

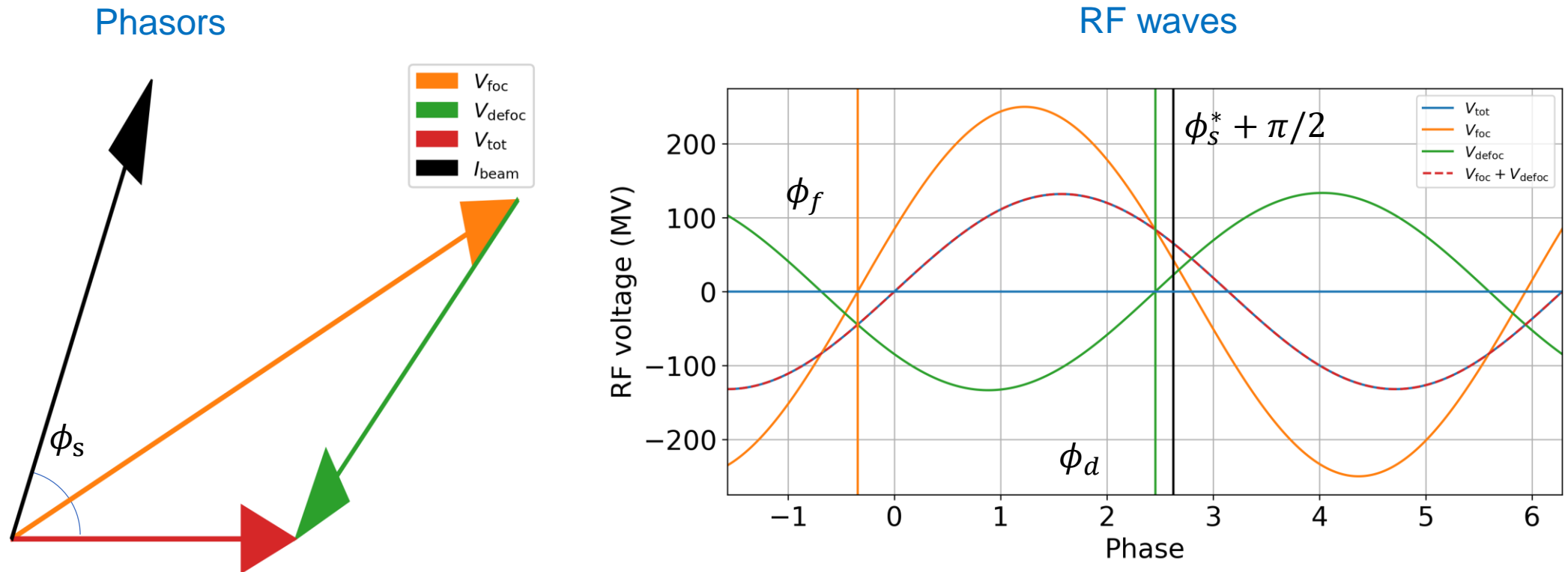
Transient beam loading for reverse phase operation mode

Ivan Karpov, Franck Peauger, Xavier Buffat with input from:
Olivier Brunner, Rama Calaga, Heiko Damerau, Jiquan Guo (JLAB), Frank Gerigk, Wolfgang Höfle,
Eric Montesinos, Igor Syratchev, Shahnam Gorgi Zadeh, Alice Vanel, Jorg Wenninger

Reverse phase operation

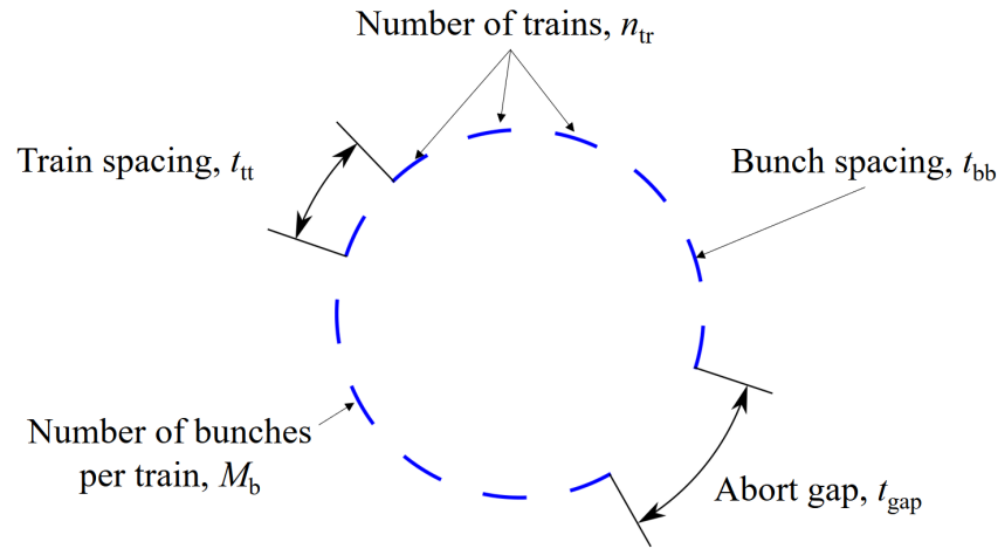
Reverse phase operation (RPO) mode allows increasing RF cavity voltage (*Y. Morita et al., SRF, 2009*)

- Experimentally verified with high beam loading in KEKB (*Y. Morita et al., IPAC, 2010*)
- Baseline solution for EIC ESR (*e.g., J. Guo et al., IPAC, 2022*)

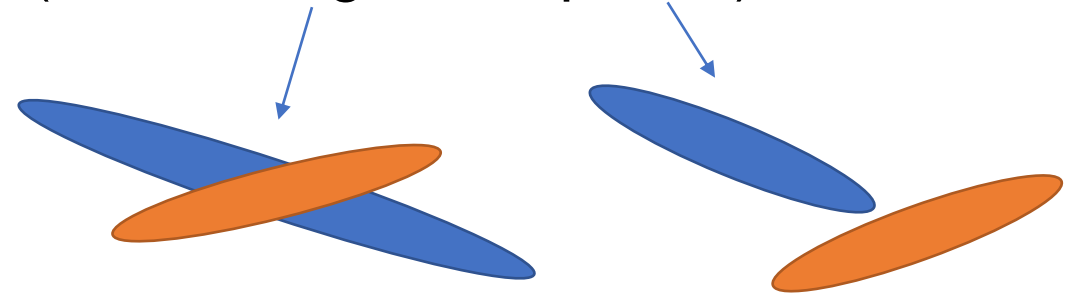


*Electron convention for synchronous phase: $\phi_{s,proton} = \frac{\pi}{2} + \phi_{s,electron}$

Transient beam loading



Gaps in machine filling will result in modulation beam parameters (bunch length and phase)



→ Modulations might impact luminosity and/or beam stability

Conventional approaches:

- Particle tracking simulations (difficult for **11200** bunches in FCC-ee Z)
- Steady-state time domain method (*J. Tückmantel, 2011*)
- Small-signal model in frequency domain (*F. Pedersen, 1992*)

→ Tracking simulation were applied for EIC (1160 bunches)

→ Small-signal model was adapted for the RPO case of FCC up to now

Reduced Pedersen model

General equations of beam-cavity interactions with reverse phase operation (RPO) mode (adaptation of formalism in [J. Tückmantel, 2011](#)):

$$I_{gf}(t) = \frac{V_f(t)}{2(R/Q)} \left(\frac{1}{Q_L} - 2i \frac{\Delta\omega_f}{\omega_{\text{rf}}} \right) + \frac{I_{b,\text{rf}}(t)}{2} + \frac{dV_f(t)}{dt} \frac{1}{\omega_{\text{rf}}(R/Q)}$$

$$I_{gd}(t) = \frac{V_d(t)}{2(R/Q)} \left(\frac{1}{Q_L} - 2i \frac{\Delta\omega_d}{\omega_{\text{rf}}} \right) + \frac{I_{b,\text{rf}}(t)}{2} + \frac{dV_d(t)}{dt} \frac{1}{\omega_{\text{rf}}(R/Q)}$$

Energy balance $V_{\text{tot}} \cos \phi_s = N_f A_f \cos(\phi_s - \phi_b + \phi_{cf} + \phi_f) + N_d A_d \cos(\phi_s - \phi_b + \phi_{cd} + \phi_d)$

To calculate beam-induced modulation we assume:

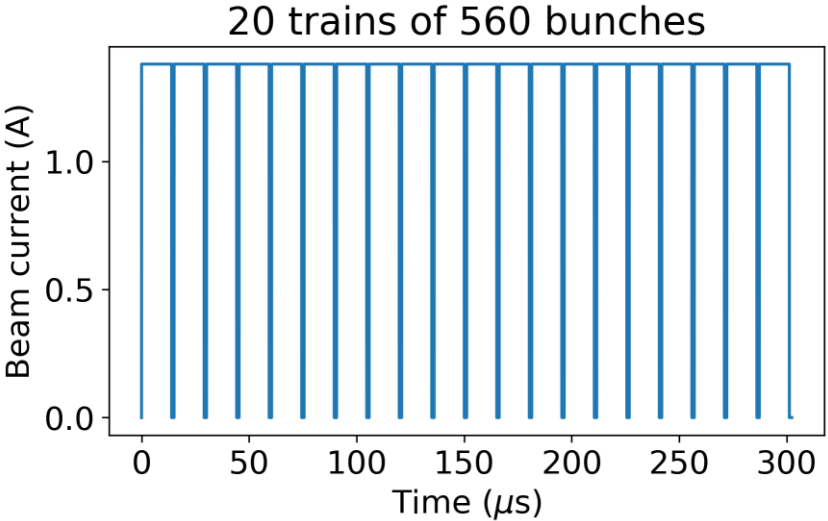
- $I_{gf,d}(t) = \text{constant}$ – no beam loading compensation
- $V_{f,d}(t) = A_{f,d}(t) e^{i\phi_f(t) + i\phi_{cf,d}}$, $I_{b,\text{rf}}(t) = A_b(t) e^{-i\phi_s + i\phi_b(t)}$

Then, system of equations is linearized to obtain transfer functions: $\frac{a_{Vf,d}}{a_b}$, $\frac{\phi_{f,d}}{a_b}$, $\frac{\phi_b}{a_b}$

$$A_{f,d} = V_{\text{cav}}(1 + a_{Vf,d}), A_b(t) = |F_b| I_{b,\text{dc}} (1 + a_b)$$

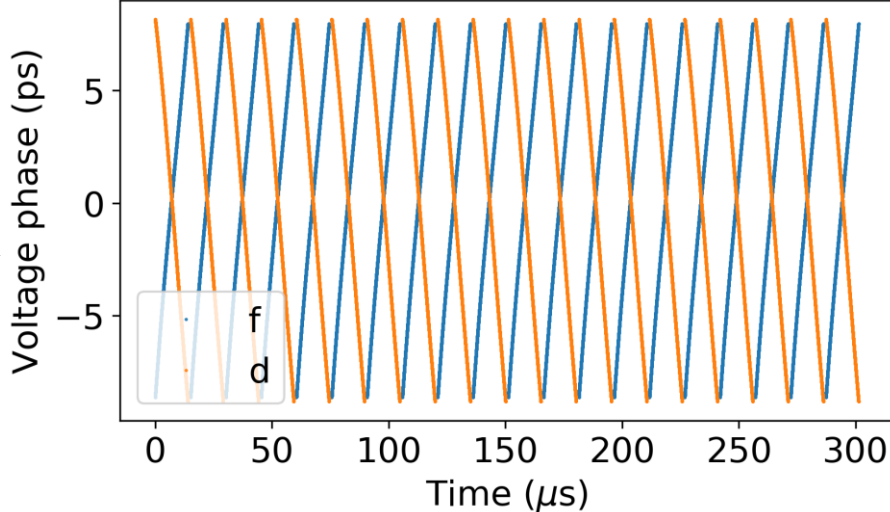
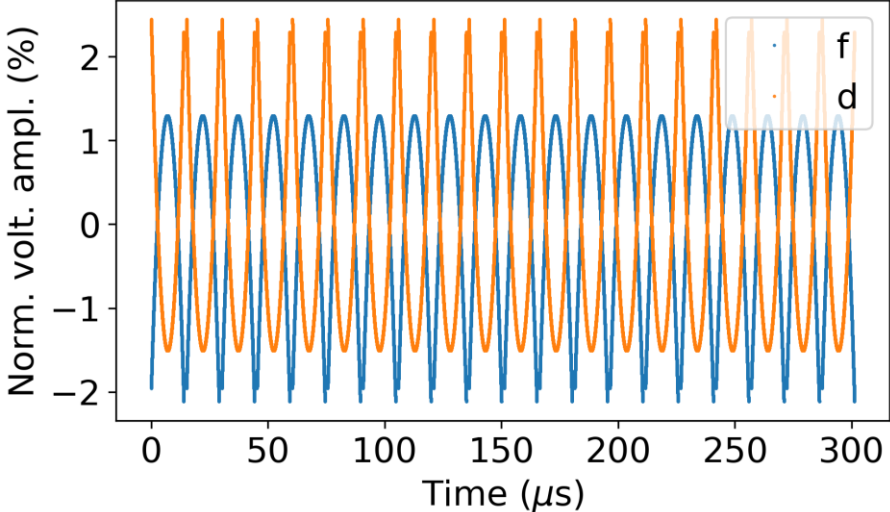
Bunch-by-bunch spread of cavity parameters

	N_f	N_d	V_{tot} Z (MV)	V_{cav} (MV)	Q_L
Current	71	61	88	7.95	9.21e5



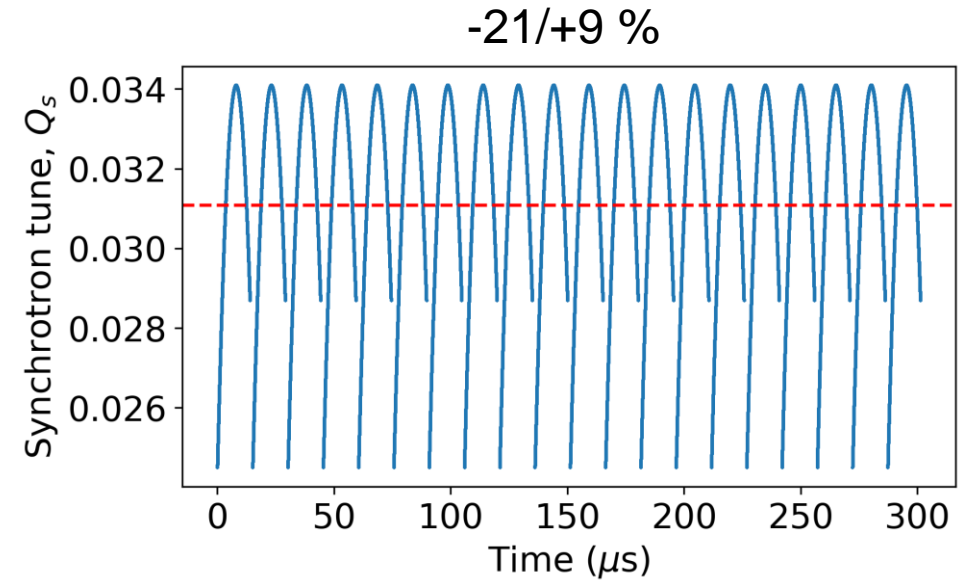
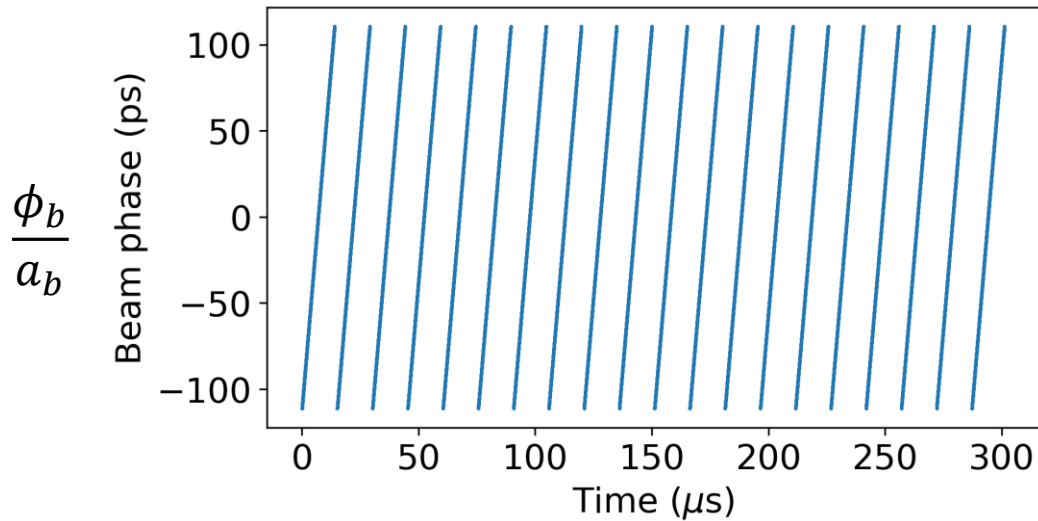
$$\frac{a_{Vf,d}}{a_b}$$

$$\frac{\phi_{f,d}}{a_b}$$



Note, the designed rms bunch length is 50 ps (with beamstrahlung)

Bunch-by-bunch spread of beam parameters

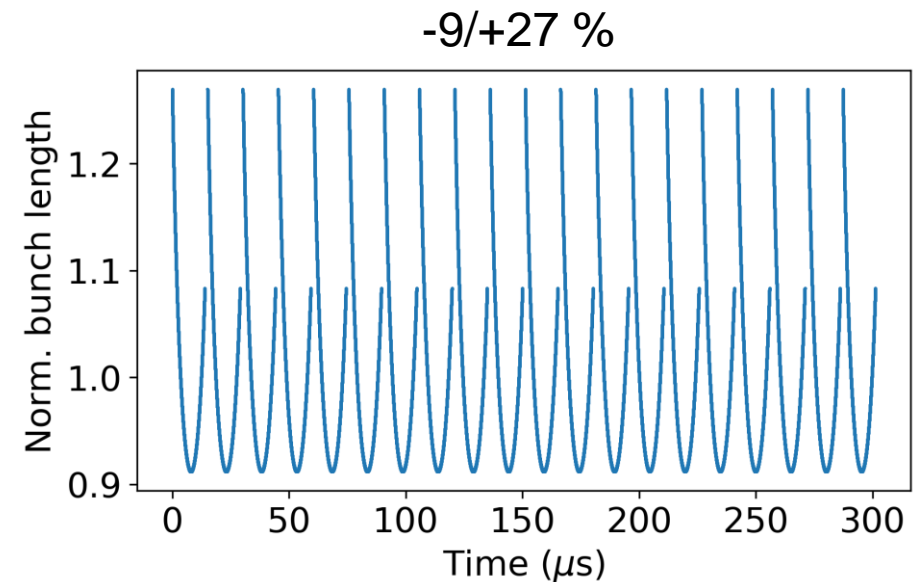


For identical rings, transients can be compensated by matching abort gaps (e.g., in PEP-II, LHC,...)

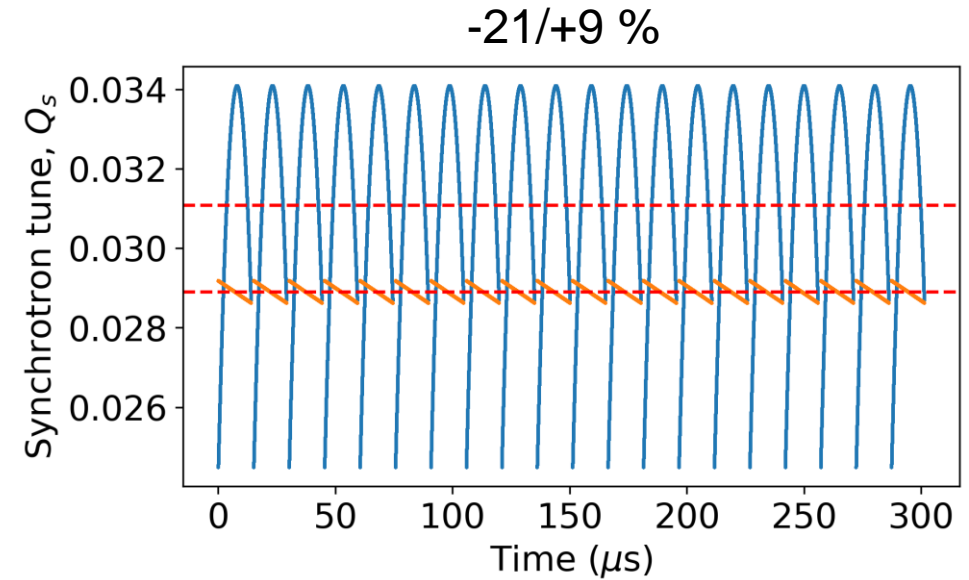
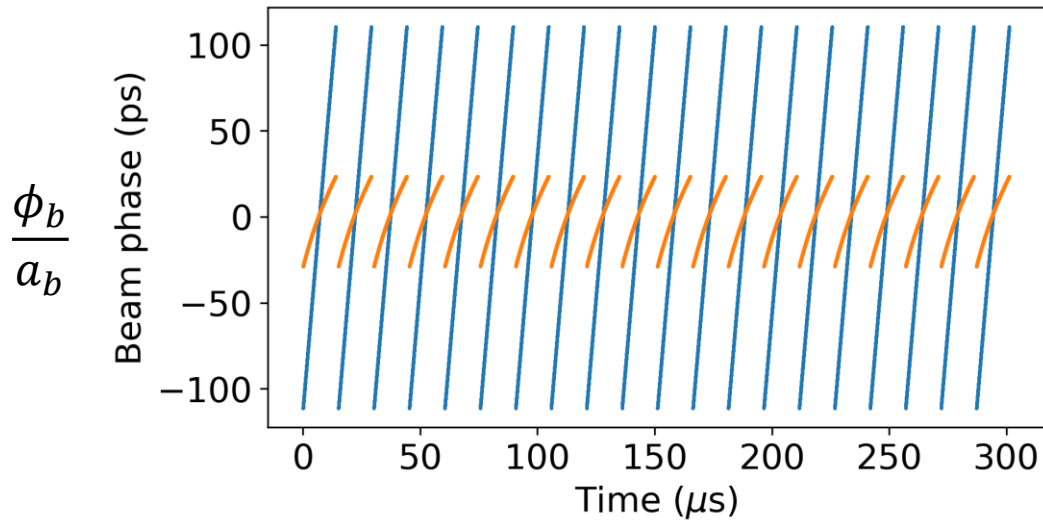
Imbalance of charge results in different detuning for electron and positron beams

→ Slightly different transients (most critical during filling)

Peak-to-peak spread of **~30%** in synchrotron tune and bunch length can have a significant impact on beam stability



Bunch-by-bunch spread of beam parameters



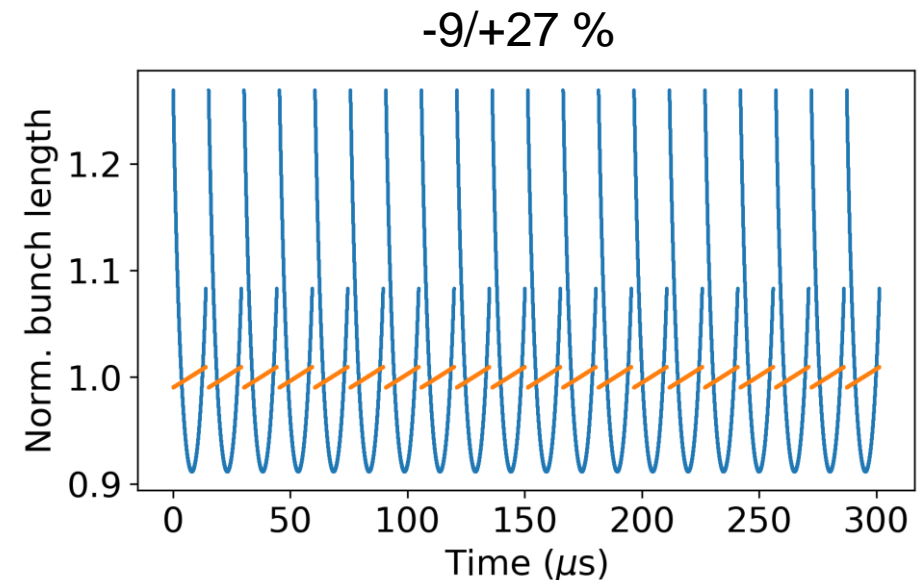
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Imbalance of charge results in different detuning for electron and positron beams

→ Slightly different transients (most critical during filling)

Peak-to-peak spread of **~30%** in synchrotron tune and bunch length can have a significant impact on beam stability

→ We lose **a factor of 15** wrt to 1-cell RF system

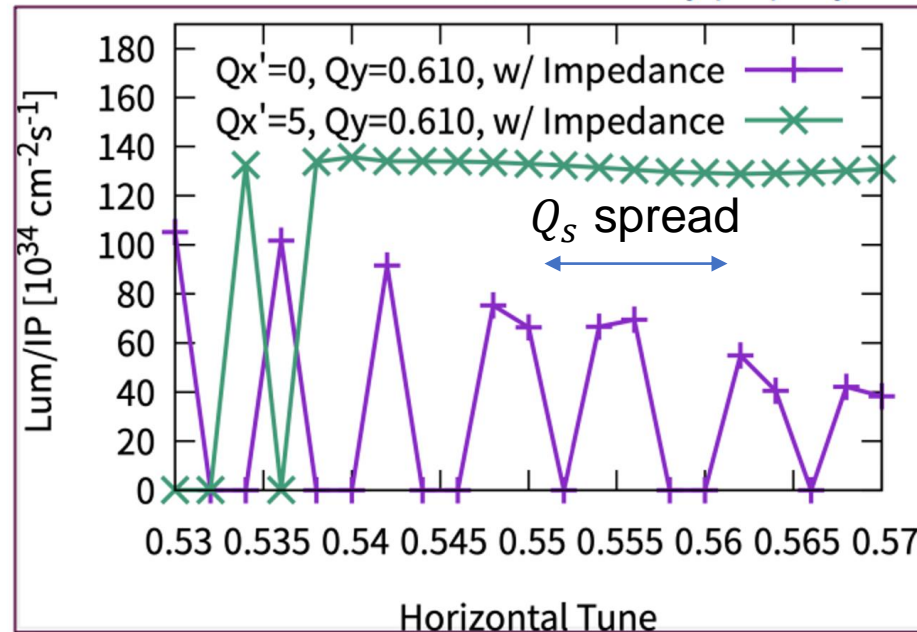


Critical impact of spread

Interplay between beam-beam and coupling impedance

A positive chromaticity has a beneficial effect on the beam-beam. Self-consistent simulations show a luminosity per IP close to the nominal value of $141 \times 10^{34} \text{ cm}^{-2}\text{s}^{-1}$ by properly choosing the collider working point.

Y. Zhang, IHEP, Beijing, China



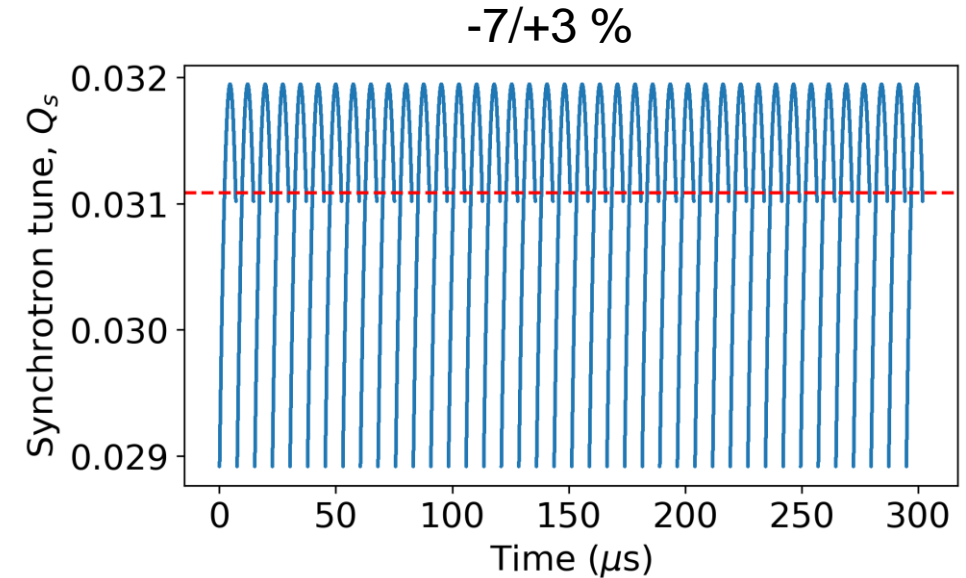
M. Migliorati, FCC Week 2024

→ No stable region for a horizontal tune can be found in presence of large Q_s spread. Possible mitigations need to be studied

Possible scenarios

Peak-to-peak beam phase spread $\propto \Delta\omega_{\text{opt}}\tau_{\text{gap}}N_{\text{tot}}/(N_f - N_d)$

1. New filling scheme (e.g., 40 trains of 280 bunches)
 - Spread is reduced by a factor of ~ 3
 - Gaps become twice shorted (~ 600 ns) - most likely unfeasible for the extraction system (1 us kicker rise time)
2. Higher total RF voltage for Z?



Optimal quality factor $Q_{L,\text{opt}} = \frac{V_{\text{cav}}^2 N_{\text{tot}}}{2P_{\text{SR}}(R/Q)}$

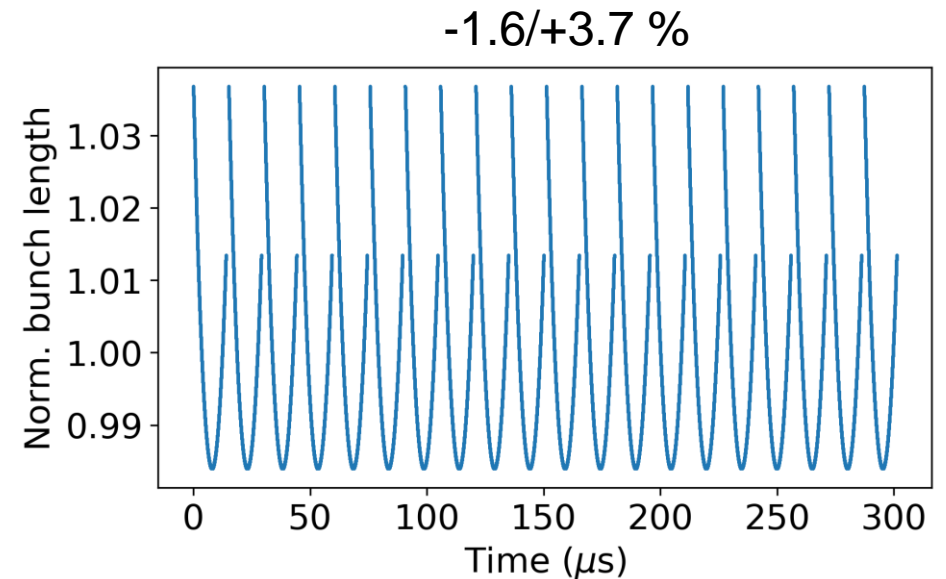
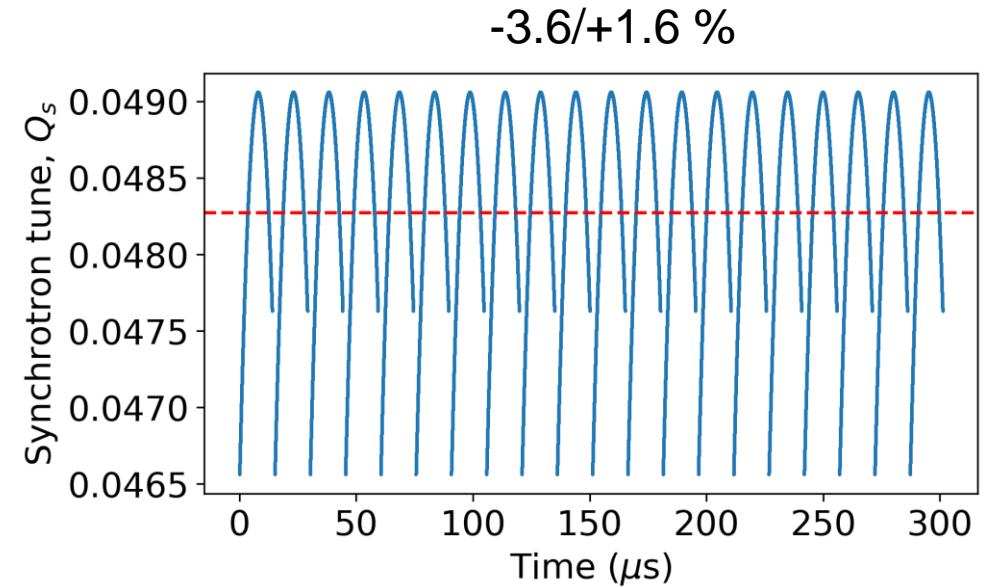
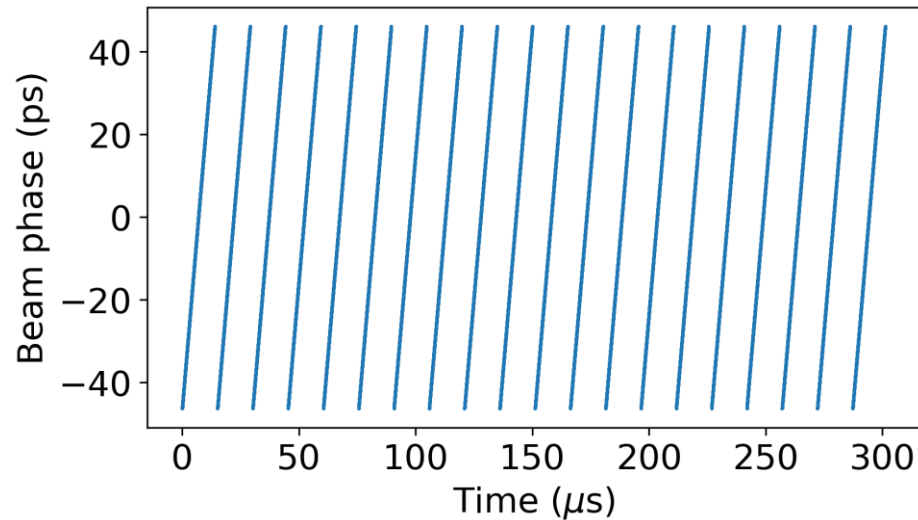
Since $Q_{L,\text{opt}}$ should be the same for Z, W, and ZH, V_{cav} cannot be changed

→ Optimal detuning is also unchanged $\Delta\omega_{\text{opt}} = -\frac{\omega_{\text{rf}}(R/Q)|F_b|I_{b,\text{dc}}}{2V_{\text{cav}}} \sqrt{1 - \frac{U_0^2}{e^2 V_{\text{cav}}^2 N_{\text{tot}}^2}}$

The only knob is to change $N_f - N_d$ by changing V_{tot} : $V_{\text{cav}} = \frac{V_{\text{tot}}}{N_{\text{tot}}} \sqrt{\frac{U_0^2}{e^2 V_{\text{tot}}^2} + \left(1 - \frac{U_0^2}{e^2 V_{\text{tot}}^2}\right) \frac{N_{\text{tot}}^2}{(N_f - N_d)^2}}$

Higher RF voltage

	N_f	N_d	$V_{\text{tot}} Z$ (MV)	V_{cav} (MV)	Q_L
Current	71	61	88	7.95	9.21e5
Option 2	78	54	195	7.95	9.21e5



Higher RF voltage reduces parameter spread to ~5%
 Peak-to-peak Q_s spread $2.5e-3$

Impact on parameters (oversimplified scaling)

FCC-ee collider parameters for the GHC lattice as of Aug. 2, 2024.

Beam energy	[GeV]	45.6	80	120	182.5
Layout			PA31-3.0		
# of IPs			4		
Circumference	[km]		90.658728		
Bend. radius of arc dipole	[km]		10.021		
Energy loss / turn	[GeV]	0.0390	0.369	1.86	9.94
SR power / beam	[MW]		50		
Beam current	[mA]	1283	135	26.8	5.0
Colliding bunches / beam		11200	1852	300	64
Colliding bunch population	[10 ¹¹]	2.16	1.38	1.69	1.48
Hor. emittance at collision ϵ_x	[nm]	0.70	2.16	0.66	1.51
Ver. emittance at collision ϵ_y	[pm]	1.9	2.0	1.0	1.36
Lattice ver. emittance $\epsilon_{y,lattice}$	[pm]	0.87	1.20	0.57	0.94
Arc cell		Long 90/90		90/90	
Momentum compaction α_p	[10 ⁻⁶]	28.67		7.52	
Arc sext families		75		146	
$\beta_{x/y}^*$	[mm]	110 / 0.7	220 / 1	240 / 1	
Transverse tunes $Q_{x/y}$		218.158 / 222.220	218.185 / 222.2	198.2	
Chromaticities $Q'_{x/y}$		0 / +5	0 / +5		
Energy spread (SR/BS) σ_δ	[%]	0.039 / 0.110	0.069 / 0.105	0.102 / 0.176	
Bunch length (SR/BS) σ_z	[mm]	5.57 / 15.6	3.46 / 5.28	1.91 / 2.32	
RF voltage 400/800 MHz	[GV]	0.079 / 0	1.00 / 0	2.1 / 9.20	
Harm. number for 400 MHz					
RF frequency (400 MHz)	MHz		40		
Synchrotron tune Q_s		0.0289	0.0809		0.0881
Long. damping time	[turns]	1171	218		19.4
RF acceptance	[%]	1.06	3.32	0.6	3.06
Energy acceptance (DA)	[%]	±1.0	±1.0	±0.9	-2.8/+2.5
Beam crossing angle at IP θ_x	[mrad]		±15		
Crab waist ratio	[%]	70	55	70	
Beam-beam ξ_x/ξ_y^a		0.0022 / 0.0977	0.013 / 0.12	0.130	
Piwinski angle $(\theta_x \sigma_{z,BS})/\sigma_x^*$		26.6	3.6		
Lifetime (q + BS + lattice)	[sec]	11800	4500	6000	1100
Lifetime (lum) ^b	[sec]	1330	960	600	
Luminosity / IP	[10 ³⁴ /cm ² s]	143	20	7.5	

^aincl. hourglass.

^bonly the energy acceptance is taken into account for the cross section, no beam size effect.

K. Oide, 2024

X-Z instability

(K. Ohmi, 2016)

Higher Q_s

→ stronger low order resonance but more space available between them

Bunches are ~50% shorter (assuming the same σ_δ)

→ stronger beamstrahlung (ξ_y increase is smaller)

→ stronger impact of longitudinal impedance?

Resonant depolarization

Figure of merit $SMI = v_s \sigma_\delta / Q_s \sim 1.3-1.4$ for baseline

→ is reduced to ~0.85 (SMI < 1 is preferred)

Many more aspects to be re-analyzed...

$$\sigma_z \propto \frac{1}{\sqrt{V_{tot}}}$$

3.3/9.4
0.195/0

0.0483
2.35

xx/0.162

$$\xi_y \propto \frac{N_p}{\sigma_z}$$

$$L \propto \xi_y$$

?

Summary

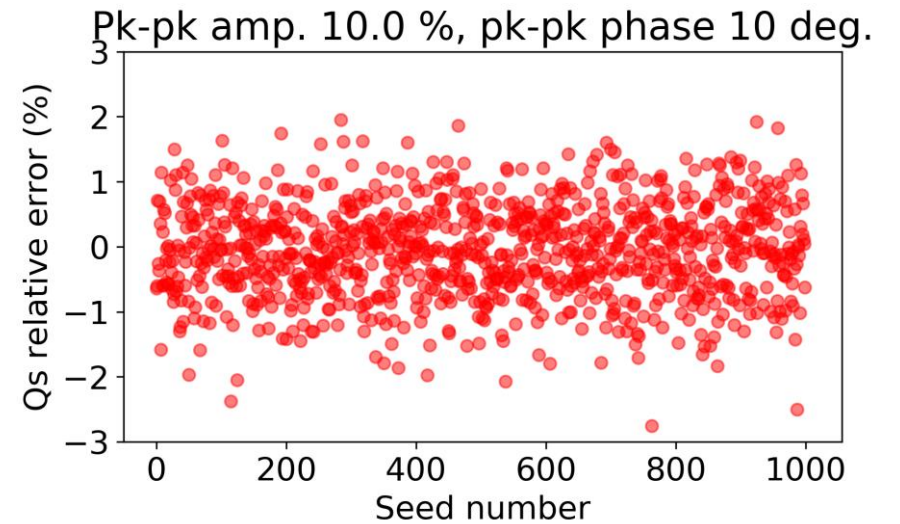
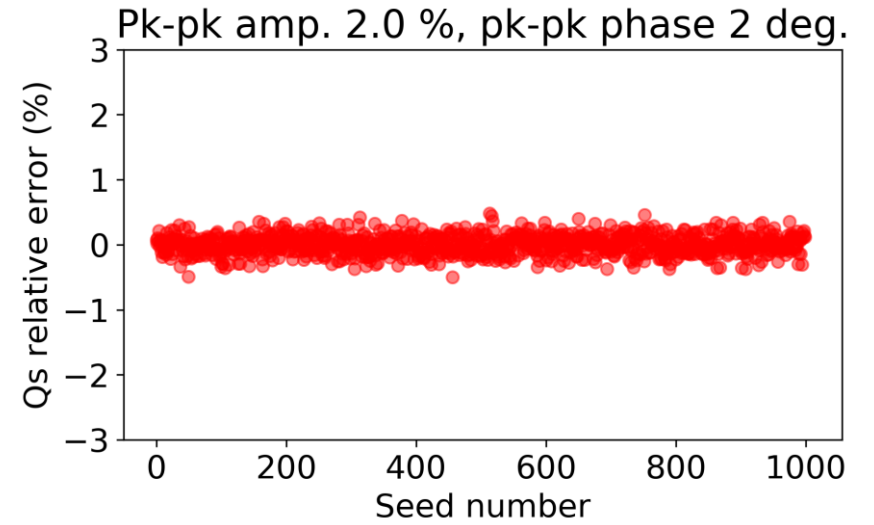
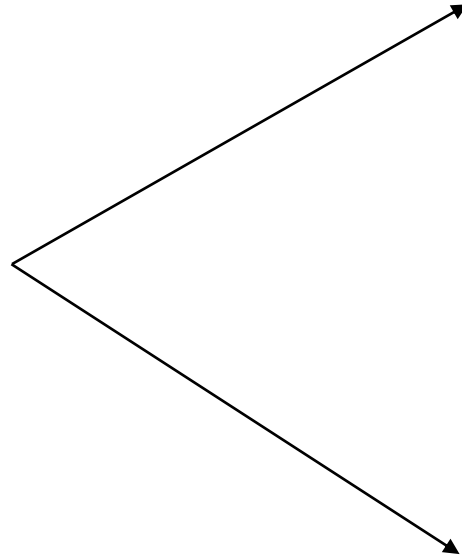
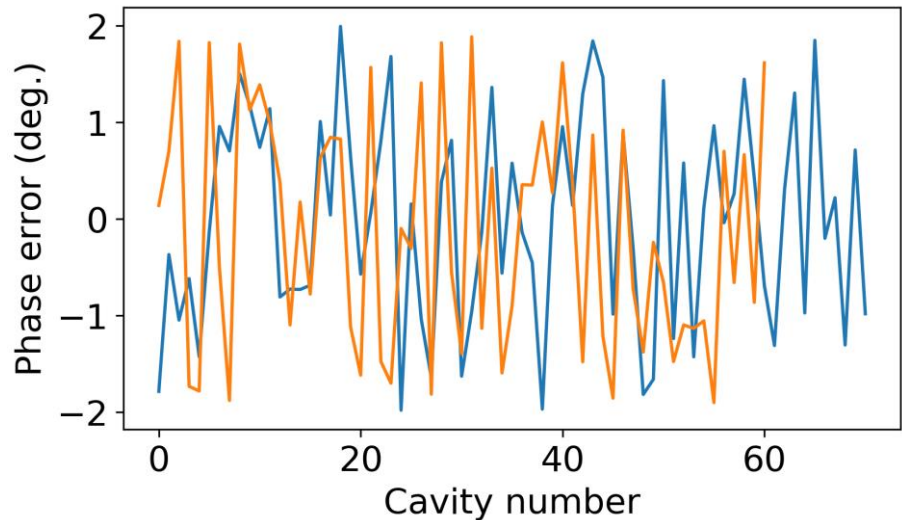
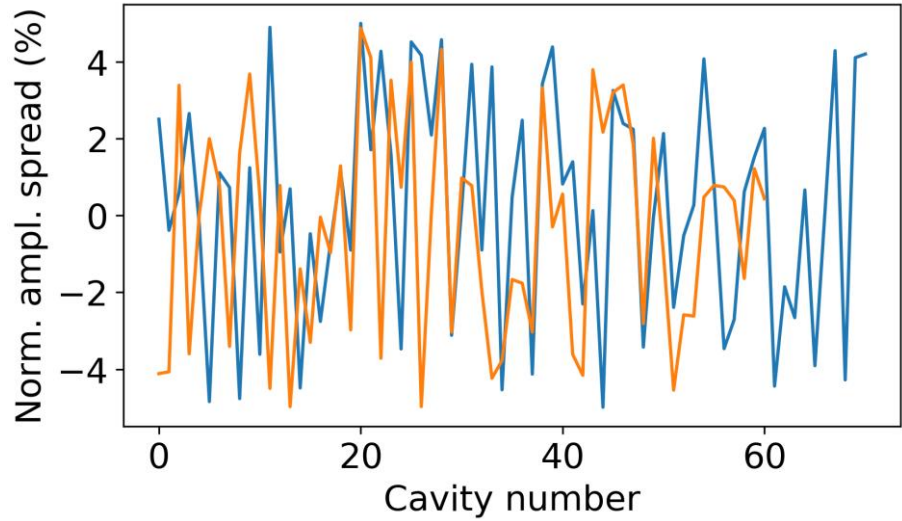
Reverse Phase Operation (RPO) mode aims to avoid hardware modification of RF system between Z, W, and ZH modes

- Synchrotron frequency and bunch length spread due to transient beam loading could be a potential showstopper.
- Possible mitigations are reduction of gap length or increase of total RF voltage
- Would it be possible to find a new parameter set before end Feasibility Study?

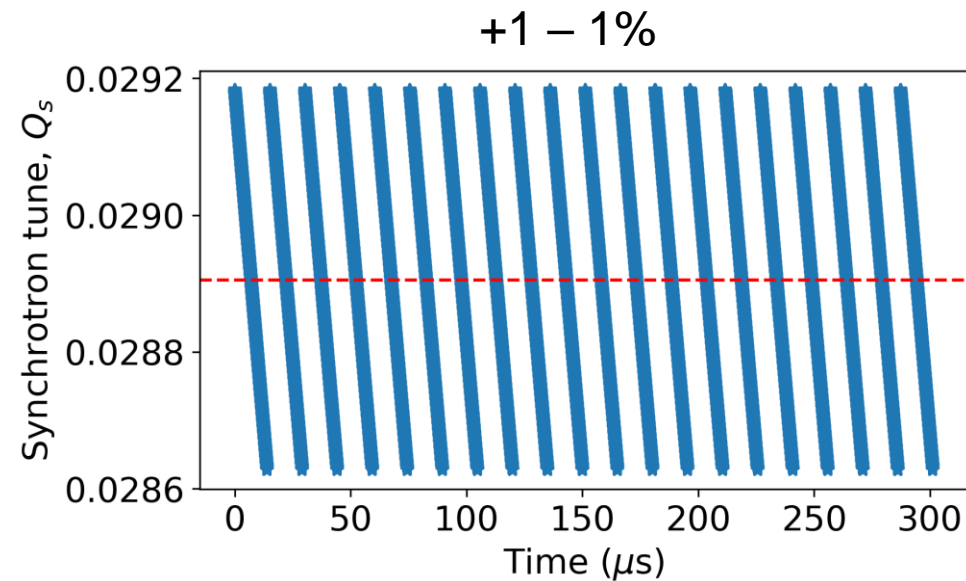
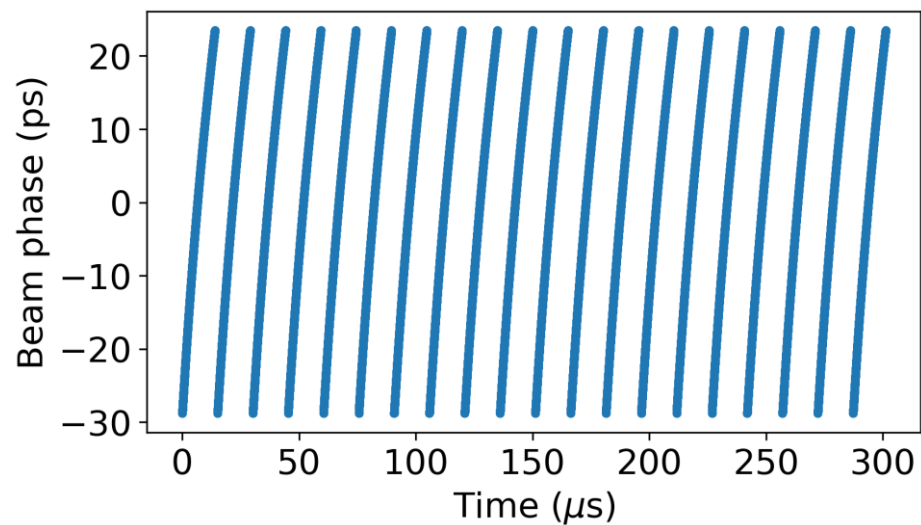
Thank you for your attention!

Backup slides

Parameter sensitivity of RPO

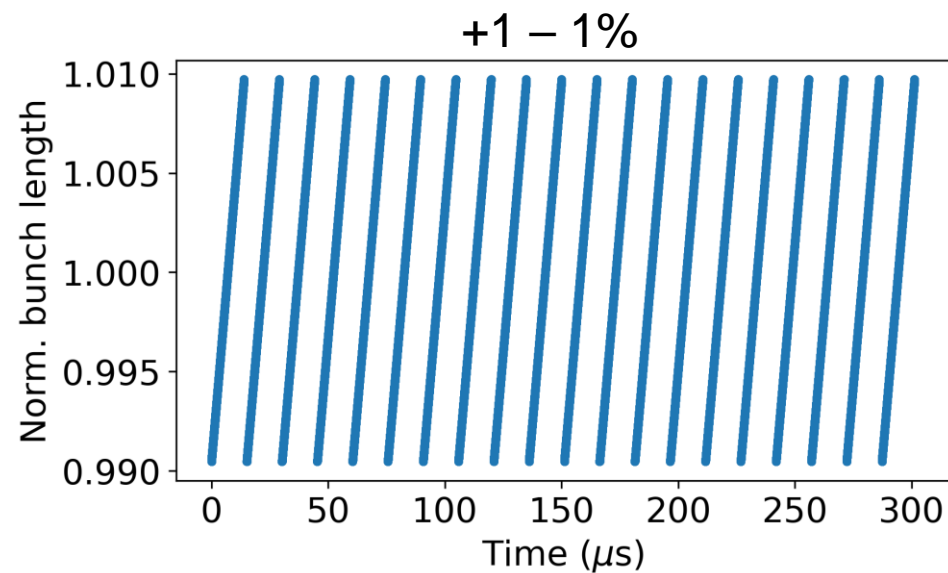


Transients for baseline scenario (56 1-cell cavities)



Q_s spread $\sim 3e4$

Almost negligible spread of bunch-by-bunch parameters



Global parameters (Déjà vu)

August 2024

March 2022

$\beta_x^* = 10 \text{ cm @ Z}$

Beam energy	[GeV]	45.6	
Layout			
# of IPs			
Circumference	[km]		
Bend. radius of arc dipole	[km]		
Energy loss / turn	[GeV]	0.0390	
SR power / beam	[MW]		
Beam current	[mA]	1283	
Colliding bunches / beam		11200	
Colliding bunch population	[10 ¹¹]	2.16	
Hor. emittance at collision ε_x	[nm]	0.70	
Ver. emittance at collision ε_y	[pm]	1.9	
Lattice ver. emittance $\varepsilon_{y,\text{lattice}}$	[pm]	0.87	
Arc cell		Long 90/90	
Momentum compaction α_p	[10 ⁻⁶]	28.67	
Arc sext families		75	
$\beta_{x/y}^*$	[mm]	110 / 0.7	
Transverse tunes $Q_{x/y}$		218.158 / 222.220	218.1
Chromaticities $Q'_{x/y}$		0 / +5	
Energy spread (SR/BS) σ_δ	[%]	0.039 / 0.110	0.0
Bunch length (SR/BS) σ_z	[mm]	5.57 / 15.6	3
RF voltage 400/800 MHz	[GV]	0.079 / 0	
Harm. number for 400 MHz			
RF frequency (400 MHz)	MHz		
Synchrotron tune Q_s		0.0289	
Long. damping time	[turns]	1171	
RF acceptance	[%]	1.06	
Energy acceptance (DA)	[%]	±1.0	
Beam crossing angle at IP θ_x	[mrad]		
Crab waist ratio	[%]	70	
Beam-beam ξ_x/ξ_y^a		0.0022 / 0.0977	0.0
Piwinski angle $(\theta_x\sigma_z, \text{BS})/\sigma_x^*$		26.6	
Lifetime (q + BS + lattice)	[sec]	11800	
Lifetime (lum) ^b	[sec]	1330	
Luminosity / IP	[10 ³⁴ /cm ² s]	143	

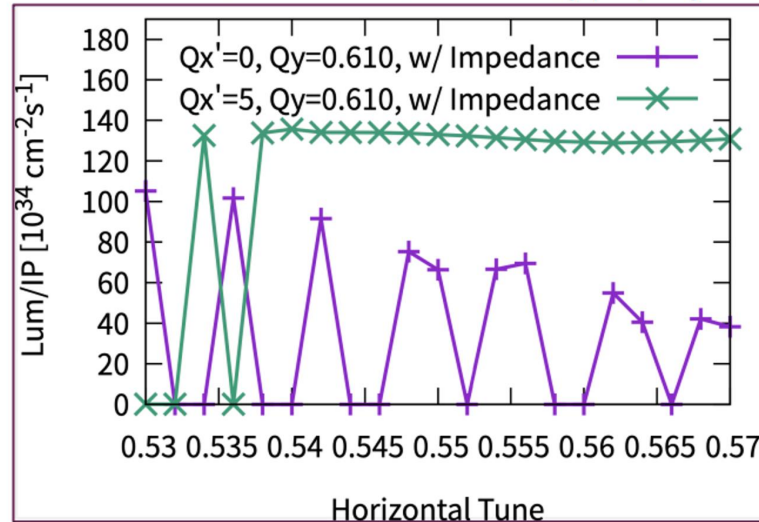
Beam energy	[GeV]	45.6	80
Layout			PA31
# of IPs			4
Circumference	[km]	91.174117	
Bending radius of arc dipole	[km]		9.9
Energy loss / turn	[GeV]	0.0391	0.370
SR power / beam	[MW]		50
Beam current	[mA]	1280	135
Bunches / beam		10000	880
Bunch population	[10 ¹¹]	2.43	2.91
Horizontal emittance ε_x	[nm]	0.71	2.16
Vertical emittance ε_y	[pm]	1.42	4.32
Arc cell		Long 90/90	
Momentum compaction α_p	[10 ⁻⁶]	28.5	
Arc sextupole families		75	
$\beta_{x/y}^*$	[mm]	100 / 0.8	200 / 1.0
Transverse tunes/IP $Q_{x/y}$		53.563 / 53.600	
Energy spread (SR/BS) σ_δ	[%]	0.038 / 0.132	0.069 / 0.154
Bunch length (SR/BS) σ_z	[mm]	4.38 / 15.4	3.55 / 8.01
RF voltage 400/800 MHz	[GV]	0.120 / 0	1.0 / 0
Harmonic number for 400 MHz			1216
RF frequency (400 MHz)	MHz	399.994581	
Synchrotron tune Q_s		0.0370	0.0801
Long. damping time	[turns]	1168	217
RF acceptance	[%]	1.6	3.4
Energy acceptance (DA)	[%]	±1.3	±1.3
Beam-beam ξ_x/ξ_y^a		0.0023 / 0.135	0.011 / 0.125
Luminosity / IP	[10 ³⁴ /cm ² s]	182	19.4
Lifetime (q + BS)	[sec]		
Lifetime (lum)	[sec]	1129	1070

Impact on beam stability (first thoughts)

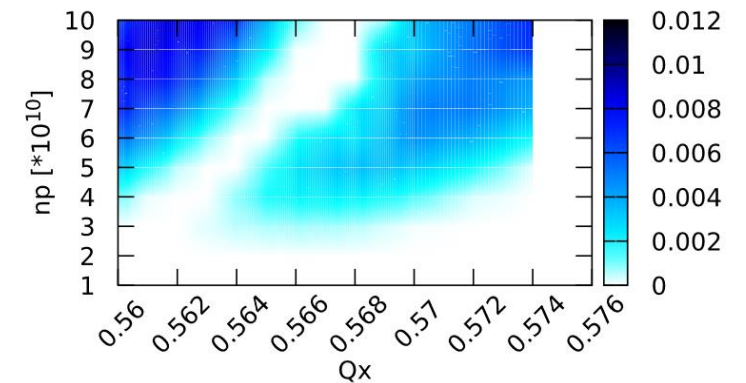
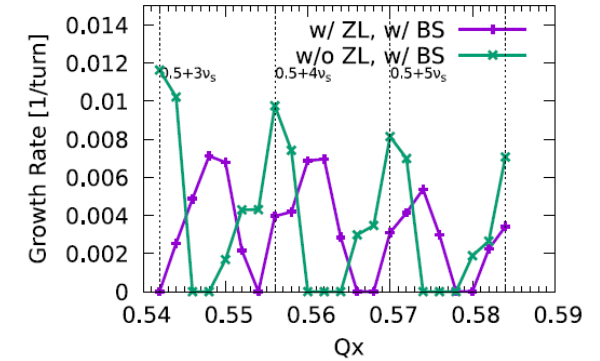
Interplay between beam-beam and coupling impedance

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Y. Zhang, IHEP, Beijing, China



CEPC CDR



X-Z instability

Higher $Q_s \rightarrow$ lower order resonance (stronger)

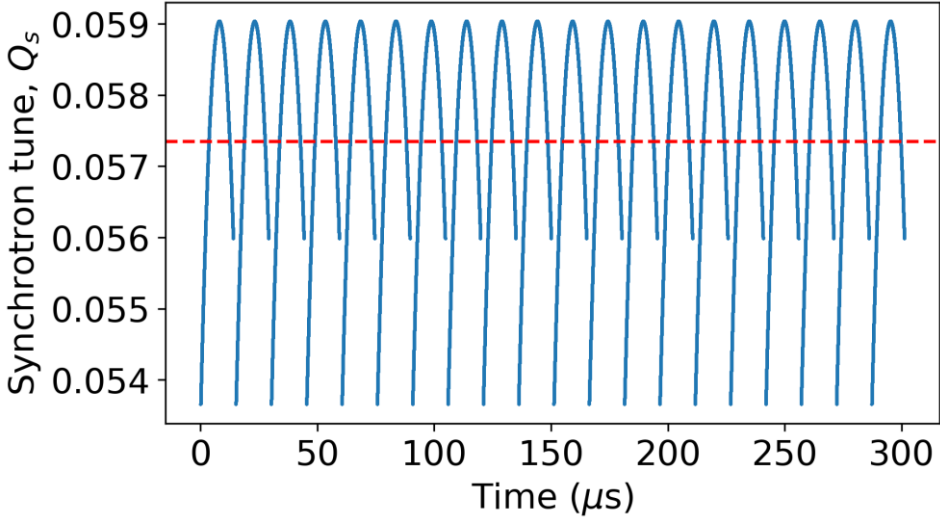
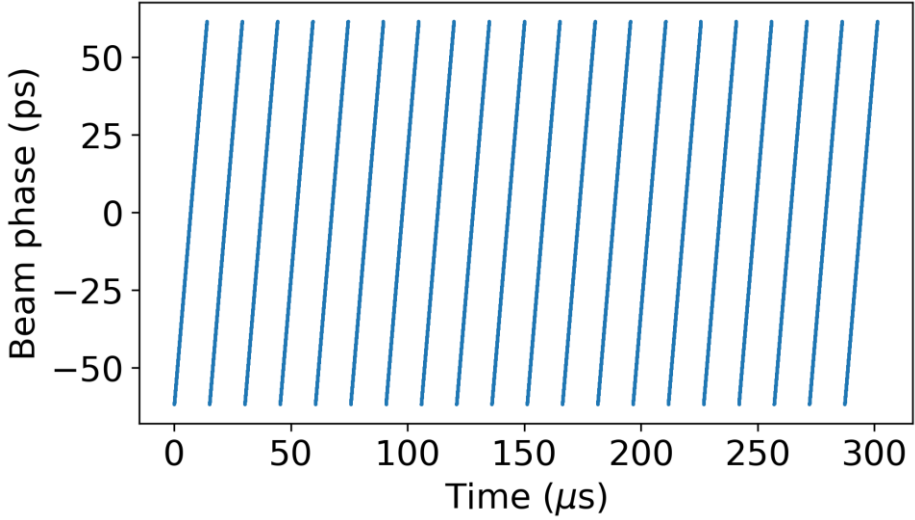
Higher $V_{tot} \rightarrow \sim 50\%$ shorter bunches (assuming the same dp/p) \rightarrow stronger beamstrahlung (ξ_y increase)

Stabilizing role of chromaticity is not fully understood

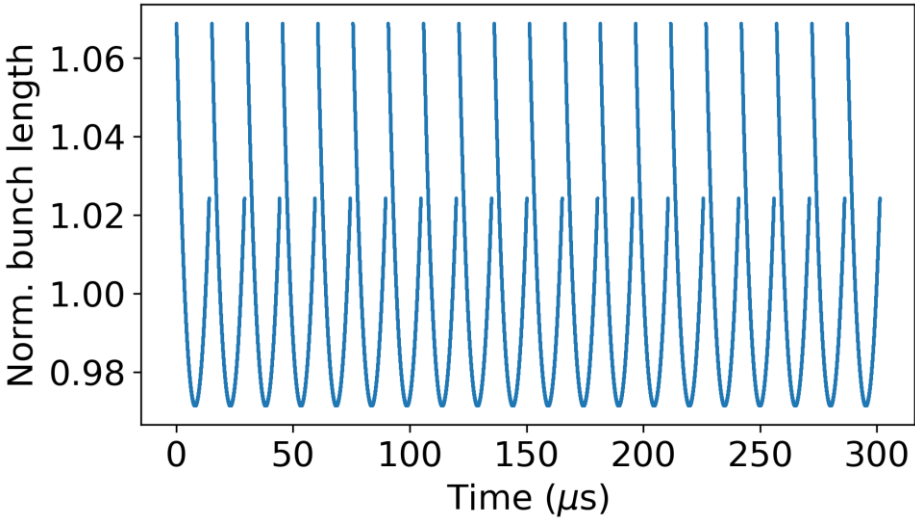
Resonant depolarization

Figure of merit $SMI = \nu_s * (\sigma_E/E) / Q_s \sim 1.3-1.4$ for baseline \rightarrow reduces ~ 0.85 ($SMI < 1$ is preferred)

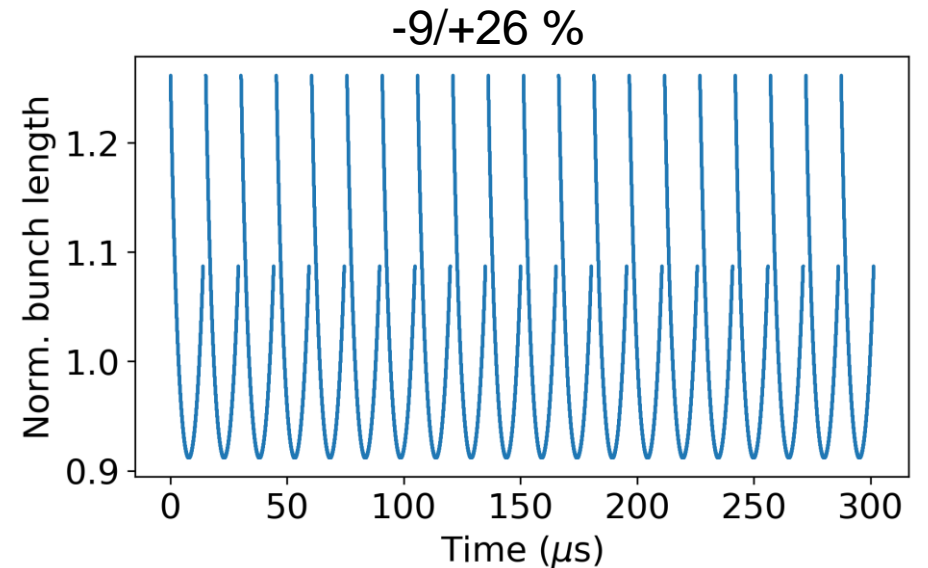
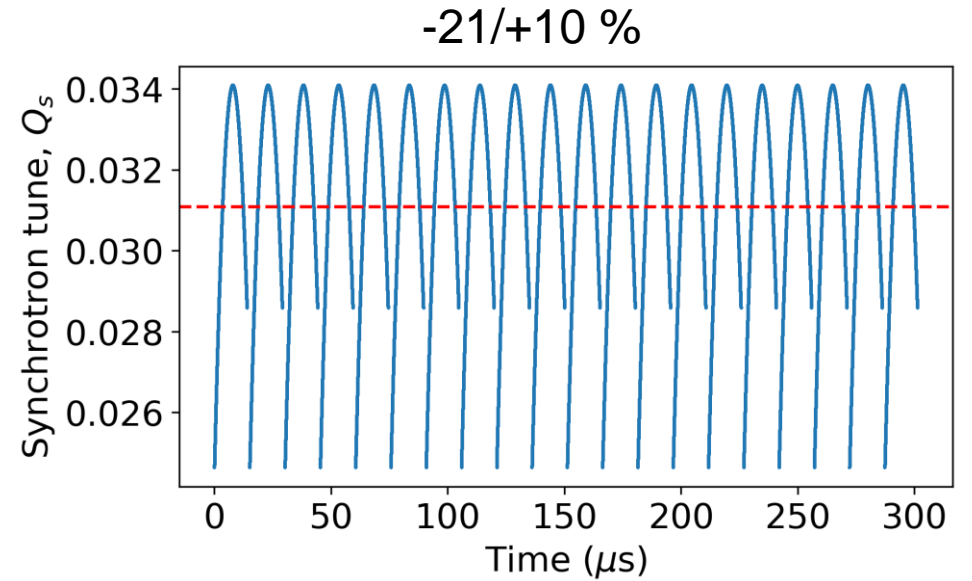
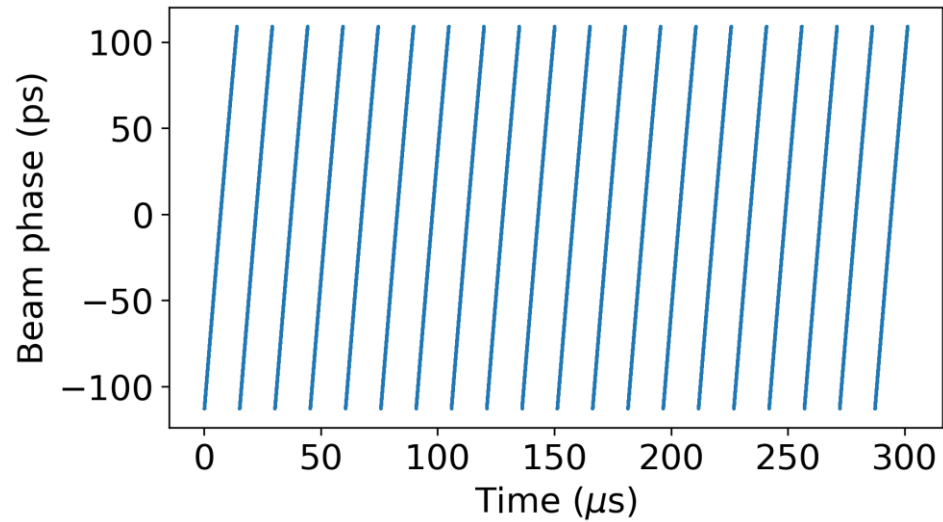
Changing optics



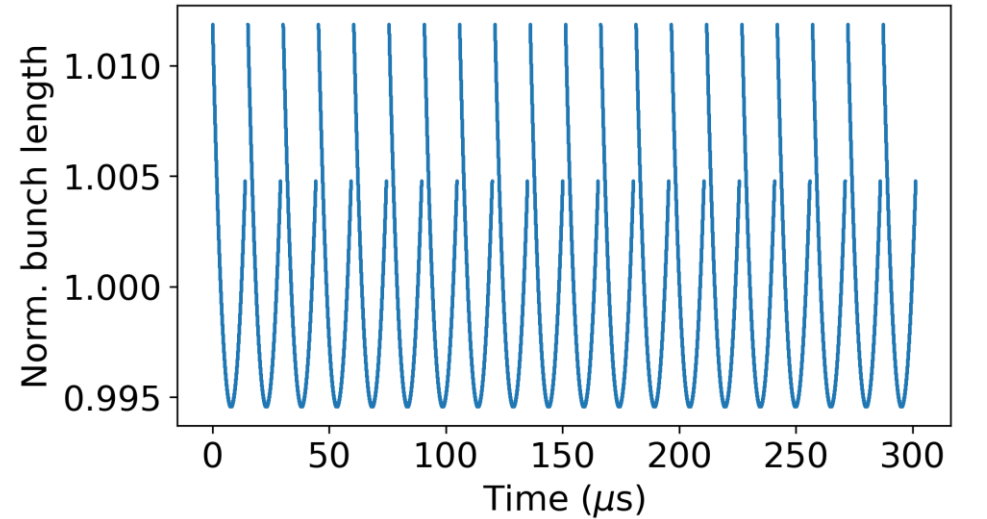
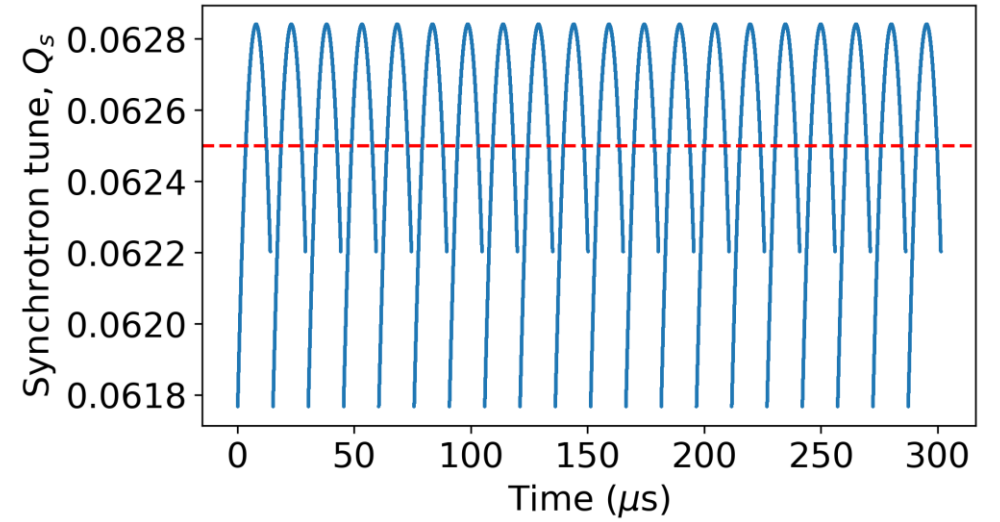
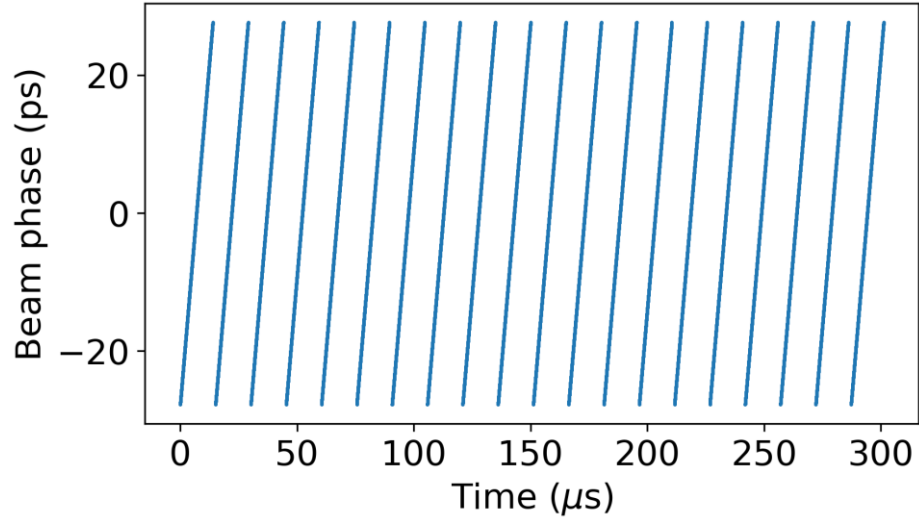
New $\alpha_p = 53.7e-6$



Reduced $Q_L=1.5e5$ for same V_{tot}



320 MV



Motivation

RF power for SRF cavities with circulators is minimized for optimal parameters:

$$\text{Optimal detuning } \Delta\omega_{\text{opt}} = -\frac{\omega_{\text{rf}}(R/Q)|F_b|I_{b,\text{dc}} \sin \phi_s}{2V_{\text{cav}}}$$

$$\text{Optimal quality factor } Q_{\text{ext,opt}} = \frac{V_{\text{cav}}}{|F_b|(R/Q)I_{b,\text{dc}} \cos \phi_s}$$

$$F_b = 2 \frac{\mathcal{F}[\lambda(t)]_{\omega=\omega_{\text{rf}}}}{\mathcal{F}[\lambda(t)]_{\omega=0}}$$

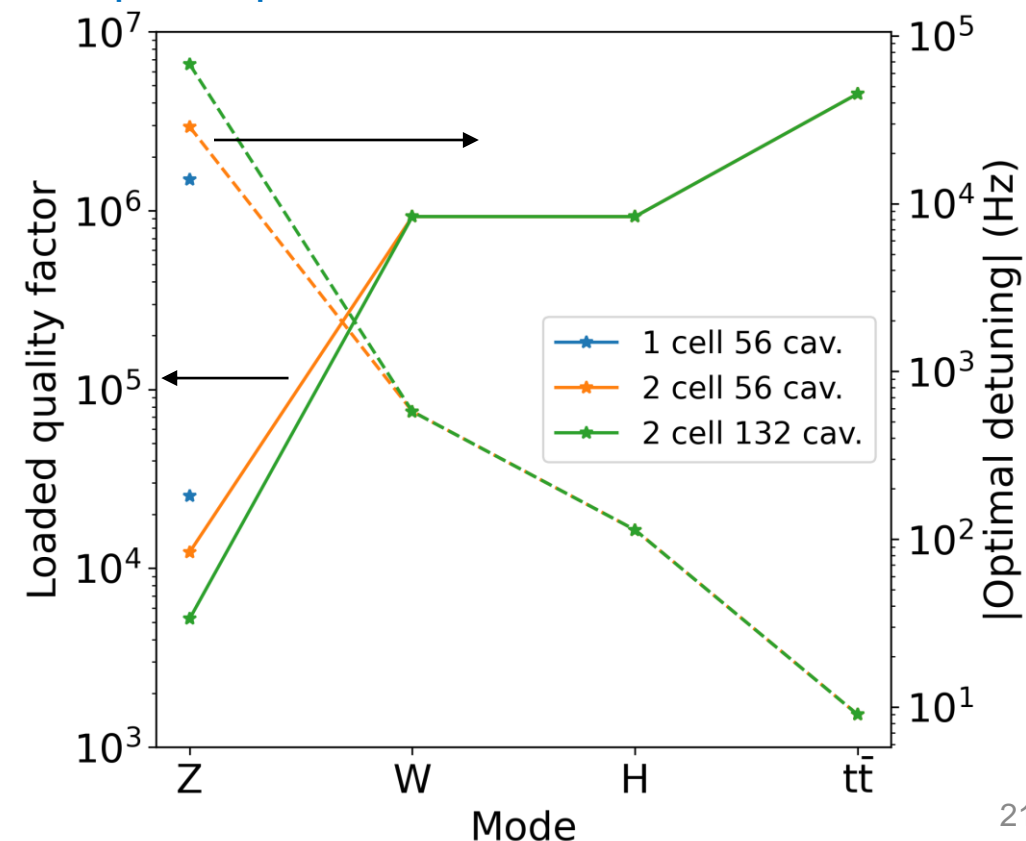
Keeping 2-cell cavities for Z, W, H, (and $t\bar{t}$):

→ Large range for $Q_{\text{ext,opt}}$ adjustment (a factor of ~75-600) starting from $\sim 5 \times 10^3$: possible FPC solutions was studied (*S. Gorgi Zadeh and E. Montesinos, CERN SRF, 2024; see also slides of F. Gerigk, FCC Week 2024*)

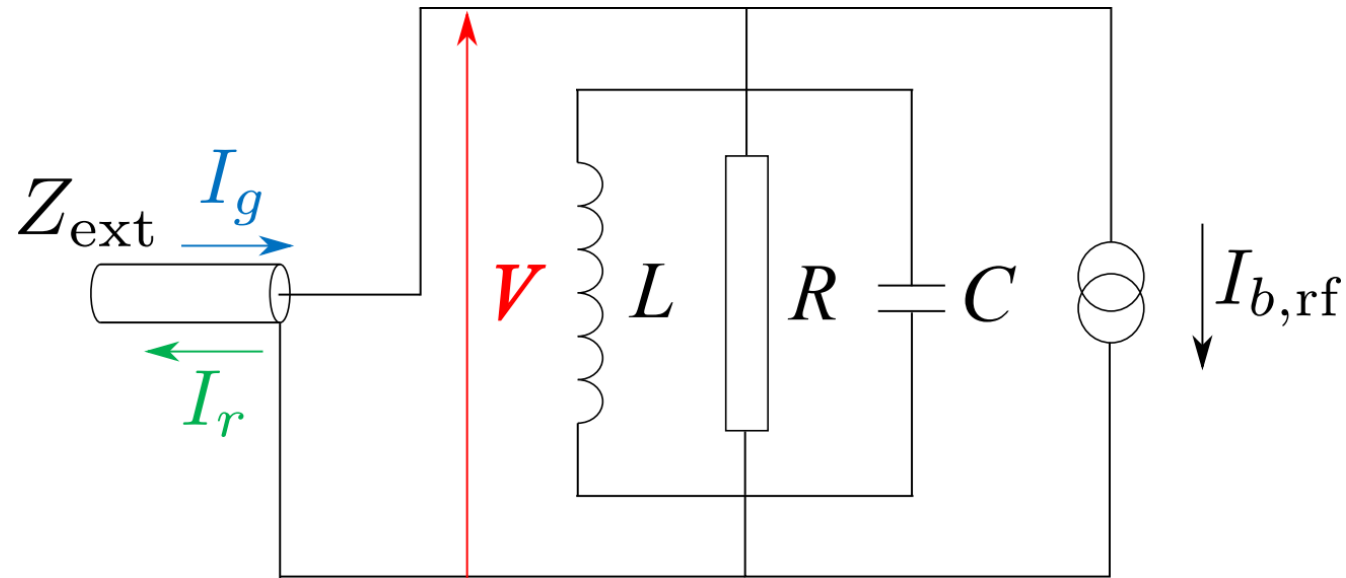
→ Increased detuning enhances instability due to fundamental mode

Can the voltage per cavity be increased for Z mode?

Optimal parameters for different scenarios



Beam loading model: main equation



Generator current

$$I_g = \frac{V}{2(R/Q)} \left(\frac{1}{Q_0} + \frac{1}{Q_{\text{ext}}} - 2i \frac{\Delta\omega}{\omega_{\text{rf}}} \right) + \frac{I_{b,\text{rf}}}{2}$$

Generator power

$$P_g = \frac{1}{2} Z_{\text{ext}} |I_g|^2 = \frac{1}{2} (R/Q) Q_{\text{ext}} |I_g|^2$$

Fixed parameters are V , (R/Q) , Q_0 , ω_{rf} , $I_{b,\text{rf}}$, while V , $\Delta\omega$, and Q_{ext} can be adjusted

RF power requirements

Constraints:

- The same $Q_{\text{ext,opt}}$ for all cavities to avoid a movable fundamental power coupler design
- The same $P_{g,\text{opt}}$ to have the identical power sources and uniform power distribution (role of variations is under study)

$$Q_{\text{ext,opt}} = \frac{|V_{\text{cav}}|}{|F_b|(R/Q)I_{b,\text{dc}} \cos(\phi_s + \phi_c)}$$

$$P_{g,\text{opt}} = \frac{|V_{\text{cav}}||F_b|I_{b,\text{dc}} \cos(\phi_s + \phi_c)}{2}$$

→ Cavity voltage must be the same for all cavities: $\cos(\phi_s + \phi_{\text{foc}}) = \cos(\phi_s + \phi_{\text{defoc}}) \rightarrow \phi_{\text{foc}} = -2\phi_s - \phi_{\text{defoc}}$

Starting with energy gain per turn

$$N_{\text{foc}}|V_{\text{cav}}| \cos(\phi_s + \phi_{\text{foc}}) + N_{\text{defoc}}|V_{\text{cav}}| \cos(\phi_s + \phi_{\text{defoc}}) = V_{\text{tot}} \cos \phi_s$$

$$N_{\text{foc}} \frac{|V_{\text{cav}}||F_b|I_{b,\text{dc}} \cos(\phi_s + \phi_{\text{foc}})}{2} + N_{\text{defoc}} \frac{|V_{\text{cav}}||F_b|I_{b,\text{dc}} \cos(\phi_s + \phi_{\text{defoc}})}{2} = \frac{|F_b|I_{b,\text{dc}}}{2} V_{\text{tot}} \cos \phi_s$$

$$N_{\text{foc}}P_{g,\text{foc}} + N_{\text{defoc}}P_{g,\text{defoc}} = I_{b,\text{dc}}U_0 = P_{\text{SR}}$$

$$P_{g,\text{opt}} = \frac{P_{\text{SR}}}{N_{\text{tot}}}$$

$$\times \frac{|F_b|I_{b,\text{dc}}}{2}$$

$$\cos \phi_s = \frac{U_0}{V_{\text{tot}}} \quad |F_b| \approx 2$$

$$P_{g,\text{foc}} = P_{g,\text{defoc}} = P_{g,\text{opt}}$$

$$N_{\text{foc}} + N_{\text{defoc}} = N_{\text{tot}}$$

→ No RF power overshoot is needed for RPO if optimal detuning and optimal quality factor are used

Reverse phasing mode equations

Preservation of energy gain

$$N_{\text{foc}}|V_{\text{cav}}| \cos(\phi_s + \phi_{\text{foc}}) + N_{\text{defoc}}|V_{\text{cav}}| \cos(\phi_s + \phi_{\text{defoc}}) = V_{\text{tot}} \cos \phi_s$$

Preservation of synchrotron tune

$$N_{\text{foc}}|V_{\text{cav}}| \sin(\phi_s + \phi_{\text{foc}}) + N_{\text{defoc}}|V_{\text{cav}}| \sin(\phi_s + \phi_{\text{defoc}}) = V_{\text{tot}} \sin \phi_s$$

→ Cavity voltage

$$|V_{\text{cav}}| = \frac{V_{\text{tot}}}{N_{\text{tot}}} \sqrt{\frac{U_0^2}{V_{\text{tot}}^2} + \left(1 - \frac{U_0^2}{V_{\text{tot}}^2}\right) \frac{N_{\text{tot}}^2}{(N_{\text{foc}} - N_{\text{defoc}})^2}}$$

Optimal detuning

$$\Delta\omega_{\text{opt}} = -\frac{\omega_{\text{rf}}(R/Q)|F_b|I_{b,\text{dc}}}{2V_{\text{cav}}} \sqrt{1 - \frac{U_0^2}{V_{\text{cav}}^2 N_{\text{tot}}^2}}$$

See, also [A. Blednykh et al, EIC-ADD-TN-33, 2022](#)

Phases

$$\phi_{\text{foc}} = -\phi_s + \arccos\left(\frac{V_{\text{tot}} \cos \phi_s}{N_{\text{tot}} V_{\text{cav}}}\right) \quad \phi_{\text{defoc}} = -\phi_s - \arccos\left(\frac{V_{\text{tot}} \cos \phi_s}{N_{\text{tot}} V_{\text{cav}}}\right)$$

The aim is to keep V_{cav} , $P_{g,\text{opt}}$, and $Q_{\text{ext,opt}}$ for Z, W, and ZH modes

→ Cavity voltage can be change in discrete steps of $N_{\text{foc}} - N_{\text{defoc}} = 2, 4, \dots$

Derivations for arbitrary cavity phase (1/2)

Generator current $I_g = \frac{V}{2(R/Q)} \left(\frac{1}{Q_0} + \frac{1}{Q_{\text{ext}}} - 2i \frac{\Delta\omega}{\omega_{\text{rf}}} \right) + \frac{I_{b,\text{rf}}}{2}$

$F_b = 2 \frac{\mathcal{F}[\lambda(t)]_{\omega=\omega_{\text{rf}}}}{\mathcal{F}[\lambda(t)]_{\omega=0}}$

Complex quantities: I_g , V , and $I_{b,\text{rf}}$ → $I_g = |I_g|e^{i\phi_L}$, $V = |V_{\text{cav}}|e^{i\phi_c}$, $I_{b,\text{rf}} = |F_b|I_{b,\text{dc}}e^{-i\phi_s}$

$$|I_g|e^{i\phi_L} = \frac{|V_{\text{cav}}|e^{i\phi_c}}{2(R/Q)} \left(\frac{1}{Q_{\text{ext}}} - 2i \frac{\Delta\omega}{\omega_{\text{rf}}} \right) + \frac{|F_b|I_{b,\text{dc}}e^{-i\phi_s}}{2} \quad \left| \quad \times e^{-i\phi_c} \right.$$

$$|I_g|e^{i\phi_L - i\phi_c} = \frac{|V_{\text{cav}}|}{2(R/Q)} \left(\frac{1}{Q_{\text{ext}}} - 2i \frac{\Delta\omega}{\omega_{\text{rf}}} \right) + \frac{|F_b|I_{b,\text{dc}}e^{-i\phi_s - i\phi_c}}{2}$$

Then splitting in real and imaginary parts:

Derivations for arbitrary cavity phase (2/2)

$$|I_g| e^{i\phi_L - i\phi_c} = \frac{|V_{\text{cav}}|}{2(R/Q)Q_{\text{ext}}} + \frac{|F_b|I_{b,\text{dc}} \cos(\phi_s + \phi_c)}{2} - i \left[\frac{|V_{\text{cav}}|}{(R/Q)} \frac{\Delta\omega}{\omega_{\text{rf}}} + \frac{|F_b|I_{b,\text{dc}} \sin(\phi_s + \phi_c)}{2} \right]$$

$$P_g = \frac{1}{2} (R/Q) Q_{\text{ext}} |I_g|^2$$

$$= \frac{1}{2} (R/Q) Q_{\text{ext}} \left[\frac{|V_{\text{cav}}|}{2(R/Q)Q_{\text{ext}}} + \frac{|F_b|I_{b,\text{dc}} \cos(\phi_s + \phi_c)}{2} \right]^2 + \frac{1}{2} (R/Q) Q_{\text{ext}} \left[\frac{|V_{\text{cav}}|}{(R/Q)} \frac{\Delta\omega}{\omega_{\text{rf}}} + \frac{|F_b|I_{b,\text{dc}} \sin(\phi_s + \phi_c)}{2} \right]^2$$

Minimized for $Q_{\text{ext,opt}} = \frac{|V_{\text{cav}}|}{|F_b|(R/Q)I_{b,\text{dc}} \cos(\phi_s + \phi_c)}$

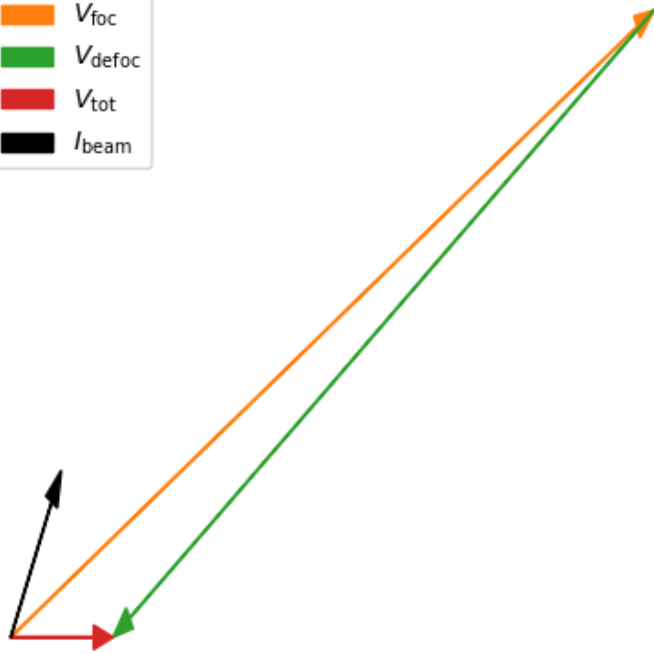
= 0 for $\Delta\omega_{\text{opt}} = -\frac{\omega_{\text{rf}}(R/Q)|F_b|I_{b,\text{dc}} \sin(\phi_s + \phi_c)}{2|V_{\text{cav}}|}$

Setting $\phi_c = 0$ recovers classical equations for optimal parameters

Adjusting ϕ_c , $Q_{\text{ext,opt}}$ can be modified to meet certain constraints

The minimum power

$$P_{g,\text{opt}} = \frac{|V_{\text{cav}}||F_b|I_{b,\text{dc}} \cos(\phi_s + \phi_c)}{2}$$

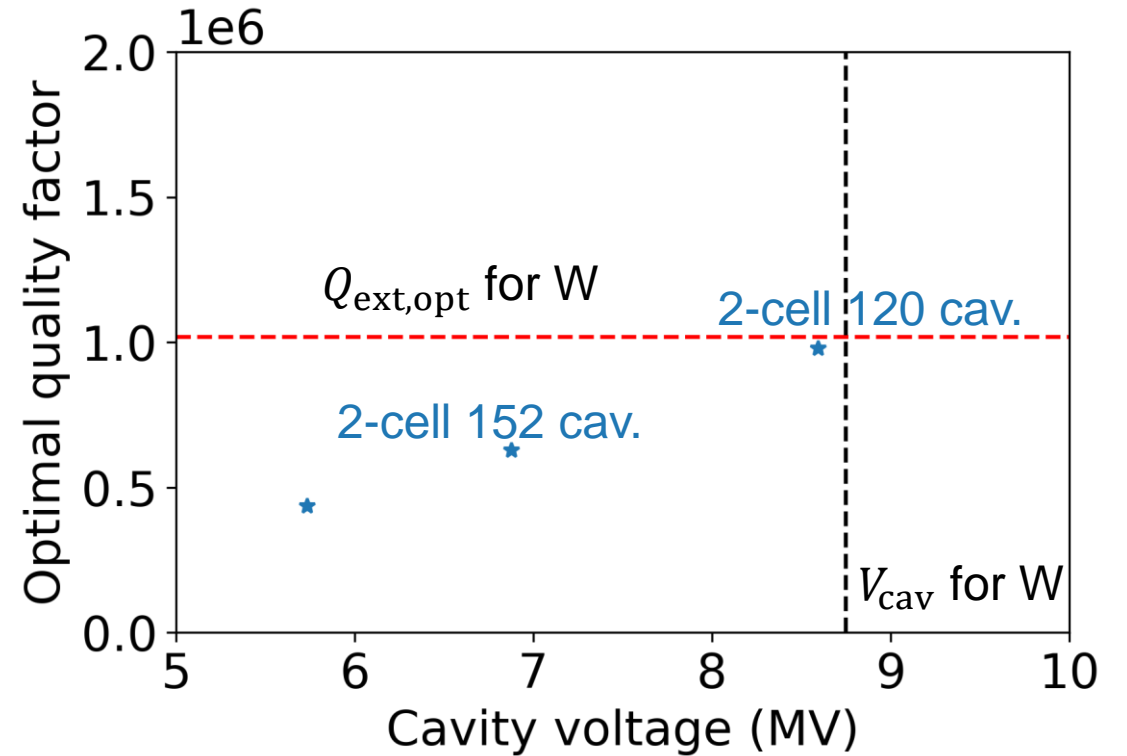


Preliminary results

$$|V_{\text{cav}}| = \frac{V_{\text{tot}}}{N_{\text{tot}}} \sqrt{\frac{U_0^2}{V_{\text{tot}}^2} + \left(1 - \frac{U_0^2}{V_{\text{tot}}^2}\right) \frac{N_{\text{tot}}^2}{(N_{\text{foc}} - N_{\text{defoc}})^2}}$$

Optimal quality factor with RPO

$$Q_{\text{ext,opt}} = \frac{V_{\text{cav}}^2 N_{\text{tot}}}{V_{\text{tot}} (R/Q) |F_b| I_{b,\text{dc}} \cos \phi_s}$$



→ RPO is under evaluation potentially allowing for the same optimal quality factor for Z, W, and H modes

Reverse phasing mode equations

Constraints: $|V_{\text{cav}}|$ and $P_{g,\text{opt}}$ are the same for focusing and defocusing cavities
 $\rightarrow \cos(\phi_s + \phi_{\text{foc}}) = \cos(\phi_s + \phi_{\text{defoc}}) \rightarrow \phi_{\text{foc}} = -2\phi_s - \phi_{\text{defoc}}$

$$P_{g,\text{opt}} = \frac{|V_{\text{cav}}| |F_b| I_{b,\text{dc}} \cos(\phi_s + \phi_c)}{2}$$

Preservation of energy gain

$$N_{\text{foc}} |V_{\text{cav}}| \cos(\phi_s + \phi_{\text{foc}}) + N_{\text{defoc}} |V_{\text{cav}}| \cos(\phi_s + \phi_{\text{defoc}}) = V_{\text{tot}} \cos \phi_s$$

Preservation of synchrotron tune

$$N_{\text{foc}} |V_{\text{cav}}| \sin(\phi_s + \phi_{\text{foc}}) + N_{\text{defoc}} |V_{\text{cav}}| \sin(\phi_s + \phi_{\text{defoc}}) = V_{\text{tot}} \sin \phi_s$$

RPO

Classical

Optimal quality factor

$$Q_{\text{ext,opt}} = \frac{V_{\text{cav}}^2 N_{\text{tot}}}{V_{\text{tot}} (R/Q) |F_b| I_{b,\text{dc}} \cos \phi_s}$$

$$Q_{\text{ext,opt}} = \frac{V_{\text{cav}}}{|F_b| (R/Q) I_{b,\text{dc}} \cos \phi_s}$$

Optimal detuning

$$\Delta\omega_{\text{opt}} = -\frac{\omega_{\text{rf}} (R/Q) |F_b| I_{b,\text{dc}}}{2V_{\text{cav}}} \sqrt{1 - \frac{\cos^2 \phi_s V_{\text{tot}}^2}{V_{\text{cav}}^2 N_{\text{tot}}^2}}$$

$$\Delta\omega_{\text{opt}} = -\frac{\omega_{\text{rf}} (R/Q) |F_b| I_{b,\text{dc}} \sin \phi_s}{2V_{\text{cav}}}$$

Reduced Pedersen model

