

Ultimate limit on beam brightness



Ilya Agapov, **Sergey Antipov**, C. Cortés Garcia

Other Science Opportunities at the FCC-ee, CERN Geneva

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Why would one want a light source at high energy?

Many contemporary light sources (will) operate at 3-6 GeV

Parameters of operating and planned 4th generation light sources

Machine	Energy	Circumference	Hor. emittance
P4	6 GeV	2.3 km	20 pm
ESRF-EBS	6 GeV	844 m	110 pm
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MAX-IV	3 GeV	528 m	200 pm

*FCC-ee Booster**

20 GeV (inj.)

90.7 km

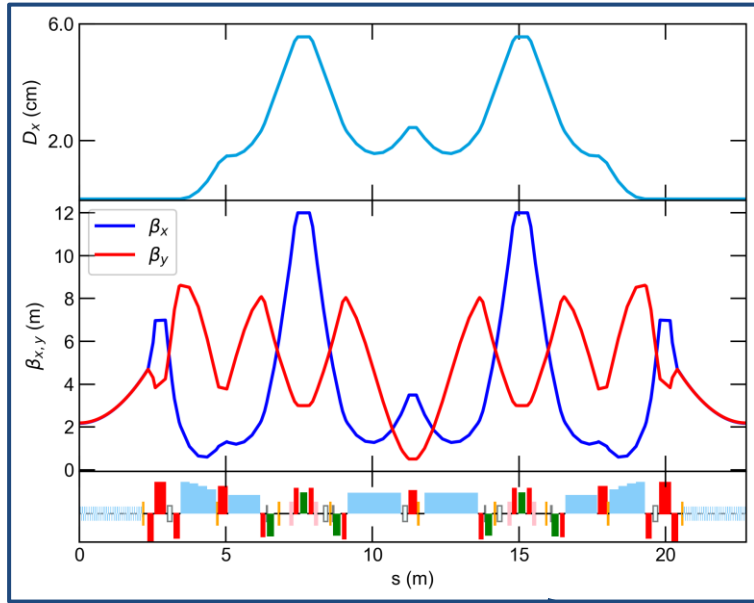
15 pm

* S. Casalbuoni, F. Zimmermann, FCC Week'21

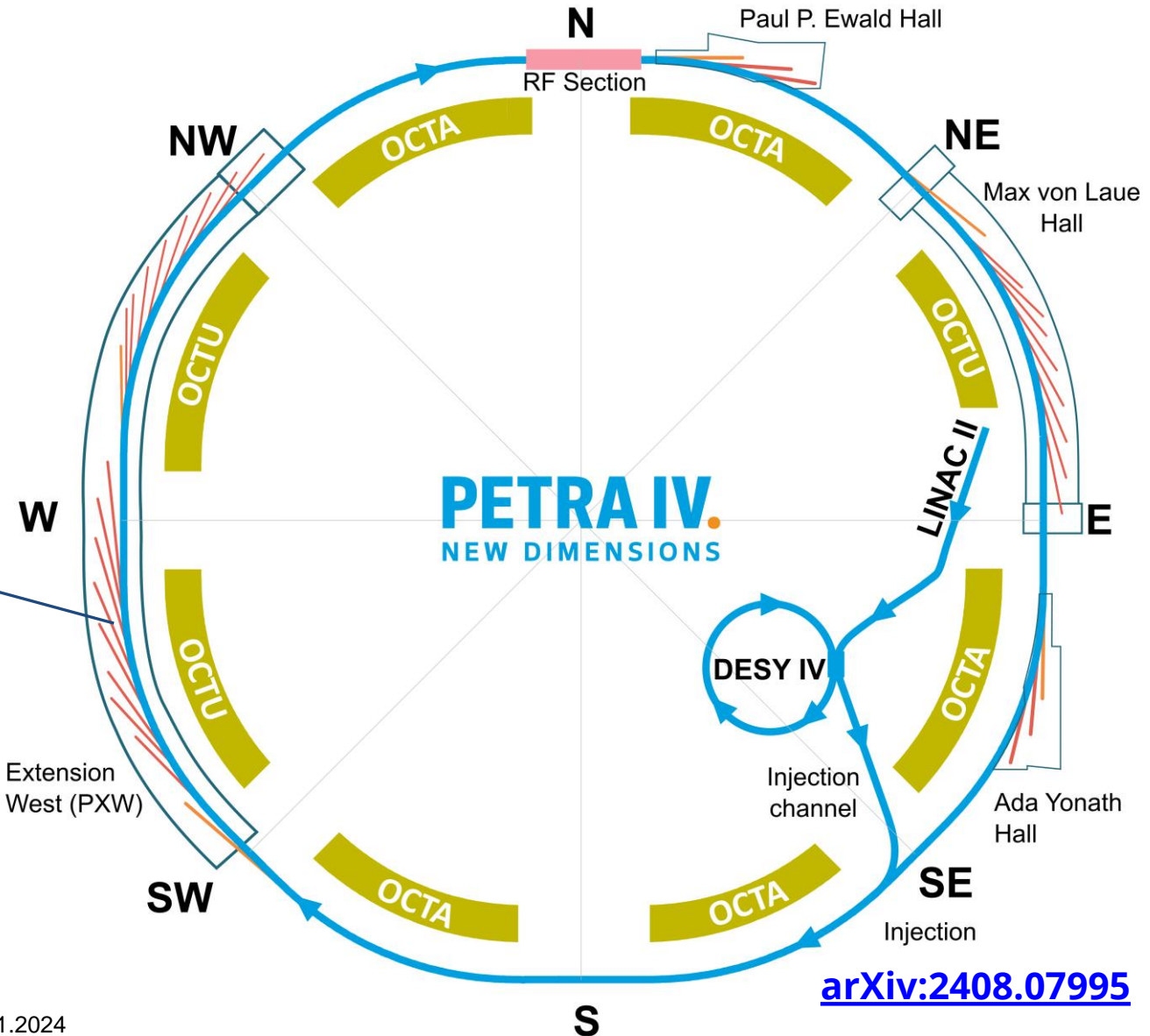


PETRA IV: Germany's future flagship light source

6 GeV, 2.3 km, 20 pm



Highly optimized achromatic cell to approach diffraction limit at 10 keV



HELMHOLTZ

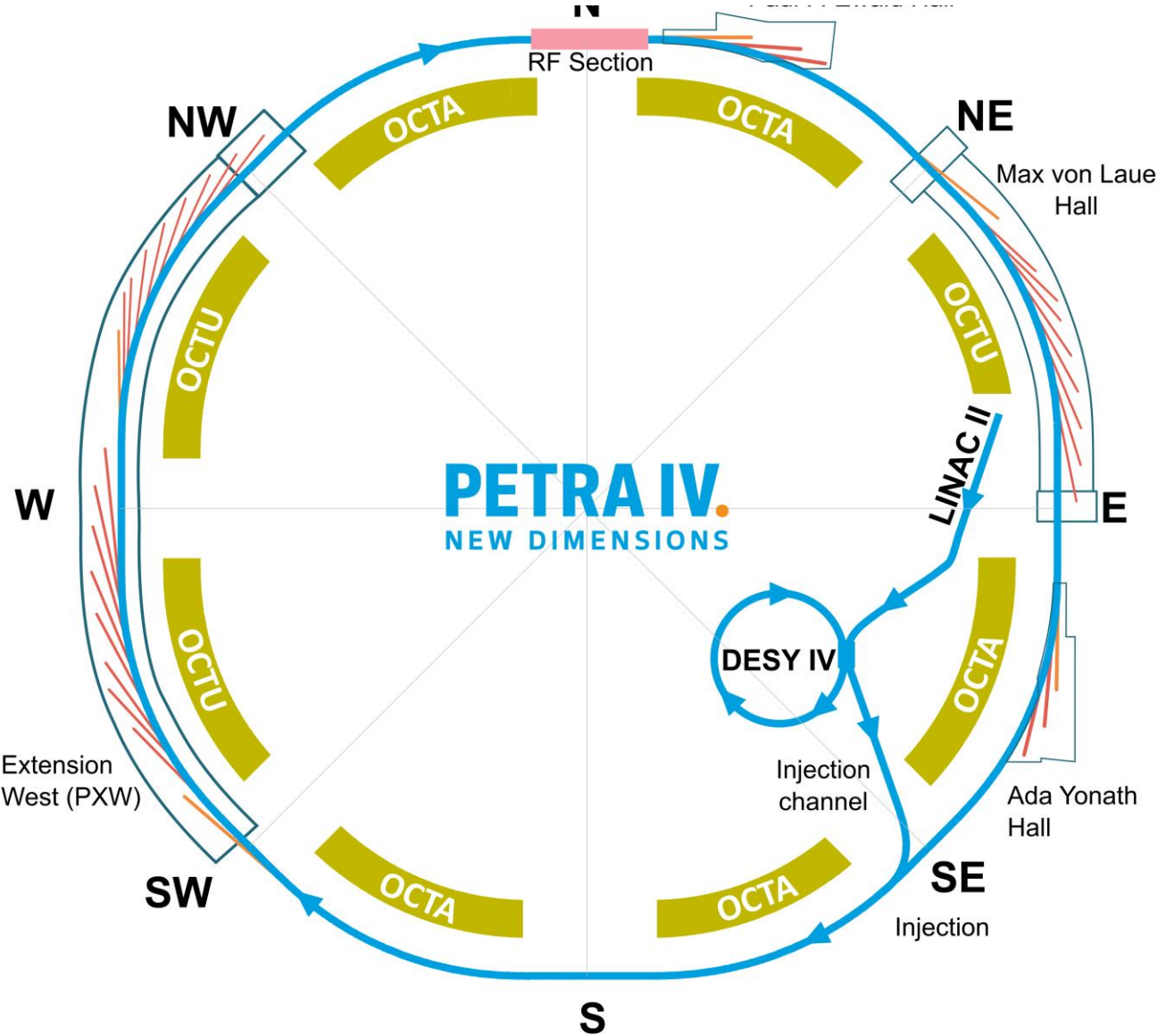


Space charge tune shift

6 GeV, 2.3 km, 20 pm

$$|\Delta\nu_{x,y}^{SC}| = \frac{N r_e C}{(2\pi)^{\frac{3}{2}} \gamma^3 \sigma_z} \left\langle \frac{\beta_{x,y}}{\sigma_{x,y} (\sigma_x + \sigma_y)} \right\rangle$$

1/10¹²

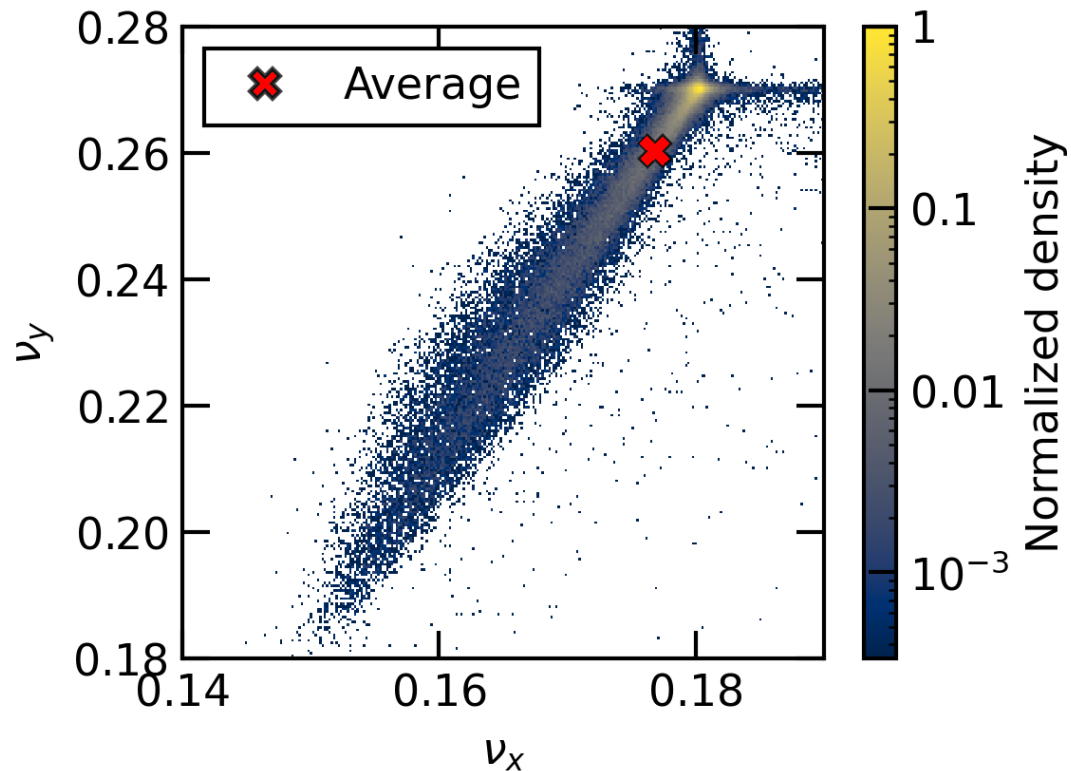


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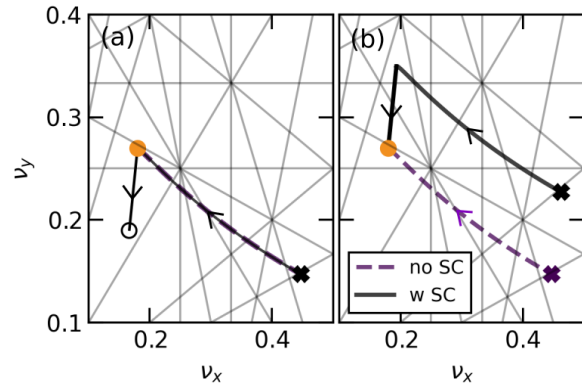


Tune footprint with SC. Simulation in XSuite.
10 nC bunch, 128 turns, long. wakes included,
single-harmonic RF system

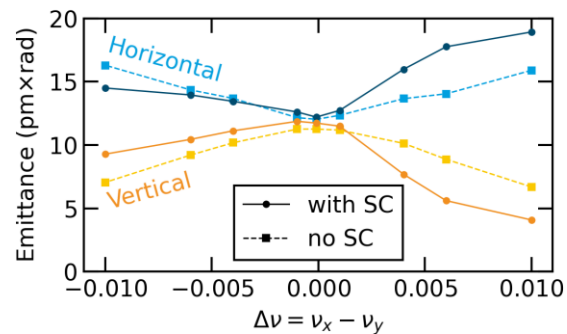
Space charge effects

Both coherent and incoherent

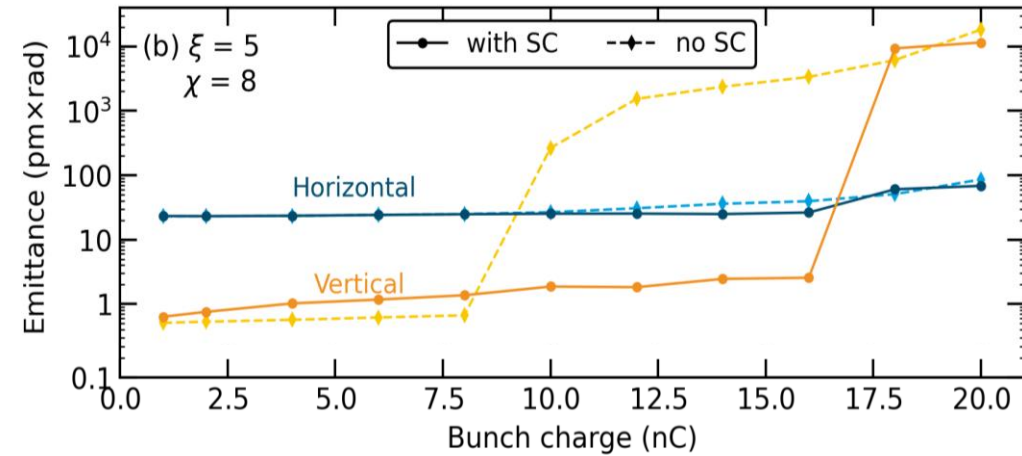
Dynamics of off-axis top-up injection



Emittance ratios at high coupling



Single-bunch intensity limit



SC will ultimately lead to crossing of the integer resonance

This shall limit the bunch intensity even if there is no beam coupling impedance

$$|\Delta\nu_{x,y}^{SC}| = \frac{N r_e C}{(2\pi)^{\frac{3}{2}} \gamma^3 \sigma_z} \left| \frac{\beta_{x,y}}{\sigma_{x,y} (\sigma_x + \sigma_y)} \right| \lesssim \frac{1}{2}$$

Key Assumptions:

- Linear SC model
- Integer resonance cannot be crossed
- The largest tune shift is in the vertical plane
- Dispersion contribution to beam size is small, on average
- Coupling is sufficiently small

Approximate SC limit

Crossing of the integer resonance

$$|\Delta v_{x,y}^{SC}| = \frac{N r_e C}{(2\pi)^2 \gamma^3 \sigma_z} \left\langle \frac{\beta_{x,y}}{\sigma_{x,y}(\sigma_x + \sigma_y)} \right\rangle \lesssim \frac{1}{2}$$

$$|\Delta v_y^{SC}| \approx \frac{N r_e C}{(2\pi)^2 \gamma^3 \sigma_z} \left\langle \frac{\beta_y}{\sqrt{\beta_y \epsilon_y} (\sqrt{\beta_y \epsilon_y} + \sqrt{\beta_x \epsilon_x})} \right\rangle$$

$$\left\langle \frac{\beta_y}{\sqrt{\beta_y \epsilon_y} (\sqrt{\beta_y \epsilon_y} + \sqrt{\beta_x \epsilon_x})} \right\rangle = \frac{1}{\epsilon_x \kappa} \left\langle \frac{\sqrt{\kappa} \sqrt{\frac{\beta_y}{\beta_x}}}{1 + \sqrt{\kappa} \sqrt{\frac{\beta_y}{\beta_x}}} \right\rangle \approx \frac{1}{\epsilon_x \kappa} \left\langle \sqrt{\kappa} \sqrt{\frac{\beta_y}{\beta_x}} \left(1 - \sqrt{\kappa} \sqrt{\frac{\beta_y}{\beta_x}} \right) \right\rangle \sim \frac{1}{4\epsilon_x \kappa} *$$

$$\kappa = \epsilon_y / \epsilon_x$$

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to be checked later

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$$\kappa = \epsilon_y / \epsilon_x$$

$$B \times \epsilon_x \lesssim \frac{8 \gamma^3 I_A}{\pi C}$$

$$|\Delta v_x^{SC}| \sim \frac{N r_e C}{4 (2\pi)^{\frac{3}{2}} \gamma^3 \sigma_z \epsilon_x \kappa} \lesssim \frac{1}{2}$$

$$B = \frac{2I}{\pi^2 \epsilon_x \epsilon_y}$$

Approximate SC limit

Crossing of the integer resonance

$$|\Delta v_{x,y}^{SC}| = \frac{N r_e C}{(2\pi)^2 \gamma^3 \sigma_z} \left\langle \frac{\beta_{x,y}}{\sigma_{x,y} (\sigma_x + \sigma_y)} \right\rangle \lesssim \frac{1}{2}$$

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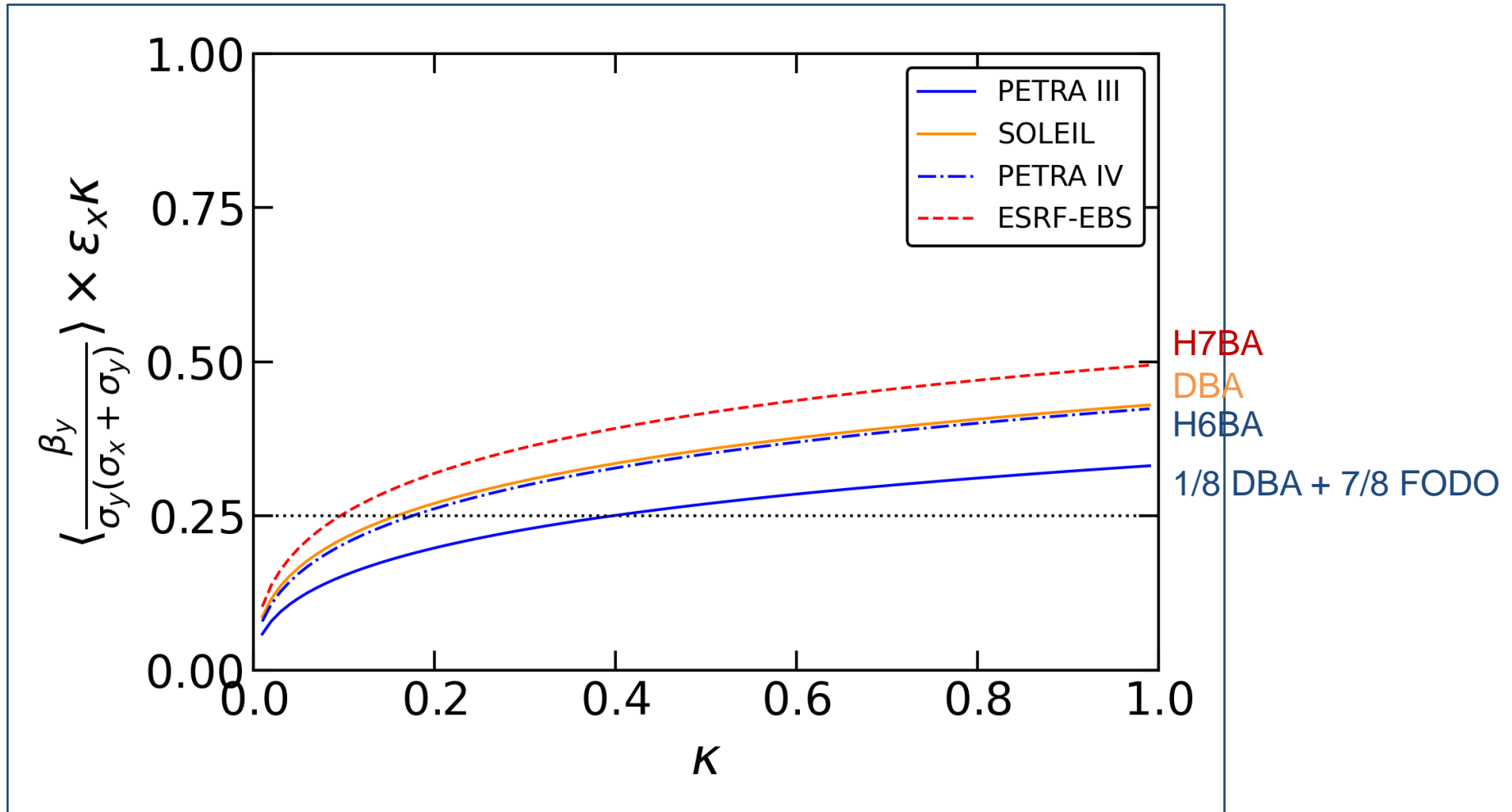
Example: PETRA IV, high charge Timing mode

Lorentz factor	γ	11.7×10^3
Circumference	C	2.3 km
Bunch charge	Ne	8 nC
Hor. emittance	ϵ_x	20 pm
Bunch length	σ_z	20 mm
Design brightness	B_{Timing}	$2.4 \times 10^{23} \text{ A/m}^2$
SC limit	B_{SC}	$1.4 \times 10^{24} \text{ A/m}^2$

$$B \times \epsilon_x \lesssim \frac{8\gamma^3 I_A}{\pi C}$$

Approximation holds in a wide parameter range

Within a factor 2 for any conceivable coupling ratio, any machine



Horizontal emittance at high energy

Governed by a balance of synchrotron radiation and quantum excitation

$$\left(\frac{d\epsilon_x}{dt}\right)_{IBS} + \left(\frac{d\epsilon_x}{dt}\right)_{SR} + \left(\frac{d\epsilon_x}{dt}\right)_{QE} = 0 \quad \text{IBS is a small correction}$$

$$\left(\frac{d\epsilon_x}{dt}\right)_{SR} + \left(\frac{d\epsilon_x}{dt}\right)_{QE} \approx 0$$

$$\epsilon_x \approx \frac{55\hbar c}{32\sqrt{3}m_e c^2} \frac{\gamma^2 I_5}{J_x I_2}$$

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$$J_x \in [1; 2]$$

$$\frac{I_5}{I_2} = \frac{\oint \mathcal{H} / \rho^3 ds}{\oint 1 / \rho^2 ds} \propto 1/C$$

$$\epsilon_x = F \frac{\gamma^2}{C}$$

F – scale-independent parameter

- defined by the ring's optics lattice,
- not its physical scale or energy.
- HMBA lattices: $F = 3 - 7 \times 10^{-16} \text{m}^2$

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$$B \lesssim \frac{8\gamma I_A}{\pi F}$$

A 20 GeV-range machine should have a higher brightness limit compared to most 4th generation light sources

Space charge imposes a limit on achievable beam brightness

The way to push further is to aim at higher energies

Space charge can no longer be ignored in state-of-the art light sources

- Emittance reaches the point where it becomes strong enough
- Generates incoherent tune shift and spread affecting beam dynamics

There is a limit on achievable beam brightness

- Due to crossing of the integer resonance
- Consequently, the photon brilliance is limited

The SC limit increases with energy as the interaction weakens

- Potentially, higher photon brilliance at higher beam energies

Extras: emittance scaling factor for different machines

Many state-of-the art HMBA machines have similar performance despite different scales

Machine	Energy	Circumference	Hor. emittance	F
P4 baseline	6 GeV	2.3 km	20 pm	$3.3 \times 10^{-16} \text{m}^2$
HEPS	6 GeV	1.36 km	35 pm	$3.5 \times 10^{-16} \text{m}^2$
P4 no DW	6 GeV	2.3 km	40 pm	$6.6 \times 10^{-16} \text{m}^2$
ESRF-EBS	6 GeV	844 m	110 pm	$6.8 \times 10^{-16} \text{m}^2$
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20 GeV (inj.)

90.7 km

15 pm

$\sim 10 \times 10^{-16} \text{m}^2$

Extras: Intra-beam scattering

A small correction at high energy

Bjorken-Mtingwa model:

$$\sigma_H = \left(\frac{1}{\sigma_p^2} + \frac{\mathcal{H}_x}{\epsilon_x} + \frac{\mathcal{H}_y}{\epsilon_y} \right)^{-1/2}$$

$$a = \frac{\sigma_H}{\gamma} \sqrt{\frac{\beta_x}{\epsilon_x}}, \quad b = \frac{\sigma_H}{\gamma} \sqrt{\frac{\beta_y}{\epsilon_y}},$$

$$g(x) \approx 2x^{0.021-0.044\ln(x)}$$

$$\alpha_x^{IBS} \approx \frac{r_e^2 c N L_C}{32 \gamma^3 \epsilon_x^{\frac{7}{4}} \epsilon_y^{\frac{3}{4}} \sigma_z \sigma_p} \left\langle \mathcal{H}_x \sigma_H g\left(\frac{a}{b}\right) \beta_x^{-1/4} \beta_y^{-1/4} \right\rangle$$

$$\sigma_z \sigma_p \propto U_0^{1/4} \propto \gamma$$

SR damping rate:

$$\alpha_x^{SR} = \gamma^3 \frac{r_e c I_2}{3C}$$

Equilibrium emittance:

$$\epsilon_x = \frac{\epsilon_{x,0}}{1 - \alpha_x^{IBS} / \alpha_x^{SR}}$$

See *K. Bane, SLAC-AP-141*
and *R. Bartolini, IPAC'22*