

# Muonic Atom Spectroscopy of $^{238}\text{U}$ for the Extraction of Nuclear Properties

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on behalf of the muX collaboration

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# muX Collaboration

## Goal

Determine the absolute nuclear charge radius of  $^{226}\text{Ra}$  for the first time

## Why $^{226}\text{Ra}$ ?

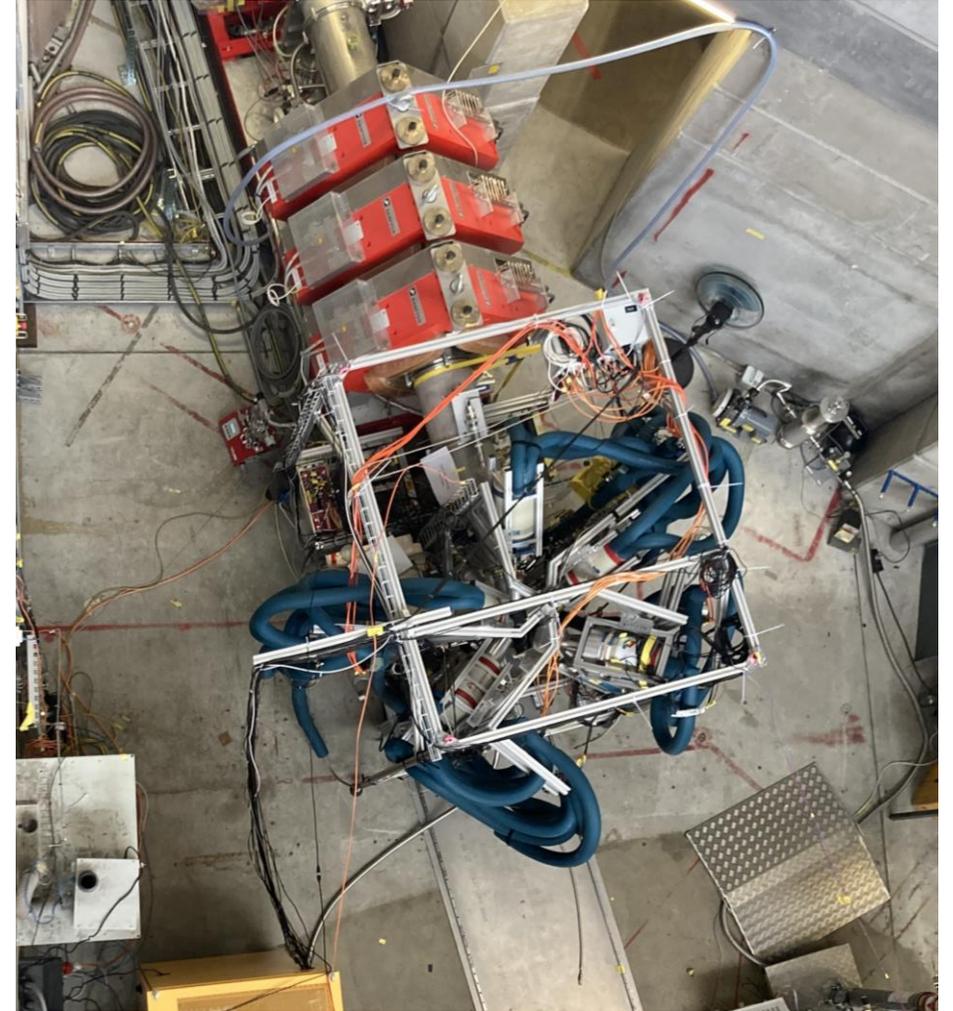
Excellent candidate for measurements of atomic parity violation → physics beyond the standard model

## BUT

Open radioactive sources available in microgram,  $\mu\text{g}$ , quantities

## WAY

Performing muonic atom spectroscopy with a novel muonic capture technique



# Muonic Atoms

Muons ( $\mu^-$ ) are leptons, with an electric charge of  $-1$ , spin- $\frac{1}{2}$  and 207 times more massive than electrons:

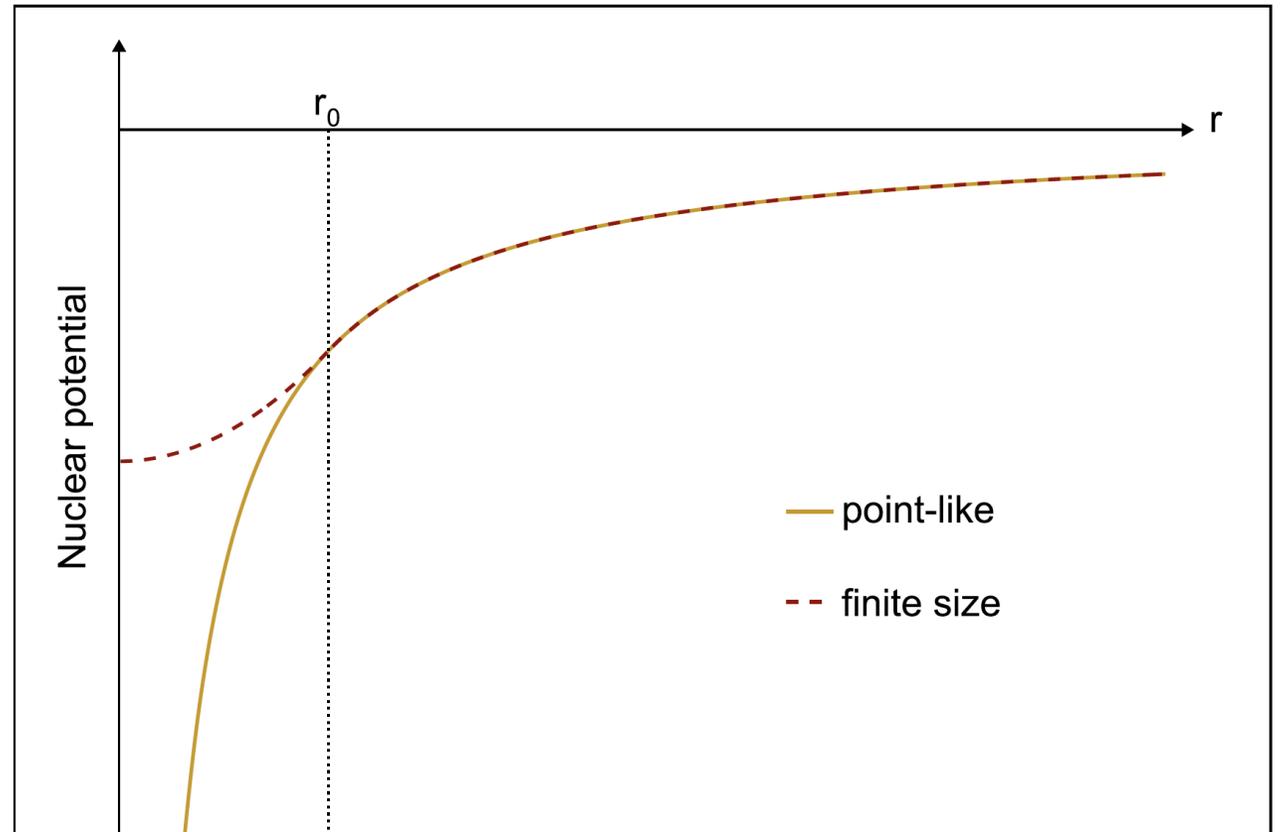
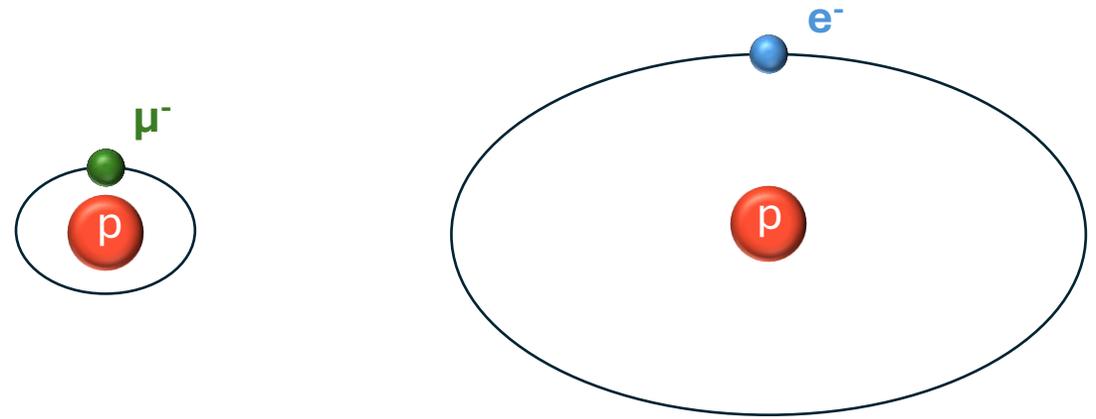
$$m_\mu \approx 207 m_e$$

Muonic atoms are formed when a muon gets captured by a nucleus

Muon orbits 207 times closer to the nucleus than electrons :

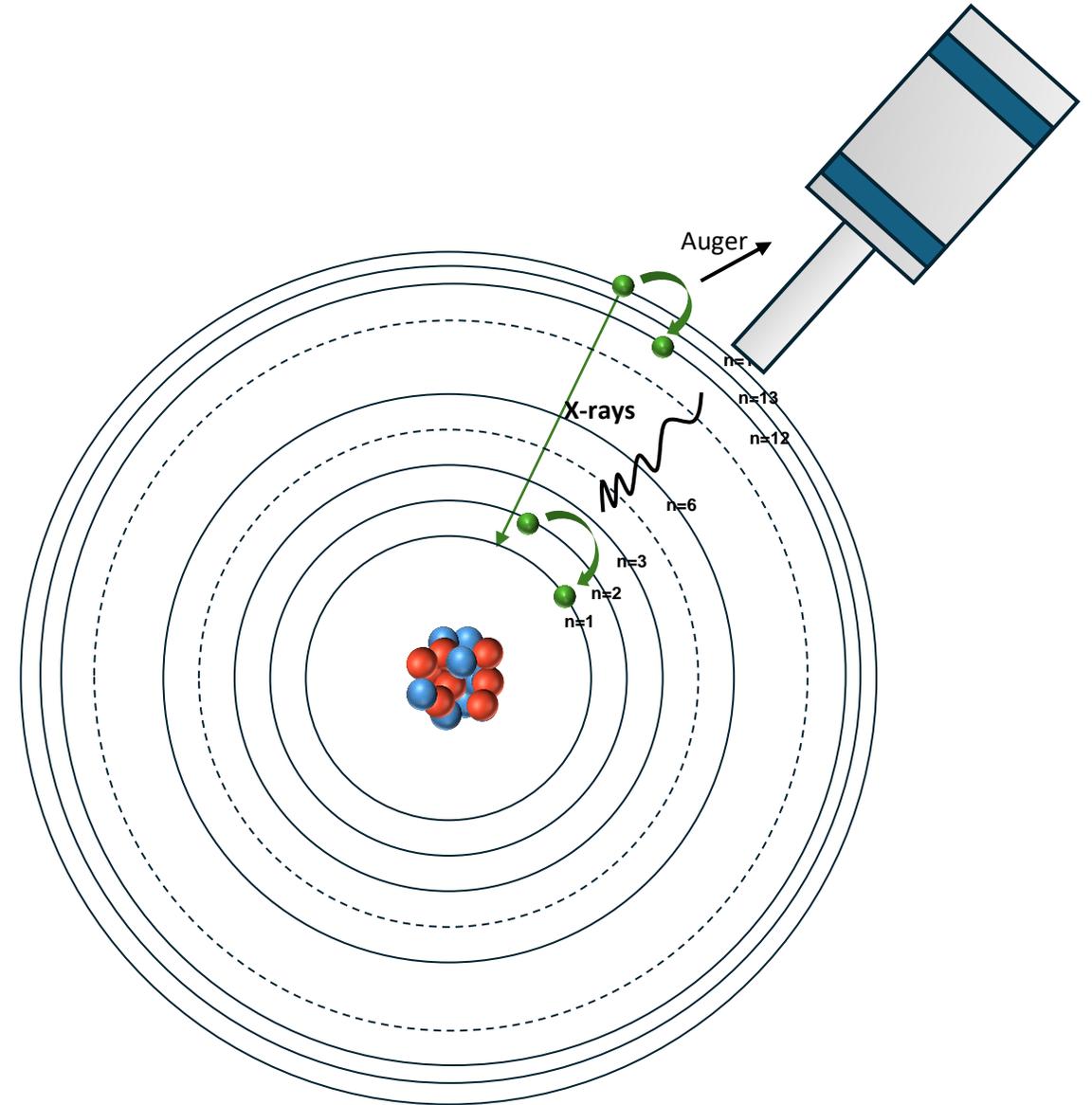
$$\text{Bohr radius: } r_\mu \approx \frac{1}{207} r_e$$

Muonic wave function overlaps with the nuclear charge distribution  $\rightarrow$  muonic energy levels show high sensitivity to the nuclear properties

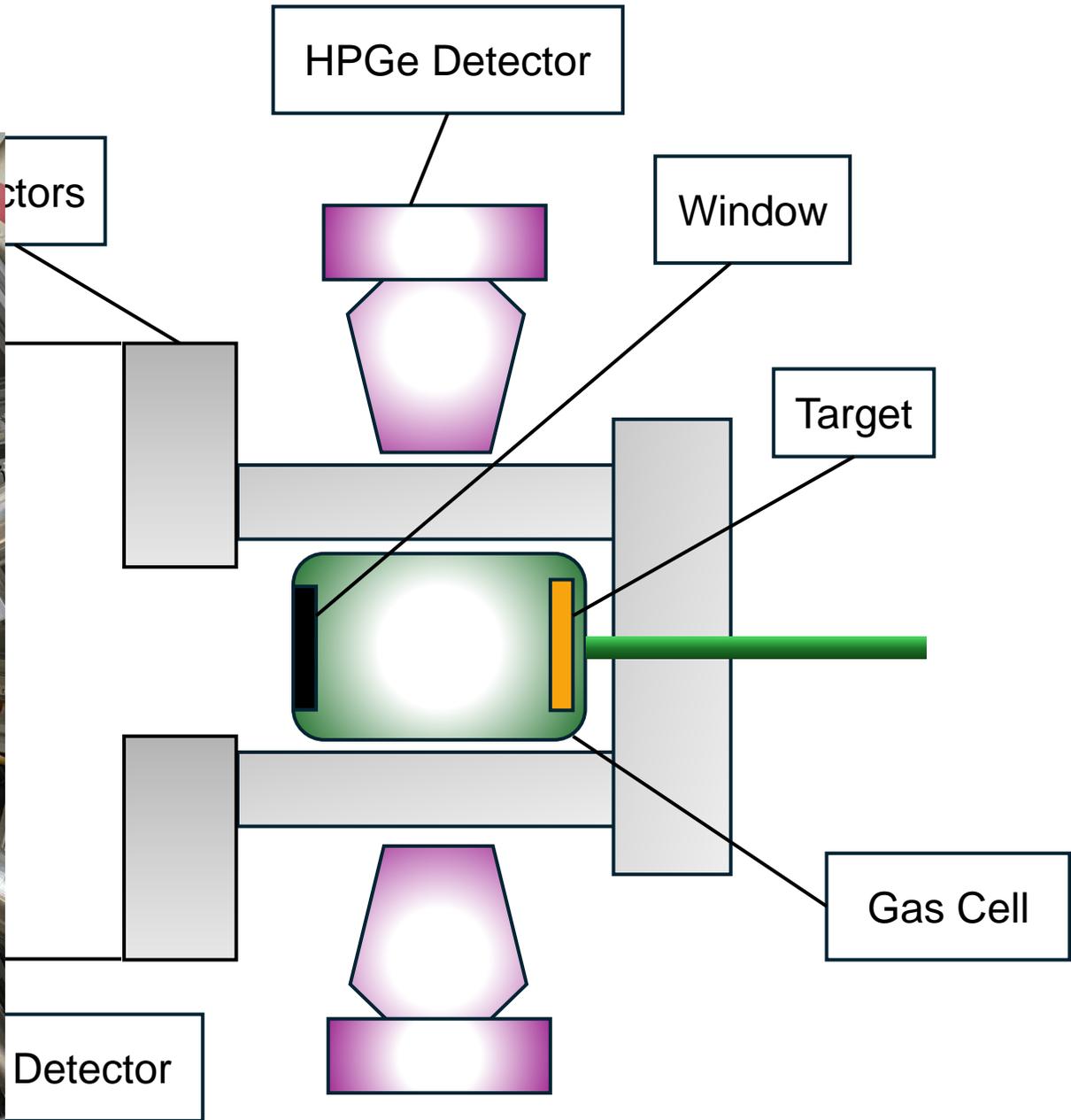


# Muonic Atom Spectroscopy

1. Muon gets captured at  $n \sim 14$
2. Cascades down to the ground,  $1s$ , state
3. Beginning of the cascade  $\rightarrow$  strong Auger effect
4. Below  $n \sim 6 \rightarrow$  radiative transitions dominant, emission of muonic X-rays
- ✓ Muonic X-rays (between keV to MeV) detected by High-Purity Germanium detectors (HPGe)
5. Reaching the ground state muon gets captured by the nucleus or decays

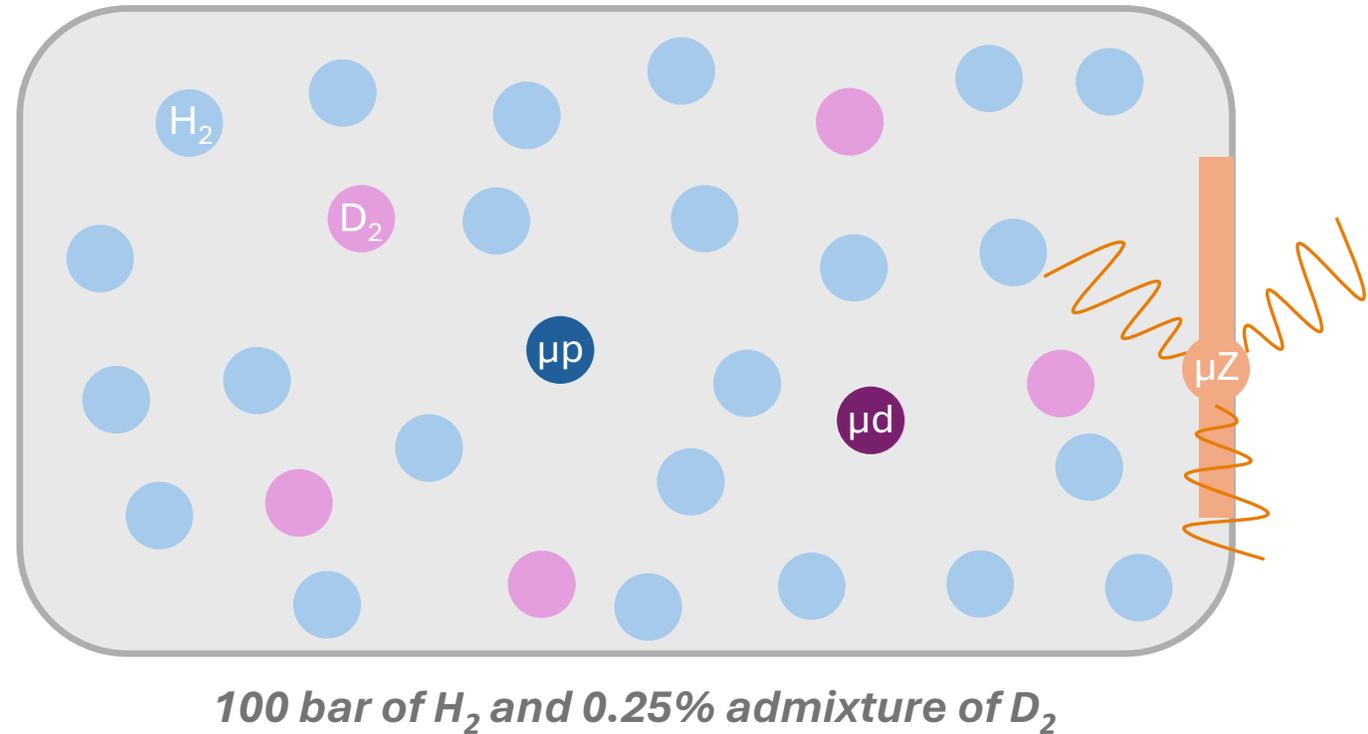
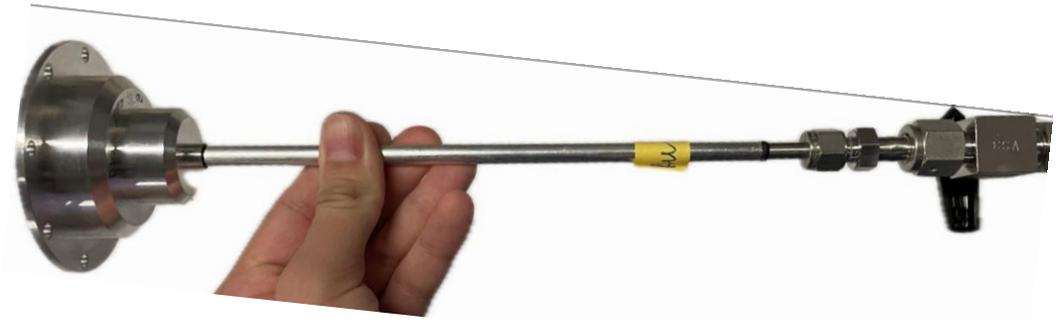


# The muX Experimental Setup



# Transfer of Muons in a Gas Cell

1. The muon ( $\mu^-$ ) enters the gas cell
2. Muon interacts with hydrogen ( $H_2$ ) and forms muonic hydrogen ( $\mu p$ )
3. Muonic hydrogen ( $\mu p$ ) interacts with deuterium ( $D_2$ ) and forms muonic deuterium ( $\mu d$ ) μ
4. Muonic deuterium ( $\mu d$ ) gains kinetic energy of 45 eV (boosted)
5. Muonic deuterium scatters and gradually loses its kinetic energy
6. At 4 eV  $\mu d$  “blind” to  $H_2$  (Ramsauer-Townsend effect)
7. Muon reaches and gets captured by the heavy-Z target material



# Muonic atom spectroscopy of $^{238}\text{U}$

Performed the muonic atom spectroscopy of  $^{238}\text{U}$  and examined the cascade behavior for direct and transfer muon capture.

## DATA SETS

**2017 data** transfer muon capture, muons go through gas mixture and to the target

**2023 data** direct muon capture, muons directly shot to the target

## APPROACH

Performed muonic  $^{238}\text{U}$  spectroscopy by:

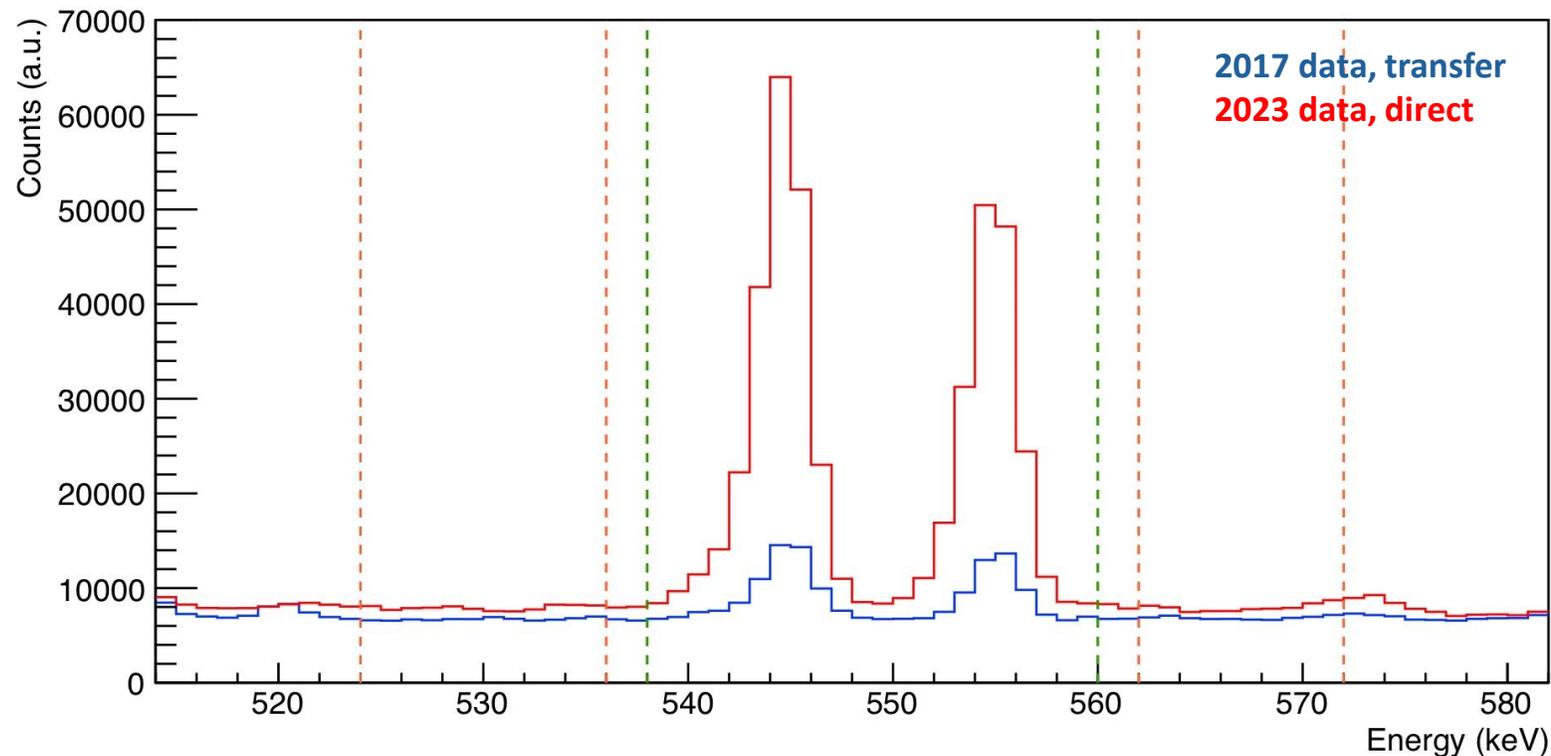
1. Removing the background and calculating the intensity of the possible transition  $2p \rightarrow 1s$ ,  $3d \rightarrow 2p$ ,  $4f \rightarrow 3d$  and  $5g \rightarrow 4f$
2. Scaling the 2017 data with respect to the 2023 data
3. Compare the cascade behavior between the two years and with the cascade simulations for late and prompt time windows



# Transition Intensity and Background Removal

1. **Transition Intensity without Background:** Going at each transition/peak of the energy spectrum, calculate the integral underneath and remove the background based on the integrals of the neighbouring regions:

$$Intensity = \sum_{j=1}^N C_j - \frac{N}{2} \frac{1}{M} \sum_{i=1}^M C_i - \frac{N}{2} \frac{1}{L} \sum_{l=i}^L C_l$$



# Data Scaling

In order to compare the energy spectrum of the two different data sets:

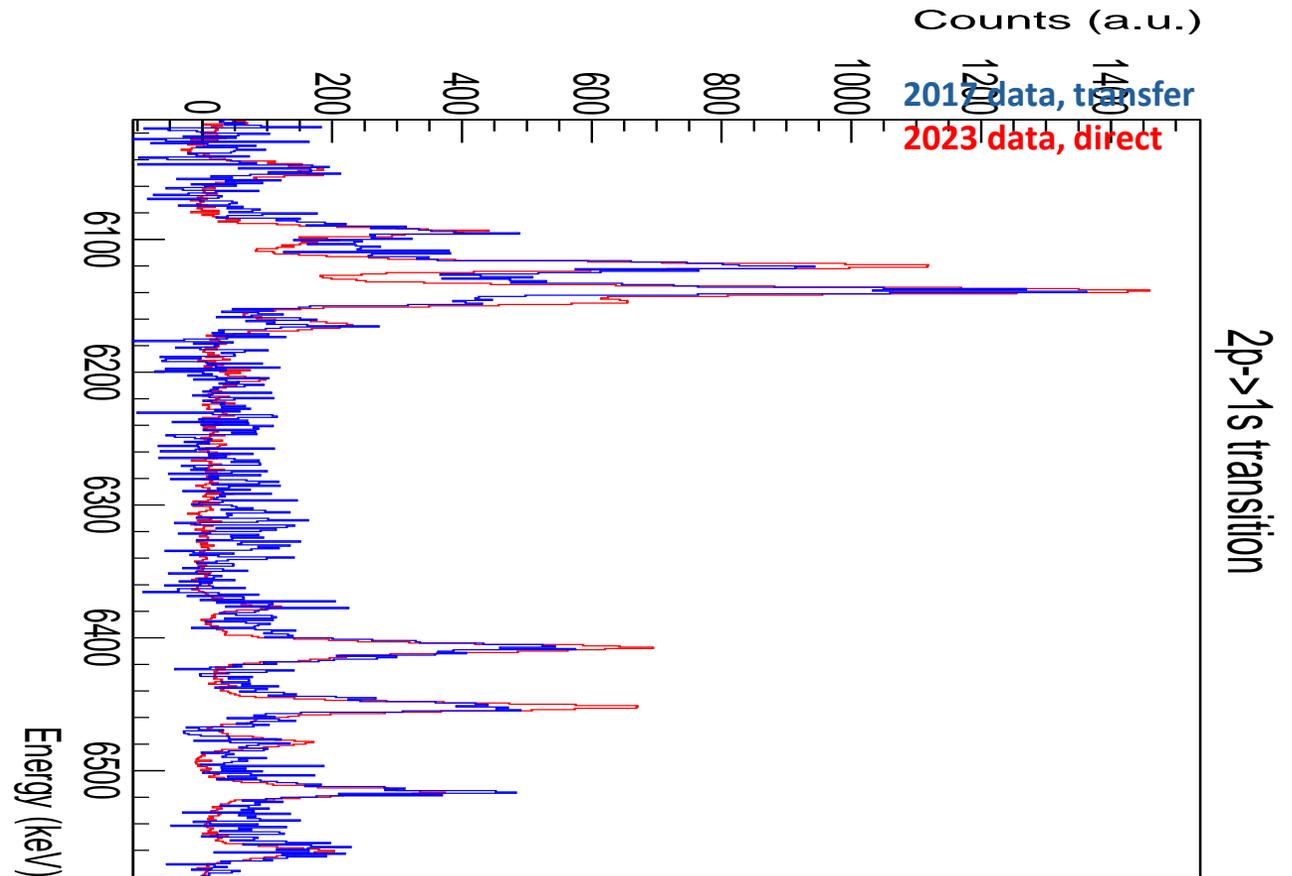
2. **Scaling data:** Scaled 2017, transfer data relative to the 2023, direct data with respect to the most prominent peaks of the  $^{238}\text{U}$  spectrum at  $2p \rightarrow 1s$  transition:

$$\text{Scaling Factor} = \frac{I_{2p \rightarrow 1s} (2023)}{I_{2p \rightarrow 1s} (2017)}$$

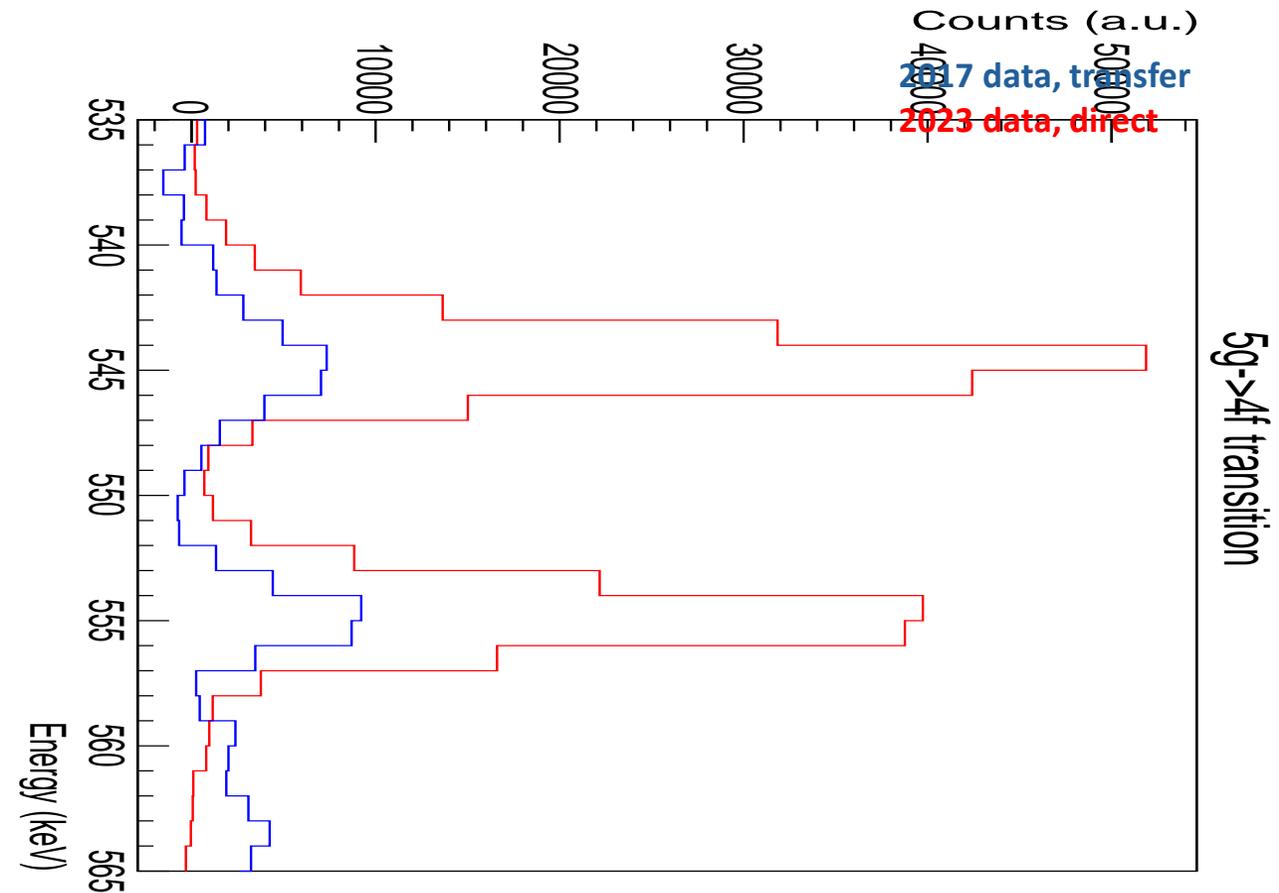
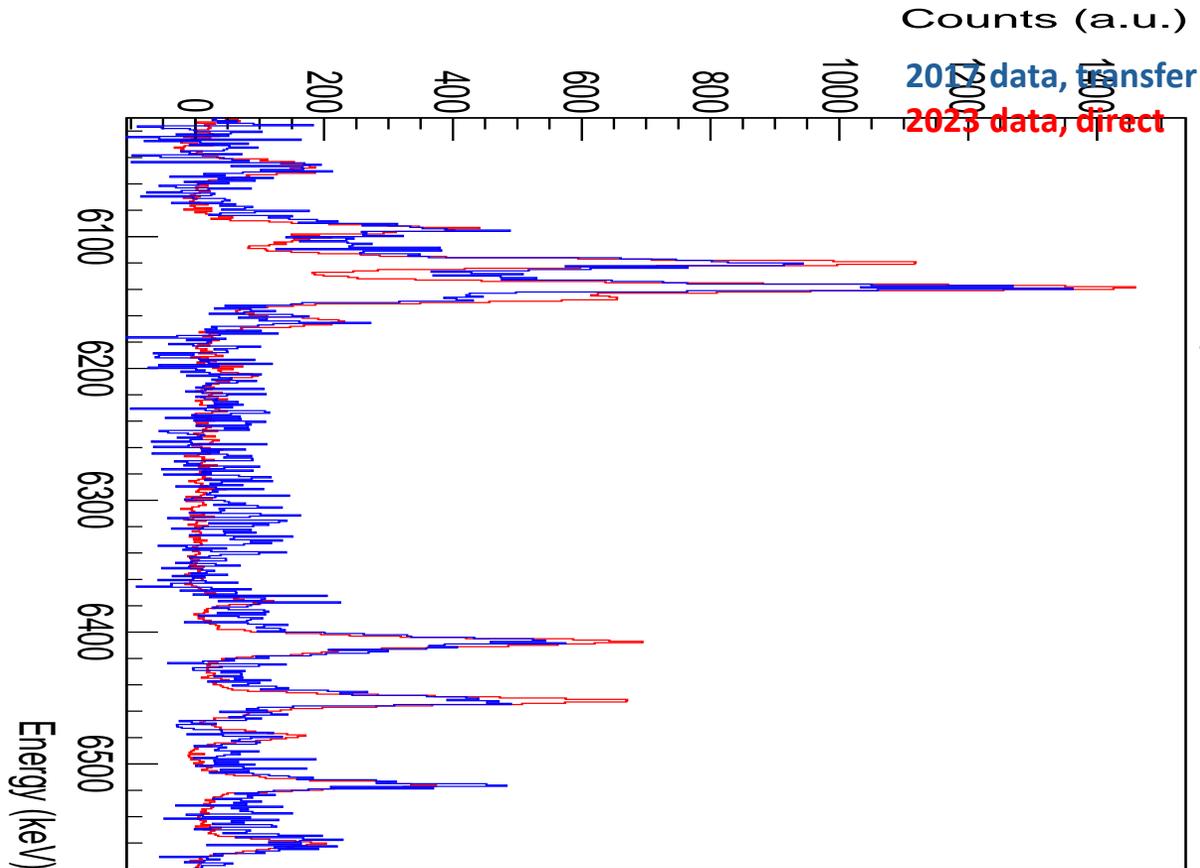
2023 data, for a prompt time interval (-50 to 50 ns)

2017 data, for a late time interval (+50 to 260 ns)

✓ After scaling the two data sets overlap



# Transition $5g \rightarrow 4f$ , Direct versus Transfer Muon Capture



2023, direct data  $\sim$  5 times more intense peaks than 2017, transfer data

✓ Matching predictions of the cascade simulations

# Conclusions

- Muonic atom spectroscopy is used to study nuclear properties such as charge radius and quadrupole moment
- muX collaboration developed a method based on transfer reactions in a gas cell to perform muonic atom spectroscopy with materials that are available in microscopic quantities, like radioactive
- Muonic atom spectroscopy of a heavy-Z target material,  $^{238}\text{U}$ , showed that circular transitions are more favored for direct than transfer muon capture
- Goal of the muX Collaboration is to determine the  $^{226}\text{Ra}$  absolute charge radius  $\rightarrow$  candidate for APV and thus physics BSM

Thank you for your  
attention!

# BACKUP SLIDES

# Muon Capture Timing

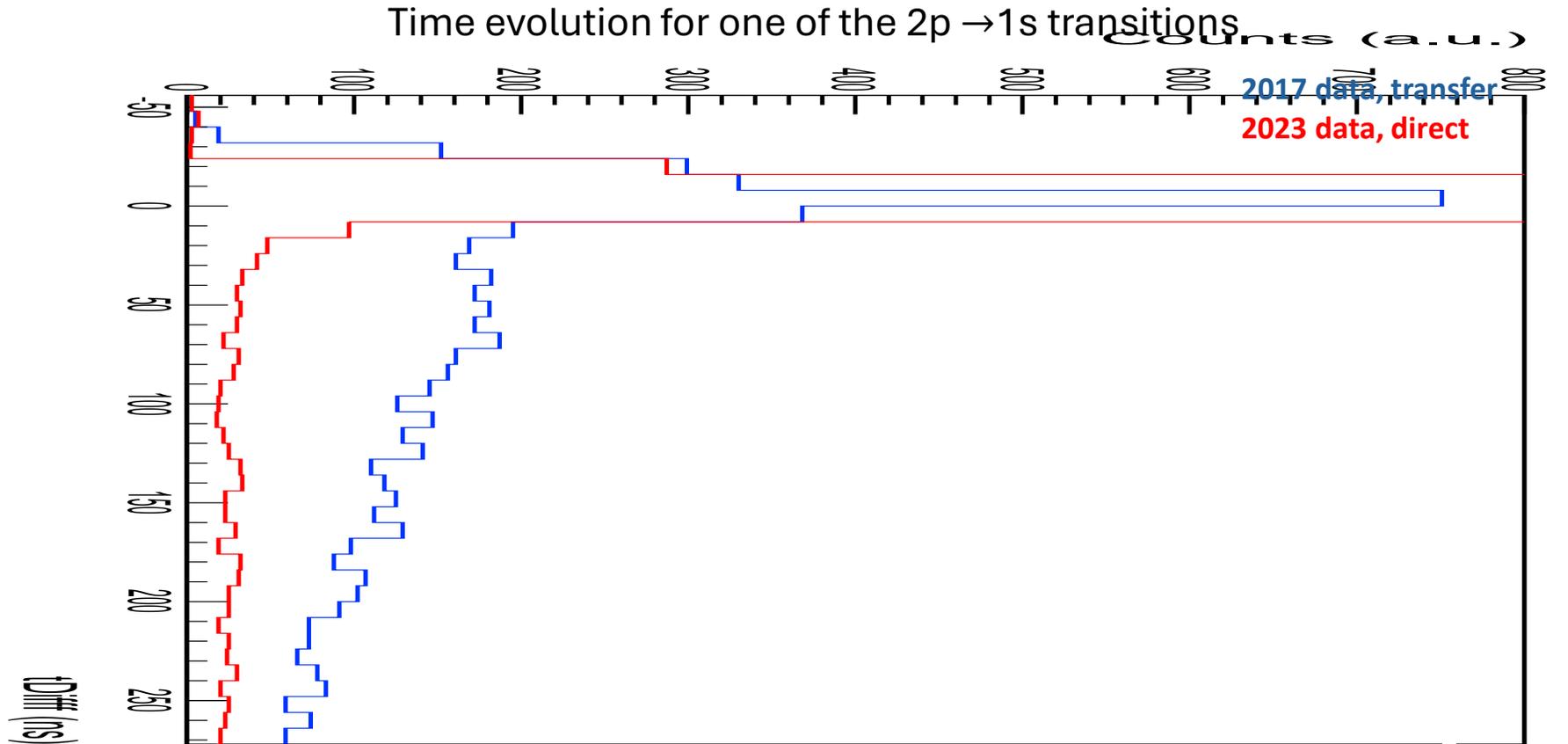
## DATA SETS

**2017 data** transfer muon capture, muons go through gas mixture

- Prompt and late muon capture
- Peak and long tail

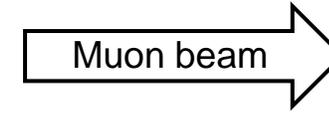
**2023 data** direct muon capture, muons directly shot to the target

- Prompt muon capture
- Sharp peak

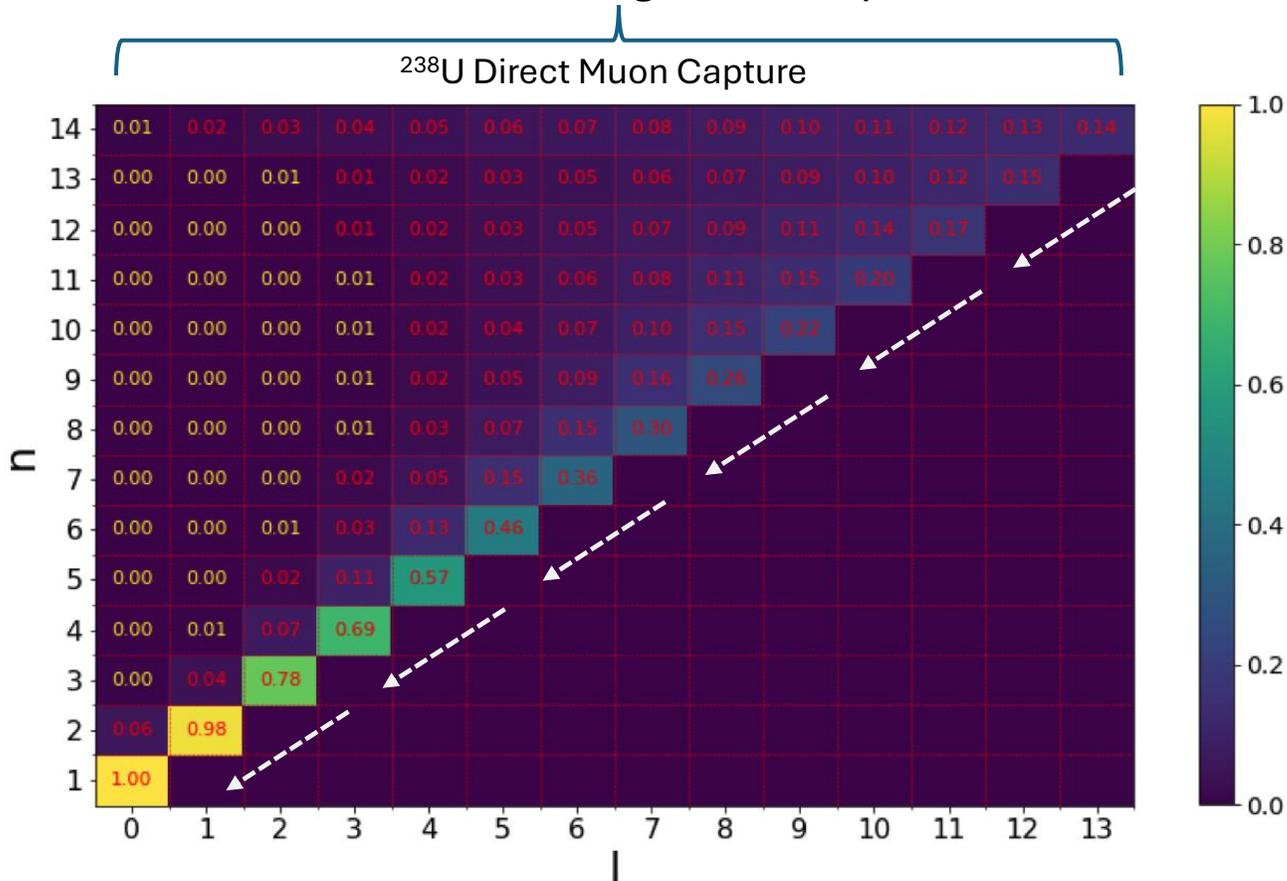


# Direct Muon Capture

Target available in macroscopic quantities, more than 100 mg. Muons shot and directly stopped at the target.



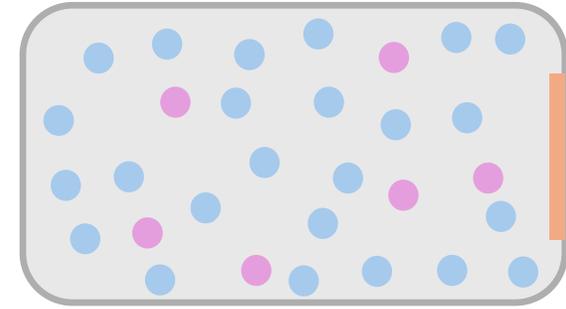
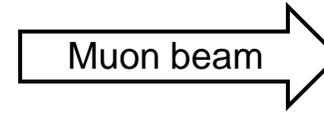
Cascade simulation calculating the rates/probabilities for the Auger and X-ray transitions



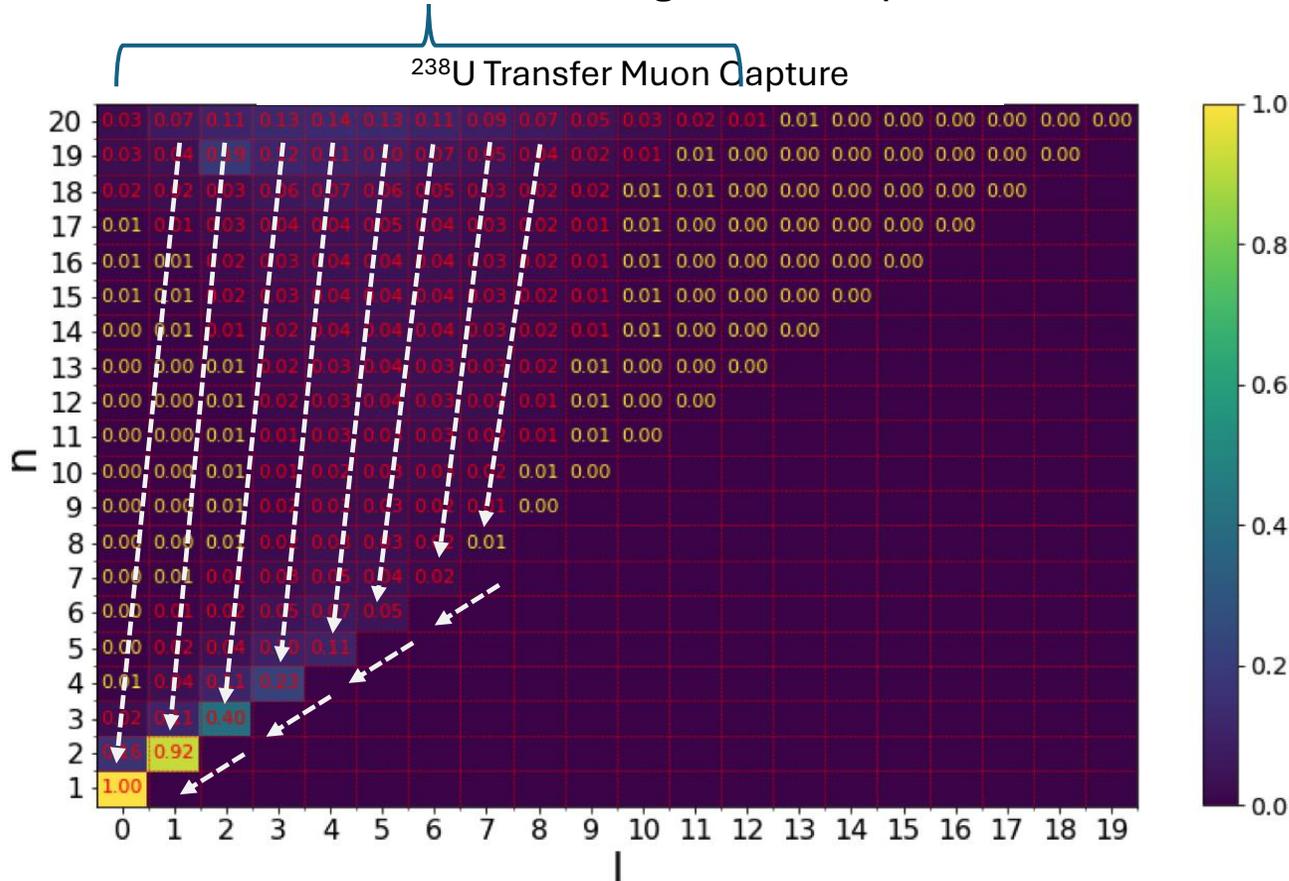
- Muons captured at  $n \sim 14$
- Angular momenta of captured muons distributed based on  $(2l+1)$  statistical distribution
- $(n, l=n-1) \rightarrow (n-1, l=n-2)$  transitions favored (circular transitions)

# Transfer Muon Capture

Target available in microscopic quantities, in the order of  $\mu\text{g}$ . Muons transverse the gas cell, undergo a three-stage transfer process and get captured by target material.

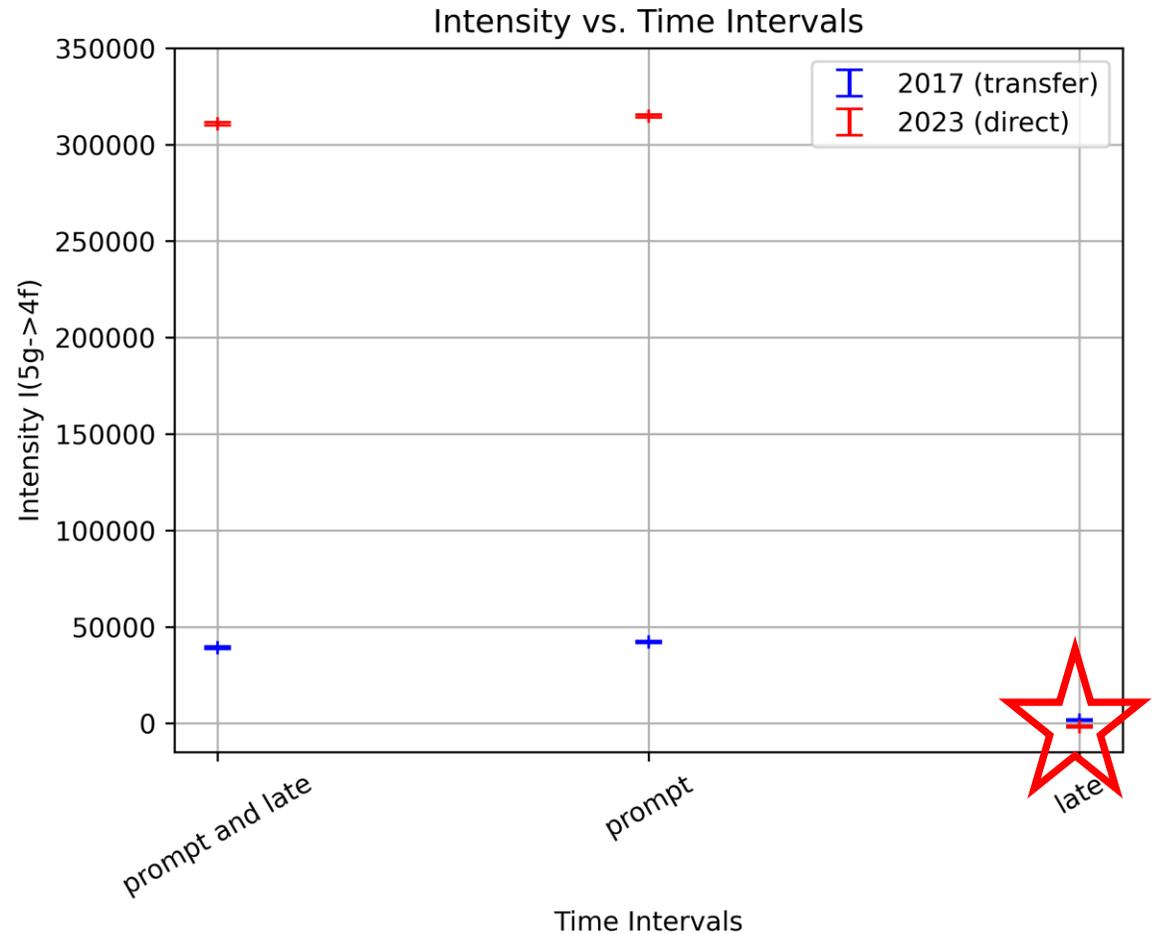
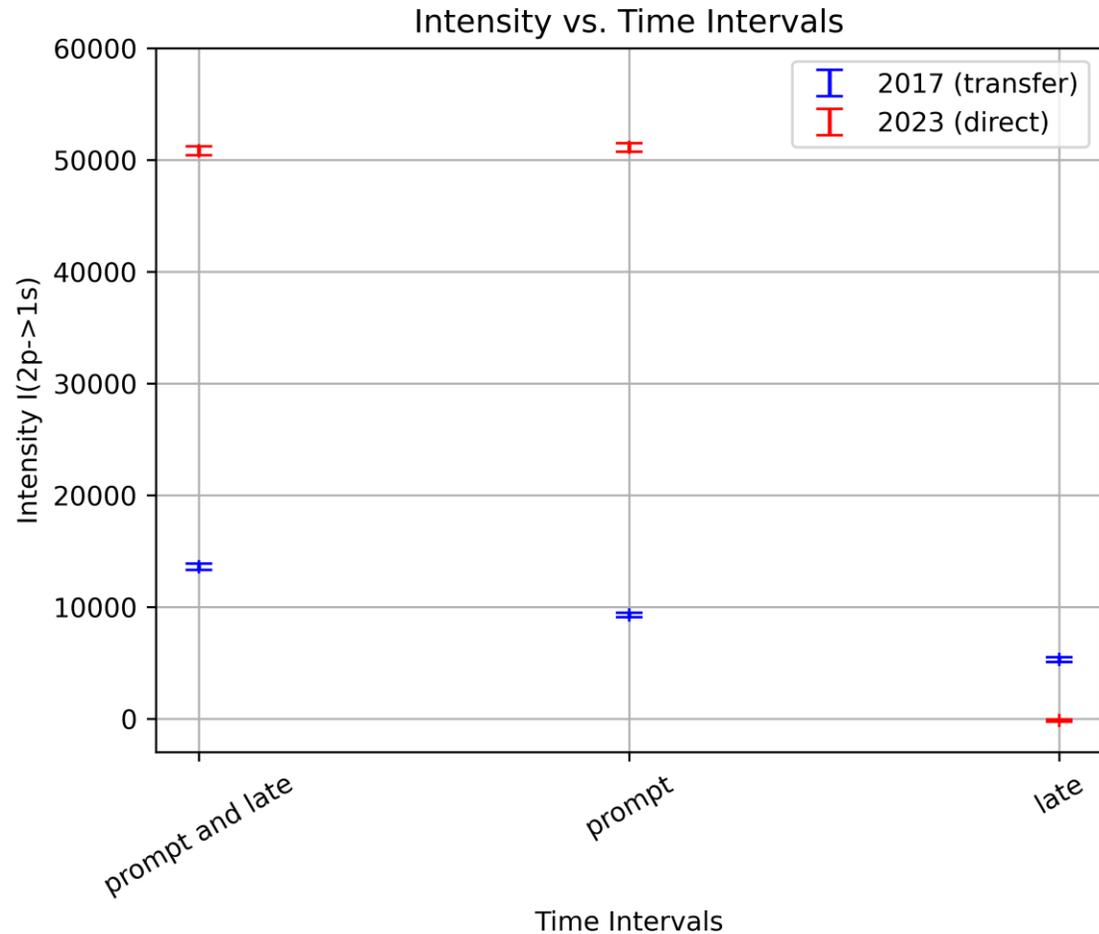


Cascade simulation calculating the rates/probabilities for the Auger and X-ray transitions



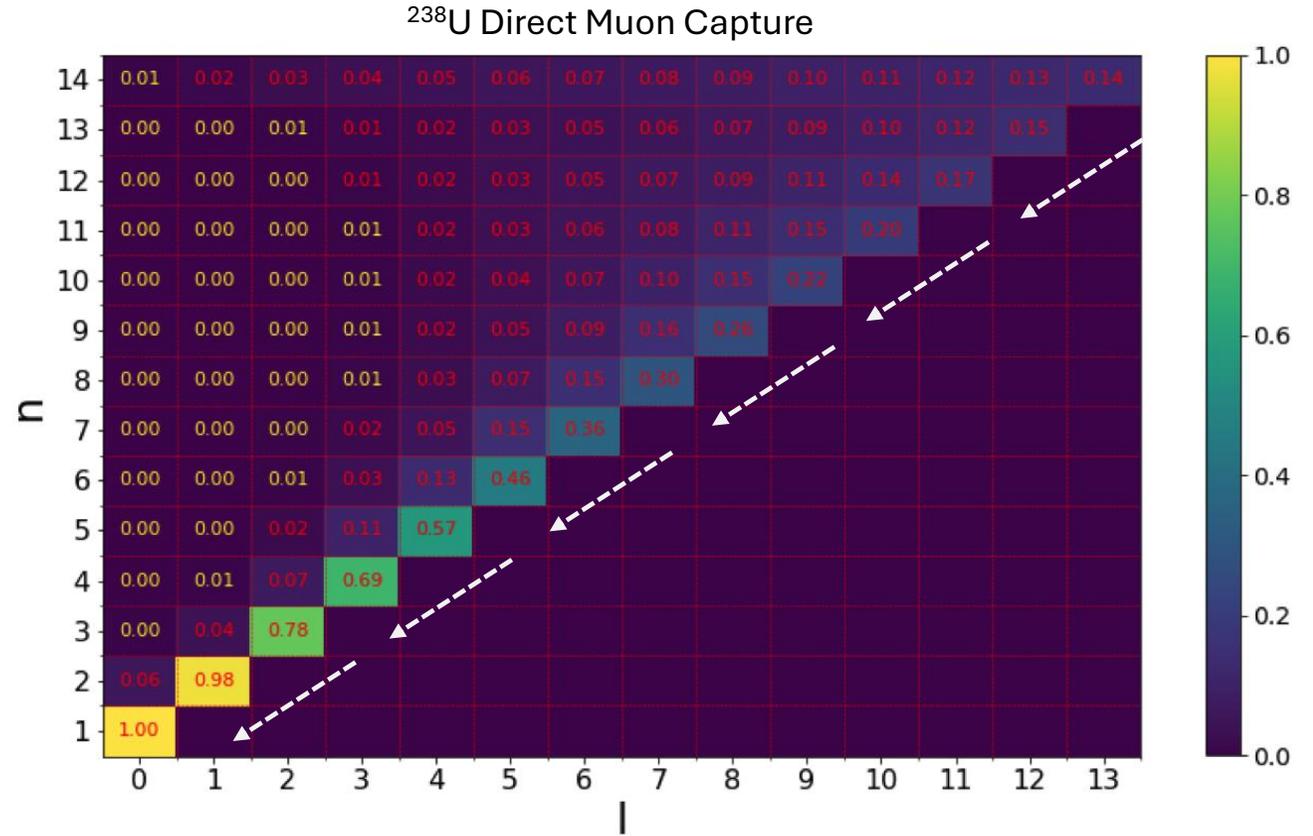
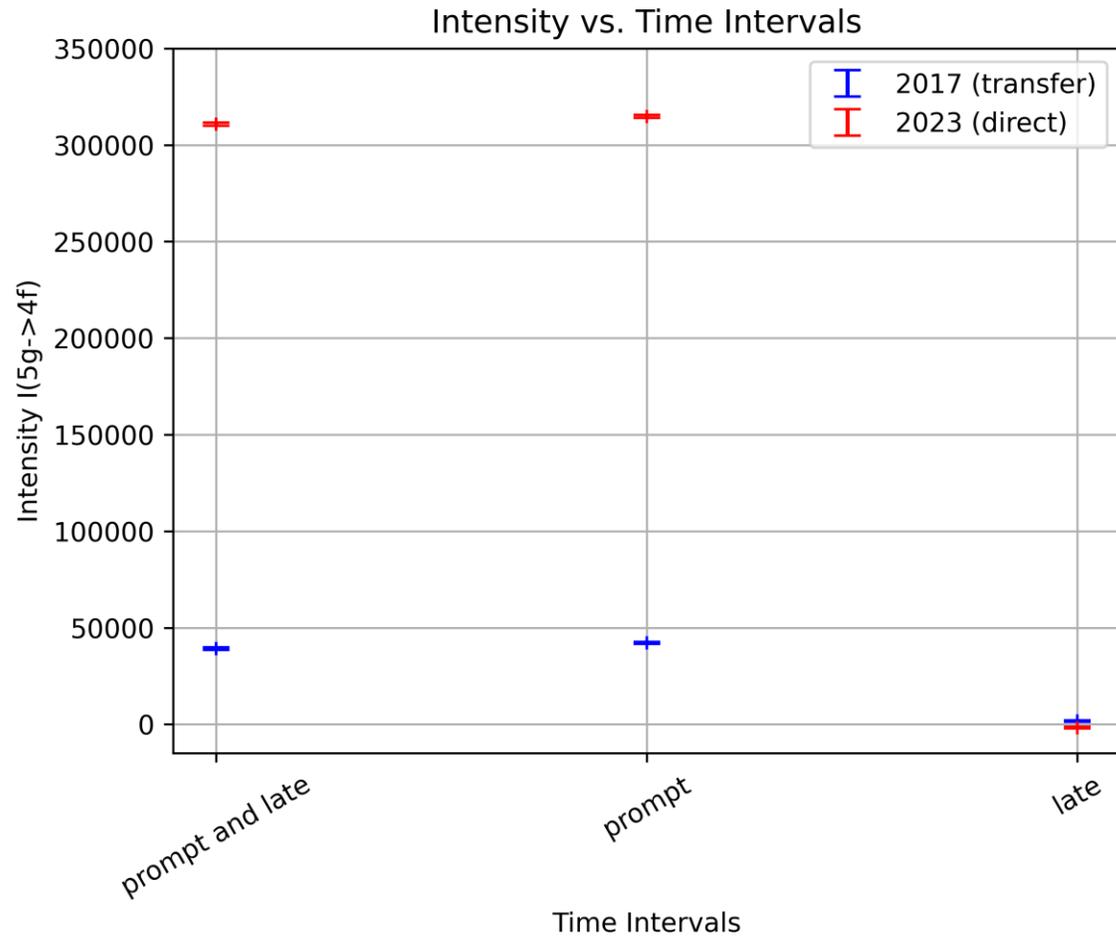
- Muons captured at  $n(Z)$ ,  $n(Z) \gg 14$  for heavy-Z atoms
- Captured muons concentrated at low angular momenta
- Starting concentrated at low angular momenta cascade down to the 1s state

# Intensity of Transitions $2p \rightarrow 1s$ and $5g \rightarrow 4f$ versus Time



- 2017 data, transfer muon capture at prompt and late time intervals
- 2023 data, direct muon capture at prompt time interval

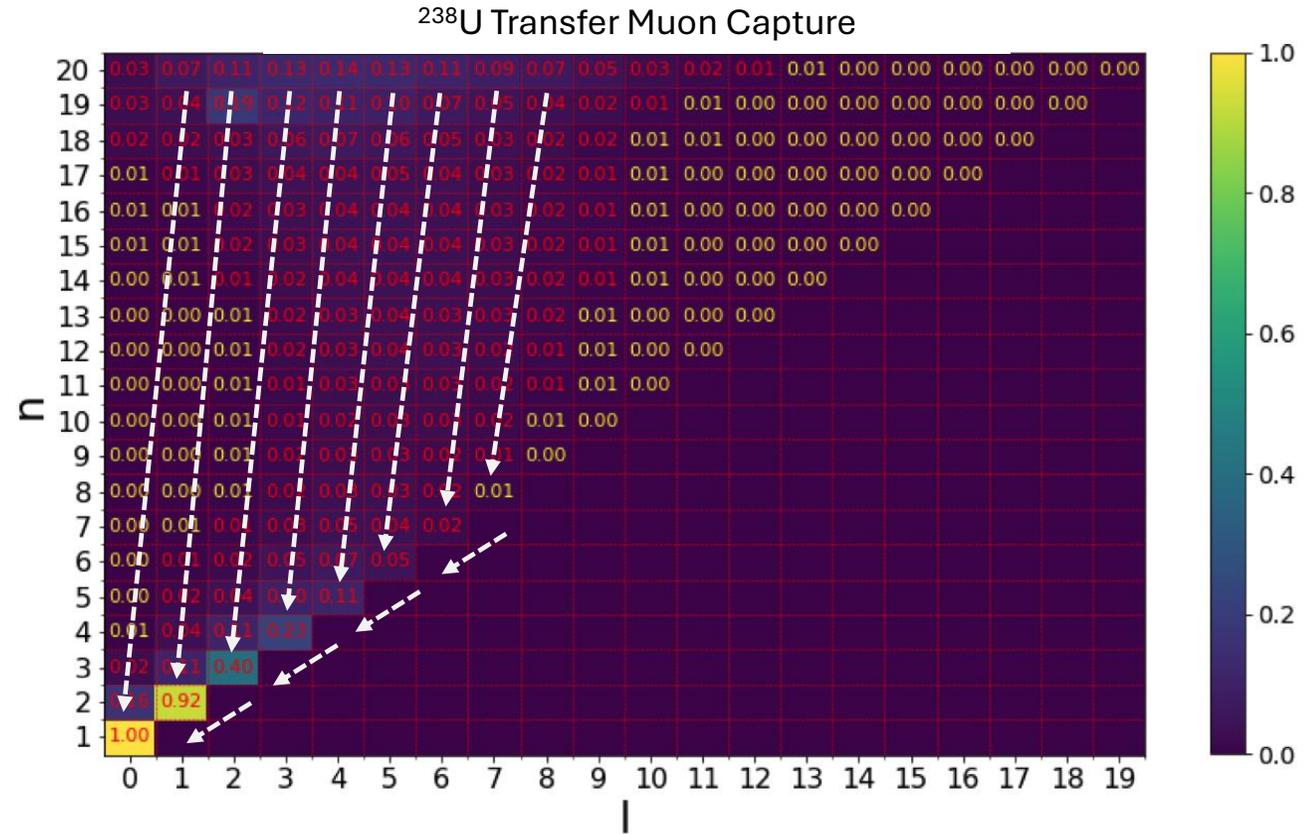
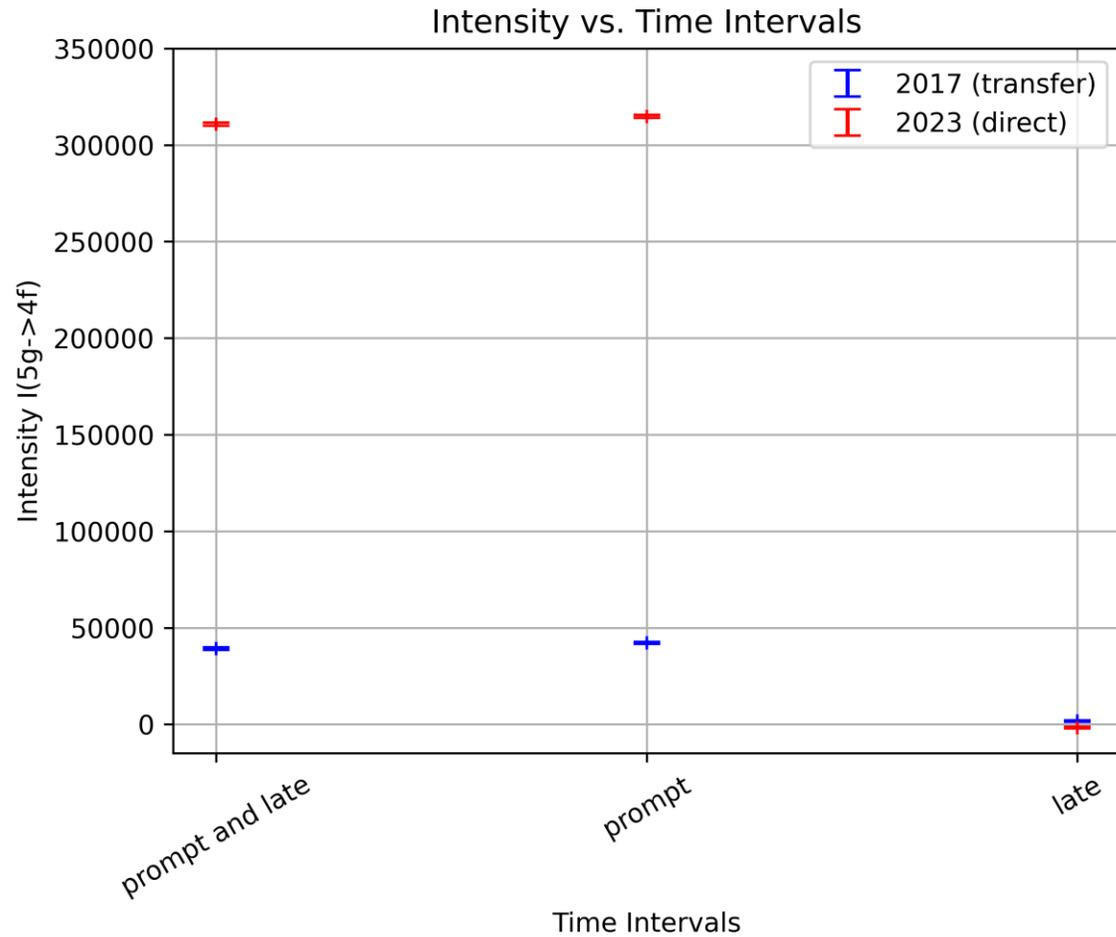
# Transition $5g \rightarrow 4f$ , Direct Muon Capture



Direct muon capture: favors  $(n, l=n-1) \rightarrow (n-1, l=n-2)$  circular transitions

➤ Circular transitions and thus  $5g \rightarrow 4f$  transition favored for prompt, direct capture

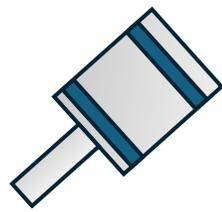
# Transition $5g \rightarrow 4f$ , Transfer Muon Capture



Transfer muon capture: starting concentrated at low angular momenta cascades down to the 1s state

➤ Circular transitions and thus  $5g \rightarrow 4f$  transition not favored for late, transfer capture

# Germanium Detectors



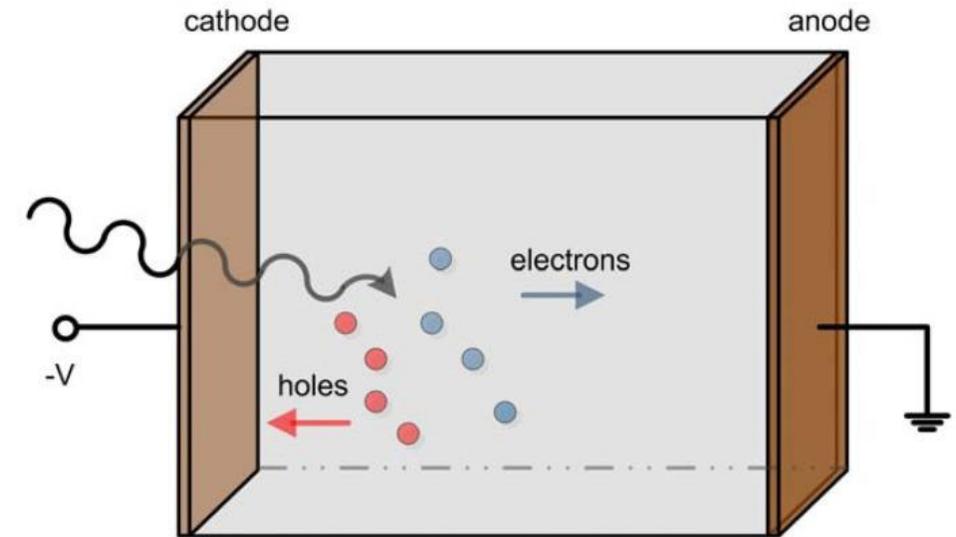
Germanium crystal under a reverse biased voltage, collecting the charge carriers produced by the ionizing radiation.

## Detection Process

- ✓ Apply reverse bias voltage on the crystal → large depletion region → little to no current flowing through the circuit
- ✓ Charged particle enters the depletion region and produces electron-hole pairs
- ✓ Collect charge carriers with the applied voltage difference
- ✓ Number of charge carriers proportional to the ionization radiation energy

Depletion thickness:  $d = \left(\frac{2\varepsilon V}{eN}\right)^{1/2}$ , where  $V$  is the reverse bias voltage and  $N$  is the impurity concentration.

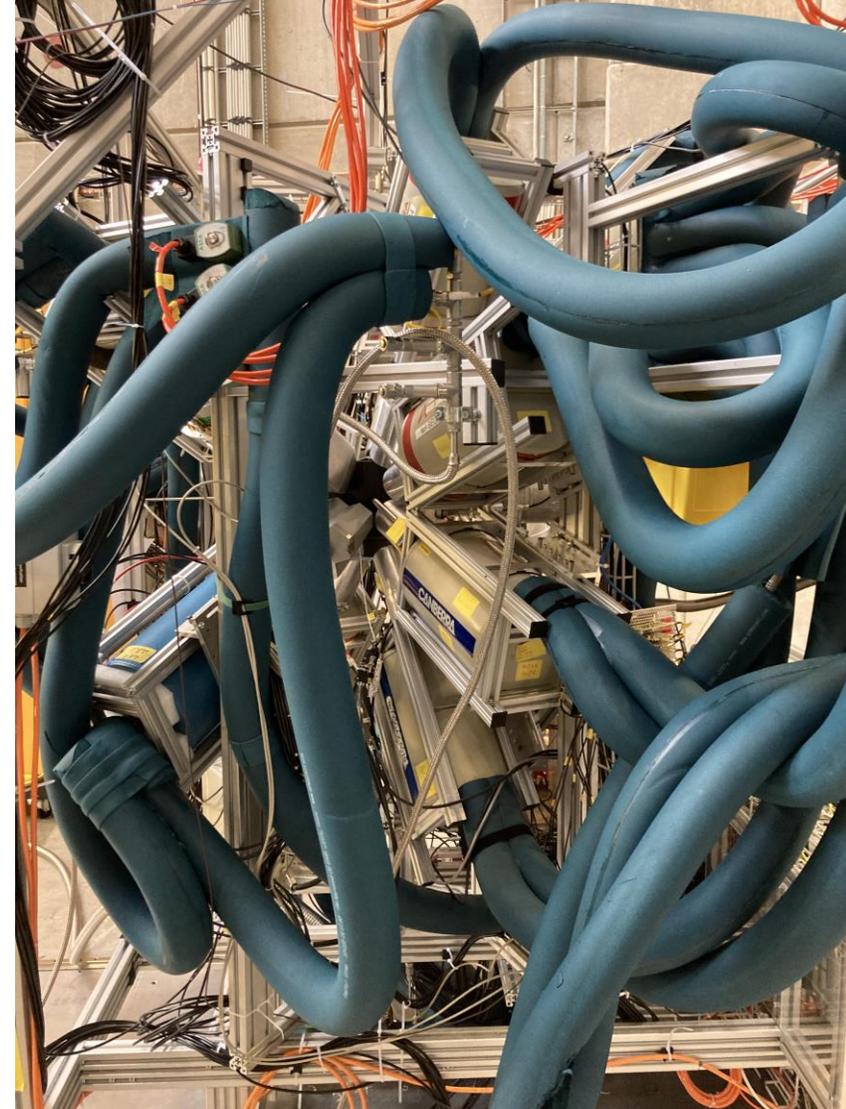
Thus for greater depletion depths, and better energy resolution the impurity concentration needs to be decreased (High purity Germanium Detectors (HPGe)).



# Germanium Detectors, Cooling

As the Germanium semiconductor has a small band gap, at room temperature electrons can move from the valence to the conduction band due to thermal excitations. This creates an unwanted current, called leakage current.

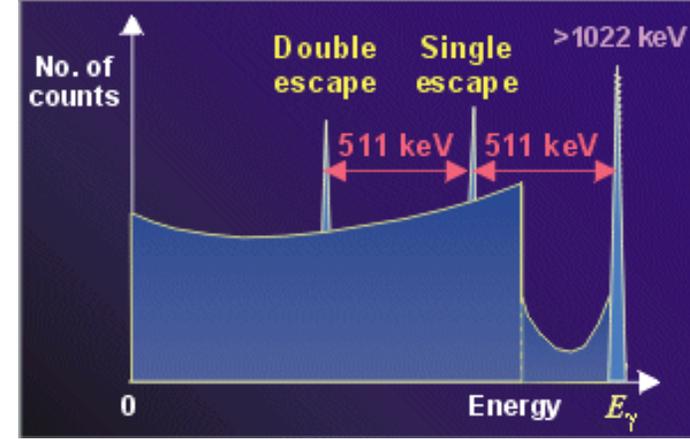
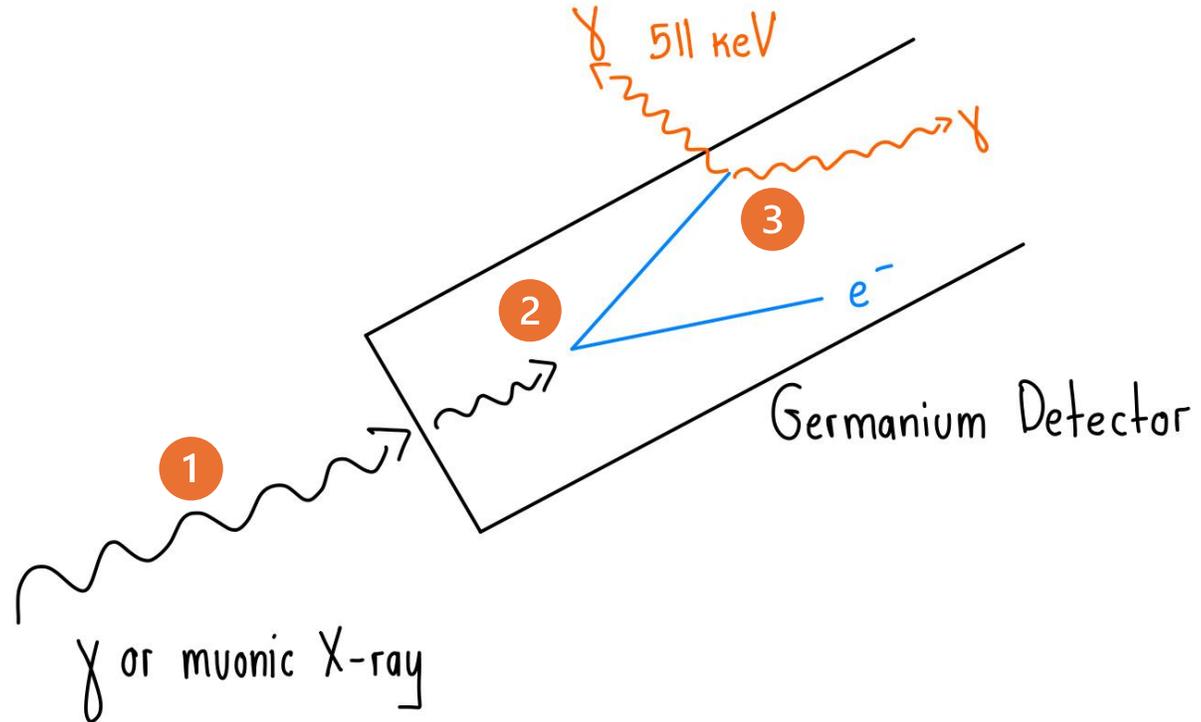
- So we cool down Germanium detectors to avoid any thermal noise, with liquid nitrogen (77K)



# Single & Double Escape

## Single & Double Escape

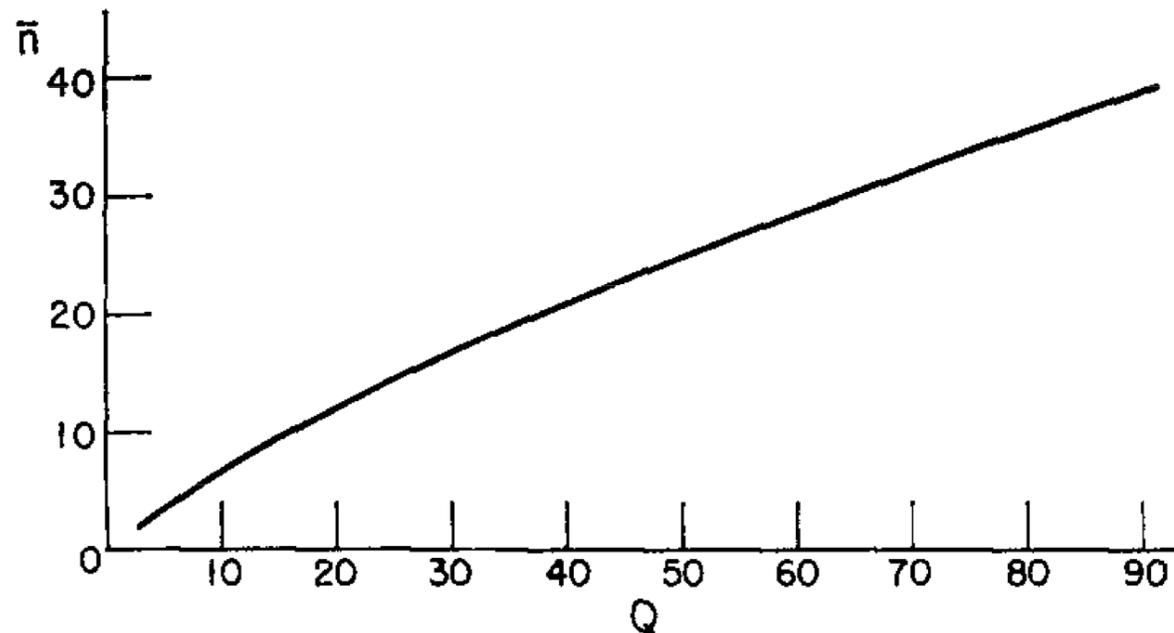
1. Gamma or muonic X-ray enters the Germanium Detector.
2. Interacts with the detector and produces an  $e^-e^+$  pair.
3. The positron interacts with one of the plenty Germanium detector electrons and produces two  $\gamma$ -rays. When one of them escapes we have *Single Escape*. When both of them escape we have *Double Escape*.



# Principle Quantum Number for Transfer Muon Capture

For a heavy-Z material ( $Z=Q$ ) the states  $n$  populated by muons after transfer capture are given by:

$$n \sim \left[ \frac{Q(1 + 2Q^{\frac{1}{2}})}{1 + 2Q^{\frac{1}{2}}} \right]^{1/2}$$



# Flushing and Filling the gas cell

## Interactions in the gas cell

- $\mu d$  interacts with hydrogen (small cross section but non-zero probability)
- $(p\mu d)^+$  decays via two muon catalyzed fusion channels

## Filling circle of the gas cell

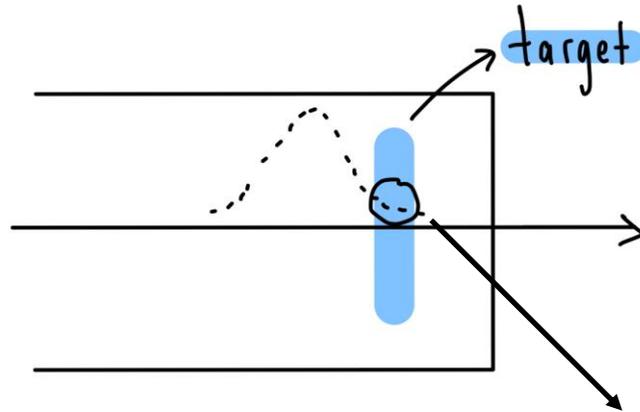
1. Pump the gas mixture from the gas cell
2. Flush with  $H_2$  (10 cycles) up to 10 bars, ensuring no remaining gases
3. Flush with 100 bars of gas mixture
4. Fill with 100 bars of the gas mixture



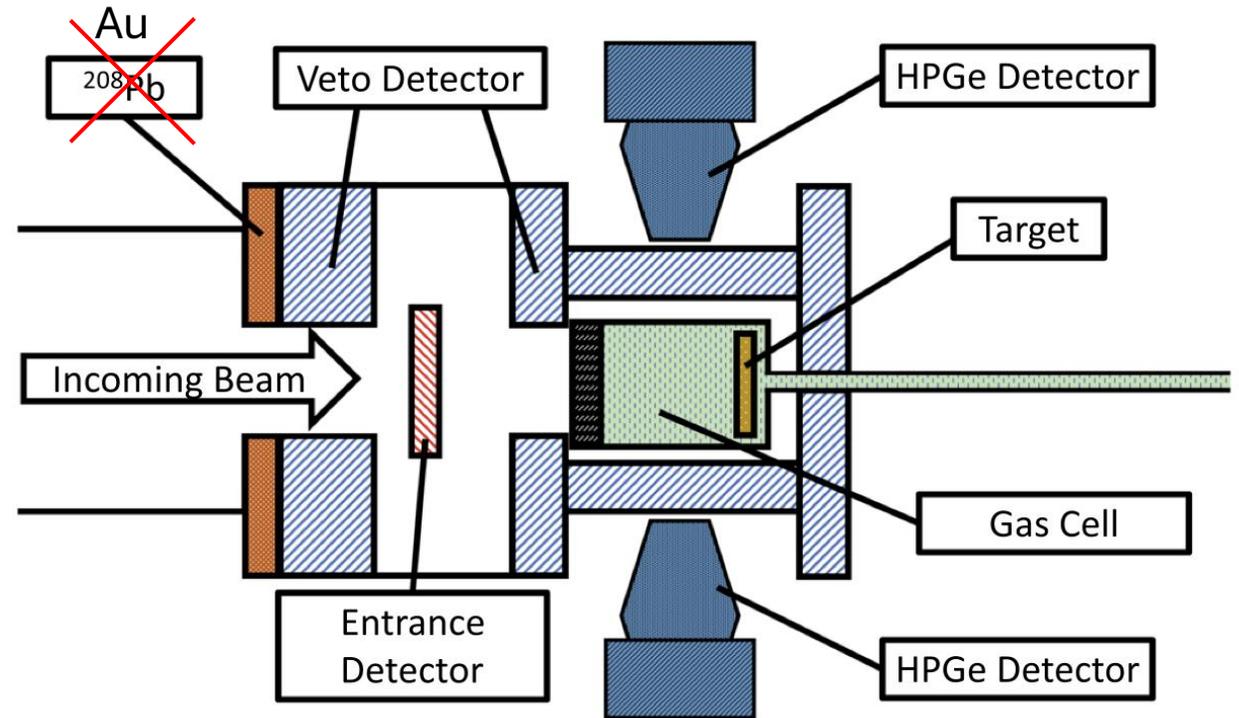
# Thick and Thin $^{137}\text{Au}$ Target

Placed a **thick gold** and a **thin gold** sample in the target position to study direct and indirect capture.

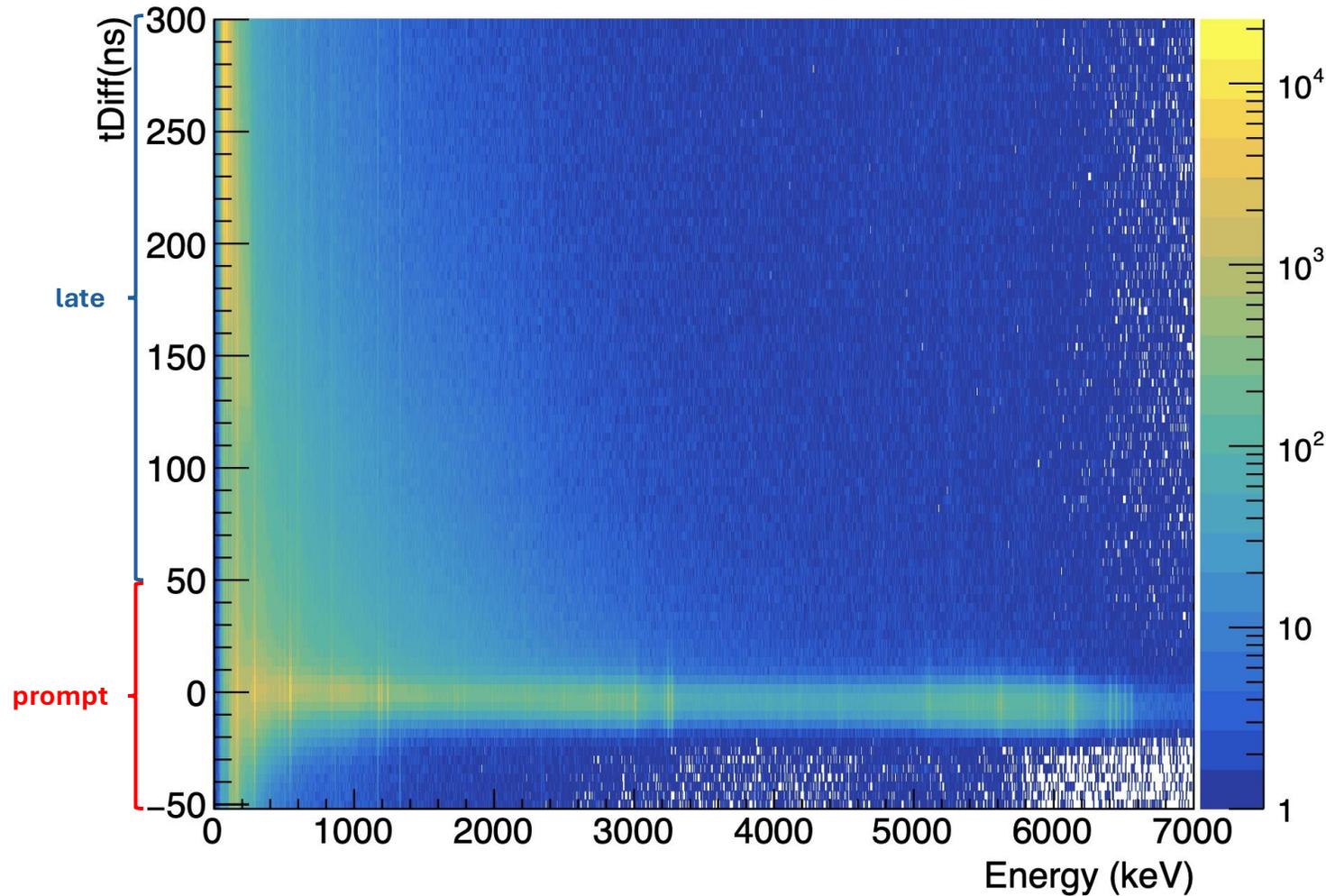
The thicker the target the more the directly detected muons (bigger part of the muon tail detected),



Directly detected muons the rest of them get captured by the gas mixture.



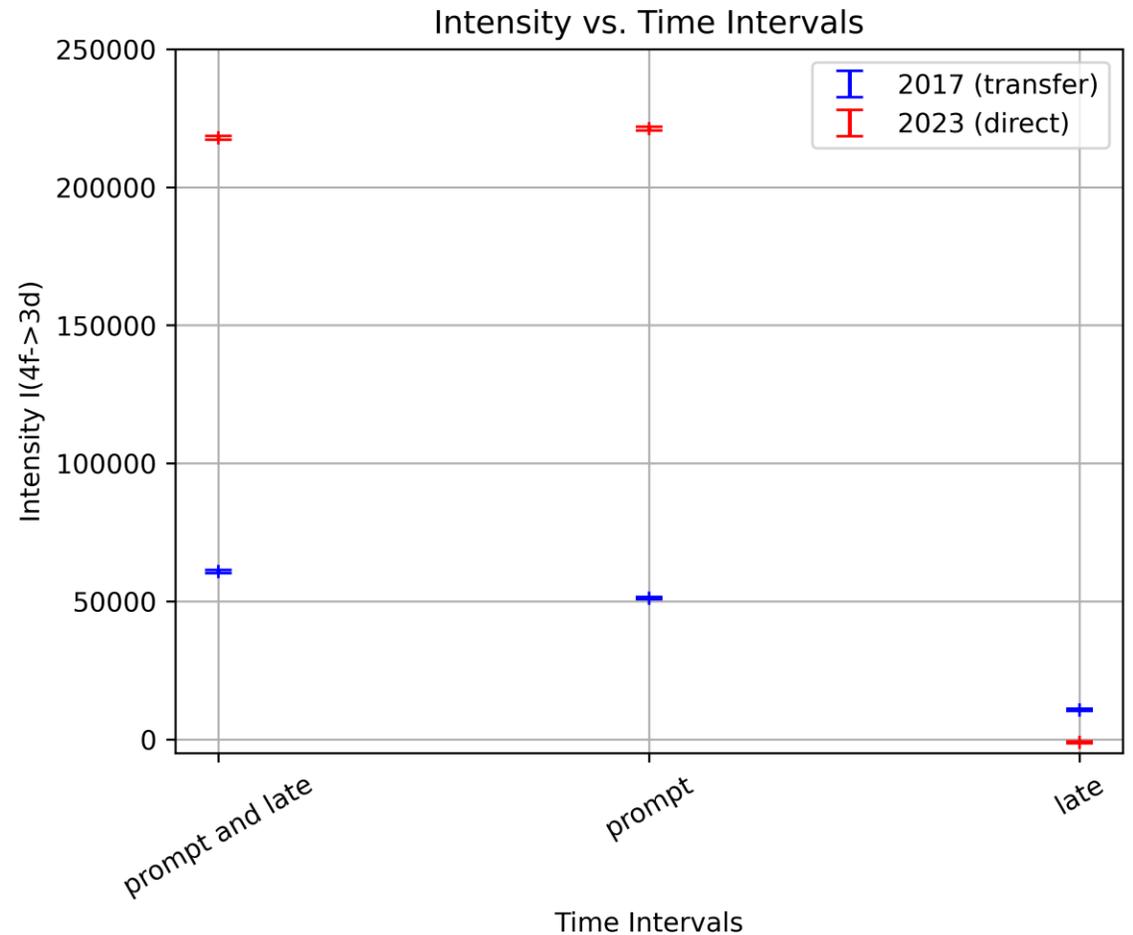
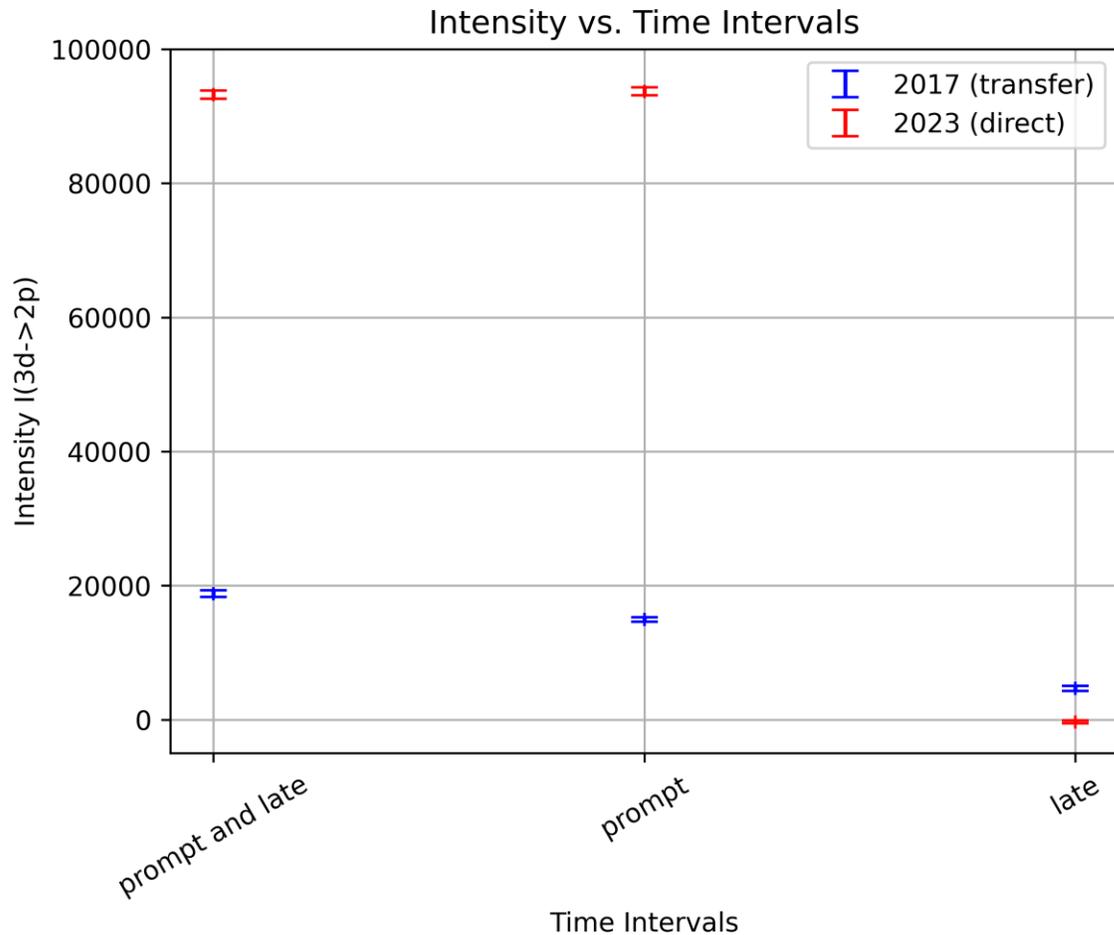
# HPGe detectors, Direct Muon Capture



High Purity Germanium Detectors (HPGe) show bad time resolution at low energies  
HPGe detectors resolution: 500-1000 eV at Energy < 100 keV



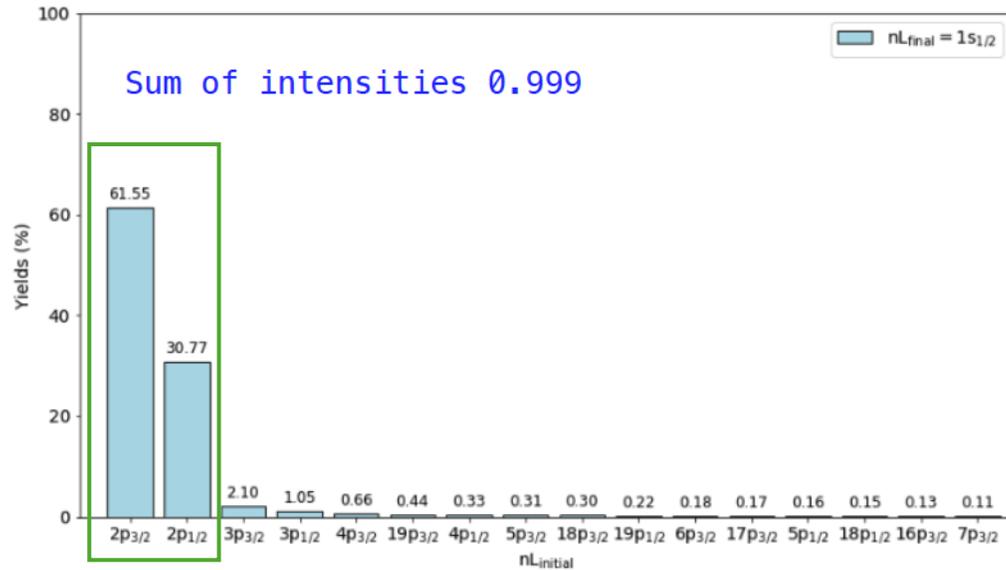
# Intensity of $3d \rightarrow 2p$ and $4f \rightarrow 3d$ transitions versus Time



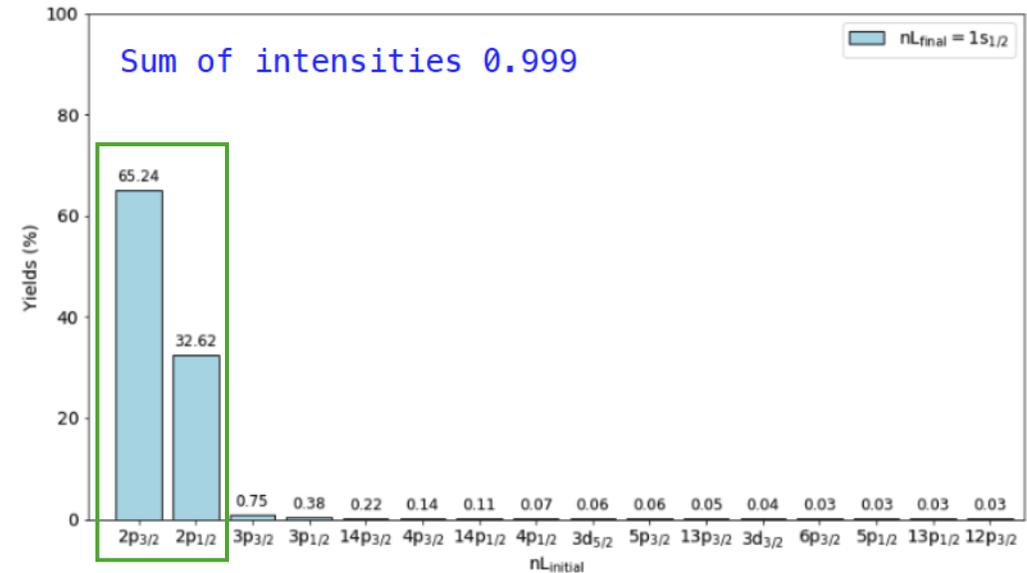
- 2017 data, transfer muon capture at a prompt and late time interval
- 2023 data, direct muon capture at a prompt time interval

# Results of Cascade Simulation for $^{238}\text{U}$ : $2p \rightarrow 1s$ Transition

After transfer



After direct capture



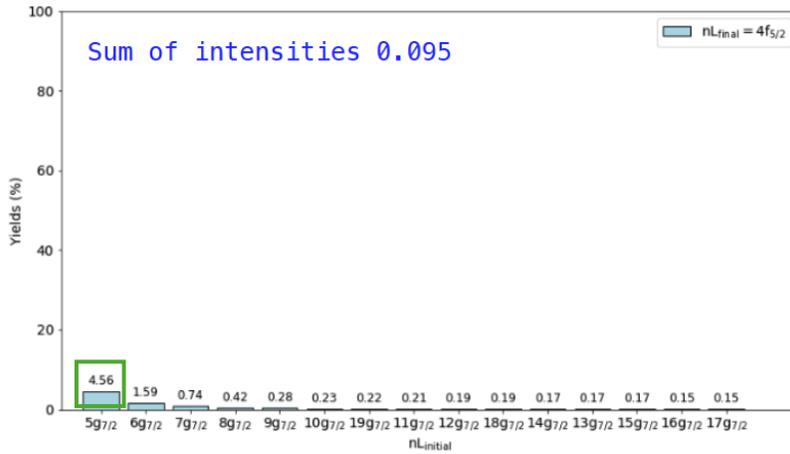
Capture	Transition	Yield (%)
Transfer	$2p_{3/2} \rightarrow 1s_{1/2}$	61.55
	$2p_{1/2} \rightarrow 1s_{1/2}$	30.77
Direct	$2p_{3/2} \rightarrow 1s_{1/2}$	65.24
	$2p_{1/2} \rightarrow 1s_{1/2}$	32.62



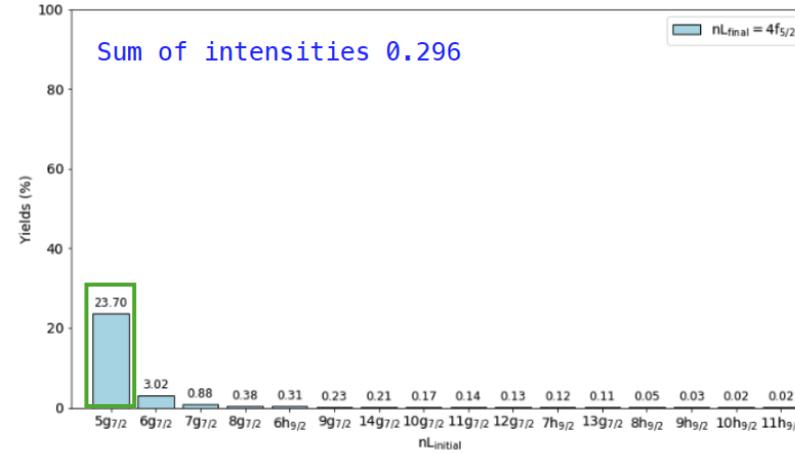
Capture	Transition	Yield (%)
Transfer	$2p \rightarrow 1s$	92.3
Direct	$2p \rightarrow 1s$	97.9

# Results of Cascade Simulation for $^{238}\text{U}$ : $5g \rightarrow 4f$ Transition

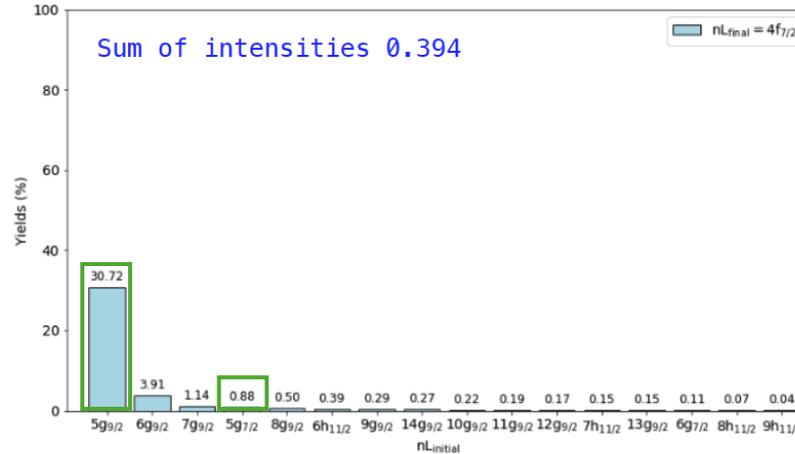
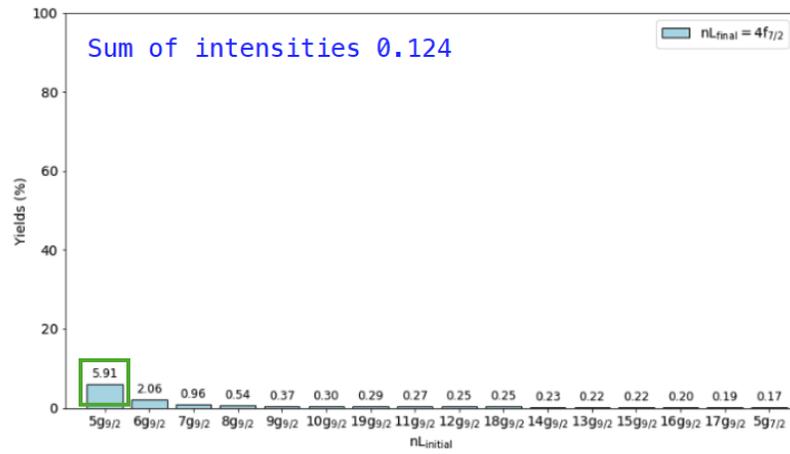
After transfer



After direct capture

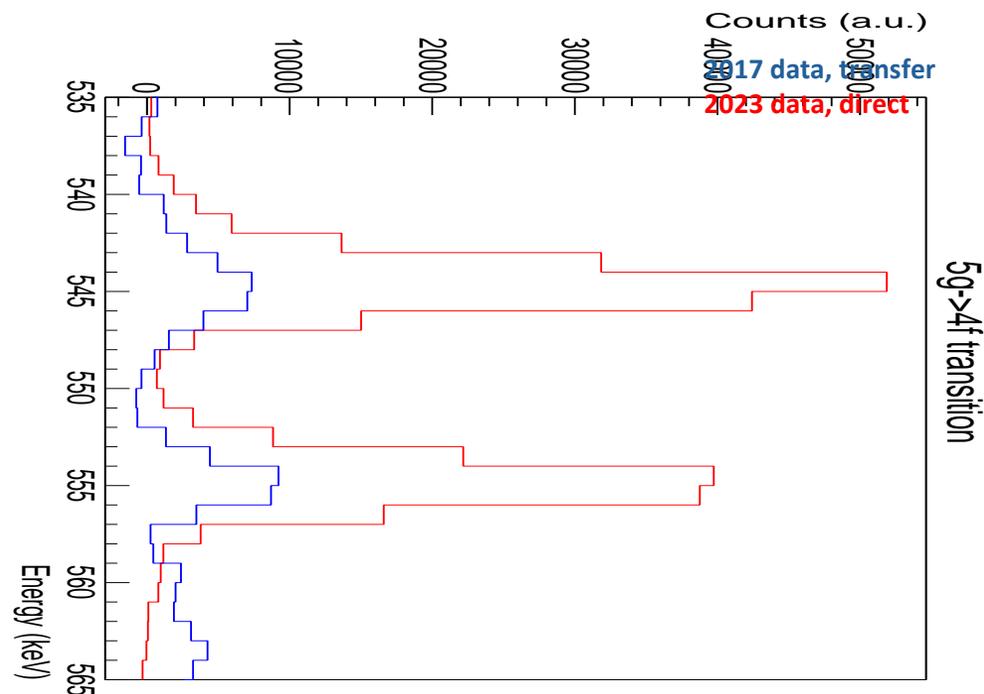
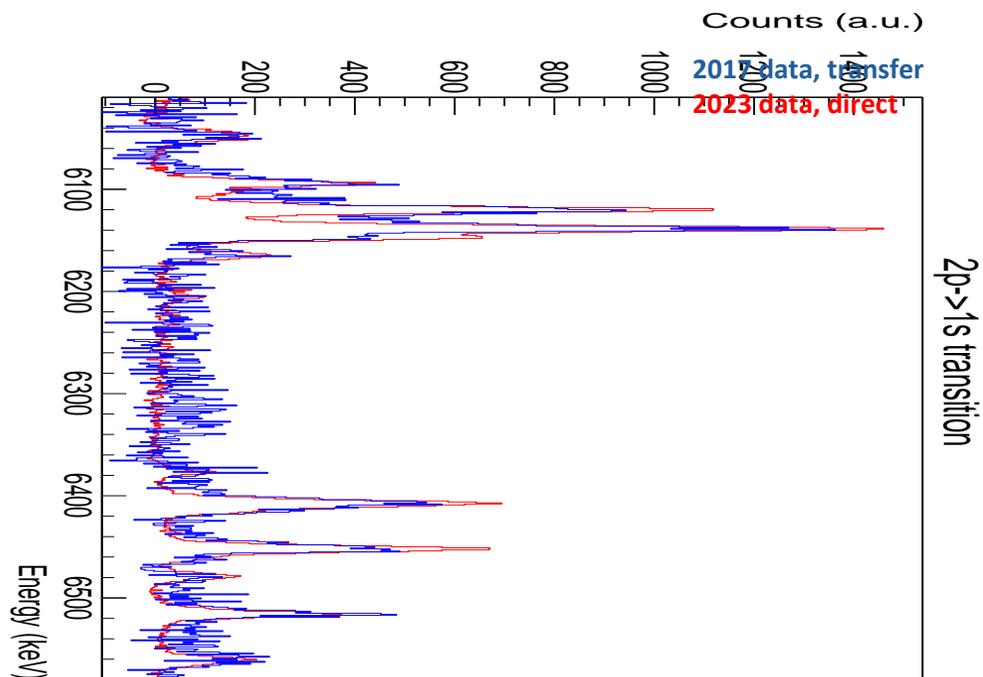


Capture	Transition	Yield (%)
Transfer	5g <sub>7/2</sub> ->4f <sub>5/2</sub>	4.56
	5g <sub>9/2</sub> ->4f <sub>7/2</sub>	5.91
Direct	5g <sub>7/2</sub> ->4f <sub>5/2</sub>	23.70
	5g <sub>9/2</sub> ->4f <sub>7/2</sub>	30.72
	5g <sub>7/2</sub> ->4f <sub>7/2</sub>	0.88



Capture	Transition	Yield (%)
Transfer	5g->4f	10.5
Direct	5g->4f	55.3

# 5g → 4f transition, Transfer vs Direct Muon capture



Capture	Transition	Yield (%)
Transfer	2p->1s	92.3
Direct	2p->1s	97.9

Capture	Transition	Yield (%)
Transfer	5g->4f	10.5
Direct	5g->4f	55.3

Normalized with respect to 2p->1s

Capture	Transition	Ratio
Transfer	(5g->4f)/(2p->1s)	0.11
Direct	(5g->4f)/(2p->1s)	0.56

x ~ 5

# 2p → 1s Transition: Removing Oxygen Background

We are breaking the 2p->1s transition into two **integrals** and we sum them up. This way we get rid of the **oxygen region**.

- **Oxygen region** there is some  $^{16}\text{O}$  leaking in the gas mixture. Muons get captured by the oxygen nucleus and produce  $^{15}\text{N}^* + n + \nu$ , excited nucleus de-excites to its ground state and produces some  $\gamma$ -rays (peaks between the two integrals).

