

Guiding center hydrodynamics with spin

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ON QCD IN EXTREME CONDITIONS



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Outline

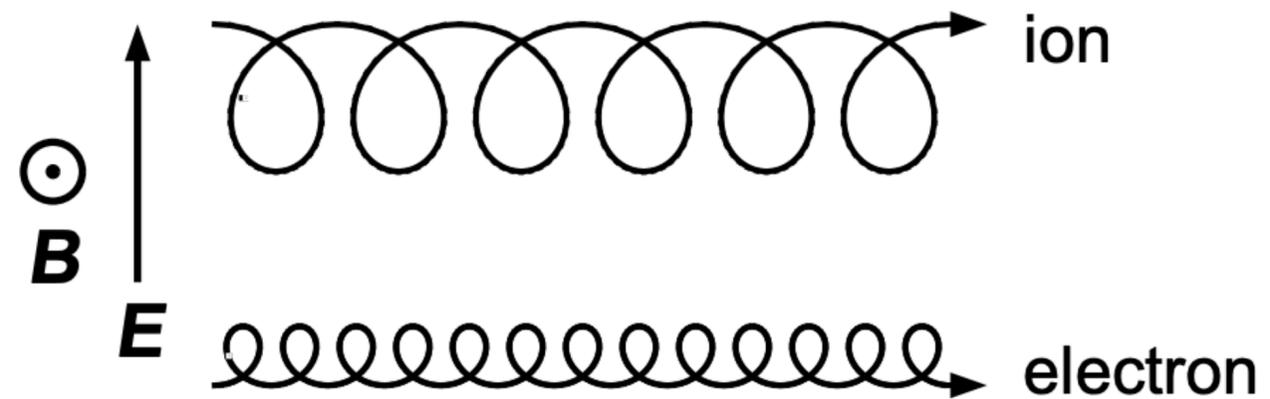
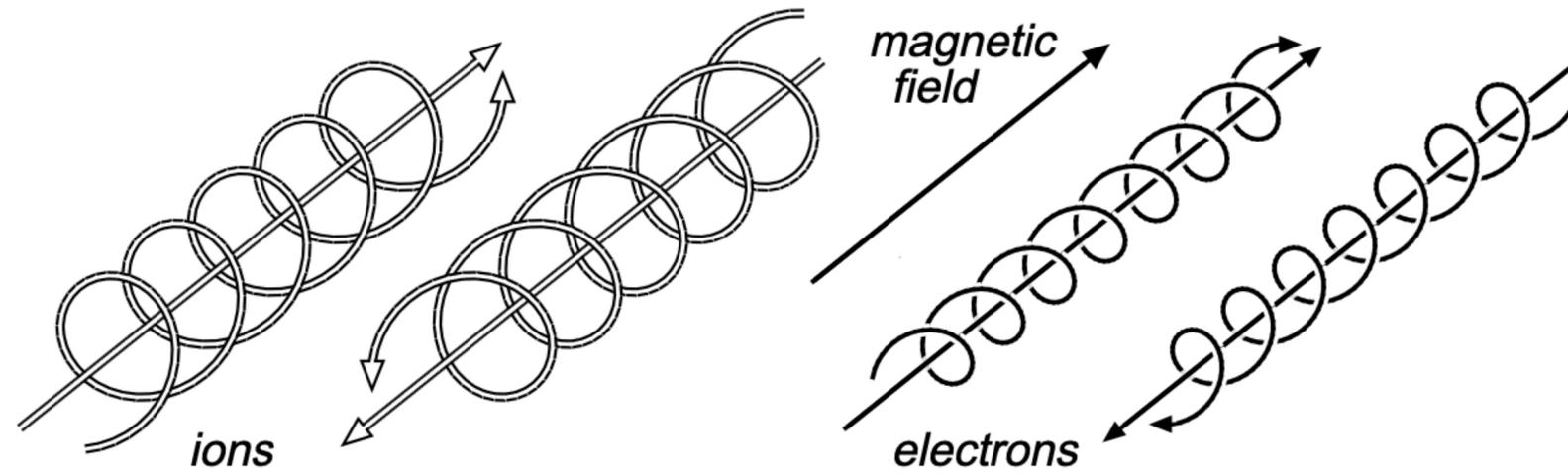
Guiding-center

Guiding-center hydrodynamics

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Guiding-center hydrodynamics with spin

Guiding-center



Assumptions

A plasma of particles immersed in a strong magnetic field can be envisioned as a collection of guiding centers. To effectively substitute the motion of particles with that of their guiding centers, the cyclotron radius should be smaller than the mean free path.

We are in a very large magnetic field and in a reference frame where $\mathbf{E} \parallel \mathbf{B}$.

We focus on a plasma composed of a single species of charged spin-half particles, with a stationary background charge density & static EM field (for simplicity)

Guiding-center hydrodynamics

→ We use classical single-particle distribution function in a phase space extended to spin

→ Identifying the collisional invariants of the local Boltzmann equation

$$f^{\pm}(x, p, s) = \exp(-\beta(x) \cdot p \pm \xi(x)) \left[1 + \frac{1}{2} \omega_{\mu\nu}(x) s^{\mu\nu} \right]$$

“+” means particle contribution

“-” means antiparticle contribution

Spin potential is assumed small (for simplicity)

$$s^{\mu\nu} = (1/m) \epsilon^{\mu\nu\alpha\beta} p_{\alpha} s_{\beta}, \quad \xi(x) = \mu_B/T, \quad \beta_{\mu}(x) = U^{\mu}/T$$

Guiding-center hydrodynamics

$$N^\mu = \int dP dS (p^\mu + v^\mu) (f^+ - f^-) = N_{\parallel}^\mu + N_{\perp}^\mu$$


Net current

$$T^{\mu\nu} = \int dP dS \left(p^\mu p^\nu + \overline{p_{\perp}^\mu p_{\perp}^\nu} \right) (f^+ + f^-) = \tilde{T}^{\mu\nu} + T_{\perp}^{\mu\nu}$$


Energy-momentum tensor

$$S^{\lambda,\mu\nu} = \int dP dS p^\lambda s^{\mu\nu} (f^+ + f^-)$$


Spin tensor

Guiding-center hydrodynamics

$$\rightarrow N^\mu = \int dP dS (p^\mu + v^\mu) (f^+ - f^-) = N_{\parallel}^\mu + \underbrace{N_{\perp}^\mu}_{\text{drift current}}$$

$$v^\mu = B_*^{-2} \left(p_\nu p^\lambda \partial_\lambda B^{\mu\nu} + \frac{\tilde{m}^2 - m^2}{2B_*} B^{\mu\nu} \partial_\nu B_* \right)$$

$$dP = \frac{d^4 p}{2\pi^2} \underbrace{\theta(p_0) \delta(p^2 - \tilde{m}^2)}_{\text{takes care of the mass-shell condition}} \delta^2(\Delta^\mu{}_\nu p^\nu) \left(B_* + \frac{p_\mu \partial_\nu B^{\mu\nu}}{B_*} \right)$$

takes care of the mass-shell condition

$$p_0^2 = p_{\parallel}^2 + \tilde{m}^2$$

$$\tilde{m}^2 = \overline{p_{\perp}^2} + m^2$$

forces the phase space to have only one momentum direction, $\mathbf{p} \parallel \mathbf{B}$

$$\Delta^\mu{}_\nu = (1/B_*^2) B^{\mu\alpha} B_{\nu\alpha}$$

projector onto the transverse plane (x,y)
perpendicular to \mathbf{B} plane in the $\mathbf{E} \parallel \mathbf{B}$ frame

$$\int dS = \frac{2m}{\sqrt{3}\pi} \int d^4 s \delta\left(s \cdot s + \frac{3}{4}\right) \delta(p \cdot s) = 2$$

normalization of s

orthogonality condition

$$N_{\parallel}^\mu = n u^\mu = 4 \sinh(\xi) n_0 u^\mu = 4 \sinh(\xi) u^\mu \frac{B_*}{(2\pi)^2} \tilde{m} K_1(\tilde{z})$$

$$\tilde{z} = \tilde{m}/T$$

$$N_{\perp}^\mu = B_*^{-2} (\tilde{T}^\lambda{}_\nu \partial_\lambda B^{\mu\nu} - M B^{\mu\nu} \partial_\nu B_*)$$

$$\partial_\mu N^\mu = 0$$

$$\tilde{\Delta}^\mu{}_\nu = \delta^\mu{}_\nu - \Delta^\mu{}_\nu \text{ projector onto the 1+1 } \mathbf{E} \parallel \mathbf{B} \text{ space, i.e., (t,z) hyperplane}$$

Guiding-center hydrodynamics

$$T^{\mu\nu} = \int dP dS \left(p^\mu p^\nu + \overline{p_\perp^\mu p_\perp^\nu} \right) (f^+ + f^-) = \tilde{T}^{\mu\nu} + T_\perp^{\mu\nu}$$

$$\tilde{T}^{\mu\nu} = 4 \cosh(\xi) \int dP p^\mu p^\nu e^{-\beta \cdot p} = (\epsilon + P) u^\mu u^\nu - P \tilde{\Delta}^{\mu\nu}$$

$$\epsilon = 4 \cosh(\xi) \int dP (u \cdot p)^2 e^{-\beta \cdot p} = 4 \cosh(\xi) \frac{B_*}{(2\pi)^2} \tilde{m} T (\tilde{z} K_0(\tilde{z}) + K_1(\tilde{z}))$$

$$T_\perp^{\mu\nu} = -2 \cosh(\xi) \int dP \Delta^{\mu\nu} p_\perp^2 e^{-\beta \cdot p} = -P_\perp \Delta^{\mu\nu}$$

$$P = 4 \cosh(\xi) \int dP \frac{1}{3} ((u \cdot p)^2 - p \cdot p) e^{-\beta \cdot p} = 4 \cosh(\xi) \frac{B_*}{(2\pi)^2} \tilde{m} T K_1(\tilde{z})$$

$$P_\perp = 2 \cosh(\xi) \frac{B_*}{(2\pi)^2} (\tilde{m}^2 - m^2) K_0(\tilde{z}) = -M B_*$$

$$\partial_\mu T^{\mu\nu} = F^{\nu\mu} J_\mu \implies \tilde{\Delta}^\sigma{}_\nu \partial_\mu T^\mu{}_\sigma = E_{\nu\lambda} N_\parallel^\lambda$$

spin potential
contribution
appears in higher-
order

$\tilde{\Delta}^\mu{}_\nu = \delta^\mu{}_\nu - \Delta^\mu{}_\nu$ projector onto the 1+1 $\mathbf{E} \parallel \mathbf{B}$ space, i.e., (t,z) hyperplane

Guiding-center hydrodynamics with spin

$$\begin{aligned}\rightarrow S^{\lambda,\mu\nu} &= \int dP dS p^\lambda s^{\mu\nu} (f^+ + f^-) \\ &= 2 \cosh(\xi) \int dP p^\lambda e^{-\beta \cdot p} \int dS s^{\mu\nu} \left(1 + \frac{1}{2} \omega_{\alpha\beta} s^{\alpha\beta} \right) \\ &= \frac{P}{4} \beta^\lambda \omega^{\mu\nu} + T^2 \left((\epsilon + P) \frac{3 T^2}{\tilde{m}^2} + \frac{P}{2} \right) \beta^\lambda \beta^\delta \beta^{[\mu} \omega^{\nu]}_\delta - \frac{T^2 (\epsilon + P)}{2\tilde{m}^2} \left\{ \tilde{\Delta}^{\lambda[\mu} \omega^{\nu]}_\delta \beta^\delta + \tilde{\Delta}^{\delta[\mu} \omega^{\nu]}_\delta \beta^\lambda + \tilde{\Delta}^{\lambda\delta} \beta^{[\mu} \omega^{\nu]}_\delta \right\}\end{aligned}$$

$$\partial_\lambda S^{\lambda,\mu\nu} = 0$$

Summary & Outlook

Incorporated spin degrees of freedom into the guiding center (ideal) hydrodynamics

In this restrictive scenario, number of EoMs are less than the usual hydro formalism

Easiest method to include EM field contribution in the hydrodynamics

Ongoing: include dissipation and dynamical EM using Maxwell equations

Ongoing: coupling between spin and background EoMs

Thank you for listening!

Entropy current

$$\rightarrow \mathfrak{S}^\mu = - \int dP dS p^\mu (f^+ \log(f^+) + f^- \log(f^-)) + \int dP dS p^\mu (f^+ + f^-)$$

$$\partial_\mu \mathfrak{S}^\mu = T^{\mu\alpha} \partial_\mu \beta_\alpha + \beta_\alpha \left(\tilde{\Delta}^{\lambda\alpha} \partial_\mu \tilde{T}^\mu{}_\lambda - E_\lambda^\alpha N_\parallel^\lambda \right) - \frac{1}{2} S^{\mu,\alpha\beta} \partial_\mu \omega_{\alpha\beta} - N^\mu (\coth(\xi))^2 \partial_\mu \xi = 0$$

$$n = \frac{\partial P}{\partial \mu}, \quad \epsilon = T \frac{\partial P}{\partial T} + \mu \frac{\partial P}{\partial \mu} - P, \quad P_\perp = P - B_* \frac{\partial P}{\partial B_*}$$