



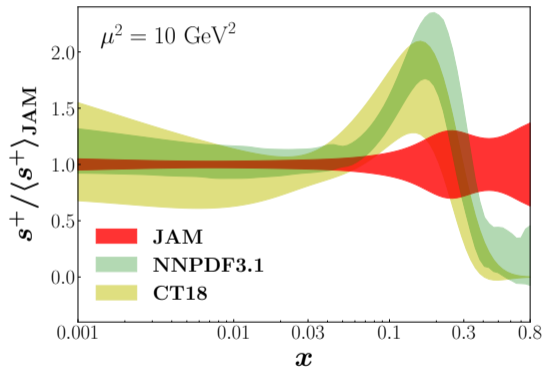
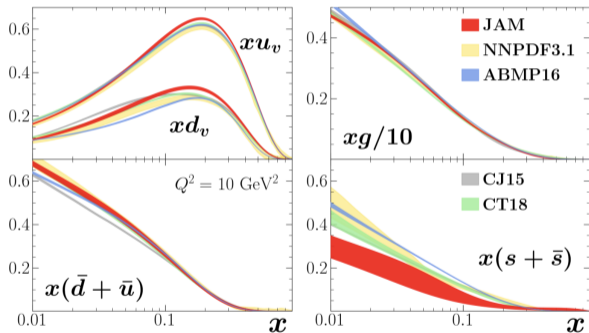
Impact of parity-violating DIS on the nucleon strangeness and weak mixing angle

Richard Whitehill

Collaborators: M. M. Dalton, T. Liu, W. Melnitchouk, J. Qiu, and N. Sato

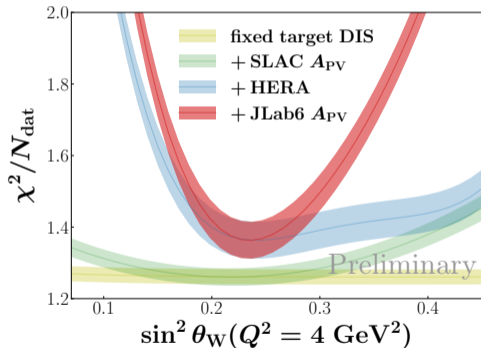
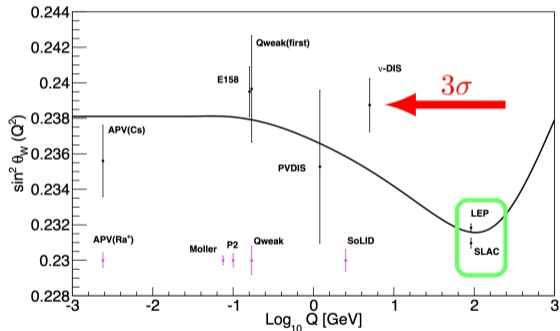


Current Status – PDFs

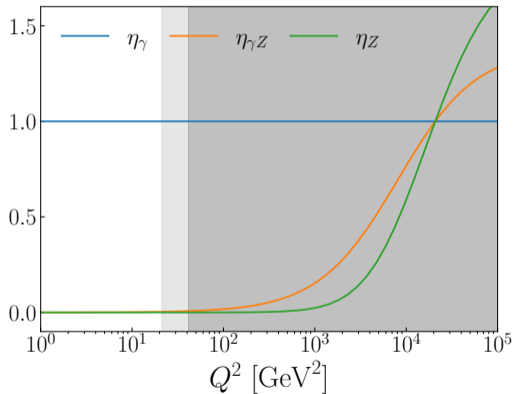
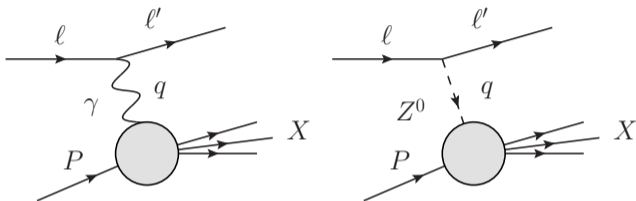


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Current Status – $\sin^2 \theta_W$



Parity-Violating Deep-Inelastic Scattering (PVDIS)



$$\frac{d\sigma_{\lambda_e}}{dx_B dy} = \frac{2\pi\alpha^2 y}{Q^4} \sum_i \eta_i C_i L_{\mu\nu}^\gamma W_{i,U}^{\mu\nu}$$

Parity-Violating Asymmetry

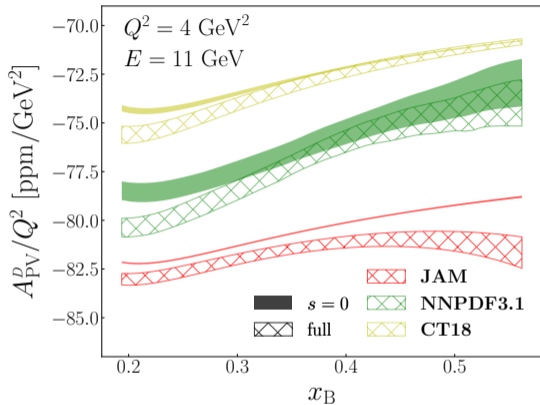
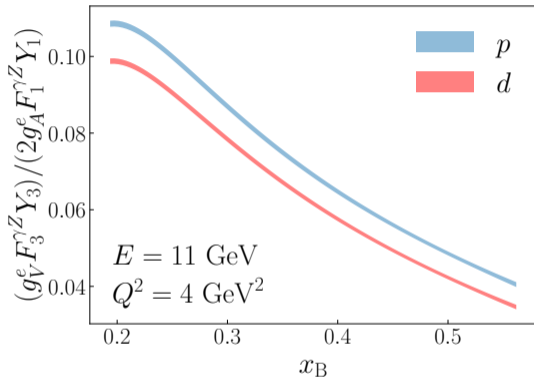
$$A_{\text{PV}} = \frac{d\sigma_+ - d\sigma_-}{d\sigma_+ + d\sigma_-} \approx \frac{G_F Q^2}{4\sqrt{2}\pi\alpha} \left[2g_A^e \frac{F_1^{\gamma Z}}{F_1^\gamma} Y_1 + g_V^e \frac{F_3^{\gamma Z}}{F_1^\gamma} Y_3 \right]$$

$$g_A^e = -1/2, \quad g_V^e = -1/2 + 2 \sin^2 \theta_W$$

$$Y_1 = \left(\frac{1 + R^{\gamma Z}}{1 + R^\gamma} \right) \frac{1 + (1 - y)^2 - \frac{y^2}{2} \left[1 + r^2 - \frac{2r^2}{1 + R^{\gamma Z}} \right]}{1 + (1 - y)^2 - \frac{y^2}{2} \left[1 + r^2 - \frac{2r^2}{1 + R^\gamma} \right]}, \quad r^2 = 1 + 4M^2 x_B^2 / Q^2$$

$$Y_3 = \left(\frac{1 + R^{\gamma Z}}{1 + R^\gamma} \right) \frac{1 - (1 - y)^2}{1 + (1 - y)^2 - \frac{y^2}{2} \left[1 + r^2 - \frac{2r^2}{1 + R^\gamma} \right]}, \quad R^i = \frac{F_2^i}{2x_B F_1^i} r^2 - 1$$

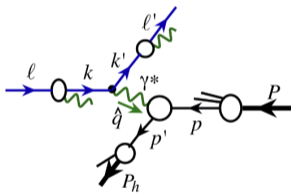
A_{PV} on a deuterium target



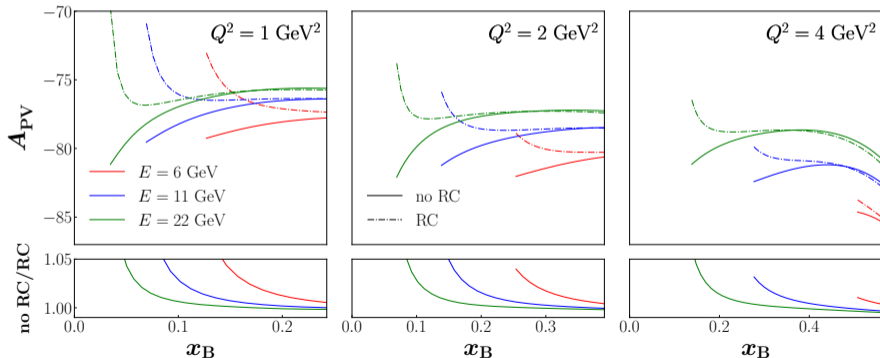
$$A_{PV}^D \approx -\frac{G_F Q^2}{4\sqrt{2}\pi\alpha} \left[\left(\frac{9}{5} - 4 \sin^2 \theta_W \right) + \frac{2}{25} \frac{s^+}{u^+ + d^+} \right]$$

QED radiative effects

$$\frac{d\sigma}{dx_B dy} = \int_{\zeta_{\min}}^1 \frac{d\zeta}{\zeta^2} \underbrace{D_{e/e}(\zeta, \mu^2)}_{\text{LFF}} \int_{\xi_{\min}}^1 d\xi \underbrace{f_{e/e}(\xi, \mu^2)}_{\text{LDF}} \left[\frac{Q^2 \hat{x}_B}{x_B \hat{Q}^2} \right] \frac{d\hat{\sigma}}{d\hat{x}_B d\hat{y}}$$



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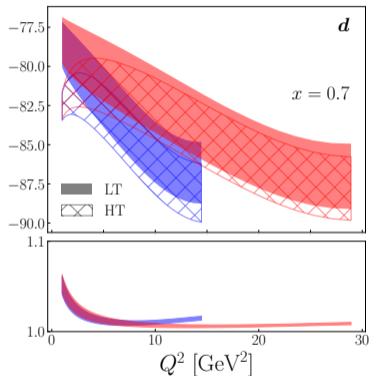
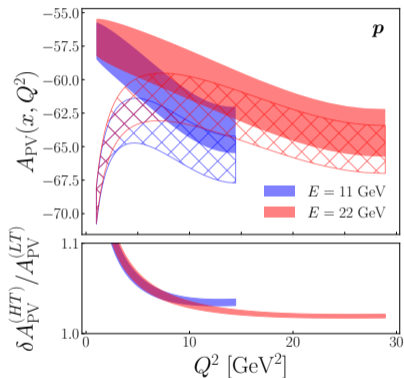
Higher Twist Corrections

$$F_i^\gamma = F_{i,LT}^\gamma \left(1 + \frac{H_i^\gamma}{Q^2} \right)$$

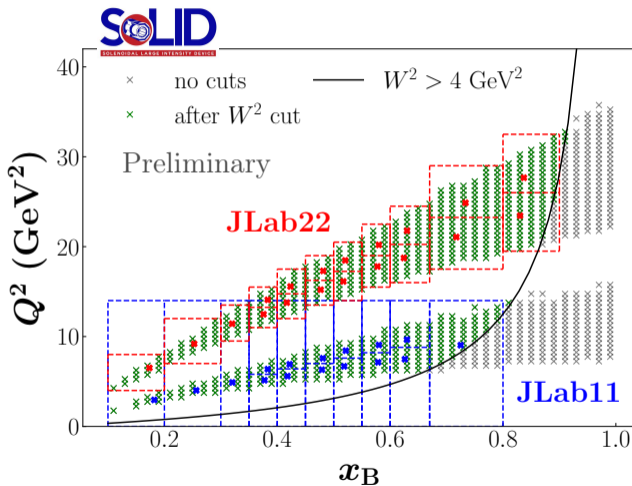
$$F_i^{\gamma Z} = F_{i,LT}^{\gamma Z} \left(1 + \frac{H_i^{\gamma Z}}{Q^2} \right)$$

Model:

$$\rightarrow F_{i,LT}^{\gamma Z} H_i^{\gamma Z} = R F_{2,LT}^\gamma H_i^\gamma$$



Simulating pseudo-data



Scenarios:

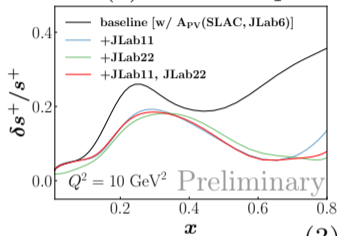
1. statistical uncertainties + experimental systematics
2. (1) + QED effects
3. (2) + HT effects

Note:

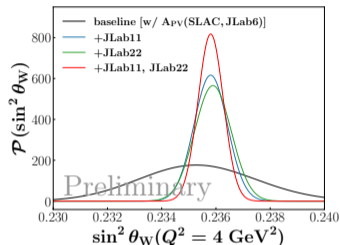
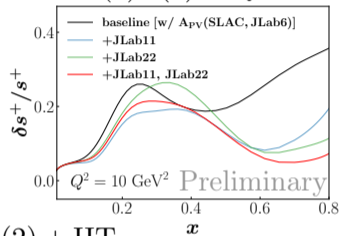
- $P = 85\%$
- $d\mathcal{L}/dt = 4.85 \times 10^{38} \text{ cm}^{-2} \text{ s}^{-1}$
- run time: 50 days/target
- $\delta^{\text{syst}} A_{\text{PV}} = 0.5\%$

Separate impact on PDFs and $\sin^2 \theta_W$

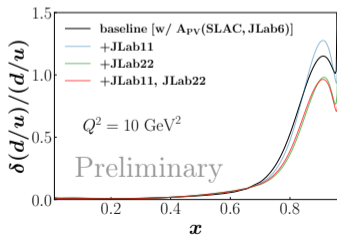
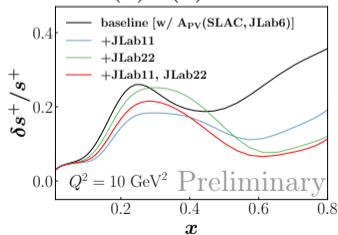
(1): stat + exp



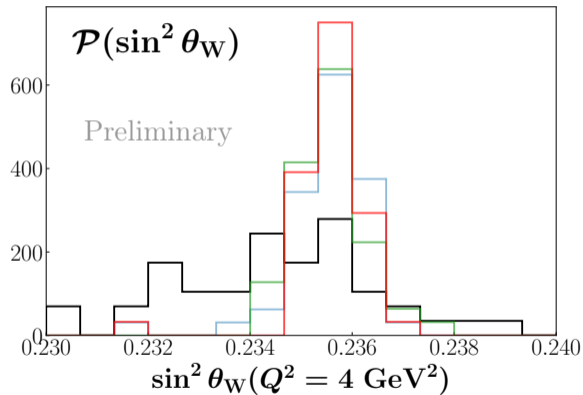
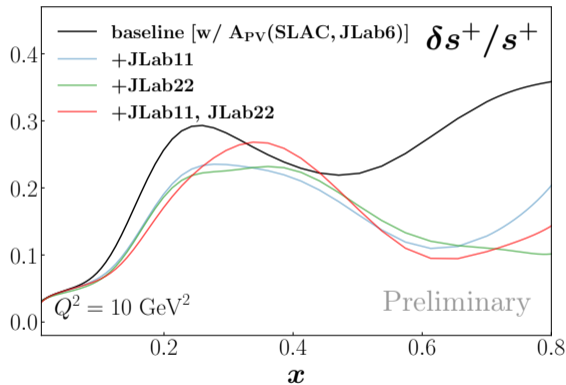
(2): (1) + QED



(3): (2) + HT



Combined impact on PDFs and $\sin^2 \theta_W$

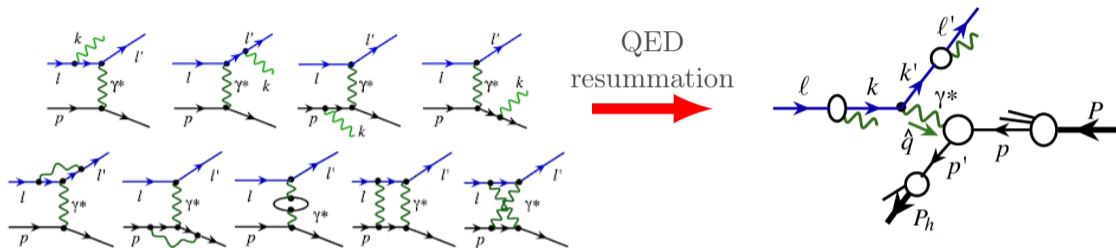


Summary and Outlook

- A_{PV} is a unique and clean observable that can be used in future global analyses to make progress toward
 - constraint of nucleon strangeness for better understanding of nucleon structure
 - tests of BSM physics through the determination of the weak mixing angle
- Future work:
 - electron/positron PVDIS for constraint of sea quark asymmetries
 - Charge symmetry violation
 - Polarized A_{PV} ?

Backup Slides

Hybrid QED+QCD factorization

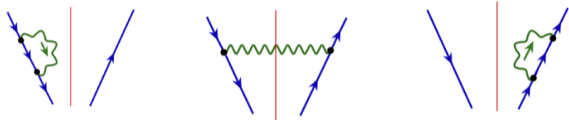


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$$f_{i/e}(\xi) = \int \frac{dz^-}{4\pi} e^{i\xi\ell^+z^-} \langle e | \bar{\psi}_i(0) \gamma^+ \Phi_{[0,z^-]} \psi_i(z^-) | e \rangle$$

$$D_{e/j}(\zeta) = \frac{\zeta}{2} \sum_X \int \frac{dz^-}{4\pi} e^{i\ell^+z^-/\zeta} \text{Tr} \left[\gamma^+ \langle 0 | \bar{\psi}_j(0) \Phi_{[0,\infty]} | e, X \rangle \langle e, X | \psi_j(z^-) \Phi_{[z^-,\infty]} | 0 \rangle \right]$$

LDF and LFF RGE



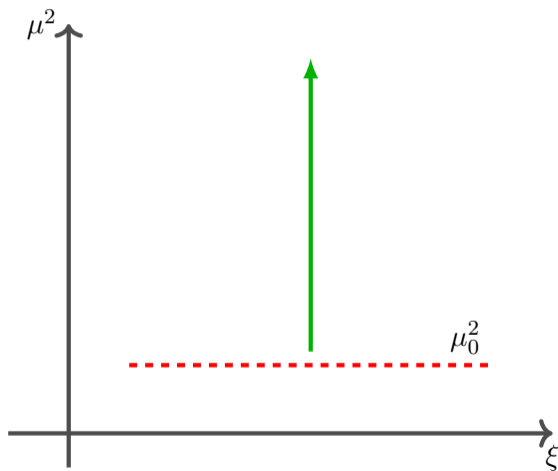
$$f_{e/e}^{(0)}(\xi, \mu_0^2) = \delta(\xi - 1)$$

$$f_{e/e}^{(1)}(\xi, \mu_0^2) = \frac{\alpha}{2\pi} \left[\frac{1 + \xi^2}{1 - \xi} \ln \frac{\mu_0^2}{(1 - \xi)^2 m_e^2} \right]_+$$

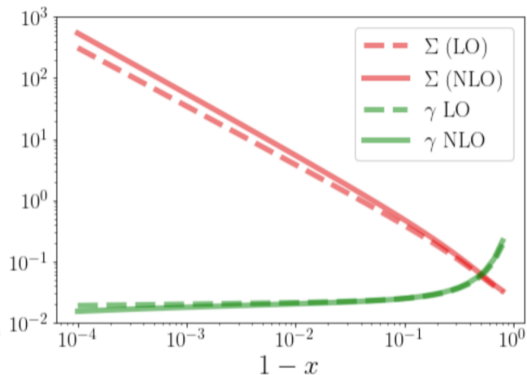
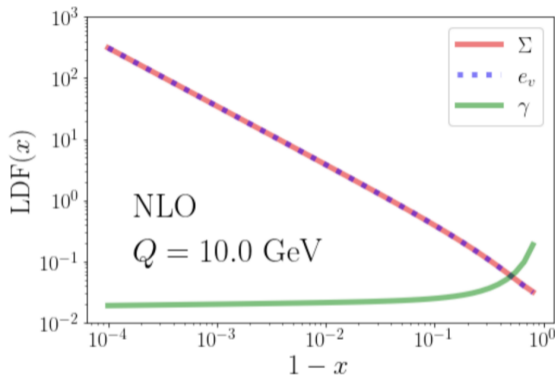
$$D_{e/e}^{(0)}(\zeta, \mu_0^2) = \delta(\zeta - 1)$$

$$D_{e/e}^{(1)}(\zeta, \mu_0^2) = \frac{\alpha}{2\pi} \left[\frac{1 + \zeta^2}{1 - \zeta} \ln \frac{\mu_0^2}{(1 - \zeta)^2 m_e^2} \right]_+$$

$$\frac{\partial f_{i/\ell}(\xi, \mu^2)}{\partial \ln \mu^2} = \sum_j \int_{\xi}^1 \frac{dz}{z} P_{ij}(z) f_{j/\ell}\left(\frac{z}{\xi}, \mu^2\right)$$



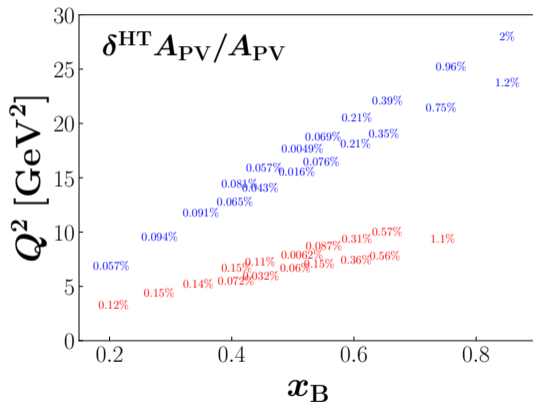
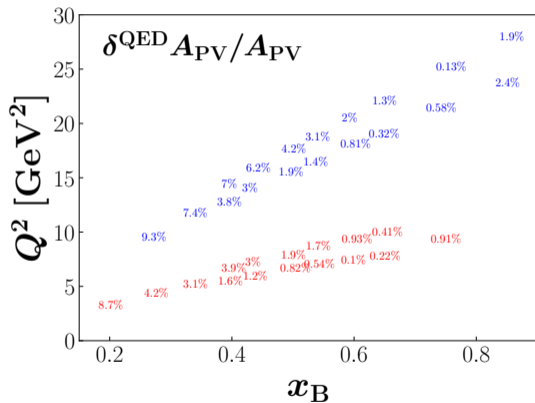
LDFs and LFFs



Subtraction “trick”:

$$\rightarrow \int_{\xi_{\min}}^1 d\xi \mathcal{H}(\xi) f(\xi) = \int_{\xi_{\min}}^1 d\xi [\mathcal{H}(\xi) - \mathcal{H}(1)] f(\xi) + \mathcal{H}(1) \frac{\xi_{\min}}{2\pi i} \int dN \xi_{\min}^{-N} \frac{F_N}{N-1}$$

Size of QED and HT systematics



Impact on s^+ with proton data

$$\mathcal{F} = 5F_{2,p}^{\gamma Z} - 4F_{2,D}^{\gamma} \sim xs^+$$

