# High energy photons from the FCC-ee complex

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Other Science Opportunities at the FCC-ee 28 November 2024









#### Outline

- FCC-ee booster: parameters used for operating it as a photon source
- Diffraction limited storage rings
  - Small emittance in high brilliance and high coherent flux
- Average and peak brilliance: FCC-ee booster versus PETRA IV and EuXFEL
- FCC-ee booster for gamma rays for nuclear physics
- FCC-ee linac operating as an XFEL
- Conclusions
- Outlook

#### Parameters used for study of FCC-ee booster as photon source\*

FCC-ee booster	
FCC-ee booster	beam of avg. bo Avg. no rms bu rms rel [10 <sup>-3</sup> ] beta at [m] wiggler wiggler magne tot. len
Technical site PH Collider RF SSS = 1400 m <sup>1</sup> FG (Experiment site) Collimation	hor. er vert. e Time f cycle
European XFEL circumference = 90.7	km

	3xU <sub>0</sub>	94xU <sub>0</sub>	Uo
beam energy [GeV]	20	20	45.6
avg. beam current [mA]	6	6	15
Avg. number of bunches	500	500	1120
rms bunch length [mm]	4	9.5	4.4
rms relative energy spread [10 <sup>-3</sup> ]	0.4	2.2	0.4
beta at wiggler /undulator [m]	1.6	1.6	1.6
wiggler field [T]	1	1	1
wiggler period [mm]	40	40	40
magnetic gap [mm]	10	10	10
tot. length wiggler [m]	6.4	264	5
hor. emittance [pm rad]	15	0.5	100
vert. emittance [pm rad]	1.5	0.05	0.2
Time for users (s) over a cycle of 4 s	2.5	2.5	0.4
without wigglors		*Courtesy of F	Zimmermann

without wiggiers

 $U_0$  = energy loss / turn = 1.33 MeV

hor. em. = 46 pm rad; vert. em. < 5 pm rad

4



## **Filling pattern**



Full-energy booster design, A. Chance et al. 2024

#### FCC-ee booster operated as photon source

Fixed-field chicane: the beam automatically moves out of the wiggler during acceleration



#### Horizontal emittance versus circumference

Storage ring photon sources upgrades to decrease the horizontal emittance FCC-ee booster small emittance as a result of large circumference + damping wigglers/undulators 10<sup>-5</sup> BESSY-II SESAME CANDLE SPEAR3 10<sup>-6</sup> tau-charm DLS  $10^{-7}$ ANKA II SPAI BA SuperKEKB ESRF ELETTRA USSREINUS APS SPRING8  $\epsilon_h$  /  $\gamma^2$  (nm) 10<sup>-8</sup> PETRA-III LHC SOLEIL NSLS-II MAX-IV 0 SuperB SPRING8-II ALS U ILSF ILC-DR LEP-1 DLS-IL SLS-II 10<sup>-9</sup> BAPS CLIC-DR PEP-X PETRA IV LEP-2 ESRF-II **FCC-ee booster** APS-U **10**<sup>-10</sup> MAX-IV - DR PEP-X DR Uo x 3 FCC-ee FCC-hh 10<sup>-11</sup> TevUSR CEPC 54 km **TevUSR DR** Uo x 94 10<sup>-12</sup> **10<sup>-13</sup>** 1000 10000 100 100000 **European XFEL** 

circumference (m)

Based on plot by R. Bartolini

A ring with horizontal emittance  $\varepsilon_x$  ( $\varepsilon_y < \varepsilon_x$ ) is diffraction limited at all photon wavelengths  $\lambda$  emitted by undulators for which



diffraction limited at photon wavelengths >100 Å (< 120 eV)  $\Rightarrow \varepsilon_x \approx 0.8$  nm rad (ESRF, APS, PETRA III) >10 Å (< 1.2 keV)  $\Rightarrow \varepsilon_x \approx 80$  pm rad (APS-U, EBS) >1 Å (< 12 keV)  $\Rightarrow \varepsilon_x \approx 8$  pm rad (PETRA IV) > 0.1 Å (< 120 keV)  $\Rightarrow \varepsilon_x \approx 0.8$  pm rad (FCC-ee booster + U<sub>0</sub> x 94)

# Small emittance high brilliance and high coherent flux

= rms beam size  $\sigma$  $\sigma'$  = rms beam divergence

High brilliance

Brilliance = 
$$\frac{Flux}{4\pi^{2}\Sigma_{x}\Sigma_{y}\Sigma'_{x}\Sigma'_{y}}$$
  
Source dimension and divergence

 $\Sigma_{x,y} = \sqrt{\sigma_{x,y}^2 + \sigma_{ph}^2} \qquad \Sigma'_{x,y} = \sqrt{\sigma'_{x,y}^2 + \sigma'_{ph}^2}$ 





Y. Hwu & G. Margaritondo, JSR 28-3, 1014 (2021)



$$\sigma_{x,y} = \sqrt{\varepsilon_{x,y} \beta_{x,y}}$$
$$\sigma'_{x,y} = \sqrt{\varepsilon_{x,y} / \beta_{x,y}}$$

Different approximations of single electron undulator emission to gaussian beam

	$\sigma_{ph}$	$\sigma_{ph}'$	$arepsilon_{ph}$	$\beta_{ph} = \sigma_{ph} / \sigma'_{ph}$
Kim (NIM 1986)	$\sqrt{\lambda/L}$	$\sqrt{\lambda L}/4\pi$	$\lambda/4\pi$	$L/4\pi$
Kim (PAC87)	$\sqrt{\lambda/2L}$	$\sqrt{2\lambda L}/4\pi$	$\lambda/4\pi$	$L/2\pi$
Ellaume (2003)	$\sqrt{\lambda/2L}$	$\sqrt{2\lambda L}/2\pi$	$\lambda/2\pi$	$L/\pi$
Lindberg & Kim (2015)	$\sqrt{\lambda/4L}$	$\sqrt{\lambda L}/2\pi$	$\lambda/4\pi$	$L/\pi$

Maximum brilliance for  $\varepsilon_{x,y} \ll \varepsilon_{ph}$  Brilliance  $= \frac{Flux}{4\pi^2 \Sigma_x \Sigma_y \Sigma'_x \Sigma'_y} \approx \frac{Flux}{4\pi^2 \varepsilon_{ph}^2} = \frac{4 Flux}{\lambda^2}$   $\sigma_{x,y} = \sqrt{\varepsilon_{x,y} \beta_{x,y}}$  $\sigma'_{x,y} = \sqrt{\varepsilon_{x,y} \beta_{x,y}}$ 

For  $\varepsilon_{x,y} \approx \varepsilon_{ph}$  brilliance maximized when the electron and photon beam phase spaces are matched



Without dispersion



High transverse coherence Coherent  $Flux = f_c \cdot Flux$ 

Fraction of X-rays transversally coherent

round beam, DLSR  

$$f_{c} = \frac{\left(\frac{\lambda}{4\pi}\right)^{2}}{\left(\varepsilon \cdot \frac{L}{\pi} + \frac{\lambda L}{4\pi^{2}}\right)\left(\varepsilon \cdot \frac{\pi}{L} + \frac{\lambda}{4L}\right)} \qquad \underbrace{ \underbrace{ \begin{array}{c} \underbrace{ \end{array}} \\ \underbrace{ \end{array}} \\ \underbrace{ \end{array}} \\ \underbrace{ \end{array}} \\ \underbrace{ \begin{array}{c} \underbrace{ \end{array}} \\ \underbrace{ \end{array}} \\ \underbrace{ \begin{array}{c} \underbrace{ \end{array}} \\ \underbrace{ \end{array}} \\ \underbrace{ \begin{array}{c} \underbrace{ \begin{array}{c} \underbrace{ \end{array}} \\ \underbrace{ \end{array}} \\ \underbrace{ \end{array}} \\ \underbrace{ \begin{array}{c} \underbrace{ \begin{array}{c} \underbrace{ \end{array}} \\ \underbrace{ \end{array}} \\ \underbrace{ \begin{array}{c} \underbrace{ \end{array}} \\ \underbrace{ \end{array}} \\ \underbrace{ \begin{array}{c} \underbrace{ \end{array}} \\ \underbrace{ \end{array}} \\ \underbrace{ \begin{array}{c} \underbrace{ \end{array}} \\ \underbrace{ \begin{array}{c} \underbrace{ \end{array}} \\ \underbrace{ \begin{array}{c} \underbrace{ \end{array}} \\ \underbrace{ \end{array}} \\ \underbrace{ \begin{array}{c} \underbrace{ \end{array}} \\ \underbrace{ \begin{array}{c} \underbrace{ \end{array}} \\ \underbrace{ \end{array}} \\ \underbrace{ \begin{array}{c} \underbrace{ \end{array}} \\ \underbrace{ \begin{array}{c} \underbrace{ \end{array}} \\ \underbrace{ \end{array}} \\ \underbrace{ \begin{array}{c} \underbrace{ \end{array}} \\ \underbrace{ \begin{array}{c} \underbrace{ \end{array}} \\ \underbrace{ \begin{array}{c} \underbrace{ \end{array}} \\ \underbrace{ \end{array}} \\ \underbrace{ \begin{array}{c} \underbrace{ \end{array}} \\ \underbrace{ \begin{array}{c} \underbrace{ \end{array}} \\ \underbrace{ \end{array}} \\ \underbrace{ \begin{array}{c} \underbrace{ \end{array}} \\ \underbrace{ \begin{array}{c} \underbrace{ \end{array}} \\ \underbrace{ \begin{array}{c} \underbrace{ \end{array}} \\ \underbrace{ \end{array}} \\ \underbrace{ \begin{array}{c} \underbrace{ \end{array}} \\ \underbrace{ \begin{array}{c} \underbrace{ \end{array}} \\ \underbrace{ \end{array}} \\ \underbrace{ \end{array}} \\ \underbrace{ \begin{array}{c} \underbrace{ \end{array}} \\ \underbrace{ } \\ \underbrace{ \end{array}} \\ \underbrace{ \begin{array}{c} \underbrace{ \end{array}} \\ \underbrace{ \end{array}} \\ \underbrace{ \begin{array}{c} \underbrace{ \end{array}} \\ \underbrace{ } \\ \\ \underbrace{ } \\ \\ \underbrace{ } \\ \underbrace{ } \\ \\ \underbrace{ } \\ \\ \underbrace{ } \\ \underbrace{ } \\ \underbrace{ } \\ \\ \underbrace{ } \\ \\ \underbrace{ } \\ \underbrace{ } \\ \\ \\ \underbrace{ } \\ \underbrace{ } \\ \underbrace{ } \\$$

flat beam, FCC-ee booster

$$f_c = f_{ch} = \frac{\binom{\lambda}{4\pi}}{\sqrt{\varepsilon_x \cdot \frac{L}{\pi} + \frac{\lambda L}{4\pi^2}} \sqrt{\varepsilon_x \cdot \frac{\pi}{L} + \frac{\lambda}{4L}}}$$

(1, 1)





#### Scientific opportunities for hard X-ray lasing (40-100 keV)

In general: Non-reproducible processes that live on time-scales of microseconds to femtoseconds and happen in difficult to access sample environments.

Two examples are (for more see workshop *Scientific opportunities with very hard x-ray FEL radiation, European XFEL Jan 18-20 2023*):

- High energy density (HED) science: Extreme conditions can be produced only for short times and once per setup (e.g. shock-compressed diamond anvil cell, pulsed magnetic fields). These set-ups are so complex that detectors have to be placed far away and see only a limited solid angle. Hard X-rays can penetrate the sample environments and compress the Q-space in scattering experiments.
- In-situ microscopy on technological processes, e.g. welding or battery research: On the best storage rings, these experiments are current flux-limited to a time resolution of seconds. With very hard X-ray lasing in burst-mode this can be pushed to microseconds and faster.



Image source: Science, doi: 10.1126/article.30069



Image source: Chinese Journal of Lasers, Volume. 46, Issue 9, 902003(2019)

**Courtesy of Harald Sinn** 



#### Above 100 keV up to few MeV

Simulations performed with SPECTRA T.Tanaka, JSR (2021). 28, 1267-1272

## **E**γ < 20 MeV

#### Applications

- Nuclear structure and astrophysics
- Nuclear applications
- Hadronic parity violation

International workshop on next generation gamma-ray sou C R Howell et al 2022 J. Phys. G: Nucl. Part. Phys. 49 01(



Vacuum gap (mm)	5
λ <sub>υ</sub> (mm)	13
B <sub>max</sub> (T)	2.2



#### 60 MeV <Eγ < 350 MeV

Low energy QCD



**Figure 18.** The appropriate degrees of freedom with which to describe nucleon structure depend on the distance scale at which it is probed: quarks and gluons at high energies (left); or nucleons, their pion clouds and nucleonic excitations at low energies (right).

International workshop on next generation gamma-ray source C R Howell et al 2022 J. Phys. G: Nucl. Part. Phys. 49 010502

	Uo
beam energy [GeV]	182.5
avg. beam current [mA]	4.9
Avg. number of bunches	56
rms bunch length [mm]	2.3
rms relative energy spread [10 <sup>-3</sup> ]	2
beta at wiggler /undulator [m]	1.6
hor. emittance [pm rad]	1400
vert. emittance [pm rad]	2.8

#### In the collider up to 100 MeV



Simulations performed with SPECTRA T.Tanaka, JSR (2021). 28, 1267-1272

#### In the collider with SCUs up to 300 MeV





Simulations performed with SPECTRA T.Tanaka, JSR (2021). 28, 1267-1272



#### **XFEL from linac**

- Possible parameters after photocathode electron RF gun + damping ring inserting bunch compressors at positions to be studied
- Doubling the number of klystrons in the high energy linac it's possible to increase the electron beam energy to 25 GeV
- No show stoppers but more detailed simulations including collective effects are necessary

Energy (GeV)	25
Bunch length (µm)	15
Slice emittance x,y (µm rad)	0.4
Bunch charge (pC)	250
Slice energy spread	0.2‰
Beta_x,y (m)	30

#### **Electron beam parameters\***

\*Courtesy of P. Craievich (PSI)

High energy photons from the FCC-ee complex

## **Booster filling patterns**

During injector commissioning: exploratory studies

Injector can be used as XFEL with WW, ZH and ttbar modes

Only few bunches can be used in the Z mode for an XFEL (not possible during accumulation at 100 Hz)

To be studied if the abort gap would allow for additional bunches for the XFEL





#### **XFEL from linac**

#### **Electron beam**

Energy (GeV)	25
Bunch length (μm)	15
Emittance x,y (µm rad)	0.4
Bunch charge (pC)	250
Energy spread	0.2‰
Beta_x,y (m)	30

4 bunches separate by 25 ns at 100 Hz when booster is not in accumulation mode

#### Superconducting Undulator

NbTi@4K	
Vacuum gap (mm)	5
λ <sub>u</sub> (mm)	18
B <sub>max</sub> (T)	1.83

L <sub>sat</sub> ( m)	85 m
Energy <sub>min</sub> (keV)	60
Pulse energy (mJ)	0.95

#### HTS@4K

Vacuum gap (mm)	5
λ <sub>U</sub> (mm)	13
B <sub>max</sub> (T)	2.2

L <sub>sat</sub> ( m)	100
Energy <sub>min</sub> (keV)	100
Pulse energy (mJ)	0.4

To be studied if possible to increase up to trains 100 ns long with > 4 bunches

**European XFEL** 

Calculations with Ming Xie'snFEL formula in SIMPLEX, T. Tanaa J Synchrotron Radiations 2015, 22(5) 1319-26

#### Conclusions

- With respect to PETRA IV, planned diffraction limited storage ring with smallest emittance, the FCC-ee booster has the potential to produce at 50-100 keV
  - a fraction of coherent X-rays larger by one order of magnitude
  - an average brilliance larger by up to two orders of magnitude
  - a peak brilliance larger by up to four orders of magnitudes

#### Conclusions

Use the booster at collision energy or the collider to produce photons up to 20 MeV for

Nuclear structure and astrophysics
 Nuclear applications
 Hadronic parity violation

and from 60 to 300 MeV for low energy QCD

The injector with superconducting undulators could be used as XFEL for producing hard to very hard X-rays with pulse energy ~mJ. This would be possible mainly for WW, ZH and ttbar modes

#### Outlook

To be further studied the compatibility of:

the use of the FCC-ee booster for the collider and requirements from possible users

the parameters of the FCC-ee injector used as a linac-based XFEL

Other Science Opportunities at the FCC-ee, Sara Casalbuoni, 28.11.2024

# Thank you for your attention !