



Exclusive π^0 muoproduction at COMPASS

Theory and Experiment in HEP, 1-4 Oct 2024, Prague

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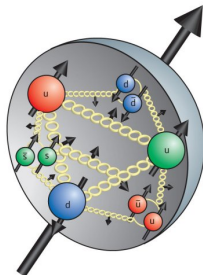


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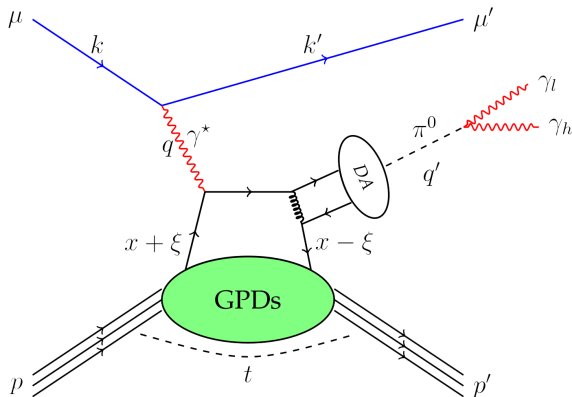
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Generalized Parton Distributions (GPDs)

- Generalization of standard PDFs \Rightarrow longitudinal momentum and transverse position
- Provide a 3D picture of how quarks and gluons build up the nucleons
- Access to GPDs via
 - Deeply Virtual Compton Scattering (DVCS)
 - **Hard Exclusive Meson production (HEMP)**
 - Hard exclusive π^0 muoproduction
 - Cross section measurement of $\mu p \rightarrow \mu' \pi^0 p'$



Hard exclusive π^0 muoproduction



■ Sensitive to the GPDs

- $\tilde{H}(x, \xi, t)$ and $\tilde{E}(x, \xi, t)$ - chiral-even (conserving the parton helicity)
- $H_T(x, \xi, t)$ and $\tilde{E}_T(x, \xi, t)$ - chiral-odd (parton helicity flip)

Reduced cross section for hard exclusive π^0 production

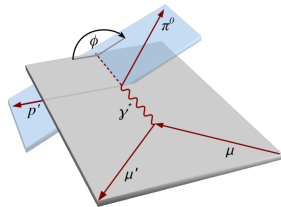
$$\frac{d^2\sigma^{\gamma^*P}}{dt d\phi_{\pi^0}} = \frac{1}{2\pi} \left[\epsilon \frac{d\sigma_L}{dt} + \frac{d\sigma_T}{dt} + \epsilon \cos(2\phi_{\pi^0}) \frac{d\sigma_{TT}}{dt} + \sqrt{2\epsilon(1+\epsilon)} \cos(\phi_{\pi^0}) \frac{d\sigma_{LT}}{dt} \right]$$

$$\frac{d\sigma_L}{dt} \propto \left[(1 - \xi^2) |\langle \tilde{H} \rangle|^2 - 2\xi^2 \text{Re} \left[\langle \tilde{H} \rangle * \langle \tilde{E} \rangle \right] - \frac{t'}{4M^2} \xi^2 |\langle \tilde{E} \rangle|^2 \right]$$

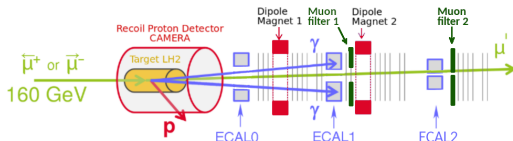
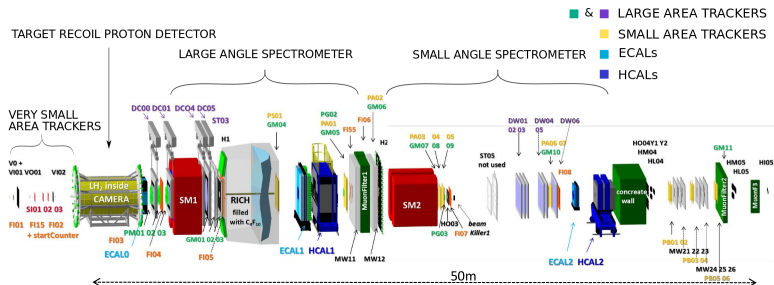
$$\frac{d\sigma_T}{dt} \propto \left[(1 - \xi^2) |\langle H_T \rangle|^2 - \frac{t'}{8M^2} |\langle \bar{E}_T \rangle|^2 \right]$$

$$\frac{d\sigma_{TT}}{dt} \propto t' |\langle \bar{E}_T \rangle|^2$$

$$\frac{d\sigma_{LT}}{dt} \propto \xi \sqrt{1 - \xi^2} \sqrt{-t'} \text{Re} \left[\langle H_T \rangle * \langle \tilde{E} \rangle \right]$$

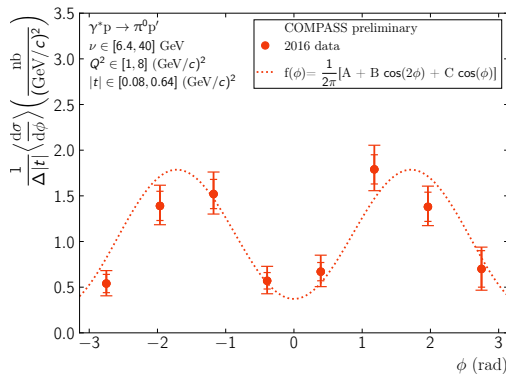
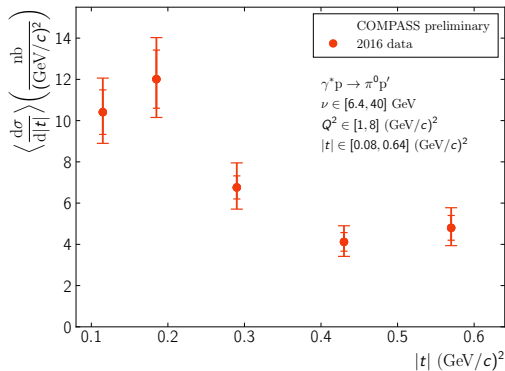


Exclusive π^0 measurement at COMPASS experiment at CERN



■ Measured in 2016

Cross section for $\nu \in (6.4, 40)$ GeV and $Q^2 \in (1, 8)$ GeV²/c², $\langle x_B \rangle = 0.134$



$$\left\langle \frac{d\sigma_T}{dt} + \epsilon \frac{d\sigma_L}{dt} \right\rangle = (6.6 \pm 0.3_{\text{stat}} \pm 0.9_{\text{sys}}) \frac{\text{nb}}{(\text{GeV}/c)^2}$$

$$\left\langle \frac{d\sigma_{TT}}{dt} \right\rangle = (-4.6 \pm 0.5_{\text{stat}} \pm 0.3_{\text{sys}}) \frac{\text{nb}}{(\text{GeV}/c)^2}$$

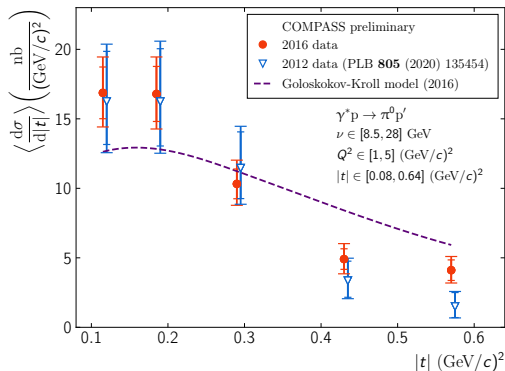
$$\left\langle \frac{d\sigma_{LT}}{dt} \right\rangle = (0.2 \pm 0.2_{\text{stat}} \pm 0.2_{\text{sys}}) \frac{\text{nb}}{(\text{GeV}/c)^2}$$

$$\langle \epsilon \rangle = 0.997$$

$$\frac{1}{2\pi} \left[\frac{d\sigma_T}{dt} + \epsilon \frac{d\sigma_L}{dt} + \epsilon \cos(2\phi_{\pi^0}) \frac{d\sigma_{TT}}{dt} + \sqrt{2\epsilon(1+\epsilon)} \cos(\phi_{\pi^0}) \frac{d\sigma_{LT}}{dt} \right]$$

- $\frac{d\sigma_{TT}}{dt}$ as large as $\frac{d\sigma_T}{dt} + \epsilon \frac{d\sigma_L}{dt} \Rightarrow$ Importance of \vec{E}_T

Cross section for $\nu \in (8.5, 28)$ GeV and $Q^2 \in (1, 5)$ GeV²/c², $\langle x_B \rangle = 0.1$



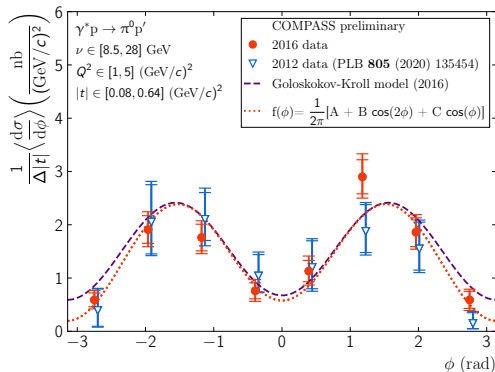
2012 data:

$$\left\langle \frac{d\sigma_T}{dt} + \epsilon \frac{d\sigma_L}{dt} \right\rangle = (8.1 \pm 0.9_{\text{stat}} \pm 1.1 |_{\text{sys}}) \frac{\text{nb}}{(\text{GeV}/c)^2}$$

$$\left\langle \frac{d\sigma_{TT}}{dt} \right\rangle = (-6.0 \pm 1.3_{\text{stat}} \pm 0.7 |_{\text{sys}}) \frac{\text{nb}}{(\text{GeV}/c)^2}$$

$$\left\langle \frac{d\sigma_{LT}}{dt} \right\rangle = (1.4 \pm 0.5_{\text{stat}} \pm 0.3 |_{\text{sys}}) \frac{\text{nb}}{(\text{GeV}/c)^2}$$

$$\langle \epsilon \rangle = 0.996$$



2016 data:

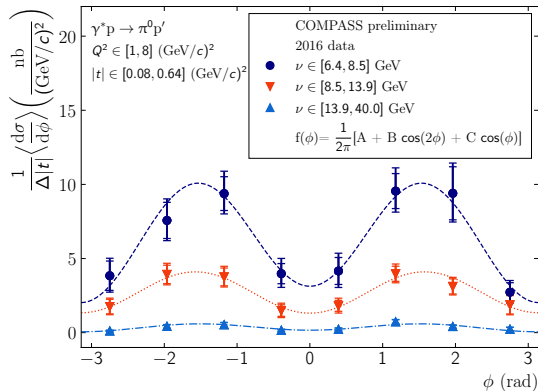
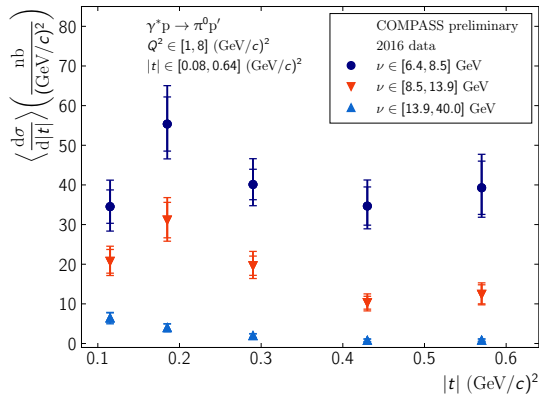
$$\left\langle \frac{d\sigma_T}{dt} + \epsilon \frac{d\sigma_L}{dt} \right\rangle = (8.7 \pm 0.5_{\text{stat}} \pm 1.0 |_{\text{sys}}) \frac{\text{nb}}{(\text{GeV}/c)^2}$$

$$\left\langle \frac{d\sigma_{TT}}{dt} \right\rangle = (-6.3 \pm 0.8_{\text{stat}} \pm 0.4 |_{\text{sys}}) \frac{\text{nb}}{(\text{GeV}/c)^2}$$

$$\left\langle \frac{d\sigma_{LT}}{dt} \right\rangle = (0.6 \pm 0.3_{\text{stat}} \pm 0.3 |_{\text{sys}}) \frac{\text{nb}}{(\text{GeV}/c)^2}$$

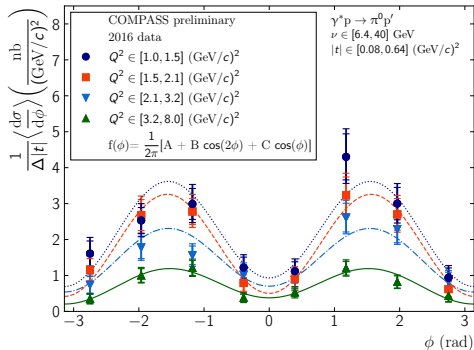
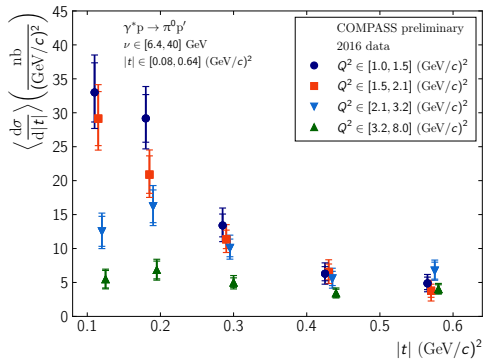
$$\langle \epsilon \rangle = 0.996$$

Evolution of the cross section with ν



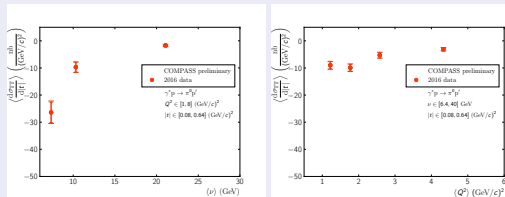
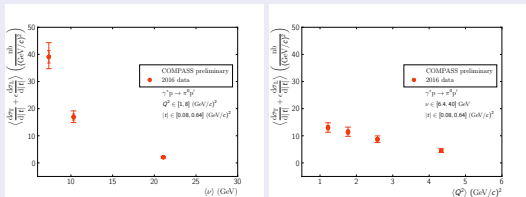
	$\langle \nu \rangle$ [GeV]	$\langle Q^2 \rangle$ [GeV ² /c ²]	$\langle x_B \rangle$	$\langle \epsilon \rangle$
$\nu \in [6.4, 8.5]$	7.35	2.15	0.156	0.999
$\nu \in [8.5, 13.9]$	10.32	2.50	0.131	0.998
$\nu \in [13.9, 40.0]$	21.08	2.09	0.057	0.989

Evolution of the cross section with Q^2



	$\langle Q^2 \rangle \text{ [GeV}^2/c^2]$	$\langle \nu \rangle \text{ [GeV]}$	$\langle x_B \rangle$	$\langle \epsilon \rangle$
$Q^2 \in [1.0, 1.5]$	1.22	10.54	0.072	0.997
$Q^2 \in [1.5, 2.1]$	1.77	9.81	0.109	0.997
$Q^2 \in [2.1, 3.2]$	2.58	9.82	0.157	0.997
$Q^2 \in [3.2, 8.0]$	4.33	10.39	0.247	0.997

Evolution of structure functions with ν and Q^2

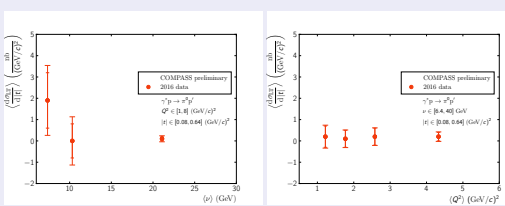


- when ν increases by a factor 3

- $\langle \frac{d\sigma_T}{dt} + \epsilon \frac{d\sigma_L}{dt} \rangle$ decreases by a factor 18
- $|\langle \frac{d\sigma_{TT}}{dt} \rangle|$ decreases by a factor 16

- when Q^2 increases by a factor 4

- $\langle \frac{d\sigma_T}{dt} + \epsilon \frac{d\sigma_L}{dt} \rangle$ decreases by a factor 3 only
- $|\langle \frac{d\sigma_{TT}}{dt} \rangle|$ decreases by a factor 3 only

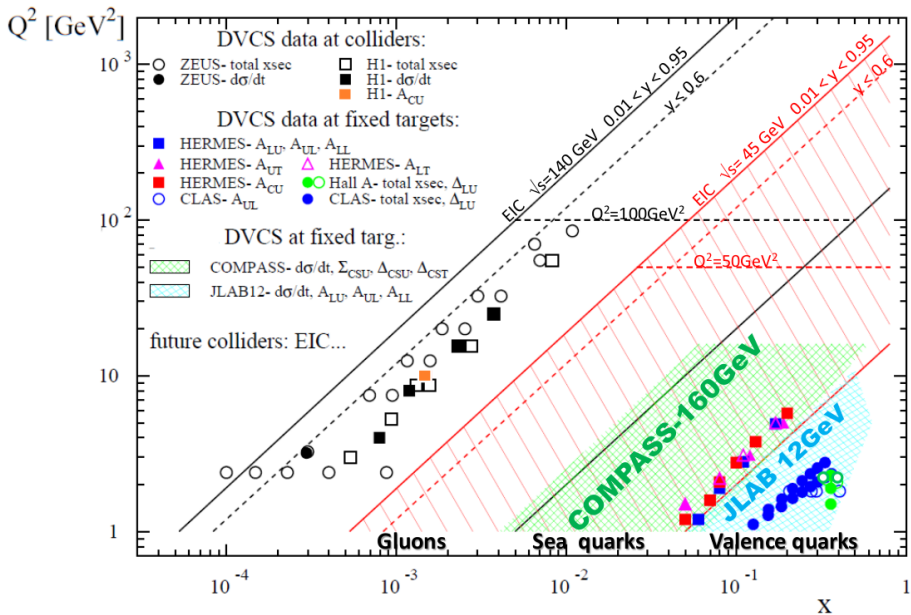


⇒ important inputs for a theoretical description of the major \bar{E}_T contribution

Summary

- Cross section from 2016 COMPASS data calculated in:
 - $\nu \in (6.4, 40)$ GeV and $Q^2 \in (1, 8)$ GeV²/c²
 - $\nu \in (8.5, 28)$ GeV and $Q^2 \in (1, 5)$ GeV²/c²
 - same as used for 2012 data (PLB 805 135454), now ~ 2.3 times larger statistics
 - Study evolution in
 - 3 bins in ν
 - 4 bins in Q^2
- Prospects:
 - New theoretical model from Peter Kroll and Kornelia Passek-Kumerički, Sergei Goloskokov, Simonetta Liuti...
 - Paper with results from 2016 data
 - 2017 data - to be analyzed

BACKUP

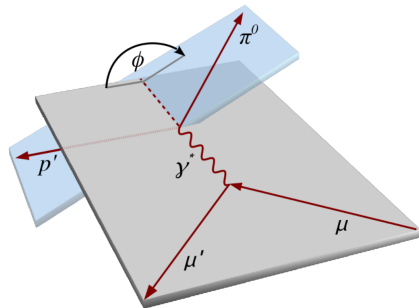


Cross section determination

- $6.4 < \nu < 40 \text{ GeV}$
- $1 < Q^2 < 8 \text{ (GeV/c)}^2$
- $0.08 < |t| < 0.64 \text{ (GeV/c)}^2$
- $-\pi < \phi < \pi \text{ rad}$

$$\frac{d^2\sigma^{\gamma^*p \rightarrow \pi^0 p'}}{dtd\phi_{\pi^0}} = \frac{1}{\Gamma(Q^2, \nu)} \frac{d^4\sigma^{\mu p \rightarrow \mu' \pi^0 p'}}{dQ^2 d\nu dtd\phi_{\pi^0}}$$

- $\Gamma(Q^2, \nu)$ - virtual photon flux



Cross section determination

$$\left\langle \frac{d^2\sigma^{\gamma^*p}}{dtd\phi} \right\rangle_{\Delta\Omega}^{\pm} = \left(\sum_{i=1}^{N_{data}^{\Delta\Omega}^{\pm}} \frac{1}{\Gamma(\Omega_i)a(\Omega_i)\mathcal{L}^{\pm}\Delta\Omega} - f^{\pm} \sum_{i=1}^{N_{LEPTO}^{\Delta\Omega}^{\pm}} \frac{1}{\Gamma(\Omega_i)a(\Omega_i)\mathcal{L}^{\pm}\Delta\Omega} \right)$$

- phase space element $\Delta\Omega = \Delta|t|\Delta\Phi\Delta Q^2\Delta\nu$
- acceptance $a(\Omega_i)$
- luminosity \mathcal{L}
- background normalization f^{\pm}

- spin-independent cross section:

$$\left\langle \frac{d^2\sigma^{\gamma^*p}}{dtd\phi} \right\rangle_{\Delta\Omega} = \frac{1}{2} \left(\left\langle \frac{d^2\sigma^{\gamma^*p}}{dtd\phi} \right\rangle_{\Delta\Omega}^{+} + \left\langle \frac{d^2\sigma^{\gamma^*p}}{dtd\phi} \right\rangle_{\Delta\Omega}^{-} \right)$$

virtual photon flux:

$$\Gamma(Q^2, \nu) = \frac{\alpha_{em}(1 - x_{Bj})}{2\pi Q^2 y E_\mu} \left[y^2 \left(1 - \frac{2m_\mu^2}{Q^2} \right) + \frac{2}{1 + Q^2/\nu^2} \left(1 - y - \frac{Q^2}{4E_\mu^2} \right) \right]$$

virtual photon polarization parameter:

$$\epsilon = \frac{1 - y - \frac{1}{4}y^2 Q^2/\nu^2}{1 - y + \frac{1}{2}y^2 + \frac{1}{4}y^2 Q^2/\nu^2}$$

reduced cross section (full version):

$$\frac{d^2\sigma^{\gamma^*p}}{dt d\phi_{\pi^0}} = \frac{1}{2\pi} \left[\frac{d\sigma_T}{dt} + \epsilon \frac{d\sigma_L}{dt} + \epsilon \cos(2\phi_{\pi^0}) \frac{d\sigma_{TT}}{dt} + \sqrt{2\epsilon(1+\epsilon)} \cos(\phi_{\pi^0}) \frac{d\sigma_{LT}}{dt} \mp |P_l| \sqrt{2\epsilon(1-\epsilon)} \sin(\phi_{\pi^0}) \frac{d\sigma'_{LT}}{dt} \right]$$