

Diffraction PDFs



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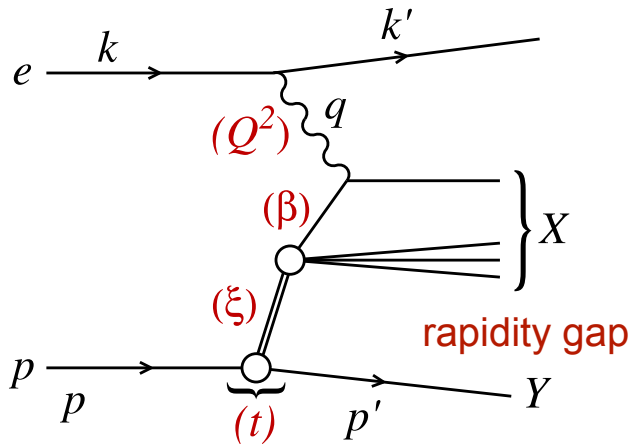
Based on Frankfurt, Guzey, Stasto, Strikman, “Selected topics in diffraction with protons and nuclei: past, present, and future”, Rep. Prog. Phys. 85 (2022) 126301 [2203.12289 [hep-ph]]

Outline:

- Collinear factorization and diffractive PDFs
- Diffractive PDFs from global fits
- Diffractive dijet photoproduction and factorization breaking
- Nuclear diffractive PDFs in eA diffractive DIS
- Outlook

Diffraction in ep DIS at HERA

- Diffractive scattering at high energies \rightarrow target intact and rapidity gap, i.e., large region in a detector with no activity.
- Present in both soft (elastic pp scattering) and hard (ep DIS) processes.
- Challenging in QCD due to enhanced HT/non-linear effects.
- Classic and most studied example: diffraction in ep DIS at HERA.



Standard DIS variables:

$$q^2 = -Q^2, \quad x = \frac{Q^2}{2p \cdot q}, \quad W^2 = (p + q)^2, \quad y = \frac{p \cdot q}{p \cdot k},$$

Additional diffraction-specific variables:

$$t = (p - p')^2, \quad \xi = \frac{Q^2 + M_X^2 - t}{Q^2 + W^2}, \quad \beta = \frac{Q^2}{Q^2 + M_X^2 - t}$$

- Reduced DIS cross section in terms of diffractive structure functions:

$$\sigma_{\text{red}}^{\text{D(4)}} = F_2^{\text{D(4)}}(\beta, \xi, Q^2, t) - \frac{y^2}{Y_+} F_L^{\text{D(4)}}(\beta, \xi, Q^2, t)$$

$$Y_+ = 1 + (1 - y)^2$$

Collinear factorization in diffractive DIS

- Similarly to inclusive DIS → collinear factorization for diffractive DIS, [Collins, PRD 57 \(1998\) 3051, PRD 61 \(2000\) 019902 \(erratum\)](#).

- Diffractive cross section given by convolution of coefficient functions (same as in inclusive case) with diffractive parton distributions (PDFs):

$$F_{2/L}^{D(4)}(\beta, \xi, Q^2, t) = \sum_i \int_{\beta}^1 \frac{dz}{z} C_{2/L,i}(\beta/z, Q^2) f_i^D(z, \xi, Q^2, t)$$

- Similarly to inclusive case, operator definition for diffractive PDFs:

$$f_i^D(z, \xi, Q^2, t) = \frac{1}{4\pi} \frac{1}{2} \sum_s \int dy^- e^{-izp^+ y^-} \sum_{X, s'} \langle p, s | \bar{\psi}(0, y^-, \mathbf{0}_T) | p', s'; X \rangle \gamma^+ \langle p', s'; X | \psi(0) | p, s \rangle$$

- Diffractive PDFs = conditional probabilities of finding partons in the proton, provided that it scatters into the final system Y with momentum p' .
- Similarly to inclusive case, diffractive PDF are universal (probed in inclusive diffraction, diffractive jet production, etc.) and obey Dokshitzer-Gribov-Lipatov-Altarelli-Parisi (DGLAP) evolution equations at fixed ξ and t .

Diffractive PDFs from global fits

- Diffractive PDFs given by long-distance matrix elements → non-perturbative and needs to be extracted from data using global QCD fits.
- Depend on 4 kinematic variables, c.f. usual PDFs → need simplifications.
- Proton vertex (Regge) factorization, [Ingelman, Schlein, PLB 152 \(1985\) 256](#), in terms of leading Pomeron and sub-leading ($\xi \geq 0.03$) Reggeon terms:

$$f_i^D(z, \xi, Q^2, t) = f_{\mathbb{P}}(\xi, t) f_i^{\mathbb{P}}(z, Q^2) + f_{\mathbb{R}}(\xi, t) f_i^{\mathbb{R}}(z, Q^2)$$

- The fluxes in the form motivated by Regge theory:

$$f_{\mathbb{P},\mathbb{R}}^P(\xi, t) = A_{\mathbb{P},\mathbb{R}} \frac{e^{B_{\mathbb{P},\mathbb{R}} t}}{\xi^{2\alpha_{\mathbb{P},\mathbb{R}}(t)-1}}, \quad \alpha_{\mathbb{P},\mathbb{R}}(t) = \alpha_{\mathbb{P},\mathbb{R}}(0) + \alpha'_{\mathbb{P},\mathbb{R}} t$$

- Simple form for sea quark (valence =0) and gluon PDFs of the Pomeron:

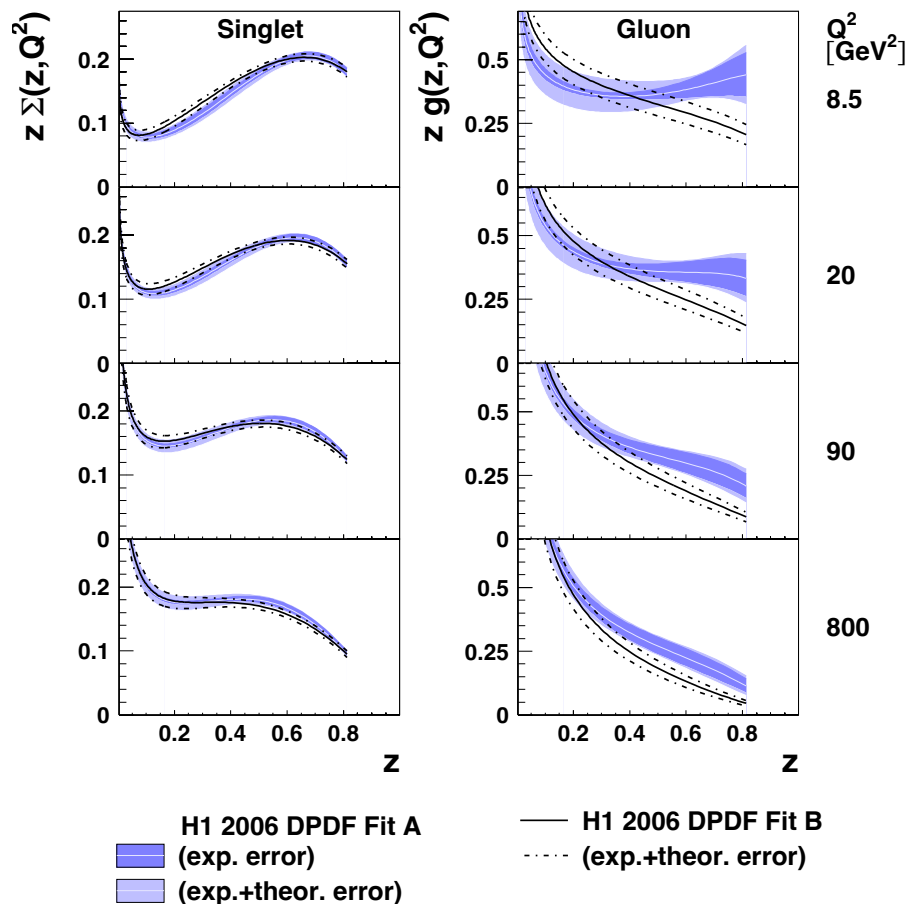
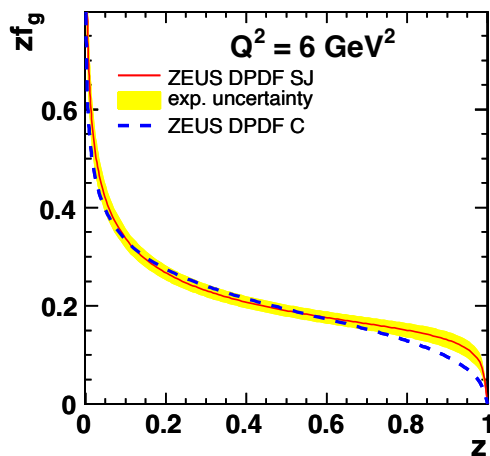
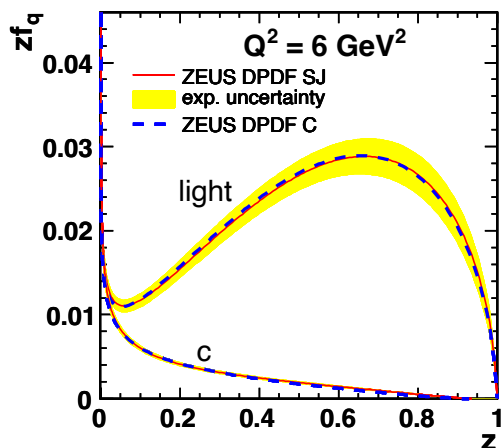
$$z f_i^{\mathbb{P}}(z, \mu_0^2) = A_i z^{B_i} (1-z)^{C_i}$$

- $f_u^{\mathbb{P}} = f_d^{\mathbb{P}} = f_s^{\mathbb{P}}$, and massless heavy flavors in the variable flavor scheme generated through DGLAP evolution.

- Reggeon $f_i^{\mathbb{R}}$ taken from pion PDFs → can be better constrained at EIC, [Armesto, Newman, Slominski, Stasto, PRD 110 \(2024\) 5, 054039](#).

Diffractive PDFs from global fits (2)

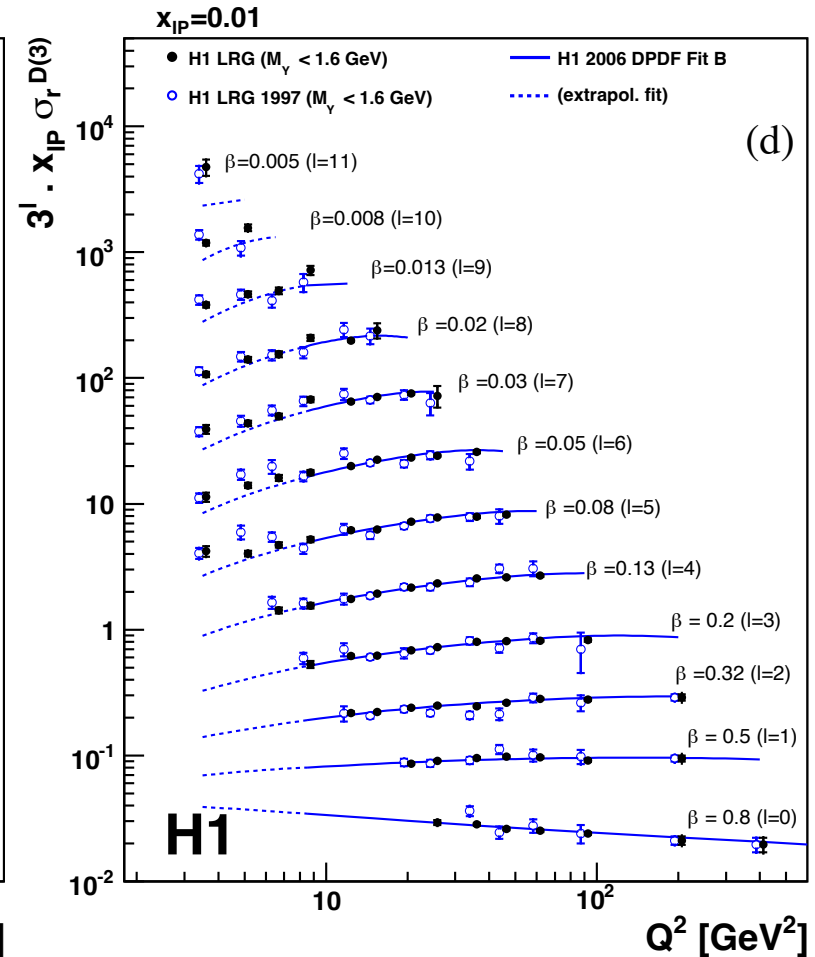
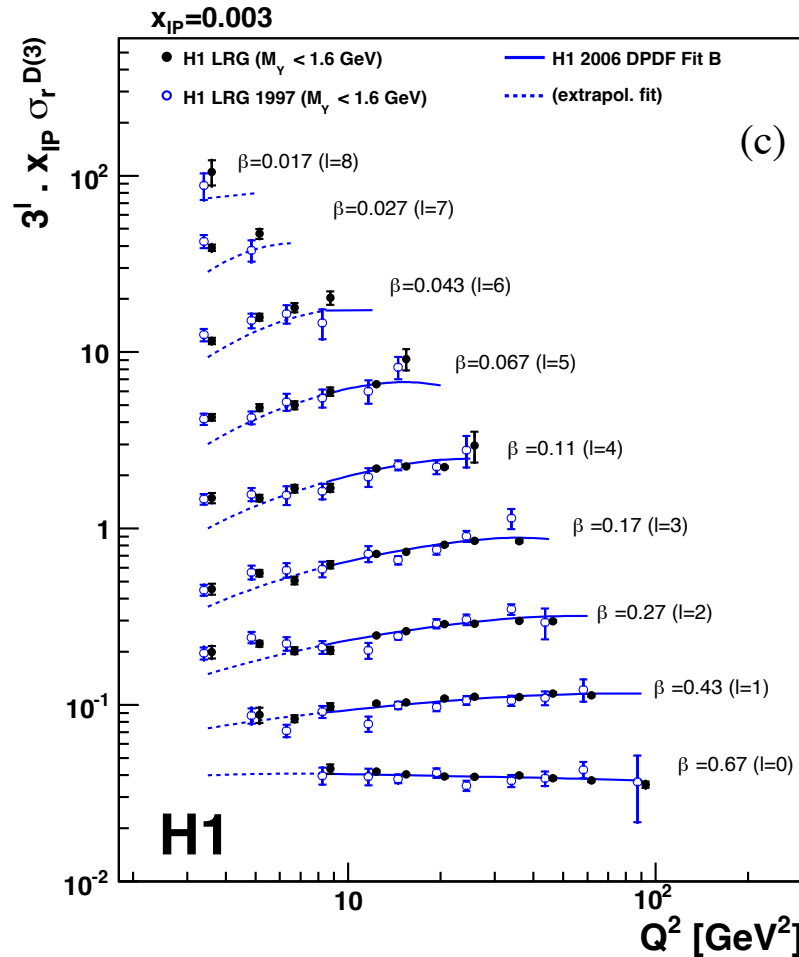
- Most notable examples are ZEUS and H1 analyses of their own data, [Chekanov et al, NPB 831 \(2010\) 1](#); [Aktas et al, EPJC 48 \(2006\) 715](#).



- Diffractive gluon PDF \gg quark PDFs, but poorly constrained at large $z \rightarrow$ need to include data on diffractive dijets \rightarrow favor ZEUS C and H1 B fits.

Diffraction PDFs from global fits (3)

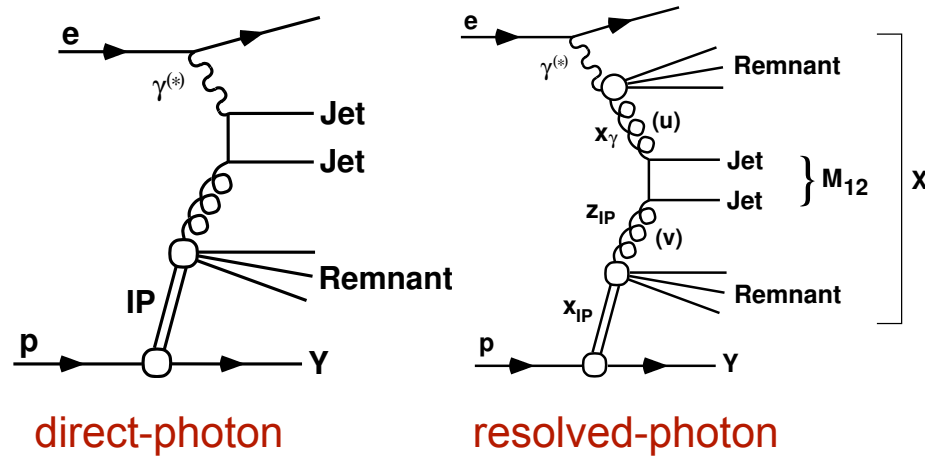
- Good description of original and more recent H1 data, [Aaron et al, EPJC 72 \(2012\) 2074](#)



- Comparison of LRG with proton-tagged cross section measurement \rightarrow $\sim 20\%$ contribution of proton dissociation.

Diffractive dijet photoproduction

- Collinear fact: same diffractive PDFs for pQCD description of various processes.
- Diffractive dijet electro- and photoproduction in ep scattering → constraints on gluon distribution.



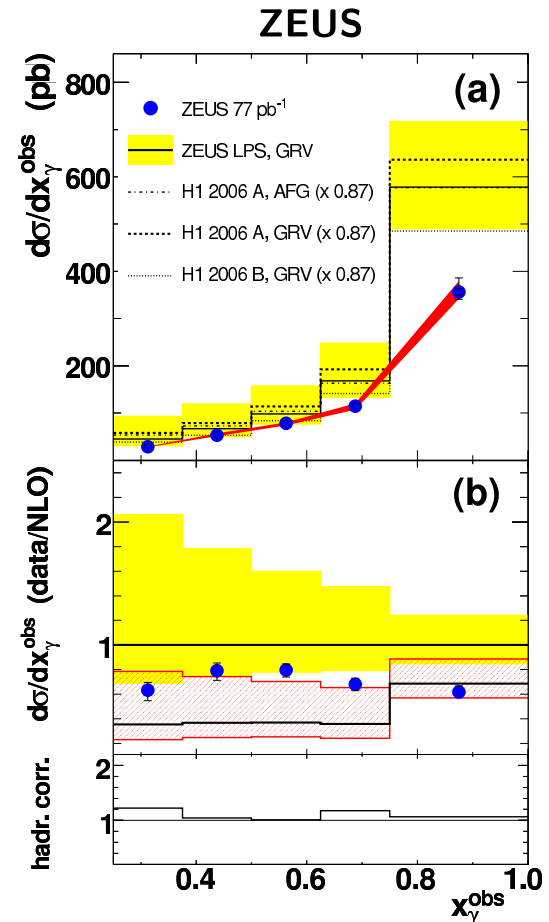
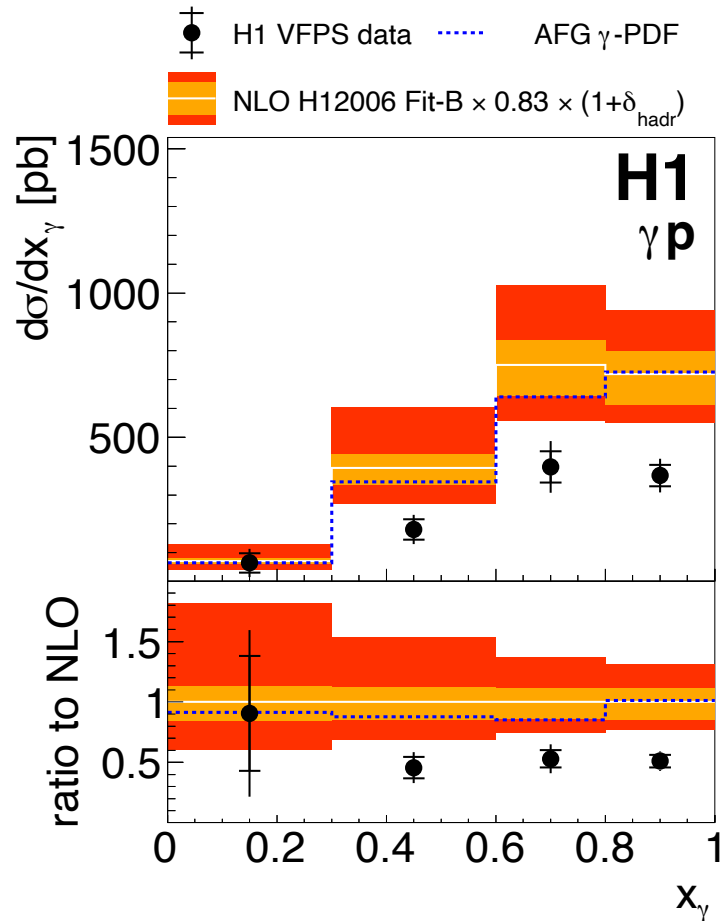
- Cross section is known to NLO accuracy [Klasen, Kramer, Salesch, Z. Phys. C 68, 113 \(1995\)](#); [Klasen, Kramer, Z. Phys. C 72, 107 \(1996\)](#), [Z. Phys. C 76, 67 \(1997\)](#); [Klasen, Rev. Mod. Phys. 74, 1221 \(2002\)](#)

$$d\sigma = \sum_{a,b} \int dy \int dx_\gamma \int dt \int dx_{\mathbb{P}} \int dz_{\mathbb{P}} f_{\gamma/e}(y) f_{a/\gamma}(x_\gamma, M_\gamma^2) f_{\mathbb{P}/p}(x_{\mathbb{P}}, t) f_{b/\mathbb{P}}(z_{\mathbb{P}}, M_{\mathbb{P}}^2) d\hat{\sigma}_{ab}^{(n)}$$

Photon flux in Weizsäcker-Williams approximation
 Photon PDFs for resolved photon; in $a=\gamma$ case, $f_{a/\gamma}=\delta(1-x_\gamma)$
 "Pomeron" flux
 Diffractive PDFs
 Hard partonic cross section

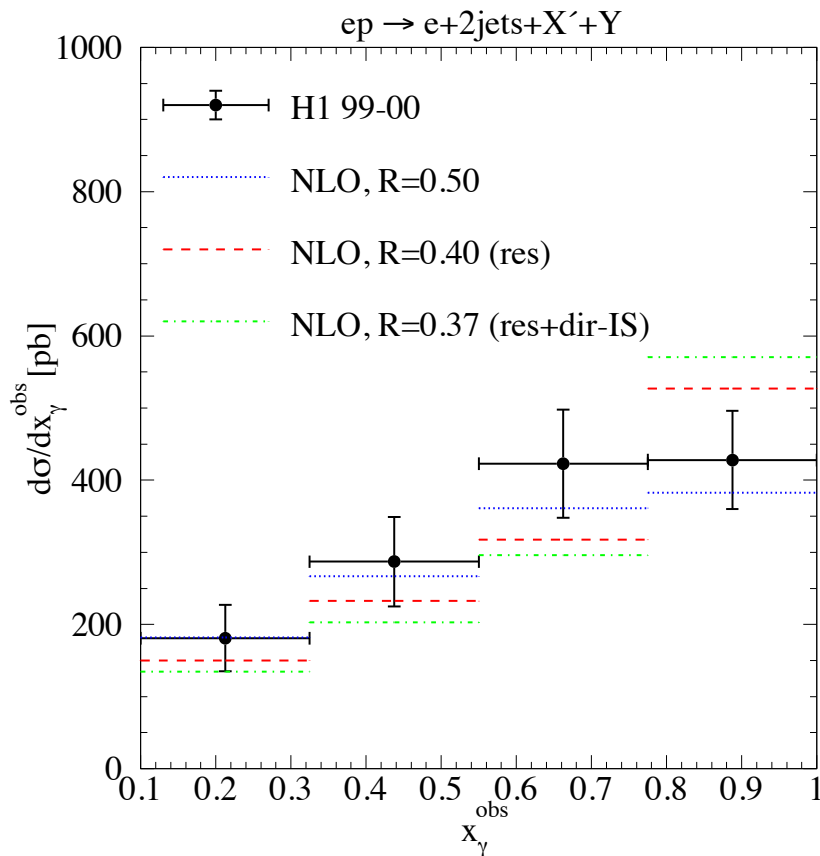
Diffractive dijet photoproduction (2)

- Universality of diffractive PDFs successfully tested in diffractive dijet and open charm product in DIS, Aktas et al. [H1 Coll.], JHEP 10, 042 (2007); EPJ C 71, 549 (2010); EPJ C 50, 1 (2007); Chekanov et al. [ZEUS Coll.], EPJ C 52, 813 (2007); Chekanov et al. [ZEUS Coll.], NPB 831, 1 (2010)
- At the same time, NLO pQCD QCD overestimates cross sections of diffractive dijet photoproduction at HERA by **factor 2** → factorization breaking, Aktas et al. [H1 Coll.], EPJ C 71, 549 (2007); Aaron et al. [H1 Coll.], EPJ C 70, 15 (2010); Andreev et al. [H1 Coll.], JHEP 05, 056 (2015); Chekanov et al. [ZEUS Coll.], EPJ C 55, 177 (2008).

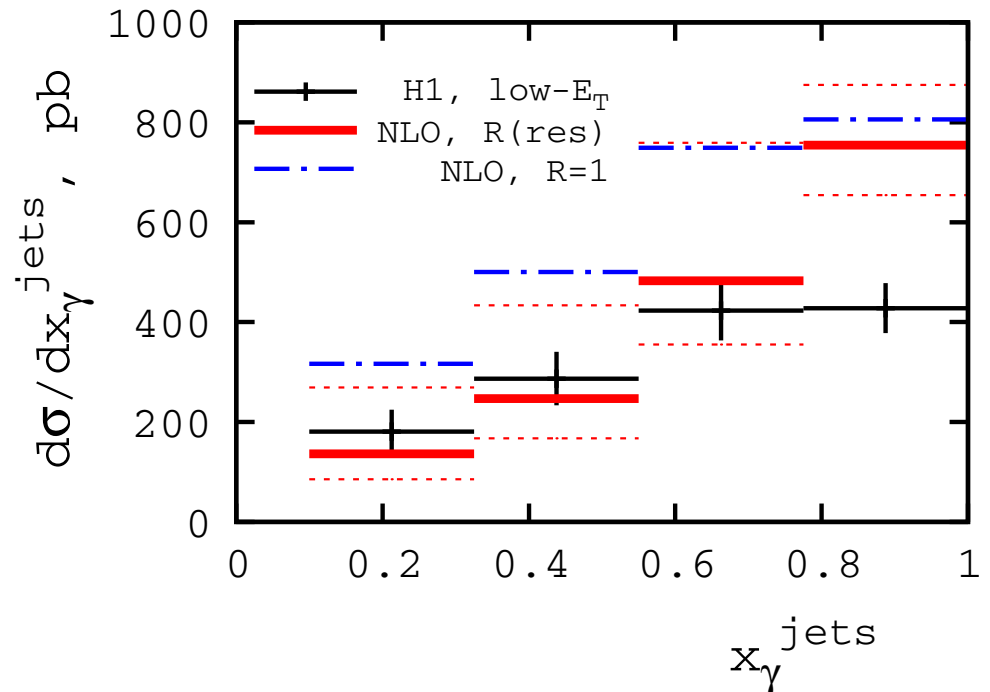


Factorization breaking

- Mechanism of this factorization breaking remains unknown:
 - global suppression factor $R \approx 0.5$
 - suppression of only resolved photon contribution by $R \approx 0.34$ as expected in hadron-hadron scattering, [Kaidalov, Khoze, Martin, Ryskin, PLB 567, 61 \(2003\)](#); [Klasen, Kramer, EPJ C 70, 91 \(2010\)](#)
 - flavor-dependent combination of these mechanisms, [Guzey, Klasen, EPJ C 76, 467 \(2016\)](#)



$$S_i^2(x_\gamma) \rightarrow \begin{cases} 1, & i = c, \\ A_q x_\gamma + 0.34, & i = u, d, s, \\ A_g x_\gamma + 0.34, & i = g, \end{cases}$$

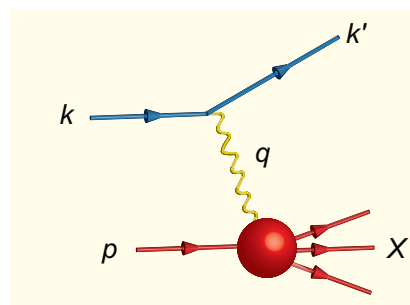


Diffraction in DIS on nuclei

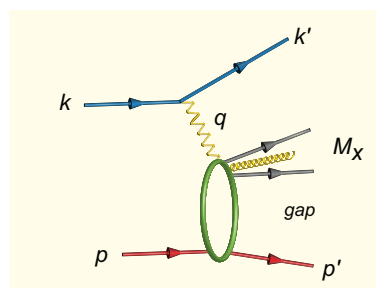
- The planned Electron-Ion Collider (EIC) in USA has potential to discriminate among approaches of NS due to:

- wide $x - Q^2$ coverage
- measurement of longitudinal structure function $F_L^A(x, Q^2)$ sensitive to gluons
- for the first time measurement of hard **diffraction** in nuclear DIS.

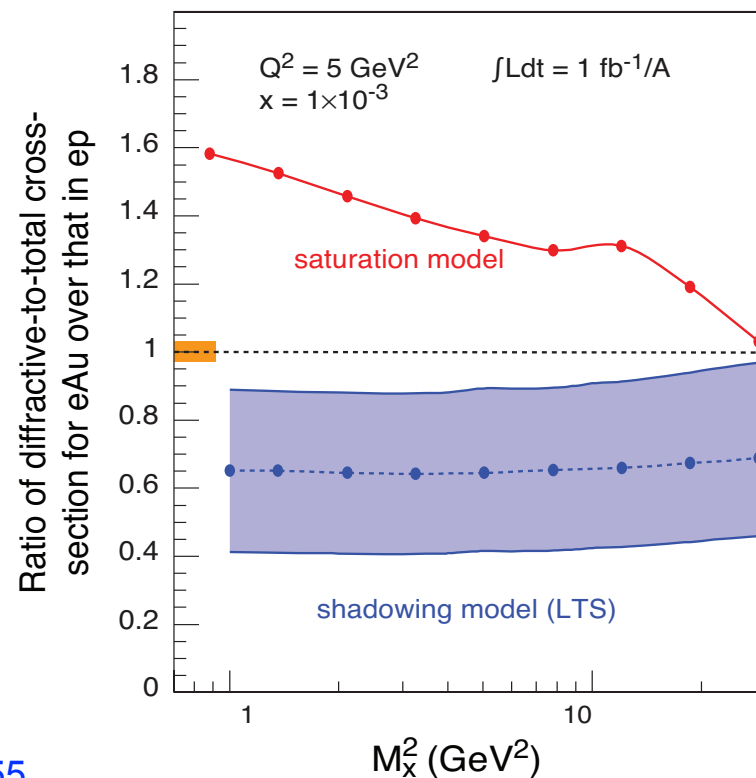
- Sensitive observable is the ratio of diffractive to total DIS cross sections for a heavy nucleus and the proton, [Accardi et al., EPJ A52 \(2016\) 9, 268 \[1212.1701 \[hep-ex\]\]](#):



Total DIS



Diffractive DIS



- Predicted to be:

- $R_{\text{diff/tot}} > 1$ due to nuclear enhancement of saturation scale $Q_{s,A}^2$, [Kowalski, Lappi, Venugopalan, PRL 100 \(2008\) 022303](#); [Lappi, Le, Mäntysaari, PRD 108 \(2023\) 114023](#)
- $R_{\text{diff/tot}} < 1$ due to strong leading twist nuclear shadowing, [Frankfurt, Guzey, Strikman, Phys. Rept. 512 \(2012\) 255](#)

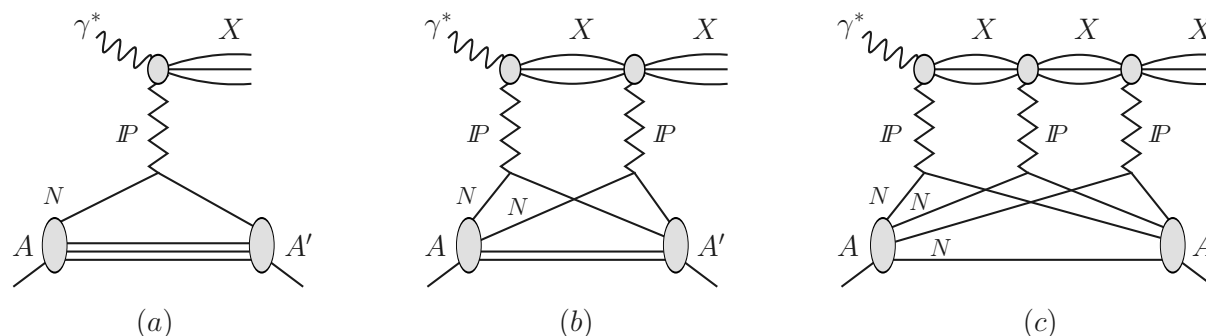
Leading twist approach to nuclear shadowing

- Method to calculate various nuclear parton distributions (usual, generalized, diffractive) as input for DGLAP evolution, Frankfurt, Strikman, EPJ A5 (1999) 293; Frankfurt, Guzey, Strikman, Phys. Rept. 512 (2012) 255 → alternative to global fits of nPDFs.

• Based on:

- Gribov-Glauber model of NS for soft hadron-nucleus scattering
- QCD factorization theorems for inclusive and diffractive DIS.

• $\gamma^* + A \rightarrow X + A'$ amplitude is a series of diffractive scattering off $i = 1, 2, \dots, A$ target nucleons:



• Coherent diffraction $A' = A$:

$$\sigma_{\gamma^*A \rightarrow XA} = \int d^2\vec{b} |\Gamma_{\gamma^*A \rightarrow XA}(\vec{b})|^2 = 4\pi \frac{d\sigma_{\gamma^*N \rightarrow XN}(t=0)}{dt} \int d^2\vec{b} \left| \int dz \rho_A(\vec{b}, z) e^{iz\Delta_{\gamma^*X}} e^{-\frac{1-in}{2}\sigma_{\text{soft}} \int_z^\infty dz' \rho_A(\vec{b}, z')} \right|^2$$

diffractive cross section measured at HERA

nuclear density

model-dependent cross section

LTA to nuclear shadowing (2)

- Apply collinear QCD factorization for diffractive DIS, [Collins, PRD 57 \(1998\); PRD 61 \(2000\) 019902](#) → from structure function to parton distributions:

$$f_{i/A}^{D(3)}(x, x_{\mathbb{P}}, Q^2) = 4\pi f_{i/p}^{D(4)}(x, x_{\mathbb{P}}, Q^2, t = 0) \int d^2\vec{b} \left| \int dz \rho_A(\vec{b}, z) e^{izx_{\mathbb{P}}m_N} e^{-\frac{1-i\eta}{2}\sigma_{\text{soft}}^i(x) \int_z^\infty dz' \rho_A(\vec{b}, z')} \right|^2$$

$$= f_{i/p}^{D(3)}(x, x_{\mathbb{P}}, Q^2) \frac{1}{\sigma_{\text{el}}^i(x)} \int d^2\vec{b} \left| 1 - e^{-\frac{1-i\eta}{2}\sigma_{\text{soft}}^i(x) T_A(\vec{b})} \right|^2$$

$$\sigma_{\text{el}}(x) = \frac{[\sigma_{\text{soft}}(x)]^2}{16\pi B_{\text{diff}}}$$

$$T_A(\vec{b}) = \int dz \rho(\vec{b}, z)$$

- Transparent interpretation: nuclear diffractive PDFs shadowed in proportion to the **nuclear elastic cross section**.

- Similarly for **quasi-elastic scattering** using completeness final states A' :

$$\sigma_{\gamma^*A \rightarrow XA'} = \int d^2\vec{b} \langle A' | \left| \Gamma_{\gamma^*A \rightarrow XA}(\vec{b}) \right|^2 | A \rangle = \sigma_{\gamma^*N \rightarrow XN} \frac{1}{\sigma_{\text{el}}} \int d^2\vec{b} \left(\left| 1 - e^{-\frac{1-i\eta}{2}\sigma_{\text{soft}} T_A(\vec{b})} \right|^2 + e^{-\sigma_{\text{in}} T_A(\vec{b})} - e^{-\sigma_{\text{soft}} T_A(\vec{b})} \right)$$

$$\tilde{f}_{i/A}^{D(3)}(x, x_{\mathbb{P}}, Q^2) = f_{i/p}^{D(3)}(x, x_{\mathbb{P}}, Q^2) \frac{1}{\sigma_{\text{el}}^i(x)} \int d^2\vec{b} \left(\left| 1 - e^{-\frac{1-i\eta}{2}\sigma_{\text{soft}}^i(x) T_A(\vec{b})} \right|^2 + e^{-\sigma_{\text{in}}^i(x) T_A(\vec{b})} - e^{-\sigma_{\text{soft}}^i(x) T_A(\vec{b})} \right)$$

$$\sigma_{\text{in}}(x) = \sigma_{\text{soft}}(x) - \sigma_{\text{el}}(x)$$

- In this case, NS is given by sum of **elastic** and **inelastic** nuclear cross sections.

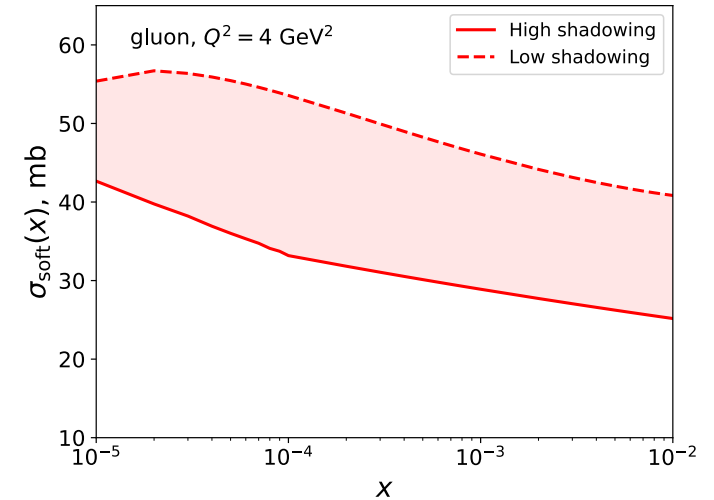
LTA predictions for nuclear diffractive PDFs

- Assumed that diffractive intermediate states X do not mix \rightarrow one free parameter $\sigma_{\text{soft}}^i(x)$ \rightarrow controls size and uncertainties of LTA predictions.

- High shadowing:** given by probability of diffraction

$$\sigma_{\text{soft}}^i(x) \approx \sigma_2(x) \equiv \frac{16\pi}{f_{ilp}(x)} \int_x^{0.1} \frac{dx_{IP}}{x_{IP}} f_{ilp}^{D(4)}(x, x_{IP}, t=0)$$

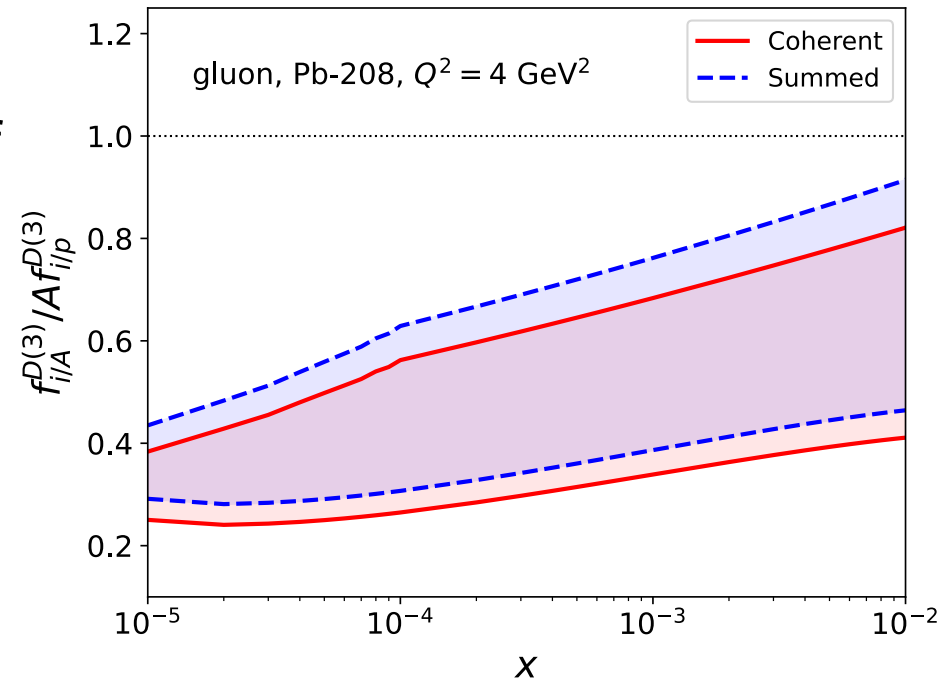
- Low shadowing:** calculated using model for hadronic structure of ρ meson.



- In LTA, nuclear shadowing driven by diffraction on proton \rightarrow 10-15% probability of diffraction in DIS@HERA leads to **large suppression** of nuclear PDFs at small x .

- Compare to impulse approximation (IA):

$$\frac{f_{ilA}^{D(3)}}{A f_{ilp}^{D(3)}} = \frac{4\pi B_{\text{diff}}}{A} \int d^2\vec{b} (T_A(\vec{b}))^2 = \frac{B_{\text{diff}}}{A} \int dt F_A^2(t) = 4.3$$

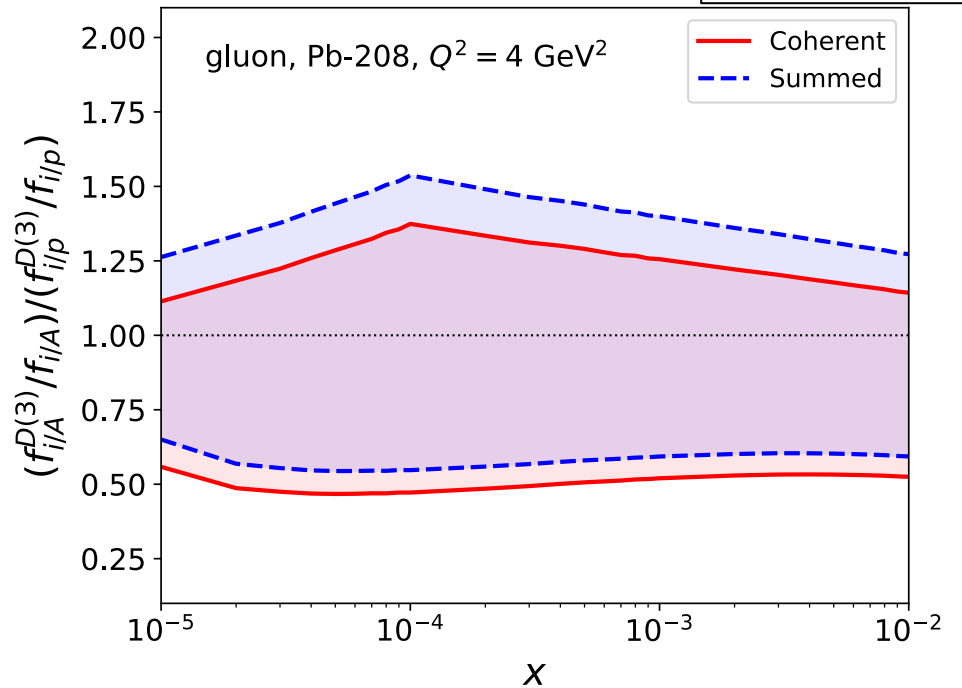
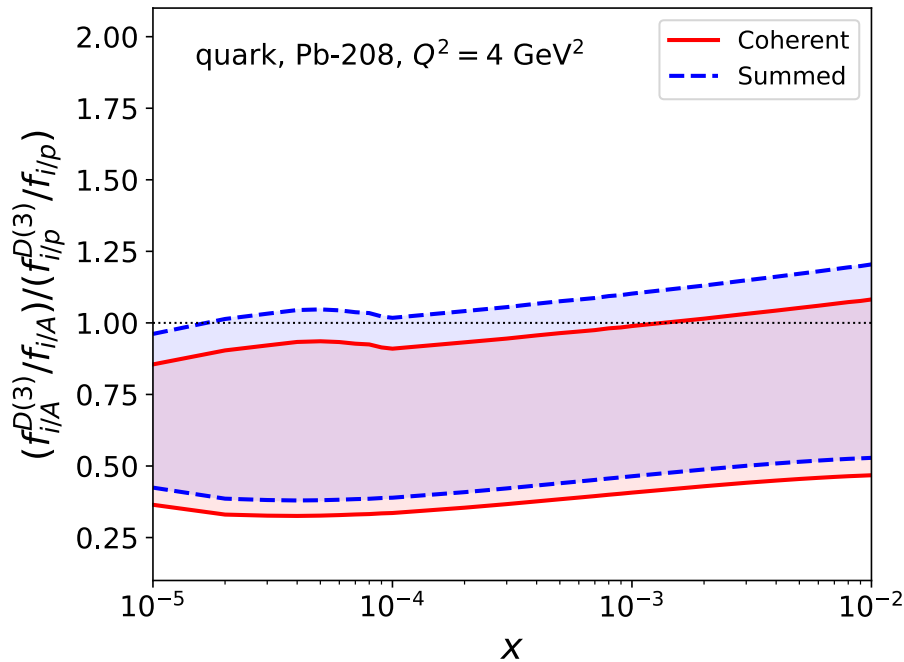


LTA predictions for $R_{\text{diff/tot}}$

- Combine LTA predictions for diffractive and usual nuclear PDFs:

$$\frac{f_{i/A}^{D(3)}(x, x_{\mathbb{P}}, Q^2)/f_{i/A}(x, Q^2)}{f_{i/p}^{D(3)}(x, x_{\mathbb{P}}, Q^2)/f_{i/p}(x, Q^2)} = \frac{\sigma_{\text{soft}}^i(x)}{\sigma_{\text{el}}^i(x)} \frac{\int d^2\vec{b} \left| 1 - e^{-\frac{1-i\eta}{2}\sigma_{\text{soft}}^i(x)T_A(\vec{b})} \right|^2}{2(1 - \lambda^i(x))\Re \int d^2\vec{b} \left(1 - e^{-\frac{1-i\eta}{2}\sigma_{\text{soft}}^i(x)T_A(\vec{b})} \right) + \lambda^i(x)A\sigma_{\text{soft}}^i(x)}$$

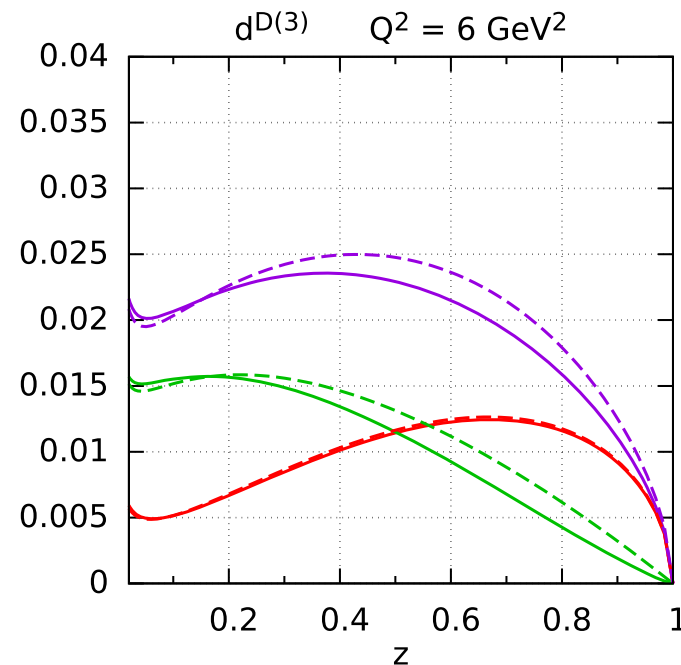
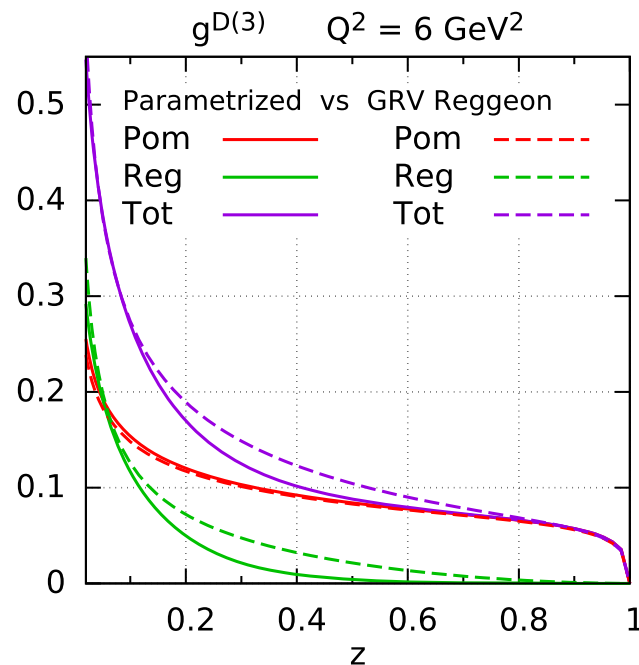
$$\lambda^i(x) = 1 - \sigma_2^i(x)/\sigma_{\text{soft}}^i(x)$$



- Suppression $R_{\text{diff/tot}} \approx 0.5 - 1$ (quarks) and $R_{\text{diff/tot}} \approx 0.5 - 1.3$ (gluons) due to interplay of large leading twist nuclear shadowing for diffractive and usual nuclear PDFs.

Outlook: diffraction at EIC

- Several recent global QCD fits for **proton** diffractive PDFs using all (inclusive+dijets) HERA data, [Salajeghen et al., PRD 107 \(2024\) 9, 093038](#); [PRD 106 \(2022\) 5, 054012](#)
- Further progress possible at the Electron-Ion Collider (EIC) $\rightarrow \sqrt{s_{NN}} \sim 100$ GeV lower than at HERA \rightarrow constrain subleading (Reggeon) contribution at large ξ , [Armesto, Newman, Slominski, Stasto, PRD 110 \(2024\) 5, 054039](#)

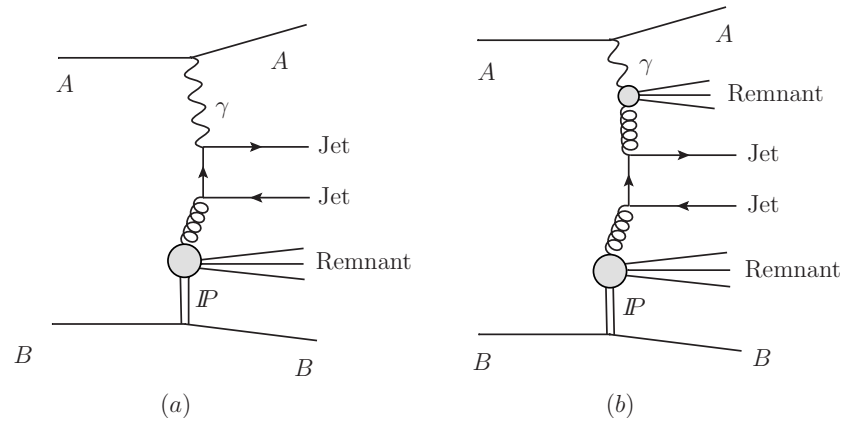


$$\xi = 0.1$$

- Similarly, NLO pQCD predicts 10-35% contribution of sub-leading Reggeon trajectory for $x_P > 0.06$ in diffractive dijet photoproduction, [Guzey, Klasen, JHEP 05 \(2020\) 074](#)

Outlook: diffraction in UPC at LHC

- Predictions for EIC can be tested in ultraperipheral collisions (UPCs) at LHC.



- Recent ATLAS measurement in $0nXn$ channel, [Aad et al, 2409.11060 \[nucl-ex\]](#) → can be extended to $0n0n$ channel probing nuclear diffractive PDFs.

- **Nucleus** can be used to suppress the resolved photon contribution → new handle on mechanism of factorization breaking, [Guzey, Klasen, JHEP 04 \(2016\) 158](#)

