

Preliminary Investigation of a Higgs Factory based on Proton-Driven Plasma Wakefield Acceleration

September 2024

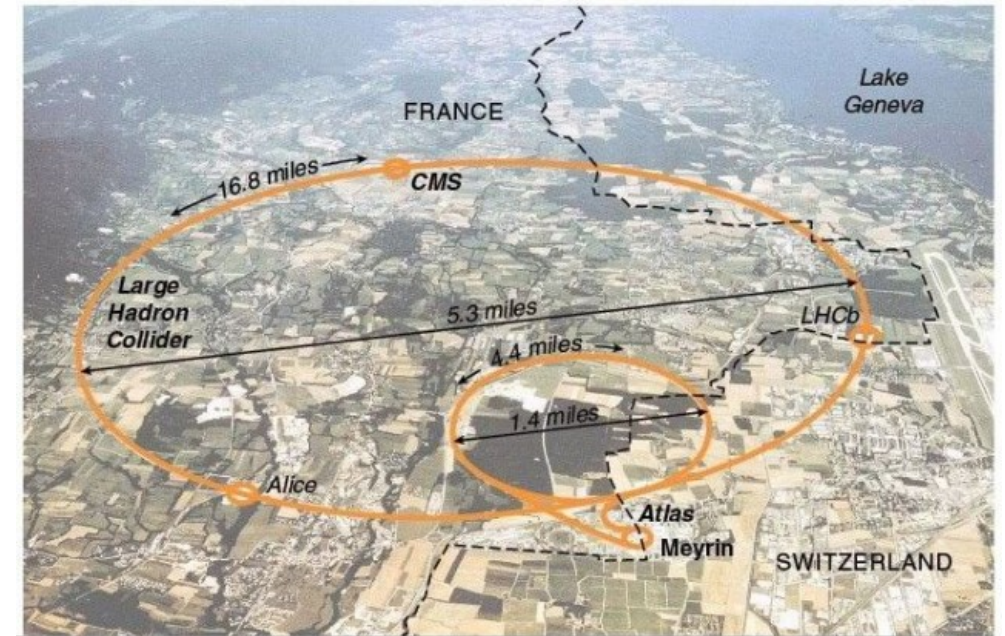
J. Farmer, A. Caldwell, and A. Pukhov



Motivation

Current state-of-the-art for accelerators is the LHC

Size determined by radiation reaction.



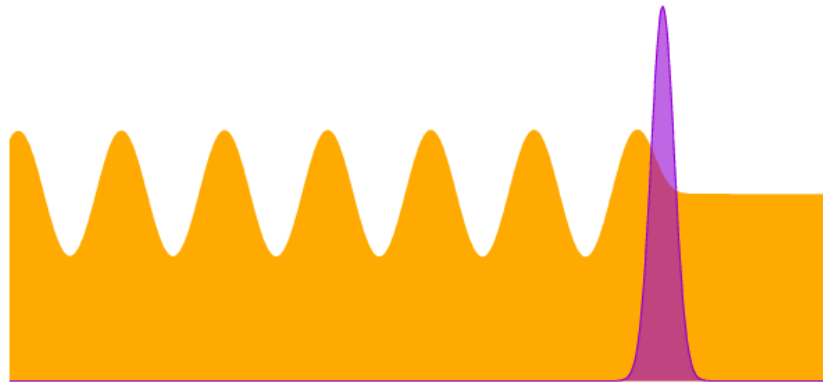
Alternatively, use a linac.

Size determined by acceleration gradient.

Proton-driven PWFA

Use proton driver for plasma wakefield acceleration

- High accelerating gradients
- Plenty of driver energy, no need for staging.
- Protons drive quasi-nonlinear wake suitable (in principle) for positron acceleration

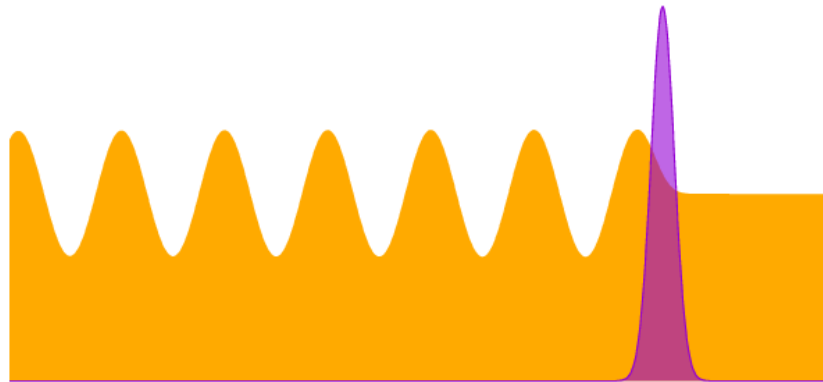


Short driver efficiently excites wakefield

Proton-driven PWFA

Use proton driver for plasma wakefield acceleration

- requires short proton driver



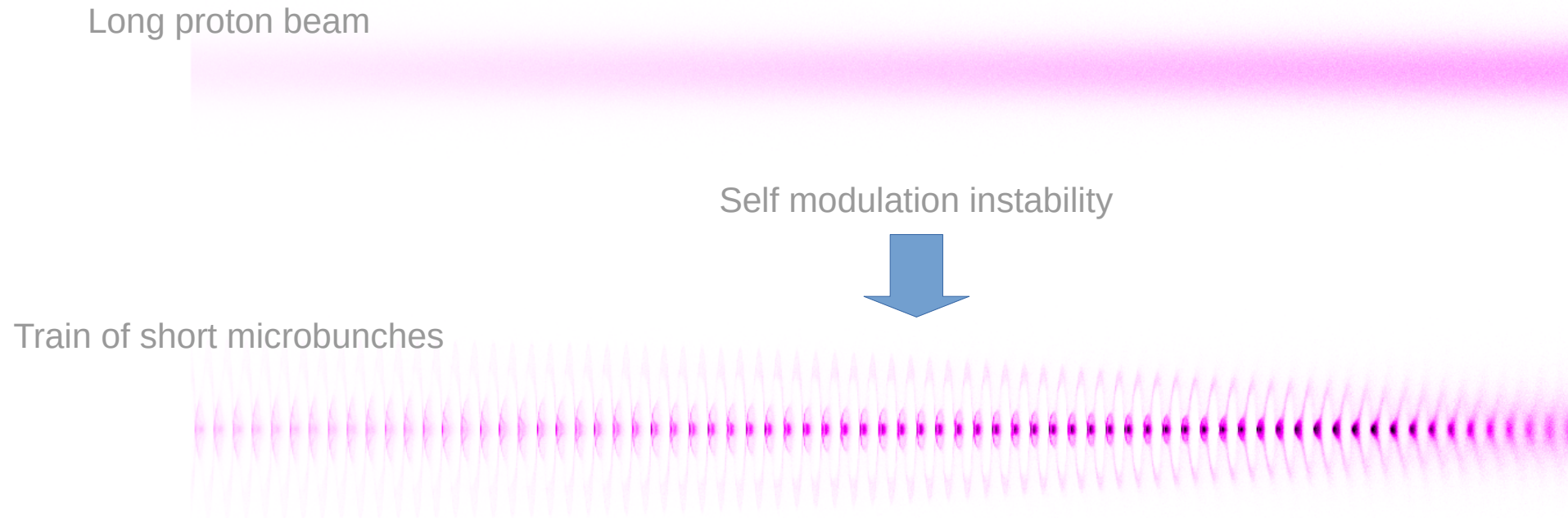
Short driver efficiently excites wakefield



Long driver suppresses its own wake

Proton-driven PWFA

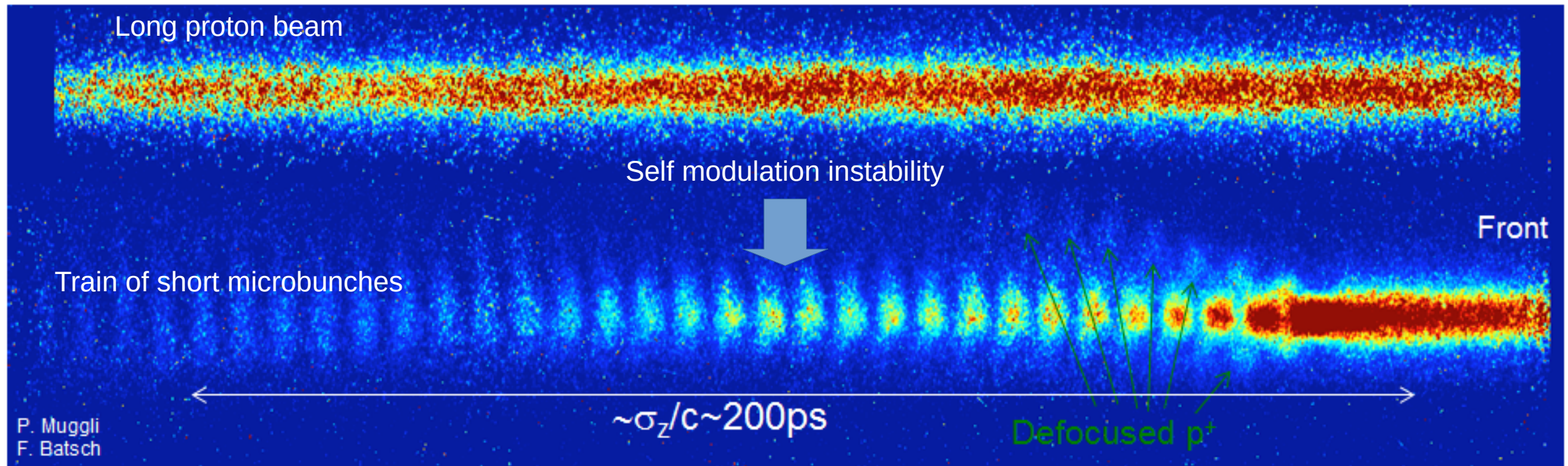
Focussing/defocussing fields in plasma



Resulting train of microbunches can drive large wakefields

Proton-driven PWFA

Not just an idea, this is exactly the basis of *AWAKE*



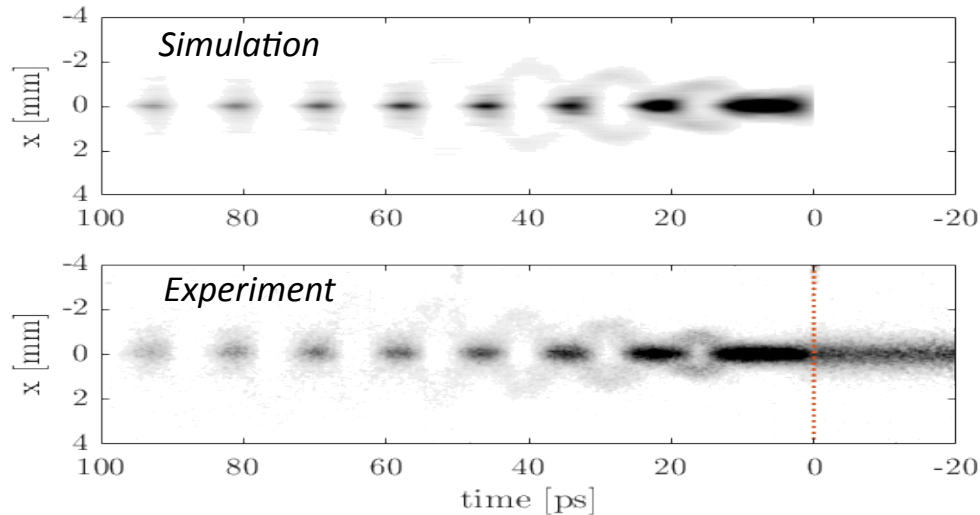
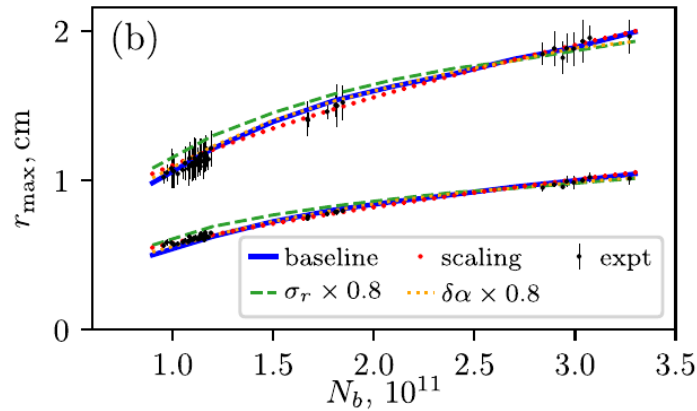
AWAKE Collaboration, Nature (2018)

F. Batsch *et al.* (AWAKE Collaboration), Phys. Rev. Lett. (2021)

Proton-driven PWFA

Beam radius after SMI at Imaging Station 2

A. Gorn, PPCF (2020)

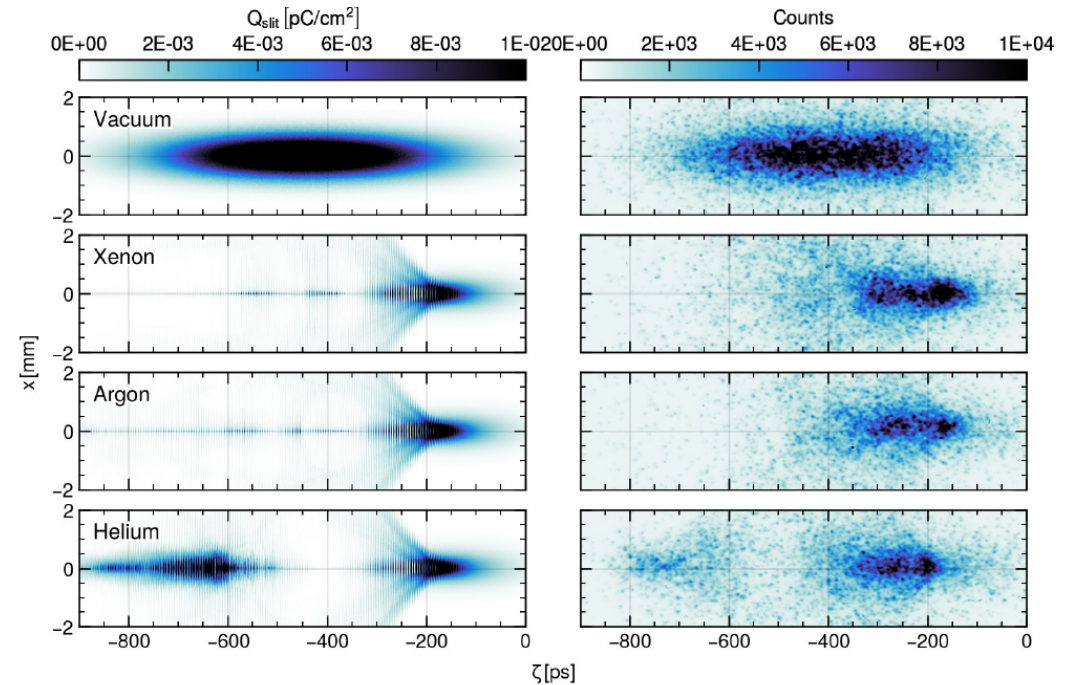


Microbunch structure at streak camera

A.-M. Bachmann, PhD thesis (2021)

John Farmer, MPP

Impact of ion motion on SMI imaged at streak camera



AWAKE preliminary

M. Turner *et al.* (AWAKE Collaboration), *in preparation*

Short proton drivers revisited



It's worth revisiting short proton drivers.

Pros:

Higher gradients

Higher efficiency

Cons:

Such drivers ($L \sim 150 \mu\text{m}$)
don't currently exist

Short proton drivers revisited

nature
physics

ARTICLES

PUBLISHED ONLINE: 12 APRIL 2009; CORRECTED ONLINE: 24 APRIL 2009 | DOI: 10.1038/NPHYS1248

Proton-driven plasma-wakefield acceleration

Allen Caldwell^{1*}, Konstantin Lotov^{2,3}, Alexander Pukhov⁴ and Frank Simon^{1,5}

[Caldwell et al. \(2009\)](#)

A short proton wakefield driver is not a new idea (2009).
Predates AWAKE! So why now?

Short proton drivers revisited

IEEE TRANSACTIONS ON APPLIED SUPERCONDUCTIVITY, VOL. 32, NO. 6, SEPTEMBER 2022

4100404

Record High Ramping Rates in HTS Based Superconducting Accelerator Magnet

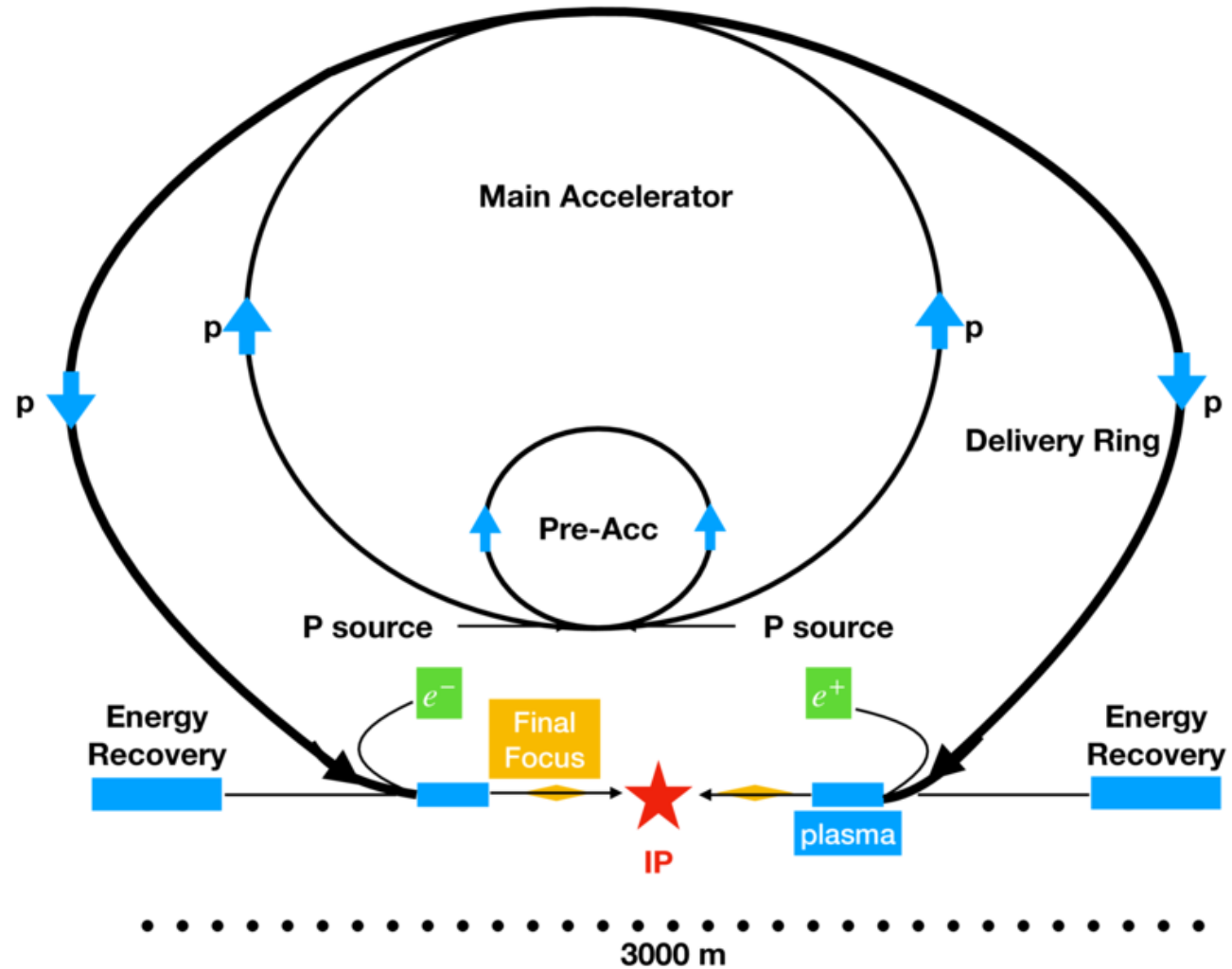
H. Piekarz , Senior Member, IEEE, S. Hays, B. Claypool, M. Kufer , and V. Shiltsev, Fellow, IEEE

[Piekarz et al. \(2022\)](#)

Developments in fast-ramping magnets would allow rapid-cycling (~5 Hz) synchrotrons.

Would allow for competitive luminosities for a proton-driven Higgs factory *if* bunch length can be achieved.

A proton-driven Higgs factory



A proton-driven Higgs factory

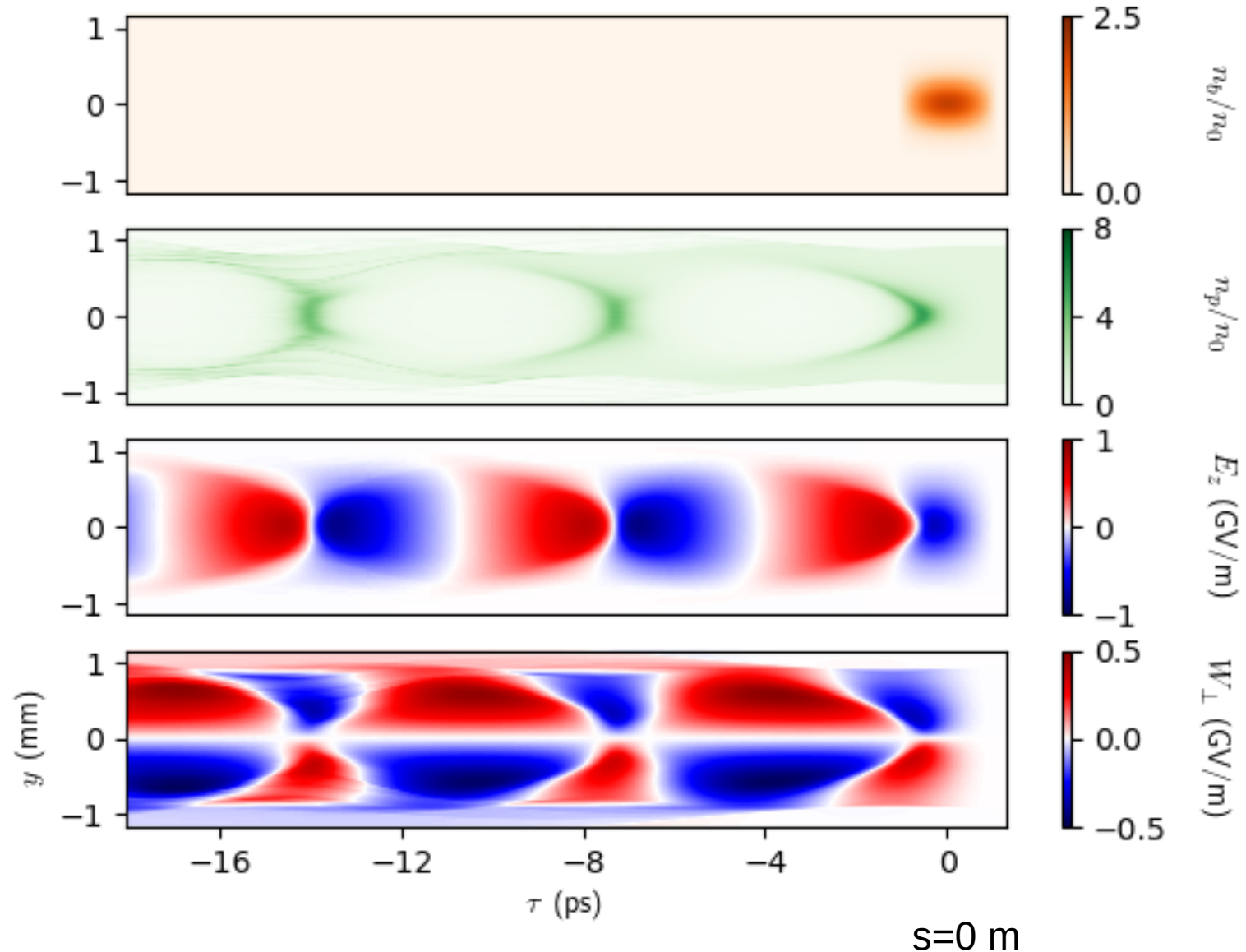
Use plasma wakefields as a transformer:
high energy protons → high energy electrons/positrons

How high? (Symmetric) Higgs factory:
125 GeV e^- colliding with 125 GeV e^+

Configuration

Initial proton driver
chosen to generate
suitable wakefields

Moderately nonlinear
wakefield allows
acceleration of both
electrons and
positrons



A word on lengths

Everything in plasma scales with a characteristic length, the plasma skin depth

$$\frac{1}{k_p} = \sqrt{\frac{c^2 \epsilon_0 m}{n e^2}}$$

We pick $1/k_p = 300 \mu\text{m}$, $n_e = 3 \times 10^{14} \text{ cm}^{-3}$

$$L_p = 150 \mu\text{m}$$

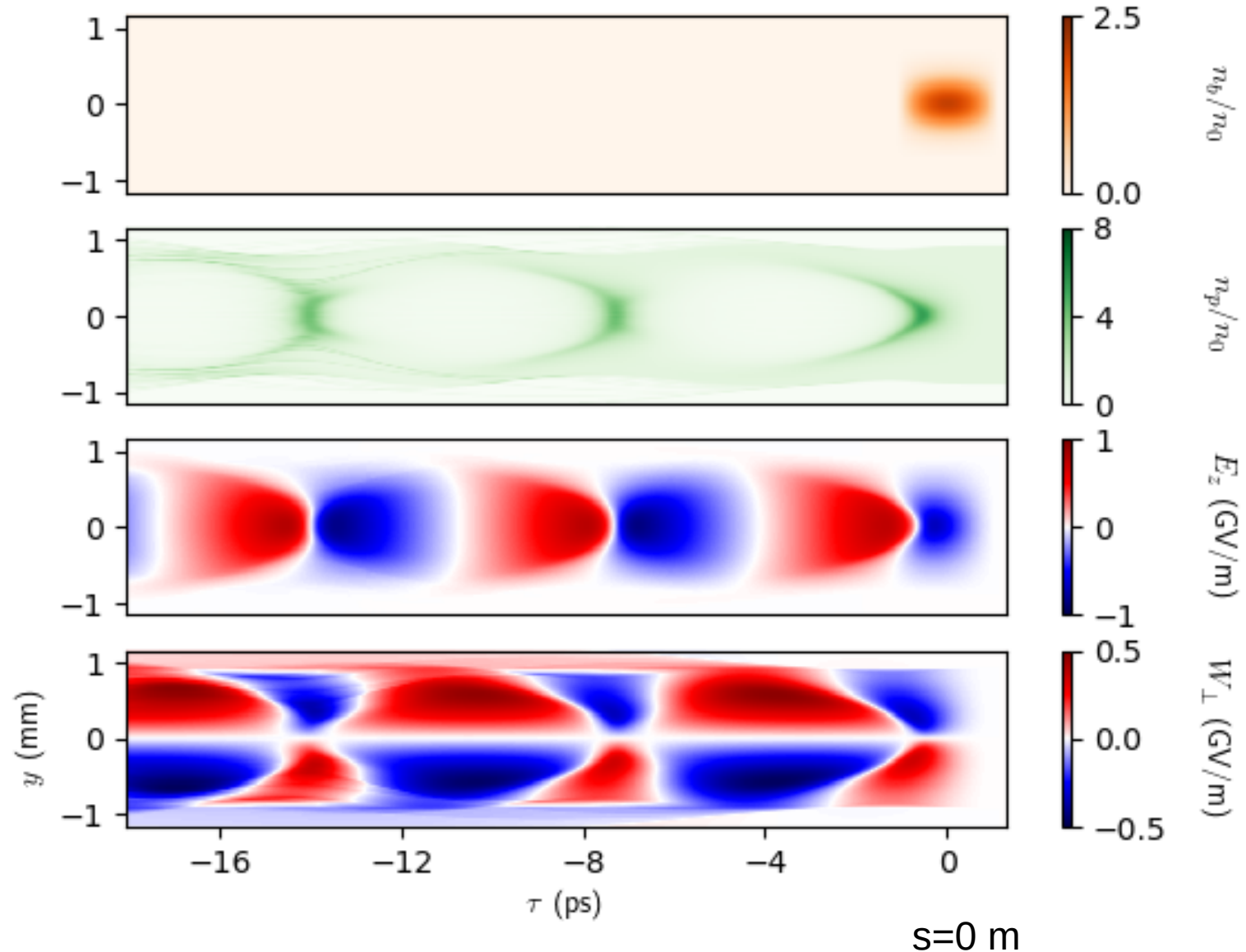
$$N_p = 1 \times 10^{11}$$

$$L_{\text{acc}} \sim 200 \text{ m}$$

Configuration

Initial proton driver
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suitable wakefields

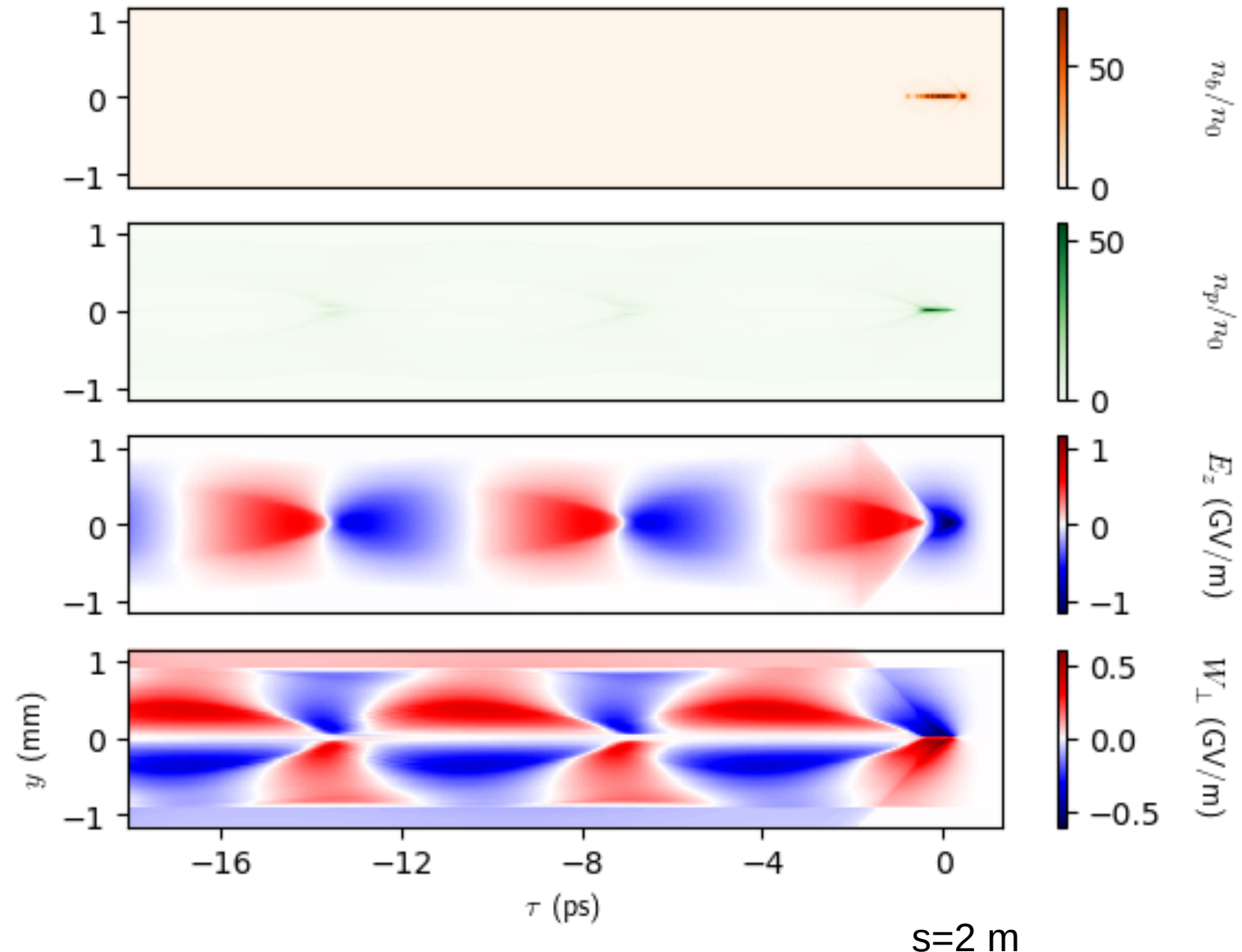
Moderately nonlinear
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Picking the driver: stability

Initial proton driver chosen to generate suitable wakefields

Driver rapidly pinches

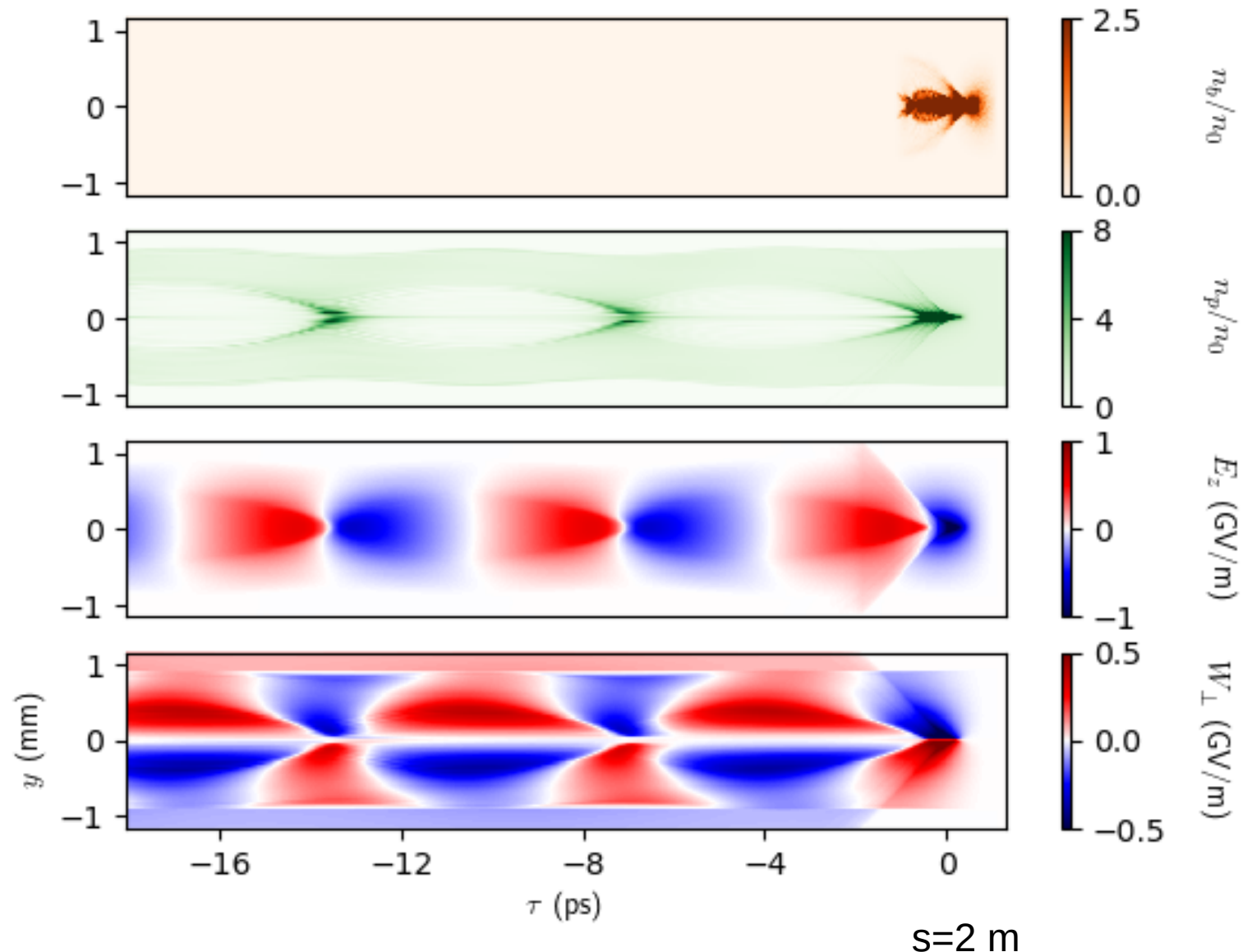


Picking the driver: stability

Initial proton driver chosen to generate suitable wakefields

Driver rapidly pinches

Highly nonlinear wakefields pose several potential problems



Picking the driver: stability

Good initial wakefields not sufficient:

- driver needs to evolve slowly
- counteract strong focussing wakefields

Picking the driver: stability

$$\sigma_z = 150 \mu\text{m}$$

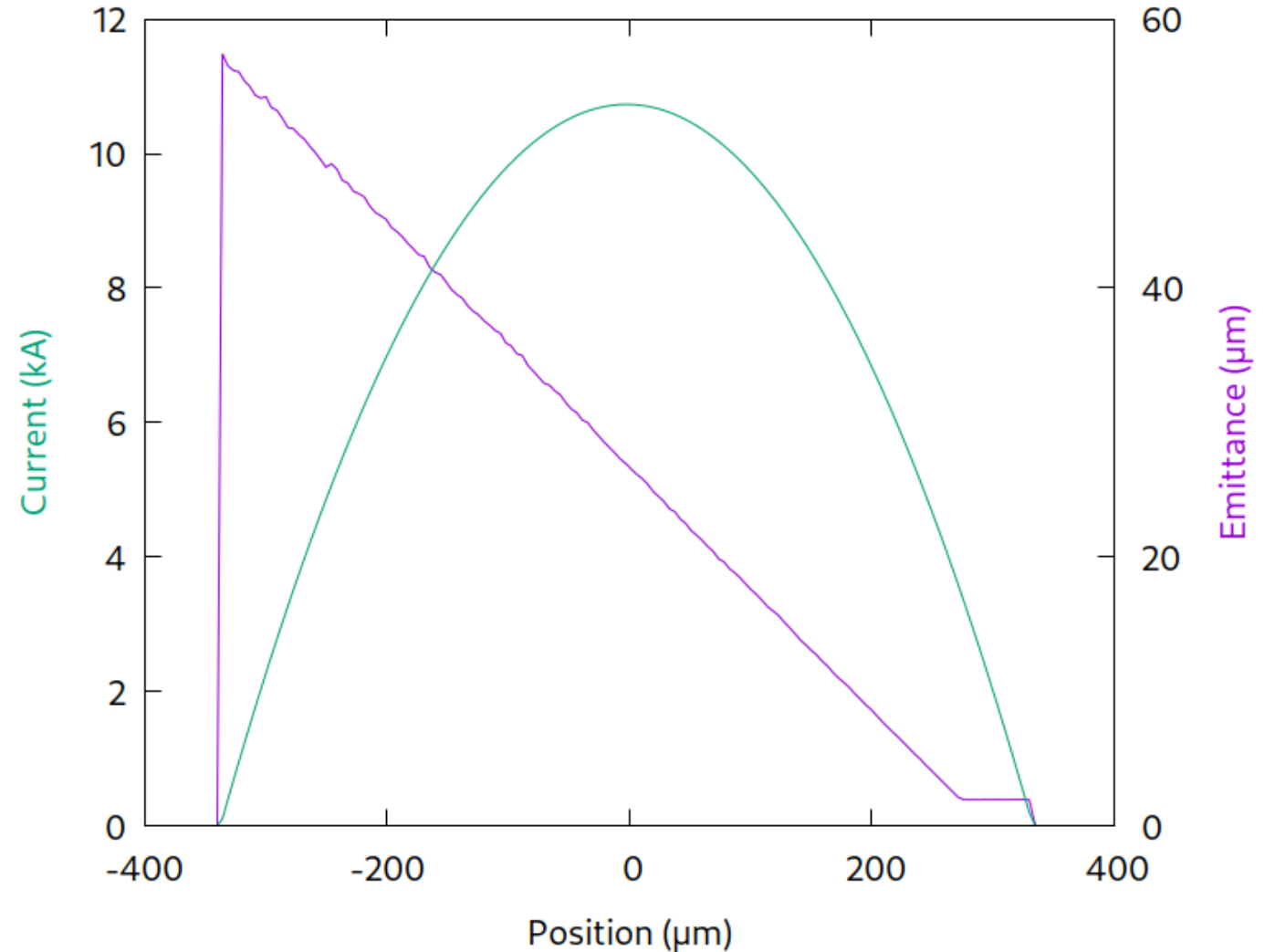
$$\sigma_r = 240 \mu\text{m}$$

$$n_b = 1 \times 10^{11}$$

$$E = 400 \text{ GeV}$$

$$\epsilon_N = \textit{tailored}$$

- 2 μm at head
- initially constant
- rises linearly to 58 μm

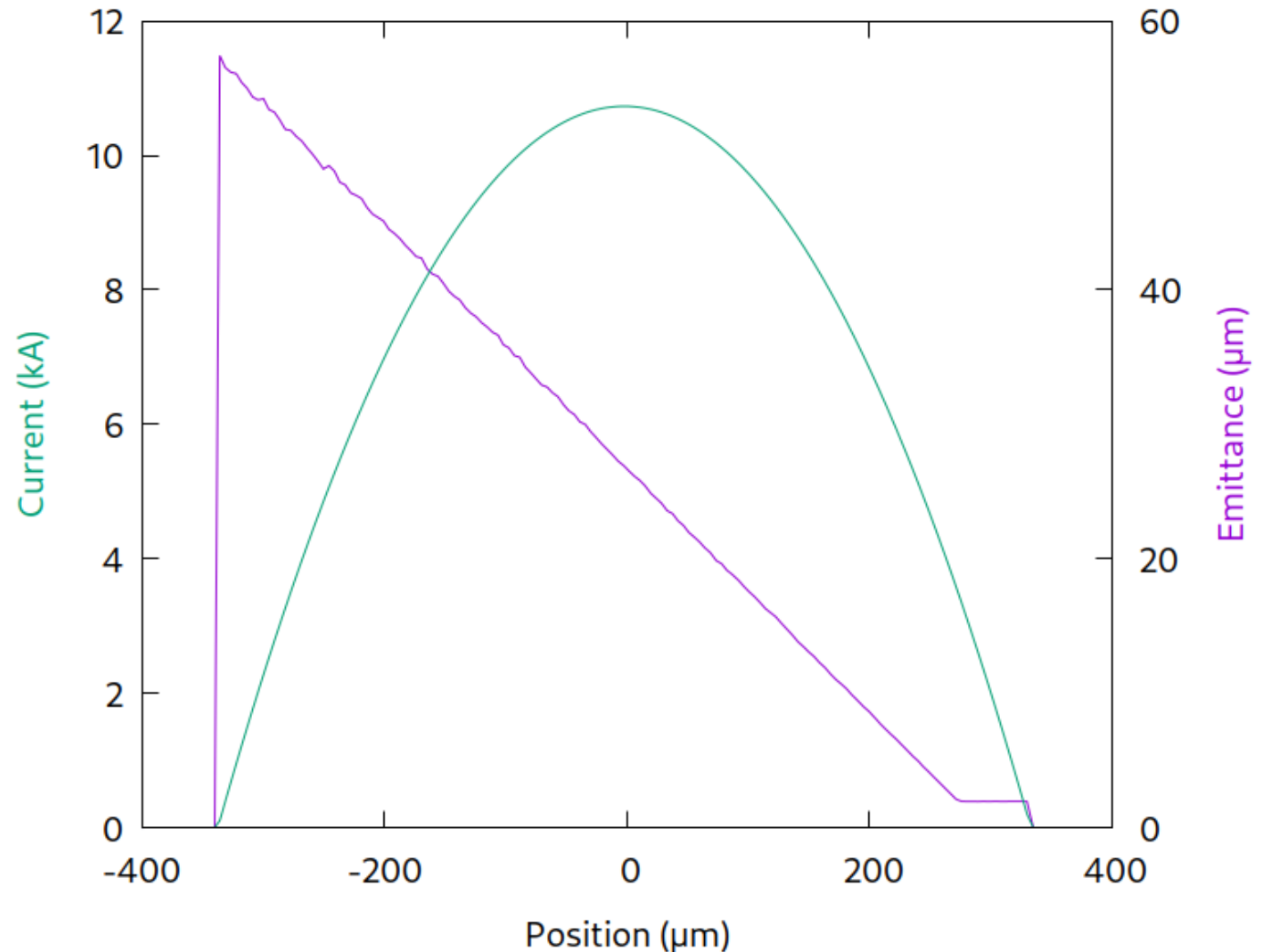


Picking the driver: stability

How can we generate a tailored emittance profile?

Most likely:
with difficulty

BUT emittance is initially constant before growing monotonically

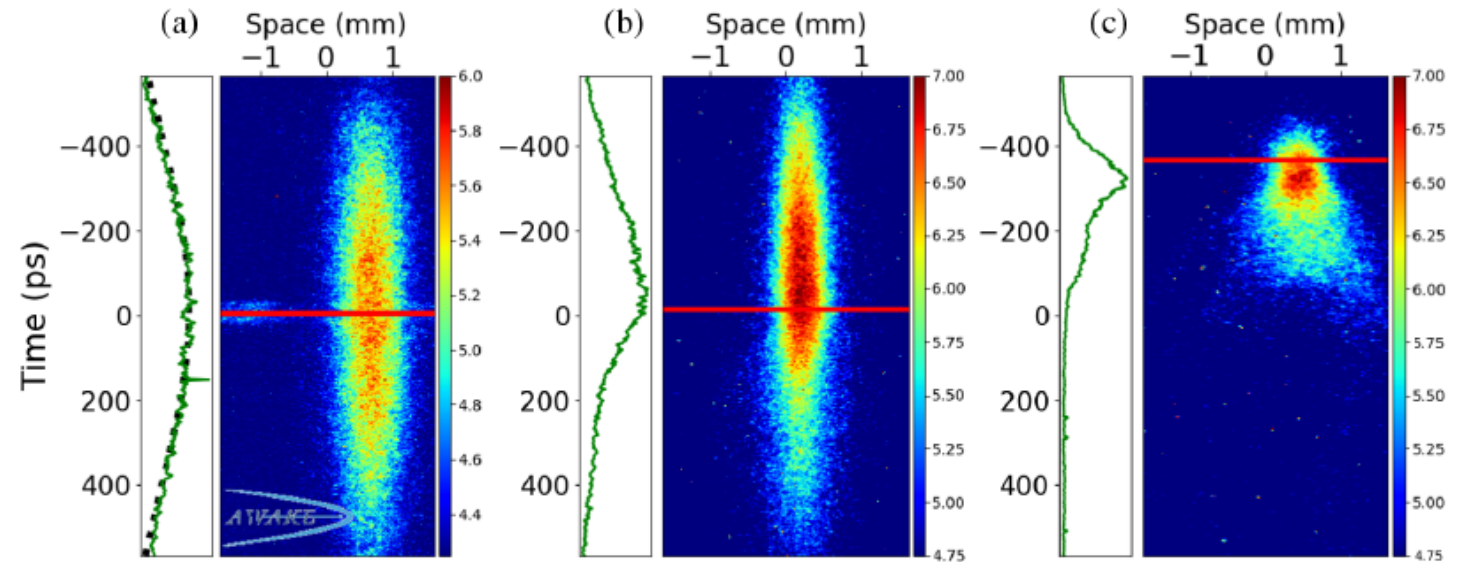


Picking the driver: stability

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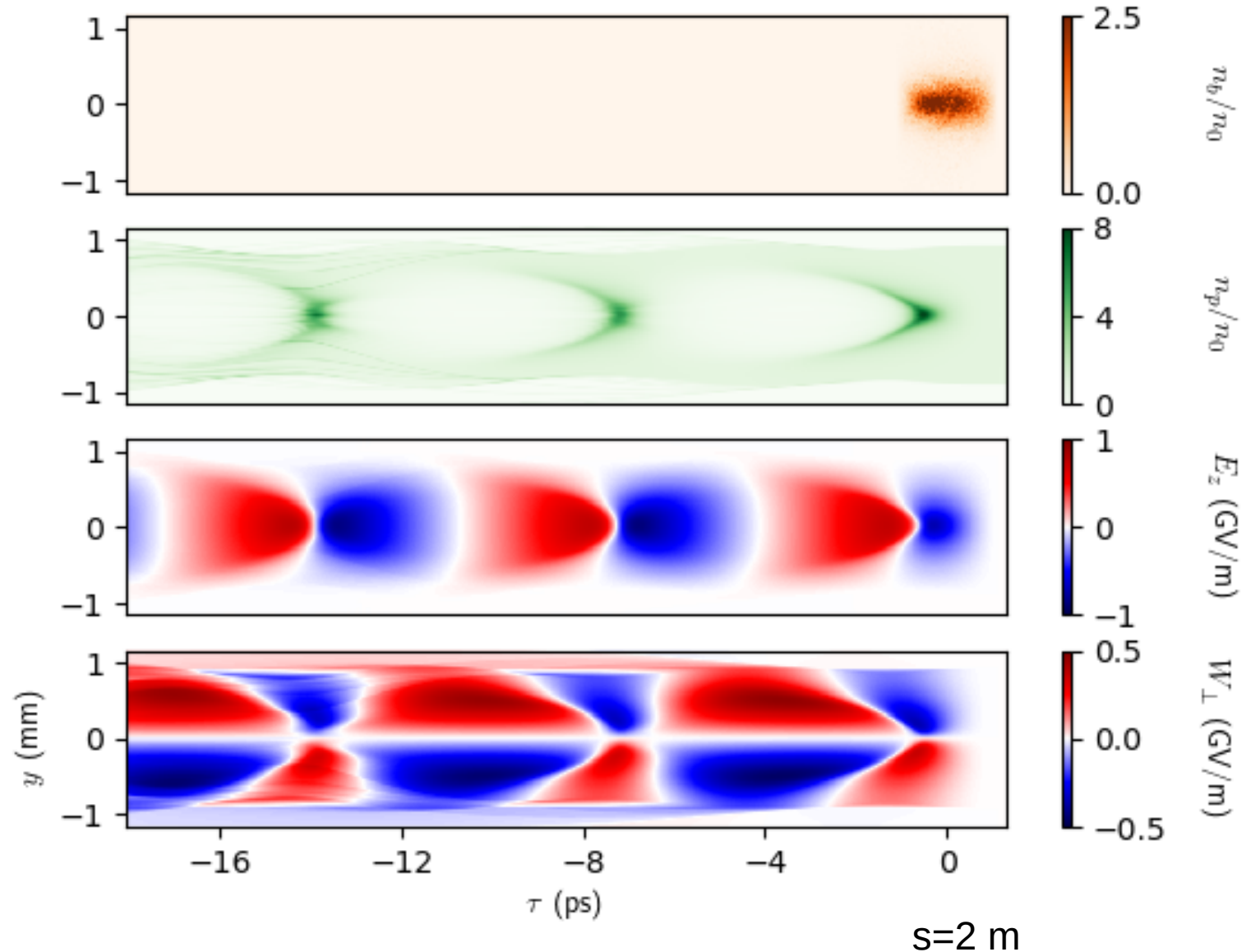
[AWAKE Collaboration, PRL \(2019\)](#)

Harness plasma instabilities?

Acceleration: dephasing

Initial proton driver
chosen to generate
suitable wakefields

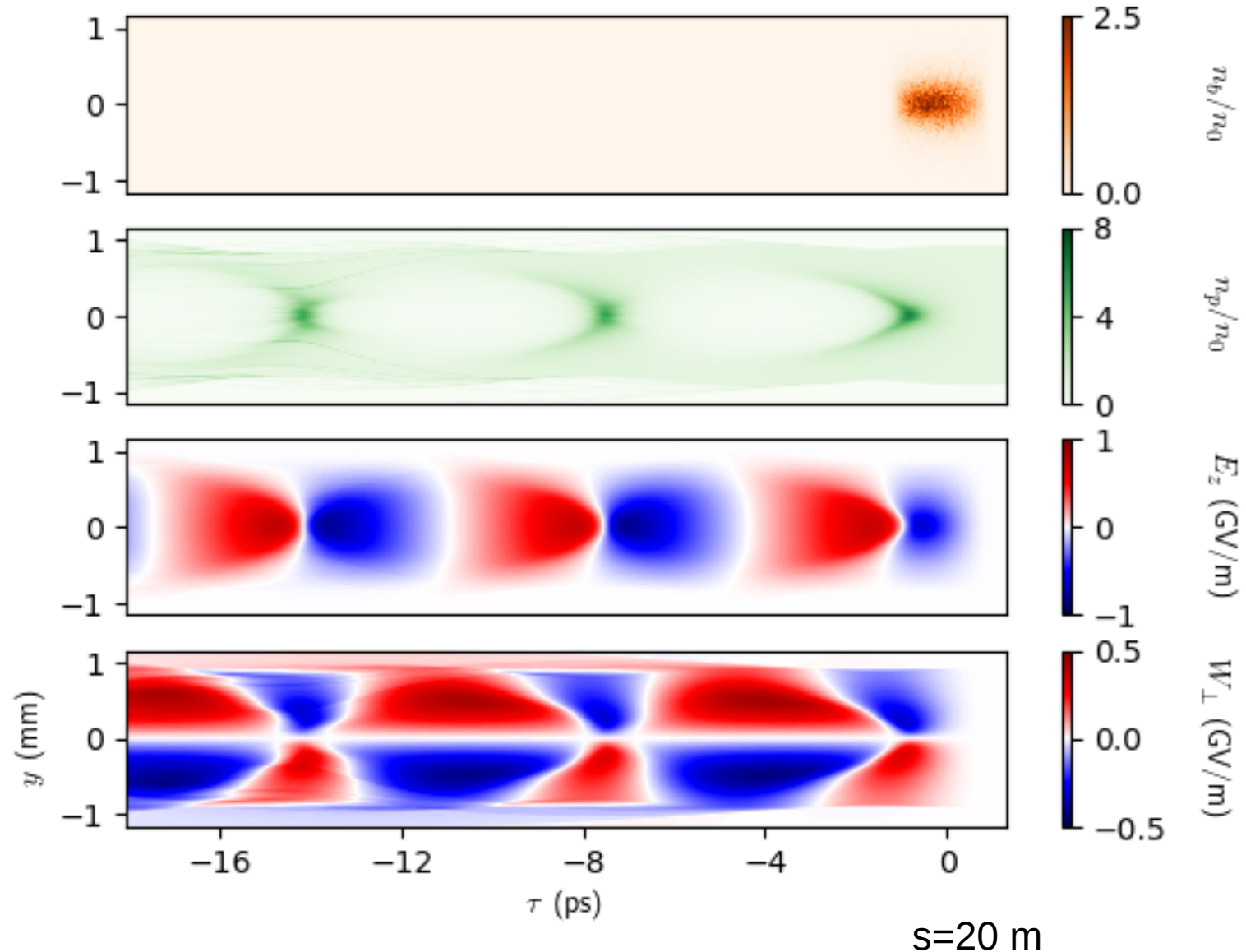
Tailored profile
stable for 2 m



Acceleration: dephasing

Initial proton driver
chosen to generate
suitable wakefields

Tailored profile
stable for 20 m

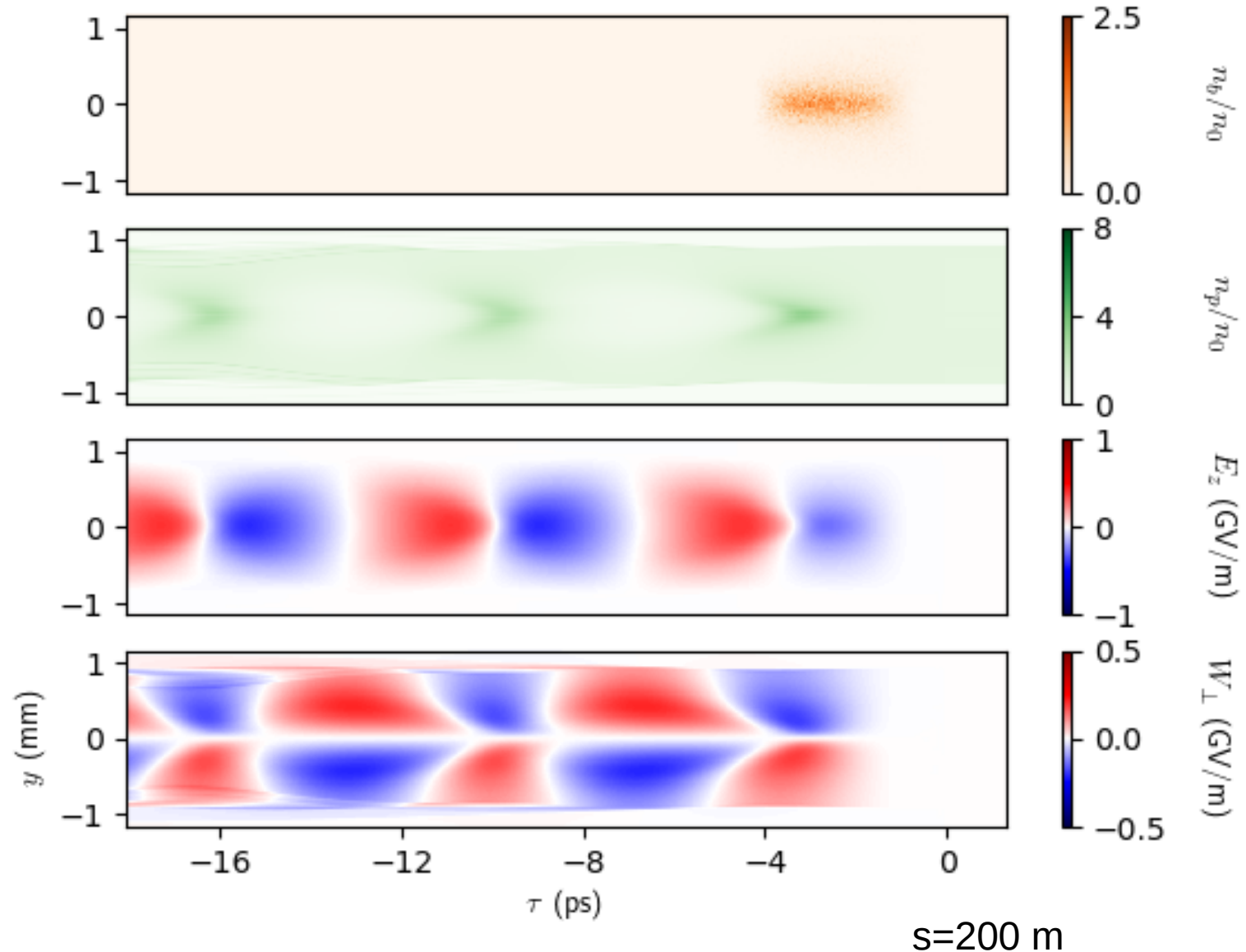


Acceleration: dephasing

Initial proton driver
chosen to generate
suitable wakefields

Tailored profile
stable for 200 m

BUT:
protons “fall back”
in the light frame

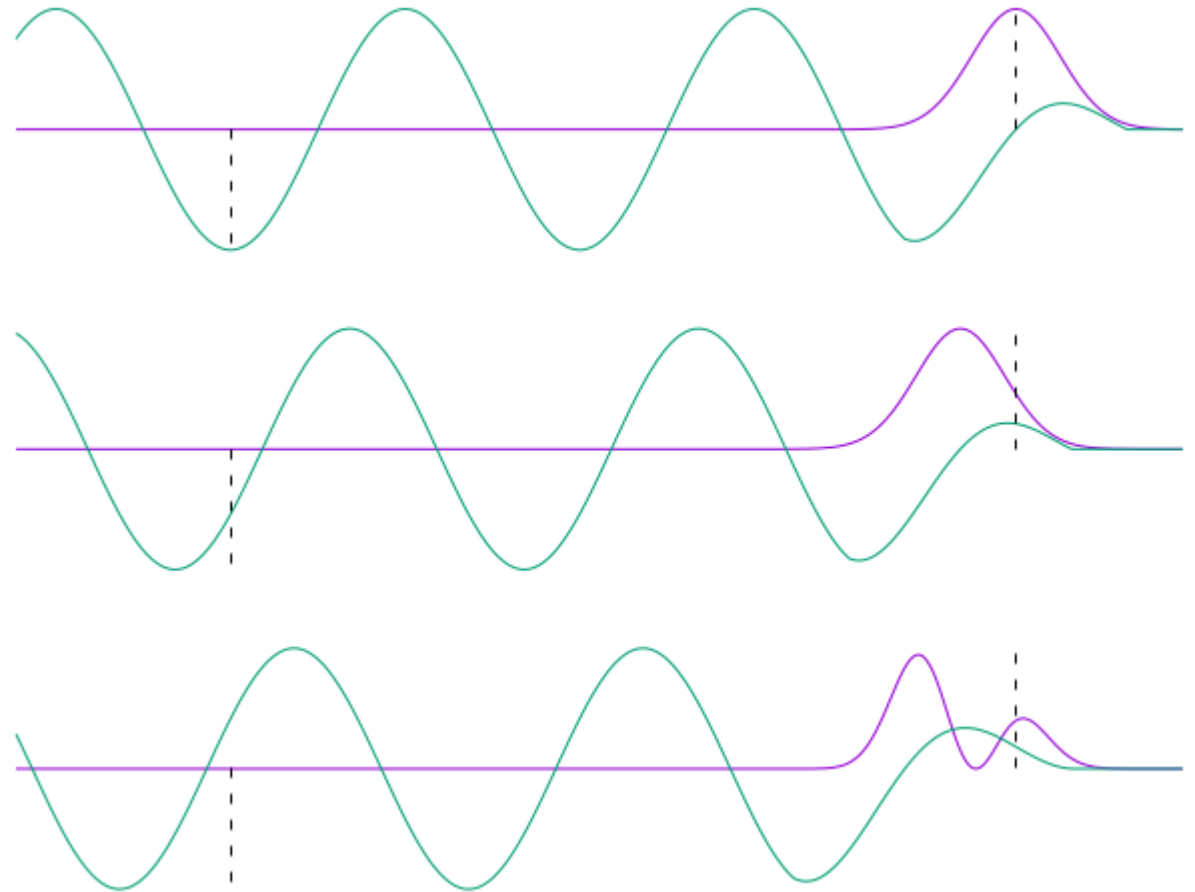


Acceleration: dephasing

Protons are fast,
but not that fast.

Driver evolution
will also modify
wakefield phase.

Witness will “catch up”
with the driver.

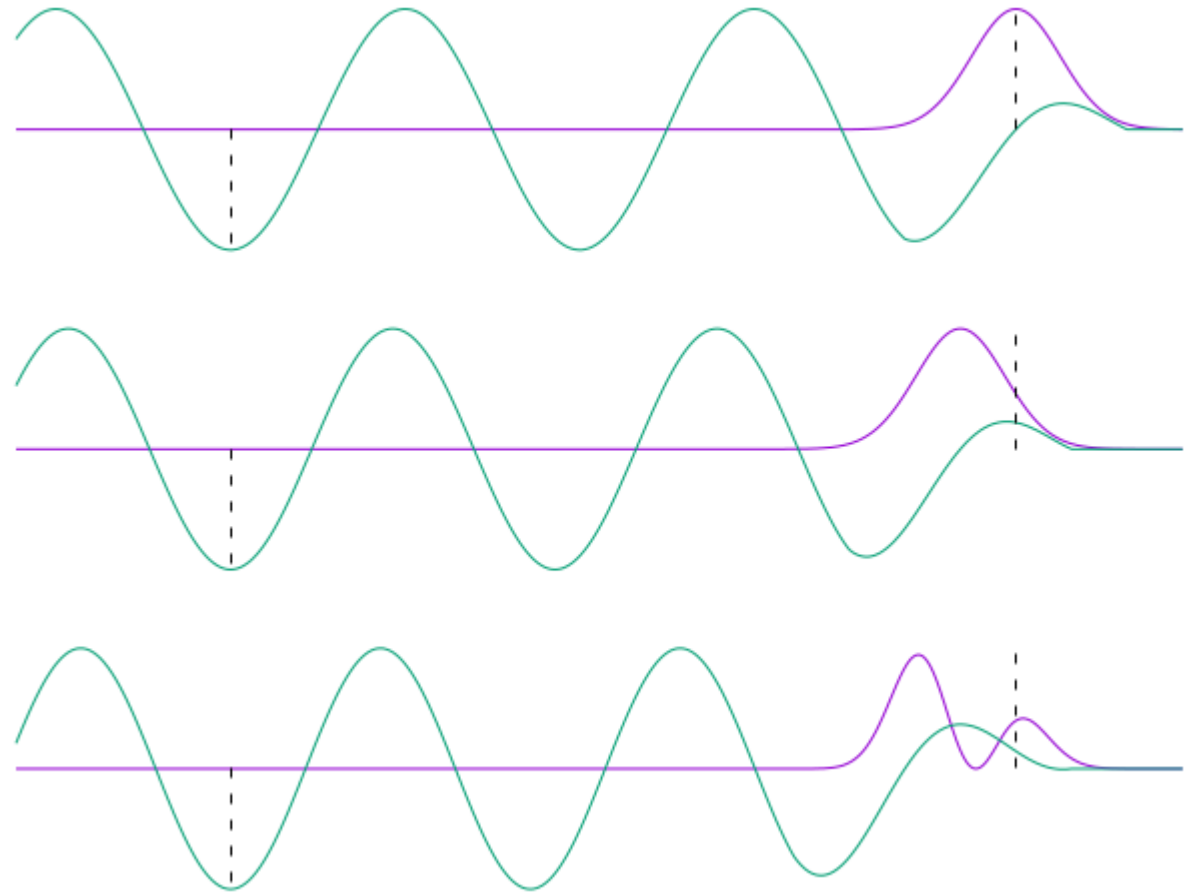


Acceleration: dephasing

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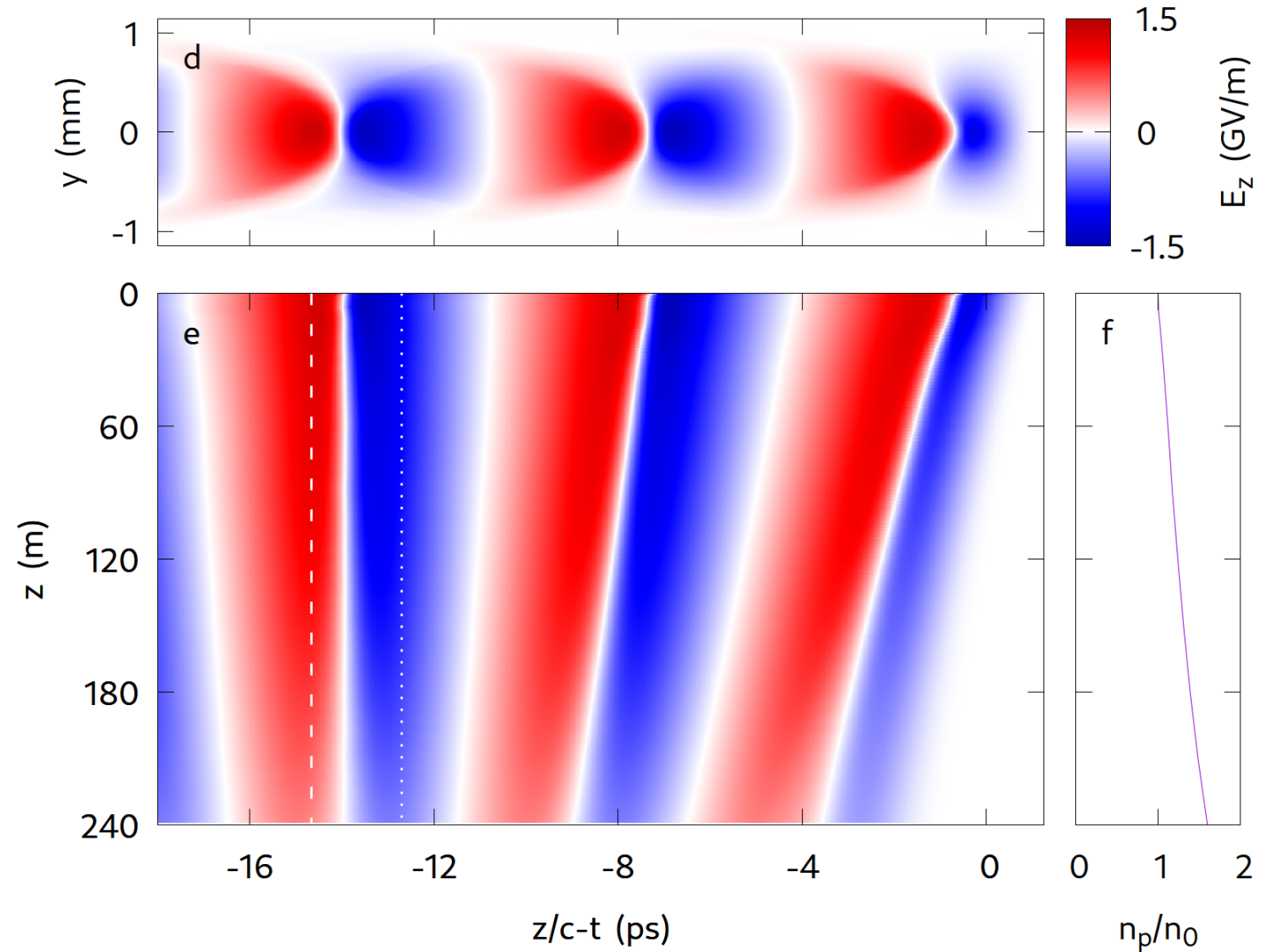
Witness will “catch up”
with the driver.



Change plasma density to keep phase constant

Acceleration: dephasing

Change plasma density to keep phase constant



Acceleration

We now have all the building blocks for Higgs factory

- Large accelerating wakefields
- Regions suitable for electron and positron acceleration
- Stable accelerating phase

Acceleration

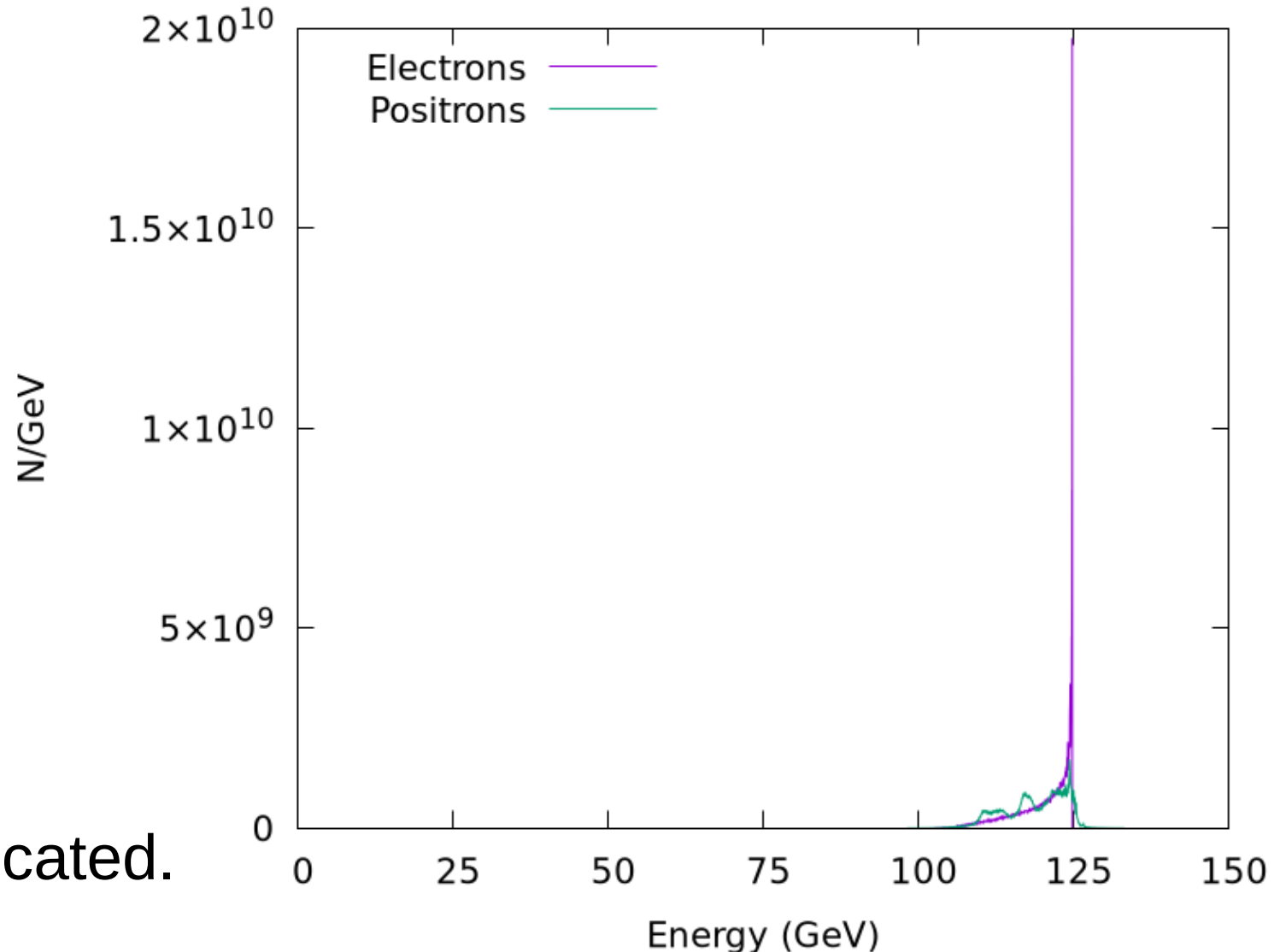
10^{11} protons at 400 GeV

10^{10} electrons/positrons injected at 1 GeV

Electron emittance and slice energy spread are conserved.

Beamloading would give a narrower peak.

Positrons are more complicated.



Acceleration: beamloading

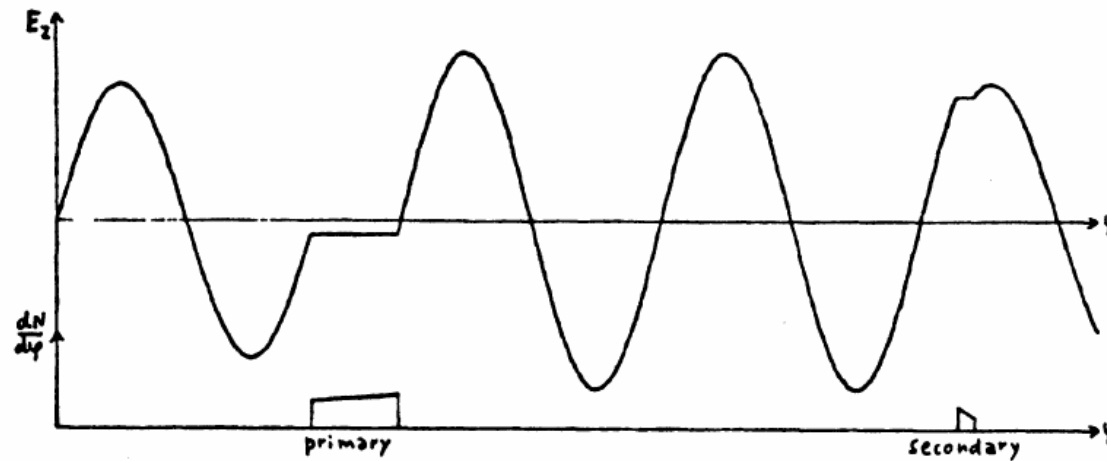


Fig. 3 - Continuous operation. Intervals between primary and secondary bunches may be changed by multiples of 2π .

Van der Meer (1985)

Witness drives its own wakefield, and is subject to the combined wakes of driver and witness

Witness evolution

Plasma provides large accelerating fields, but also large focussing fields.

Emittance preservation relies on choosing the witness radius such that the emittance pressure compensates the focussing.

$$\sigma_r = \sqrt[4]{\frac{2}{\gamma}} \sqrt{\frac{\epsilon_N}{k_p}}$$

- extremely small witness radius at high energy:
0.27 μm for electrons at 125 GeV with 0.1 μm emittance.

Witness evolution

The transverse fields from such a tightly focussed beam are extreme. For our ~ 1.5 kA peak current, $E_r = 150$ GV/m.

Leads to some headaches

- secondary ionization
- ion motion

Witness evolution: ionization

We can use the PPT (1966) ionization model for tunnel ionization in a static field:

$$w_{\text{stat}}(F) = \omega_0 |C_{xl}|^2 (2l + 1) \frac{(l + |m|)!}{2^{|m|} (|m|)! (l - |m|)!} \times \left(\frac{2F_c}{F}\right)^{2\lambda - |m| - 1} \exp\left\{-\frac{2}{3} \frac{F_0}{F}\right\},$$

Witness evolution: ionization

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$$w_{\text{stat}}(F) = \omega_0 |C_{xl}|^2 (2l + 1) \frac{(l + |m|)!}{2^{|m|} (|m|)! (l - |m|)!} \times \left(\frac{2F_c}{F}\right)^{2\lambda - |m| - 1} \exp\left\{-\frac{2}{3} \frac{F_0}{F}\right\},$$

F – applied field

F_0 – critical field of ionization

Everything else: $O(1)$ (C_{xl} from ADK (1985))

Witness evolution: ionization

Critical field varies for atom, ionization level, state:

$$F_{\text{He II}} = 3.6 \text{ TV/m}$$

$$F/F_{\text{He II}} = 0.042$$

Characteristic tunnelling time 0.3 ps.

the length of our witness!

Not impossible, could use higher emittance (lowers luminosity) or lower-density plasma (lowers gradient).

Witness evolution: ionization

Critical field varies for atom, ionization level, state:

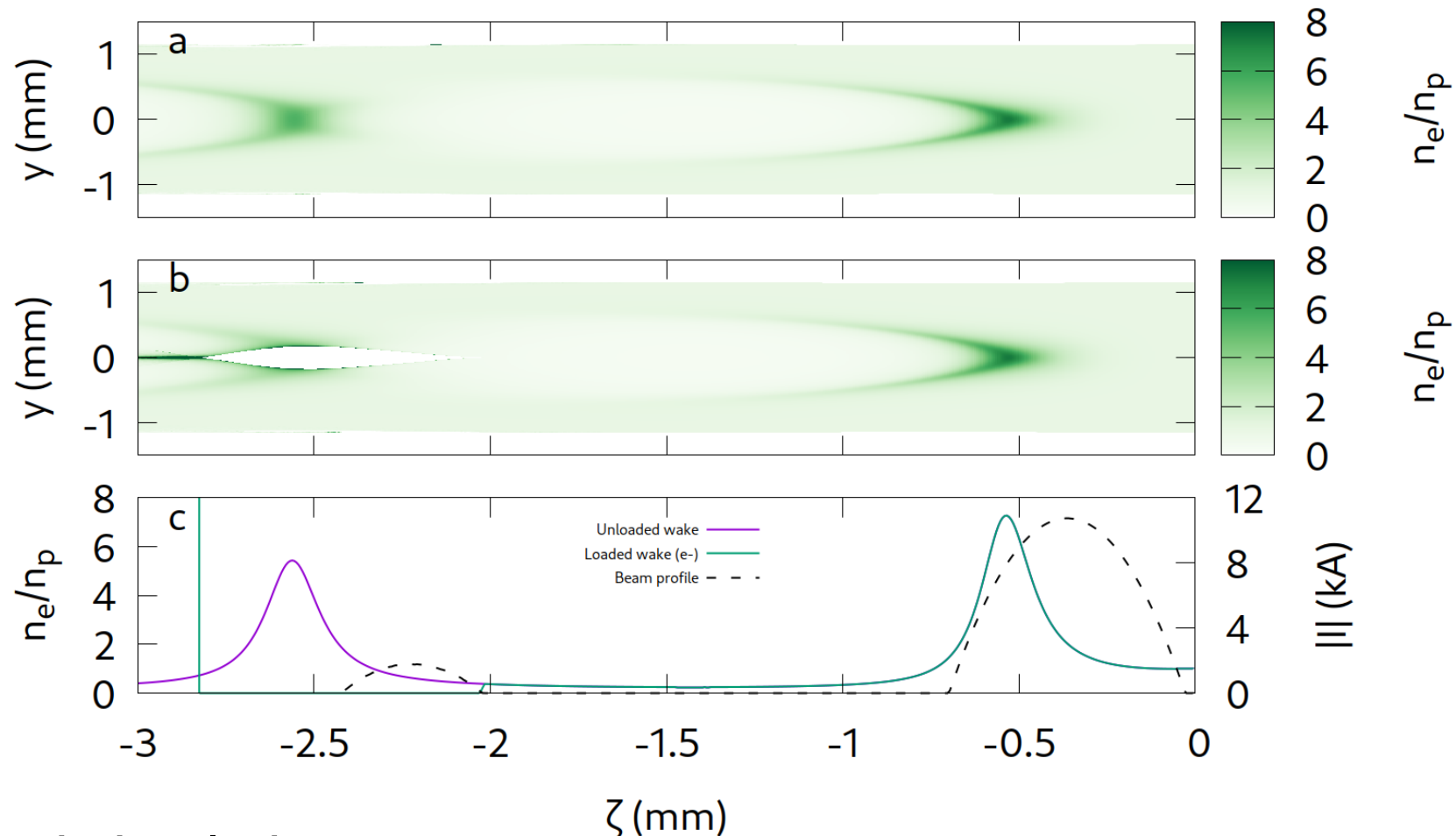
$$F_{\text{Li II}} = 6.7 \text{ TV/m}$$

$$F/F_{\text{He II}} = 0.023$$

Characteristic tunnelling time $0.4 \mu\text{s}$.

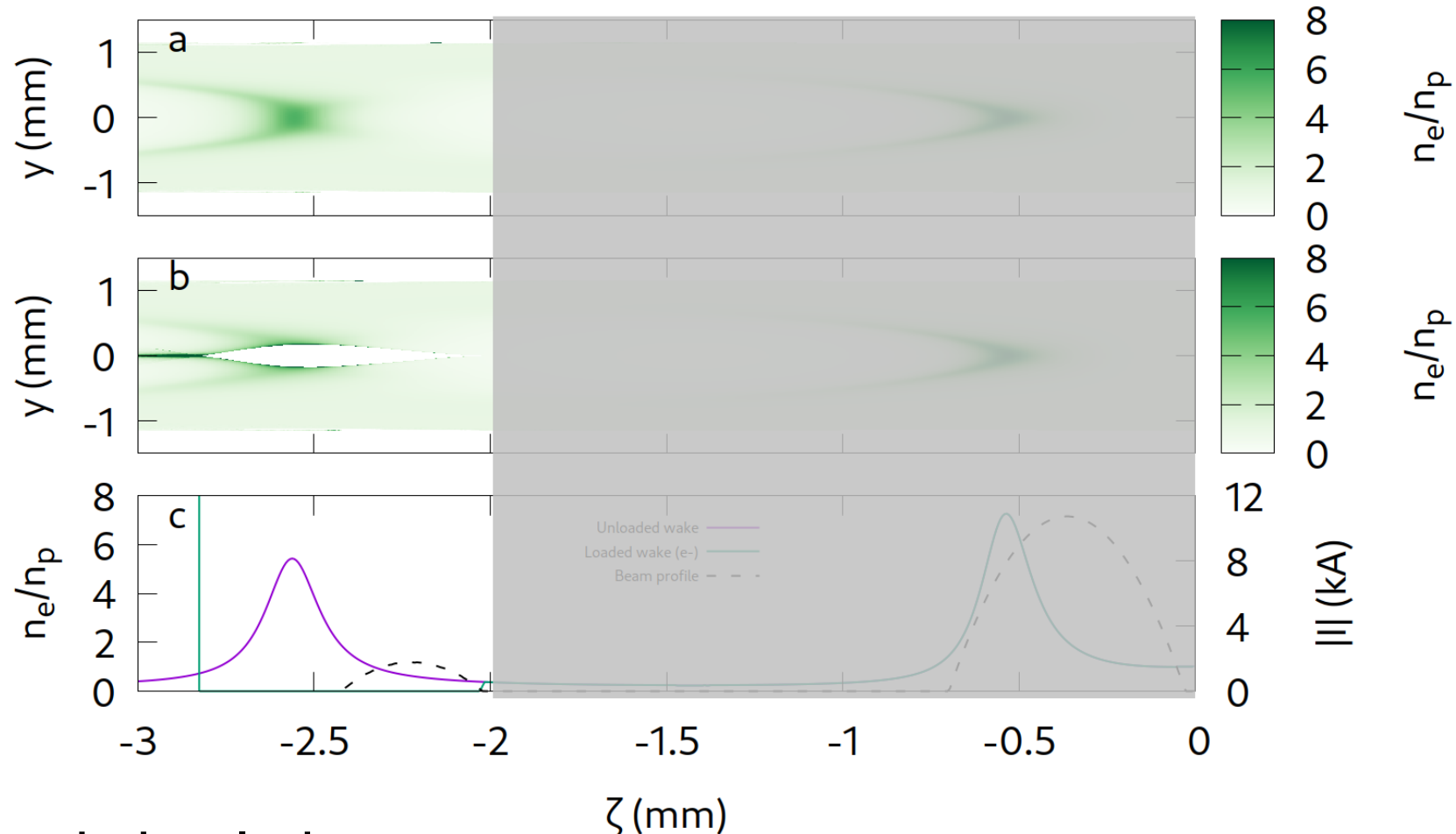
Certainly some challenges to create a lithium discharge, but also need to consider the motion of ions in this field.

Acceleration: ion motion



First resolved simulations:
2D geometry (LCODE), frozen driver, electron witness.

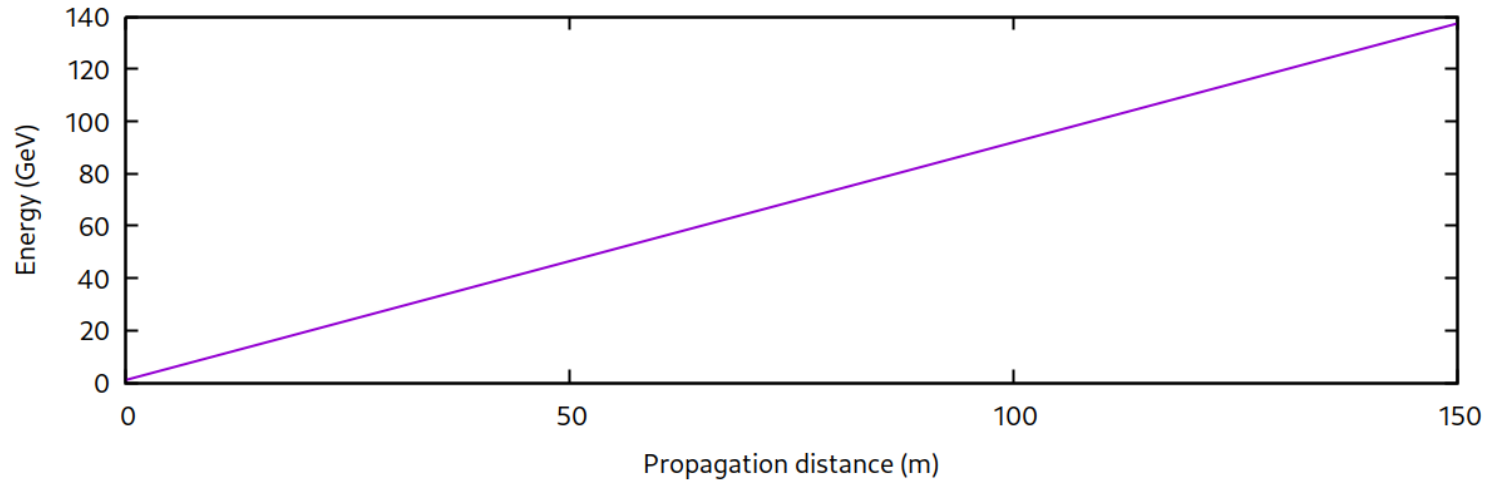
Acceleration: ion motion



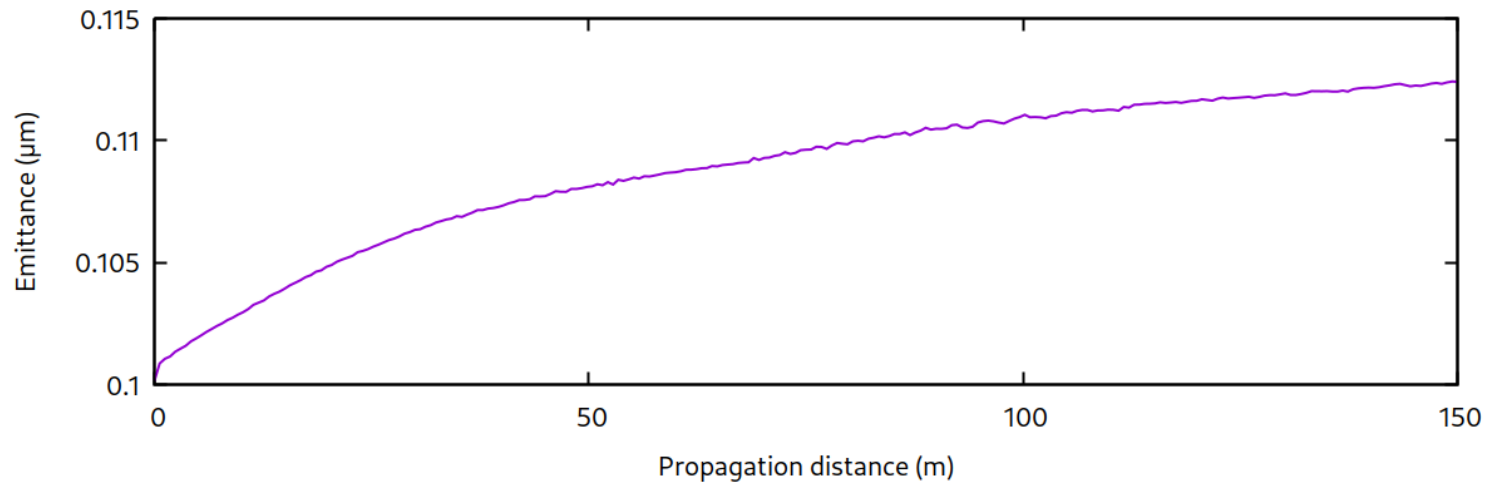
First resolved simulations:
2D geometry (LCODE), frozen driver, electron witness.

Acceleration: ion motion

Energy gain
(trivial for frozen driver)



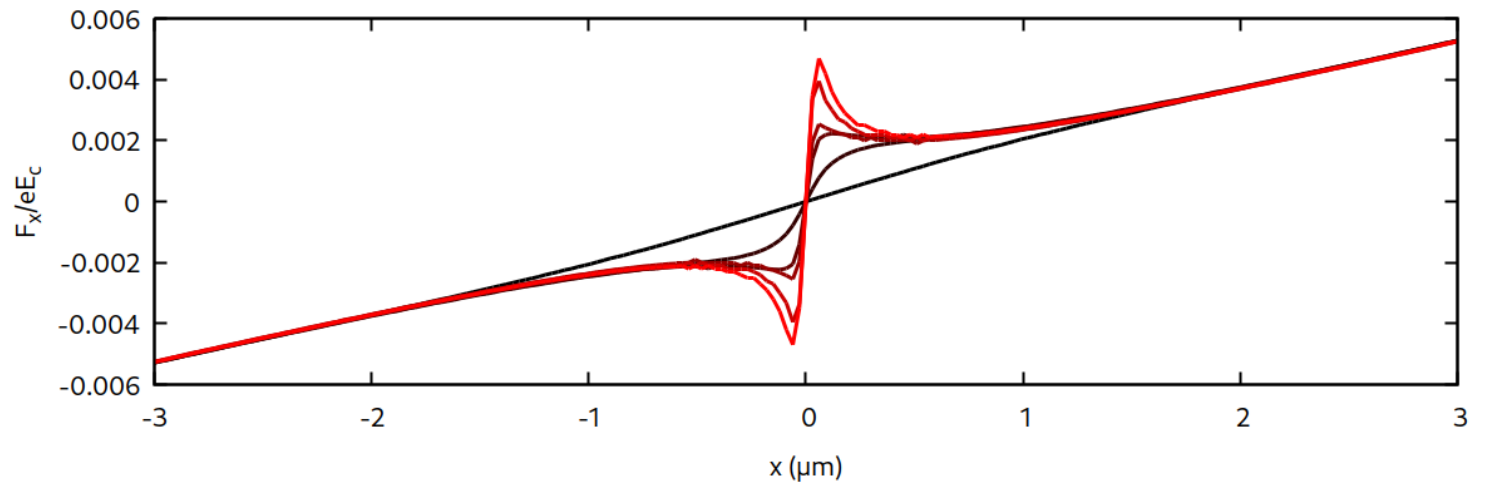
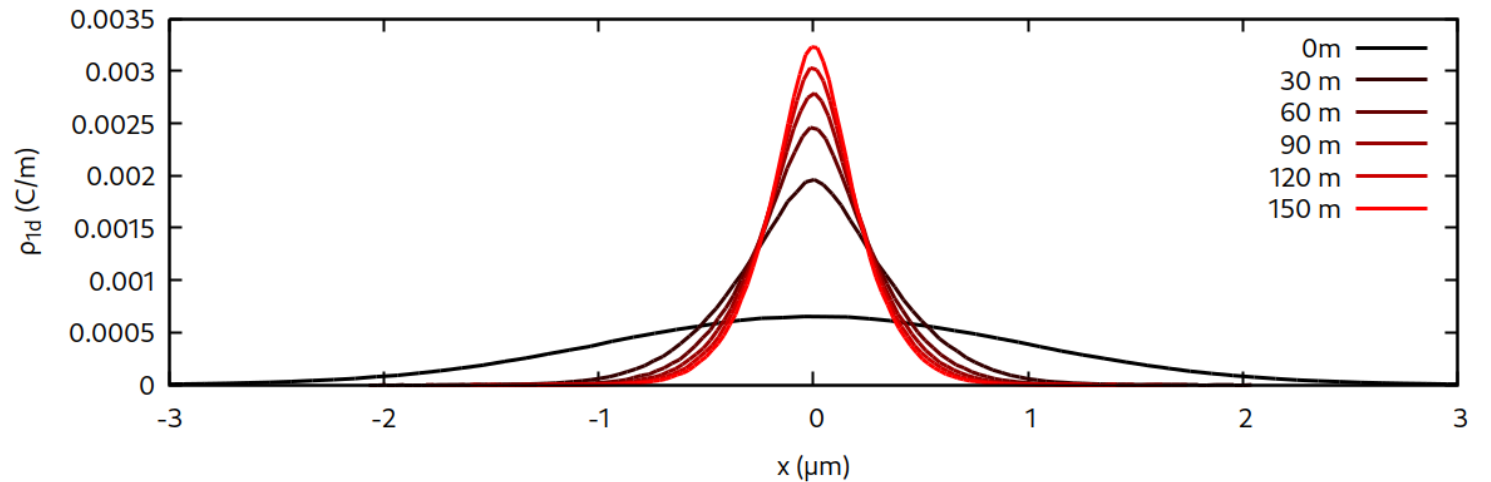
Emittance growth
due to ion motion
(lithium)



Acceleration: ion motion

Adiabatic focussing of witness during acceleration
 $1\mu\text{m} \rightarrow 0.23\mu\text{m}$

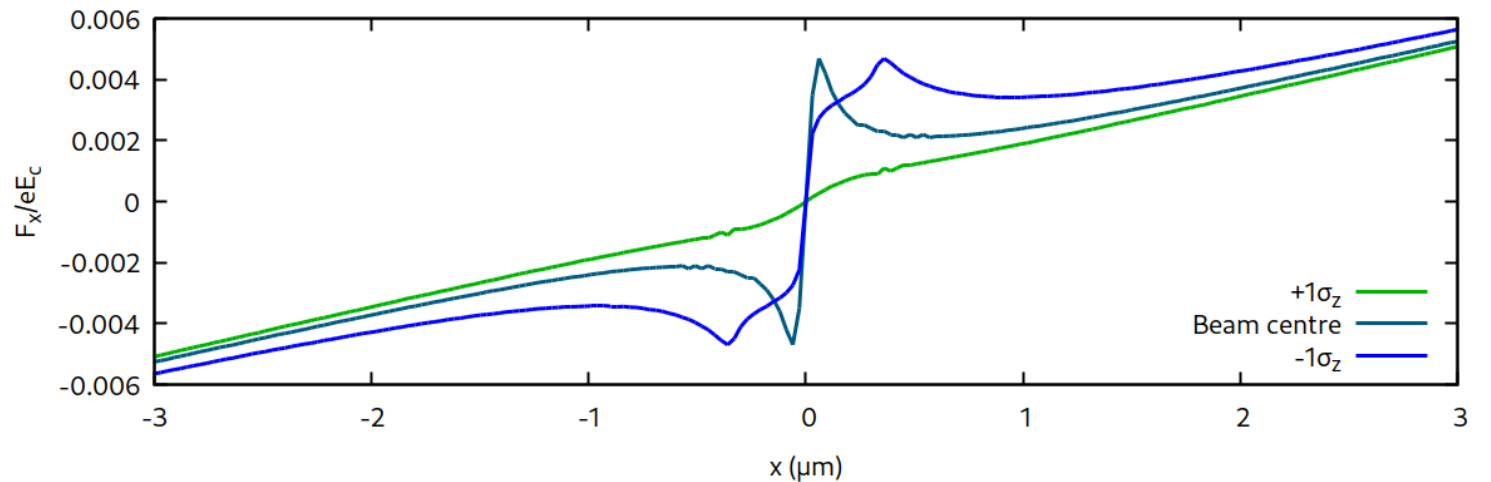
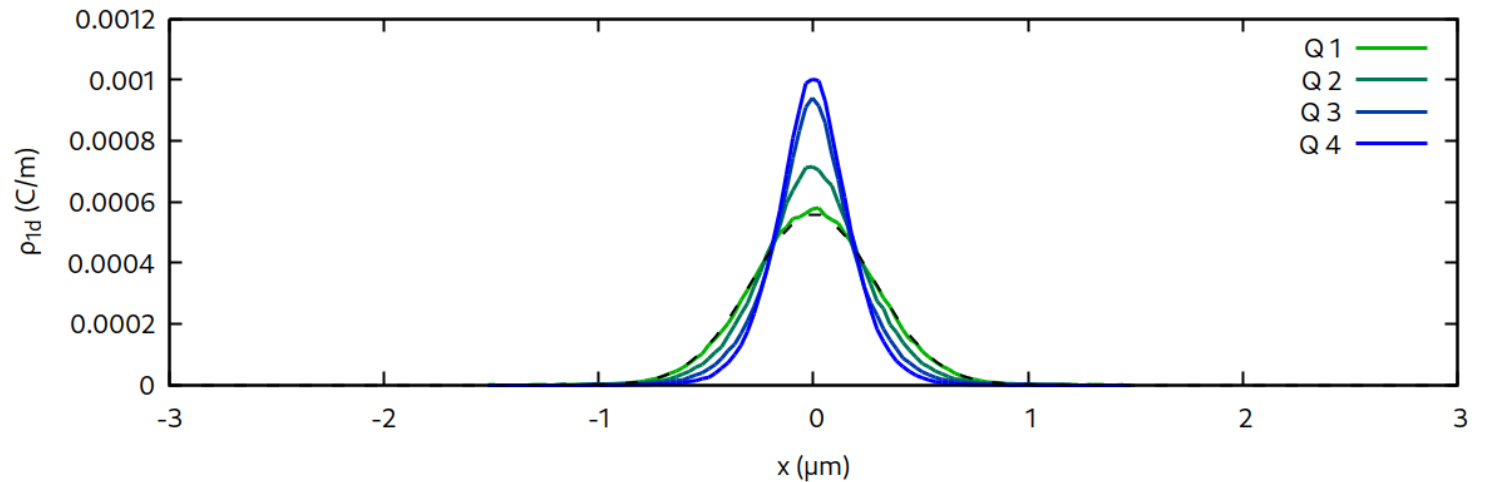
Focussing field becomes increasingly nonlinear



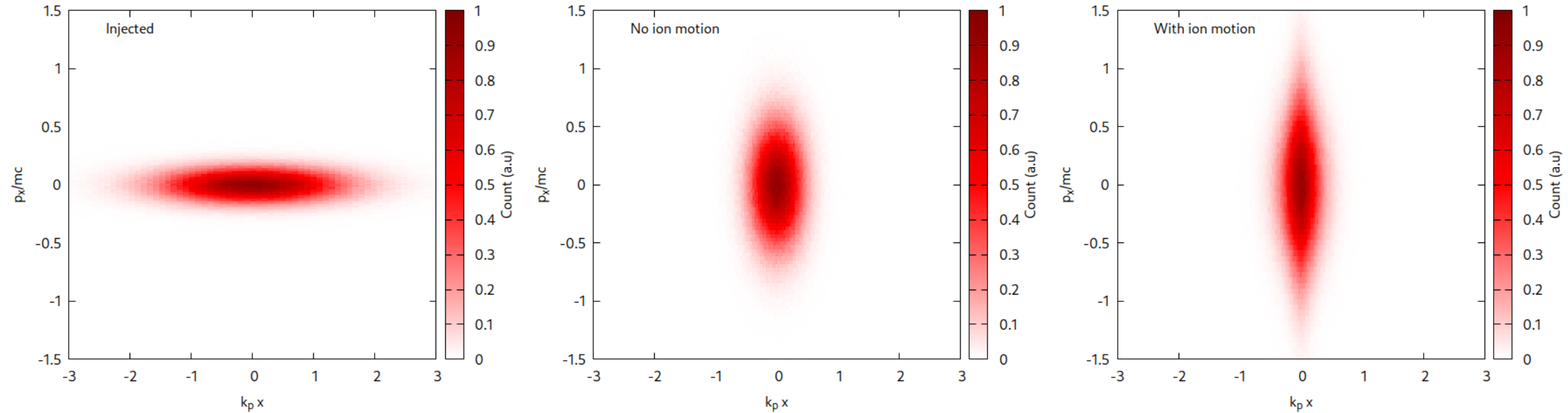
Acceleration: ion motion

Different longitudinal slices of the witness have different profiles

Head-to-tail variation of focussing fields.



Acceleration: ion motion



Suggests adiabatic focussing allows witness to self-match to nonlinear focussing fields

Luminosity

Combine everything:

- Assume proton beams at 5 Hz, with 1000 bunches per beam
- Assume witness beams with 20% driver charge, 100 nm emittance*, ILC optics, and negligible energy spread

*Flat beams should be investigated

and this scheme is competitive:

$$1.7 \times 10^{34} \text{ cm}^{-2} \text{ s}^{-1}$$

Proton Accelerator Parameter	Symbol	Unit	Value
Proton energy	E_p	GeV	400
Refill Time	τ	s	0.2
Bunch population	N_p	10^{10}	10
Number of bunches	n		1000
Longitudinal RMS	σ_z	μm	150
Transverse RMS	$\sigma_{x,y}$	μm	240
Normalized transverse emittance	$\epsilon_{T,p}$	μm	3 – 75 μm
Power Usage	P	MW	150
Plasma Parameters	Symbol	Unit	Value
e^- cell Length	L_{e^-}	m	240
e^+ cell Length	L_{e^+}	m	240
density - upstream	n_p	10^{14} cm^{-3}	3.2
density - downstream	n_p	10^{14} cm^{-3}	5.2
e^\pm Bunch Parameters	Symbol	Unit	Value
Injection Energy	$E_{e,in}$	GeV	1
Final Energy	E_e	GeV	125
Bunch population	N_{e^\pm}	10^{10}	2
Normalized transverse emittance	$\epsilon_{T,e}$	nm	100
Hor. beta fn.	β_x^*	mm	13
Ver. beta fn.	β_y^*	mm	0.41
Hor. IP size.	σ_x^*	nm	73
Ver. IP size.	σ_y^*	nm	13
e^-e^+ Collider Parameter	Symbol	Unit	Value
Center-of-Mass Energy	E_{cm}	GeV	250
Average Collision Rate	f	kHz	5
Luminosity	\mathcal{L}	$\text{cm}^{-2} \text{ s}^{-1}$	1.7×10^{34}

Luminosity

Combine everything:

- Assume proton beams at 5 Hz, with 1000 bunches per beam
- Assume witness beams with 20% driver charge, 100 nm bunch length, and 100 pA current

*Flux

Preliminary Investigation of a Higgs Factory
 based on Proton-Driven Plasma Wakefield Acceleration
 J. Farmer, A. Caldwell, and A. Pukhov
 Now on arXiv

and the resulting luminosity is:

$$1.7 \times 10^{34} \text{ cm}^{-2} \text{ s}^{-1}$$

Proton Accelerator Parameter	Symbol	Unit	Value
Proton energy	E_p	GeV	400
Refill Time	τ	s	0.2
Bunch population	N_p	10^{10}	10
Number of bunches			1000
Longitudinal FWHM		um	150
Transverse FWHM		um	240
Beam size at IP		um	3 – 75 um
Beam size at IP		um	150

Parameter	Value
Witness beam energy	240
Witness beam current	240
Witness beam bunch length	3.2
Witness beam bunch spacing	5.2

Parameter	Value
Witness beam energy	1
Witness beam current	125
Witness beam bunch length	10^{10}
Witness beam bunch spacing	2
Witness beam transverse emittance	$\epsilon_{T,e}$
Hor. beta fn.	β_x^*
Ver. beta fn.	β_y^*
Hor. IP size.	σ_x^*
Ver. IP size.	σ_y^*

e^-e^+ Collider Parameter	Symbol	Unit	Value
Center-of-Mass Energy	E_{cm}	GeV	250
Average Collision Rate	f	kHz	5
Luminosity	\mathcal{L}	$\text{cm}^{-2} \text{s}^{-1}$	1.7×10^{34}

Upgrade path

Witness energy gain limited by dispersion of driver.

Witness energy gain scales as

$$\gamma_W \sim \gamma_D^{3/2}$$

$t\bar{t}$ collider with 525 GeV driver.

HALHF-like 500 GeV electron witness with 1 TeV driver.

Conclusions and outlook

Proof-of-principle simulations for stability and energy gain
(evolving driver, 3D simulations)

Proof-of-principle simulations for emittance control
(electron witness, frozen driver, 2D simulations)

Acceleration to high energy is possible,
moderate driver-to-witness efficiency, clear path to improve this.

Conclusions and outlook

Key challenges:

- Short proton bunches with high rep rate
 - Ferdinand, Jake, today, working session today
- Energy recovery
 - Working sessions today, tomorrow
- PWFA acceleration of positron bunches
 - Alexander, Severin, tomorrow
- Suitable plasma source
 - Nelson, tomorrow
- Beam delivery
 - Vera, tomorrow

~ **Fin** ~

~ Backups ~

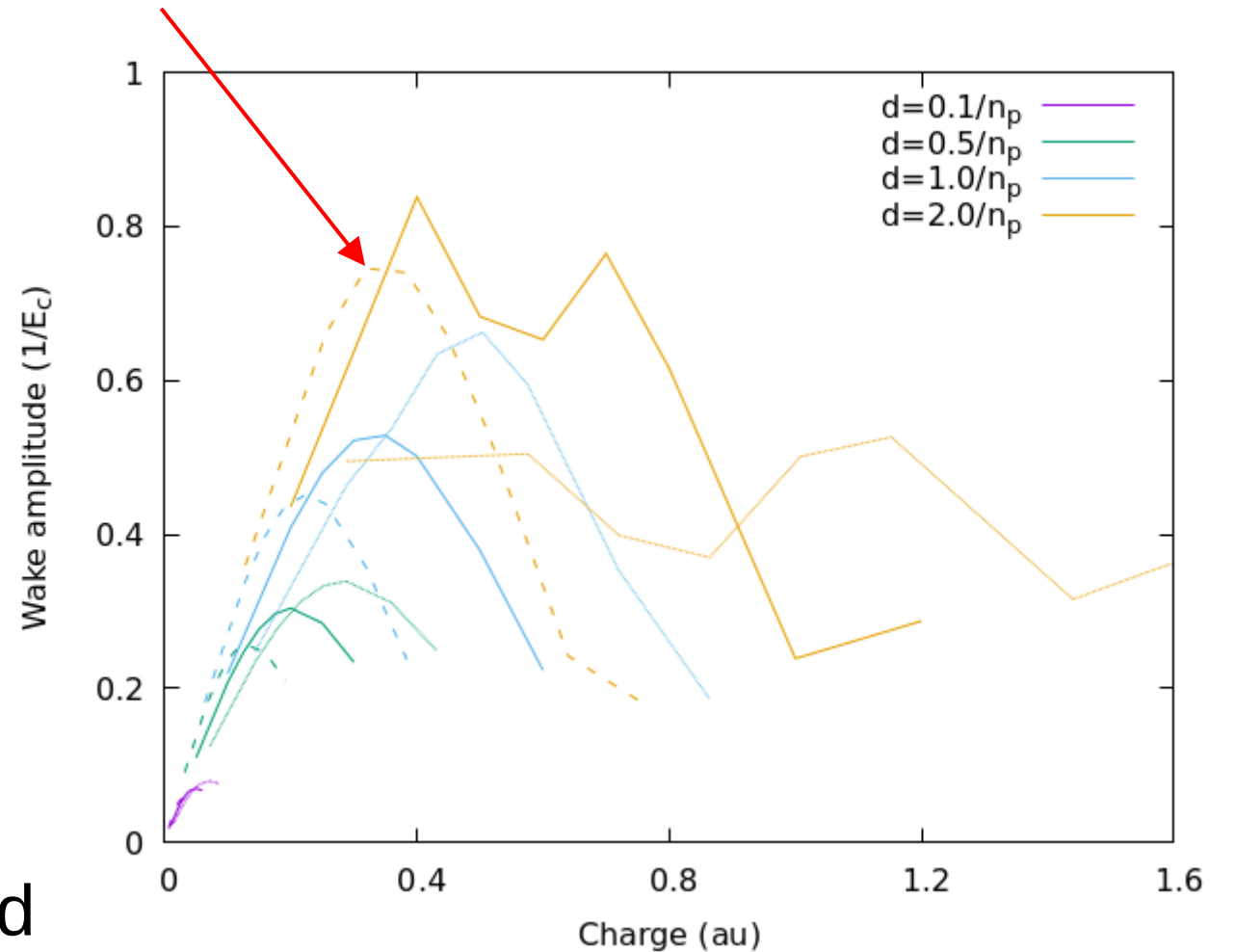
Picking the driver: efficiency

Everything scales with plasma frequency

1×10^{11} protons gives

- plasma density $3 \times 10^{14} \text{ cm}^{-3}$
- driver length $150 \text{ }\mu\text{m}$
- Initial wakefields $\sim 0.8 \text{ GV/m}$

Pick 10% driver energy spread for “realistic” longitudinal emittance



Dashed line: $k_p r = 0.8$
Solid line: $k_p r = 1.0$
Dotted line: $k_p r = 1.2$

Cooling

Witness with 10% driver charge
absorbs ~20% of wakefield energy

Witness with 20% driver charge
absorbs ~40% of wakefield energy

Assume acceleration over 240m,
gives required cooling as 12.5 kW/m

Cooling

Moderately nonlinear wakefields retain their structure after loading.

Could use a second witness bunch to “mop up” excess wakefield