

April 1, 2024



Overview

- Where were T2K and NOvA when we started joint analysis and what were we trying to do
 - For how we ended up where we did see Ryan's talk
- Some key choices we made and why
- What might we do differently if circumstances were different
- Caveat: These are my personal opinions having been through the process and not official views



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Background



Long-baseline oscillation experiments **Next Generation Current Generation Previous Generation** DEEP UNDERGROUND NEUTRINO EXPERIMENT MINOS K2K **Hyper-Kamiokande OPERA** 2010 2020 2030 **NOvA collects first** beam data. T2K collects first beam data. [1] T2K: Eur. Phys. J. C (2023) 83:782 (2023) Published dataset^[1,2] until 2020 by both [2] NOvA: Phys. Rev D 106, 032004 (2022) (Frequentist) and arXiv:2311.07835 (Bayesian)

experiments. Our results use this dataset!

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Flux Model Cross Section ND Model Data ND Detector Model External ND Fit Data Constraints/ These fits can also be combined into a single ND simultaneous fit Constrained Model FD FD Detector Data Model Oscillation Fit Oscillation **Parameters**

Analysis Strategy

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 The experiments have different analysis approaches driven by contrasting detector designs.



Construction of the individual analyses

- Central value models and systematic uncertainties around them are not the same
- Different MC generators are used by both experiments (T2K: NEUT, NOvA: Genie)
- Analysis codes used are not the same and do not simultaneously compile against each other





Why NOvA-T2K joint fit?

- The complementarity between the experiments provides the power to break degeneracies and improve oscillation sensitivity
- Full implementation of:

Energy reconstruction and detector response
 Detailed likelihood from each experiment
 Consistent statistical inference across the full dimensionality

• In-depth review of:

Models, systematic uncertainties and possible correlations

Different analysis approaches driven by contrasting detector designs



Constructing the NOvA-T2K joint analysis



Constructing the joint-analysis

- A joint fit needs a joint likelihood to interrogate
- Components of this are:
 - Poisson likelihood comparing data to predictions as a function of model parameters
 - Penalty terms from the priors on those parameters
 - 3. External constraints on θ_{13} , θ_{12} , Δm_{21}^2 from solar and reactor neutrino experiments





Constructing the joint-analysis

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solar and reactor neutrino experiments





2 and 3 are easy, 1 is hard IMPERIAL

Sharing the sample likelihood

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- Each experiment has spent a lot of time and personpower making their 'fitter' able to predict their FD event rates
- This is not implemented in a way that makes it easily portable to another codebase
- These codebases do not compile against each other simultaneously due to dependency versioning etc.



Red represents must be T2K codebase & blue shows must be NOvA codebase.



Constructing the joint-analysis

- Solution is to compartmentalise analysis using containers
- Both T2K and NOvA's Bayesian analysis frameworks were packaged into containers
- 'Bridge' into and out of the container was constructed that returned LLH when given parameter values
 - Link to bifrost: <u>https://github.com/nova-t2k/bifrost/tree/main</u>
- Existing analyses then able to interrogate other experiment's container



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 $\mathbf{NO}_{\mathbf{VA}}$

Was this a good idea?

Pros:

- Existing analyses can be used as is
 - No arguing over who overhauls code to be compatible
- Internals are hidden from the other experiment easing political issues around data sharing

Cons:

- Internals are hidden from the other experiment making validation harder
 - Eg checking all data spectra being fit for a given study cannot be done by one person

Conclusion:

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- Right thing for T2K+NOvA:
 - We started after experiments had been running for years, had existing analyses, no common/compatible fitting frameworks
 - Political will was not there for full openness at start of effort and personpower was not
 present for full code overhaul
 - This approach let us surmount real obstacles we had no other solution for



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Pros:

- Existing analyses can be used as is
 - No arguing over who overhauls code to be compatible
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Conclusion:

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- Maybe not the right thing if you start earlier
 - Feedback and validation loops being long caused real issues taking up analyser time and making studies take longer to complete
 - We never had problems with one collaboration trying to reverse engineer the other's result and data sharing worries relaxed fairly quickly once we were working together

Cons:

- Internals are hidden from the other experiment making validation harder
 - Eg checking all data spectra being fit for a given study cannot be done by one person



Challenge: When? What? How? to correlate common physics parameters between the two experiments.



• **Challenge:** No direct mapping between the cross-section systematics parameters

*Phys. Rev. D 86, 053003





- Challenge: No direct mapping between the cross-section systematics parameters
- Strategy: Explore a range of artificially crafted scenarios to bracket the impact of possible correlations
 - Example: Fabricated systematics equal in size to total statistical uncertainty, causing a correlated bias in the oscillation dip across both experiments.
 - Uncorrelated and correctly correlated (full correlation) credible intervals agree with negligible differences, while incorrectly correlating systematics shows a bias.





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Lessons: Shared models with ability to describe all targets and energy ranges of T2K and NOvA not currently mature so these studies were key to joint fit feasibility
Caveat: All of these studies are specific to T2K and NOvA at current exposure. Assumption will break down at some point





Cross-section: Impact of alternate models

- Evaluate the robustness of the fit against various alternate models
- Generated simulated fake data using reweighting to alternate models for both the near and far detector, then analyze the credible intervals of the full joint-fit
- Compare results against nominal fit and make sure certain criteria are met



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*Phys. Rev. D 100, 072005 (2019)

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- Generated simulated fake data using reweighting to alternate models for both the near and far detector, then analyze the credible intervals of the full joint-fit
- Compare results against nominal fit and make sure certain criteria are met
 - Today's focus is on ability to do these studies in a joint fit, many people present who can discuss pros/cons of these studies in oscillation analyses generally
 - Key problem is ability to actually look at the same model in both experiments at the same time



*Phys. Rev. D 100, 072005 (2019)

Alternate model tests

- T2K uses Neut, NOvA uses Genie. Neither experiment has run a production with the other generator for oscillation analysis for a long time
- Each experiment's analysis and simulation toolchain cannot easily take input from the other's generator
- How then do you actually fit data generated with the same model in both experiments at the same time?



Patrick Dunne

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Alternate model tests

- Answer is reweighting with a careful choice of variables and binning to ensure fit variables are properly described
- Long process requiring much care and hard validations as answer isn't obvious



Analysis	Analysis Variable	Truth Proxy
	$E_{ u}$	$E_{ u}$
T2K ND + FD	$p_\ell^{ m reco}$	p_ℓ
	$ heta_\ell^{ m reco}$	$ heta_\ell$
	$p_{{\scriptscriptstyle \mathrm{T}},\ell}^{\mathrm{reco}}$	$p_{ extsf{tr},\ell}$
NOvA ND + FD	$E_{ u}^{ m reco}$	$E_{ u}$
	$y^{ m Vis.} = E_{ m Had}^{ m reco}/E_{ u}^{ m reco}$	$E_{ m Av.} = \sum_{p,\pi^{\pm}} T + \sum_{e,\gamma,\pi^0,{ m K},{ m Other}} E$
		$+\sum_{B_{ ext{strange}}}{(E-m_p)}+\sum_{ar{p}}{(E+m_p)}$
2		

A better way! - NuHEPMC

- If you were starting from scratch why not agree on the format for generator output/analysis simulation input to directly generate data with same model
- Collider physics world has had this for a long time starting with the Les Houches accords ~2001 and now leading to LHE files/HEPMC3
- NuHEPMC effort uses HEPMC3 format from collider physics with agreed extra neutrino information to make a common generator output format

common /HEPRUP/ IDBMUP(2), EBMUP(2), PDFGUP(2), PDFSUP(2),
+ IDWTUP, NPRUP, XSECUP(MAXPUP), XERRUP(MAXPUP),
+ XMAXUP(MAXPUP), LPRUP(MAXPUP)
common /HEPEUP/ NUP, IDPRUP, XWGTUP, SCALUP, AQEDUP, AQCDUP,
+ IDUP(MAXNUP), ISTUP(MAXNUP), MOTHUP(2,MAXNUP),
+ ICOLUP(2,MAXNUP), PUP(5,MAXNUP), VTIMUP(MAXNUP),
+ SPINUP(MAXNUP)

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Computer Physics Communications Volume 260, March 2021, 107310



The HepMC3 event record library for Monte Carlo event generators

x, x



NuHEPMC

A better way! - NuHEPMC



- If all experiments going forward agree to be able to take in this format, alternate model studies can be done much faster and much more robustly
 - With another hat on the DUNE Phase II ND reco can already use NuHEPMC
- Not only useful for alternate model studies in joint fits but also in standalone analyses
- Also opens up the possibility for shared central value models and uncertainties

- T2K+NOvA analysis used several innovative methods that allowed us to combine already designed analyses
- Choices driven by context when analysis started that one might revisit if starting from scratch with new experiments:



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- Choices driven by context when analysis started that one might revisit if starting from scratch with new experiments:
- 1. Containerised LLHs
 - Let you simultaneously call code that isn't compatible
 - Meaningfully increase validation cycles and make some studies hard to run





- T2K+NOvA analysis used several innovative methods that allowed us to combine already designed analyses
- Choices driven by context when analysis started that one might revisit if starting from scratch with new experiments:
- 2. Correlation impact checks
 - Systematics are sub-dominant for our input analyses so it was plausible that correlations could be negligible
 - This made it worth testing if they were to avoid needing unified xsec model
 - This assumption will break at some point plus the next generation of experiments might need the more sophisticated unified models anyway







- T2K+NOvA analysis used several innovative methods that allowed us to combine already designed analyses
- Choices driven by context when analysis started that one might revisit if starting from scratch with new experiments:
- 3. Parameterised reconstruction mapping
 - Allowed models not simulated by other experiment to be tested
 - Not exact and needed careful thought for each study on whether variables used were right
 - Newer techniques like NuHEPMC would allow other generator models to go through exact reco





- Resulting analysis lifted degeneracies and meaningfully improved oscillation parameter constraints
- Process itself was valuable: I certainly learned a lot about both experiments and the working group continues to be a great source of new ideas for oscillation analyses





Backup



• Challenge: When? What? How? to correlate common physics

Flux Model

parameters between the two experiments.

• Strategy:

Detector Model

- □ Is the overall impact negligible on the result?
- Do we expect any correlations between the experiments?

Cross Section Model

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□ Is the impact of the correlations negligible on the result?



• Different energies

- Different tuning to external data
 - thin target vs thick target data
- Enters the analysis differently

No significant correlations between the experiments



• Different energies

Flux Model

Detector Model

- Different tuning to external data
 - thin target vs thick target data
- Enters the analysis differently
- Different detector design and targets
- Different selections
 - inclusive vs exclusive outgoing pions
- Different energy reconstruction
 - calorimetric vs lepton kinematics

No significant correlations between the experiments

 Explored possible correlations
 between leptonic energy scales; pion and neutron secondary interactions



• Different energies

Flux Model

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Mode	els & Systematics	٦
Flux Model	 Different tuning to external data thin target vs thick target data Enters the analysis differently 	No significant correlations between the experiments
Detector Model	 Different detector design and targets Different selections inclusive vs exclusive outgoing pions Different energy reconstruction calorimetric vs lepton kinematics 	No significant correlations between the experiments
Cross Section Model	 As the underlying physics is fundamentally the same, we expect correlations Different neutrino interaction models optimized for different energy ranges Systematics are designed for individual models and analysis strategies 35 	Investigate the impact of models and correlations on the joint analysis

- Challenge: No direct mapping between the cross-section systematics parameters
 - Exception: Uncertainties in ν_e / ν_μ and $\overline{\nu}_e / \overline{\nu}_\mu$ crosssection have identical origin* and similar treatment
 - Fully correlated in the joint fit.

*Phys. Rev. D 86, 053003


Cross-section: Impact of correlations

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TZK

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Cross-section: Impact of alternate models

- Evaluate the robustness of the fit against various alternate models
- Generated simulated fake data using reweighting to alternate models for both the near and far detector, then analyze the credible intervals of the full joint-fit
- Pre-decided thresholds for bias:
 - Change in the width of the 1D intervals <10%
 - Change in central value < 50% of systematic uncertainty
- Example: Suppression in single pion channel based on tune to the MINERvA data*



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*Phys. Rev. D 100, 072005 (2019)

Cross-section: Impact of alternate models

- Example: Suppression in single pion channel based on the tune to the MINERvA data*
- Additional tests:
 - Cross-experiment models after the ND constraint
 - Impact of alternative nuclear response model: HF-CRPA**
 - Full list available in backup
- No alternate model tests failed the preset threshold bias criteria. *Phys. Rev. D 100, 072005 (2019)



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Mode	els & Systematics	
	Different energies	No significant correlations between
Flux Model	 Different tuning to external data 	the experiments
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	Enters the analysis differently	J
	 Different detector design and targets 	ן
Detector Model	 Different selections 	No significant correlations between
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	 Different energy reconstruction 	
	 calorimetric vs lepton kinematics 	
	 As the underlying physics is fundamentally 	Impact of correlations is negligible on the
Cross Section Model	the same, we expect correlations	results at the current statistical
	 Different neutrino interaction models 	significance.
	 optimized for different energy ranges 	Merits continued investigations for higher
	 Systematics are designed for individual 	data exposures.
	models and analysis strategies 41	

Why NOvA-T2K joint fit?

- The complementarity between the experiments provides the power to **break degeneracies**.
- Full implementation of:



Energy reconstruction and detector response

Detailed likelihood from each experiment

Consistent statistical inference across the full dimensionality

• In-depth review of:



Models, systematic uncertainties and possible correlations



Different analysis approaches driven by contrasting detector designs.



Data Results





FD Data Samples

- The joint-fit uses the data collected by each experiment up until 2020.
- Using both experiments data roughly doubles the total statistics at the far detectors.

Channel	NOvA	T2K
v _e	82	94 (ν _e 0π)
		14 (ν _e 1π)
$\overline{\nu}_{e}$	33	16
$ u_{\mu}$	211	318
$\overline{ u}_{\mu}$	105	137



Compatibility of datasets

- Posterior predictive p-values (PPP)*
 - Compare likelihood best fit to data and fluctuated predictions
 - A good PPP is around 0.5
- The data from both experiments is described well by the joint fit.

Channel	NOvA	T2K	Combined
v _e	0.90	0.19 (ν _e) 0.79 (ν _e 1π)	0.62
$\overline{\nu}_{e}$	0.21	0.67	0.40
$ u_{\mu}$	0.68	0.48	0.62
$\overline{ u}_{\mu}$	0.38	0.87	0.72
Total	0.64	0.72	0.75
nosterior predictive p-value			

posterior predictive p-value

*Statistica Sinica, vol. 6, no. 4, 1996, pp. 733-60. JSTOR



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Mixing angles: $\theta_{23} \& \theta_{13}$

• Without any external constraint from reactor experiments, long-baseline measurements have a degeneracy in $\sin^2 \theta_{23}$ and $\sin^2 2\theta_{13}$ parameters.





Mixing angles: $\theta_{23} \& \theta_{13}$

• Without any external constraint from reactor experiments, long-baseline measurements have a degeneracy in $\sin^2 \theta_{23}$ and $\sin^2 2\theta_{13}$ parameters.

• Using the average constraint on $\sin^2 2\theta_{13} = 0.085 \pm 0.0027$ [PDG 2020], restricts us to a narrow posterior in θ_{13} and lifts this degeneracy.



Mixing angles: $\theta_{23} \& \theta_{13}$

- No significant preference for either octant from the joint-analysis.
- This preference shifts to a small but still insignificant preference for the upper octant when the reactor constraint on θ_{13} is applied.

	NOvA - T2K w/o reactor	NOvA – T2K – w/ reactor
Bayes factor	1.17 Lower Octant/Upper Octant ~54% : ~46% posterior	3.59 Upper Octant/Lower Octant ~78% : 22% posterior

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Δm^2_{32} and $sin^2\,\theta_{23}$

- Marginalizing over each mass ordering, we note a small but distinct difference in the $\sin^2 \theta_{23}$ and Δm_{32}^2 phase space.
- Measurements remain consistent with the maximal mixing hypothesis for θ_{23} mixing angle.



Mass Ordering





CP Violation

- For both mass orderings, $\delta_{CP} = +\pi/2$ lies outside 3-sigma credible interval.
- Normal Ordering allows for a broad range of permissible $\delta_{\rm CP}$
- For the Inverted Ordering, CP conserving values of δ_{CP} (0, π) lie outside the 3-sigma credible interval.





CP Violation: Jarlskog

 Jarlskog-invariant is a parameterization independent way* to measure CP violation.

 $J = \sin \theta_{13} \cos^2 \theta_{13} \sin \theta_{12} \cos \theta_{12} \sin \theta_{23} \cos \theta_{23} \sin \delta_{\rm CP}$

J=0: CP-Conservation $J \neq 0$: CP-Violation

- For Normal Ordering, a considerably wider range of probable values for J
- J = 0 lies outside the 3σ interval for the Inverted Ordering
 - for priors that are both uniform in δ_{CP} and uniform in sin δ_{CP} *Phys. Rev. D 100, 053004 (2019)





Comparisons





Comparison with NOvA-only & T2K-only fits

- The joint-fit prefers the region in the normal ordering where the two individual experiment's preference overlaps as you'd expect.
- There is a tighter constraint in the Inverted
 Ordering where NOvA-only and T2K-only had the same best fit point.



Comparison with NOvA-only & T2K-only fits

- The 1D posterior in Δm_{32}^2 highlights the switch in the mass ordering preference when NOvA and T2K are combined.
- The joint-fit enhances the precision of Δm_{32}^2 over individual experiments.



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Bayes factor

Global Comparisons - Δm_{32}^2

• This analysis has the smallest uncertainty on $|\Delta m_{32}^2|$ as compared to other previous measurements.





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Global Comparisons - $\delta_{ ext{CP}}$

- The δ_{CP} measurements are
 consistent across all
 experiments and their
 combinations.
- The uncertainty on δ_{CP} remains large.



Global Comparisons – θ_{13}

61

- Daya Bay leads the precision on the measurement of θ_{13} with 2.8% uncertainty.
- Overall, the long-baseline
 measurements are consistent with
 reactor experiments, with larger
 consistency in the normal ordering
 than the inverted ordering.

Normal mass ordering $10.03^{+1.23}_{-1.24}$ 12.3% T2KNOvA + T2K $8.92^{+1.58}_{-1.25}$ 15.9% $8.4 \ ^{+1.9}_{-1.6}$ NOvA 20.8%Dava Bay nGd 8.51 ± 0.24 2.8%RENO nGd 8.92 ± 0.63 7.1%Daya Bay nH $7.1 \pm 1.1 \quad 15.5\%$ RENO nН $8.6 \ \pm 1.2 \ 14.0\%$ Double CHOOZ $10.2 \pm 1.2 \quad 11.8\%$ 10 11 6 $\sin^2 2\theta_{13}, \ 10^{-2}$ Inverted mass ordering T2K $11.09^{+0.55}_{-1.34}$ 8.5%NOvA + T2K $10.08^{+1.62}_{-1.25}$ 14.2% $9.2 \ ^{+2.3}_{-1.4}$ NOvA 20.1%Daya Bay nGd 8.51 ± 0.24 2.8%RENO nGd 8.92±0.63 7.1% Dava Bay nH $7.1 \pm 1.1 \quad 15.5\%$ RENO nH $8.6 \hspace{0.2cm} \pm 1.2 \hspace{0.2cm} 14.0\%$ Double CHOOZ $10.2 \pm 1.2 11.8\%$ 10 8 11 6 9 $\sin^2 2\theta_{13}, 10^{-2}$

Summary & Outlook





Summary

- The joint analysis of NOvA and T2K demonstrates simultaneous compatibility with both datasets.
- The joint analysis shows:
 - Very strong constraint on $|\Delta m_{32}^2|$.
 - Mass Ordering preference remains inconclusive.
 - Small, not significant, preference for Inverted Ordering in the joint fit whereas individual experiments prefer Normal Ordering.
 - $\delta_{CP} = +\pi/2$ lies outside 3-sigma credible interval for both mass orderings.
 - Normal ordering permits a wide range of permissible $\delta_{CP, v}$ while CP conserving values for the Inverted Ordering fall outside the 3-sigma range.
 - Similar conclusions for Jarlskog.





Outlook

- Both experiments continue to collect high quality data and improve their analyses -
 - Data is expected to double, plus updated systematic models, detector response, and new data samples
- Collaboration and information exchange has resulted in a deeper understanding of the analyses
- We are actively exploring the scope and timeline for the next round of this work







Backup



NOvA+T2K+Daya Bay

- Enhanced precision in Δm_{32}^2 presents a "new" lever on measuring neutrino mass-ordering*.
- In the true mass ordering, reactor and long-baseline measurements of Δm_{32}^2 would be consistent but in the incorrect mass ordering would be wrong by different amounts.

Also see: Stephen Parke W&C, 2023

<u>*Phys. Rev. D 72: 013009, 2005</u>

Another possible way to determine

the Neutrino Mass Hierarchy

Hiroshi Nunokawa¹, * Stephen Parke², † and Renata Zukanovich Funchal^{3‡}



NOvA+T2K+DayaBay

- Including the Δm_{32}^2 constraint from the Daya Bay*, reverse the mass ordering preference back to the Normal Ordering.
- Overall, this analysis does not show a significant preference for either mass ordering.

2.47

Inverted/Normal

~71% : ~29% posterior



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Bayes factor

NOvA+T2K Disappearance Data samples:



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T2K

NOVA





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NOVA



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NOVA
Post-fit Correlation Matrix





Priors

Parameter	ARIA sampling prior	MaCh3 sampling Prior	Priors used for the analysis
θ ₂₃	Uniform in $\theta_{_{23}}$	Uniform in $sin^2\theta_{_{23}}$	Uniform in $sin^2\theta_{_{23}}$
θ ₁₃	Uniform in θ_{13}	Uniform in $sin^2 2\theta_{13}$	Uniform in sin²2θ ₁₃ & Gaussian reactor constraint
$ \Delta m^2_{32} $	Uniform in ∆m² ₃₂	Uniform in ∆m² ₃₂	Uniform in $ \Delta m^2_{32} $
MO	Uniform in MO with a 50% switch probability	Uniform in MO with a 50% switch probability	Uniform in MO with a 50% switch probability
δ _{CP}	Uniform in $\boldsymbol{\delta}_{_{\mathrm{CP}}}$	Uniform in $\boldsymbol{\delta}_{_{\mathrm{CP}}}$	Uniform in S_{CP} & Uniform in sin S _{CP} (for J)



Without reactor constraint





No reactor constraint: CP Phase - $\delta_{\rm CP}$





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No reactor constraint Δm^2_{32} and $sin^2\,\theta_{23}$





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Fitter: ARIA No reactor: Bayesian Cred. Int. **NO Conditional** 0.04 \vdash No reactor constraint Posterior density CP Phase - δ_{CP} 1σ 0.03 2σ 3σ 0.02 Fitter: ARIA Bayesian Cred. Int. **Both MO** NOvA-0.04 No reactor constraint 0.01 Posterior density 1σ 0.03 $\underline{0}_{\pi}$ T2K $-\frac{\pi}{2}$ $\frac{\pi}{2}$ 0 2σ δ_{CP} 3σ 0.02 Preliminary Fitter: ARIA Bayesian Cred. Int. **IO Conditional** No reactor constraint 0.01 Posterior density 1σ 0.04 2σ $\underline{0}_{\pi}$ 3σ $\frac{\pi}{2}$ $\frac{\pi}{2}$ π 0 δ_{CP} 0.02 <u>0</u><u></u>π

N

Preliminary

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Pre

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π

 $\frac{\pi}{2}$

0

NOVA

 δ_{CP}

 $-\frac{\pi}{2}$

π

No reactor constraint









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With reactor constraint







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 $\sum_{\text{For both mass orderings, } \delta_{\text{CP}} = \pi/2 \text{ lies outside 3-sigma credible interval.}}$ •

- Normal Ordering allows for a broad range of permissible δ_{CP} •
- For the Inverted Ordering, CP conserving values of $\delta_{CP}(0, \pi)$ lie outside the 3-sigma credible interval. •



Jarlskog



NO Conditional

Flat in $sin(\delta_{CP})$

3σ

0.04

3σ

0.04

2σ

1σ

0.02

1σ

0.02

IO Conditional

Flat in $sin(\delta_{CP})$

2σ

0

Jarlskog

0

Jarlskog

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NOvA-T2K

Preliminary

NOV

 \supset

T2K

Preliminary



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Comparisons





Fitter comparisons All 3 sets of data-fits are consistent with each other.







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Comparison with NOvA-only & T2K-only fits





Comparison with NOvA-only & T2K-only fits





Comparison with NOvA-only & T2K-only fits





Global Comparisons __ A Normal mass ordering IceCube 0.51 ± 0.05 9.8% $0.43 \ _{-0.04}^{+0.20}$ MINOS+ 27.9% $0.559^{+0.023}_{-0.045}$ $0.465^{+0.031}_{-0.011}$ NOvA 6.1% $0.561\substack{+0.021\\-0.039}$ $NOvA + T2K = 0.470^{+0.016}_{-0.008}$ 5.3%------ $\begin{array}{c} 0.45 \ {}^{+0.06}_{-0.03} \ 0.0\% \ {}^{+0.06}_{-0.028} \\ 0.549 \ {}^{+0.019}_{-0.028} \ 4.3\% \ {}^{+0.025}_{-0.053} \\ 0.555 \ {}^{+0.025}_{-0.053} \ 7.0\% \ {}^{+1}_{-1} \end{array}$ SuperK SuperK+T2K $0.485^{+0.020}_{-0.026}$ T2K0.350.400.450.500.550.600.65 $\sin^2 \theta_{23}$





Comparisons with global fitters





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Impact of correlations & alternate models





Asimov Oscillation noints used for testing

	δ_{CP}	$\Delta m_{32}^2 \left[\mathrm{eV}^2 \right]$	$\sin^2 heta_{23}$
Asimov0 (NOvA best-fit like)	2.576	2.41×10^{-3}	0.57
Asimov1 (T2K best-fit like)	-1.60	2.51×10^{-3}	0.53
Asimov4 (NuFit like)	-1.60	-2.45×10^{-3}	0.55

• Other parameters were kept constant at:

$$\Delta m_{21}^2 = (7.53 \pm 0.18) \times 10^{-5} \text{eV}^2$$
$$\sin^2 \theta_{12} = 0.307 \pm 0.013$$
$$\sin^2 \theta_{13} = (2.18 \pm 0.07) \times 10^{-2}$$



Impact of correlations • Δm_{32}^2 bias mock data (nightmare) study











Oscillation	Largest NOvA	Largest T2K	
Parameter	Systematic	Systematic	
δ_{CP}	second class currents	$\sigma_{\nu_e}/\sigma_{\nu_{\mu}}$ cross section	
	and radiative corrections	and antineutrino equivalents	
$\sin^2 heta_{23}$	neutron visible energy	2p2h C-O scaling	
Δm^2_{32}	calibration	7% SK energy-scale [*]	



(a) δ_{CP}

2σ Credible Intervals

0.55

0.5

(c) $\sin^2 \theta_{23}$

0.6

Simulated Valu

Correlation

100% Correlation

0.4

0.45



(b) Δm_{32}^2



NOVA

 Correlating largest systematics on Δm_{32}^2 across both experiments.

IMPERIAL

0.35

posterior probability

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Oscillation	Largest NOvA	Largest T2K
Parameter	Systematic	Systematic
δ_{CP}	second class currents	$\sigma_{\nu_e}/\sigma_{\nu_{\mu}}$ cross section
	and radiative corrections	and antineutrino equivalents
$\sin^2 heta_{23}$	neutron visible energy	2p2h C-O scaling
Δm^2_{32}	calibration	7% SK energy-scale [*]

 Correlating largest systematics on $\sin^2 \theta_{23}$ across both experiments.



NOVA

×10⁻³

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Oscillation	Largest NOvA	Largest T2K	
Parameter	Systematic	Systematic	
δ_{CP}	second class currents	$\sigma_{\nu_e}/\sigma_{\nu_{\mu}}$ cross section	
	and radiative corrections	and antineutrino equivalents	
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Δm^2_{32}	calibration	7% SK energy-scale [*]	

• Correlating largest systematics on $\sin^2 \theta_{23}$ across both experiments.



• Various Placente of alternation at the marges timpact on T2K's 2020-era fit and

- the two cross-experiment model checks were done for the joint analysis:
 - Non-QE: ND280 CC0π data are under-predicted by the T2K pre-fit prediction. This difference can be taken accounted for by the large freedom in the CCQE model. To check this large freedom does not cause bias, an alternate model where this under-prediction is attribution to only non-QE processes is produced.
 - Minerva1Pi: suppression of CC and NC resonant pion production at low-Q² to describe for GENIE v2 implementation of Rein-Seghal model to describe the data.
 - Pion SI: replaced GEANT4 model* was replaced with NEUT's Salcedo–Oset model**

* S. Agostinelli et al., (The GEANT4 collaboration), Nucl. Instrum. Meth. A 506 (2003) 250–303 SLAC-PUB-9350

** L. L. Salcedo, E. Oset, M. J. Vicente-Vacas, and C. Garcia-Recio, Nucl. Phys. A 484 (1988) 557–592 Print-87-1084 (Valencia)



• Test for bias 1D posteriors of Δm_{32}^2 and $\sin^2 \theta_{23}$:

- Change in the width of the 1D intervals <10%
- Change in central value < 50% of systematic uncertainty

Alternate Model	Δm^2_{32} Change in 1D contour < 10%	Δm^2_{32} Bias in central value < 50%	sin² θ₂₃ Change in 1D contour < 10%	sin ² θ ₂₃ Bias in central value < 50%
Non-QE	\checkmark	\checkmark	\checkmark	\checkmark
Minerva1p	\checkmark	\checkmark	\checkmark	\checkmark
Pion-SI	\checkmark	\checkmark	\checkmark	\checkmark
NOvA-like	\checkmark	\checkmark	\checkmark	\checkmark
T2K-like	\checkmark	\checkmark	\checkmark	\checkmark



- Discrete ot od flatte other ties: models
 - Fractional change in Bayes factor for mass ordering and octant should not change any conclusions.
- Additional test on whether alternate models change our conclusion on the significance of CP violation.

Alternate Model	Conclusion on δ_{CP}	Conclusions on J	Mass Ordering Fractional change in BF	Octant Fractional change in BF
Non-QE	\checkmark	\checkmark	1.02	0.88
Minerva1p	\checkmark	\checkmark	1.03	0.92
Pion-SI	\checkmark	\checkmark	0.94	1.11
NOvA-like	\checkmark	\checkmark	1.10	1.00
T2K-like	\checkmark	\checkmark	1.08	1.16



Impact of alternate^{*} models: Minerva1p



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Impact of alterna • Minerva1Pi:

suppression of CC and NC resonant pion production at low-Q² to describe for GENIE v2 implementation of Rein-Seghal model to describe the data.





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 $\mathbf{NO} \mathbf{VA}$
Impact of altern • Pion SI: replaced GEANT4 model* was replaced with NEUT's Salcedo–Oset model**

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 $\mathcal{N}O_{\mathcal{N}}\Lambda$

• Harther Approximation (CRPA)*

- Applies modifications to the nuclear models (Spectral Function for T2K, Local Fermi Gas for NOvA)
- Recent T2K analyses have included an additional ٠ smearing on Δm^2_{32} based on variations seen when considering the HF-CRPA nuclear model.
 - Both NOvA and T2K independently studied the impact ٠ of this alternate nuclear model on their 2020-era analyses.
 - When taken together in the context of the joint fit, the bias is no larger than the thresholds set for any of the fake data metrics.



Impact of cross-experiment model: T2K-like

- The purpose of the cross-experiment test is to verify each analysis is not broken by pseudo-data made using a representative allowed point from the other experiment's model.
- Unlike the FDS, which test a single variation, these tests also include modification of multiple processes.
- Extrapolating the individual processes consistently to the to other experiment's energy at the same time consistently is not possible.



Impact of cross-experiment model: NOvA-like

 NOvA near-to-far extrapolation method does not produce a 'post-fit' cross section model; we choose the NOvA prefit tune as the test model, but this is an arbitrary choice from a large space of valid choices.



CCOπ E_{v,reconstructed}-E_{lepton} in T2K vs NOvA

- In T2K the p_T of the lepton is used to measure the recoiling energy by two body quasielastic kinematics.
- In NOvA, the visible recoil is measured.
- In this T2K-NOvA analysis, we are not relying on a single model to simultaneously describe these variables, but we may in the future
- MINERvA compares the two types of energy measures: recoil in bins of q₀^{QE} (the energy T2K adds to the muon energy)
 - Agreement with this model is poor
 - Events where the QE hypothesis says there should be lots of proton energy added, but MINERvA does not see that energy!
- T2K and NOvA naturally continue to investigate improvements in their cross section models. We appreciate the continued theoretical and experimental effort in the community



Category	NOvA Parameters ZNormCCQE ZExpAxialFFSyst2020_EV1 ZExpAxialFFSyst2020_EV2 ZExpAxialFFSyst2020_EV3 ZExpAxialFFSyst2020_EV4 RPAShapeenh2020 RPAShapesupp2020 MECEnuShape2020Nu MECEnuShape2020AntiNu	 Models and systematics used for 2020 analysis [NOVA: PhysRevD.106.032004, T2K:arXiv:2303.032222v1] will be used in the joint fit. The base-models are tuned to internal (NOVA-ND data by NOVA) and external datasets. The tuning modifies the underlying models drastically (eg: NOvA's 2p2h tune.) 							
MEC	MECShape2020Nu MECShape2020AntiNu MECInitStateNPFrac2020Nu	2p2h Shape O 2p2h Edep low Enu 2p2h Edep high Enu	Experiment	Generator	QE	MEC/2p2h	RES	DIS	FSI
RES	MECInitStateNPFrac2020AntiNu MaCCRES MvCCRES MaNCRES MvNCRES	2p2h Edep low Enubar 2p2h Edep high Enubar CA5 MA RES ISO Bkg Low PPi	NOvA	GENIE v3.0.6	Local Fermi Gas Z-expansion axial form factor	Valencia* (*with NOvA 2020 tune)	Berger- Sehgal	Bodek-Yang	hN Semi Classical Cascade (*fit to pion scattering data)
FSI	LowQ2RESSupp2020 hNFSI_MFP_2020 hNFSI_FateFracEV1_2020	ISO Bkg FEFQE FEFQEH FEFINEL FEFABS FEFCX	T2K	NEUT 5.4	Spectral Function M ^{A^{QE} form factor}	Valencia	Rein- Sehgal	Bodek-Yang	Semi- Classical Cascade



2D constraints from Daya Bay





θ_{13} measurements



NOvA+T2K+Daya Bay



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NOVA

Fitter: ARIA 0.05 DayaBay 2023 • In Netrue Aats Treeking, reatyran Bargebaseline ^Dosterior Density NOvA-T2K w/o reac. 0.04 0.03 measurements of Δm^2_{32} would be consistent but 0.02 the incorrect mass ordering would be wrong by different 0.01 amounts 2.2 2.4 2.6 $|\Delta m_{32}^2| \times 10^{-3} \, eV^2$ ee : reactor disappearance channel \rightarrow Daya Bay* Fitter: ARIA $\mu\mu$: long-baseline disappearance channel \rightarrow NOvA+T2K DayaBay 2023 0.05



NO Conditional

2.8

Baselines

- More sensitivity to mass ordering for higher neutrino energy for the longer baseline.
- Opposite impact of matter effect and δ_{CP} for $v_e vs \overline{v}_e$ appearance probability.

	T2K	NOvA
L (baseline)	295 km	810 km
Energy (beam peak)	0.6 GeV	2 GeV
Matter effect	~ ±9%	~ ±19%
CP effect	~ ±30%	~ ±25%



NOVA

Long-baseline landscape in 2020



- T2K saw an asymmetry in their v_e and v_e appearance while NOvA did not.
- T2K's data favored large CP violation and normal mass ordering while NOvA data lies close to the degenerate δ_{CP} -MO phase space.



- Opposite impact of matter effect and δ_{CP} for v_e vs \overline{v}_e appearance probability.
- Larger matter effect for higher neutrino energy \rightarrow higher sensitivity to mass ordering.
 - Therefore, associated asymmetry is higher for the longer baseline.



