



XIV International Conference  
on New Frontiers in Physics  
17-31 July 2025, OAC, Kolymbari, Crete, Greece

# Recent results on CPV from $b$ -hadron to charmonium decays at LHCb

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XIV ICNFP

30 JUL 2025

# Outline

- Introduction
  - The CKM Matrix
  - Penguin pollution
  - Constraining penguin using  $b \rightarrow c\bar{c}d$
- Branching fraction:  $B^0 \rightarrow J/\psi\pi^0$  [JHEP 05 (2024) 065]
- CPV and branching fraction:  $B_s^0 \rightarrow J/\psi\bar{K}^*(892)^0$  [arXiv: 2506.22090]
- **Evidence of direct CPV:**  $B^+ \rightarrow J/\psi\pi^+$  [PRL 134, 101801 (2025)]
- **Evidence of direct CPV:**  $\Lambda_b^0 \rightarrow J/\psi p\pi^-$  [in preparation]

# CKM Matrix

- The transformation between mass & weak-interaction eigenstates:

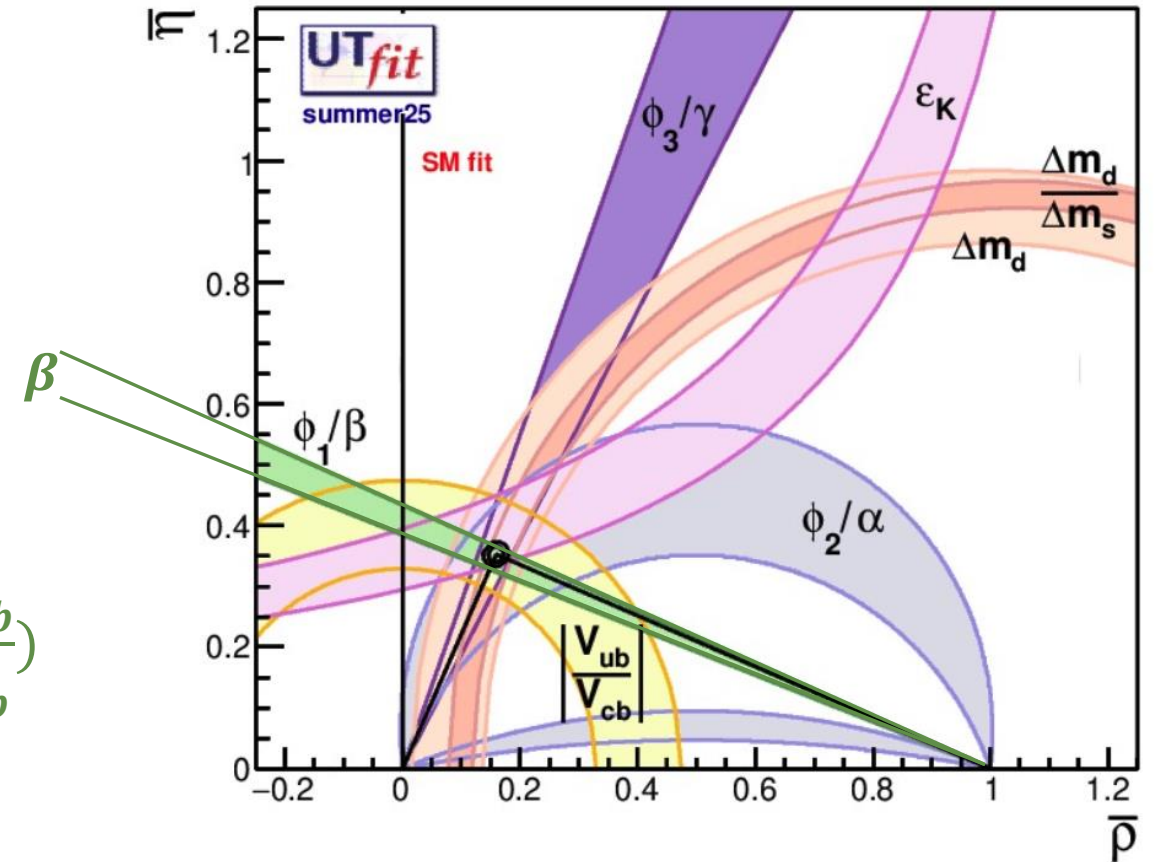
$$-\mathcal{L}_{W^\pm} = \frac{g}{\sqrt{2}} \bar{u}_{Li} \gamma^\mu (V_{CKM})_{ij} d_{Lj} W_\mu^\pm + \text{h.c.}$$

$$V_{CKM} = \begin{pmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{pmatrix}$$

- Unitary triangle angles:

$$\alpha \equiv \arg\left(-\frac{V_{td}V_{tb}^*}{V_{ud}V_{ub}^*}\right), \quad \beta \equiv \arg\left(-\frac{V_{cd}V_{cb}^*}{V_{td}V_{tb}^*}\right)$$

$$\gamma \equiv \arg\left(-\frac{V_{ud}V_{ub}^*}{V_{cd}V_{cb}^*}\right)$$



# Precision test of the Standard Model: CKM angle $\beta$

- Measure  $\beta$  using  $B_d^0$  decays to CP eigenstates with  $b \rightarrow c\bar{c}s$  quark transition
  - Benchmark channel:  $B^0 \rightarrow J/\psi K^0$
  - Amplitude:  $A(B_d^0 \rightarrow f) = A_{tree} + A_{penguin} \equiv A_{tree} \times (1 - b_f e^{i\rho_f} e^{i\gamma})$

- When measuring the time-dependent CP asymmetry:

$$\begin{aligned} \mathcal{A}_{CP}(t) &\equiv \frac{|A(B_d^0(t) \rightarrow f)|^2 - |A(\bar{B}_d^0(t) \rightarrow f)|^2}{|A(B_d^0(t) \rightarrow f)|^2 + |A(\bar{B}_d^0(t) \rightarrow f)|^2} \\ &= \frac{\mathcal{A}_{CP}^{dir} \cos(\Delta mt) + \mathcal{A}_{CP}^{mix} \sin(\Delta mt)}{\cosh(\Delta\Gamma t/2) + \mathcal{A}_{\Delta\Gamma} \sinh(\Delta\Gamma t/2)} \end{aligned}$$

- Assuming no penguin contribution:  $b_f = 0 \Rightarrow \mathcal{A}_{CP}^{mix} = \eta_f \sin(2\beta + \phi^{NP})$
- Penguin pollution:  $b_f \neq 0 \Rightarrow \mathcal{A}_{CP}^{mix} = \eta_f \sin(2\beta + \phi^{NP} + \Delta\phi_f)$  ( $\eta_f = \pm 1$ , CP eigenvalue)
- **The penguin contribution  $b_f$  is hard to calculate theoretically!** Needs experimental methods

# Enhanced penguin diagram in $b \rightarrow c\bar{c}d$ processes

Wolfenstein parameter:  $\lambda \sim 0.23$

$b \rightarrow c\bar{c}s$ :

$$A_f \sim (V_{cs}^* V_{cb})T + (V_{us}^* V_{ub})P^u$$

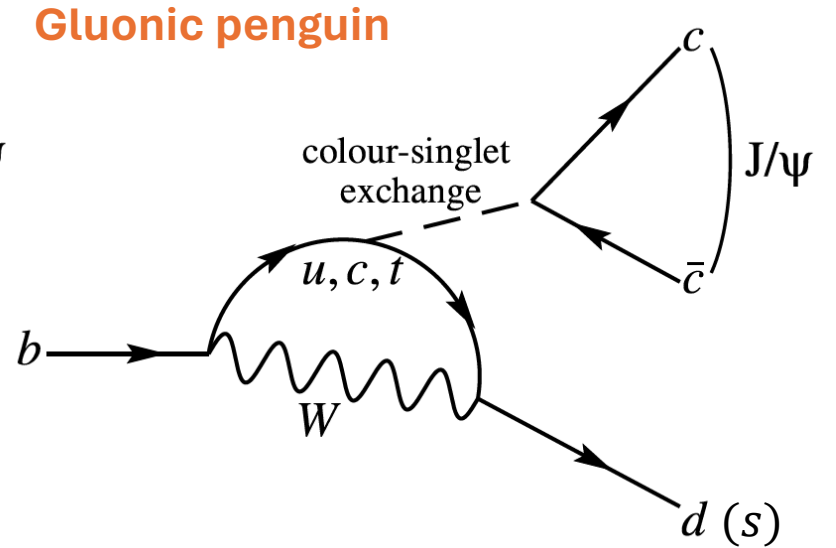
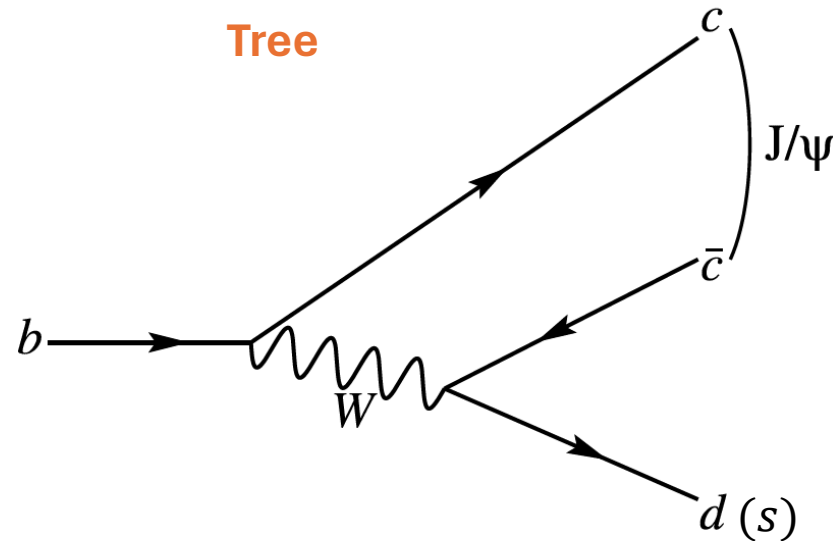
$$A\lambda^2 \quad (\rho - i\eta)A\lambda^4$$

$b \rightarrow c\bar{c}d$ :

$$A_f \sim (V_{cd}^* V_{cb})T + (V_{td}^* V_{tb})P^t$$

$$-A\lambda^3 \quad (1 + \rho + i\eta)A\lambda^3$$

**Enhanced penguin contribution!**



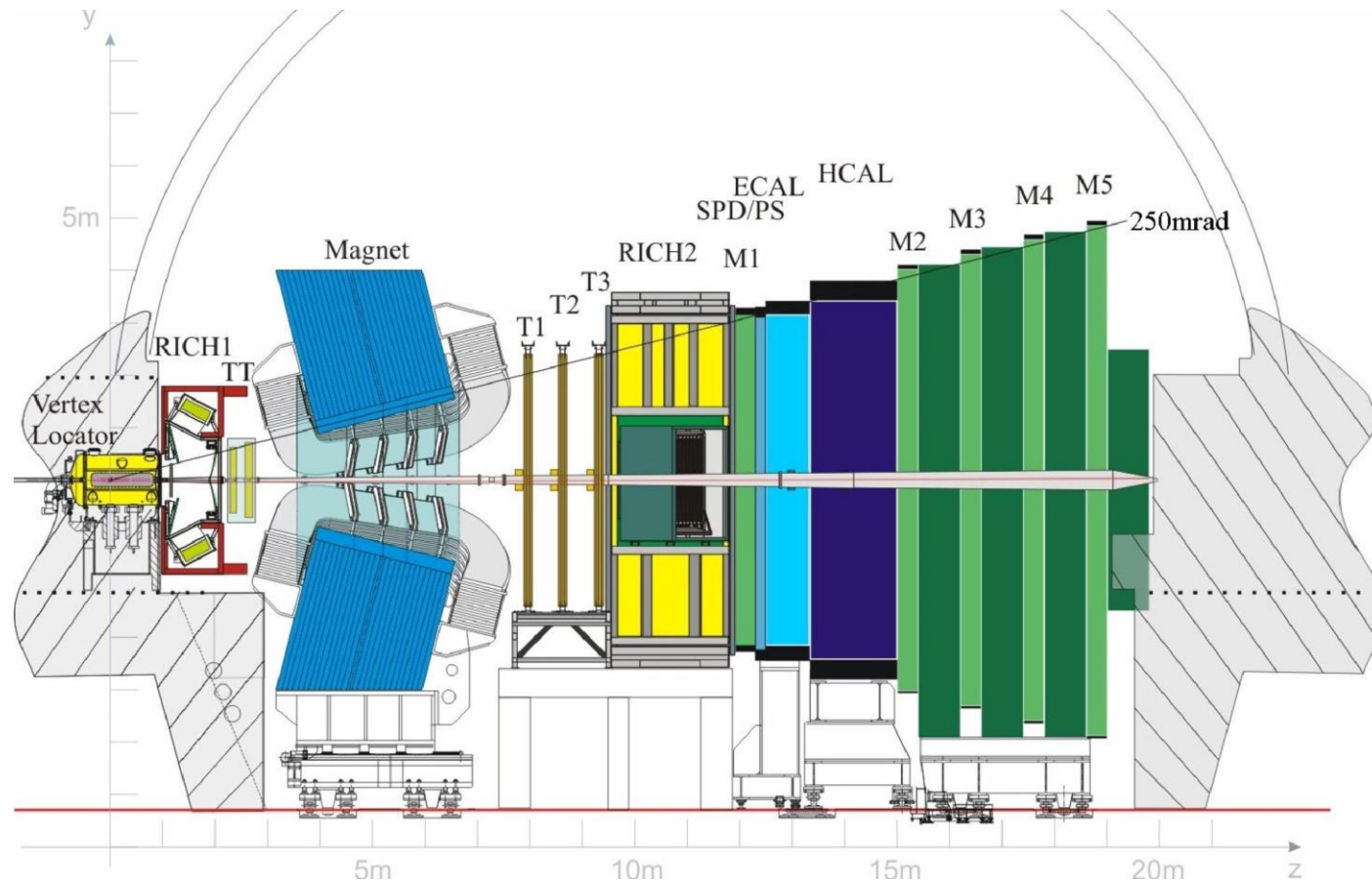
Two experimental observables in  $b \rightarrow c\bar{c}d$  to constrain penguin in  $b \rightarrow c\bar{c}s$

- Branching fraction
- CP asymmetry from decay

**$SU(3)$  symmetry assumed**

# The LHCb Detector

- The LHCb detector is a single-arm forward spectrometer covering the pseudorapidity range  $2 < \eta < 5$ , designed for the study of particles containing  $b$  or  $c$  quarks.



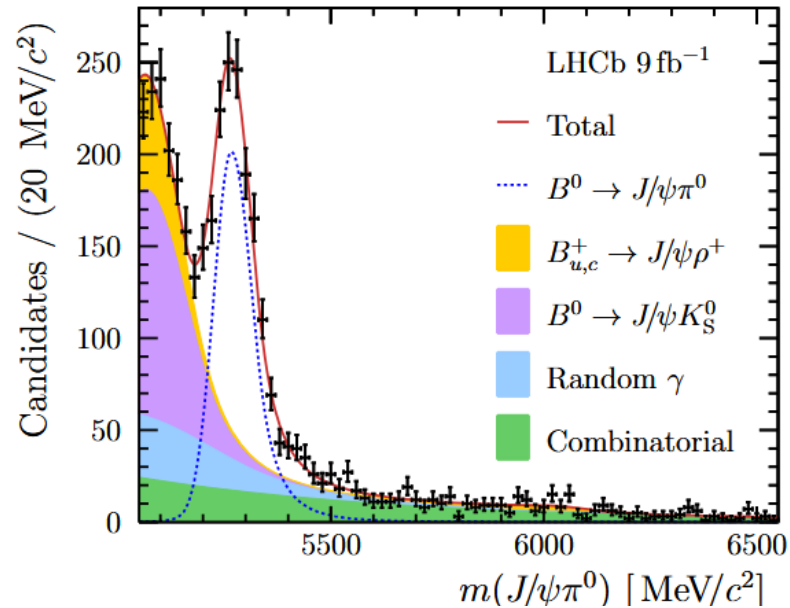
LHCb Detector in Run 2

Xiaofan Hu

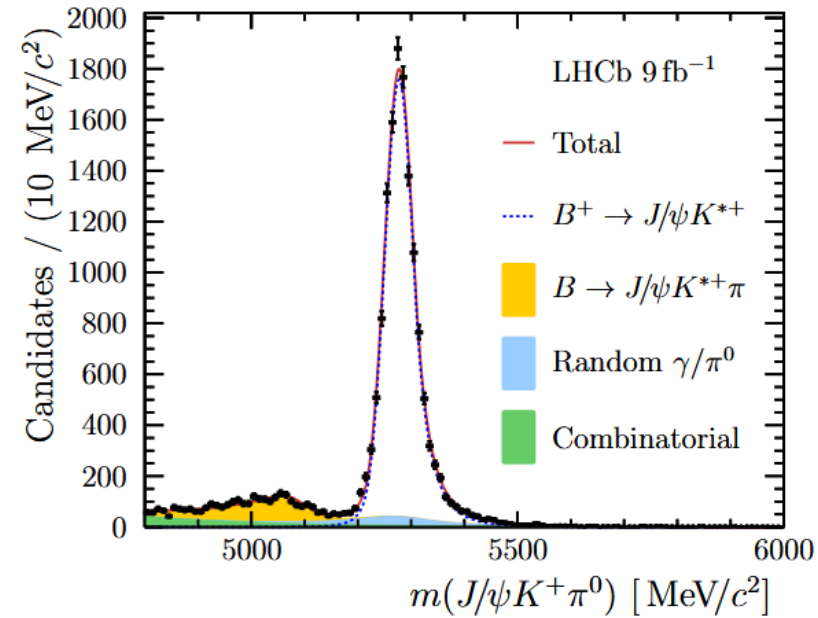
# Branching fraction of the $B^0 \rightarrow J/\psi\pi^0$ decays

- LHCb Run1 (2011–12) + Run2 (2015–2018)

**$1232 \pm 55 B^0 \rightarrow J/\psi\pi^0$  signal yields**



**$13052 \pm 115 B^+ \rightarrow J/\psi K^{*+}$  signal yields**



- Results:

$$\mathcal{R} = \frac{N_{B^0 \rightarrow J/\psi\pi^0}}{N_{B^+ \rightarrow J/\psi K^{*+}}} \times \frac{\epsilon_{B^+ \rightarrow J/\psi K^{*+}}}{\epsilon_{B^0 \rightarrow J/\psi\pi^0}} \times \mathcal{B}_{K^{*+} \rightarrow K^+\pi^0}$$

$$= (1.153 \pm 0.053 \text{ (stat.)} \pm 0.048 \text{ (syst.)}) \times 10^{-2}.$$

$$\mathcal{B}_{B^0 \rightarrow J/\psi\pi^0} = (\underbrace{1.670}_{\text{stat.}} \pm \underbrace{0.077}_{\text{sys.}} \pm \underbrace{0.069}_{\text{external}} \pm 0.095) \times 10^{-5}$$

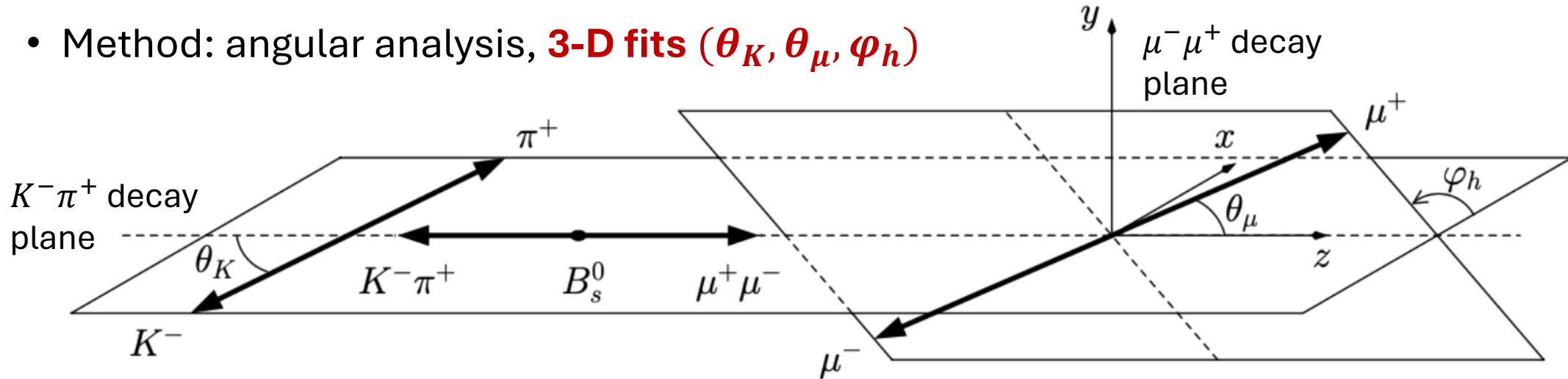
**Total uncertainty:  $0.14 \times 10^{-5}$**   
**Competitive with the most precise single measurement!**

- **Belle:  $(1.62 \pm 0.13) \times 10^{-5}$**

# CPV in the $B_s^0 \rightarrow J/\psi \bar{K}^*(892)^0 (\rightarrow K^- \pi^+)$ decay

Arxiv: [2506.22090](https://arxiv.org/abs/2506.22090)  
Submitted to JHEP

- Dataset: LHCb Run-2 data
- Method: angular analysis, **3-D fits** ( $\theta_K, \theta_\mu, \varphi_h$ )



**Angular distribution:**

$$\text{PDF}(\theta_K, \theta_\mu, \varphi_h) = \sum_{\alpha_\mu = \pm 1} \left| \sum_{\lambda, J} \sqrt{\frac{2J+1}{4\pi}} \mathcal{H}_\lambda^J e^{-i\lambda\varphi_h} d_{\lambda, \alpha_\mu}^1(\theta_\mu) d_{-\lambda, 0}^1(\theta_K) \right|^2$$

**Helicity  $\lambda = 0, \pm 1$  for  $\bar{K}^{*0}$  or  $J/\psi$**

**Angular momentum between  $K^- \pi^+$ :  $L = 1$  (p-wave)**

**Amplitudes:**

$$\begin{aligned} A_0 &= \mathcal{H}_0^1, \\ A_{\parallel} &= \frac{1}{\sqrt{2}}(\mathcal{H}_{+1}^1 + \mathcal{H}_{-1}^1), \\ A_{\perp} &= \frac{1}{\sqrt{2}}(\mathcal{H}_{+1}^1 - \mathcal{H}_{-1}^1), \\ A_S &= \mathcal{H}_0^0. \end{aligned}$$

**Integrate  $|A|^2$  over phase-space to get widths  $\Gamma$**

# $B_S^0 \rightarrow J/\psi \bar{K}^* (892)^0$ : CPV observables

Arxiv: [2506.22090](https://arxiv.org/abs/2506.22090)  
Submitted to JHEP

- Observables:

$$\mathcal{A}_k^{CP} = \frac{\bar{\Gamma}_k - \Gamma_k}{\bar{\Gamma}_k + \Gamma_k}, \quad k = 0, \parallel, \perp, S$$

- Results:

$$\mathcal{A}_0^{CP} = 0.014 \pm 0.029 \text{ (stat)} \pm 0.007 \text{ (syst)},$$

$$\mathcal{A}_{\parallel}^{CP} = -0.055 \pm 0.065 \text{ (stat)} \pm 0.007 \text{ (syst)},$$

$$\mathcal{A}_{\perp}^{CP} = 0.060 \pm 0.057 \text{ (stat)} \pm 0.016 \text{ (syst)}.$$

- After combination with LHCb Run-1 results: **JHEP 11 (2015) 082**

$$\mathcal{A}_0^{CP} = 0.021 \pm 0.026 \text{ (stat)} \pm 0.007 \text{ (syst)},$$

$$\mathcal{A}_{\parallel}^{CP} = -0.073 \pm 0.060 \text{ (stat)} \pm 0.007 \text{ (syst)},$$

$$\mathcal{A}_{\perp}^{CP} = 0.057 \pm 0.049 \text{ (stat)} \pm 0.014 \text{ (syst)}.$$

**No evidence of CP violation found**

# $B_s^0 \rightarrow J/\psi \bar{K}^* (892)^0$ : Branching fraction

Arxiv: [2506.22090](https://arxiv.org/abs/2506.22090)  
Submitted to JHEP

$$\frac{\mathcal{B}(B_s^0 \rightarrow J/\psi \bar{K}^{*0})}{\mathcal{B}(B^0 \rightarrow J/\psi K^{*0})} = \frac{N_{B_s^0}}{N_{B^0}} \times \frac{\varepsilon'_{B^0}}{\varepsilon'_{B_s^0}} \times \frac{\zeta_{B^0}}{\zeta_{B_s^0}} \times \frac{f_d}{f_s},$$

$N$  – Signal yield

$\varepsilon'$  – detection efficiency obtained from simulation

$\zeta$  – correction from data & simulation discrepancy and angular fraction

$f$  – fragmentation fraction

- Results: 
$$\frac{\mathcal{B}(B_s^0 \rightarrow J/\psi \bar{K}^{*0})}{\mathcal{B}(B^0 \rightarrow J/\psi K^{*0})} = (3.08 \pm 0.11 \text{ (stat)} \pm 0.06 \text{ (syst)} \pm 0.10 \left(\frac{f_d}{f_s}\right))\%.$$

$$\begin{aligned} \mathcal{B}(B_s^0 \rightarrow J/\psi \bar{K}^{*0}) &= \frac{\mathcal{B}(B_s^0 \rightarrow J/\psi \bar{K}^{*0})}{\mathcal{B}(B^0 \rightarrow J/\psi K^{*0})} \times C \times \mathcal{B}(B^0 \rightarrow J/\psi K^{*0})_{\text{Belle}} && \text{Correction to Belle result:} \\ & && \mathbf{C = 1.026 \pm 0.016} \\ &= (4.07 \pm 0.15 \text{ (stat)} \pm 0.07 \text{ (syst)} \pm 0.13 \left(\frac{f_d}{f_s}\right) \pm 0.45 (\mathcal{B}_{B^0})) \times 10^{-5}. \end{aligned}$$

- After combination with LHCb Run-1 result: *(JHEP 11 (2015) 082)* **With the best precision!**

$$\begin{aligned} \frac{\mathcal{B}(B_s^0 \rightarrow J/\psi \bar{K}^{*0})}{\mathcal{B}(B^0 \rightarrow J/\psi K^{*0})} &= (3.12 \pm 0.09 \text{ (stat)} \pm 0.06 \text{ (syst)} \pm 0.10 \left(\frac{f_d}{f_s}\right))\%. && \text{LHCb Run 1: } \frac{\mathcal{B}(B_s^0 \rightarrow J/\psi \bar{K}^{*0})}{\mathcal{B}(B^0 \rightarrow J/\psi K^{*0})} \\ & && = (2.99 \pm 0.14 \text{ (stat)} \pm 0.12 \text{ (syst)} \pm 0.17 (f_d/f_s))\%, \end{aligned}$$

$$\mathcal{B}(B_s^0 \rightarrow J/\psi \bar{K}^{*0}) = (4.13 \pm 0.12 \text{ (stat)} \pm 0.07 \text{ (syst)} \pm 0.14 \left(\frac{f_d}{f_s}\right) \pm 0.45 (\mathcal{B}_{B^0})) \times 10^{-5}.$$

# CP asymmetry difference of $B^+ \rightarrow J/\psi\pi^+$ and $B^+ \rightarrow J/\psi K^+$ decays

PRL 134, 101801 (2025)

- LHCb Run2 data

- $\Delta\mathcal{A}_{CP} = \Delta\mathcal{A}_{\text{raw}} - \Delta\mathcal{A}_{\text{det}} - \Delta\mathcal{A}_{\text{PID}}$

- $\Delta\mathcal{A}_{\text{raw}} = \mathcal{A}_{\text{raw}}(B^+ \rightarrow J/\psi\pi^+) - \mathcal{A}_{\text{raw}}(B^+ \rightarrow J/\psi K^+)$

$$\Delta\mathcal{A}^{CP} = \begin{cases} (1.43 \pm 0.87 \pm 0.09) \times 10^{-2} & (2016), \\ (0.81 \pm 0.87 \pm 0.11) \times 10^{-2} & (2017), \\ (1.58 \pm 0.80 \pm 0.11) \times 10^{-2} & (2018), \end{cases}$$

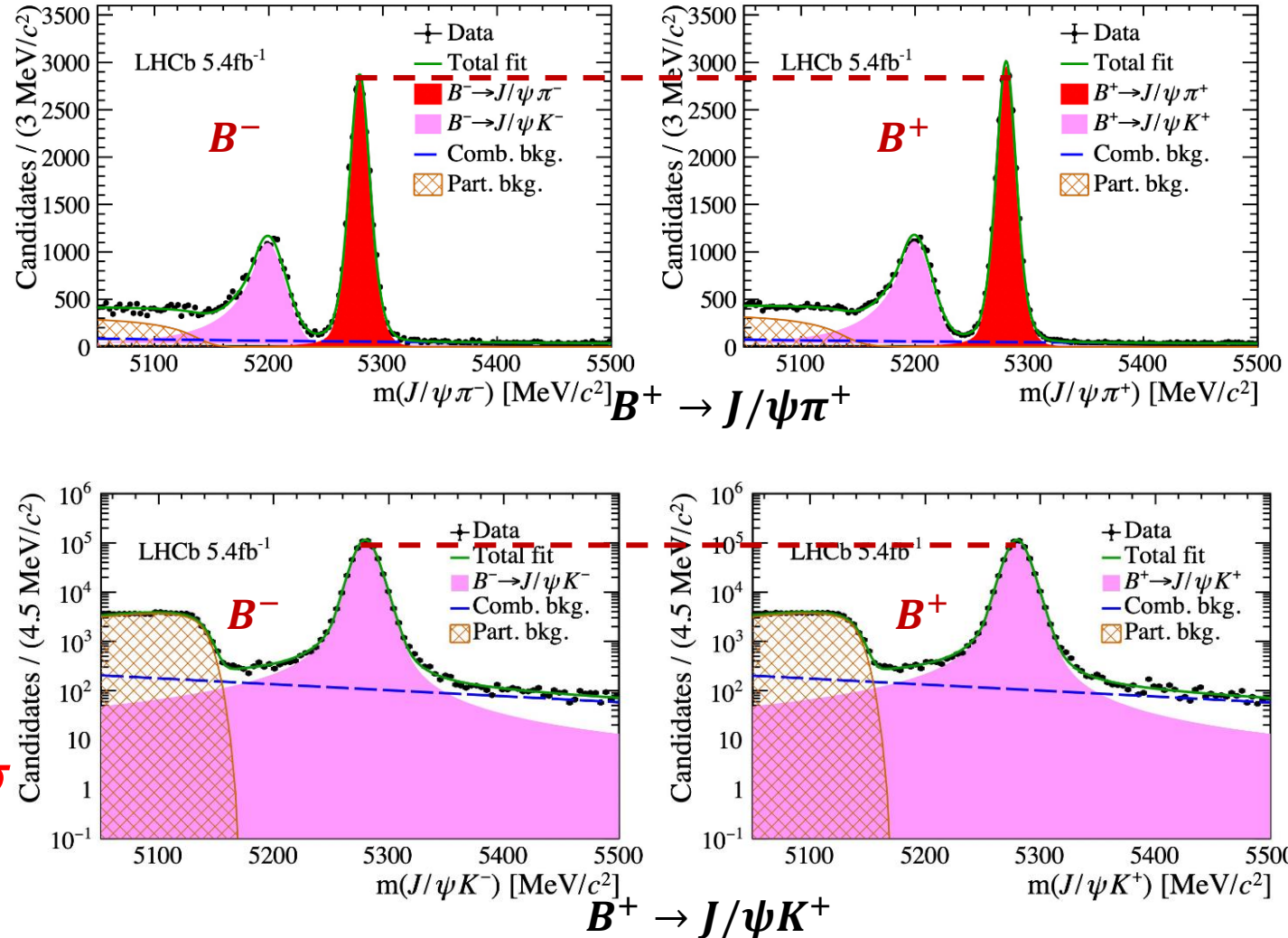
- Run-2 result:

$$\Delta\mathcal{A}_{CP} = (1.29 \pm 0.49 \pm 0.08)\%$$

Combined with Run 1:

$$\Delta\mathcal{A}_{CP} = (1.42 \pm 0.43 \pm 0.11)\%$$

**First evidence of CP violation in B hadron to charmonium decays!**



# Branching fraction ratio

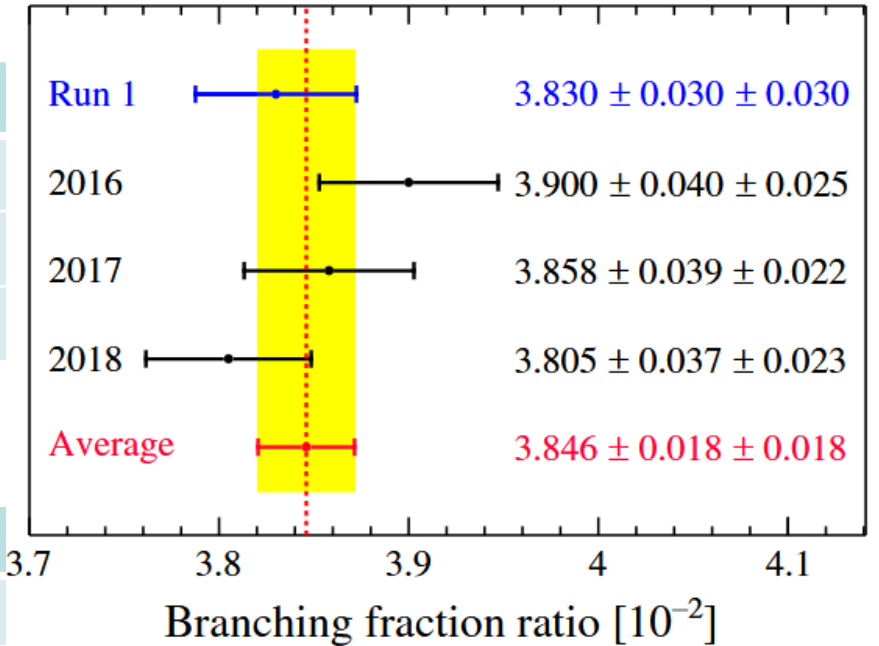
PRL 134, 101801 (2025)

- Signal Yields:

	2016	2017	2018
$N_\pi$	$15500 \pm 140$	$15140 \pm 140$	$18130 \pm 150$
$N_K$	$371700 \pm 600$	$367300 \pm 600$	$454100 \pm 700$
$N_\pi/N_K$ (%)	$4.170 \pm 0.038$	$4.121 \pm 0.039$	$3.993 \pm 0.034$

- Efficiencies:

	2016	2017	2018
$\epsilon_\pi/\epsilon_K$	$0.935 \pm 0.004$	$0.936 \pm 0.004$	$0.953 \pm 0.005$



$$\mathcal{R}_{\pi/K} \equiv \frac{\mathcal{B}(B^+ \rightarrow J/\psi\pi^+)}{\mathcal{B}(B^+ \rightarrow J/\psi K^+)} = \frac{N_\pi}{N_K} \times \frac{\epsilon_K}{\epsilon_\pi} = (3.852 \pm 0.022 \pm 0.018) \times 10^{-2}$$

**Combined with Run 1:  $\mathcal{R}_{\pi/K} = (3.846 \pm 0.018 \pm 0.018) \times 10^{-2}$**

Important for constraining penguin contribution in  $b \rightarrow c\bar{c}s$  processes!

# Constraints to penguin parameters $a$ and $\theta$

PRL 134, 101801 (2025),  
supplementary

- Decay amplitudes:  $\frac{\text{penguin}}{\text{tree}} \equiv -b_f e^{i\rho_f} e^{i\gamma}, \quad = a e^{i\theta} e^{i\gamma} (b \rightarrow c\bar{c}d); \quad = \epsilon a' e^{i\theta'} e^{i\gamma} (b \rightarrow c\bar{c}s)$
- $A(B^+ \rightarrow J/\psi\pi^+) = -\lambda A_{tree}(1 + a e^{i\theta} e^{i\gamma}), \quad A(B^+ \rightarrow J/\psi K^+) = \left(1 - \frac{\lambda^2}{2}\right) A'_{tree}(1 + \epsilon a' e^{i\theta'} e^{i\gamma})$
- Assuming SU(3) symmetry:  $a = a', \theta = \theta' \quad \lambda \sim 0.23, \epsilon \equiv \frac{\lambda^2}{1-\lambda^2} = 0.056$

**68% CL:**

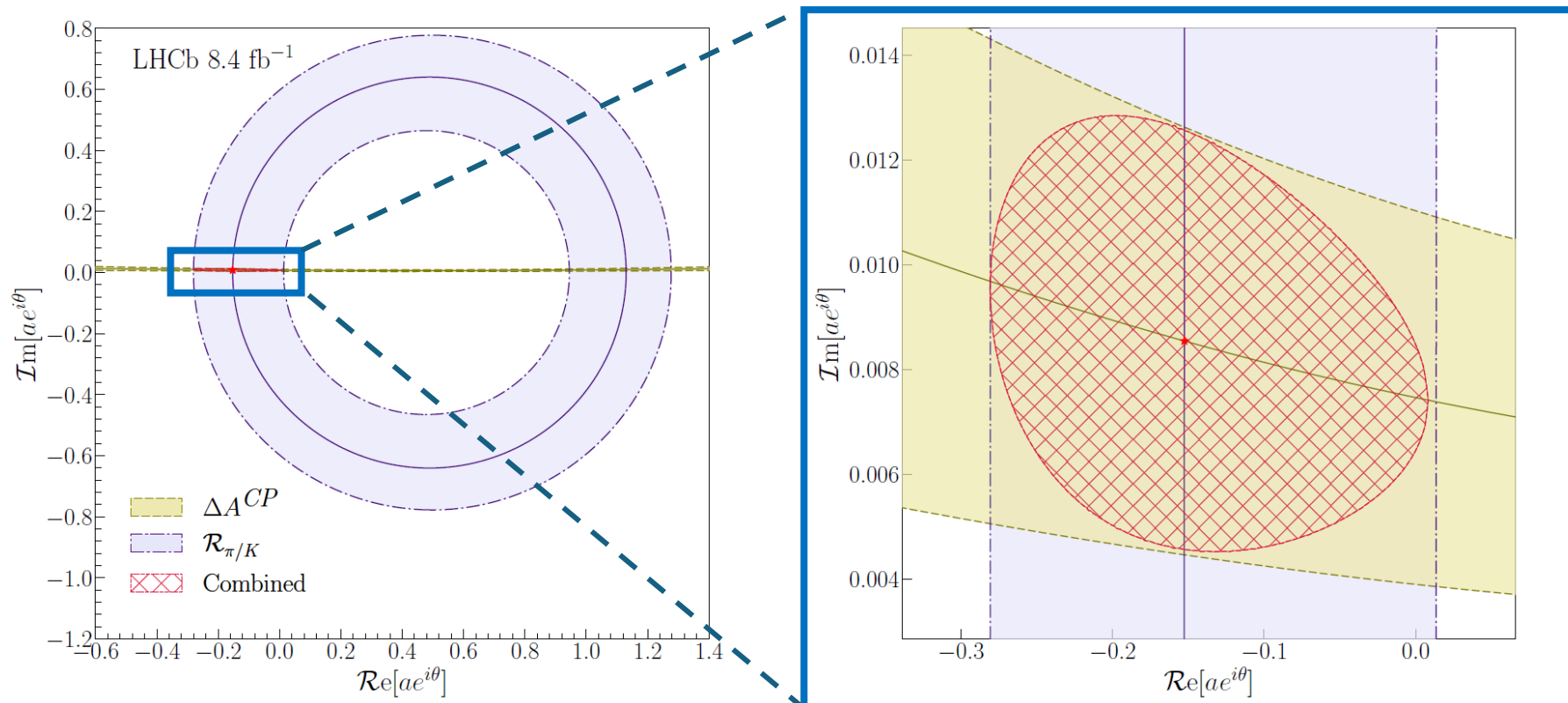
External results:

$$\frac{A'_{tree}}{A_{tree}} = 1.32 \pm 0.07$$

(*J. Phys. G*: **48** (2021) 065002)

$$\gamma = (64.6 \pm 2.8)^\circ$$

(LHCb-CONF-2024-004)



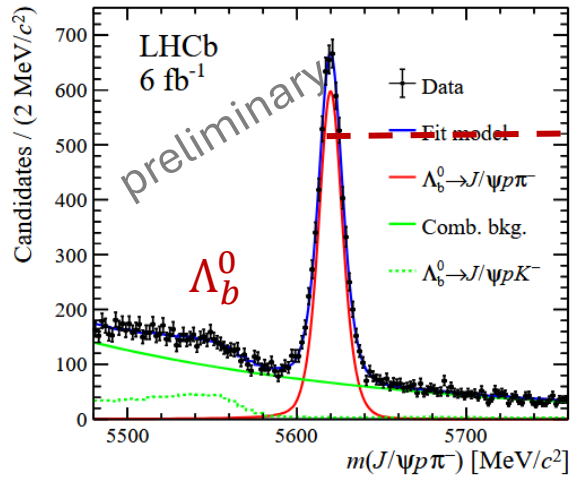
# CPV difference in $\Lambda_b^0 \rightarrow J/\psi p \pi^-$ and $\Lambda_b^0 \rightarrow J/\psi p K^-$ decays

[LHCb-PAPER-2025-021], **[New!]**  
in preparation

- Measurement of the CP asymmetry difference  $\Delta\mathcal{A}_{CP}$

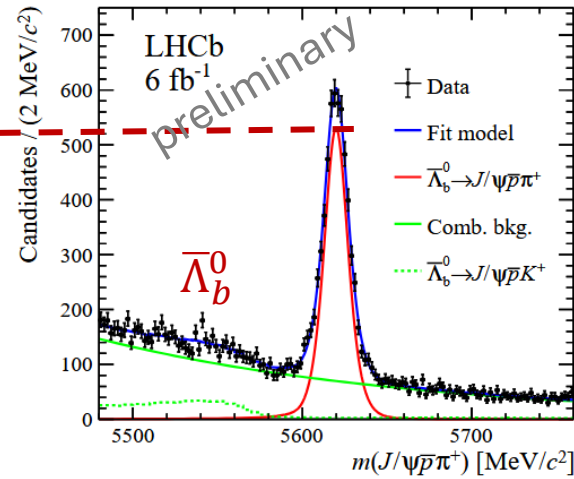
- $\Delta\mathcal{A}_{CP} = \mathcal{A}_{\text{raw}}(\Lambda_b^0 \rightarrow J/\psi p \pi^-) - \mathcal{A}_{\text{raw}}(\Lambda_b^0 \rightarrow J/\psi p K^-) + A_D(K^-) - A_D(\pi^-)$

- LHCb Run2:



(10853 ± 134) yields

$$\mathcal{A}_{\text{raw}}(\Lambda_b^0 \rightarrow J/\psi p \pi^-) = (5.94 \pm 1.14) \%$$



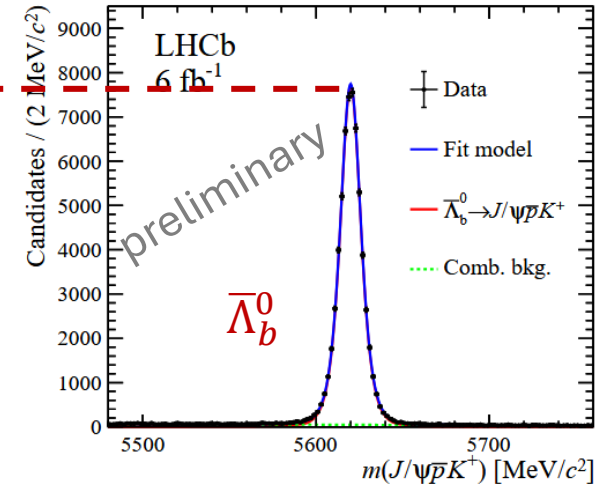
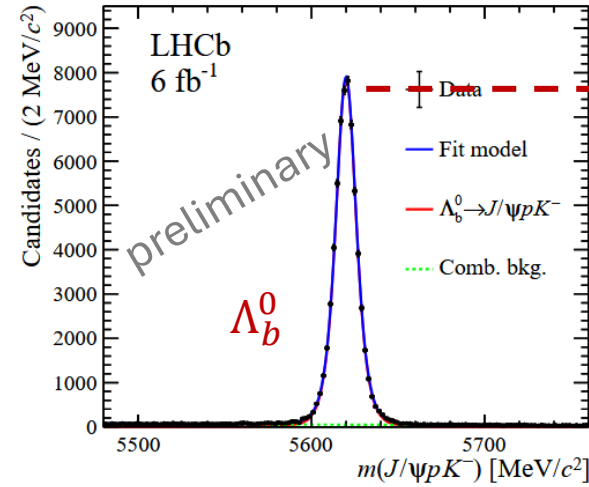
(125380 ± 372) yields

$$\mathcal{A}_{\text{raw}}(\Lambda_b^0 \rightarrow J/\psi p K^-) = (0.98 \pm 0.30) \%$$

- $\Delta\mathcal{A}_{CP} = (4.03 \pm 1.18 \pm 0.23)\% \quad \mathbf{3.3 \sigma}$

- Combined with Run 1:  $\Delta\mathcal{A}_{CP} = (4.31 \pm 1.06 \pm 0.28)\% \quad \mathbf{3.9 \sigma}$

**First evidence of CP violation in b baryon to charmonium decays!**



Run 1 results: **JHEP 07 (2014) 103**  
 $\Delta\mathcal{A}_{CP} = (5.7 \pm 2.4 \pm 1.2)\%$   
 $\mathcal{R}_{\pi/K} = (8.24 \pm 0.25 \pm 0.42) \times 10^{-2}$

# Triple-product asymmetry (TPA) of the $\Lambda_b^0 \rightarrow J/\psi p \pi^-$ decay

[LHCb-PAPER-2025-021], **[New!]**  
in preparation

- Triple product:

$$\Lambda_b^0: \quad C_T \equiv \vec{p}_{\mu^+} \cdot (\vec{p}_p \times \vec{p}_{\pi^-})$$

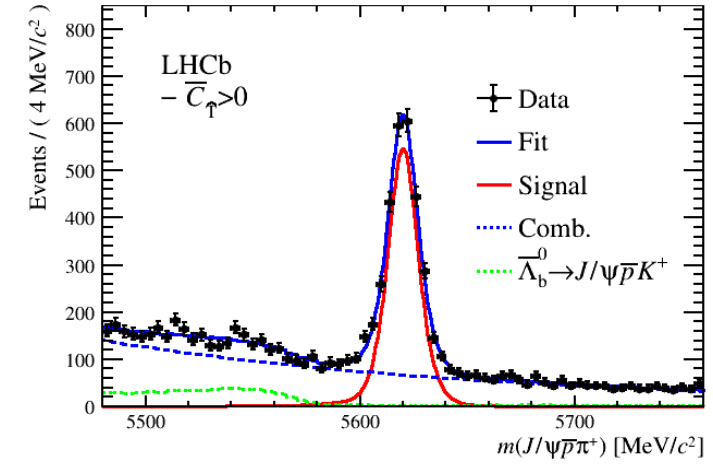
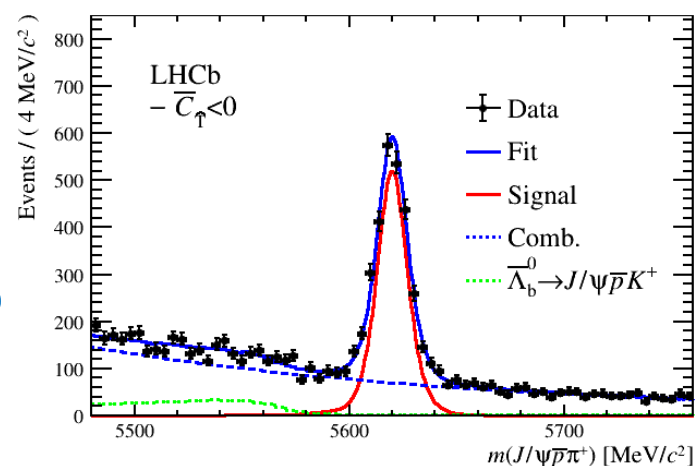
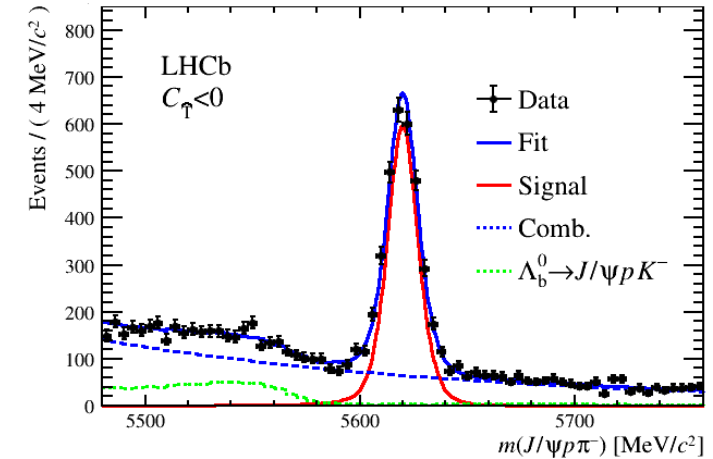
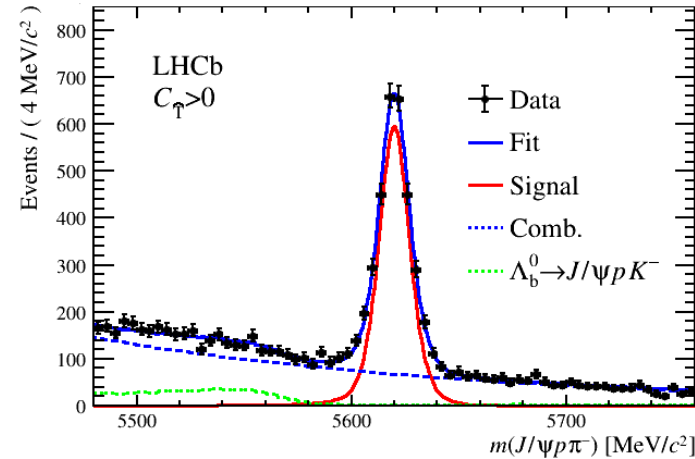
$$\bar{\Lambda}_b^0: \quad \bar{C}_T \equiv \vec{p}_{\mu^-} \cdot (\vec{p}_{\bar{p}} \times \vec{p}_{\pi^+})$$

- Asymmetry  $\mathcal{A}_{\hat{T}} = \frac{\Gamma(C_T > 0) - \Gamma(C_T < 0)}{\Gamma(C_T > 0) + \Gamma(C_T < 0)}$
- Under CP:  $C_T \rightarrow -\bar{C}_T$ ,  $\mathcal{A}_{\hat{T}} \rightarrow \bar{\mathcal{A}}_{\hat{T}}$
- CP violation variable:

$$\mathcal{A}_{T\text{-odd}} = \frac{1}{2} (\mathcal{A}_{\hat{T}} - \bar{\mathcal{A}}_{\hat{T}})$$

(no evidence of TPA found)

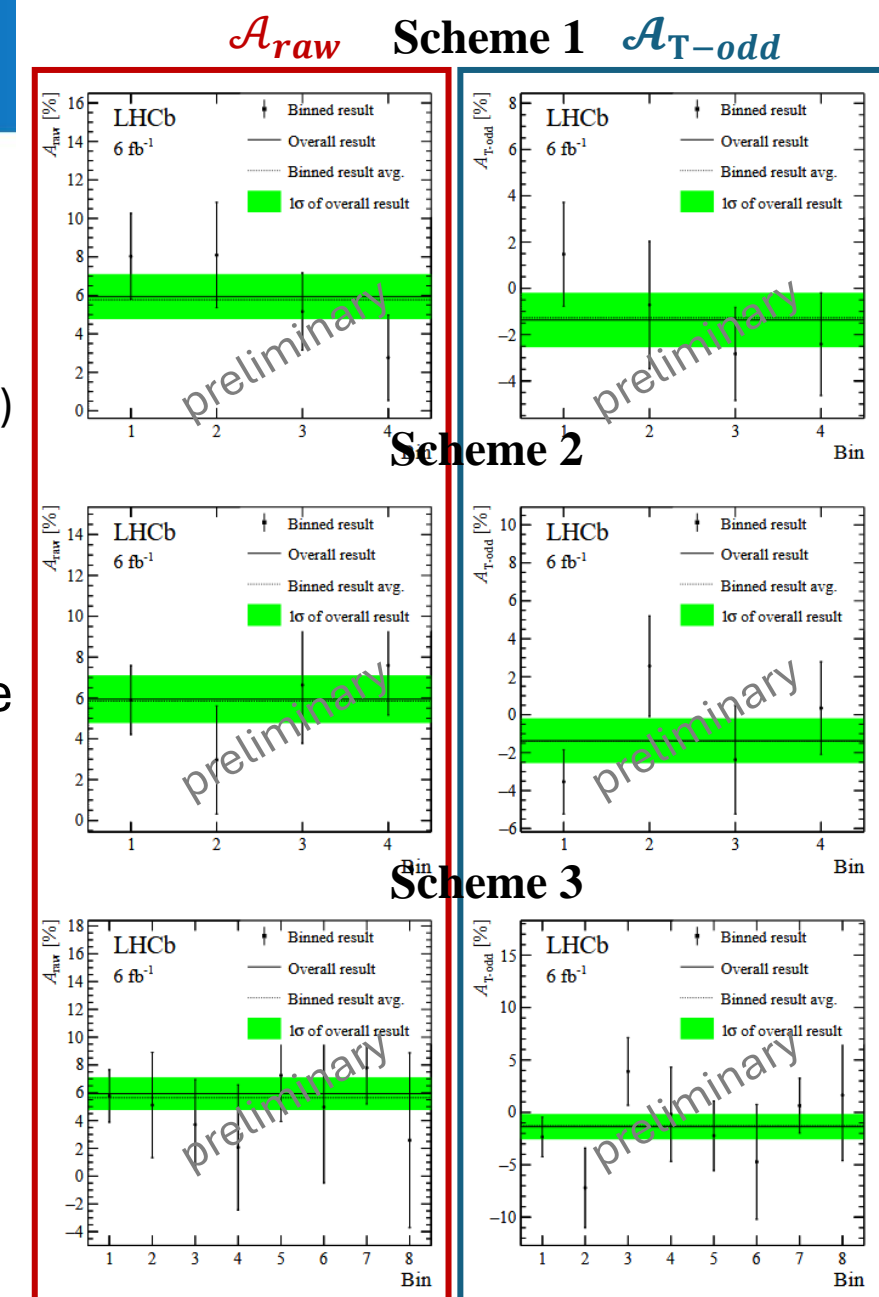
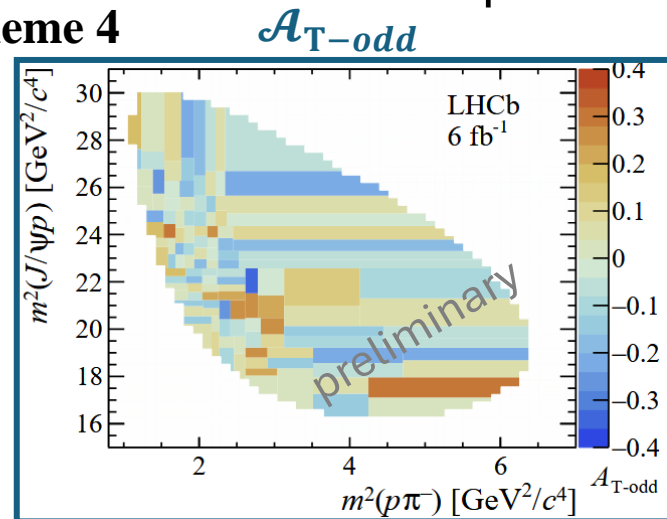
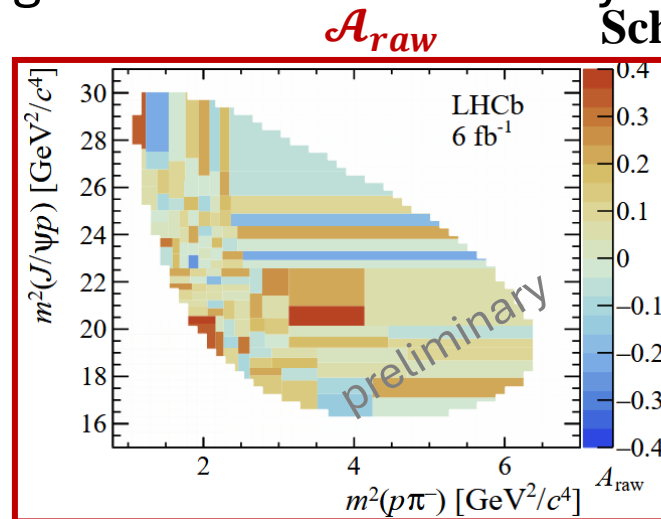
$$\mathcal{A}_{T\text{-odd}} = (-1.37 \pm 1.15 \pm 0.21) \%$$



# Binned analysis

[LHCb-PAPER-2025-021],  
in preparation **[New!]**

- Measurement of local CP asymmetry with **4 binning schemes**
  - 1: Split the  $J/\psi p\pi^-$  data sample evenly in Dalitz-plot (4 bins)
  - 2: Split according to resonances in  $M(p\pi^-)$  spectrum (4 bins)
  - 3: Split binning scheme 2 further, according to boosting angle  $\theta_{p\pi}$  (8 bins)
  - 4: Split evenly into 128 bins
- Both raw asymmetry  $A_{raw}$  and TPA  $A_{T-odd}$  are measured
- No significant variation of asymmetries across the phase space



# Summary

- Conclusion
  - Branching fraction of  $B^0 \rightarrow J/\psi\pi^0$  : a first step towards its CPV analysis at LHCb
  - Branching fraction and CPV of  $B_s^0 \rightarrow J/\psi\bar{K}^*(892)^0$ : most precise values to date
  - Branching ratio and CPV of  $B^+ \rightarrow J/\psi\pi^+$  : first evidence of CPV, constraints to penguin parameters
  - CPV of  $\Lambda_b^0 \rightarrow J/\psi p\pi^-$ : first evidence of CPV
- Prospects
  - LHCb Run3 (started in 2023)
    - Higher integrated luminosity (25–30 fb<sup>-1</sup>), significantly larger dataset

*Thanks for your attention!*

# BACKUPS

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# Wolfenstein parametrization

$$V_{\text{CKM}} = \begin{pmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{pmatrix}$$

$$V_{\text{CKM}} = \begin{pmatrix} 1 - \frac{1}{2}\lambda^2 - \frac{1}{8}\lambda^4 & \lambda & A\lambda^3(\rho - i\eta) \\ -\lambda + \frac{1}{2}A^2\lambda^5[1 - 2(\rho + i\eta)] & 1 - \frac{1}{2}\lambda^2 - \frac{1}{8}\lambda^4(1 + 4A^2) & A\lambda^2 \\ A\lambda^3[1 - (1 - \frac{1}{2}\lambda^2)(\rho + i\eta)] & -A\lambda^2 + \frac{1}{2}A\lambda^4[1 - 2(\rho + i\eta)] & 1 - \frac{1}{2}A^2\lambda^4 \end{pmatrix}.$$

- Take  $B^0 \rightarrow J/\psi K^0$  as example ( $b \rightarrow c\bar{c}s$ ):

$$A(B_d^0 \rightarrow J/\psi K_S) = \lambda_c^{(s)} (A_{cc}^{c'} + A_{pen}^{c'}) + \lambda_u^{(s)} A_{pen}^{u'} + \lambda_t^{(s)} A_{pen}^{t'},$$

$$\lambda_q^{(s)} \equiv V_{qs} V_{qb}^* \quad (2)$$

are the usual CKM factors. Making use of the unitarity of the CKM matrix and applying the Wolfenstein parametrization [9], generalized to include non-leading terms in  $\lambda$  [10], we obtain

$$A(B_d^0 \rightarrow J/\psi K_S) = \left(1 - \frac{\lambda^2}{2}\right) \mathcal{A}' \left[1 + \left(\frac{\lambda^2}{1 - \lambda^2}\right) a' e^{i\theta'} e^{i\gamma}\right], \quad (3)$$

where

$$\mathcal{A}' \equiv \lambda^2 A (A_{cc}^{c'} + A_{pen}^{ct'}), \quad (4)$$

with  $A_{pen}^{ct'} \equiv A_{pen}^{c'} - A_{pen}^{t'}$ , and

$$a' e^{i\theta'} \equiv R_b \left(1 - \frac{\lambda^2}{2}\right) \left(\frac{A_{pen}^{ut'}}{A_{cc}^{c'} + A_{pen}^{ct'}}\right). \quad (5)$$

- For  $B_S^0 \rightarrow J/\psi \bar{K}^0$  ( $b \rightarrow c\bar{c}d$ ):

$$A(B_S^0 \rightarrow J/\psi K_S) = -\lambda \mathcal{A} [1 - a e^{i\theta} e^{i\gamma}],$$

$$\mathcal{A} \equiv \lambda^2 A (A_{cc}^c + A_{pen}^{ct})$$

$$a e^{i\theta} \equiv R_b \left(1 - \frac{\lambda^2}{2}\right) \left(\frac{A_{pen}^{ut}}{A_{cc}^c + A_{pen}^{ct}}\right)$$

- For  $B^+ \rightarrow J/\psi \pi^+$  or  $J/\psi K^+$ :

- $a = a', \theta = \theta'$

- Normalization factors

$$\frac{\mathcal{A}'}{\mathcal{A}} = f_{B_d \rightarrow K} / f_{B_d \rightarrow \pi} = 1.32 \pm 0.07$$

Form factors:

$$f_{B_d \rightarrow \pi}^+(m_{J/\psi}^2) = 0.487 \pm 0.018,$$

$$f_{B_d \rightarrow K}^+(m_{J/\psi}^2) = 0.645 \pm 0.022,$$

- Measure  $\beta_s$  using  $B_s^0$  decays to CP eigenstates with  $b \rightarrow c\bar{c}s$  quark transition

- Benchmark channel:  $B_s^0 \rightarrow J/\psi\phi$

$$\beta_s \equiv \arg\left(-\frac{V_{ts}V_{tb}^*}{V_{cs}V_{cb}^*}\right)$$

- When measuring the time-dependent CP asymmetry:

$$\mathcal{A}_{CP}(t) = \frac{\mathcal{A}_{CP}^{dir} \cos(\Delta mt) + \mathcal{A}_{CP}^{mix} \sin(\Delta mt)}{\cosh(\Delta\Gamma t/2) + \mathcal{A}_{\Delta\Gamma} \sinh(\Delta\Gamma t/2)}$$

- Amplitude:  $A(B_s^0 \rightarrow f) = A_{tree} + A_{penguin} \equiv A_{tree} \times (1 - b_f e^{i\rho_f} e^{i\gamma})$

- Assuming no penguin contribution:  $b_f = 0 \Rightarrow \mathcal{A}_{CP}^{mix} = \eta_f \sin(2\beta_s + \phi^{NP})$

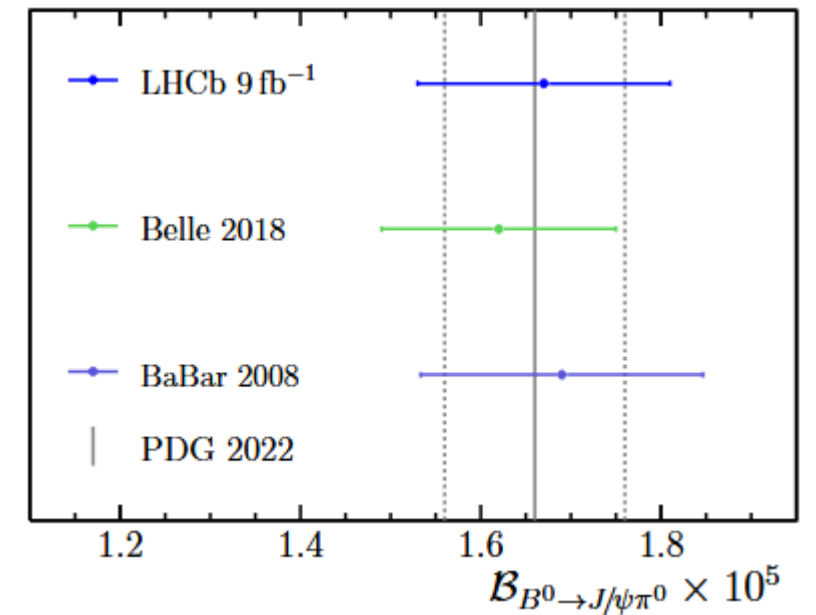
- Penguin pollution:  $b_f \neq 0 \Rightarrow \mathcal{A}_{CP}^{mix} = \eta_f \sin(2\beta_s + \phi^{NP} + \Delta\phi_f)$  ( $\eta_f = \pm 1$ ), CP eigenvalue

- **Constrain penguin using  $B_s$  decays with  $b \rightarrow c\bar{c}d$  transition!**

	$\sin 2\beta$	$\sin 2\beta_s$
Global fit result	$0.7155^{+0.0079}_{-0.0071}$ (CKMFitter)	$0.03757^{+0.00057}_{-0.00054}$ (CKMFitter)
Average from $\mathcal{A}_{CP}^{mix}$	$0.699 \pm 0.015$ (UTfit 2025)	$0.052 \pm 0.013$ (HFLAV Winter 2024)

# B to Jpsi pi0

- Measurement by Belle, Babar, and global average
- Global fit of  $B^+ \rightarrow J/\psi K^{*+}$



The current world-average for the normalisation branching fraction of  $(1.43 \pm 0.08) \times 10^{-3}$  [12] includes seven measurements performed using  $K^{*+}$  candidates with a reconstructed mass within a  $100 \text{ MeV}/c^2$  mass window or smaller of the known  $K^{*+}$  mass [40–46]. The Belle analysis [42] is the only to measure the yield for resonant decays but estimates the contributions from nonresonant decays at the level of 5.2% in the  $K^*$  mass region. Including the nonresonant contributions, the uncertainty-weighted average of the normalisation branching fraction increases to

$$\mathcal{B}_{B^+ \rightarrow J/\psi K^{*+}} = (1.449 \pm 0.083) \times 10^{-3}.$$

# Angular correction to $B_s^0 \rightarrow J/\psi \bar{K}^* (892)^0$

- Angular correlation factor

$$\zeta \equiv c_{\text{decay}} / F_{\text{decay}}^{\text{P}}$$

$$F_{\text{decay}}^{\text{P}} = \frac{\sum_{k=1}^6 \xi_k a_k}{\sum_{k=1}^{10} \xi_k a_k C_k},$$

$$c_{\text{decay}} = \frac{\sum_{k=1}^6 \xi_k a_k (\vec{A}_{\text{data}})}{\sum_{k=1}^6 \xi_k a_k (\vec{A}_{\text{sim}})}.$$

$$\frac{\zeta_{B^0}}{\zeta_{B_s^0}} = 1.015 \pm 0.034 (\text{stat}) \pm 0.017 (\text{syst}).$$

P-wave

acceptance correction  $\xi_k = \int \varepsilon(\Omega_{\text{hel}}) G_k(\Omega_{\text{hel}}) d\Omega_{\text{hel}}$

Index $k$	Amplitude product $a_k$	Angular function $G_k(\Omega_{\text{hel}})$
1 (0)	$ A_0 ^2$	$\frac{9}{16\pi} \cos^2 \theta_K (1 - \cos^2 \theta_\mu)$
2 ( $\parallel$ )	$ A_\parallel ^2$	$\frac{9}{32\pi} (1 - \cos^2 \theta_K) (1 - (1 - \cos^2 \theta_\mu) \cos^2 \phi_h)$
3 ( $\perp$ )	$ A_\perp ^2$	$\frac{9}{32\pi} (1 - \cos^2 \theta_K) ((1 - \cos^2 \theta_\mu) \cos^2 \phi_h + \cos^2 \theta_\mu)$
4	$\text{Im}(A_\parallel^* A_\perp)$	$\frac{9}{16\pi} (1 - \cos^2 \theta_K) (1 - \cos^2 \theta_\mu) \sin(2\phi_h)$
5	$\text{Re}(A_0^* A_\parallel)$	$\frac{9\sqrt{2}}{16\pi} \cos \theta_K \cos \theta_\mu \sqrt{(1 - \cos^2 \theta_K) (1 - \cos^2 \theta_\mu)} \cos \phi_h$
6	$\text{Im}(A_0^* A_\perp)$	$\frac{9\sqrt{2}}{16\pi} \cos \theta_K \cos \theta_\mu \sqrt{(1 - \cos^2 \theta_K) (1 - \cos^2 \theta_\mu)} \sin \phi_h$
7 (S)	$ A_S ^2$	$\frac{3}{16\pi} (1 - \cos^2 \theta_\mu)$
8	$\text{Re}(A_S^* A_\parallel)$	$\frac{3\sqrt{6}}{16\pi} \cos \theta_\mu \sqrt{(1 - \cos^2 \theta_K) (1 - \cos^2 \theta_\mu)} \cos \phi_h$
9	$\text{Im}(A_S^* A_\perp)$	$\frac{3\sqrt{6}}{16\pi} \cos \theta_\mu \sqrt{(1 - \cos^2 \theta_K) (1 - \cos^2 \theta_\mu)} \sin \phi_h$
10	$\text{Re}(A_S^* A_0)$	$\frac{3\sqrt{3}}{8\pi} \cos \theta_K (1 - \cos^2 \theta_\mu)$

Index $k$	Amplitude product $a_k$	Normalised weight $\xi_k / \xi_1$
1	$ A_0 ^2$	1
2	$ A_\parallel ^2$	$1.4249 \pm 0.0073$
3	$ A_\perp ^2$	$1.4489 \pm 0.0074$
4	$ A_\parallel   A_\perp  \sin(\delta_\parallel - \delta_\perp)$	$-0.0013 \pm 0.0043$
5	$ A_0   A_\parallel  \cos(\delta_\parallel)$	$-0.0189 \pm 0.0029$
6	$ A_0   A_\perp  \sin(\delta_\perp)$	$0.0027 \pm 0.0028$
7	$ A_S ^2$	$1.2485 \pm 0.0062$
8	$ A_S   A_\parallel  \cos(\delta_S - \delta_\parallel)$	$-0.0412 \pm 0.0040$
9	$ A_S   A_\perp  \sin(\delta_S - \delta_\perp)$	$0.0072 \pm 0.0040$
10	$ A_S   A_0  \cos(\delta_S)$	$-0.6919 \pm 0.0096$

# $B_S^0 \rightarrow J/\psi \bar{K}^*(892)^0$ : correction for Belle results

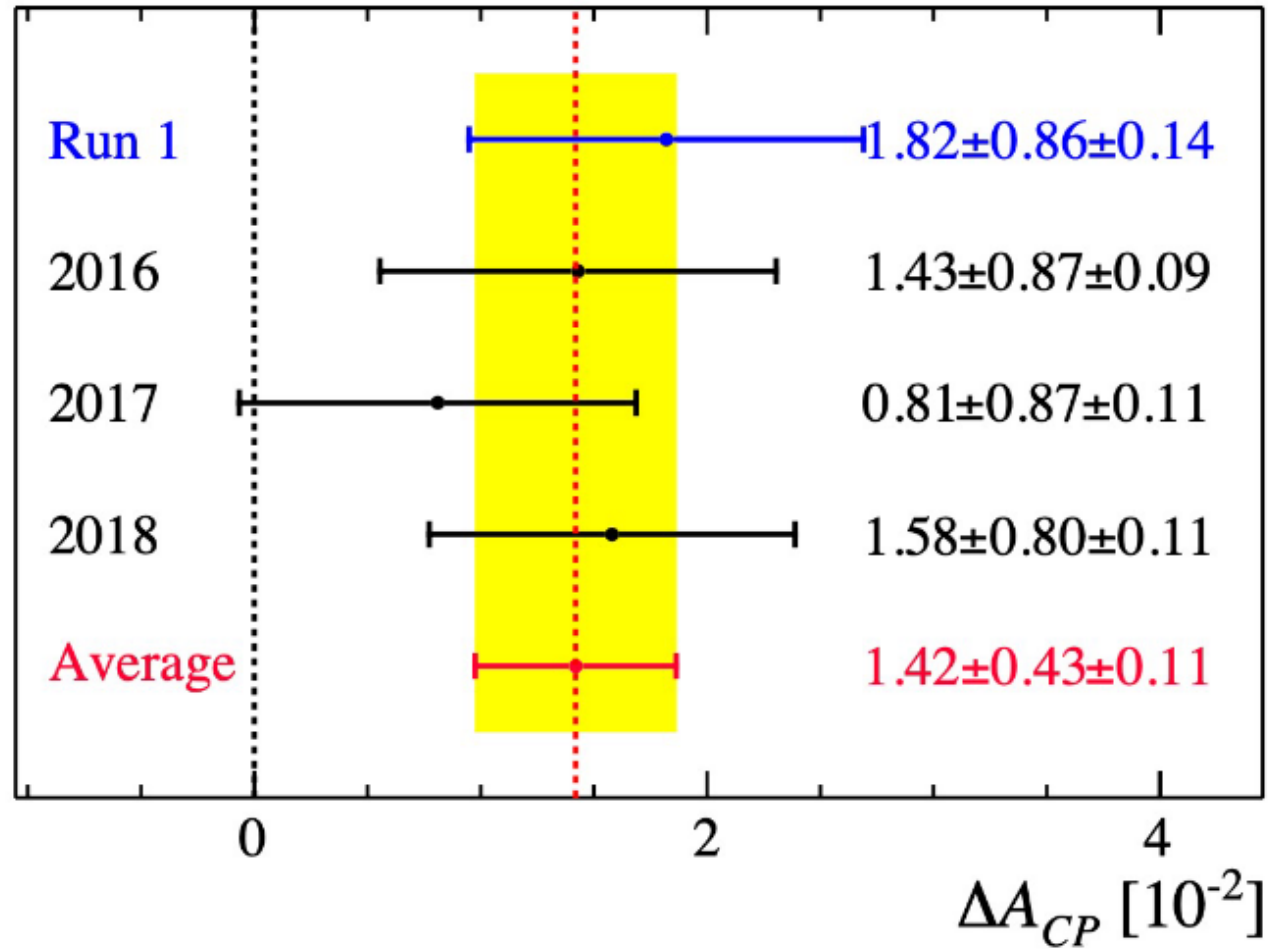
- Belle uses B mesons from  $\Upsilon(4S) \rightarrow B^0 \bar{B}^0$  to measure **Phys. Lett. B538 (2002) 11**

$$\mathcal{B}(B^0 \rightarrow J/\psi K^{*0})_{\text{Belle}} = (1.29 \pm 0.05 \text{ (stat)} \pm 0.13 \text{ (syst)}) \times 10^{-3}.$$

- In that analysis, the production ratio  $R^{+/0} \equiv \Gamma(\Upsilon(4S) \rightarrow B^+ B^-) / \Gamma(\Upsilon(4S) \rightarrow B^0 \bar{B}^0)$  was assumed to be 1.
- Measured ratio:  $R^{+/0} = 1.052 \pm 0.031$
- Correction factor:  $C = (1 + R^{+/0})/2 = 1.026 \pm 0.016$

**HFLAV 2023**

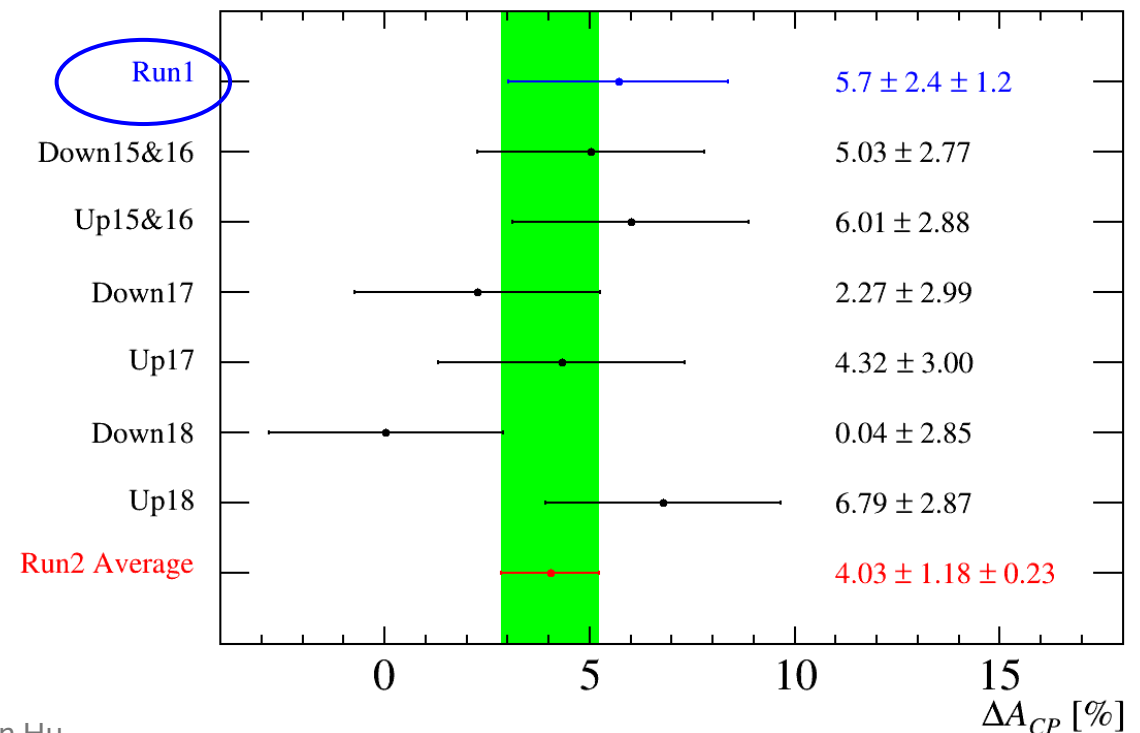
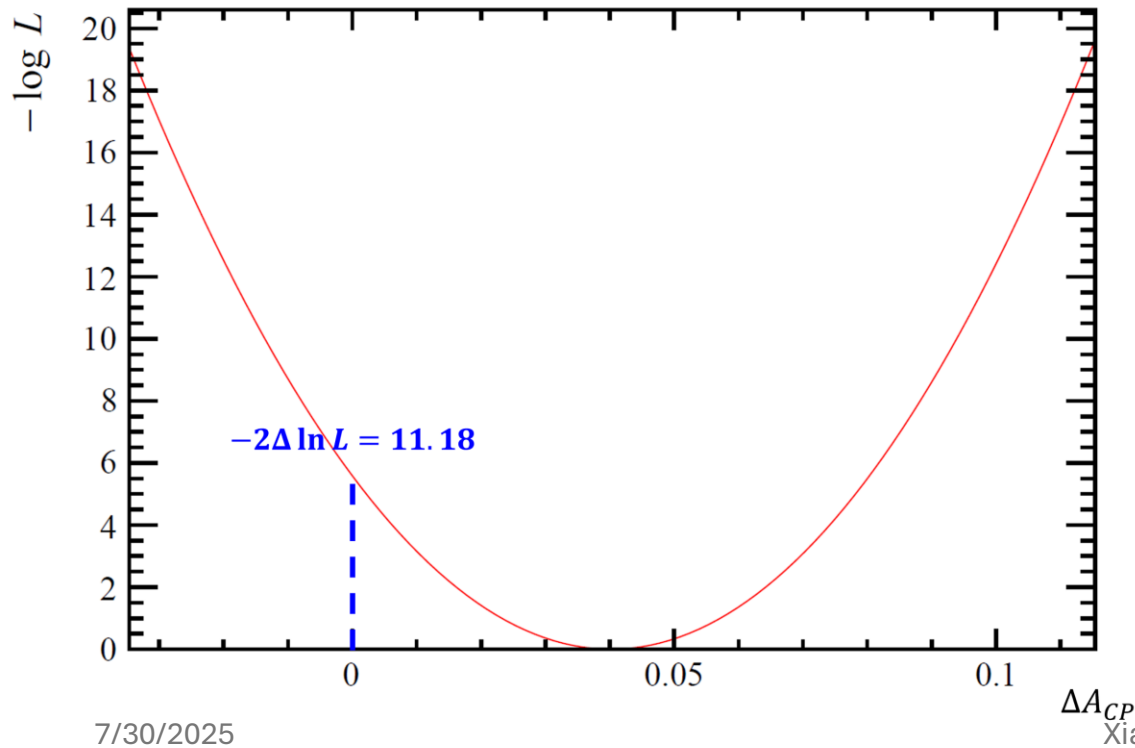
# $B^+ \rightarrow J/\psi\pi^+$ : CP asymmetry of each year



# CPV difference in $\Lambda_b^0 \rightarrow J/\psi p \pi^-$ and $\Lambda_b^0 \rightarrow J/\psi p K^-$

[LHCb-PAPER-2025-021], **[New!]**  
in preparation

- Significance level from Wilks's theorem:  $\sqrt{-2\Delta \ln \mathcal{L}} = 3.3\sigma$
- Results in different years/polarities compatible with overall results, as well as **run-1 result**
- Combined with the LHCb Run1 results with the Best Linear Unbiased Estimate (BLUE) method
  - $\Delta\mathcal{A}_{CP} = (4.31 \pm 1.06 \pm 0.28)\%$  **3.9 $\sigma$**  **First evidence of CP violation in b-baryon to charmonium decays!**



# Combination with Run-1 results

- Combination of  $\Delta\mathcal{A}_{CP}$
- Run-1 result (JHEP 07 (2014) 103):  **$(5.7 \pm 2.4 \pm 1.2)\%$**
- Best Linear Unbiased Estimation (BLUE) method is used for sys. uncertainty

Source	Run 1 [%]	Run 2 [%]	Correlation assumption
Stat.	2.4	1.18	0%
signal shape	0.14	0.016	100%
comb. bkg	0.71	0.012	100%
Detection asymmetry	0.85	0.040	0%
Mis-id $J/\psi K^+ \pi^-$ background	–	0.013	0%
Binding $\Lambda_b^0$ peak parameters	–	0.21	0%
Reweighting	–	0.017	0%
PID asymmetry	–	0.074	0%

Combined result:  **$\Delta\mathcal{A}_{CP} = (4.31 \pm 1.06(\text{stat.}) \pm 0.28(\text{sys.}))\%$**   **$3.9 \sigma$**

# The BLUE method

- Total result  $\hat{a}$  and total variance  $\hat{\sigma}^2$ :

$$\hat{a} = \sum_i^n w_i a_i,$$
$$\hat{\sigma}^2 = \sum_{i,j} w_i V_{ij} w_j.$$

- The total **covariance matrix** is calculated as

$$V_{ij} = V_{ij}^{\text{stat}} + V_{ij}^u + V_{ij}^c$$

- ▶ Diagonal:  $V_{ii} = \sigma_{i,\text{tot}}^2$
- ▶ Non-diagonal:  $V_{ij} = V_{ji} = \rho_{ij} \sigma_{i,c} \sigma_{j,c}$   
( $\rho \in [-1, 1]$  is the correlation coefficient)

- Weights:

$$w_i = \frac{\sum_j (V^{-1})_{ij}}{\sum_{ij} (V^{-1})_{ij}}$$

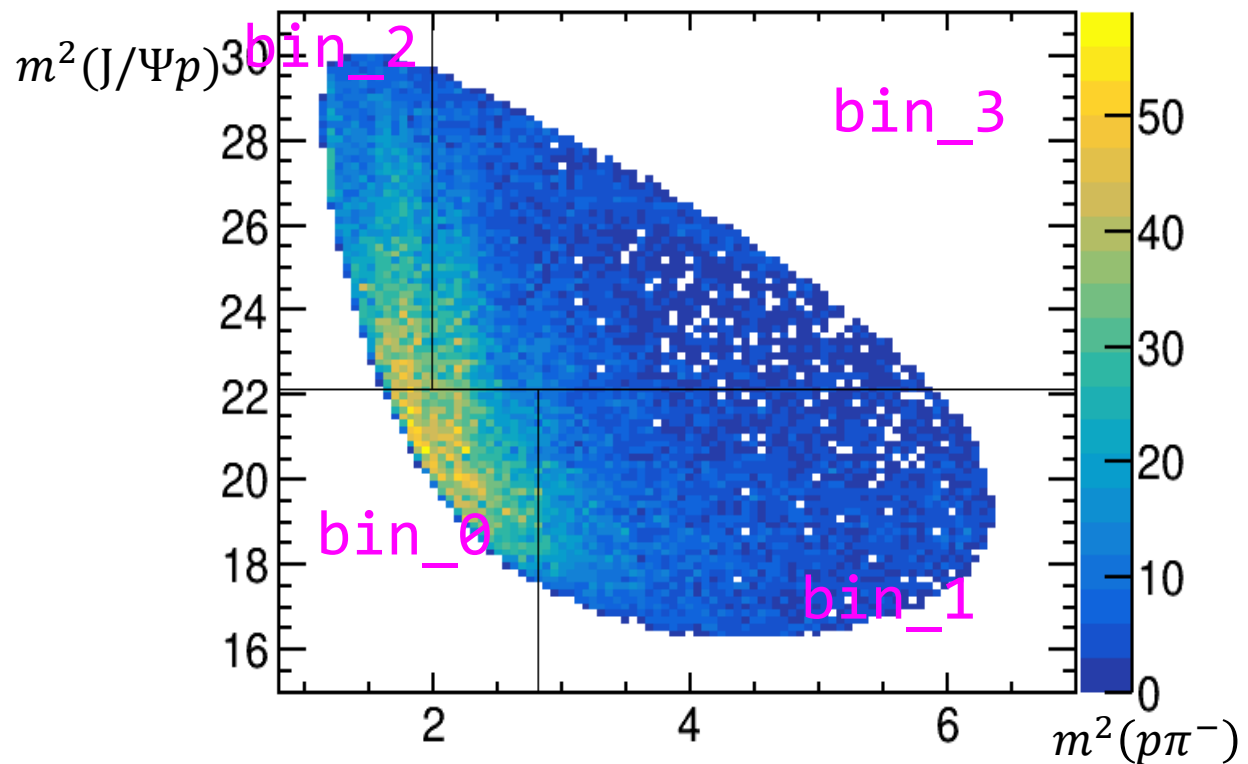
## Summary of combination:

1. Construct the covariance matrix, calculate its inverse, and then the weights.
2. Average the measured values using the weights.
3. Calculate each component of uncertainties:

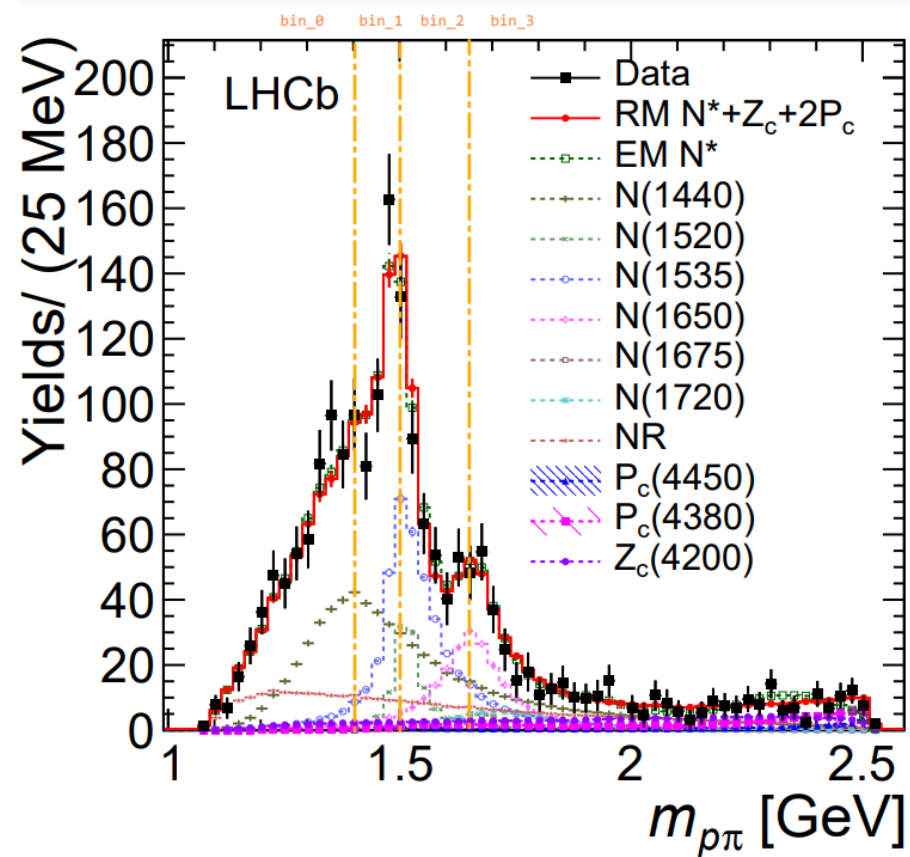
- ▶ Statistical:  $\hat{\sigma}_{\text{stat}} = \sqrt{\sum_i w_i^2 V_{ii}^{\text{stat}}} = \sqrt{\sum_i w_i^2 \sigma_{i,\text{stat}}^2}$
- ▶ Systematic, uncorrelated:  $\hat{\sigma}_u = \sqrt{\sum_i w_i^2 V_{ii}^u} = \sqrt{\sum_i w_i^2 \sigma_{i,u}^2}$
- ▶ Systematic, correlated:  $\hat{\sigma}_c = \sqrt{\sum_{ij} w_i w_j V_{ij}^c}$

# Binning schemes

- Scheme 1



- Scheme 2



# Significance

- Look-elsewhere effects for 128 bins:

Global significance	$3\sigma$	$5\sigma$
Local significance requirement	$4.3\sigma$	$5.9\sigma$

$$P(1 \text{ or more bins } p\text{-value} \leq t; 128 \text{ bins})$$

$$= 1 - P(128 \text{ bins } p\text{-value} > t)$$

$$= 1 - [P(p\text{-value} > t)]^{128} = 1 - [1 - t]^{128},$$

