

QCD axion dark matter from parametric resonance

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CLUSTER OF EXCELLENCE
QUANTUM UNIVERSE



Initial Field Displacement

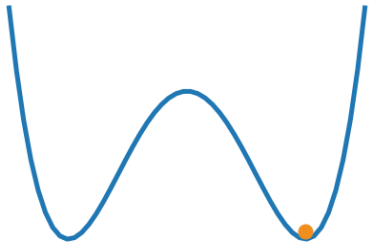
$$\Phi = \frac{\phi}{\sqrt{2}} e^{i\theta}$$

$$S = \int d^4x \sqrt{-g} (\partial_\mu \Phi \partial^\mu \Phi^* + cH^2 |\Phi|^2 - V_{PQ}(|\Phi|))$$

Peccei-Quinn Field

Hubble-induced
mass

*Peccei-Quinn Symmetry
Breaking Potential*



During Inflation the Hubble-induced mass provides a minimum at high field values.

Initial Field Displacement

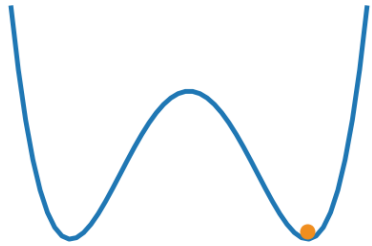
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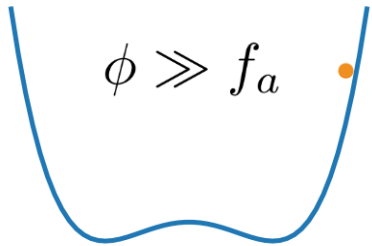
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At the end of Inflation the Peccei-Quinn potential is restored, but the field is frozen at **large VEVs**, because of Hubble friction.

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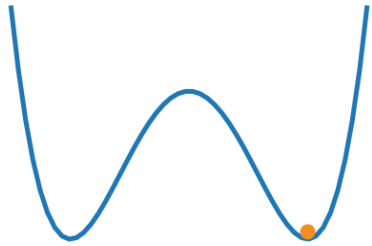
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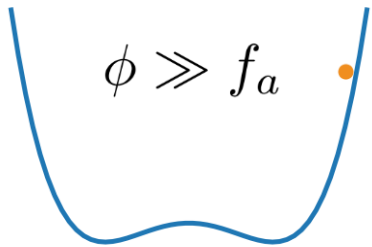
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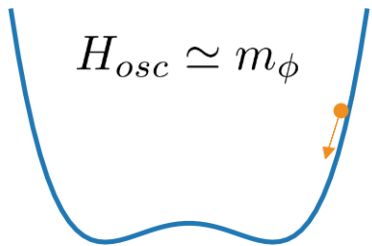
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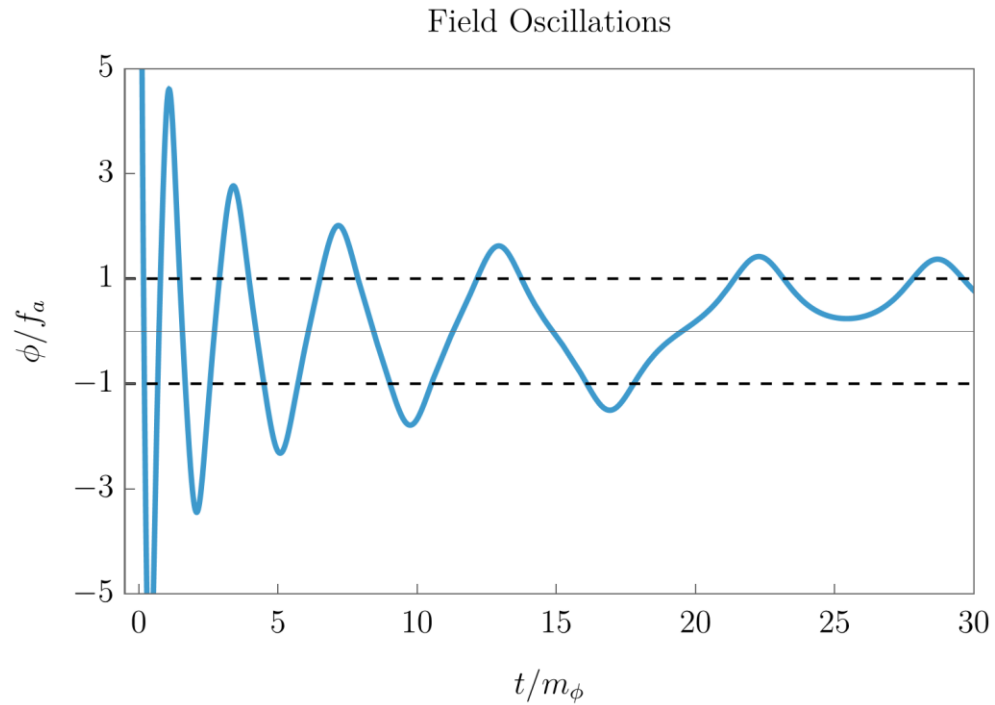


At the end of Inflation the Peccei-Quinn potential is restored, but the field is frozen at **large VEVs**, because of Hubble friction.



When the Hubble friction is weak enough, around $H_{osc} \simeq m_\phi$, the field **starts oscillating**

Parametric Resonance



We can decompose the Peccei-Quinn field in a **homogenous component** + **spatial perturbations**

$$\phi(t, x) = \bar{\phi}(t) + \delta\phi(t, x)$$

Letting the homogeneous field oscillating and treating it as a background field, we can now study the dynamics of **perturbations**

Decomposing perturbations into Fourier modes we notice the arising of a **Hill differential equation**

$$\delta\ddot{\phi}_k + [k^2 + F(t)]\delta\phi_k = 0$$

$$F(t + T) = F(t) \text{ periodic function}$$

This equation admits **exponentially growing solutions** for those modes whose wave numbers resonate with the oscillation frequency of the background field

Parametric Resonance (Quartic Potential Example)

Let's take the case of a **quartic potential** and decompose the field into its **real** and **imaginary** components

$$\Phi = X + iY \quad \frac{\lambda}{4} (|\Phi|^2 - f_a^2)^2$$

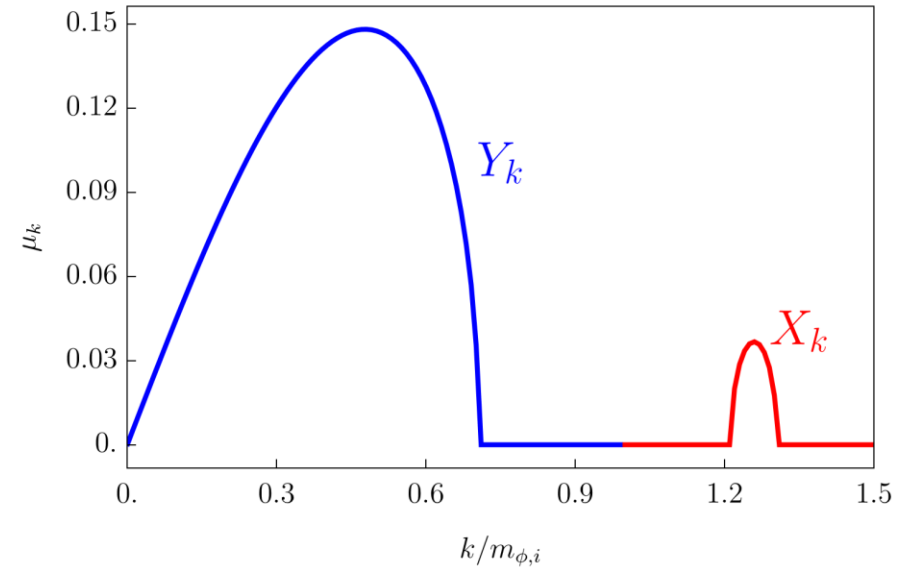
Perturbations follow a particular case of Hill equation, called Mathieu equation:

$$\ddot{y}_k + [k^2 + \lambda \phi_i^2 \cos(\omega_\phi \tau)] y_k = 0$$

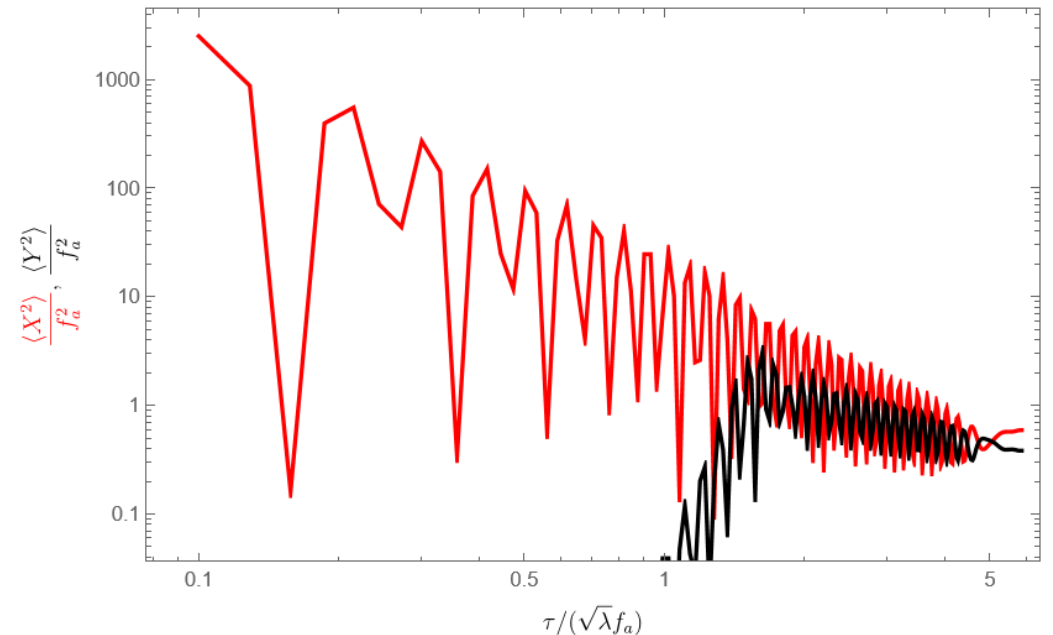
This equation has **bands of instability** in which **solutions grow exponentially**

$$y_k \sim \exp(\mu_k \tau)$$

Floquet Exponents Quartic Potential



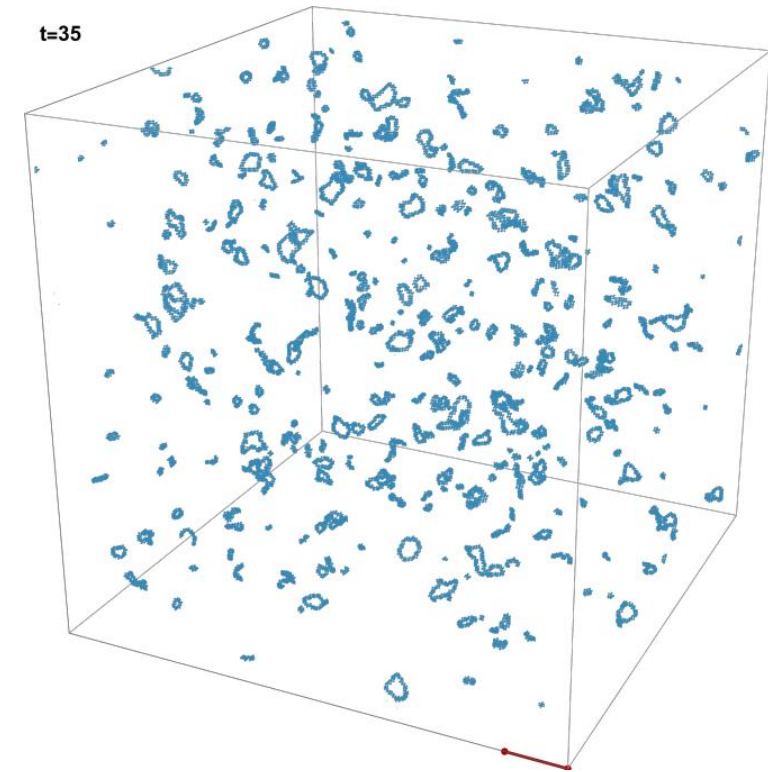
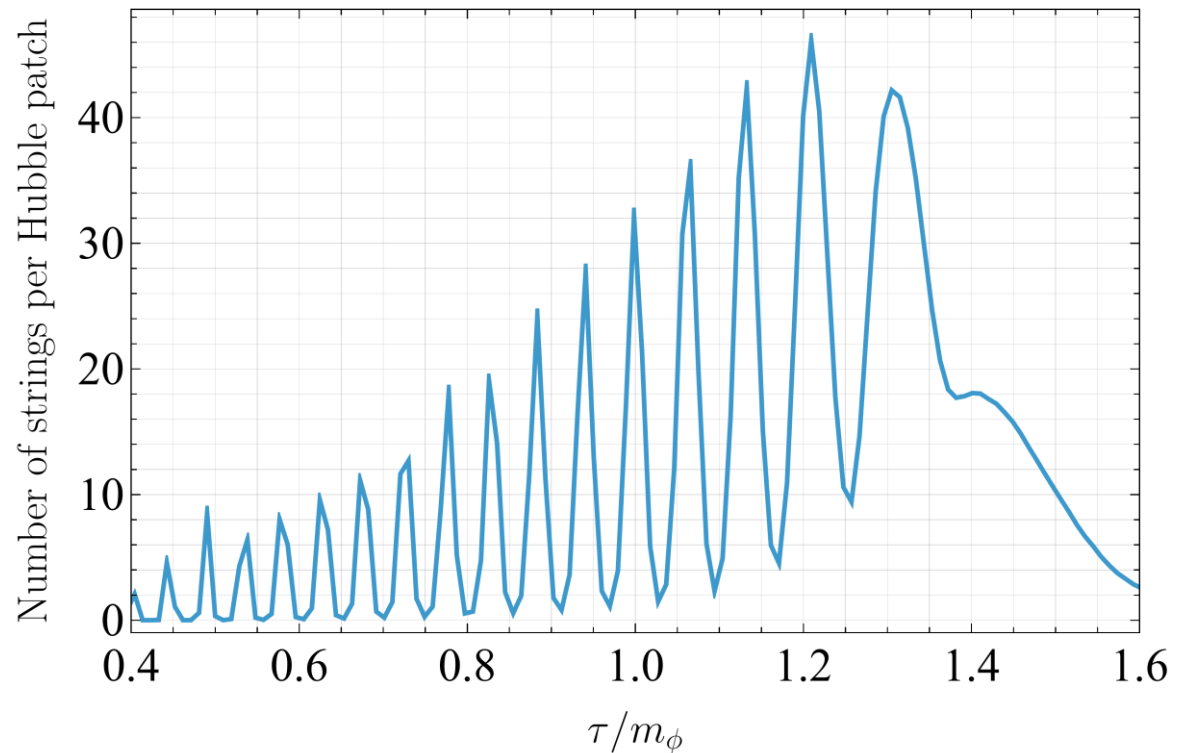
Dynamics of X and Y



Axion Strings Formation

The growth of perturbations leads soon to a non-linear dynamics which has to be tackled via **lattice simulations**.

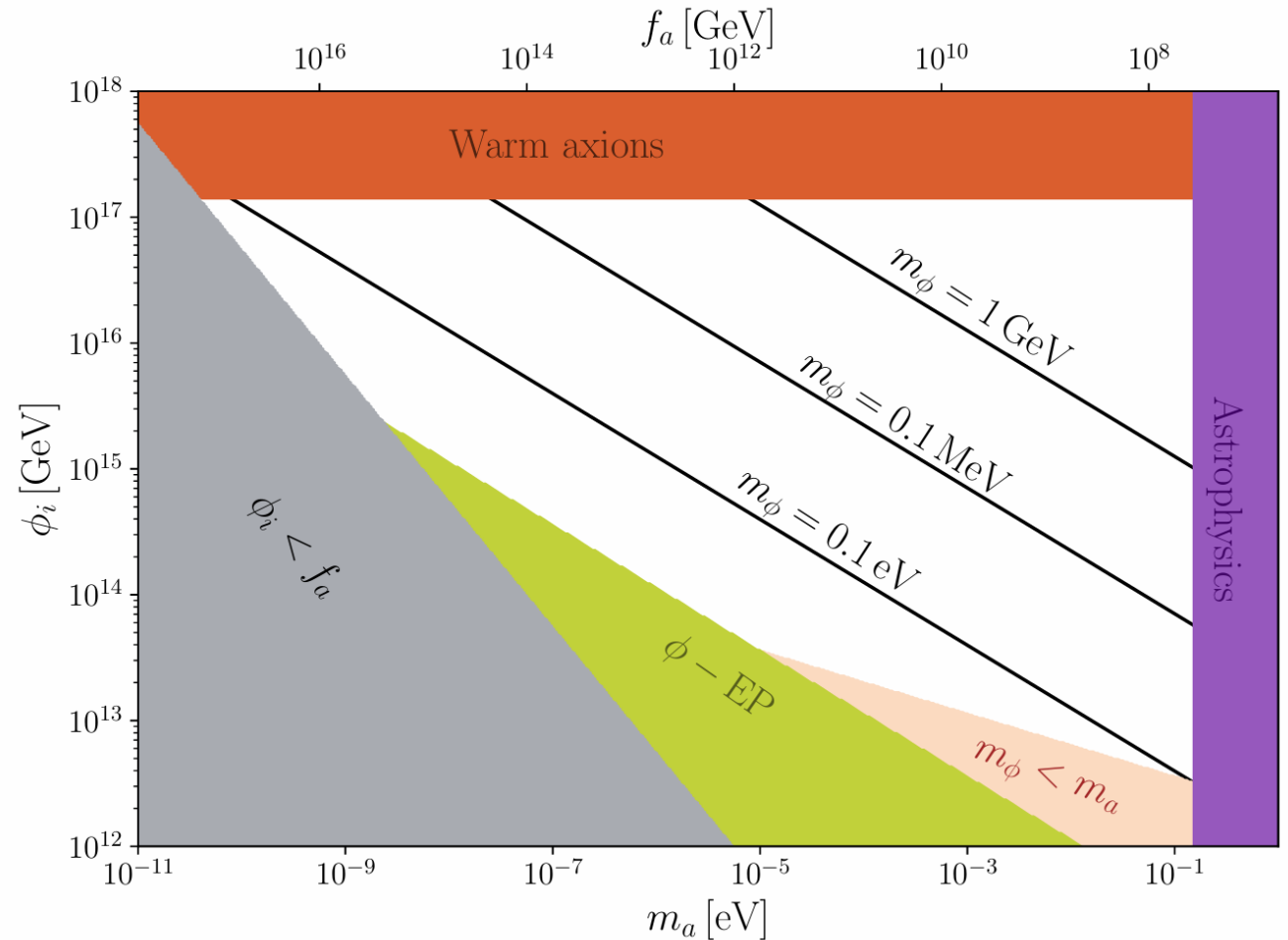
What we observe is the **periodic formation of axion string loops**, with a period fixed by the radial mode oscillation:



Axion Dark Matter

In conclusion, the relic axion abundance, produced by parametric resonance can account to cold dark matter via the relation:

$$h^2\Omega_{axion} \simeq 0.12 \left(\frac{\phi_i}{10^{-2} M_{Pl}} \right)^2 \left(\frac{m_a}{10^{-3} eV} \right) \left(\frac{10^2 GeV}{m_{\phi_i}} \right)^{1/2}$$

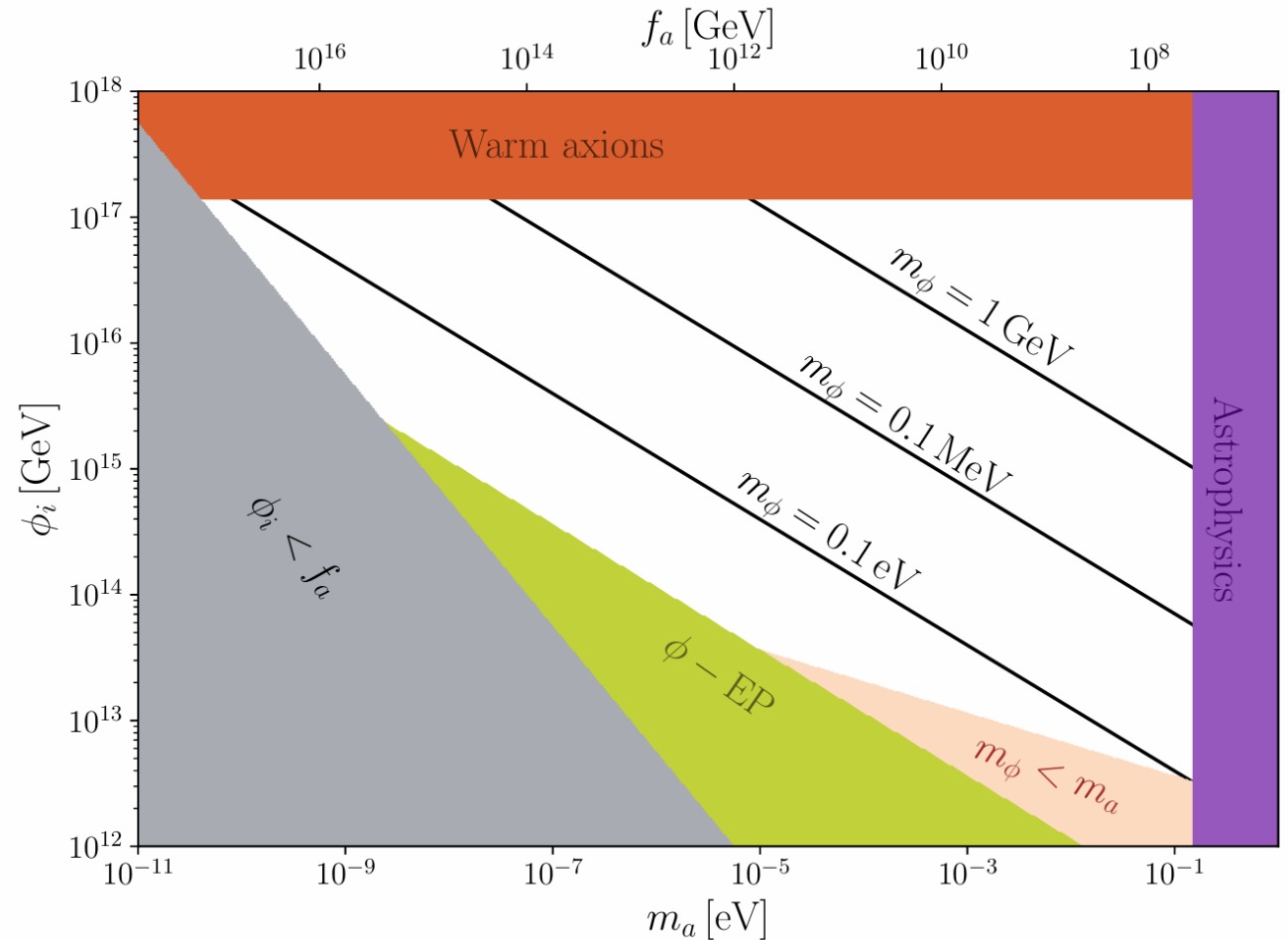


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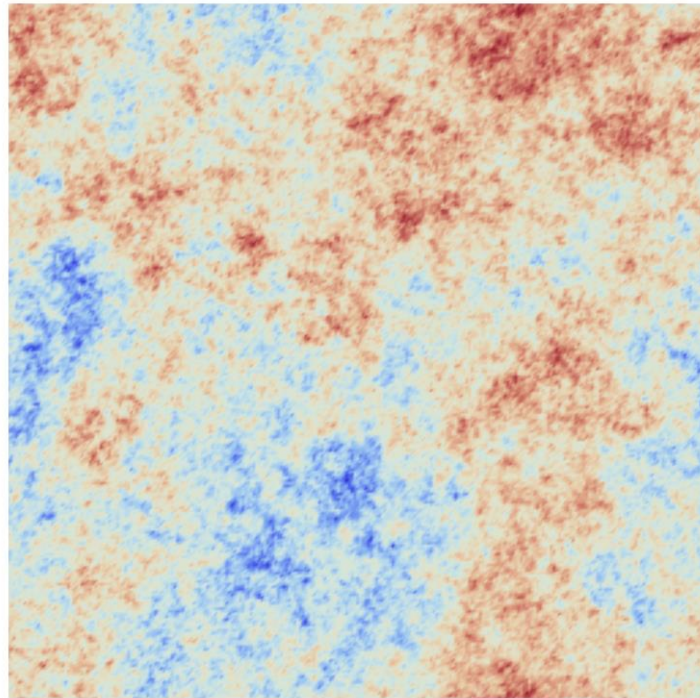
Thank you!



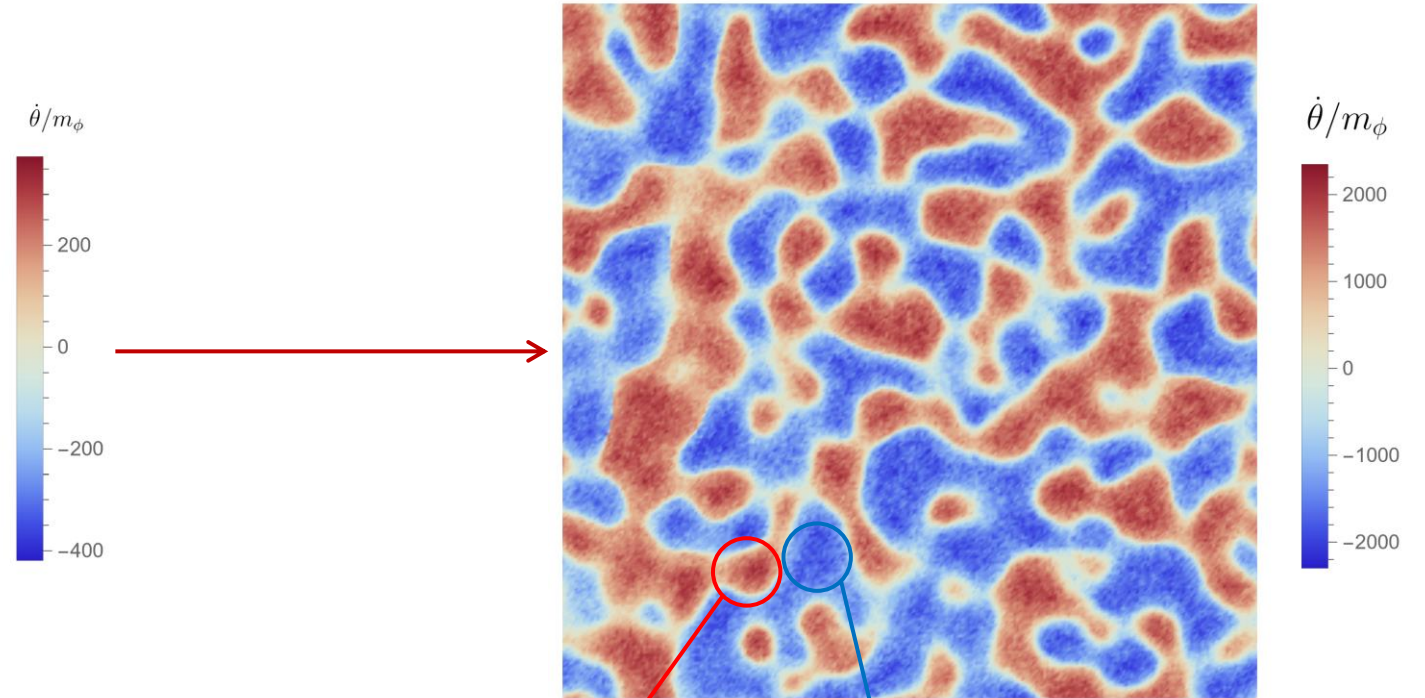
Backup slides

Axion Strings Formation

What we observe is the formation of many well-defined **domains**, in which the axion field is either rotating **clockwise** or **counterclockwise**.



Initial randomized axion field



Counterclockwise
rotation

Clockwise rotation