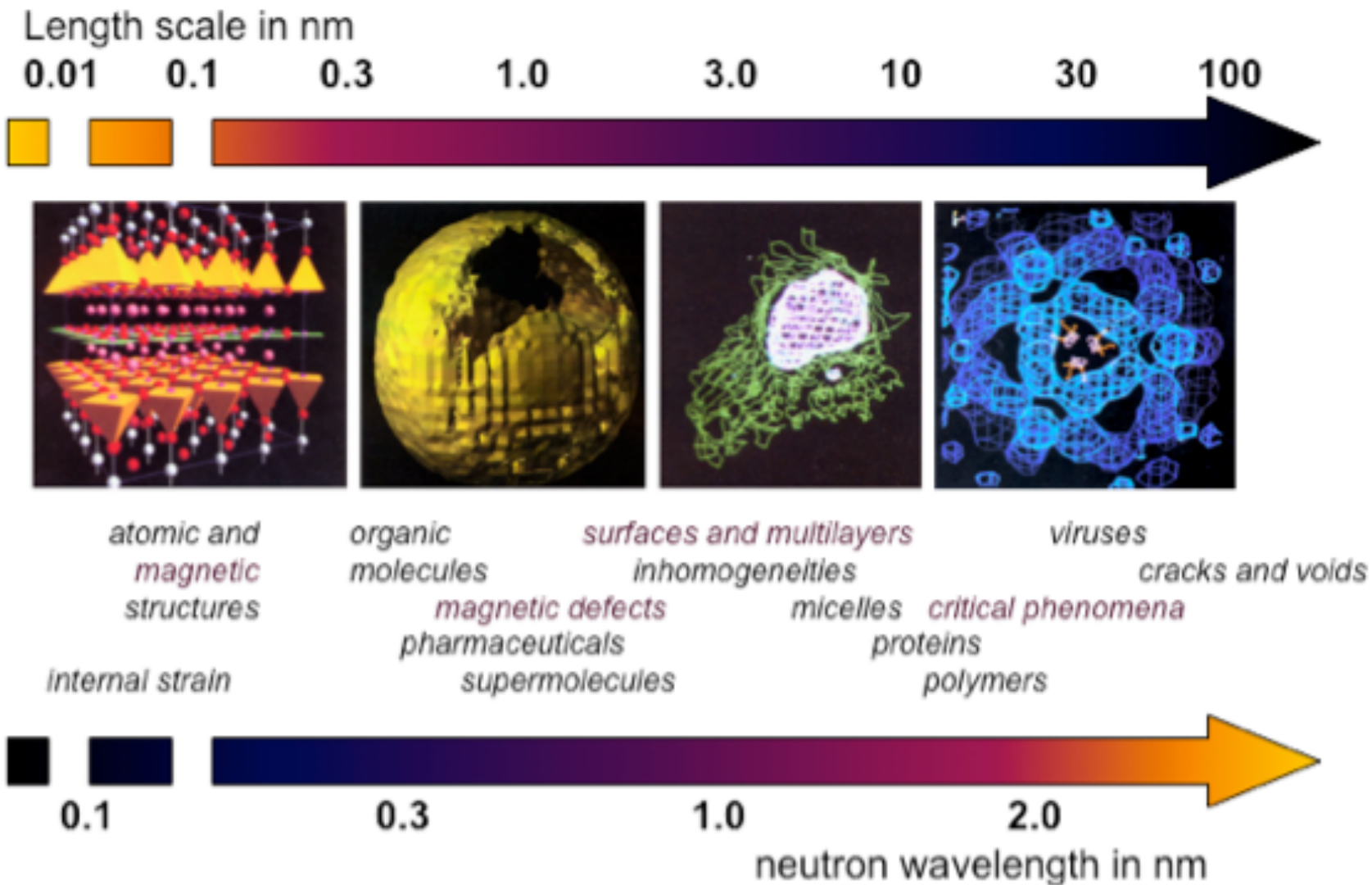
Frequency (Hz)



Temperature of objects at which this radiation is the most intense wavelength emitted



Neutron Microscope – Length scales

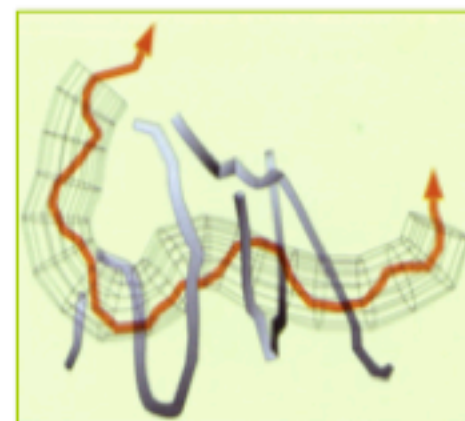
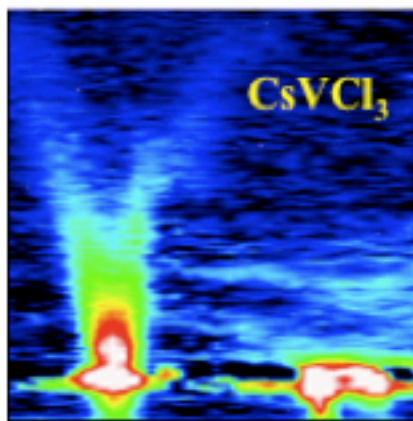
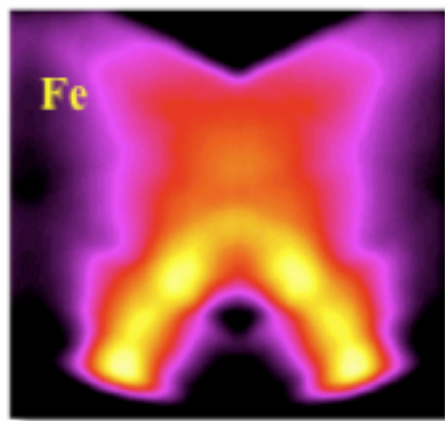


Neutron Microscope – Time and energy scales

Time scale (seconds)

10^{-13}

10^{-7}



Crystal fields

*single particle
excitations
molecular excitations*

magnons and phonons

spin fluctuations

*tunneling
diffusion*

spin relaxation

*polymer reptation
glassy dynamics
libration*



Excitation energy (eV)

1

10^{-1}

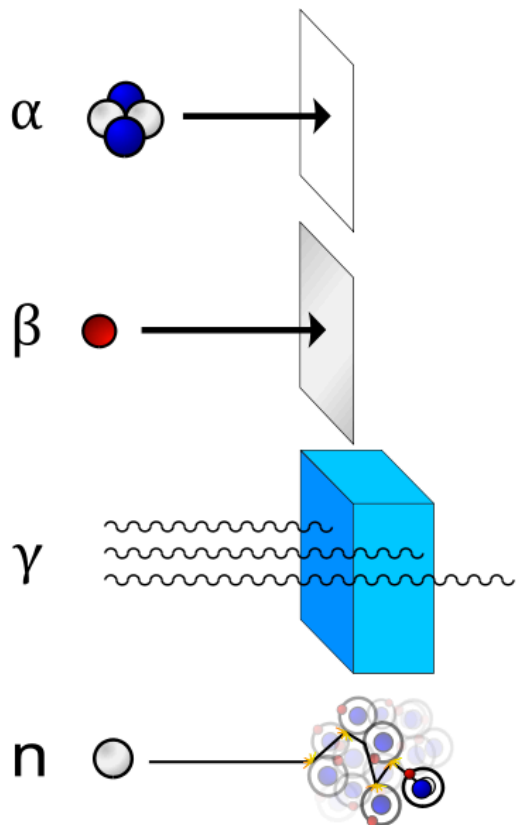
10^{-2}

10^{-3}

10^{-6}

10^{-9}

Ionizing radiation is radiation composed of particles that individually carry enough energy to liberate an electron from an atom or molecule without raising the bulk material to ionization temperature.



When ionizing radiation is emitted by or absorbed by an atom, it can liberate a particle.

Such an event can alter chemical bonds and produce ions, usually in ion-pairs, that are especially chemically reactive.

Note: Neutrons, having zero electrical charge, **do not interact electromagnetically** with electrons, and so they cannot *directly* cause ionization by this mechanism.

➔ High precision non-destructive probe ... why ?



- 1920 **E. Rutherford**: disparity between the atomic number of an atom and its atomic mass could be explained by a neutral particle within atomic nucleus.
- 1930 **V. Ambartsumian and D. Ivanenko**: proved that the nucleus cannot consist of protons and electrons only.
- 1931 **W. Bothe and H. Becker**: interactions of α -particles with light nuclei (Li, Be, B) produce an unusually penetrating radiation.
- 1932 **Irène Joliot-Curie and F. Joliot**: interactions of *beryllium radiation* with hydrogen-containing compounds ejects energetic protons (moderation).



80th Birthday – June 1st, 2012

----- Original Message -----

Subject: The neutron's 80th birthday
Date: Fri, 01 Jun 2012 12:29:01 +0200
From: Francoise Vauquois <vauquois@ill.fr>
To: gen_l@ill.fr

Dear all,

It is 80 years today since the publication by the Royal Society of Cambridge Physicist James Chadwick's famous paper proving the existence of the neutron - on 1 June 1932. Chadwick was awarded the Nobel Prize for his paper two years later.

Chadwick's discovery launched the development of neutron research, which has since made incontestably decisive breakthroughs for modern science. And the ILL has been world-leading in this domain! 2012 is also the fortieth anniversary of the start of experiments with neutrons at ILL, in 1972.

You will find more on the history of Chadwick's discovery on the ILL website:

<http://www.ill.eu/fr/news-events/presse/communiqués-de-presse/ill-fete-les-80-ans-du-neutron-1062012/>

You can also hear the ILL director, Andrew Harrison, interviewed for the occasion by the BBC, by visiting the BBC's site, on <http://www.bbc.co.uk/programmes/b01j6t0n>

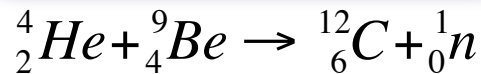
The French newspaper Le Figaro has also marked the occasion:

<http://www.lefigaro.fr/sciences/2012/05/24/01008-20120524ARTFIG00844-le-neutron-fete-ses-80-ans-cette-annee.php>

and so did Spain's Muy interesante:

[http://www.muyinteresante.es/el-neutron-cumple-80-anos?
utm_source=twitter&utm_medium=socialomph&utm_campaign=muy-interesante-twitter689](http://www.muyinteresante.es/el-neutron-cumple-80-anos?utm_source=twitter&utm_medium=socialomph&utm_campaign=muy-interesante-twitter689)

James Chadwick 1932 (α, n) reaction



“Whatever the radiation from Be may be, it has most remarkable properties”

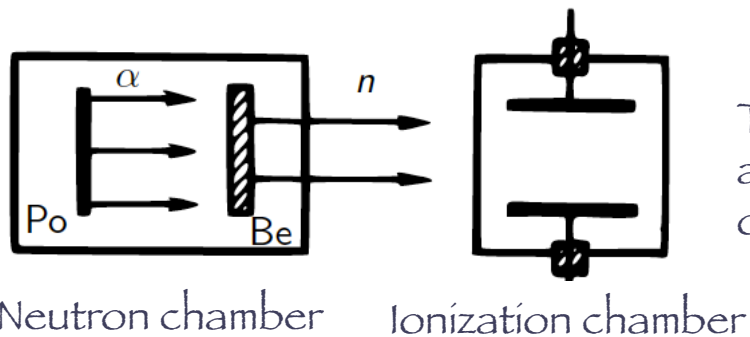
Cambridge Laboratory,
Cambridge,
24 February 1932.

Dear Bohr,

I enclose the proof of a letter I have written to "Nature" and which will appear either this week or next. I thought you might like to know about it beforehand.

The suggestion is that α particles eject from beryllium (and also from boron) particles which have no net charge, and which probably have a mass about equal to that of the proton. As you will see, I put this forward rather cautiously, but I think the evidence is really rather strong. Whatever the radiation from Be may be, it has most remarkable properties. I have made many experiments which I do not mention in the

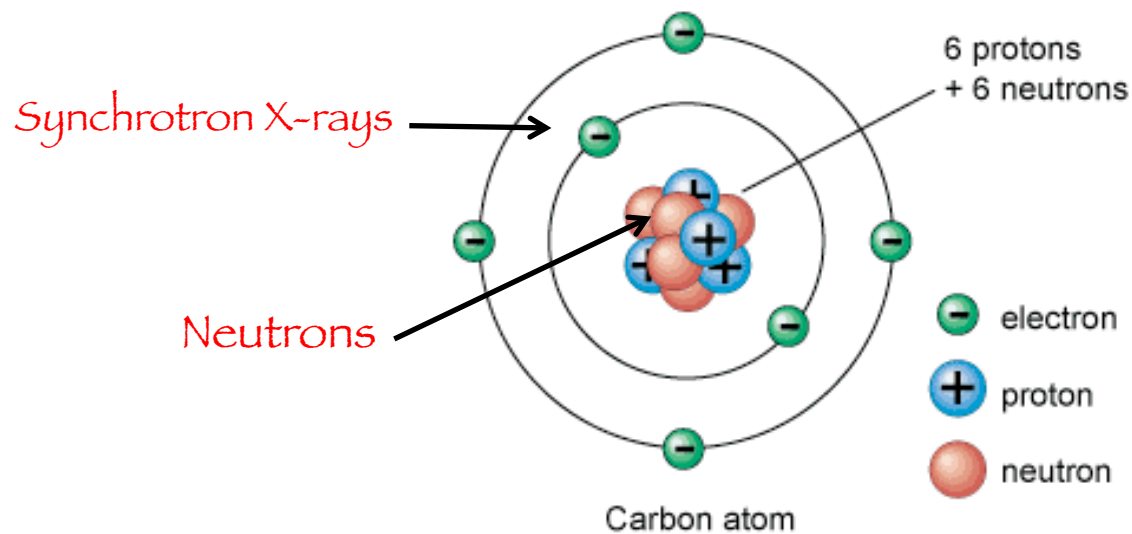
letter to "Nature" and they can all be interpreted readily on the assumption that the particles are neutrons. Feather has taken some pictures in the dispersion chamber and we have already found about 20 cases of recoil atoms. About 4 of these show an abrupt end (and it is almost certain that ~~the~~ one arm of this fork represents a recoil atom and the other some other particle, probably an α particle. They are disintegrated due to the capture of the neutron Np or O_{16} . I enclose two photographs one of which shows a single recoil atom, and the other shows a pair of recoil atoms. The ~~above~~ were printed



To amplifier and oscilloscope



Neutrons Properties



Note: X-rays are emitted by electrons outside the nucleus, while gamma rays are emitted by the nucleus.

1.675×10^{-27} kg	Mass	939.57 MeV	$m_n \approx m_p + 2.5m_e$
	Mean lifetime	15 min	$n \rightarrow p + e^- + \bar{\nu}_e$
	Composition	udd	hadron
	Electric charge	0	high penetration
	Magnetic moment	$-1.04 \mu_B$	feels the nucleus

$$E = k_B T$$

Boltzmann distribution

$$E = k_B T = \frac{1}{2} m v^2 = \frac{h}{2m\lambda^2}$$

$$\lambda = \frac{h}{mv}$$

De Broglie

$$E [meV] = 0.0862 T [K] = 5.22 v^2 [km/s] = 81.81 \frac{1}{\lambda^2} [A]$$

Source	Energy	Temperature	Wavelength
cold	0.1-10	1-120	30-3
thermal	5-100	60-1000	4-1
hot	100-500	1000-6000	1-0.4

- Wavelengths comparable to interatomic spacings (1-5 Å)
- Energies comparable to structural and magnetic excitations (1-100 meV)
- Neutrons interact only weakly with matter
- Neutrons are deeply penetrating (bulk samples can be studied)
- Neutrons are scattered with a strength that varies randomly from element to element (and isotope to isotope) – tuning?
- Neutrons have a magnetic moment : $\mu_n \approx -1.913 \mu_N$

→ Neutron scattering is therefore an ideal probe of magnetic and atomic structures and excitations

Why neutrons?

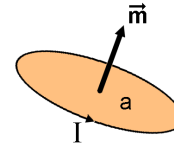
Wave



Particle



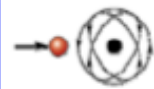
Magnetic moment



Neutral



Neutron properties are used to understand the nature of the solid and liquid states of matter, as an analytical tool to aid the development of materials and as a tool to examine curiosity-driven research that spans from cosmology, superconductivity to the dynamics of the molecules of life.



1. Neutrons see the Nuclei



2. Neutrons see Elementary Magnets



3. Neutrons see light Atoms next to Heavy Ones



4. Neutrons measure the Velocity of Atoms



5. Neutrons penetrate deep into Matter



6. Neutrons are Elementary Particles

Electrically Neutral – neutrons are non-destructive and can penetrate deep into matter. This makes them an ideal probe for biological materials and samples under extreme conditions of pressure, temperature, magnetic field or within chemical reaction vessels.

Microscopically Magnetic – they possess a magnetic dipole moment which makes them sensitive to magnetic fields generated by unpaired electrons in materials. Precise information on the magnetic behavior of materials at atomic level can be collected. In addition, the scattering power of a neutron off an atomic nucleus depends on the orientation of the neutron and the spin of the atomic nuclei in a sample. This makes the neutron a powerful instrument for detecting the nuclear spin order.

Angstrom wavelengths – neutron wavelengths range from 0.1 Å to 1000 Å, making them an ideal probe of atomic and molecular structures, be they single atomic species or complex biopolymers.

Energies of millielectronvolts – their energies are of the **same magnitude as the diffusive motion in solids and liquids**, the coherent waves in single crystals (phonons and magnons), and the vibrational modes in molecules. It is easy to detect any exchange of energy between a sample of between 1 microeV (even 1 neV with spin-echo) and 1eV and an incoming neutron.

Randomly sensitive – with neutrons the variation in scattering power from one nucleus to another within a sample varies in a quasi-random manner. This means that lighter atoms are visible despite the presence of heavier atoms, and neighboring atoms may be distinguished from each other. In addition, **contrast** can be varied in certain samples **using isotopic substitution (for example D for H, or one nickel isotope for another)**; specific structural features can thus be highlighted. The neutron is particularly sensitive to hydrogen atoms; it is therefore a powerful probe of hydrogen storage materials, organic molecular materials, and biomolecular samples or polymers.

The Nobel Prize in Physics 1994



Clifford G. Shull, MIT, Cambridge, Massachusetts, USA, receives one half of the 1994 Nobel Prize in Physics for development of the neutron diffraction technique.



S Shull made use of elastic scattering i.e. of neutrons which change direction without losing energy when they collide with atoms.

Because of the same nature of neutrons, a diffraction pattern can be recorded which indicates where in the sample the atoms are situated. Even the placing of light elements such as hydrogen in metallic hydrides, or hydrogen, carbon and oxygen in organic substances can be determined.

The pattern also shows how atomic dipoles are oriented in magnetic materials, since neutrons are affected by magnetic forces. Shull also made use of this phenomenon in his neutron diffraction technique.



An early neutron scattering experiment at the University of Chicago, 1940s. Shull is on the right.

Neutrons see more than X-rays

X-rays are scattered by electrons, neutrons by nuclei. With X-rays it is easier to see atoms that have many electrons. Neutrons, for example, which see only one electron, can see to see. With neutrons, all kinds of atoms are visible.



But in neutron diffraction, unlike X-rays, the scattering is not only by the nucleus. It is also by the magnetic moments of the atoms. In a neutron scattering experiment, a beam of neutrons is directed at a sample. Some neutrons are scattered, and some are absorbed. The scattered neutrons are then detected by a detector.

Neutrons reveal inner stresses

It has been pointed out in an important recent article that neutrons can be used to study the internal stresses in materials. This is done by measuring the change in the distance between the atoms in a crystal lattice when the material is stressed.



The stress causes the lattice to expand or contract. This change in the lattice spacing is detected by a detector. The detector is placed at a certain angle, and the intensity of the scattered neutrons is measured. The change in intensity is proportional to the change in the lattice spacing.

Neutrons show what atoms remember

Of these rather problems when they were conducted in solution or in a solid state, neutrons can be used to study the internal stresses in materials. This is done by measuring the change in the distance between the atoms in a crystal lattice when the material is stressed.



The detector is placed at a certain angle, and the intensity of the scattered neutrons is measured. The change in intensity is proportional to the change in the lattice spacing.

Neutrons behave as particles and as waves

Neutrons reveal structure and dynamics

Neutrons show where atoms are

When the neutrons collide with atoms in the sample material, they change direction (see scattering) - elastic scattering.



Crystal that sorts out forwards neutrons of a certain wavelength (range) - more orientational neutrons

Neutrons bounce against atomic nuclei. They also react to the magnetism of the atoms.

Research reactor



Neutrons show what atoms do

3-axis spectrometer with movable crystals and rotatable sample

Atoms in a crystalline sample

Crystal that sorts out forwards neutrons of a certain wavelength (range) - more orientational neutrons

When the neutrons generate the sample they start to oscillate in the atoms. If the neutrons transfer momentum or energy they transfer it to the atoms - inelastic scattering.

Changes in the energy of the neutrons are first analysed in an analyser crystal...

...and the neutrons then scattered in a detector.

The Royal Swedish Academy of Sciences has awarded the 1994 Nobel Prize in Physics for pioneering contributions to the development of neutron scattering techniques for studies of condensed matter.

Bertalan N. Brockhouse, McMaster University, Hamilton, Ontario, Canada, receives one half of the 1994 Nobel Prize in Physics for the development of neutron spectroscopy.



B Brockhouse made use of inelastic scattering i.e. of neutrons, which change both direction and energy when they collide with atoms. They then start or cause atomic oscillations in crystals and record movements in liquids and solids. Neutrons can also interact with spin waves in magnets.

With his 3-axis spectrometer Brockhouse measured energies of phonons (atomic vibrations) and magnons (magnetic waves). He also studied how atomic structures in liquids change with time.

Where it started

Brockhouse and Shull made their pioneering contributions at the first nuclear reactors in the USA and Canada back in the 1940s and 1950s. It was then that the neutrons of the reactor became available for scientific research.

... here it continues

Thousands of research centres are now working at the many neutron research centres throughout the world. New and new advanced neutron scattering installations have been built and more are planned in Europe, the USA and Asia. At these super-installations the researchers are studying the structure of new materials, superconductors, molecular membranes on surfaces of atoms by analysis without changing their structure and the connection between the structure and the atomic properties of polymers.



Information on the 1994 Nobel Prize in Physics can be found at the Nobel Prize website: <http://www.nobelprize.org>. The Royal Swedish Academy of Sciences, Box 1307, S-171 21 Solna, Sweden. Tel: +46 8 737 9300. Fax: +46 8 737 9301. E-mail: info@nobelprize.org. The Royal Swedish Academy of Sciences, Box 1307, S-171 21 Solna, Sweden. Tel: +46 8 737 9300. Fax: +46 8 737 9301. E-mail: info@nobelprize.org.

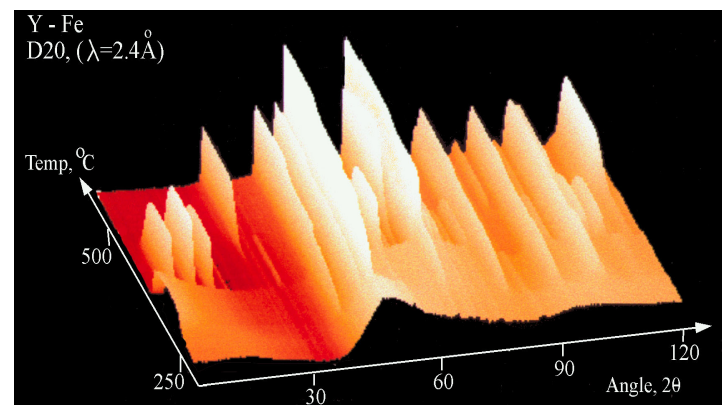
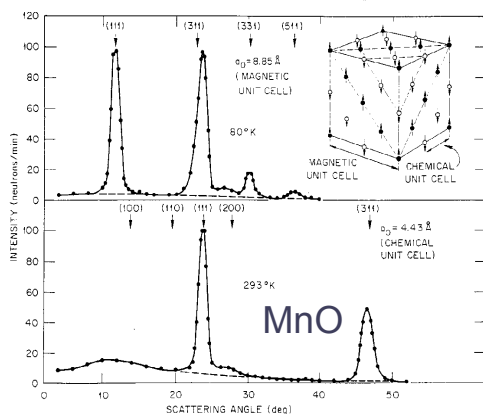
Further reading:

- <http://www.nobelprize.org> - The Nobel Prize in Physics 1994
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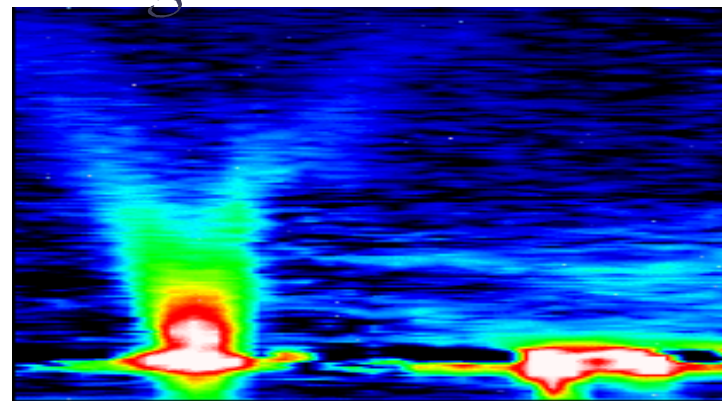
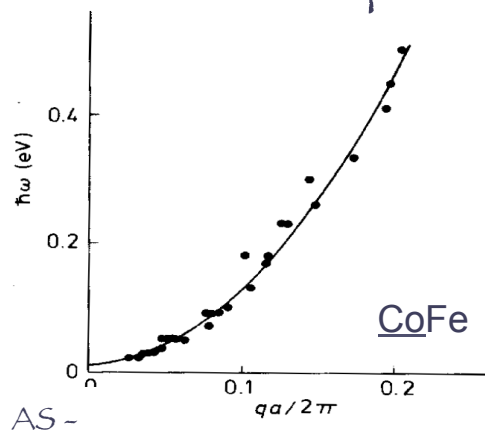
Historic - Why neutrons?

In half a century we have developed neutron scattering science enormously with an effective gain in source performance of **only a factor of 4!**

Cliff Shull - Neutron diffraction - showing where atoms are:

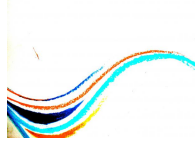


Bert Brockhouse - Spectroscopy - showing what atoms do:



Using nowadays technique..

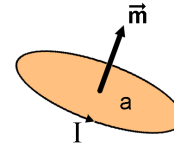
Wave



Particle



Magnetic moment

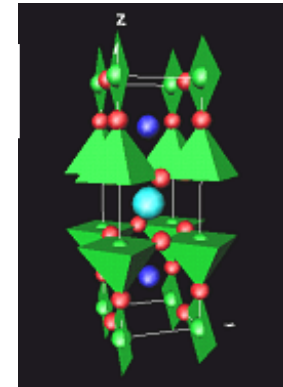


Neutral



Diffraction meters - Measure structures

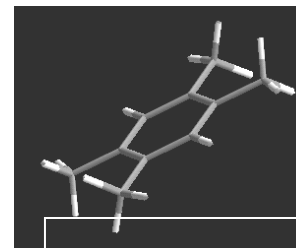
- Where atoms and molecules are 1 - 10 Ångström



→ To analyze the structure of a material from the scattering pattern produced when a beam of radiation or particles (such as X-rays or neutrons) interacts with it

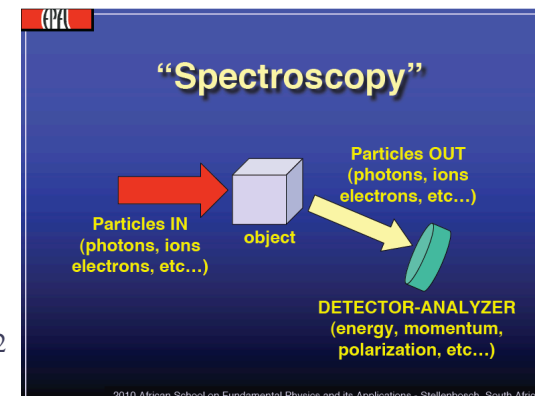
Spectrometers - Measure dynamics

- What atoms and molecules do

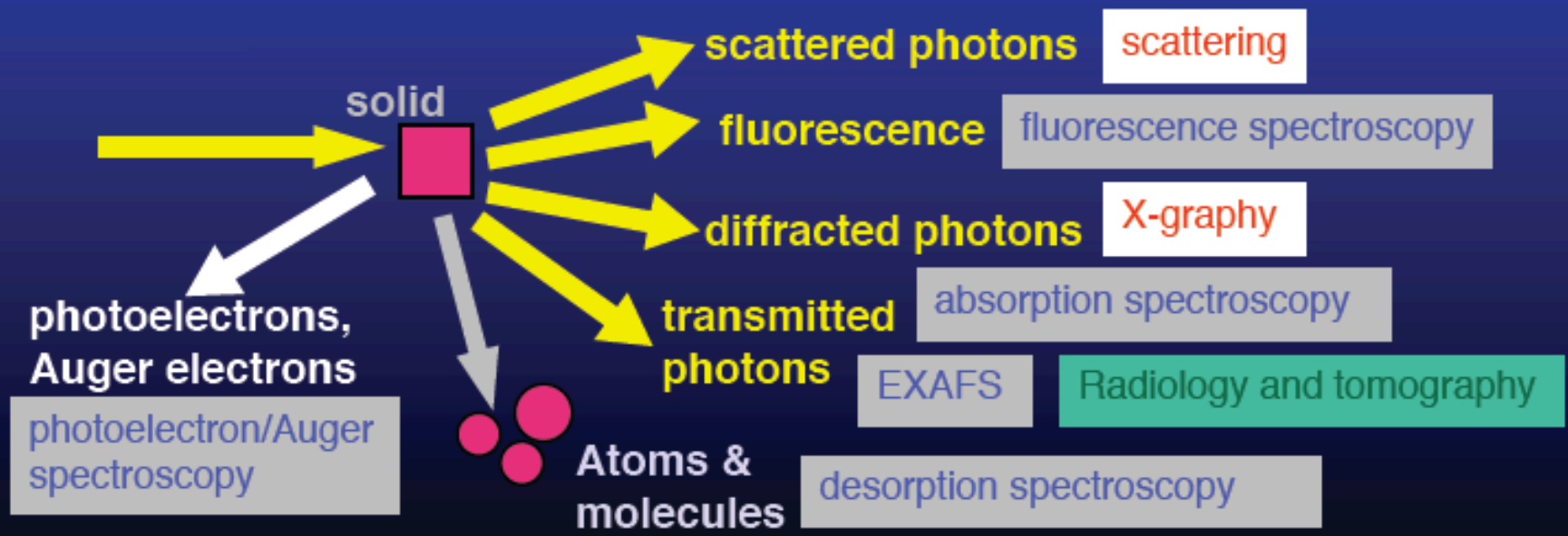
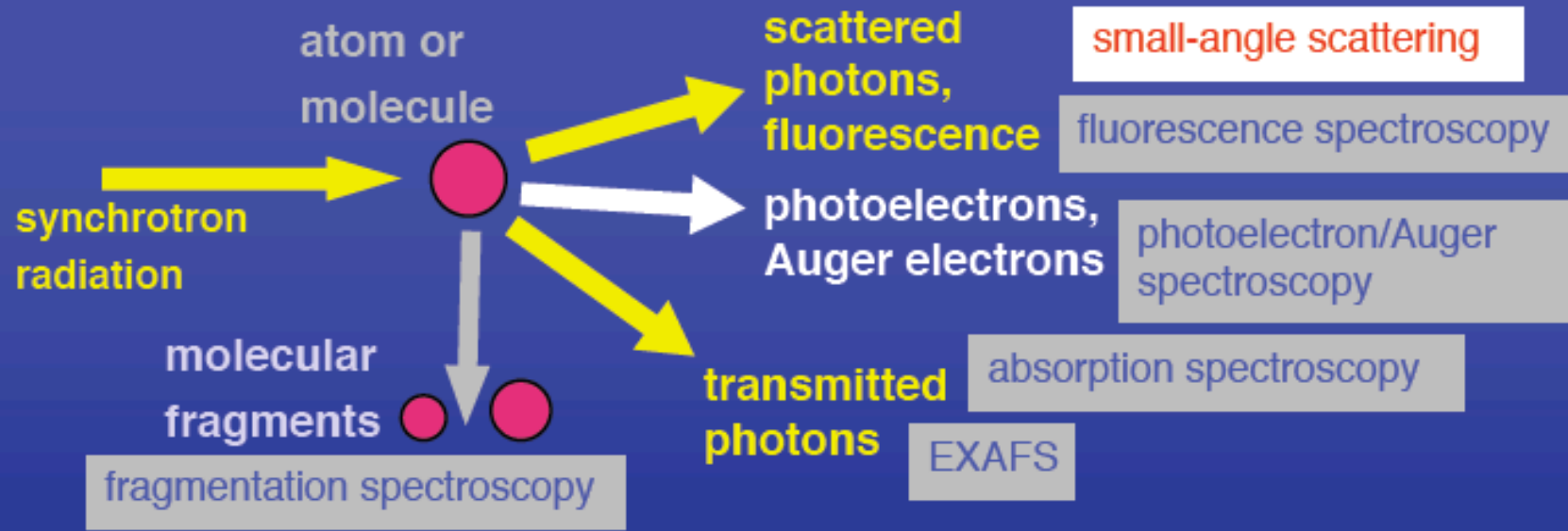


1 - 80 meV

→ To measure properties of light over a specific portion of the electromagnetic spectrum

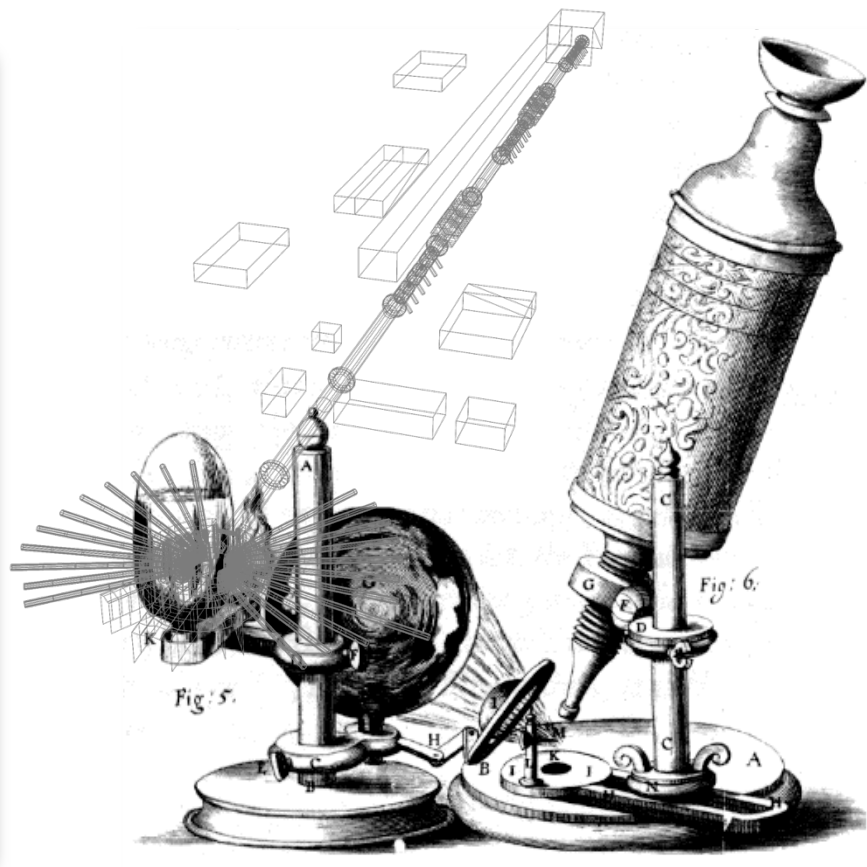
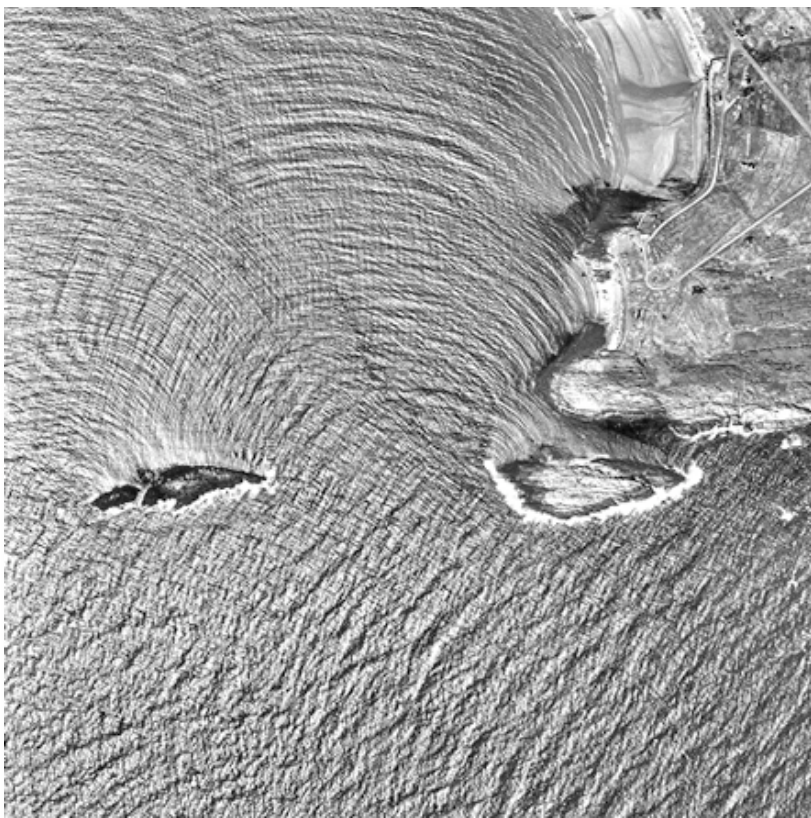


Three classes of experiments:





X-Ray and Neutron beam



Consider ESS/MaxIV equal in terms of functionality

Neutrons

- Particle beam (neutral subatomic particle)
- Interactions with the nuclei and the magnetic moment of unpaired electrons (in the sample)
- Scattered by all elements, also the light ones like the hydrogen isotopes
- Deep penetration depth (bulk studies of samples)
- Less intense beam measuring larger samples

Synchrotron radiation

- Light beam (electromagnetic wave)
- Interactions with the electrons surrounding the nuclei (in the sample)
- Mainly scattered by heavy elements
- Small penetration depth (surface studies of samples)
- Very intense beam measuring small or ultra-dilute samples

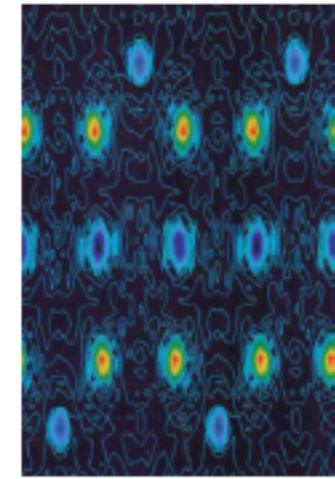
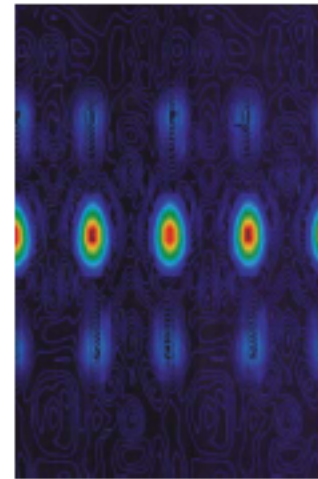
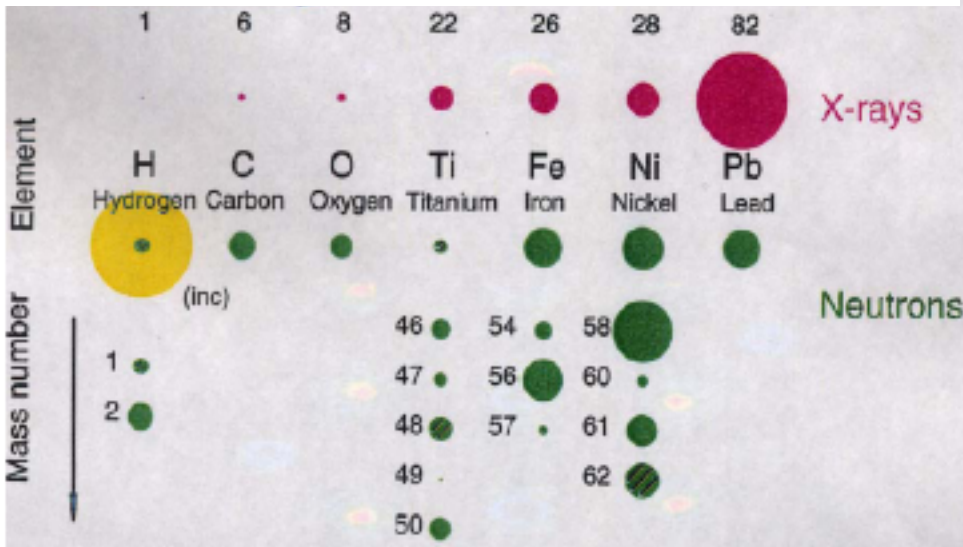
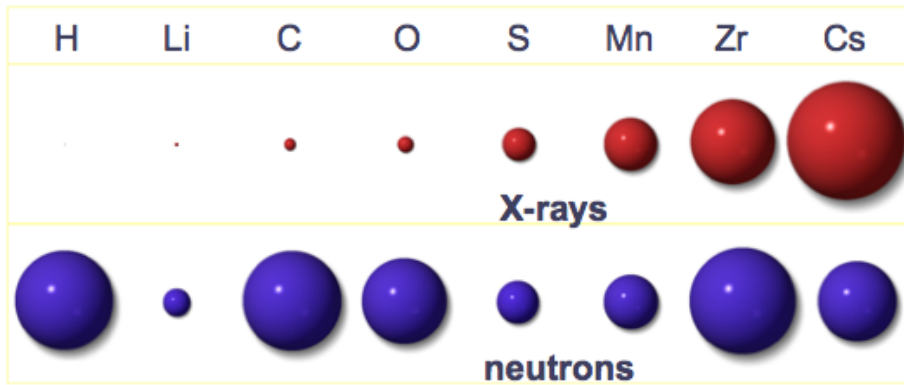
Neutrons Applications

- Magnetic structures & excitations
- Organic structures using the H-D isotope effect
- Bulk studies (strains, excitations)
- Low-energy spectroscopy e.g. molecular vibrations

SR Applications

- Protein-crystal structures
- Fast chemical reactions
- Surface studies (defects, corrosion)
- High-energy spectroscopy e.g. measurements of electron energy-levels

Complementarity between X-rays and Neutrons



— Li
— O
— Mn
— O
— Li

Material for Li-battery seen by
X rays (left) and
Neutrons (right)

T. Kamiyama, et al.

X-rays interact with electrons.

→ X-rays see high-Z atoms.

Neutrons interact with nuclei.

→ Neutrons see low-Z atoms.



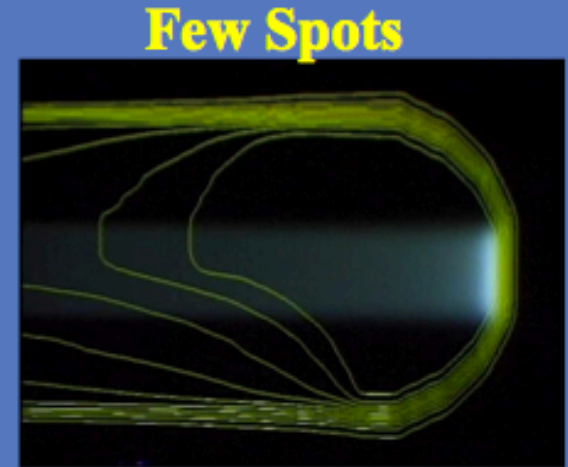
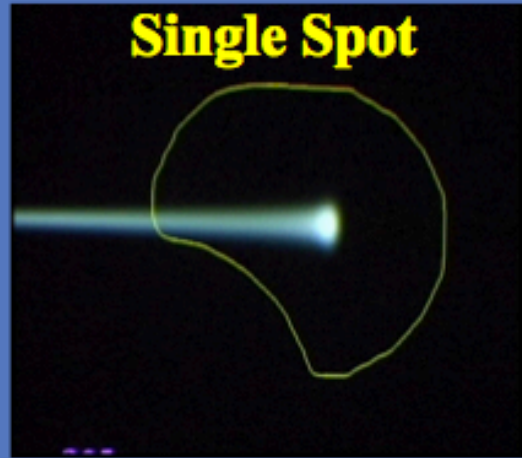
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Neutron beam and X-Ray for Medical Applications

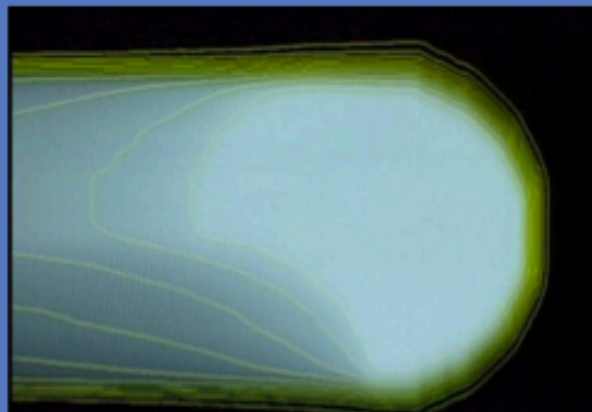
iThemba Particle Therapy Centre [iTPTC]



Spot Scanning Principle

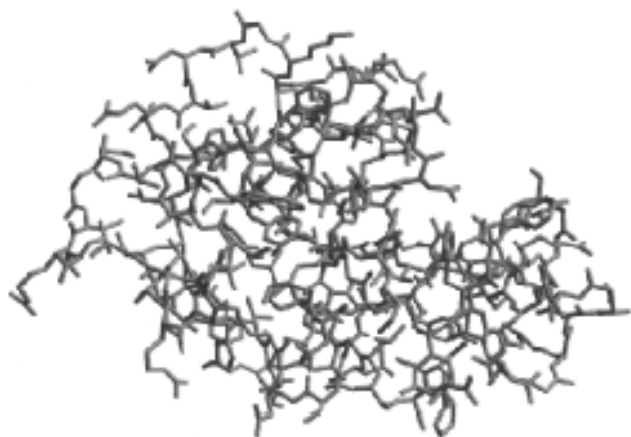


Final Dose Distribution

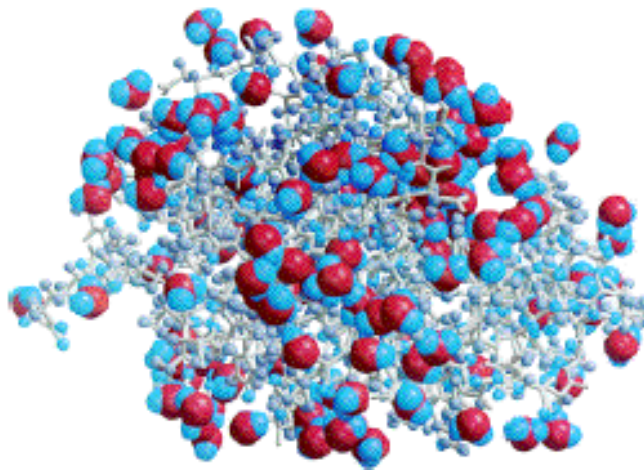


Courtesy of Zeblon Vilakazi

Hen Egg-White Lysozyme



X-rays



Neutrons

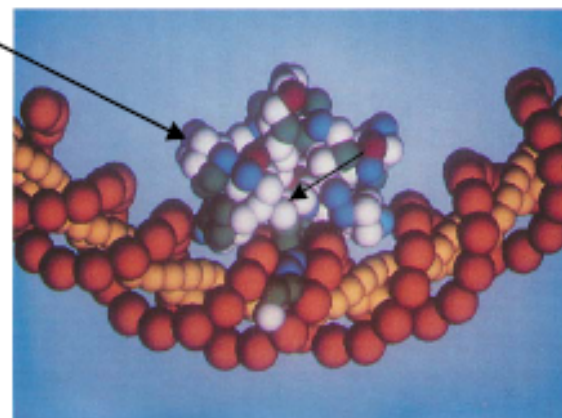
Water molecules
Observed with
neutrons

N. Niimura, et al.



From structure to function

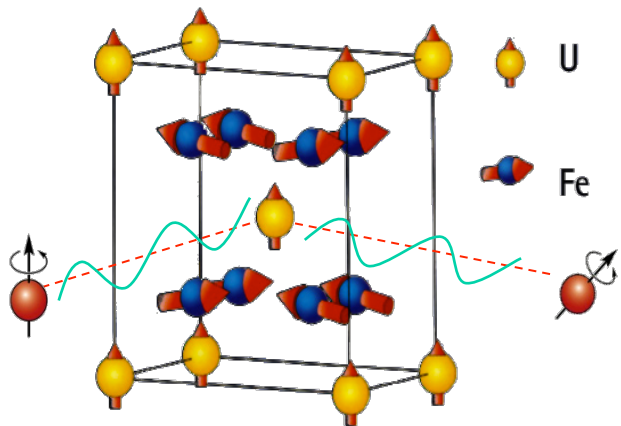
Protein



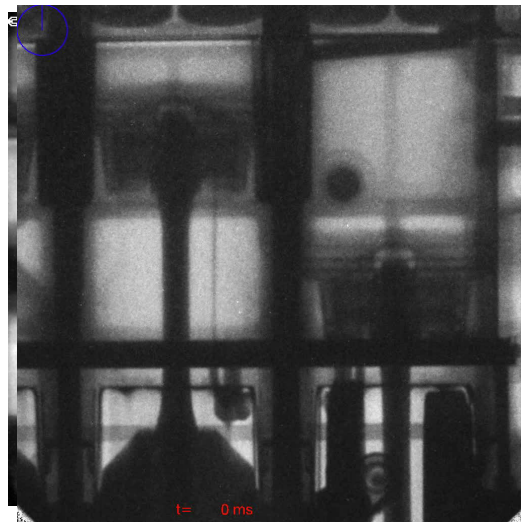
DNA

A protein
molecule
moving along
the DNA chain

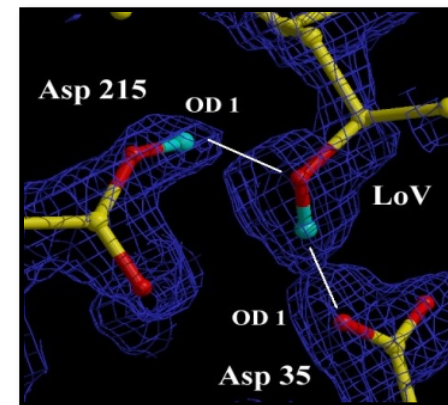
Neutrons and X-rays are complementary



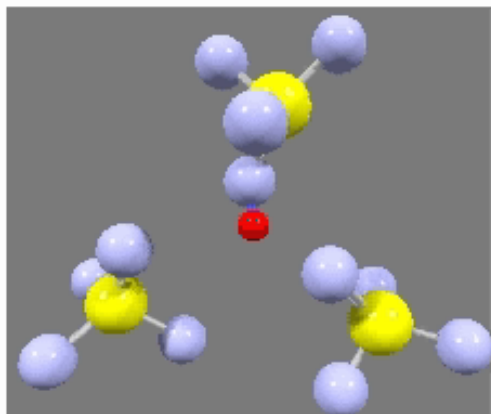
..see magnetic atoms



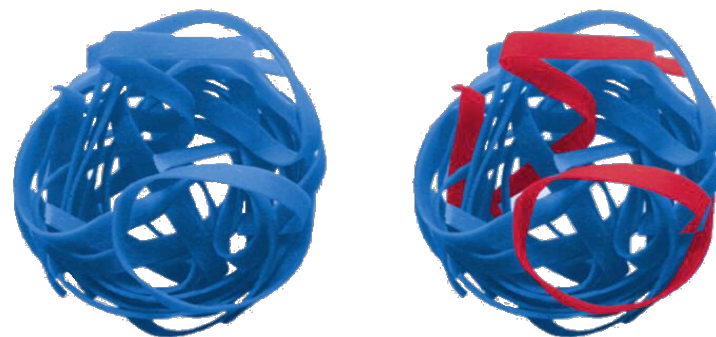
..see inside materials



..see light atoms



..see atoms move



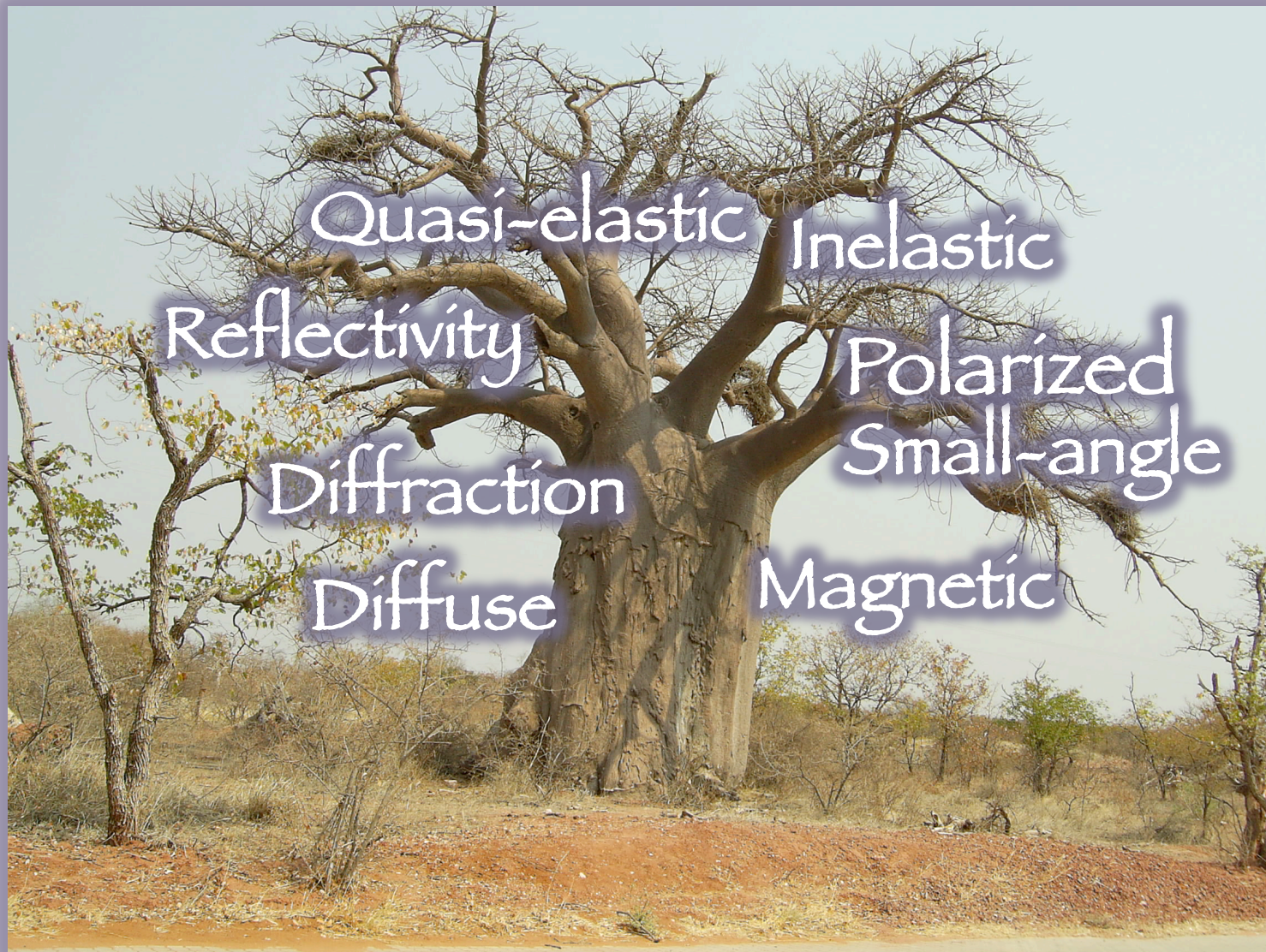
..see isotopes



Scattering and Diffraction



Neutron Scattering Techniques



Isotope specific contrast

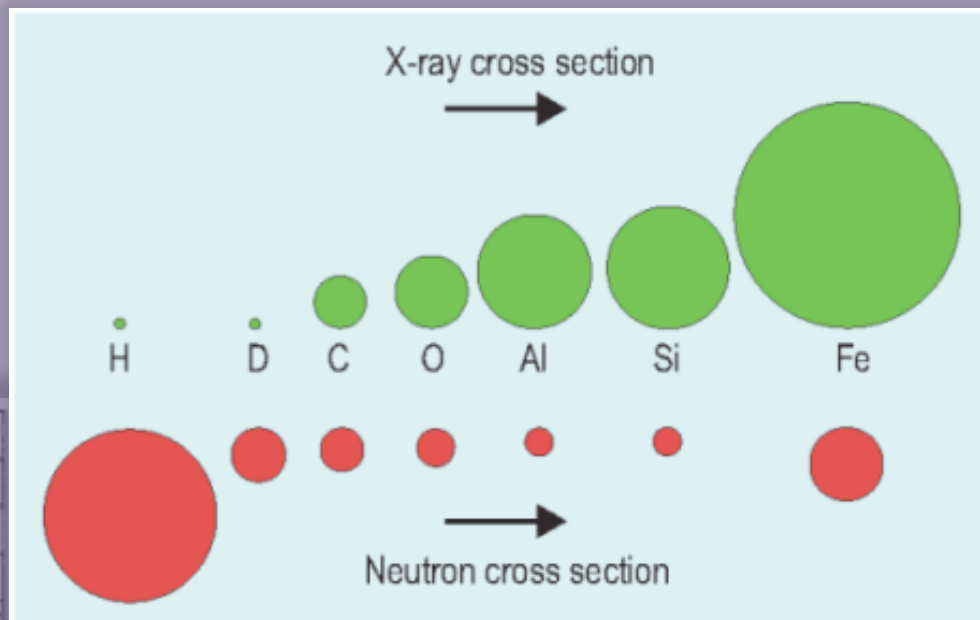
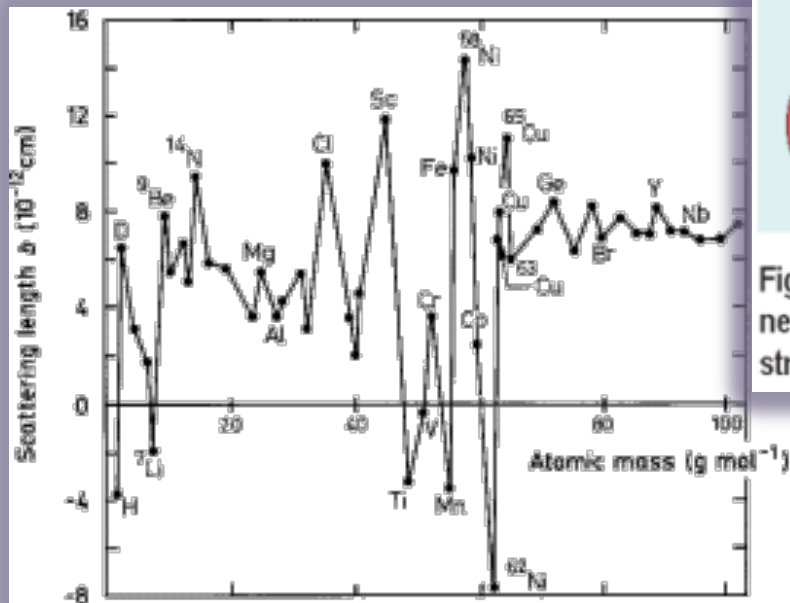
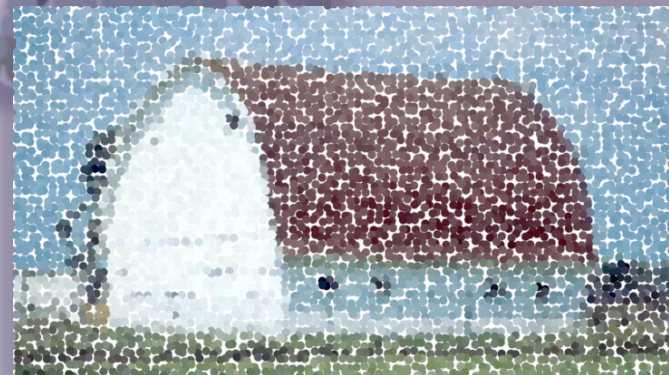
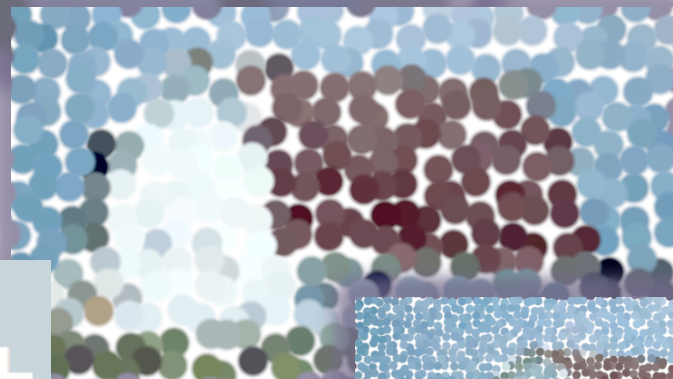
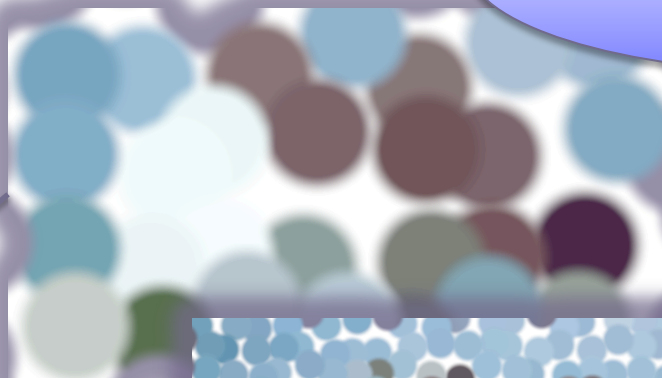
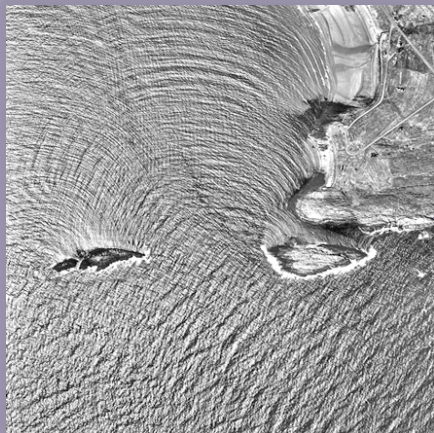


Fig. 2. Neutron and x-ray scattering cross-sections compared. Note that neutrons penetrate through Al much better than x rays do, yet are strongly scattered by hydrogen.



$$\sigma_{coh} = 4\pi \langle b \rangle^2$$

$$\sigma_{incoh} = 4\pi \left(\langle b^2 \rangle - \langle b \rangle^2 \right)$$



WHY NEUTRON SCATTERING IS USEFUL

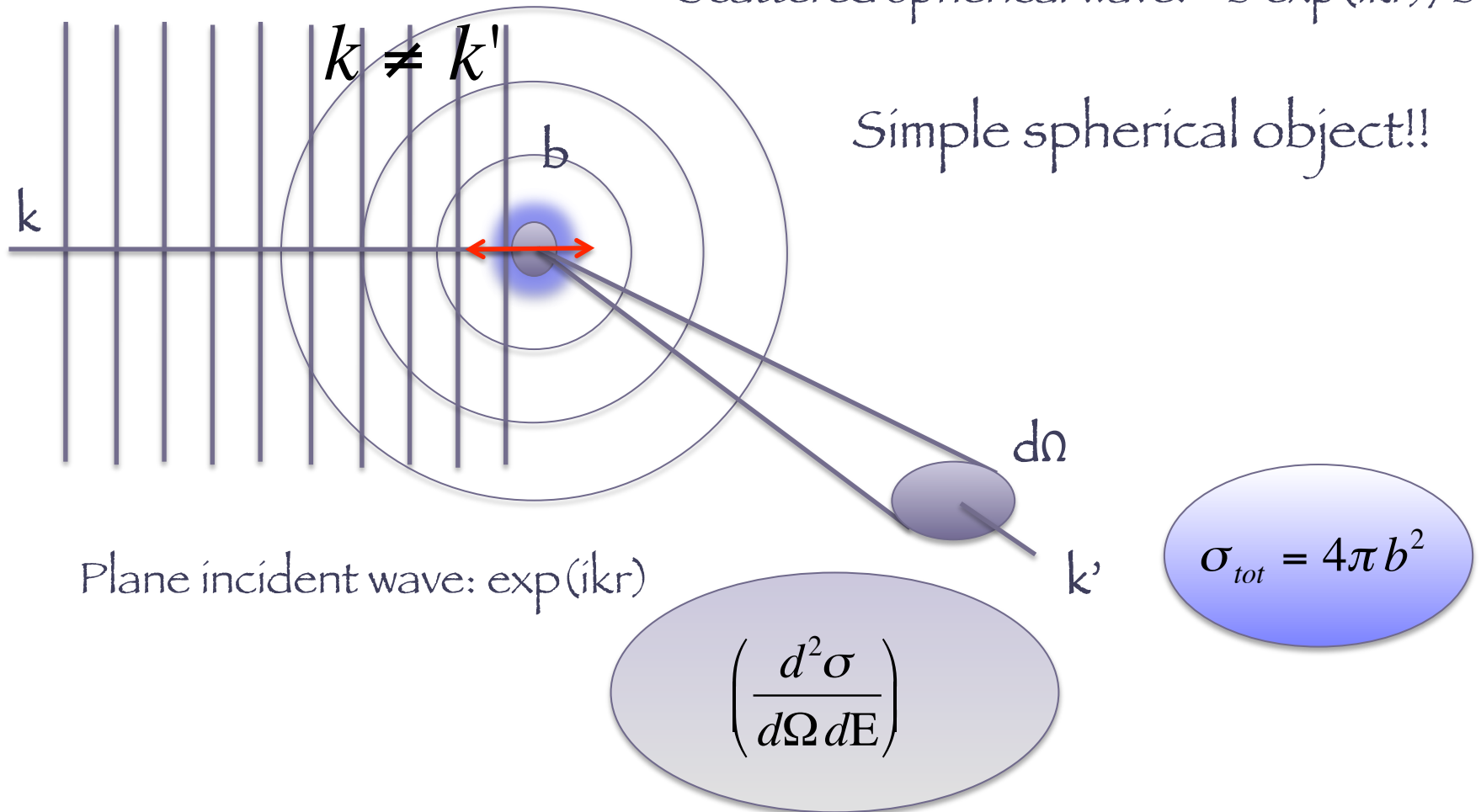
When used as a probe for small samples of materials, neutron beams have the power to reveal what is invisible using other radiations. Neutrons can appear to behave either as particles or as waves or as microscopic magnetic dipoles, and it is these specific properties which enable them to uncover information which is often impossible to access using other techniques.

Isotopes and spin?

Inelastic Scattering of plane wave

σ_{tot} = number of neutrons scattered in all directions per sec/incident flux

Scattered spherical wave: $-b \exp(ikr)/b$



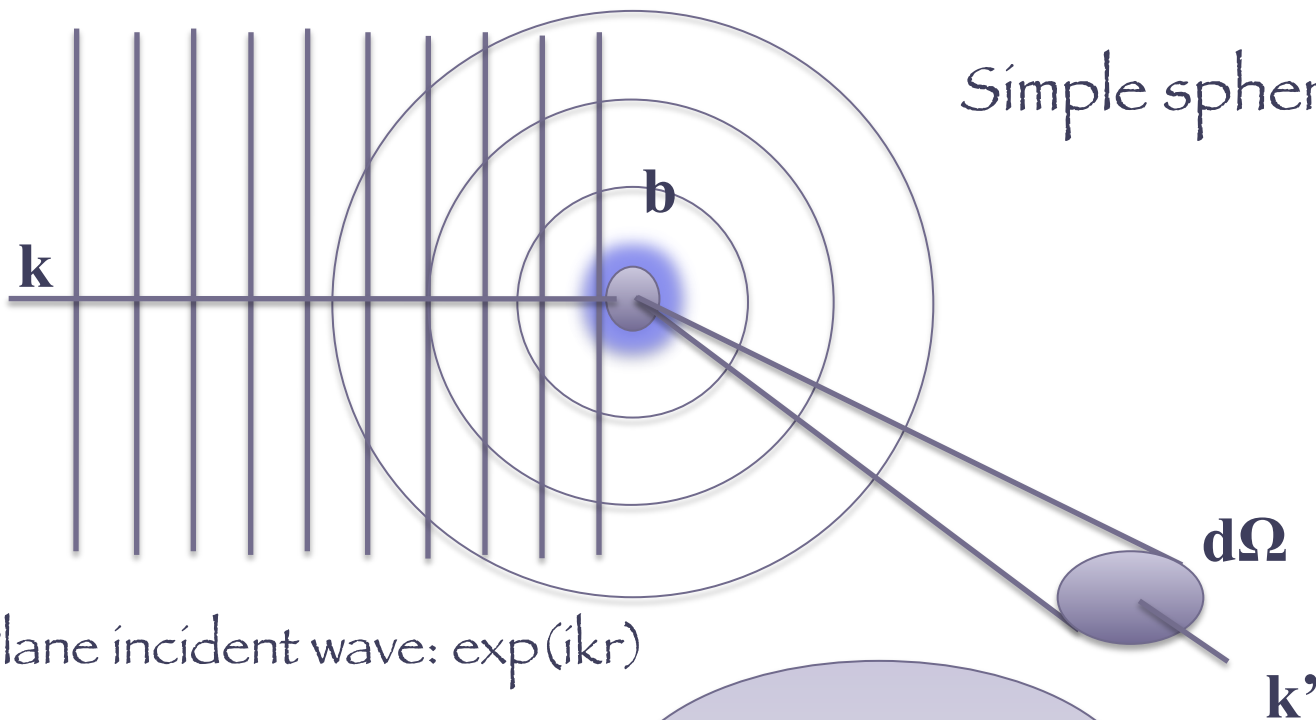


Diffraction = Coherent Elastic Scattering of plane wave

σ_{tot} = number of neutrons scattered in all directions per sec/incident flux

Scattered spherical wave: $-b \exp(ikr)/b$

Simple spherical object!!



Plane incident wave: $\exp(ikr)$

$$\left(\frac{d\sigma}{d\Omega} \right), \left(\frac{d^2\sigma}{d\Omega dE} \right)$$

$$\sigma_{tot} = 4\pi b^2$$

- **Cross section:** $d\sigma = \frac{(\# \text{ particles scattered into solid angle } \Delta\Omega/s)}{(\# \text{ particles incident/sec})(\# \text{ scattering centers/area})}$

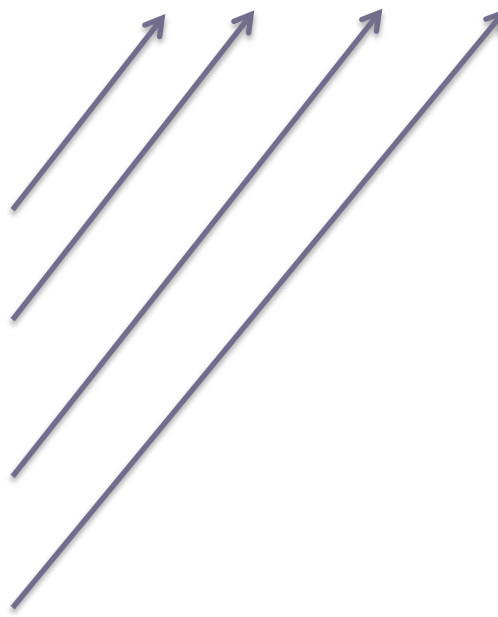
$$\frac{d\sigma(Q)}{d\Omega} = N_p V_p \bar{\rho} P(Q) S(Q)$$

Volume fraction

Contrast

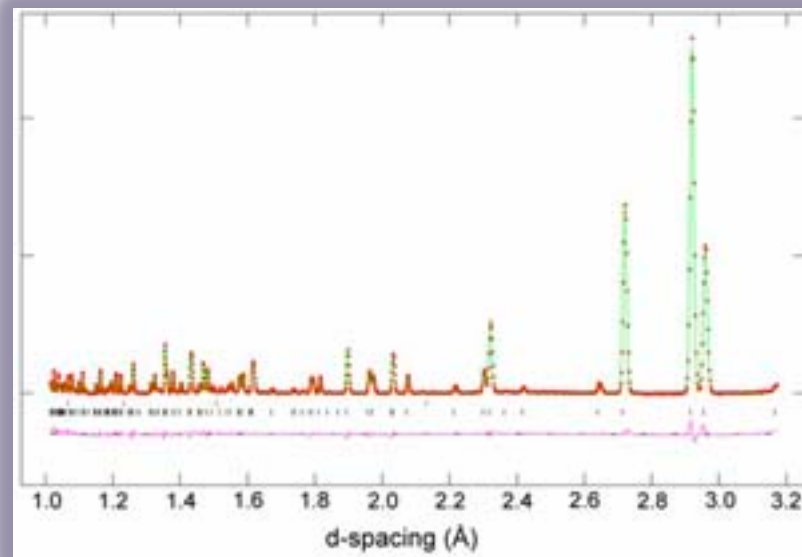
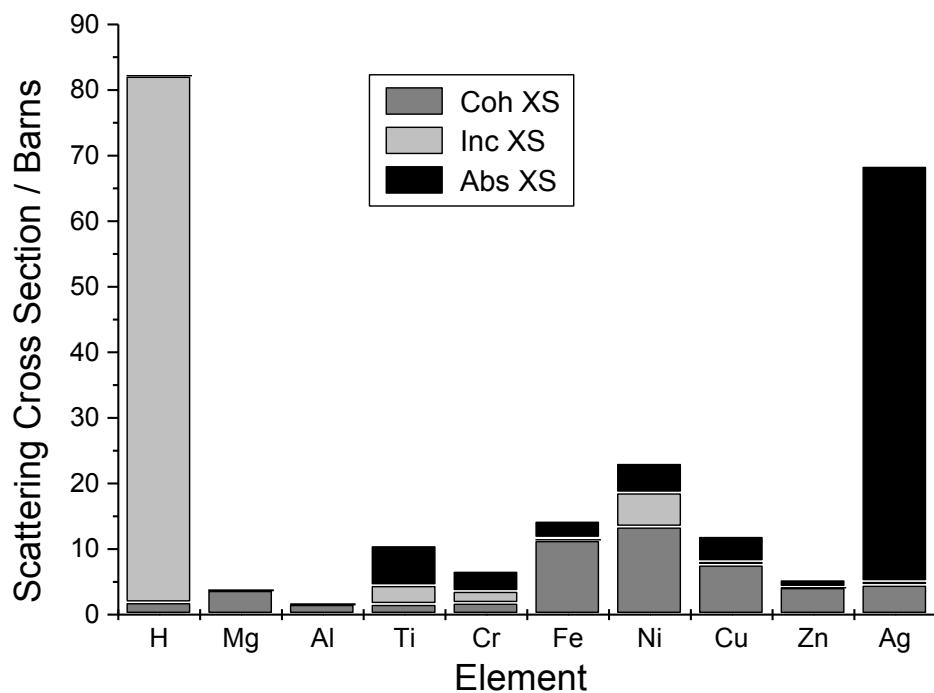
Shape

Interaction



Neutron Scattering XS

Coherent XS ~ Signal
 Incoherent XS ~ Background
 Absorption XS ~ 1/Intensity



The neutron diffraction pattern of sulfuric acid tetrahydrate at 4.2 K as seen by the 90 degree detectors on HRPD, and fitted with the existing X-ray derived structural model for the deuterated species.



NIST
National Institute of
Standards and Technology

NIST Center for Neutron Research

Home ICP Experiments UserProposal Instruments SiteMap

Neutron scattering lengths and cross sections

Neutron scattering lengths and cross sections							
Isotope	conc	Coh b	Inc b	Coh xs	Inc xs	Scatt xs	Abs xs
Fe	---	9.45	---	11.22	0.4	11.62	2.56
54Fe	5.8	4.2	0	2.2			
56Fe	91.7	9.94	0	12.4			
57Fe	2.2	2.3	---	0.66			
58Fe	0.3	15.(7.)	0	28			

				He
C	N	O	F	Ne
Si	P	S	Cl	Ar

Neutron scattering lengths and cross sections							
Isotope	conc	Coh b	Inc b	Coh xs	Inc xs	Scatt xs	Abs xs
Gd	---	6.5-13.82i	---	29.3	151.(2.)	180.(2.)	49700.(125.)
152Gd	0.2	10.(3.)	0	13.(8.)	0	13.(8.)	735.(20.)
154Gd	2.1	10.(3.)	0	13.(8.)	0	13.(8.)	85.(12.)
155Gd	14.8	6.0-17.0i	(+/-)5.(5.)-13.16i	40.8	25.(6.)	66.(6.)	61100.(400.)
156Gd	20.6	6.3	0	5	0	5	1.5(1.2)
157Gd	15.7	-1.14-71.9i	(+/-)5.(5.)-55.8i	650.(4.)	394.(7.)	1044.(8.)	259000.(700.)
158Gd	24.8	9.(2.)	0	10.(5.)	0	10.(5.)	2.2
160Gd	21.8	9.15	0	10.52	0	10.52	0.77

NOTE: The above are only thermal neutron dependent cross sections please go to...

Select the element, and you will get a list of...
Feature section of neutron scattering lengths...
3, 1992, pp. 29-37.



PRL **100**, 250404 (2008)

PHYSICAL REVIEW LETTERS

week ending
27 JUNE 2008

Measurements of the Vertical Coherence Length in Neutron Interferometry

D. A. Pushin,^{1,*} M. Arif,² M. G. Huber,³ and D. G. Cory¹

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²*National Institute of Standards and Technology, Gaithersburg, Maryland, USA*

³*Department of Physics, Tulane University, New Orleans, Louisiana, USA*

(Received 19 March 2008; published 26 June 2008)

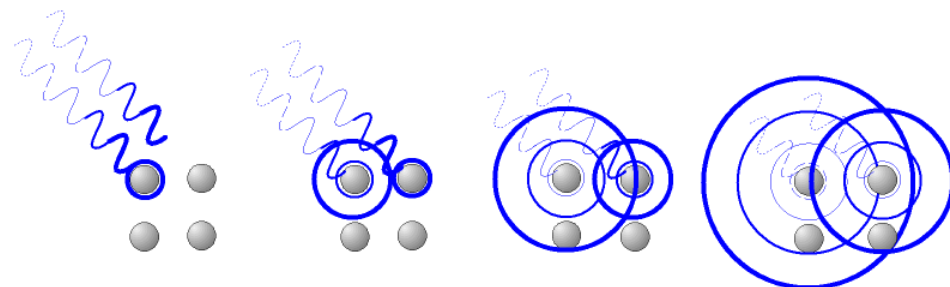
The study and use of macroscopic quantum coherence requires long coherence lengths. Here we describe an approach to measuring the vertical coherence length in neutron interferometry, along with improvements to the NIST interferometer that led to a measured coherence length of 790 Å. The measurement is based on introducing a path separation and measuring the loss in contrast as this separation is increased. The measured coherence length is consistent with the momentum distribution of the neutron beam. Finally, we demonstrate that the loss in contrast with beam displacement in one leg of the interferometer can be recovered by introducing a corresponding displacement in the second leg.

DOI: [10.1103/PhysRevLett.100.250404](https://doi.org/10.1103/PhysRevLett.100.250404)

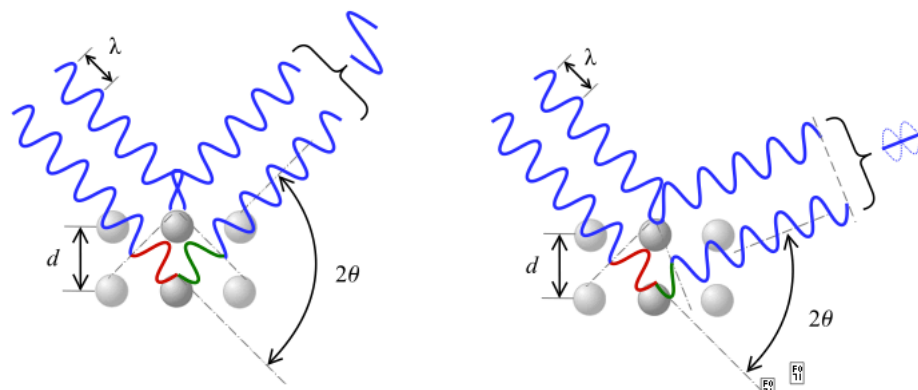
PACS numbers: 03.75.Dg, 03.65.-w, 42.50.-p

Diffraction - Bragg

Diffraction of X-rays or neutrons by polycrystalline samples is one of the most important, powerful and widely used analytical techniques available to materials scientists. For most crystalline substances of technological importance, the bulk properties of a powder or a polycrystalline solid, averaged throughout the sample, are required; in general a single-crystal data, even if they can be obtained, are usually of little interest except for determination of the crystal structure or for studying some other fundamental physical property. By *J Ian Langford and Daniel Louer*



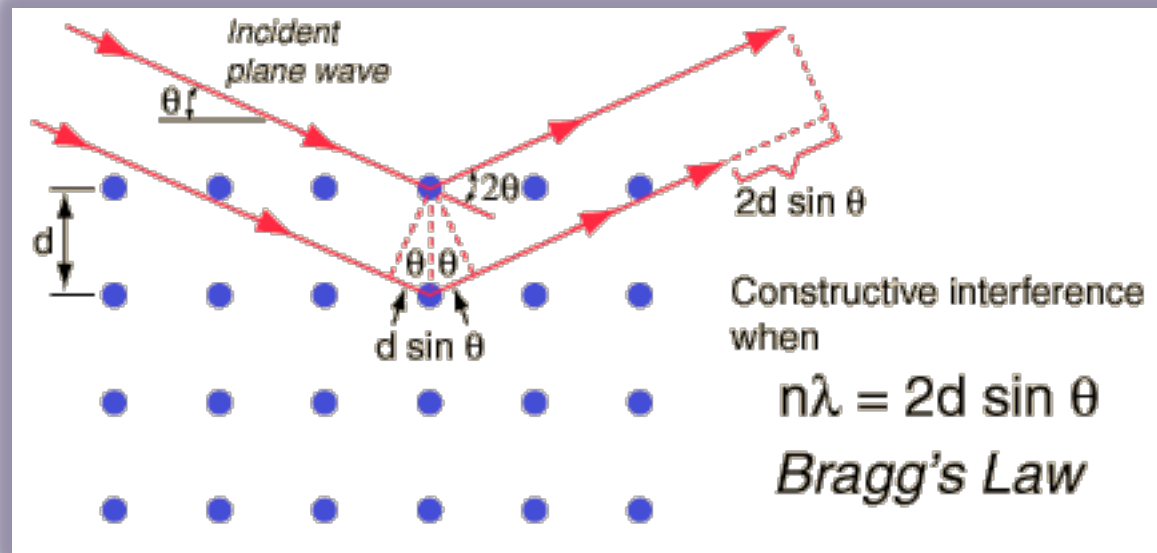
X-rays interact with the atoms in a crystal.



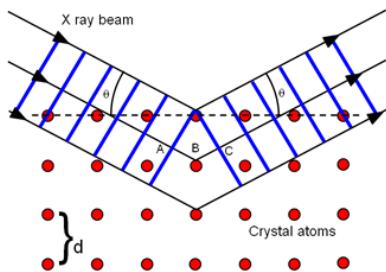
According to the 2θ deviation, the phase shift causes constructive (left figure) or destructive (right figure) interferences.

- Bragg / Laue scattering
- Coherent elastic scattering

Two beams with identical wavelength and phase approach a crystalline solid and are scattered off two different atoms within it. The lower beam traverses an extra length of $2d \sin \theta$. **Constructive interference** occurs when this length is equal to an integer multiple of the wavelength of the radiation.



Using the grains as internal strain gauges

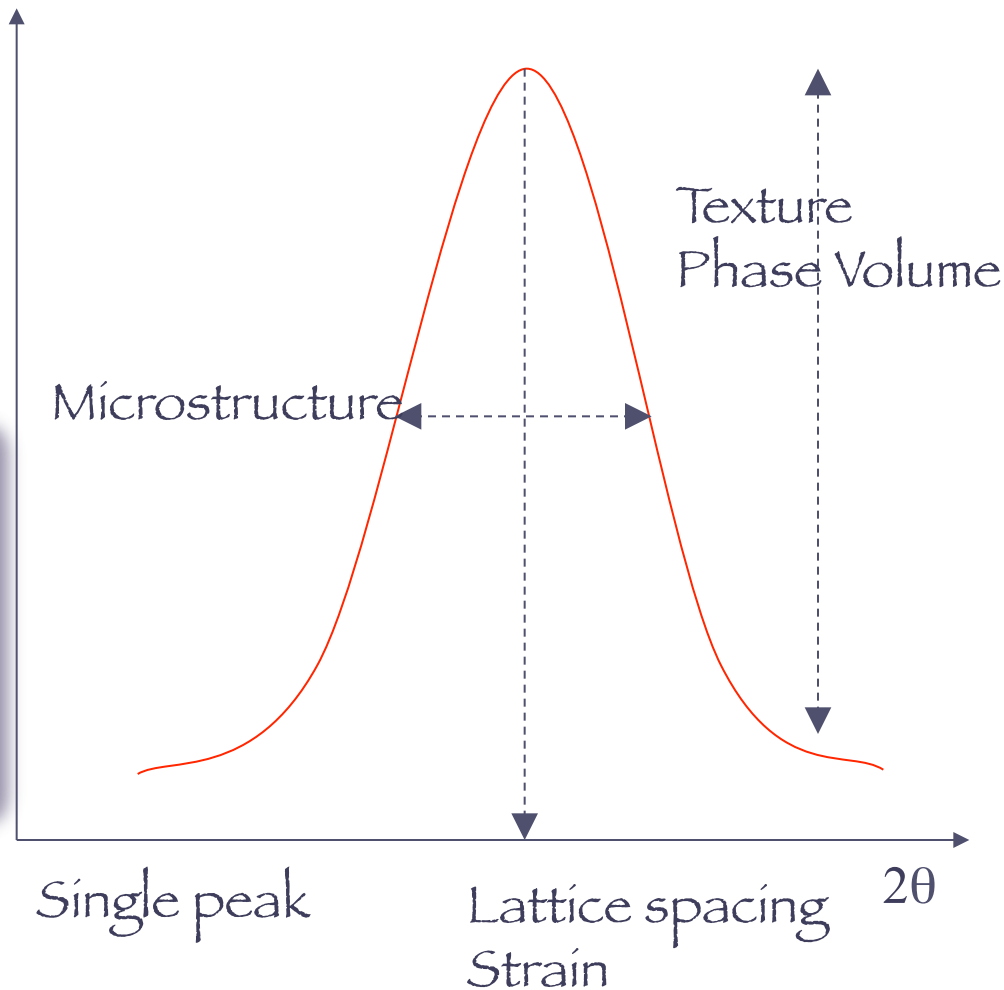


$$\lambda = 2d \sin \theta$$

$$\varepsilon = -\cot(\theta) (\theta - \theta_0)$$

Two ways to measure d :

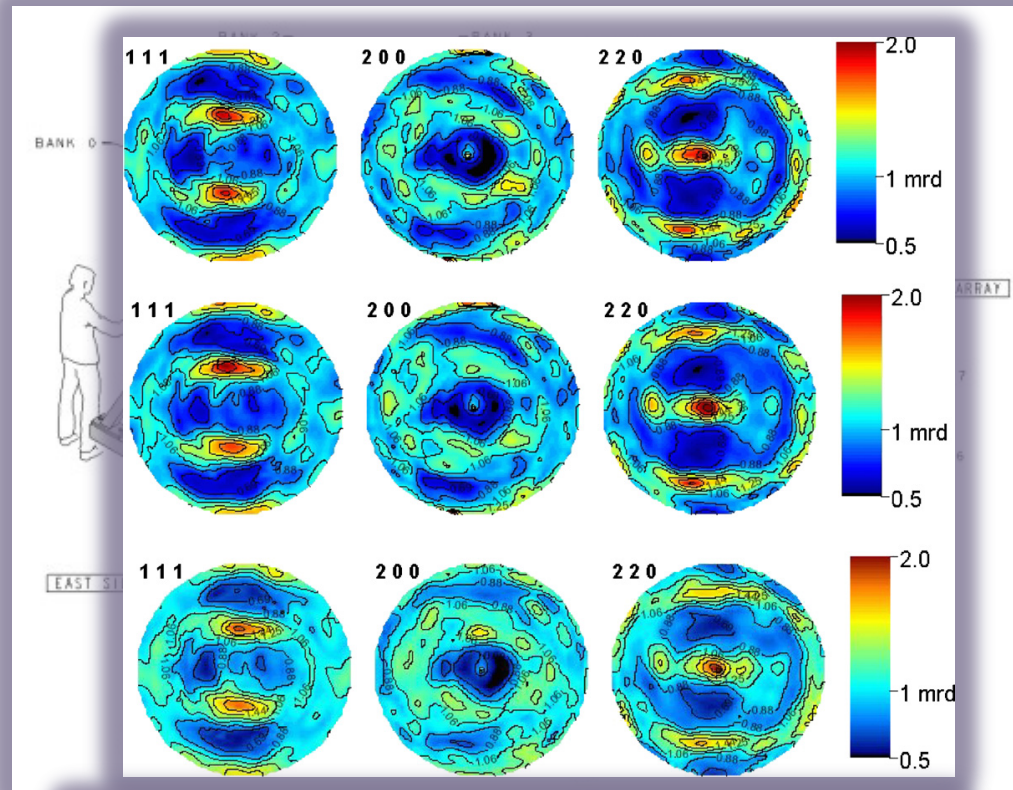
- keep λ fixed and measure θ constant wavelength
- keep θ fixed and measure λ time-of-flight



- ISIS: GEM instrument
- Near 4π coverage

Example:

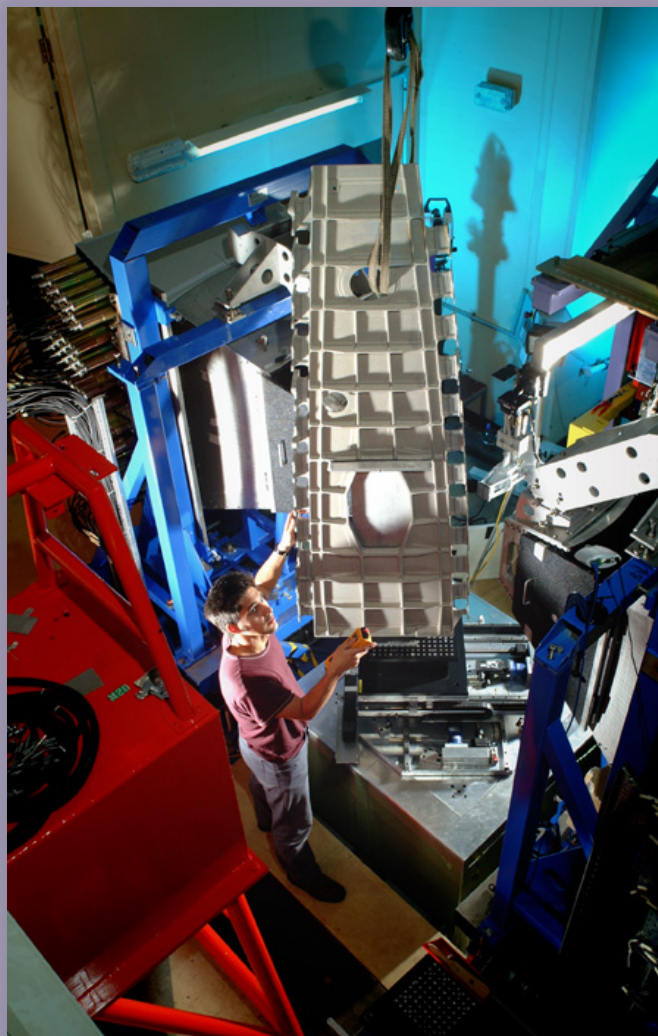
- Cold rolled copper, simulating manufacturing process for archeometry
- 2mm thick disks
- $20 \times 20 \text{ mm}^2$ beam
- 2 min counting times



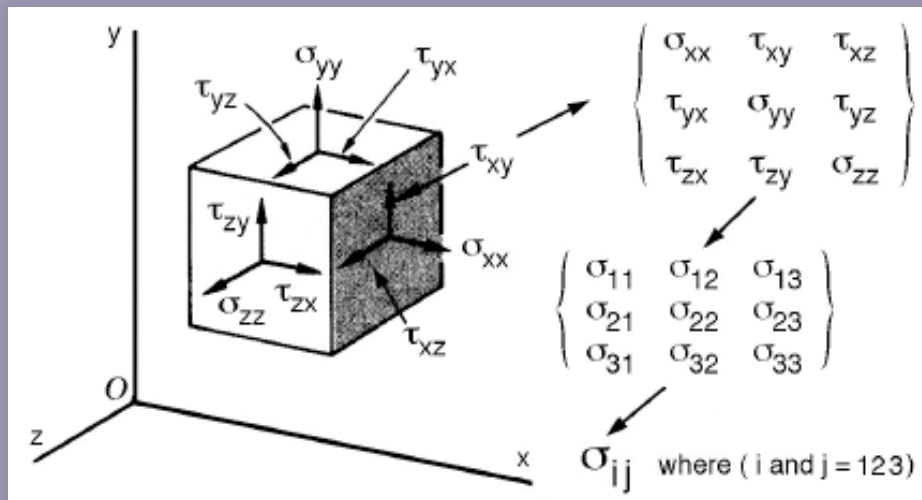
$$\frac{P_{hkl}}{P_0} = \frac{\lambda^3 l}{4\pi r} \frac{h\rho'}{\rho} \frac{e^{-\mu_h \sec\theta}}{\sin^2 2\theta} j_{hkl} N^2 F_{hkl}^2, \quad (15)$$

Intensity of diffraction peak

Diffraction: Stress and Strain



CD - Neutrons Sources



Applications:

- Residual stresses
- Fatigue/Structural Integrity
- Welds
- Alloy development
- Microstructure/Texture
- Phase transformation
- In-situ experiments

Example - Diffraction

PHYSICAL REVIEW

VOLUME 73, NUMBER 8

APRIL 15, 1948

The Diffraction of Neutrons by Crystalline Powders

E. O. WOLLAN AND C. G. SHULL

Oak Ridge National Laboratory, Oak Ridge, Tennessee

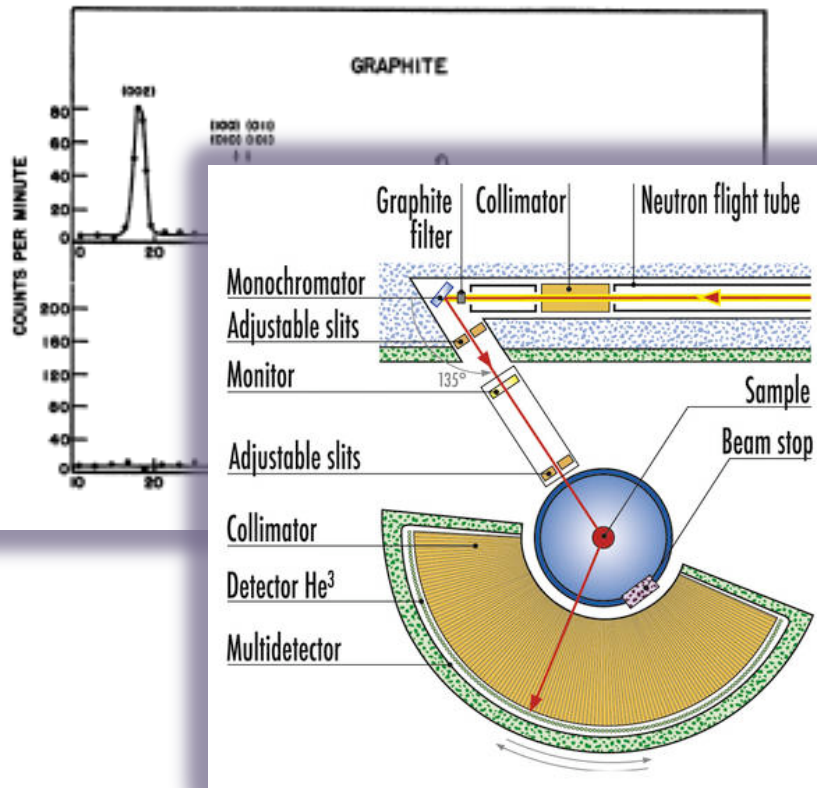


FIG. 3. Powder diffraction patterns for diamond and graphite. The major part of the diffuse scattering in these patterns arises from multiple scattering in the samples.

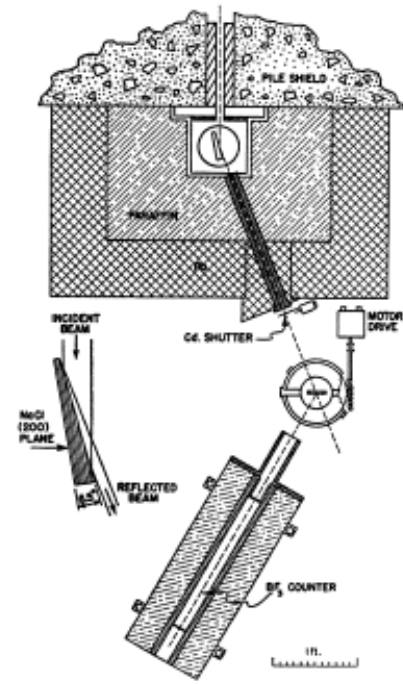


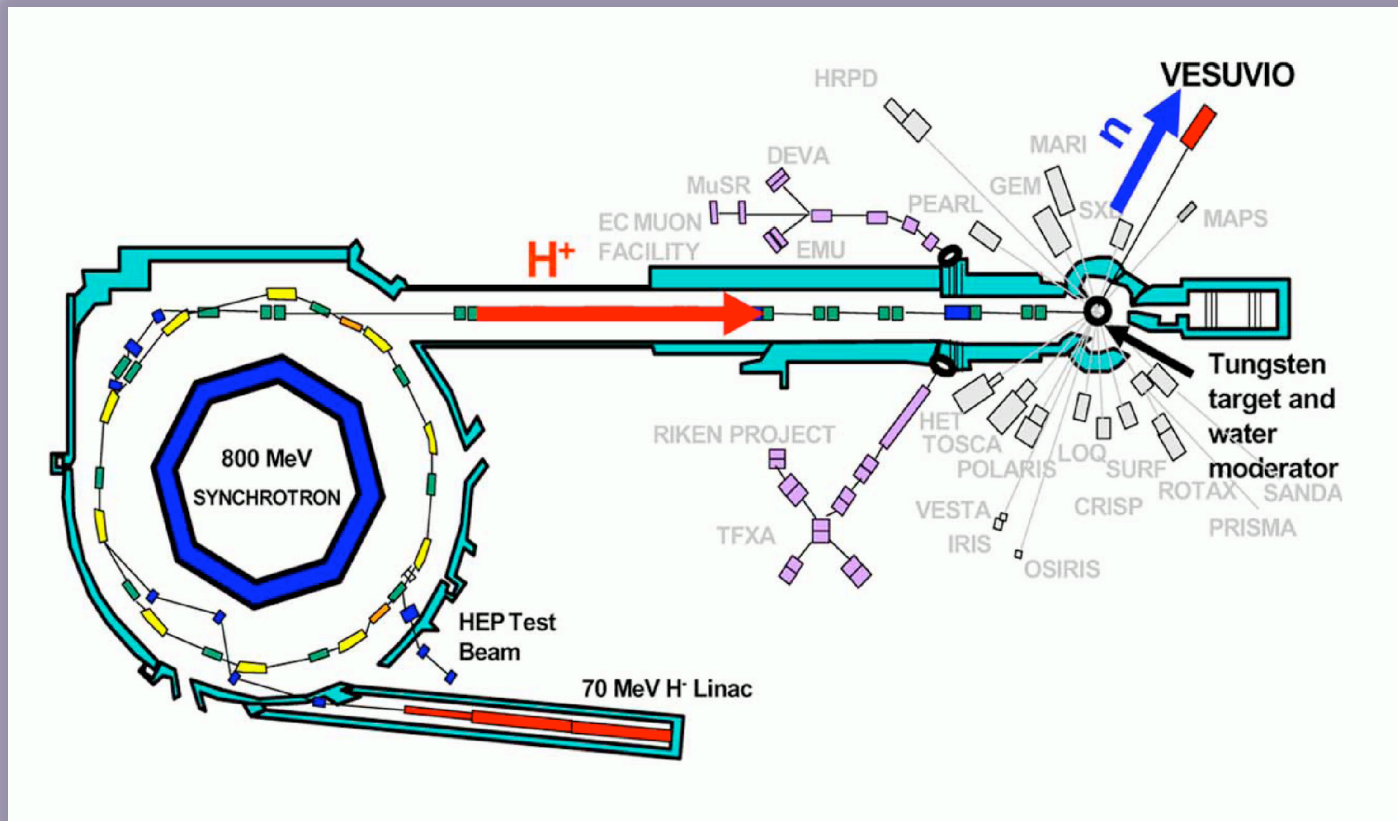
FIG. 1. Arrangement of apparatus, showing the monochromating crystal (detailed in left center) collimating slits, shielding, second spectrometer with location of powder specimen and counter.

$$\frac{P_{hkl}}{P_0} = \frac{\lambda^3 l}{4\pi r} \frac{h\rho'}{\rho} \frac{e^{-\mu h \sec\theta}}{\sin^2 2\theta} j_{hkl} N^2 F_{hkl}^2, \quad (15)$$

Intensity of diffraction peak
D2B at the ILL, Grenoble, 50 years later - July 31th, 2012



Spallation Sources: Time of Flight



$$\lambda = \frac{h}{p} = \frac{h}{mv} = \frac{ht}{mL}$$

De Broglie
+

$$\lambda = 2d \sin \theta$$

Bragg =

Time-of-Flight:

$$d = \frac{h}{2mL \sin \theta} t$$

Fast, short-wavelength neutrons arrive earlier at detector!

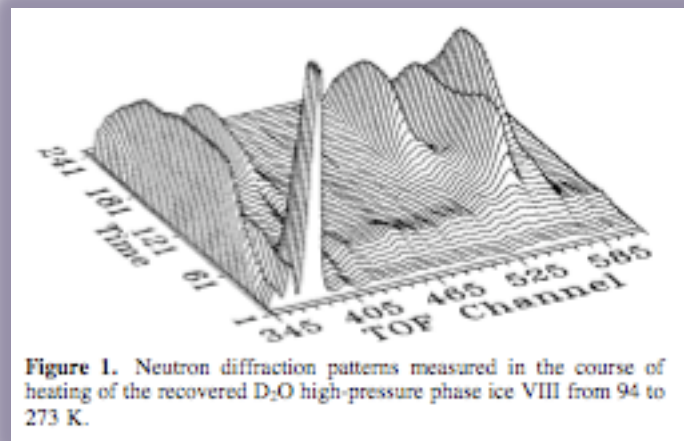
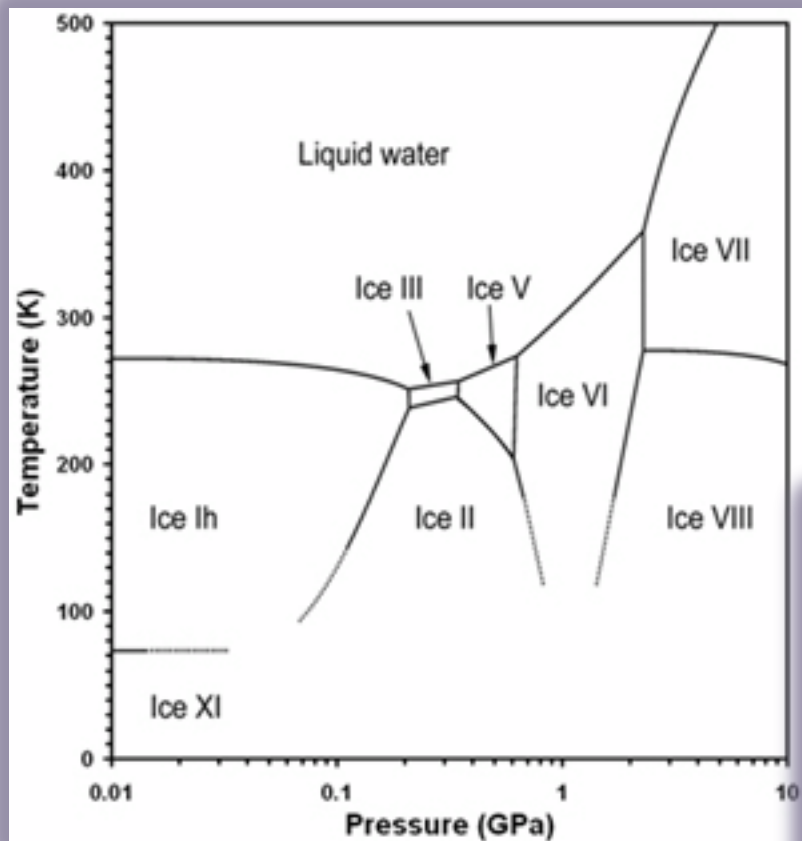


Figure 1. Neutron diffraction patterns measured in the course of heating of the recovered D₂O high-pressure phase ice VIII from 94 to 273 K.

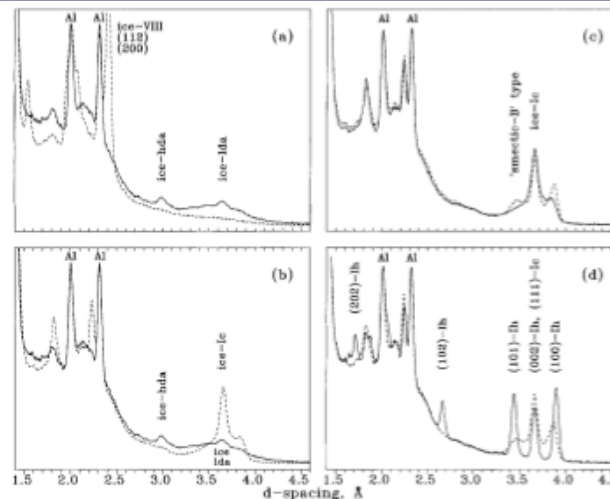


Figure 2. Neutron diffraction patterns of D₂O ices before and after phase transitions: (a) from ice VIII (dashed line) to a mixture of Ih + Ic ices (solid line) at 125 K; (b) from the mixture of Ih + Ic ices (solid line) to ice Ic (dashed line) at 160 K; (c) from ice Ic (solid line) to the hexagonal phase, "smectic B" type structure (dashed line), at 190 K; (d) from "smectic B" type ice (dashed line) to ice Ih (solid line) at 230 K. The diffraction patterns also contain peaks from the aluminum sample can and cryostat.

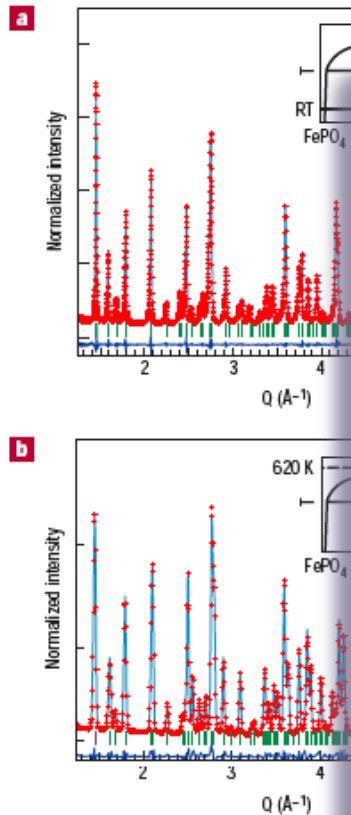


Figure 2 Neutron diffraction patterns measured at the $\text{FePO}_4\text{-LiFePO}_4$ binary phase diagram. a,b, Rietveld fit and time-of-flight neutron diffraction profile measured for LiFePO_4 at room temperature (a) and the angle-dispersive neutron diffraction profile measured at 620 K (b). Two different neutron diffractometers were used for each measurement as explained in the Supplementary Information. The data points are plotted using the common scale $Q = 4\pi \sin \theta / \lambda$ in the Q range for VEGA and HERMES for comparison. Specific composition and temperature are given in the inset phase diagram. Observed intensity I_{obs} and the fit I_{fit} are represented by red plus signs and the green curve at the bottom represents the residual difference between the fit and the data. The refined parameters are summarized in Supplementary Information. An impurity phase was identified, and the crystal structure was determined with the space group $Pnma$.

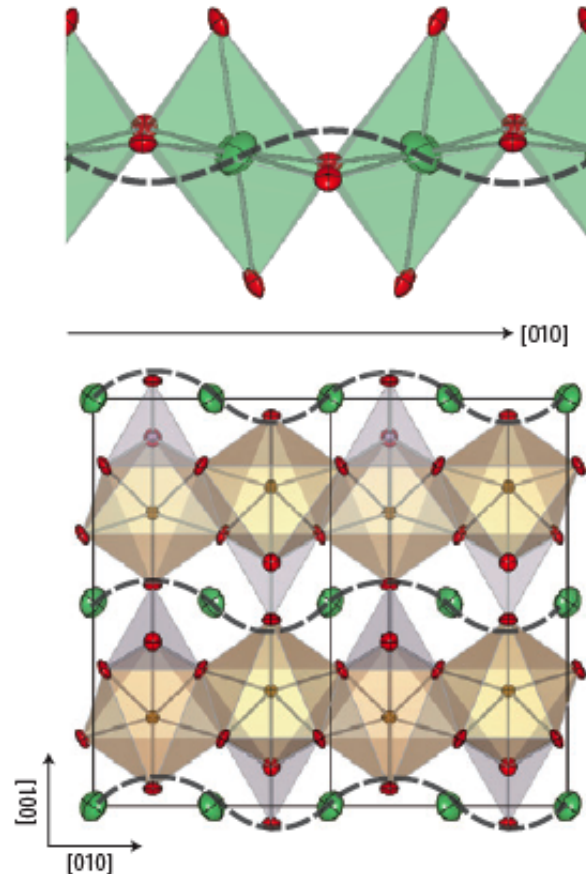


Figure 3 Anisotropic harmonic lithium vibration in LiFePO_4 , shown as green thermal ellipsoids and the expected diffusion path. The ellipsoids were refined with 95% probability by Rietveld analysis for room-temperature neutron diffraction data. Expected curved one-dimensional continuous chains of lithium motion are drawn as dashed lines to show how the motions of Li atoms evolve from vibrations to diffusion.

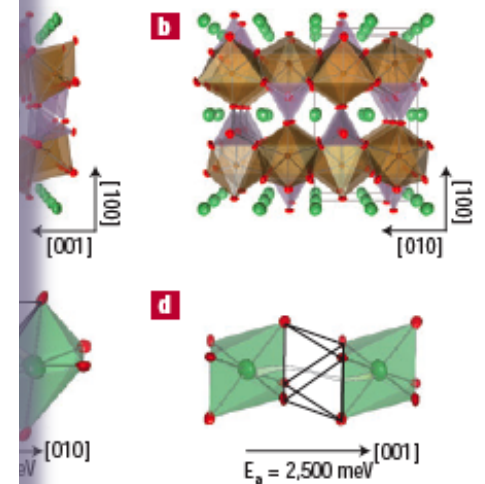


Figure 4 Crystal structure of LiFePO_4 and possible lithium pathways. a,b, The structure is projected along the $[010]$ (a) and $[001]$ (b) directions. The lithium pathways are parallel to these directions. The structure parameters obtained through this work and the crystallographic data are given in Supplementary Information, Table S1. The structure can be described as a layered structure with a hexagonal close-packed oxygen sub-array, in which Li, Fe and P occupy the octahedral sites to form (1) corner-sharing FeO_6 octahedra that form a distorted two-dimensional square lattice perpendicular to the c axis, and (2) LiO_6 octahedra aligned in parallel chains along the c axis, with the LiO_6 groups connecting neighbouring planes or arrays. The thermal ellipsoids indicate Li, Fe, P and O atoms, respectively. The dashed lines indicate the lithium migration paths: c, along the $[010]$ direction through the octahedral sites; and d, along the $[001]$ direction through the tetrahedral sites. One-dimensional diffusion along the $[010]$ direction was observed experimentally and confirmed by the computational method^{15,16}.

Electric Circuit
(40% - 60% efficiency)



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Solid State Ionics 177 (2006) 2357–2362

**SOLID
STATE
IONICS**

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Proton co



Table 1

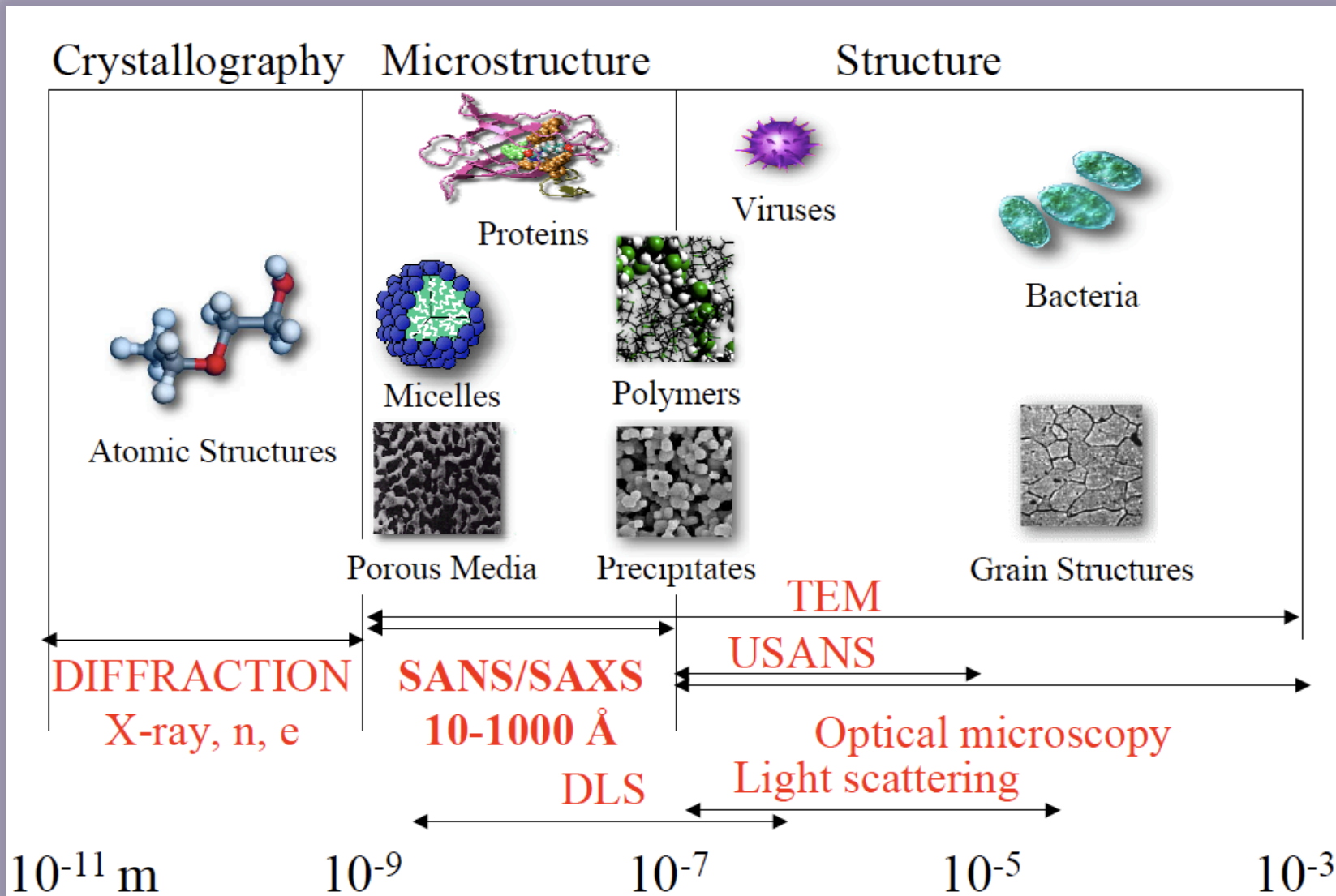
Summary of results obtained from Rietveld analysis of neutron powder diffraction data for $\text{BaZr}_{1-x}\text{In}_x\text{O}_{3-0.5x}$, collected at 10K on the NPD diffractometer

	$x=0.0$ (as-prepared)	$x=0.0$ (deuterated)	$x=0.25$ (as-prepared)	$x=0.25$ (deuterated)	$x=0.50$ (as-prepared)	$x=0.50$ (deuterated)
Space group	Pm-3m	Pm-3m	Pm-3m	Pm-3m	Pm-3m	Pm-3m
a (Å)	4.1879(1)	4.1880(1)	4.1916(1)	4.1983(1)	4.1942(2)	4.2260(4)
Thermal parameters, B_{iso} (Å ²)						
Ba on 1(b) 1/2,1/2,1/2	0.09(3)	0.08(2)	0.10(3)	0.21(4)	0.48(6)	0.97(10)
Zr on 1(a) 0,0,0	0.14(2)	0.11(2)	0.19(3)	0.23(3)	0.43(5)	0.37(9)
In on 1(a) 0,0,0	–	–	0.19(3)	0.23(3)	0.43(5)	0.37(9)
O on 3(d) 1/2,0,0	0.24(2)	0.26(2)	0.63(2)	0.56(2)	1.16(4)	1.17(5)
Oxygen site occupancy	2.98(1)	2.99(1)	2.82(1)	2.90(2)	2.68(2)	2.98(3)
χ^2	1.96	1.99	1.85	1.64	1.72	3.16
Weighted R_{wp}	4.17%	5.16%	4.7%	4.74%	5.28%	6.62%
Bragg R_{B}	3.97%	3.86%	4.63%	4.40%	8.16%	6.62%
Fitted parameters	16	16	16	16	16	16

2θ (°)

Fig. 2. Low temperature (10K) neutron powder diffraction patterns of as-prepared and respective deuterated (marked with D) $\text{BaZr}_{1-x}\text{In}_x\text{O}_{3-\delta}$ ($x=0.00, 0.25$ and 0.50) samples.

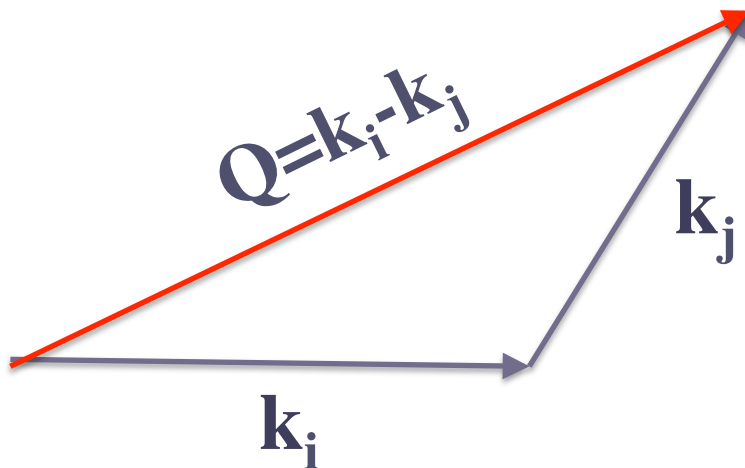
Small Angle Neutron Scattering



- Scattering Vector

$$\sin \theta = Q \lambda / 4\pi = \lambda / 2d$$

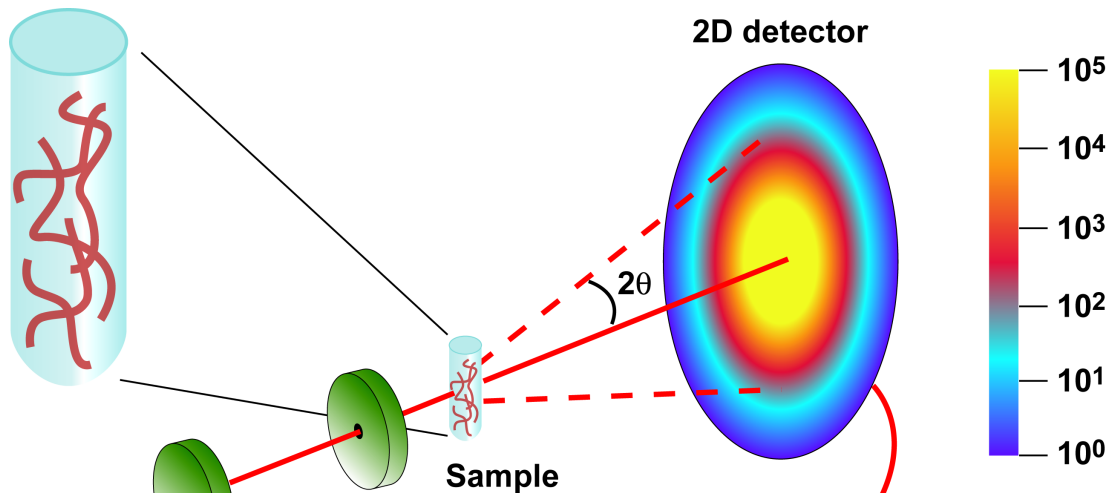
$$\rightarrow d = 2\pi / Q$$



Small angle \rightarrow small $Q \rightarrow$ large distances
 long λ (cold neutrons) $\approx 6 \text{ \AA}$, $d = 200 \text{ \AA}$, $2\theta \approx 1.7^\circ$



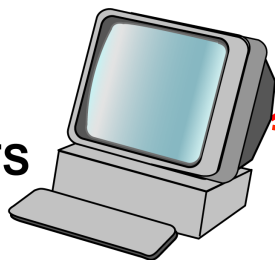
SANS: Experimental Setup



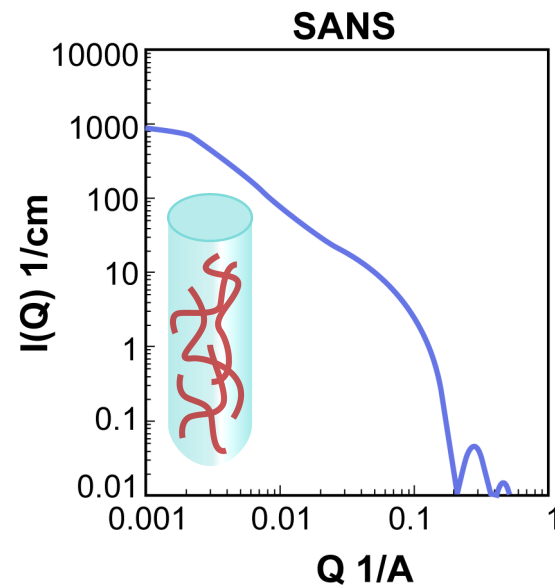
Measures structures at length scales from 1 nm to 100 nm

Beam:
Neutron (SANS)
or X-ray (SAXS)

RESULTS

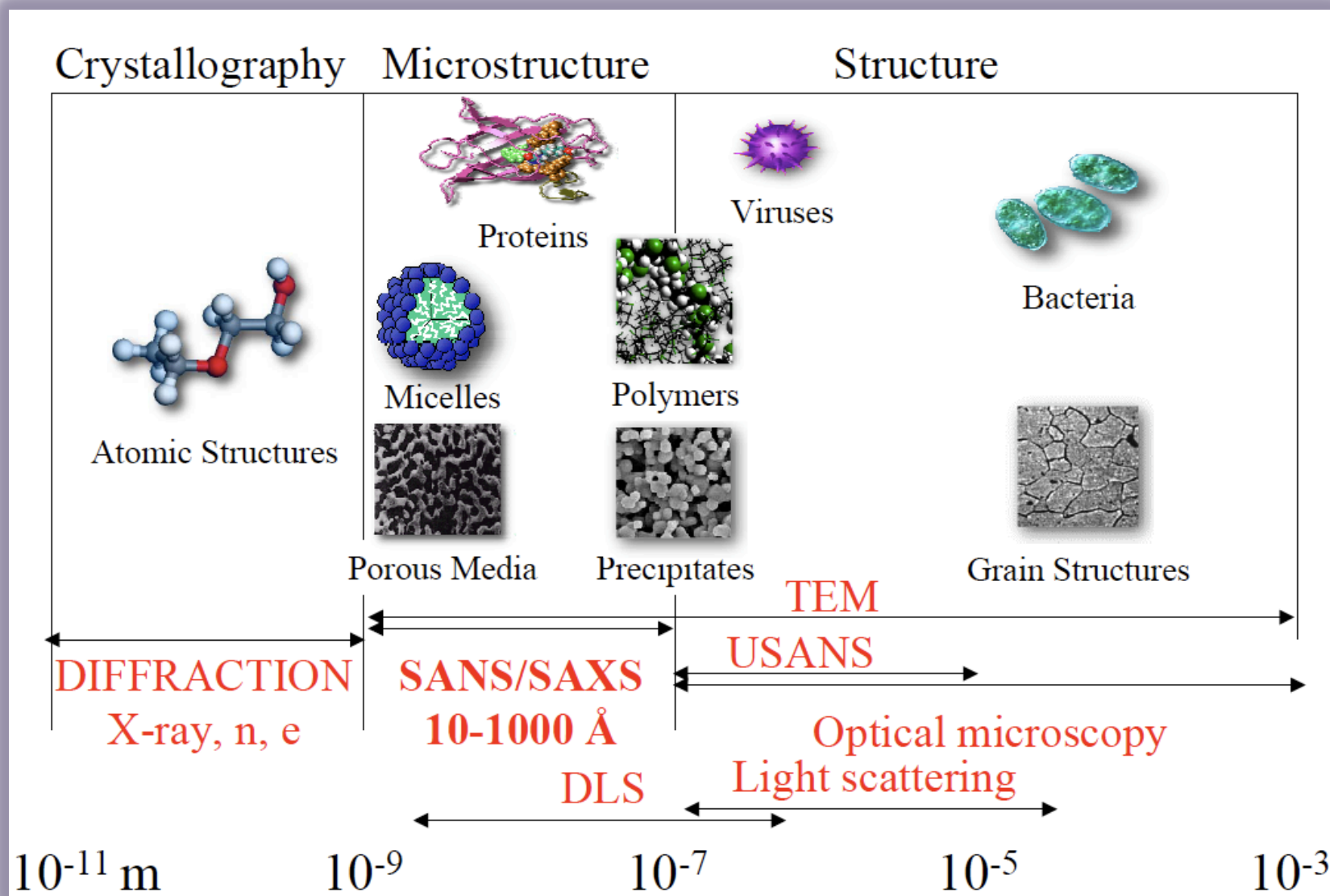


$$|Q| = \frac{4\pi \sin \theta}{\lambda}$$



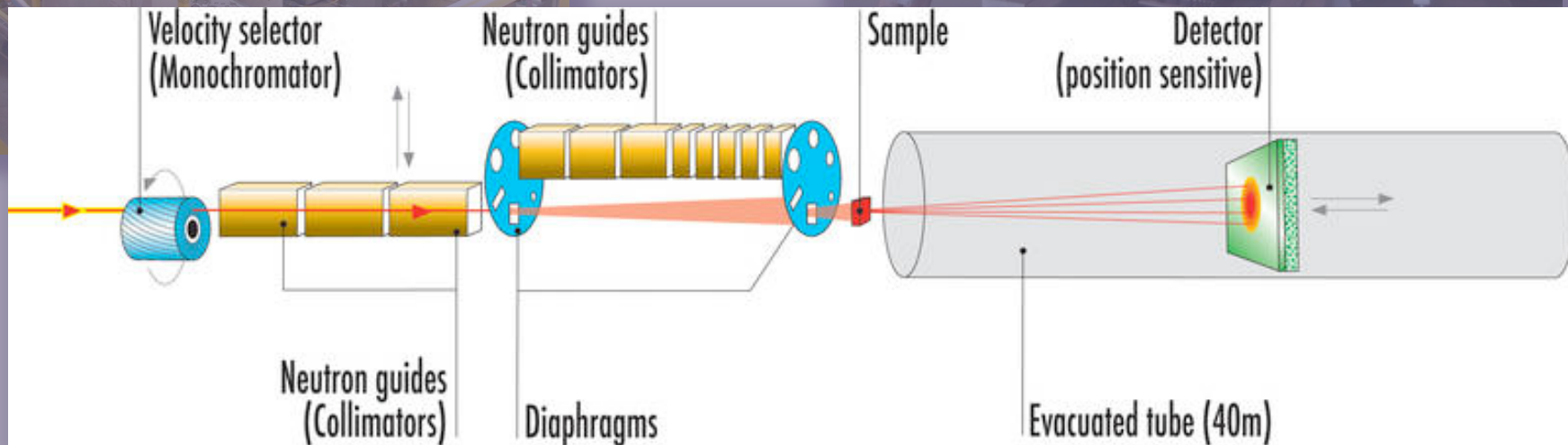
$$\Theta_{\min} \approx 0.03^\circ, \quad \Theta_{\max} \approx 3^\circ$$

SANS: applications





Small Angle Instruments



- **Cross section:** $d\sigma = \frac{(\# \text{ particles scattered into solid angle } \Delta\Omega/s)}{(\# \text{ particles incident/sec})(\# \text{ scattering centers/area})}$

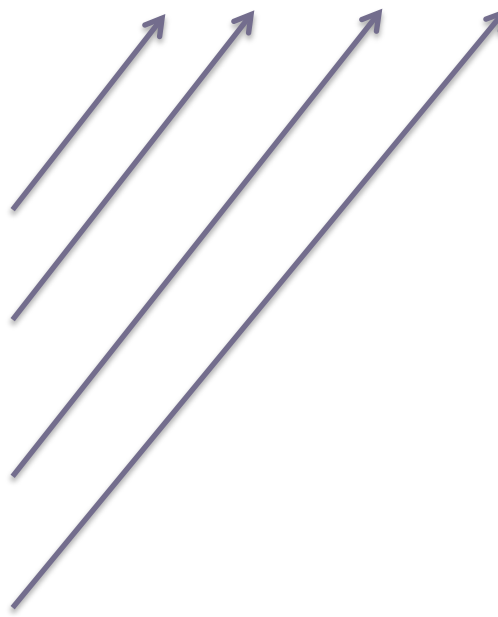
$$\frac{d\sigma(Q)}{d\Omega} = N_p V_p \bar{\rho} P(Q) S(Q)$$

Volume fraction

Contrast

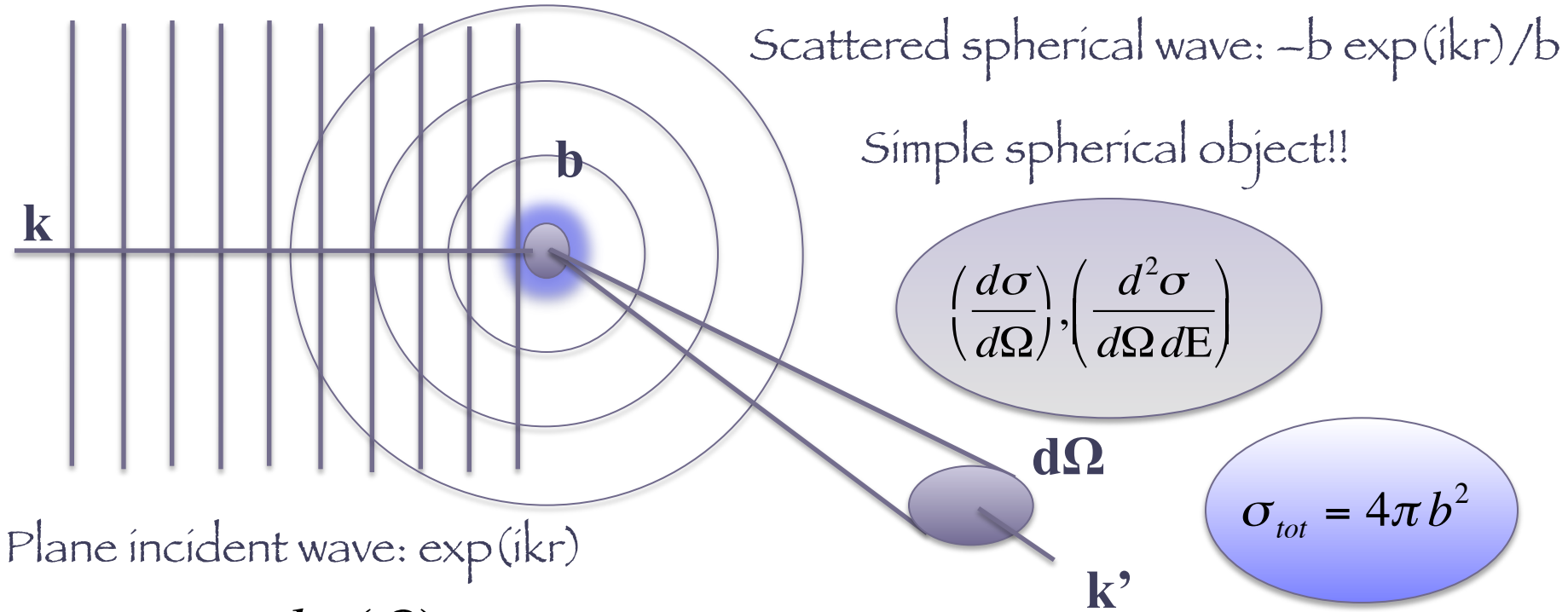
Shape

Interaction



SANS: Scattering of plane wave

σ_{tot} = number of neutrons scattered in all directions per sec/incident flux

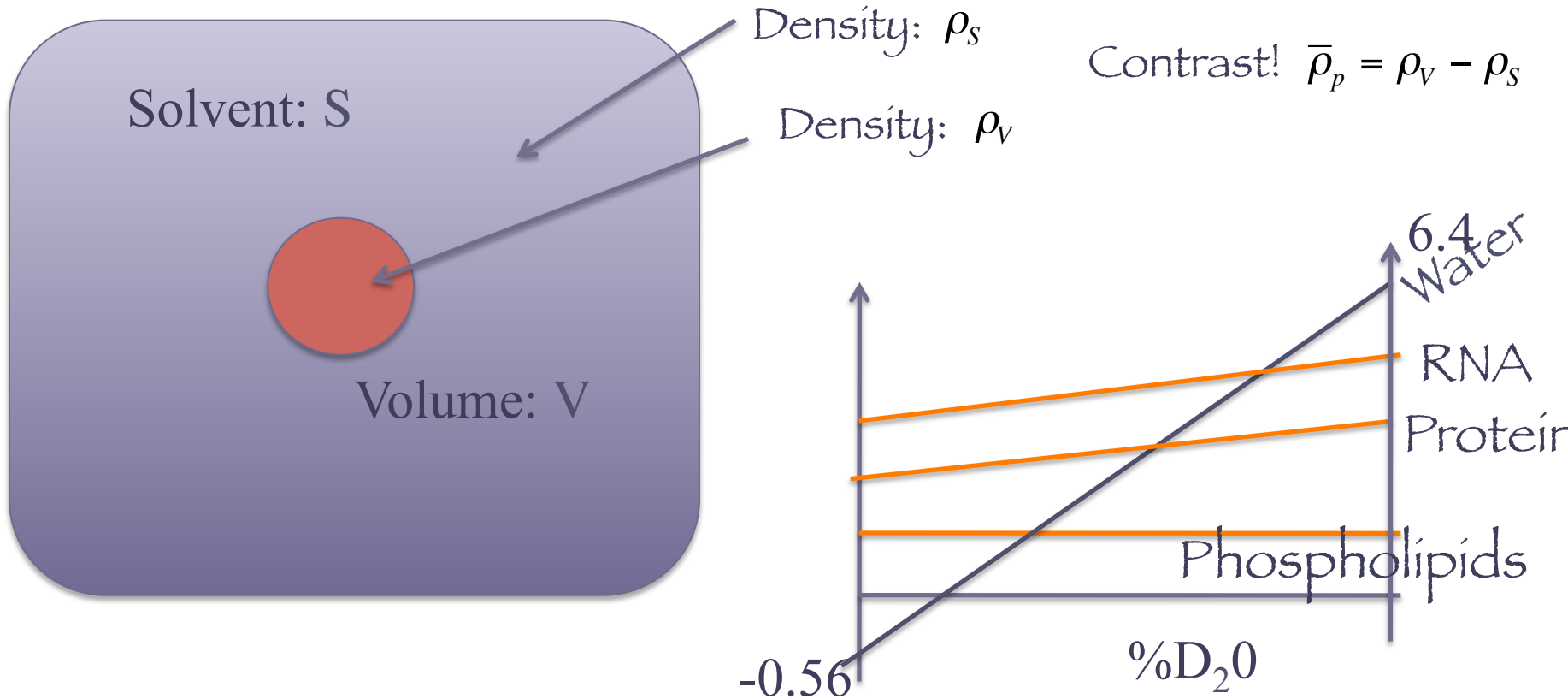


$$\frac{d\sigma(Q)}{d\Omega} = N_p V_p \bar{\rho} P(Q) S(Q)$$

Volume fraction

Contrast

SANS: Particles: contrast!

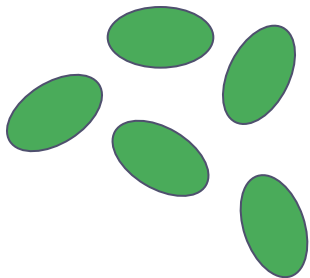


SANS: quick example


Protein based drugs:

- Typically proteins in solution to be injected
- Long shelf life (up to 2 years)
- Control of release profile is desirable

Fast action: Monomeric and dimeric insulin



10 nm

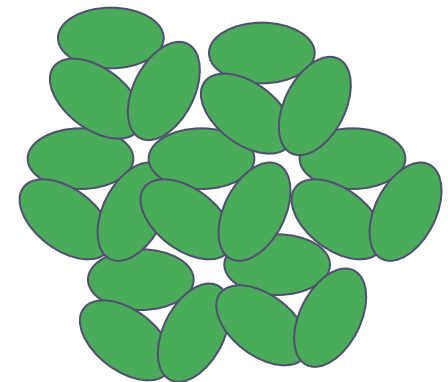
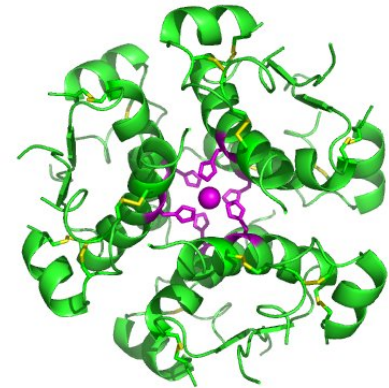


Medium action: Hexameric insulin



Slow action: Large complexes of hexameric insulin

Insulin Hexamer



Knowledge and control of solution properties of the proteins are crucial

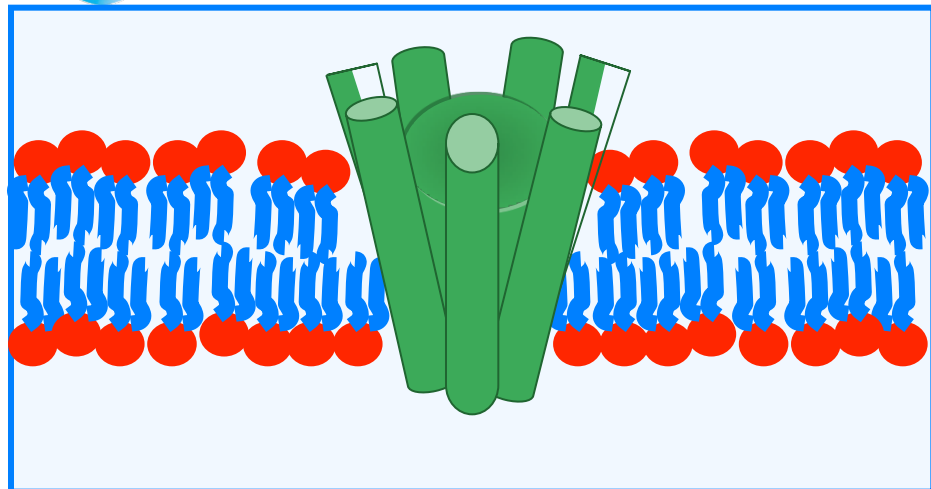
Source: L Arleth, Uni Copenhagen



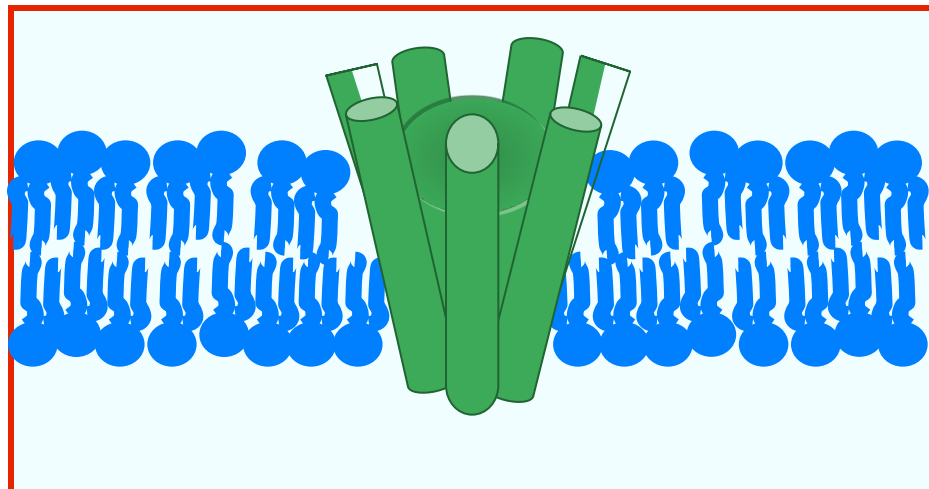
EUROPEAN
SPALLATION
SOURCE

SANS versus SAXS

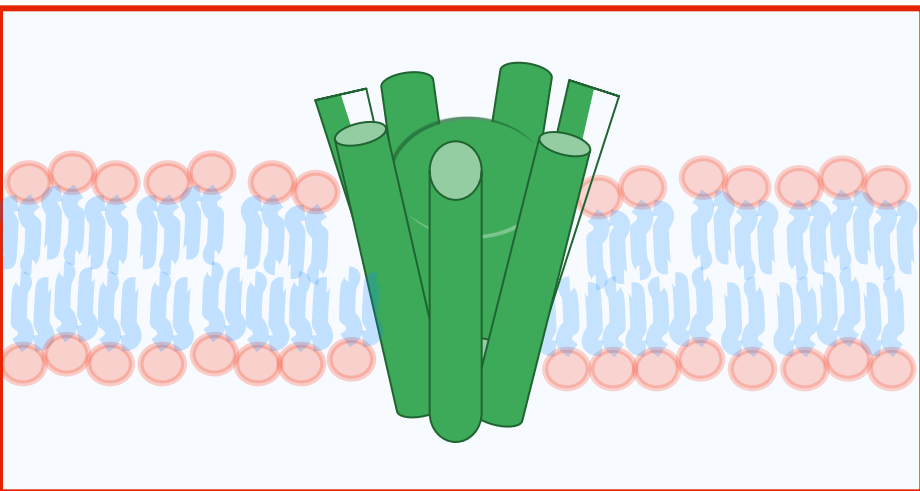
SAXS contrast



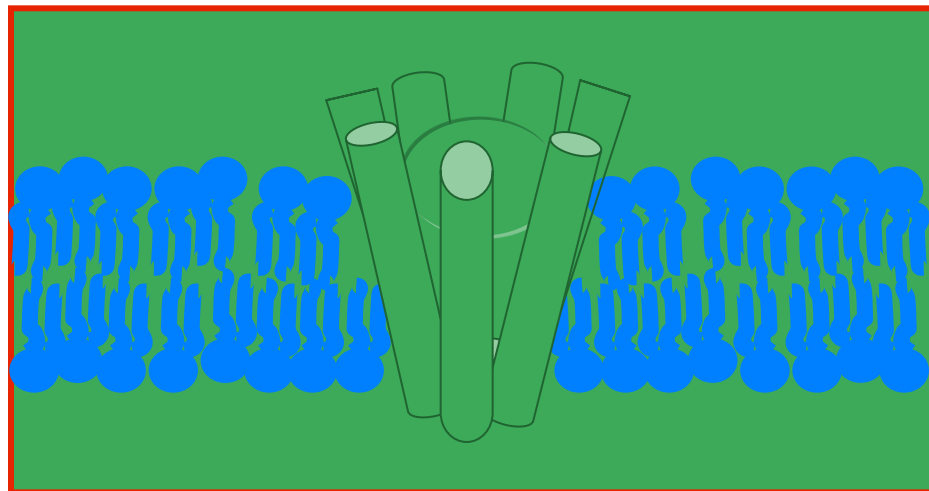
SANS contrast 1



SANS contrast 2

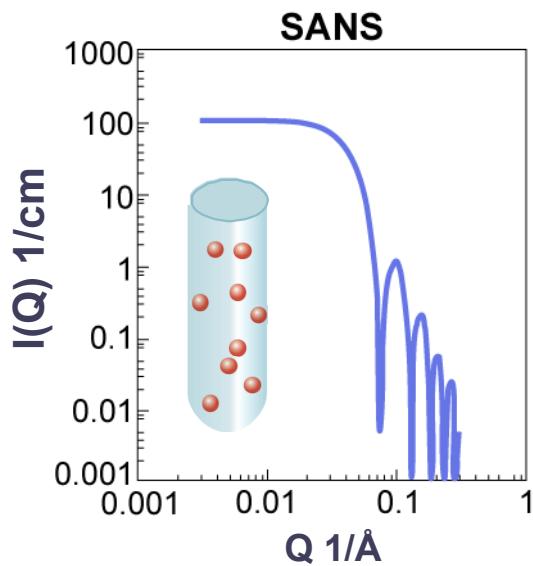


SANS contrast 3

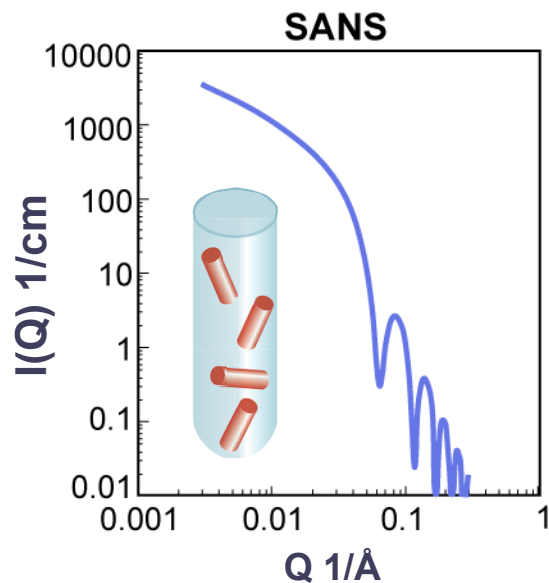


SANS gives the possibility of **not** seeing everything at the same time....

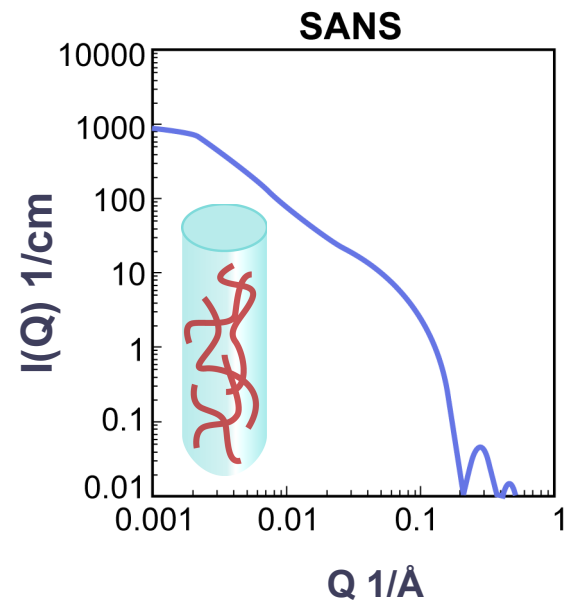
- **Spheres:**
 $R = 60 \text{ \AA}$



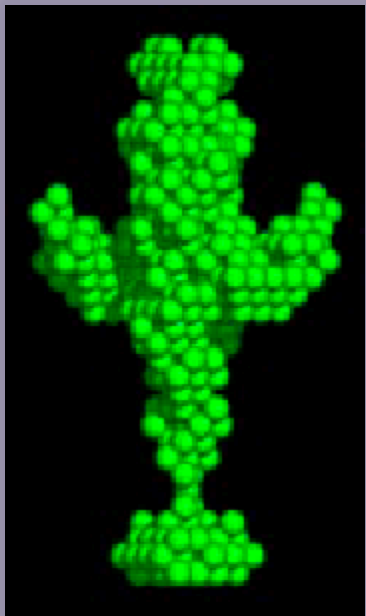
- **Rods: $R = 60 \text{ \AA}$**
 $L = 1200 \text{ \AA}$



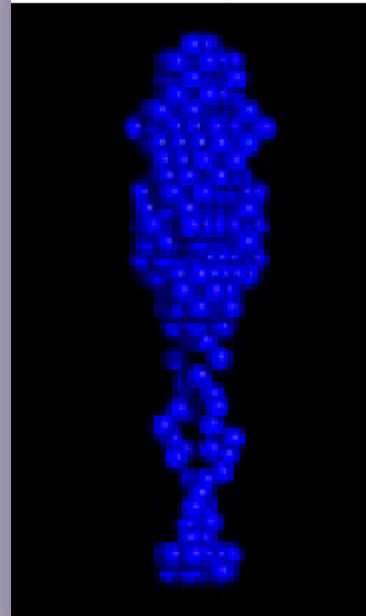
- **Worms: $R = 18 \text{ \AA}$**
 $L = 5000 \text{ \AA}$,
Kuhn Length = 300 \AA



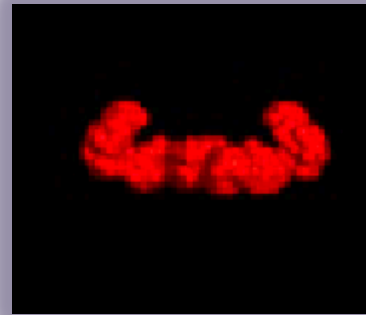
sensitivity and selectivity
isotopic substitution/contrast variation



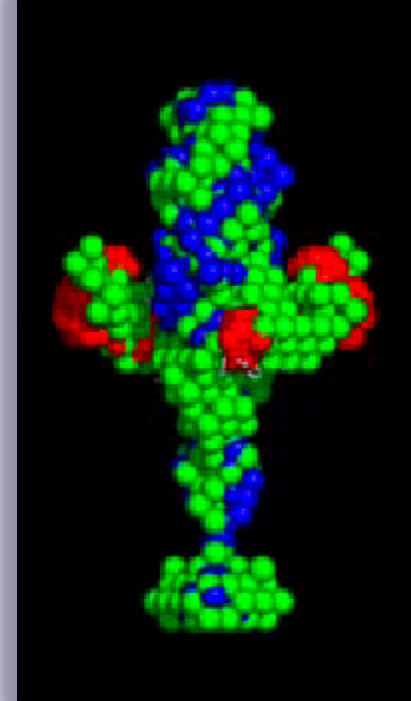
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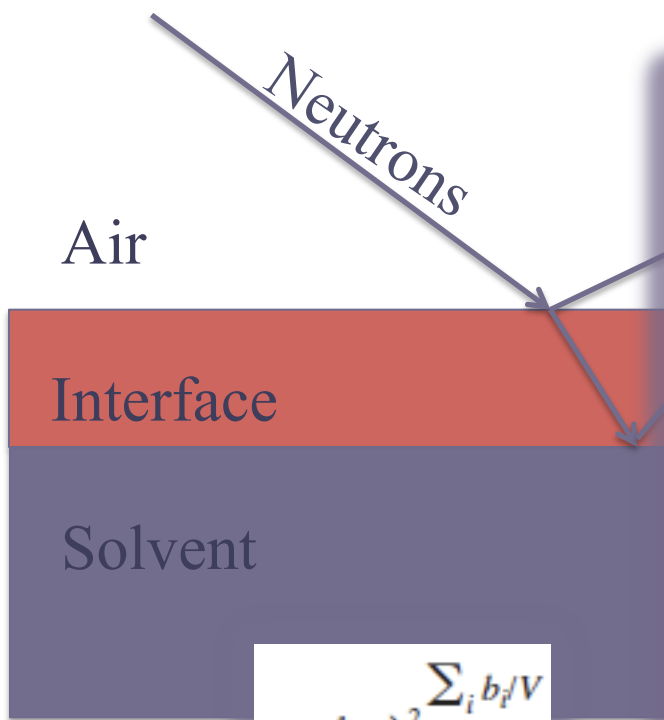
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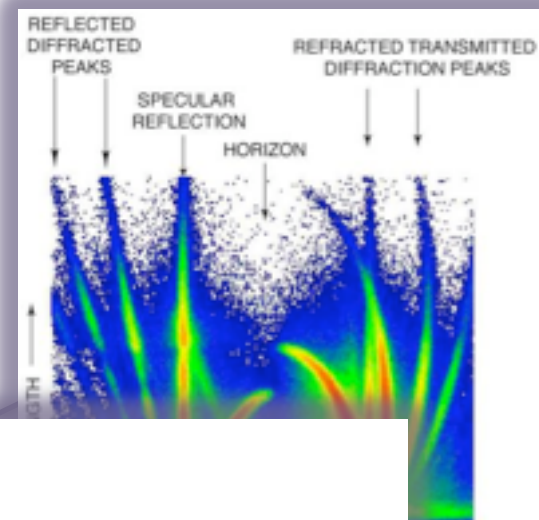
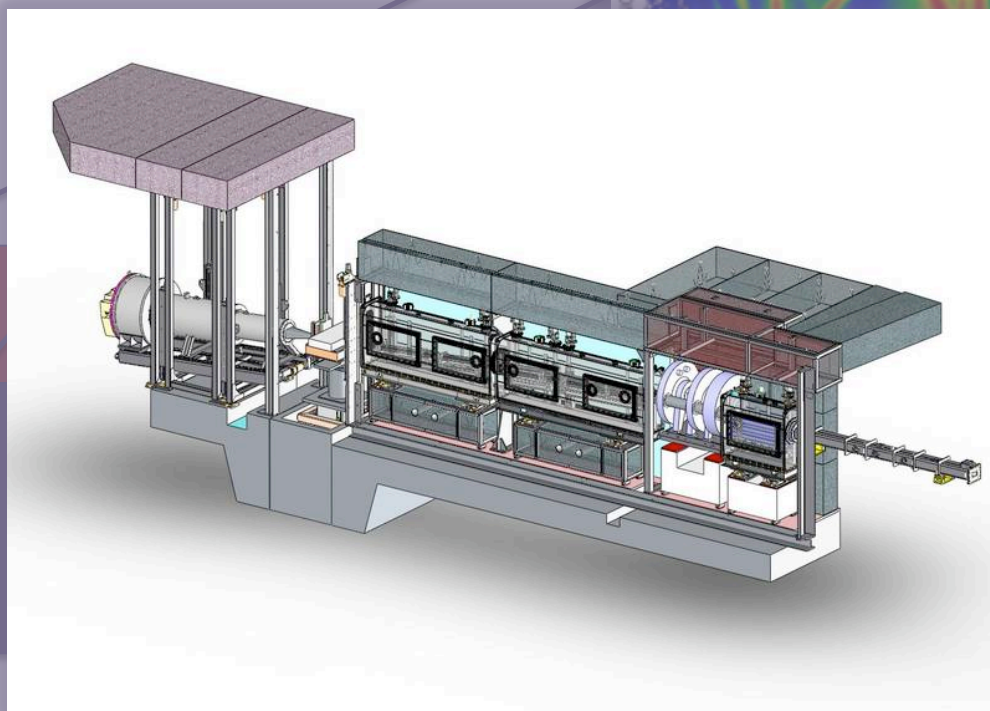
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- Basic principle



$$n = 1 - \lambda^2 \frac{\sum_i b_i / V}{2\pi}$$



10 μm 63 μÅ⁻¹

Neutron reflectometry to investigate the delivery of lipids and DNA to interfaces (Review)

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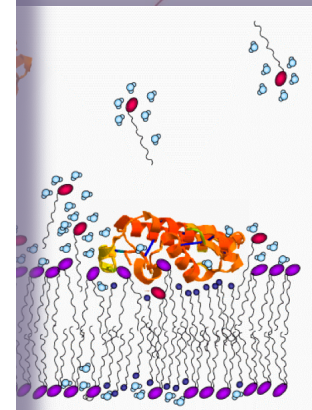
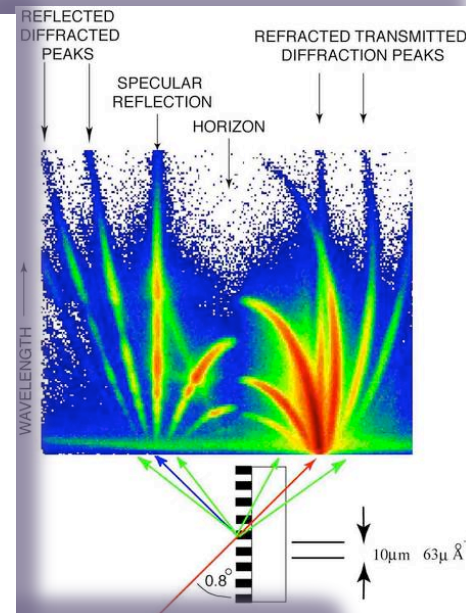
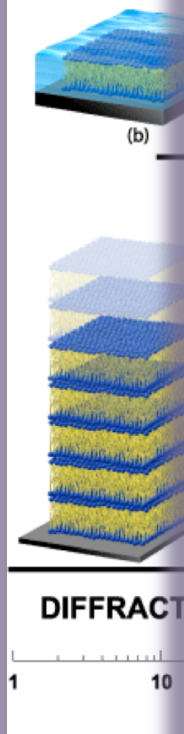
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(Received 6 May 2008; accepted 1 August 2008; published 19 December 2008)

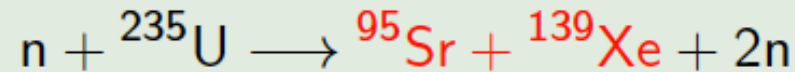
The application of scattering methods in the study of biological and biomedical problems is a field of research that is currently experiencing fast growth. In particular, neutron reflectometry (NR) is a technique that is becoming progressively more widespread, as indicated by the current commissioning of several new reflectometers worldwide. NR is valuable for the characterization of biomolecules at interfaces due to its capability to provide quantitative structural and compositional information on relevant molecular length scales. Recent years have seen an increasing number of applications of NR to problems related to drug and gene delivery. We start our review by summarizing the experimental methodology of the technique with reference to the description of biological liquid interfaces. Various methods for the interpretation of data are then discussed, including a new approach based on the lattice mean-field theory to help characterize stimulus-responsive surfaces relevant to drug delivery function. Recent progress in the subject area is reviewed in terms of NR studies relevant to the delivery of lipids and DNA to surfaces. Lastly, we discuss two case studies to exemplify practical features of NR that are exploited in combination with complementary techniques. The first case concerns the interactions of lipid-based cubic phase nanoparticles with model membranes (a drug delivery application), and the second case concerns DNA compaction at surfaces and in the bulk solution (a gene delivery application). © 2008 American Vacuum Society. [DOI: 10.1116/1.2976448]





Fission and Spallation

Example



Why incident neutron?

- ▶ Zero electric charge (\Rightarrow no Coulomb's repulsion) allows the neutrons of very low energy to approach the nucleus at $l \ll L_{nucl}$
- ▶ Very high fission cross-section with ${}^{235}\text{U}$.

Why ${}^{235}\text{U}$?

- ▶ Most of heavy nuclei can undergo the fission reaction initiated by neutrons, but ${}^{235}\text{U}$ has very high cross section with thermal neutrons.
- ▶ No fission threshold energy \Rightarrow incident neutrons can be very slow.

Prompt neutrons

- ▶ Average energy: ~ 1 MeV

Fission fragments

- ▶ Beta radioactive

\rightarrow Chain Reaction



The spallation process

Spallation Neutron Yield (i.e. multiplicity of emitted neutrons)

determines the requirement in terms of the accelerator power
(current and energy of incident proton beam).

Spallation Neutron Spectrum (i.e. energy distribution of emitted neutrons)

determines the damage and activation of the structural materials
(design of the beam window and spallation target)

Spallation Product Distributions

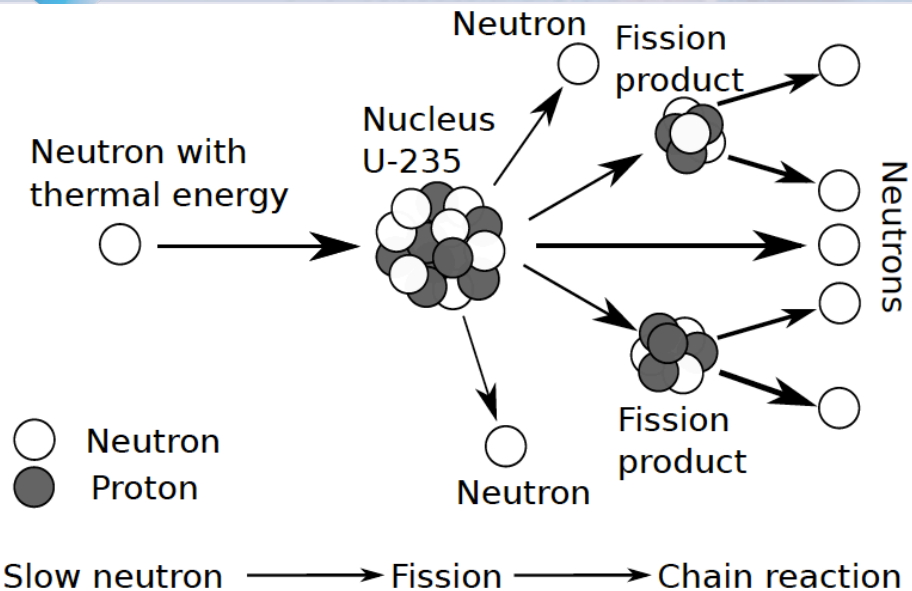
determines the radiotoxicity of the residues (radioprotection requirements).

Energy Deposition

determines the thermal-hydraulic requirements (cooling capabilities and nature of the spallation target).

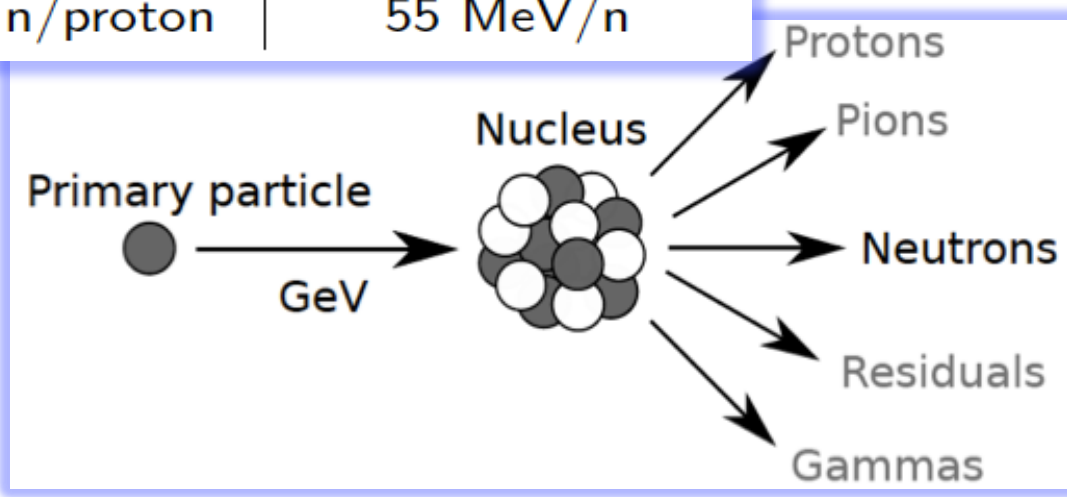
→ Sub-critical Reaction

Fission and Spallation

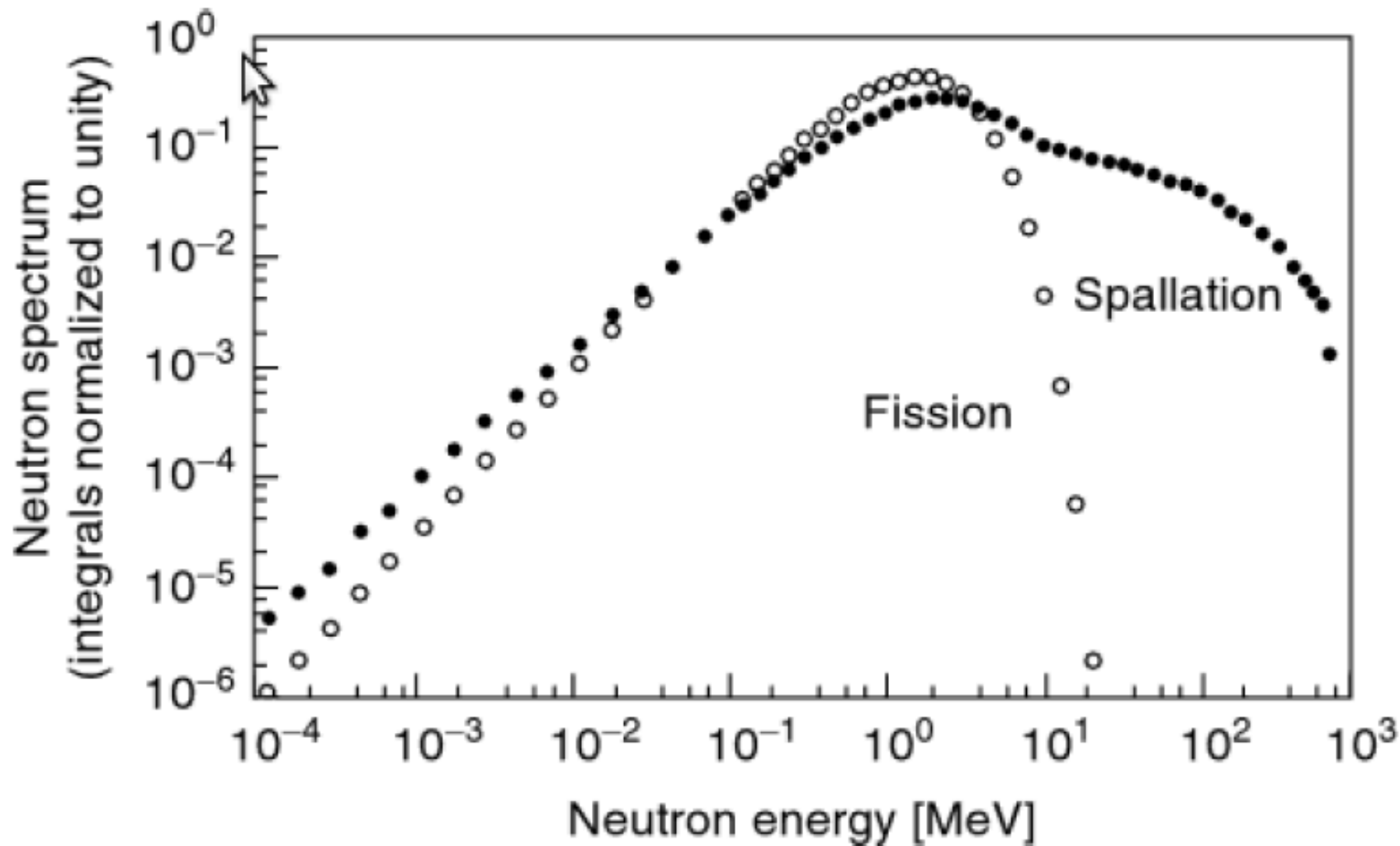


Spallation is a non-elastic nuclear interaction induced by a high-energy particle producing numerous secondary particles

Process	Reaction	Neutron yield	Energy deposition
Fission	$^{235}\text{U}(n,f)$	3 n/fission	190 MeV/n
Spallation	$p\ 1\ \text{GeV} \rightarrow \text{Hg}$	30 n/proton	55 MeV/n



Fission and Spallation



Energy efficiency is key for high intensity neutron beam production

Fast neutrons produced / joule **heat deposited** in target station

Fission reactors:	$\sim 10^9$	(in ~ 50 liter volume)
Spallation:	$\sim 10^{10}$	(in ~ 2 liter volume)
Fusion:	$\sim 1.5 \times 10^{10}$	(in ~ 2 liter volume) (but neutron slowing down efficiency reduced by ~ 20 times)
Photo neutrons:	$\sim 10^9$	(in ~ 0.01 liter volume)
Nuclear reaction (p, Be):	$\sim 10^8$	(in ~ 0.001 liter volume)
Laser induced fusion:	$\sim 10^4$	(in $\sim 10^{-9}$ liter volume)

Spallation: most favourable for the foreseeable future