Modified Alfvén wave of multi-ion species in the upper ionosphere of Mars

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- Plasma State
- Plasma Waves
- Magnetosonic Waves
- Alfven Waves
- Theoretical Model
- Application to Observed Structure
- Results
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- Conclusions
- Future work
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Plasma State

- plasma is a quasineutral gas of charged and neutral particles which exhibits collective behavior.
- Conditions: 1. λ_D << L.
 2. N_D>>>1.
 3. ωτ > 1.









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Plasma Waves

- An interconnected set of particles and fields which propagate in a periodically repeating manner.
- These particle are (e⁻)and (i⁺) in the simplest case.



- May contain multiple positive/negative ions and neutral particles.
- The electromagnetic fields in a plasma are assumed to have two parts, one static/equilibrium part and one oscillating/perturbation part.

Plasma Waves

- Continuity Equation: $\frac{\partial n}{\partial t} + \nabla \cdot (nv) = 0$ Eq(1)
- Momentum Equation:

$$mn\left(\frac{\partial \boldsymbol{v}}{\partial t} + \boldsymbol{v} \cdot \nabla \boldsymbol{v}\right) = -\boldsymbol{\nabla} \mathsf{P} + qn(\mathbf{E} + \boldsymbol{v} \times \boldsymbol{B}) \mathsf{Eq(2)}$$



Types of Plasma Waves:

- *Electrostatic* Waves: Governed by *electric fields*; particles oscillate along field lines.
- *Electromagnetic* Waves: Governed by both *electric* and *magnetic* fields.

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Magnetosonic Waves



- Magnetosonic Waves:
 - Compressible
 - Parallel propagation (**Slow** and **Fast**)
 - Perpendicular propagation (Fast)

 Magnetosonic Waves (compressional Alfven waves)

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Alfven Waves



- Alfven Waves:
 - Incompressible
 - Parallel and oblique propagation



Compressional Alfven waves

Shear Alfven waves

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Theoretical Model

- $\mathbf{E} = \mathbf{E}_1 \, \hat{x}$,
- $\mathbf{B} = (B_0 + B_1) \hat{z}$,
- $\mathbf{K} = k \ \hat{y}$.

where(0) and (1) stand for the unperturbed and the perturbed quantities.

• Using the multi-fluid model, the set of governing equations for each species *j* are

$$m_{j}n_{j}\frac{\partial v_{j}}{\partial t} = q_{j}n_{j}(\mathbf{E} + v_{j} \times B), \qquad \text{Eq(6)}$$

$$\nabla \times \mathbf{E} = -\frac{\partial B}{\partial t}, \qquad \text{Eq(7)}$$

$$\nabla \times \mathbf{B} = \mu_{0}\mathbf{J}_{j} + \mu_{0}\epsilon_{0}\frac{\partial E}{\partial t}, \qquad \text{Eq(8)}$$

The dependent variables are adopted as: $n_j = n_{j0} + n_{j1}e^{i(ky-\omega t)}$, Eq(9) $\mathbf{E} = \mathbf{E}_1 e^{i(ky-\omega t)} \hat{x}$, Eq(10) $\mathbf{B} = (\mathbf{B}_0 + \mathbf{B}_1 e^{i(ky-\omega t)}) \hat{z}$, Eq(11) $\mathbf{J} = \mathbf{J}_1 e^{i(ky-\omega t)} \hat{x}$, Eq(12)

Eq(3)

Eq(4)

Eq(5)

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k

Bo

→ E1

Theoretical Model

The Dispersion Relation:

$$\sum_{j} \frac{\Omega_{pj}^2}{1 - \frac{\omega_{cj}^2}{\omega^2}} - c^2 k^2 \sum_{j} \frac{\frac{\Omega_{pj}^2}{\omega^2}}{1 - \frac{\omega_{cj}^2}{\omega^2}} - \left(\sum_{j} \frac{\Omega_{pj}^2}{1 - \frac{\omega_{cj}^2}{\omega^2}}\right) \left(\sum_{j} \frac{\frac{\Omega_{pj}^2}{\omega^2}}{1 - \frac{\omega_{cj}^2}{\omega^2}}\right) = 0, \qquad \text{Eq(13)}$$

where Ω_{pj} and ω_{cj} are the plasma frequency and the cyclotron frequency of the species *j*.

This is the dispersion relation of a modified Alfven waves in collisionless, cold plasma medium that propagate perpendicularly to the ambient magnetic field.

The derived mathematical model could be applied to a wide range of application in space and astrophysical plasmas.

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Application to Observed Structure



$$j = CO_2^+, O_2^+, O^+, He^{+2}, H_2^+, e^-$$

The Dispersion Relation becomes:

$$c_{14}\omega^{14} + c_{12}\omega^{12} + c_{10}\omega^{10} + c_{8}\omega^{8} + c_{6}\omega^{6} + c_{4}\omega^{4} + c_{2}\omega^{2} + c_{0}\omega^{0} = 0, \quad \text{Eq(14)}$$

Where $c_{14}, c_{12}, ..., c_0$ are the coefficients of ω .

Solving this equation numerically relying on the plasma parameters observed on the Martian upper ionosphere, results in 7 symmetrical roots.

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Results

- A root corresponds to each species.
- Each root is divided into two wave modes.
- Three different wave modes:
 - Alfven mode.
 - Whistler mode.
 - Magnetized Plasma Analog of Langmuir mode.
- For each species' mode, there are:
 - Long wavelength (Magnetized mode).
 - Short wavelength (cyclotron mode).



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• Effects of varying the magnetic field and the total plasma density and the mixing ratios on the Alfven Mode:



• Whistler Mode:



• Effects of varying the magnetic field and the total plasma density and the mixing ratios on the Whistler Mode:



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Magnetized Plasma Analog of Langmuir mode 6×10^{6} 5×10^{6} Magnetized 4×10^{6} **Plasma Analog** $\omega(s^{-1})$ of Langmuir 3×10^{6} mode 2×10^{6} 1×10^{6} 0 40×10^{-6} 60×10^{-6} 20×10^{-6} 80×10^{-6} $k(m^{-1})$

• Effects of varying the magnetic field and the total plasma density and the mixing ratios on the Whistler Mode:



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Conclusions

- The derived mathematical model could be applied to a wide range of application in space and astrophysical plasmas.
- A Modified Alfven wave has been studied in conditions relevant to the upper ionosphere of Mars.
- Every species has a propagating wave mode that saturates at the specie's cyclotron frequency.
- An Alfven mode exists in the ion species with the heaviest mass.
- A Whistler mode exists in all of the less massive species.
- A Magnetized Analog of Langmuir mode exists due to the displacement current.

Conclusions

- Increasing the magnetic field increases the resonance frequency of the Alfven mode.
- Increasing the magnetic field increases the difference between the cutoff and resonance frequencies of the Whistler mode and causes nothing to the magnetized analog of Langmuir mode.
- Varying the total number density affects the value of the critical wavenumber at the Alfven and Whistler modes.
- Varying the total number density causes a shift in the magnetized analog of Langmuir mode.

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Future Work

- The observation expects a Magnetosonic wave to be steepened.
- We suggest further parameters to be considered:
 - Inhomogeneity of the medium in the magnetic field, density or temperature.
 - Oblique propagation of Magnetosonic waves. (on going)
 - Streaming conditions relevant to the effects of the solar wind.

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