

# Diffractive PDFs



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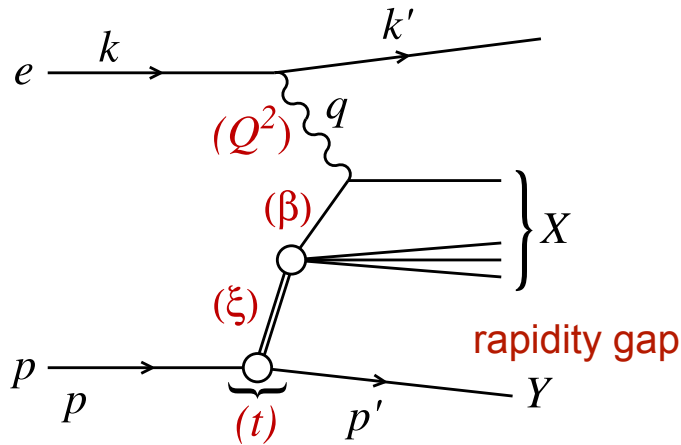
Based on Frankfurt, Guzey, Stasto, Strikman, “Selected topics in diffraction with protons and nuclei: past, present, and future”, Rep. Prog. Phys. 85 (2022) 126301 [2203.12289 [hep-ph]]

## Outline:

- Collinear factorization and diffractive PDFs
- Diffractive PDFs from global fits
- Diffractive dijet photoproduction at HERA and factorization breaking
- Nuclear diffractive PDFs in eA diffractive DIS
- Outlook and summary

# Diffraction in ep DIS at HERA

- Diffractive scattering at high energies → target intact and rapidity gap, i.e., large region in a detector with no activity.
- Present in both soft [elastic pp scattering] and hard [ep deep inelastic scattering (DIS)] processes.
- Challenging in QCD due to enhanced HT/non-linear effects.
- Classic and most studied example: diffraction in ep DIS at HERA → one of main HERA results that diffraction ~10-15% of total DIS cross section.



Standard DIS variables:

$$q^2 = -Q^2, \quad x = \frac{Q^2}{2p \cdot q}, \quad W^2 = (p + q)^2, \quad y = \frac{p \cdot q}{p \cdot k},$$

Additional diffraction-specific variables:

$$t = (p - p')^2, \quad \xi = \frac{Q^2 + M_X^2 - t}{Q^2 + W^2}, \quad \beta = \frac{Q^2}{Q^2 + M_X^2 - t}$$

- Reduced DIS cross section in terms of diffractive structure functions:

$$\sigma_{\text{red}}^{\text{D}(4)} = F_2^{\text{D}(4)}(\beta, \xi, Q^2, t) - \frac{y^2}{Y_+} F_L^{\text{D}(4)}(\beta, \xi, Q^2, t).$$

$$Y_+ = 1 + (1 - y)^2$$

# Collinear factorization in diffractive DIS

- Similarly to inclusive DIS → collinear factorization for diffractive DIS, [Collins, PRD 57 \(1998\) 3051, PRD 61 \(2000\) 019902 \(erratum\)](#).
- Diffractive cross section given by convolution of coefficient functions (same as in inclusive case) with diffractive parton distributions (PDFs):

$$F_{2/L}^{D(4)}(\beta, \xi, Q^2, t) = \sum_i \int_{\beta}^1 \frac{dz}{z} C_{2/L,i}(\beta/z, Q^2) f_i^D(z, \xi, Q^2, t)$$

- Similarly to inclusive case, operator definition for diffractive PDFs:

$$f_i^D(z, \xi, Q^2, t) = \frac{1}{4\pi} \frac{1}{2} \sum_s \int dy^- e^{-izp^+ y^-} \sum_{X, s'} \langle p, s | \bar{\psi}(0, y^-, \mathbf{0}_T) | p', s'; X \rangle \gamma^+ \langle p', s'; X | \psi(0) | p, s \rangle$$

- **Diffractive PDFs = conditional probabilities** of finding partons in the proton, provided that it scatters into the final system  $Y$  with momentum  $p'$ .
- Similarly to inclusive case, diffractive PDF are universal (probed in inclusive diffraction, diffractive jet production, etc.) and obey **Dokshitzer-Gribov-Lipatov-Altarelli-Parisi (DGLAP)** evolution equations at fixed  $\xi$  and  $t$ .

# Diffractive PDFs from global fits

- Diffractive PDFs given by long-distance matrix elements → non-perturbative and need to be extracted from data using global QCD fits.
- Depend on 4 kinematic variables, c.f. usual PDFs → need simplifications.
- Proton vertex (Regge) factorization, [Ingelman, Schlein, PLB 152 \(1985\) 256](#), for the leading Pomeron and sub-leading ( $\xi \geq 0.03$ ) Reggeon contributions:

$$f_i^D(z, \xi, Q^2, t) = f_{\mathbb{P}}(\xi, t) f_i^{\mathbb{P}}(z, Q^2) + f_{\mathbb{R}}(\xi, t) f_i^{\mathbb{R}}(z, Q^2)$$

- The fluxes motivated by Regge theory:

$$f_{\mathbb{P},\mathbb{R}}^P(\xi, t) = A_{\mathbb{P},\mathbb{R}} \frac{e^{B_{\mathbb{P},\mathbb{R}} t}}{\xi^{2\alpha_{\mathbb{P},\mathbb{R}}(t)-1}}, \quad \alpha_{\mathbb{P},\mathbb{R}}(t) = \alpha_{\mathbb{P},\mathbb{R}}(0) + \alpha'_{\mathbb{P},\mathbb{R}} t$$

- Simple form for sea quark (valence =0) and gluon PDFs of the Pomeron:

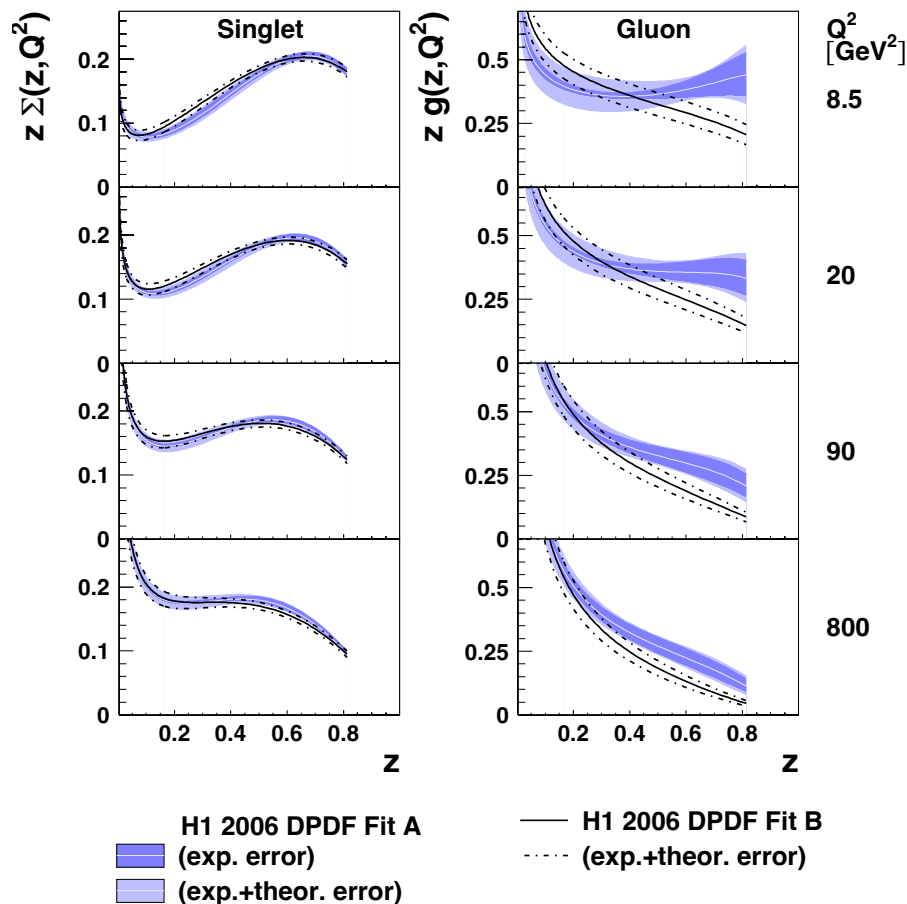
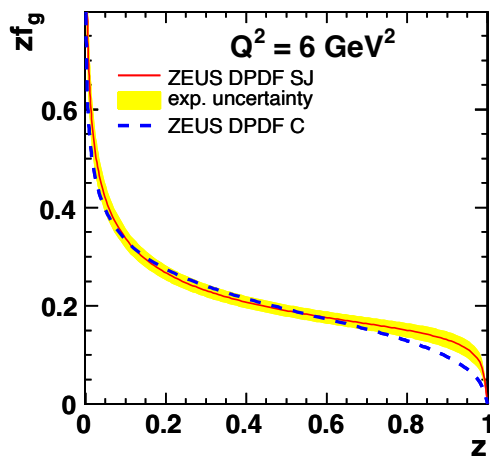
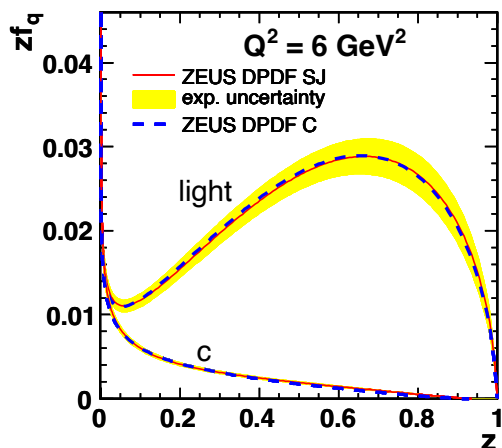
$$z f_i^{\mathbb{P}}(z, \mu_0^2) = A_i z^{B_i} (1-z)^{C_i}$$

- $f_u^{\mathbb{P}} = f_d^{\mathbb{P}} = f_s^{\mathbb{P}}$ , and massless heavy flavors in the variable flavor scheme generated through DGLAP evolution.

- Reggeon  $f_i^{\mathbb{R}}$  taken from pion PDFs → can be better constrained at EIC, [Armesto, Newman, Slominski, Stasto, PRD 110 \(2024\) 5, 054039](#).

# Diffractive PDFs from global fits (2)

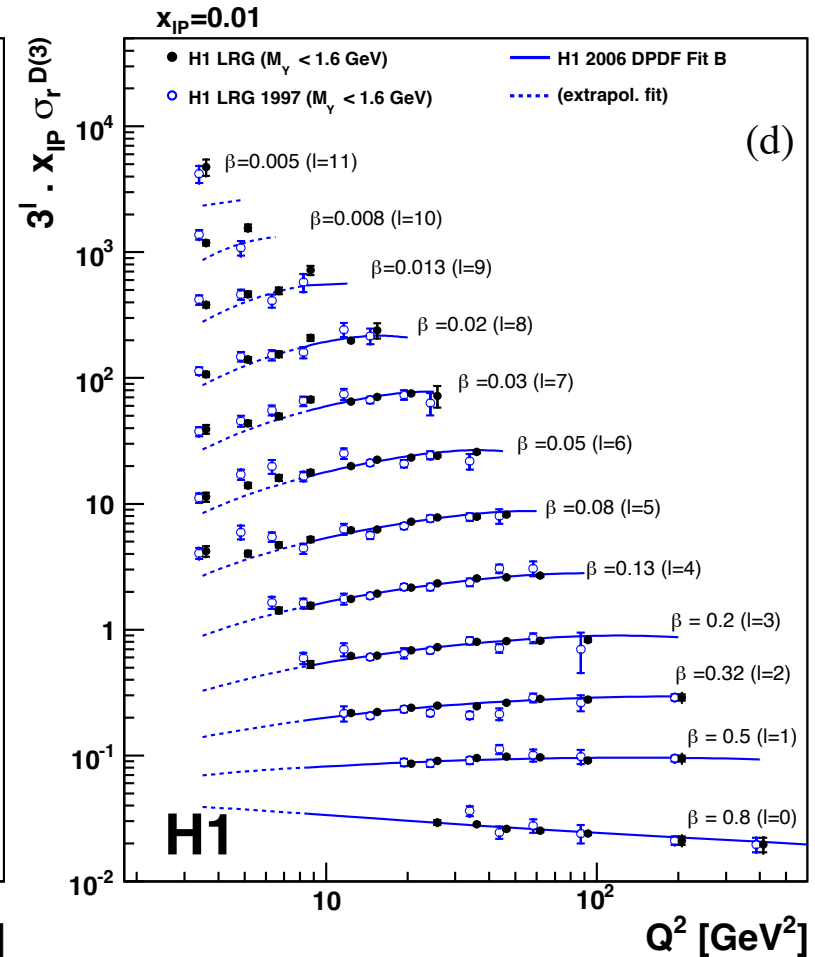
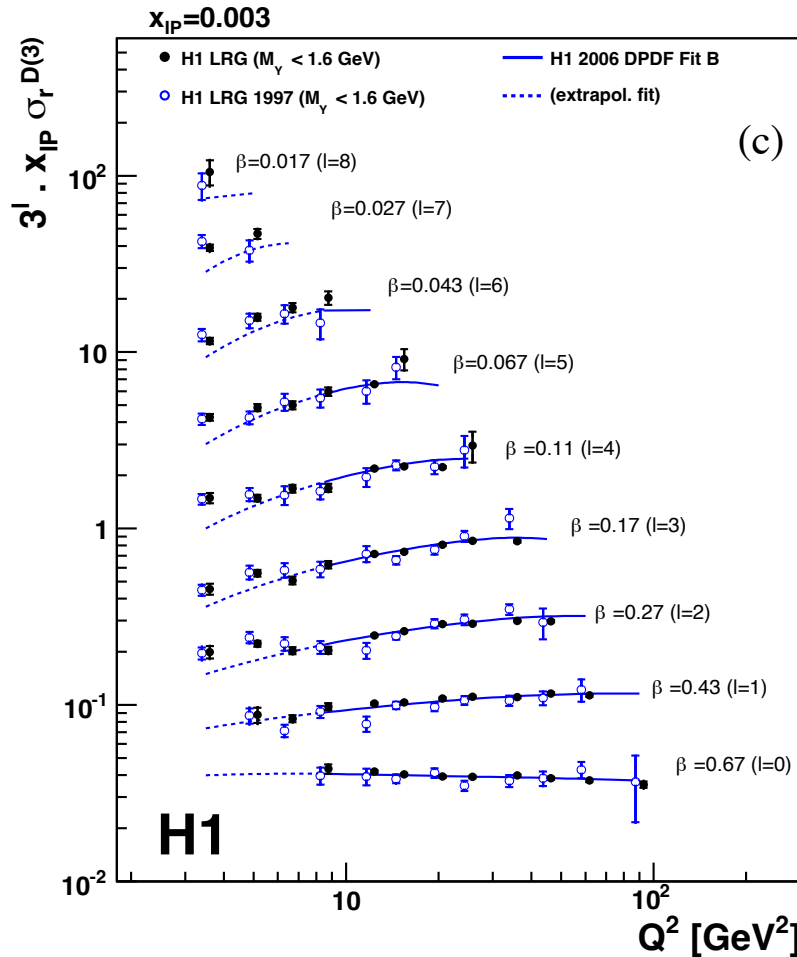
- Most notable examples are ZEUS and H1 analyses of their own data, [Chekanov et al, NPB 831 \(2010\) 1](#); [Aktas et al, EPJC 48 \(2006\) 715](#).



- Diffractive gluon PDF  $\gg$  quark PDFs, but poorly constrained at large  $z \rightarrow$  need to include data on diffractive dijets  $\rightarrow$  favor ZEUS C and H1 B fits.

# Diffraction PDFs from global fits (3)

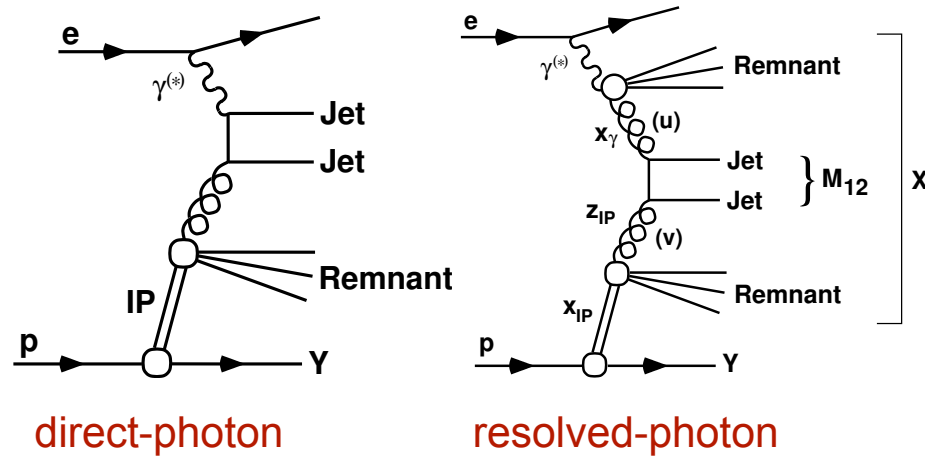
- Good description of original and more recent H1 data, [Aaron et al, EPJC 72 \(2012\) 2074](#)



- Comparison of LRG with proton-tagged cross section measurement  $\rightarrow$   **$\sim 20\%$**  contribution of proton dissociation.

# Diffractive dijet photoproduction

- Collinear fact: same diffractive PDFs for pQCD description of various processes.
- Diffractive dijet electro- and photoproduction in ep scattering → constraints on gluon distribution.



- Cross section is known to NLO accuracy [Klasen, Kramer, Salesch, Z. Phys. C 68, 113 \(1995\)](#); [Klasen, Kramer, Z. Phys. C 72, 107 \(1996\)](#), [Z. Phys. C 76, 67 \(1997\)](#); [Klasen, Rev. Mod. Phys. 74, 1221 \(2002\)](#)

$$d\sigma = \sum_{a,b} \int dy \int dx_\gamma \int dt \int dx_{\mathbb{P}} \int dz_{\mathbb{P}} f_{\gamma/e}(y) f_{a/\gamma}(x_\gamma, M_\gamma^2) f_{\mathbb{P}/p}(x_{\mathbb{P}}, t) f_{b/\mathbb{P}}(z_{\mathbb{P}}, M_{\mathbb{P}}^2) d\hat{\sigma}_{ab}^{(n)}$$

Photon flux in Weizsäcker-Williams approximation

Photon PDFs for resolved photon; in  $a=\gamma$  case,  $f_{a/\gamma}=\delta(1-x_\gamma)$

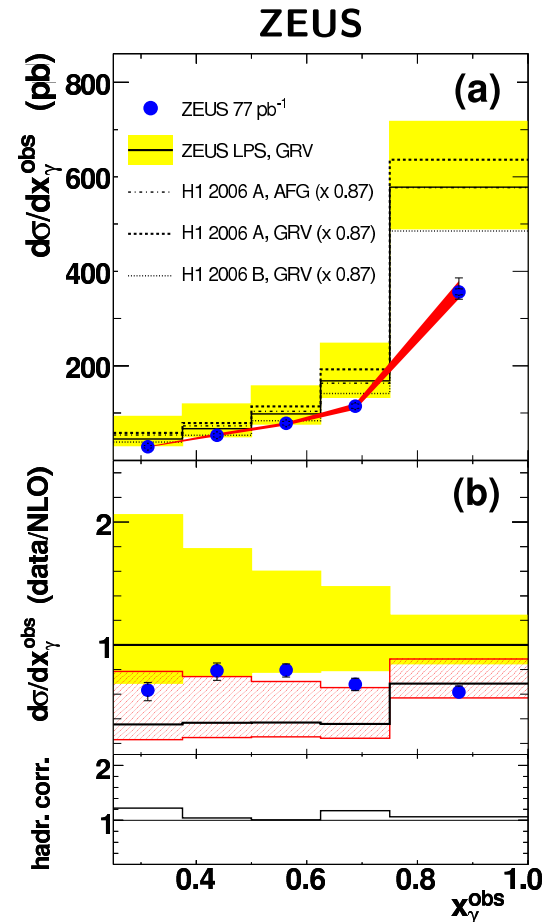
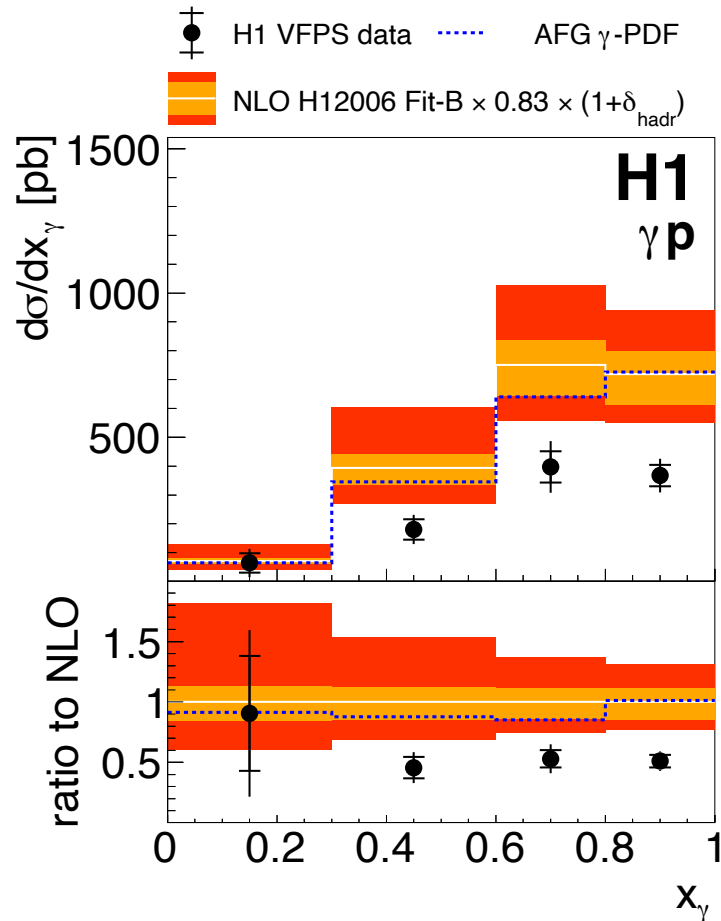
“Pomeron” flux

Diffractive PDFs

Hard partonic cross section

# Diffractive dijet photoproduction (2)

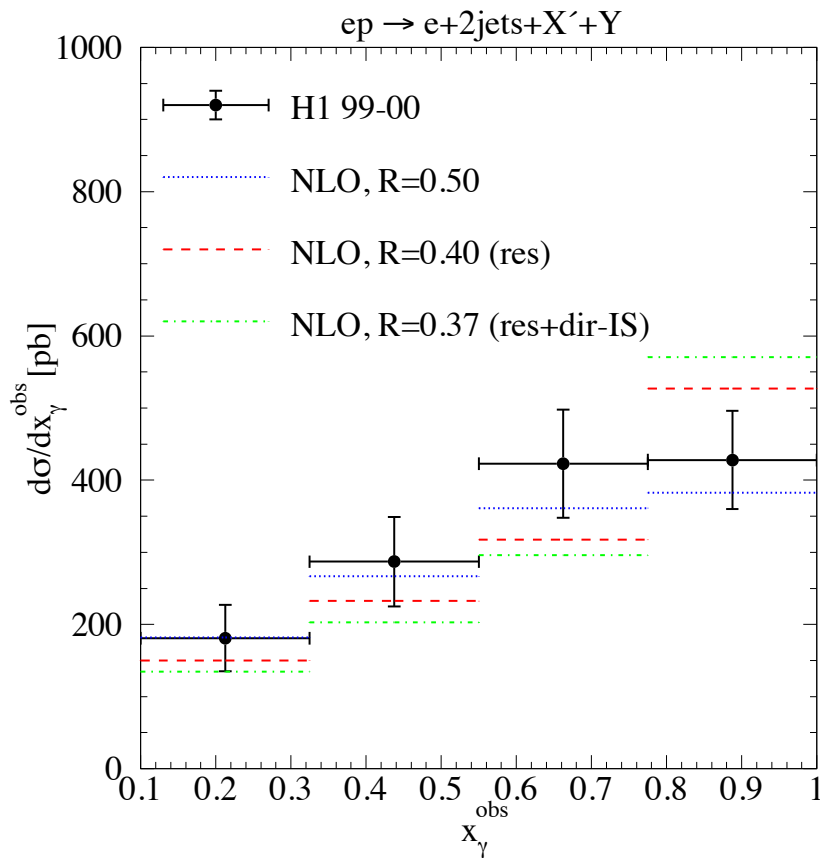
- Universality of diffractive PDFs successfully tested in diffractive dijet and open charm product in DIS, Aktas et al. [H1 Coll.], JHEP 10, 042 (2007); EPJ C 71, 549 (2010); EPJ C 50, 1 (2007); Chekanov et al. [ZEUS Coll.], EPJ C 52, 813 (2007); Chekanov et al. [ZEUS Coll.], NPB 831, 1 (2010)
- At the same time, NLO pQCD QCD overestimates cross sections of diffractive dijet photoproduction at HERA by **factor 2** → factorization breaking, Aktas et al. [H1 Coll.], EPJ C 71, 549 (2007); Aaron et al. [H1 Coll.], EPJ C 70, 15 (2010); Andreev et al. [H1 Coll.], JHEP 05, 056 (2015); Chekanov et al. [ZEUS Coll.], EPJ C 55, 177 (2008).



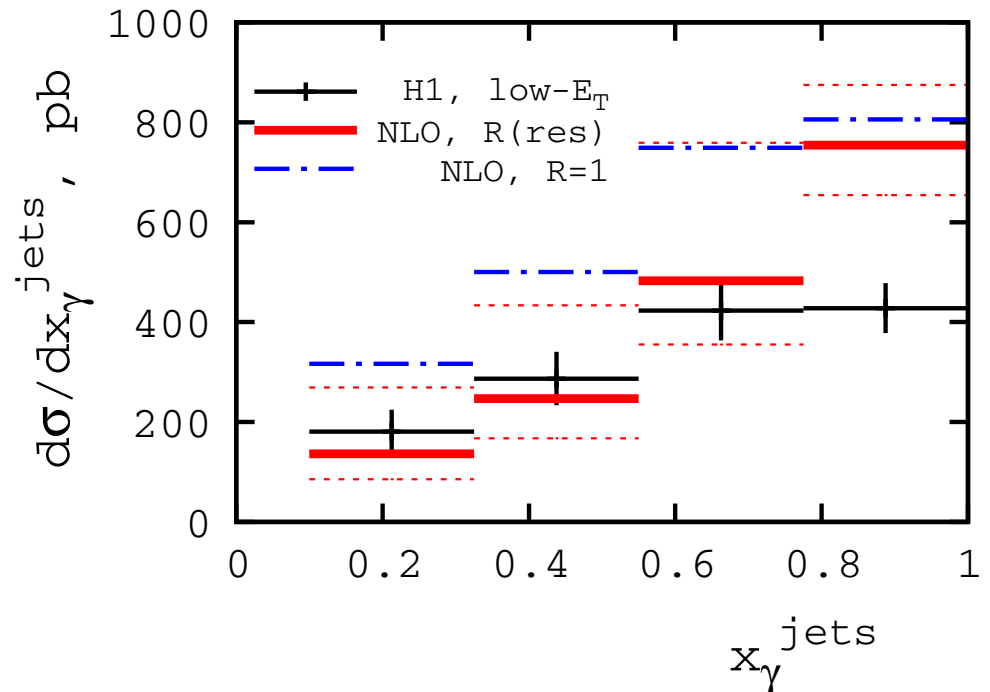


# Factorization breaking

- Mechanism of this factorization breaking remains unknown:
  - global suppression factor  $R \approx 0.5$
  - suppression of only resolved photon contribution by  $R \approx 0.34$  as expected in hadron-hadron scattering, [Kaidalov, Khoze, Martin, Ryskin, PLB 567, 61 \(2003\)](#); [Klasen, Kramer, EPJ C 70, 91 \(2010\)](#)
  - flavor-dependent combination of these mechanisms, [Guzey, Klasen, EPJ C 76, 467 \(2016\)](#)

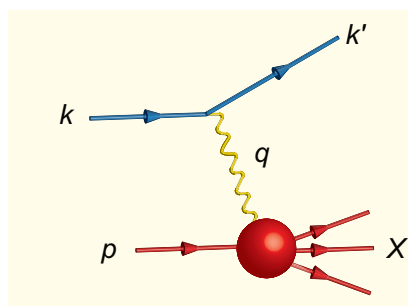


$$S_i^2(x_\gamma) \rightarrow \begin{cases} 1, & i = c, \\ A_q x_\gamma + 0.34, & i = u, d, s, \\ A_g x_\gamma + 0.34, & i = g, \end{cases}$$

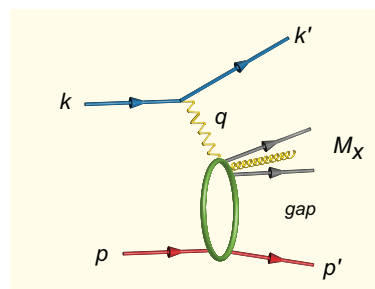


# Diffraction in DIS on nuclei

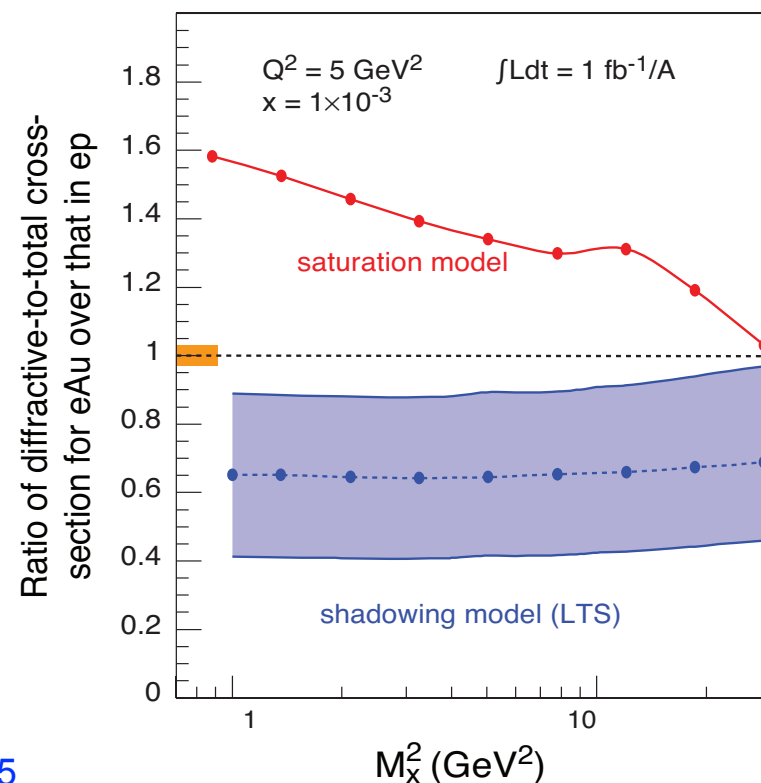
- Never been measured at HERA, but will be at the Electron-Ion Collider (EIC).
- Besides accessing **nuclear diffractive PDFs** for the first time, enhanced sensitivity to non-linear effects, e.g., **gluon saturation** in heavy nuclei.
- Sensitive observable is the ratio of diffractive to total DIS cross sections for a heavy nucleus and the proton, [Accardi et al., EPJ A52 \(2016\) 9, 268 \[1212.1701 \[hep-ex\]\]](#):



Total DIS



Diffractive DIS

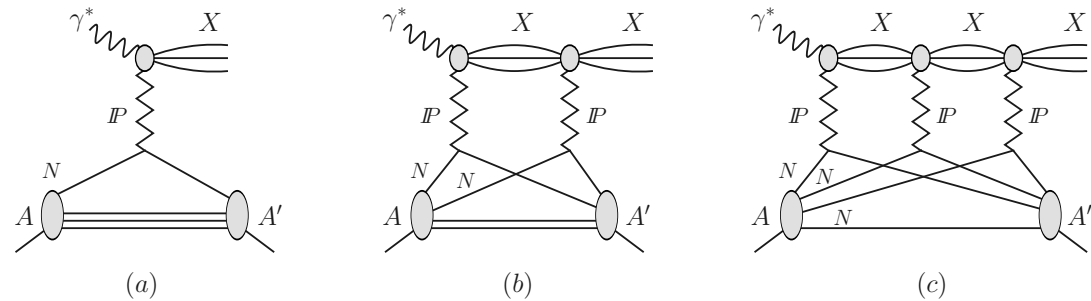


- Predicted to be:

- $R_{\text{diff/tot}} > 1$  due to nuclear enhancement of saturation scale  $Q_{s,A}^2$ , [Kowalski, Lappi, Venugopalan, PRL 100 \(2008\) 022303](#); [Lappi, Le, Mäntysaari, PRD 108 \(2023\) 114023](#)
- $R_{\text{diff/tot}} < 1$  due to strong leading twist nuclear shadowing, [Frankfurt, Guzey, Strikman, Phys. Rept. 512 \(2012\) 255](#)

# Leading twist approach to nuclear shadowing

- Nuclear shadowing: suppression of nuclear PDFs  $f_{i/A}(x, Q^2)/[Af_{i/p}(x, Q^2)] < 1$
- LTA = method to calculate various nuclear parton distributions at small x (usual, generalized, diffractive), Frankfurt, Strikman, EPJ A5 (1999) 293; Frankfurt, Guzey, Strikman, Phys. Rept. 512 (2012) 255 → alternative to global fits of PDFs.
- Based on:
  - Gribov-Glauber model of NS for soft hadron-nucleus scattering
  - QCD factorization theorems for inclusive and diffractive DIS.
- $\gamma^* + A \rightarrow X + A'$  amplitude is a series of diffractive scattering off  $i = 1, 2, \dots, A$  target nucleons:



- Coherent diffraction  $A' = A$ :

$$\sigma_{\gamma^*A \rightarrow XA} = \int d^2\vec{b} |\Gamma_{\gamma^*A \rightarrow XA}(\vec{b})|^2 = 4\pi \frac{d\sigma_{\gamma^*N \rightarrow XN}(t=0)}{dt} \int d^2\vec{b} \left| \int dz \rho_A(\vec{b}, z) e^{iz\Delta_{\gamma^*X}} e^{-\frac{1-i\eta}{2}\sigma_{\text{soft}} \int_z^\infty dz' \rho_A(\vec{b}, z')} \right|^2$$

diffractive cross section on proton measured at HERA

nuclear density

model-dependent cross section

# LTA to nuclear shadowing (2)

- Apply collinear QCD factorization for diffractive DIS, Collins, PRD 57 (1998); PRD 61 (2000) 019902 → from structure function to parton distributions:

$$f_{i/A}^{D(3)}(x, x_{\mathbb{P}}, Q^2) = 4\pi f_{i/p}^{D(4)}(x, x_{\mathbb{P}}, Q^2, t = 0) \int d^2\vec{b} \left| \int dz \rho_A(\vec{b}, z) e^{izx_{\mathbb{P}}m_N} e^{-\frac{1-i\eta}{2}\sigma_{\text{soft}}^i(x) \int_z^\infty dz' \rho_A(\vec{b}, z')} \right|^2$$

$$= f_{i/p}^{D(3)}(x, x_{\mathbb{P}}, Q^2) \frac{1}{\sigma_{\text{el}}^i(x)} \int d^2\vec{b} \left| 1 - e^{-\frac{1-i\eta}{2}\sigma_{\text{soft}}^i(x) T_A(\vec{b})} \right|^2$$

$$\sigma_{\text{el}}(x) = \frac{[\sigma_{\text{soft}}(x)]^2}{16\pi B_{\text{diff}}}$$

$$T_A(\vec{b}) = \int dz \rho(\vec{b}, z)$$

- Transparent interpretation: nuclear diffractive PDFs suppressed (shadowed) in proportion to the nuclear elastic cross section.
- Similarly for quasi-elastic scattering using completeness final states  $A'$ :

$$\sigma_{\gamma^*A \rightarrow XA'} = \int d^2\vec{b} \langle A | \left| \Gamma_{\gamma^*A \rightarrow XA}(\vec{b}) \right|^2 | A \rangle = \sigma_{\gamma^*N \rightarrow XN} \frac{1}{\sigma_{\text{el}}} \int d^2\vec{b} \left( \left| 1 - e^{-\frac{1-i\eta}{2}\sigma_{\text{soft}} T_A(\vec{b})} \right|^2 + e^{-\sigma_{\text{in}} T_A(\vec{b})} - e^{-\sigma_{\text{soft}} T_A(\vec{b})} \right)$$

$$\tilde{f}_{i/A}^{D(3)}(x, x_{\mathbb{P}}, Q^2) = f_{i/p}^{D(3)}(x, x_{\mathbb{P}}, Q^2) \frac{1}{\sigma_{\text{el}}^i(x)} \int d^2\vec{b} \left( \left| 1 - e^{-\frac{1-i\eta}{2}\sigma_{\text{soft}}^i(x) T_A(\vec{b})} \right|^2 + e^{-\sigma_{\text{in}}^i(x) T_A(\vec{b})} - e^{-\sigma_{\text{soft}}^i(x) T_A(\vec{b})} \right)$$

$$\sigma_{\text{in}}(x) = \sigma_{\text{soft}}(x) - \sigma_{\text{el}}(x)$$

- In this case, NS is given by sum of elastic and inelastic nuclear cross sections.

# LTA predictions for nuclear diffractive PDFs

- Assumed that diffractive intermediate states  $X$  do not mix  $\rightarrow$  one free parameter  $\sigma_{\text{soft}}^i(x)$   $\rightarrow$  controls size and uncertainties of LTA predictions.

- High shadowing:** given by probability of diffraction

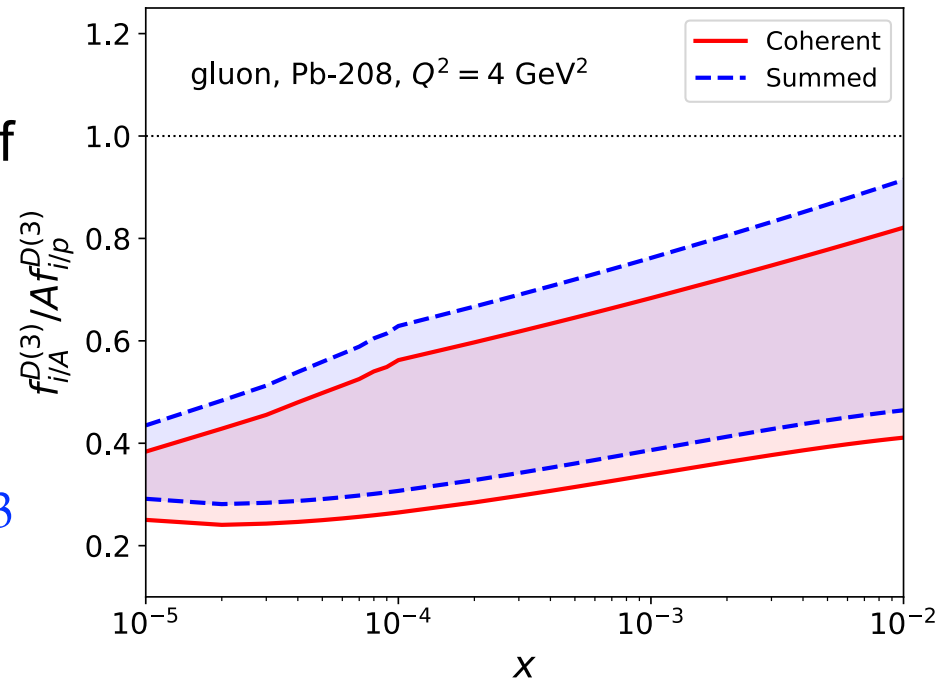
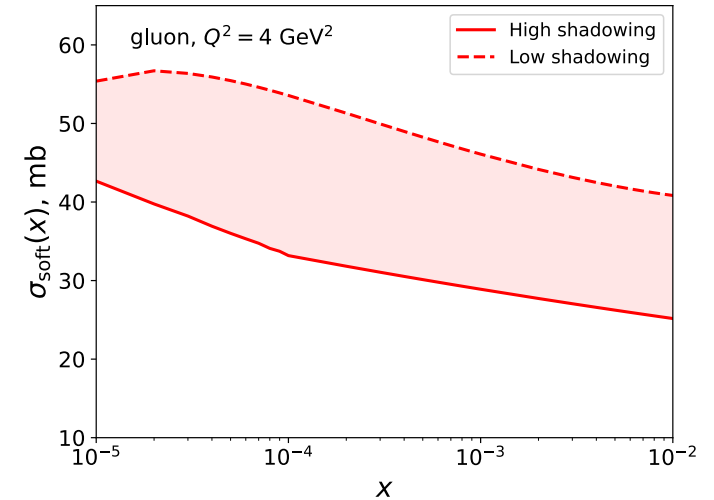
$$\sigma_{\text{soft}}^i(x) \approx \sigma_2(x) \equiv \frac{16\pi}{f_{i/p}(x)} \int_x^{0.1} \frac{dx_{\mathbb{P}}}{x_{\mathbb{P}}} f_{i/p}^{D(4)}(x, x_{\mathbb{P}}, t=0)$$

- Low shadowing:** calculated using model for hadronic structure of  $\rho$  meson.

- In LTA, nuclear shadowing driven by diffraction on proton  $\rightarrow$  10-15% probability of diffraction in DIS@HERA leads to **large suppression** of nuclear PDFs at small  $x$ .

- Compare to impulse approximation (IA):

$$\frac{f_{i/A}^{D(3)}}{A f_{i/p}^{D(3)}} = \frac{4\pi B_{\text{diff}}}{A} \int d^2\vec{b} (T_A(\vec{b}))^2 = \frac{B_{\text{diff}}}{A} \int dt F_A^2(t) = 4.3$$

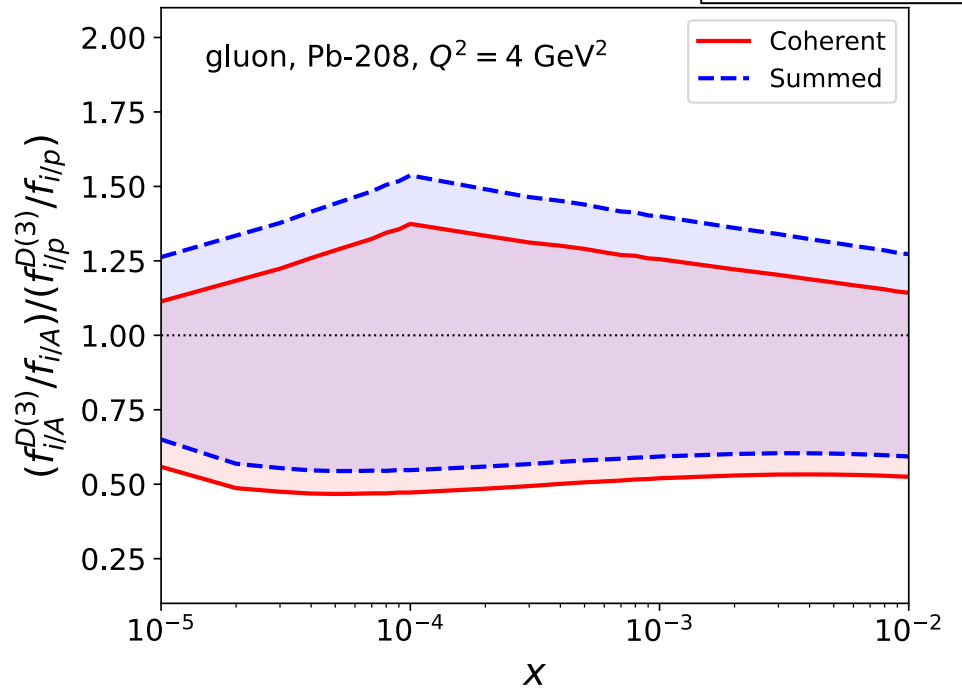
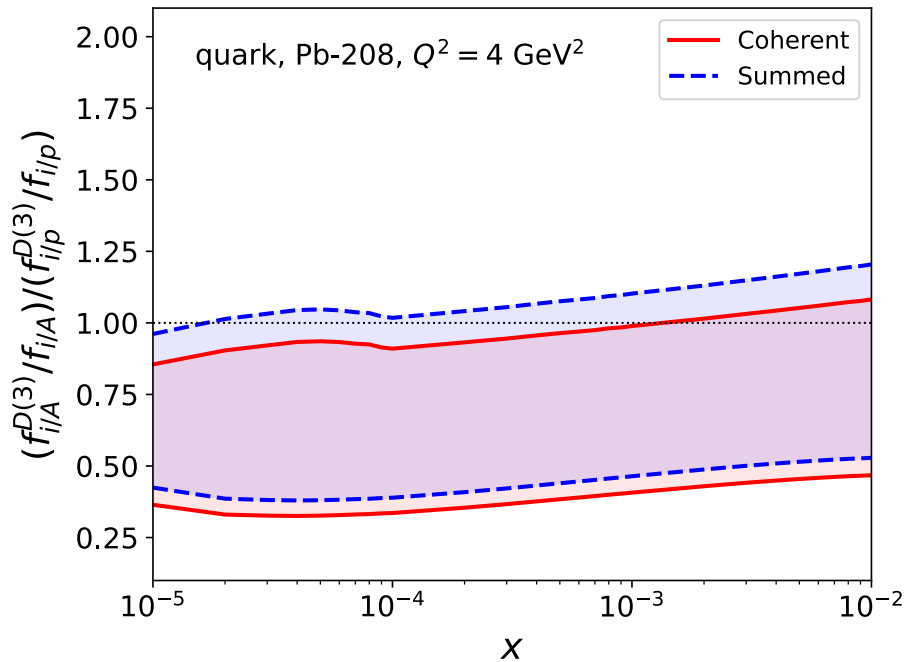


# LTA predictions for $R_{\text{diff/tot}}$

- Combine LTA predictions for diffractive and usual nuclear PDFs:

$$\frac{f_{i/A}^{D(3)}(x, x_{\mathbb{P}}, Q^2)/f_{i/A}(x, Q^2)}{f_{i/p}^{D(3)}(x, x_{\mathbb{P}}, Q^2)/f_{i/p}(x, Q^2)} = \frac{\sigma_{\text{soft}}^i(x)}{\sigma_{\text{el}}^i(x)} \frac{\int d^2\vec{b} \left| 1 - e^{-\frac{1-i\eta}{2}\sigma_{\text{soft}}^i(x)T_A(\vec{b})} \right|^2}{2(1 - \lambda^i(x))\Re \int d^2\vec{b} \left( 1 - e^{-\frac{1-i\eta}{2}\sigma_{\text{soft}}^i(x)T_A(\vec{b})} \right) + \lambda^i(x)A\sigma_{\text{soft}}^i(x)}$$

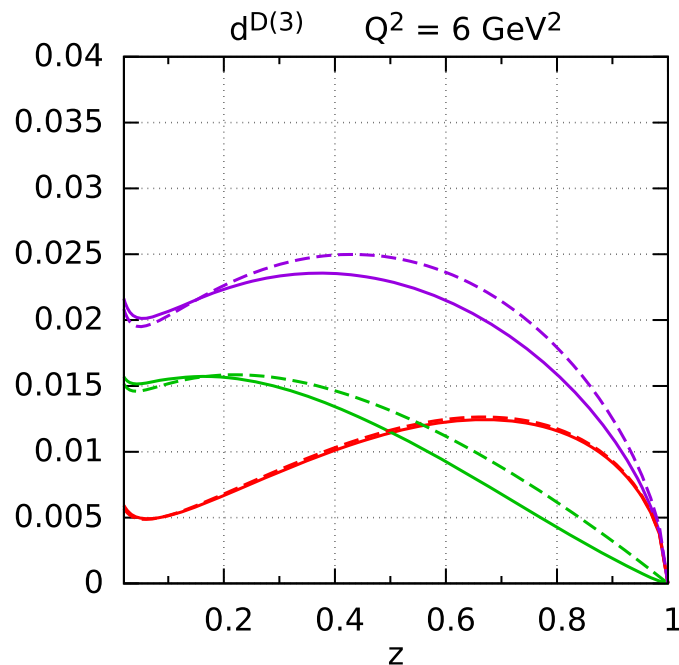
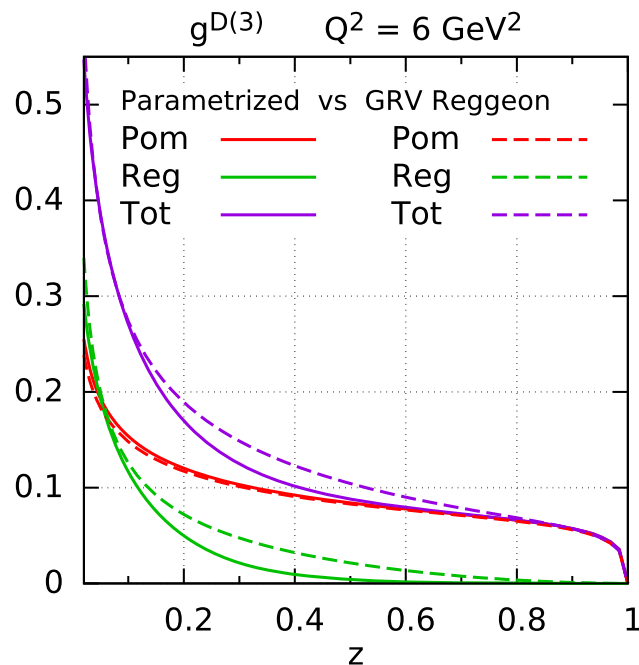
$$\lambda^i(x) = 1 - \sigma_2^i(x)/\sigma_{\text{soft}}^i(x)$$



- Suppression  $R_{\text{diff/tot}} \approx 0.5 - 1$  (quarks) and  $R_{\text{diff/tot}} \approx 0.5 - 1.3$  (gluons) due to interplay of large leading twist nuclear shadowing for diffractive and usual nuclear PDFs.

# Outlook: diffraction at EIC

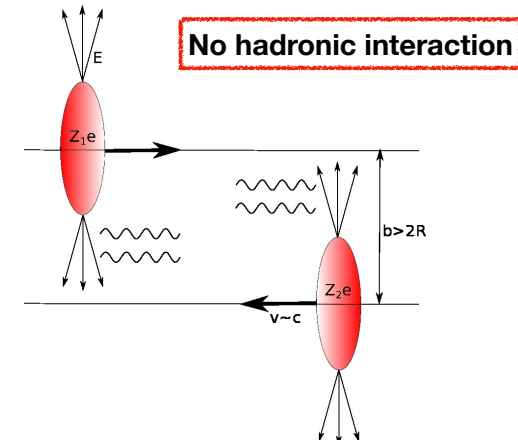
- Several recent global QCD fits for **proton** diffractive PDFs using all (inclusive+dijets) HERA data, [Salajeghen et al., PRD 107 \(2024\) 9, 093038](#); [PRD 106 \(2022\) 5, 054012](#)
- Further progress possible at the Electron-Ion Collider (EIC)  $\rightarrow \sqrt{s_{ep}} \sim 100$  GeV lower than at HERA  $\rightarrow$  constrain sub-leading (Reggeon) contribution at large  $\xi$ , [Armesto, Newman, Slominski, Stasto, PRD 110 \(2024\) 5, 054039](#)



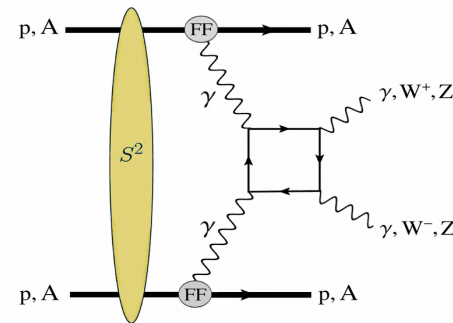
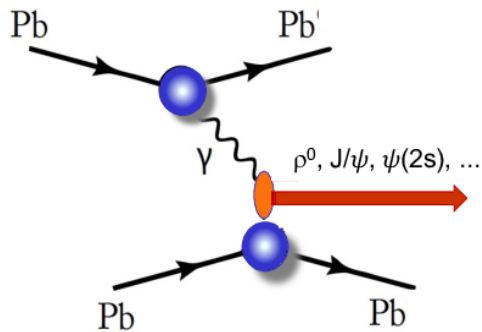
- Similarly, NLO pQCD predicts **10-35%** contribution of sub-leading Reggeon trajectory for  $x_P > 0.06$  in diffractive dijet photoproduction at EIC, [Guzey, Klasen, JHEP 05 \(2020\) 074](#)

# Outlook: diffraction in UPC at LHC

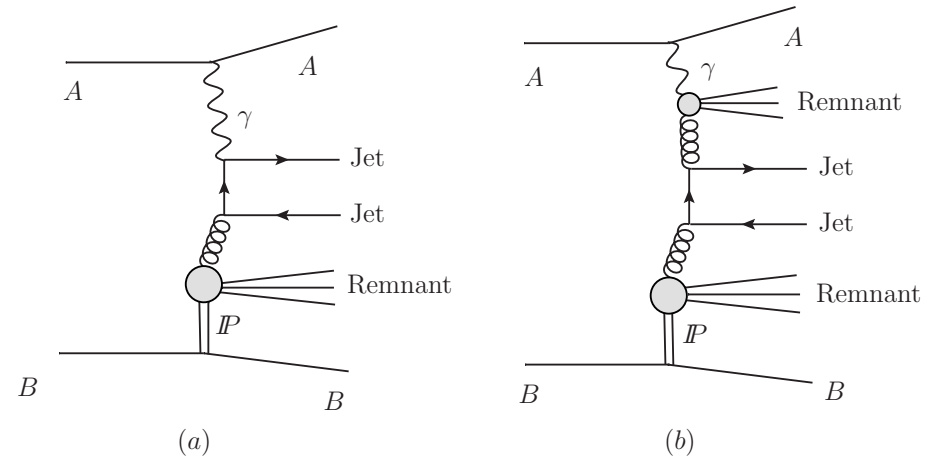
- **Ultrapерipheral collisions (UPCs)** at large transverse distances  $|\vec{b}| \gg 2R_A \rightarrow$  interaction via quasi-real photons  $\rightarrow$  LHC is a high-energy *photon collider*  $\rightarrow$  can be used to test predictions for EIC.



- Photon-nucleus (proton) and photon-photon scattering in UPCs:



- Recent ATLAS measurement in  $0nXn$  channel, [Aad et al, 2409.11060 \[nucl-ex\]](#)  $\rightarrow$  can be extended to  $0n0n$  channel probing **nuclear diffractive PDFs**.





# Outlook: diffraction in UPC at LHC (2)

- Besides nuclear diffractive PDFs, **heavy nucleus** can be used to suppress the resolved photon contribution → new handle on mechanism of factorization breaking, [Guzey, Klasen, JHEP 04 \(2016\) 158](#)

- Suppression factor for resolved photon: 
$$R(\text{res.}) = \frac{\int d^2b |\mathcal{A}_{\gamma T \rightarrow VT}(W, b)|^2 P_{VT}(W, b)}{\int d^2b |\mathcal{A}_{\gamma T \rightarrow VT}(W, b)|^2}$$

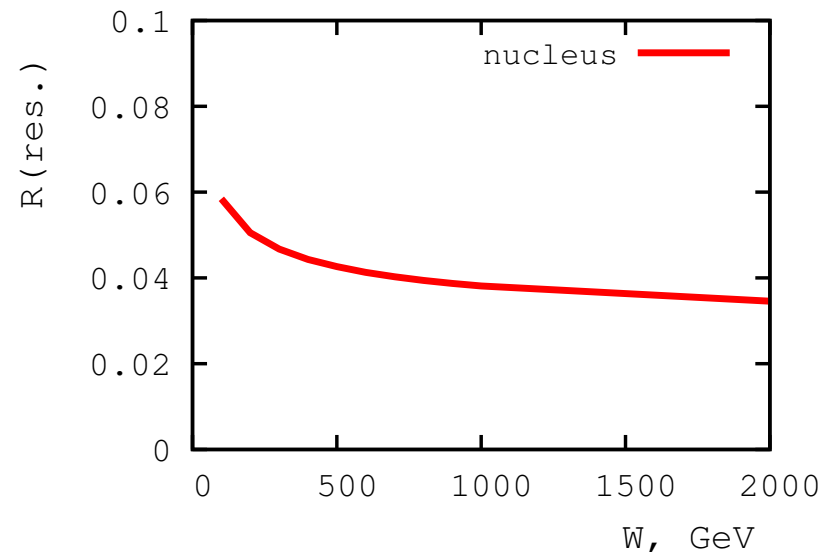
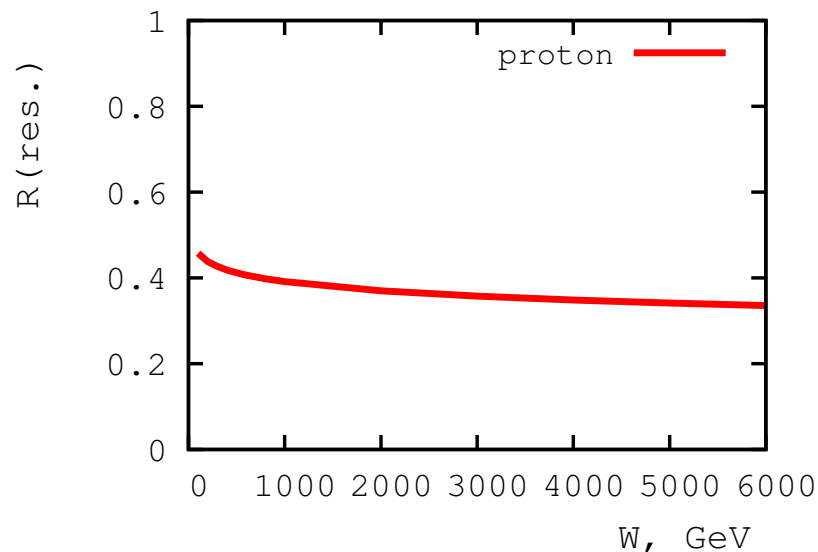
VM photoproduction  
in Glauber model

$$\mathcal{A}_{\gamma A \rightarrow VA}(W, b) = \frac{e}{f_V} \left( 1 - e^{-\frac{\sigma_{\rho N}(W)}{2} T_A(b)} \right)$$

Probability to not have  
strong interactions

$$P_{VA}(W, b) = e^{-\sigma_{\rho N}(W) T_A(b)}$$

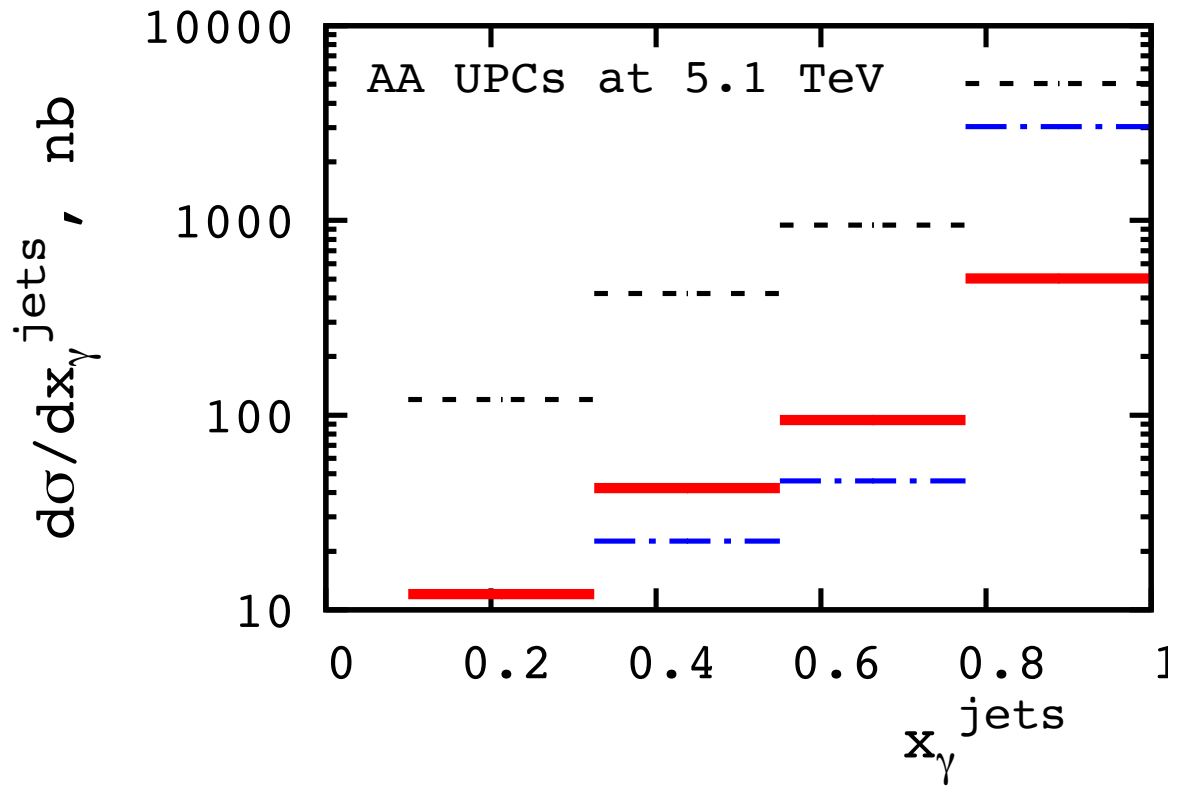
- Suppression factor: proton vs. nucleus



# Outlook: diffraction in UPC at LHC (3)

- It is much easier to break up nucleus and fill the rapidity gap  $\rightarrow R(\text{res})_A \ll R(\text{res})_p$
- NLO pQCD predictions for diffractive dijet photoproduction in Pb-Pb UPCs at LHC, [Guzey, Klasen, JHEP 04 \(2016\) 158](#)

—  $R(\text{glob.})=0.1$   
- · -  $R(\text{res.})=0.04$   
- - -  $R=1$



# Summary

- Diffractive PDFs is a standard tool of perturbative QCD.
- In the proton case, they are extracted using global fits to HERA data.
- Further progress at EIC for the sub-leading contribution for large  $\xi$ .
- In the nuclear case, diffractive PDFs never been measured  $\rightarrow$  will be at EIC and can be in Pb-Pb UPCs at LHC.
- Open question: mechanism of factorization breaking in diffractive dijet photoproduction  $\rightarrow$  can be addressed in UPCs at LHC.
- Diffraction in ep and eA can be alternatively addressed in the dipole model, where the emphasis on the gluon saturation.