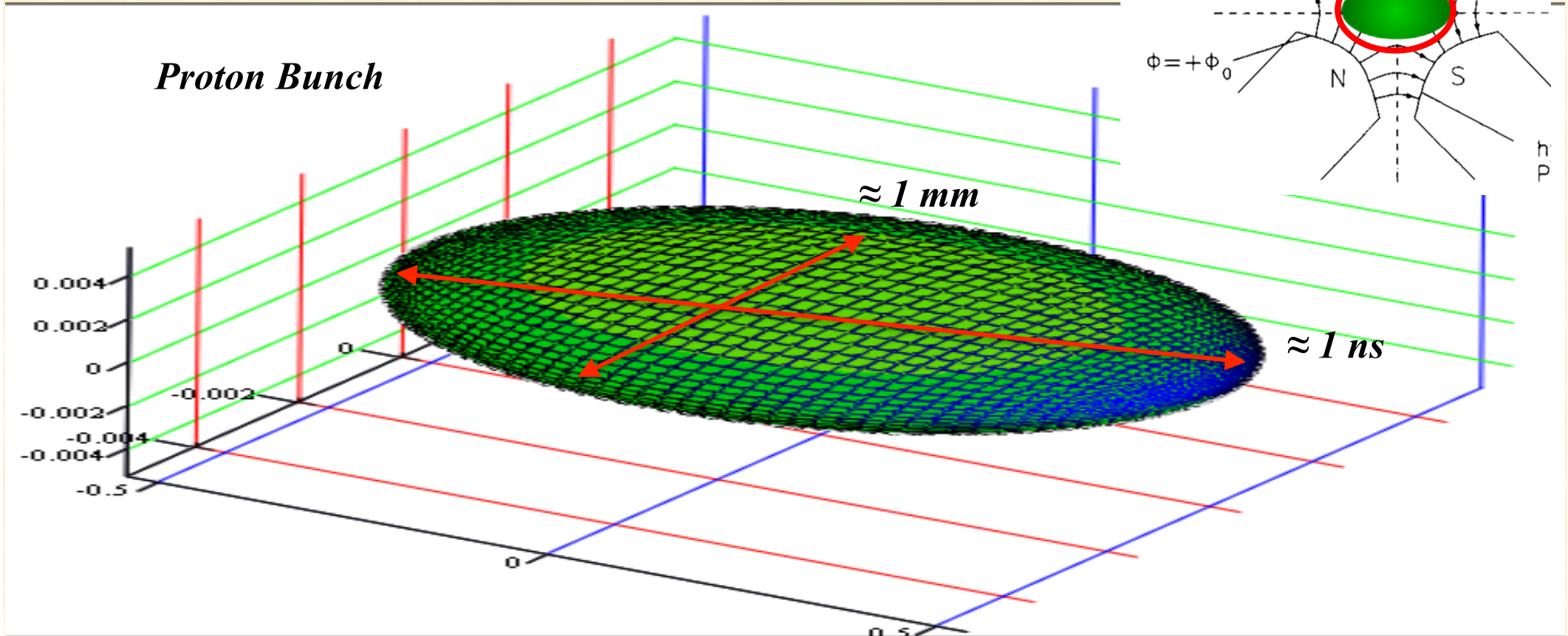


Transverse Beam Optics II

Bernhard Holzer,
CERN

The Ideal World: Particle Trajectories & Beams



Reminder of Part I

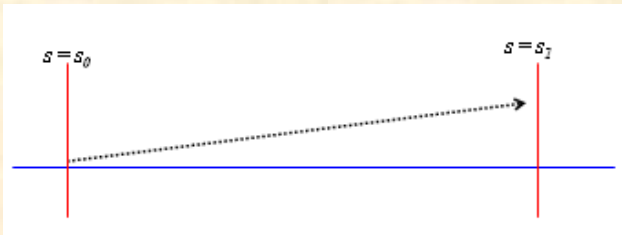
Equation of Motion:

$$\mathbf{x}'' + \mathbf{K} \mathbf{x} = 0 \quad K = \left(\frac{1}{\rho^2} - k\right) = 0 \quad \dots \text{ hor. plane:}$$

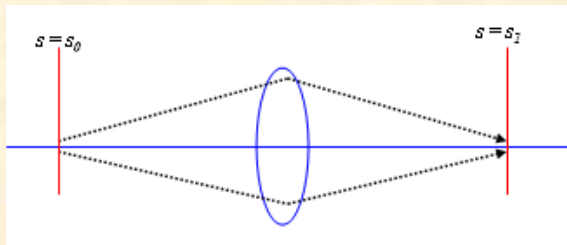
$$K = k \quad \dots \text{ vert. Plane:}$$

Solution of Trajectory Equations

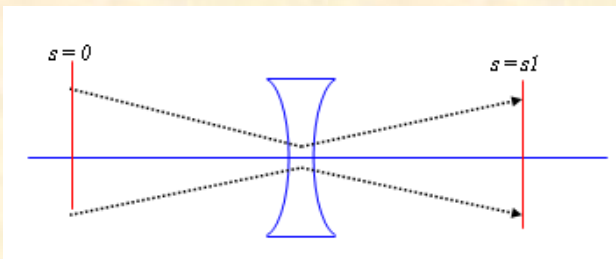
$$\begin{pmatrix} \mathbf{x} \\ \mathbf{x}' \end{pmatrix}_{s_1} = \mathbf{M} * \begin{pmatrix} \mathbf{x} \\ \mathbf{x}' \end{pmatrix}_{s_0}$$



$$\mathbf{M}_{drift} = \begin{pmatrix} 1 & l \\ 0 & 1 \end{pmatrix}$$



$$\mathbf{M}_{foc} = \begin{pmatrix} \cos(\sqrt{|K|}l) & \frac{1}{\sqrt{|K|}} \sin(\sqrt{|K|}l) \\ -\sqrt{|K|} \sin(\sqrt{|K|}l) & \cos(\sqrt{|K|}l) \end{pmatrix}$$



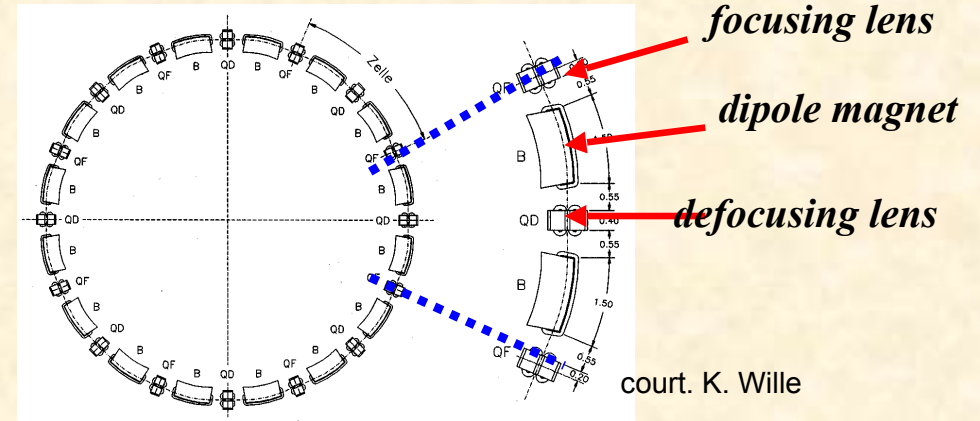
$$\mathbf{M}_{defoc} = \begin{pmatrix} \cosh(\sqrt{|K|}l) & \frac{1}{\sqrt{|K|}} \sinh(\sqrt{|K|}l) \\ \sqrt{|K|} \sinh(\sqrt{|K|}l) & \cosh(\sqrt{|K|}l) \end{pmatrix}$$

Transformation through a system of lattice elements

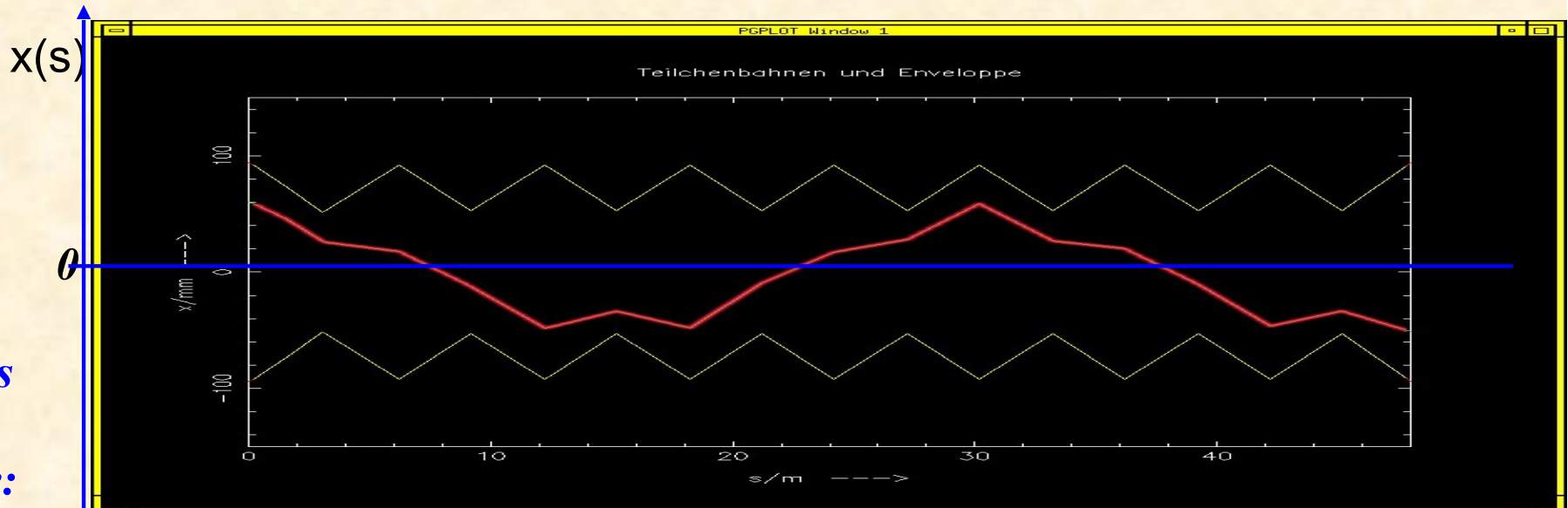
combine the single element solutions by multiplication of the matrices

$$M_{total} = M_{QF} * M_D * M_{QD} * M_{Bend} * M_{D*} * \dots$$

$$\begin{pmatrix} x \\ x' \end{pmatrix}_{s2} = M_{1 \rightarrow 2} \begin{pmatrix} x \\ x' \end{pmatrix}_{s1}$$



in each accelerator element the particle trajectory corresponds to the movement of a harmonic oscillator ,



typical values
in a strong
foc. machine:

$$x \approx \text{mm}, x' \lesssim \text{mrad}$$

Quadrupole Magnets:

required: **focusing forces** to keep trajectories in vicinity of the ideal orbit

linear increasing Lorentz force

linear increasing magnetic field

$$B_y = g x \quad B_x = g y$$

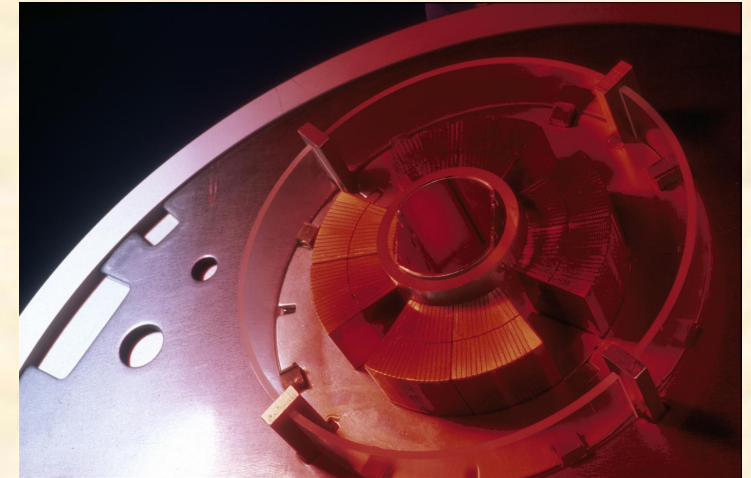
normalised quadrupole field:

gradient of a quadrupole magnet: $g = \left| \frac{\partial B_y}{\partial x} \right|$

normalised gradient $k = \frac{g}{p/e}$

LHC main quadrupole magnet

$$g \approx 25 \dots 220 \text{ T/m}$$

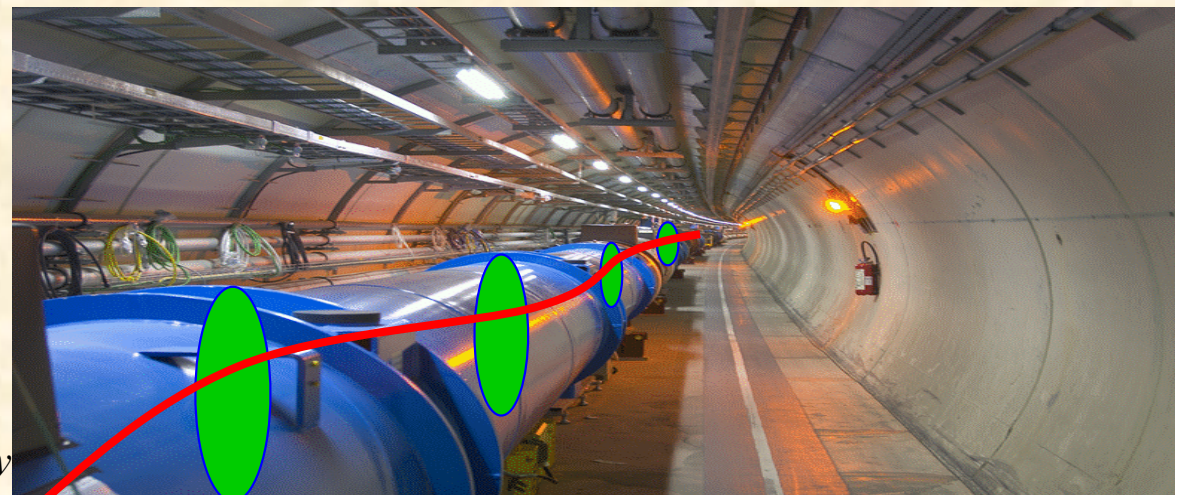


$$\frac{B}{p/q} = \frac{B}{B\rho} = \frac{1}{\rho} \quad k := \frac{g}{p/q}$$

what about the vertical plane:
... Maxwell

$$\vec{\nabla} \times \vec{B} = \cancel{j} + \cancel{\frac{\partial \vec{E}}{\partial t}} = 0$$

$$\Rightarrow \frac{\partial B_y}{\partial x} = \frac{\partial B_x}{\partial y}$$



9.) Orbit & Tune:

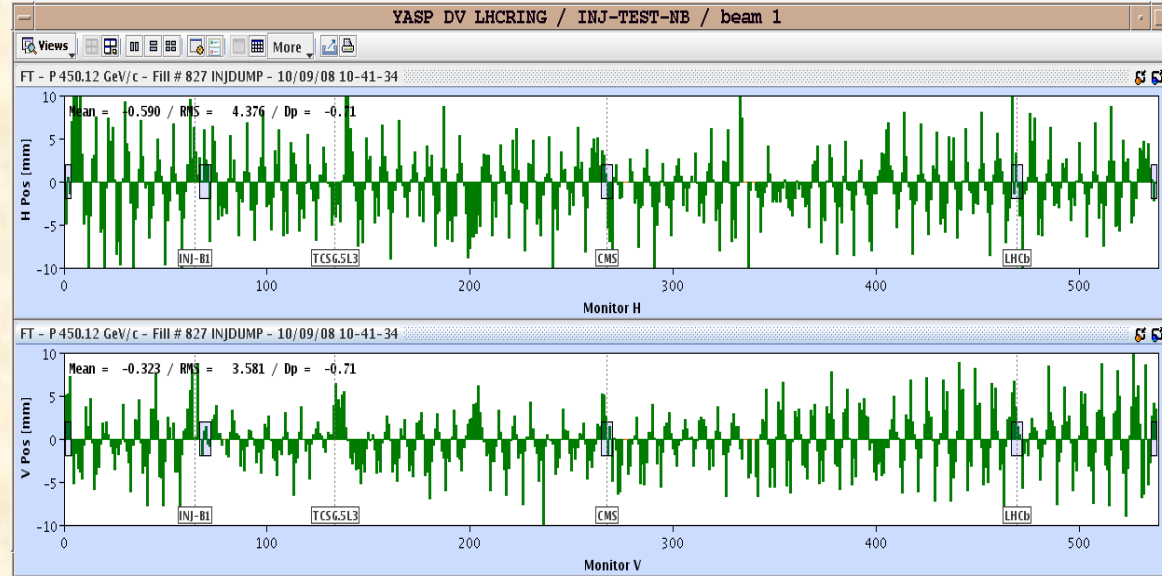
Tune: number of oscillations per turn

64.31

59.32

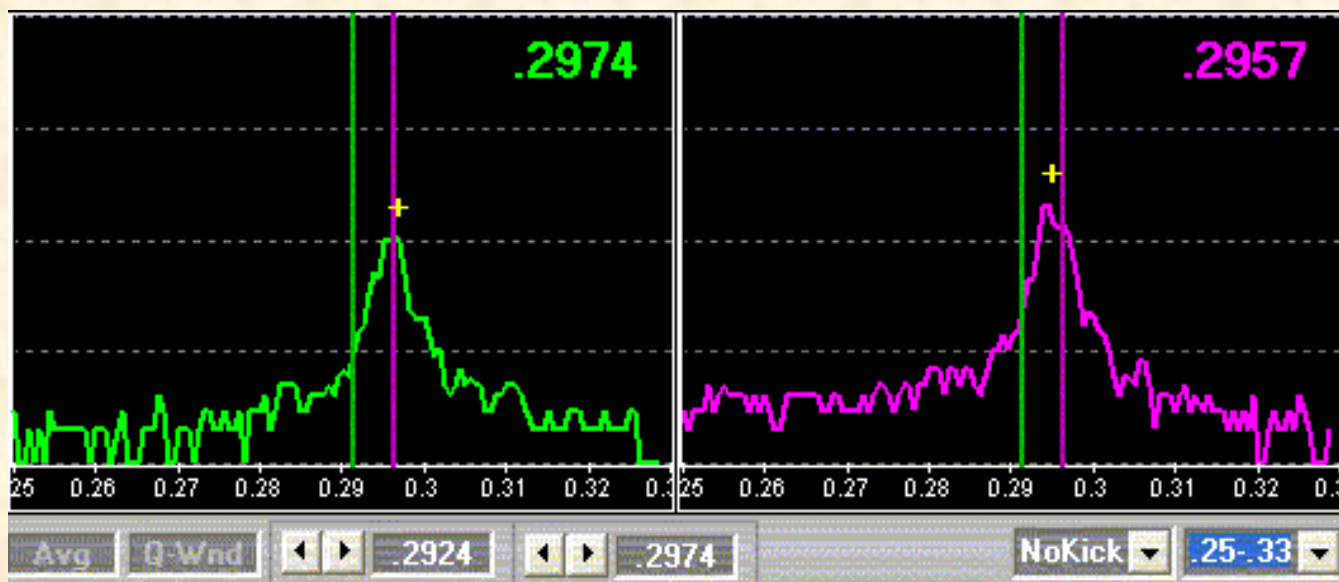
Relevant for beam stability:

non integer part



LHC revolution frequency: 11.3 kHz

$$0.31 * 11.3 = 3.5 \text{ kHz}$$



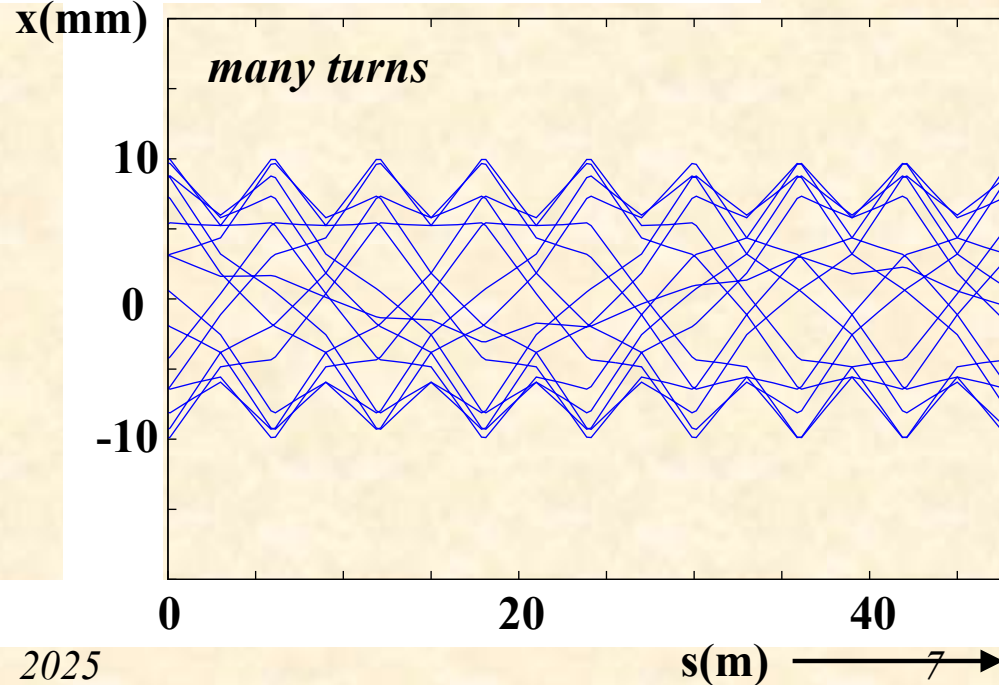
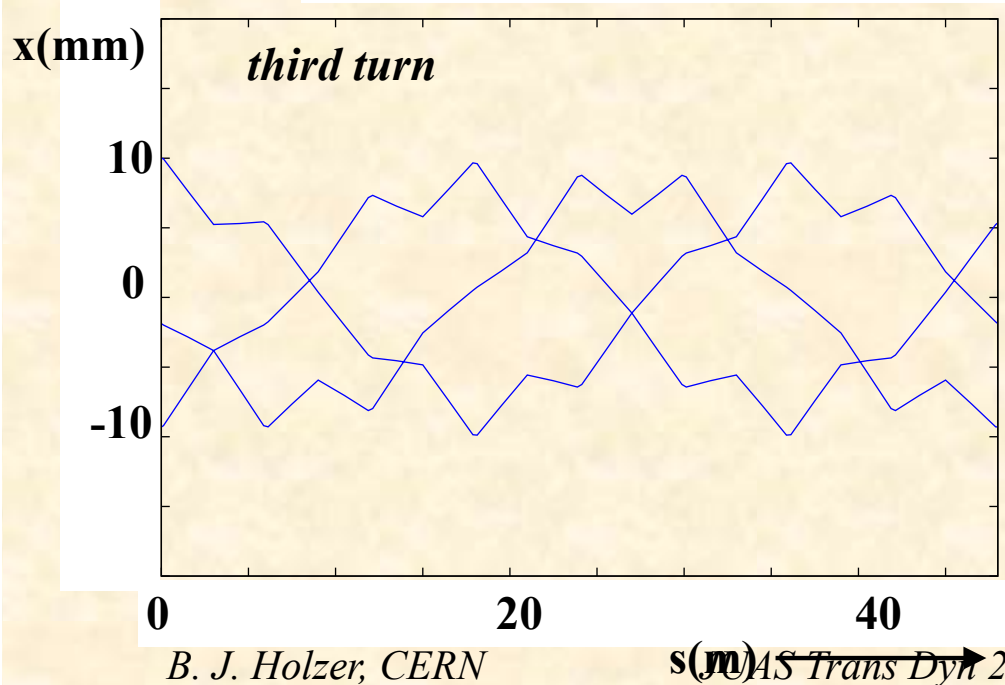
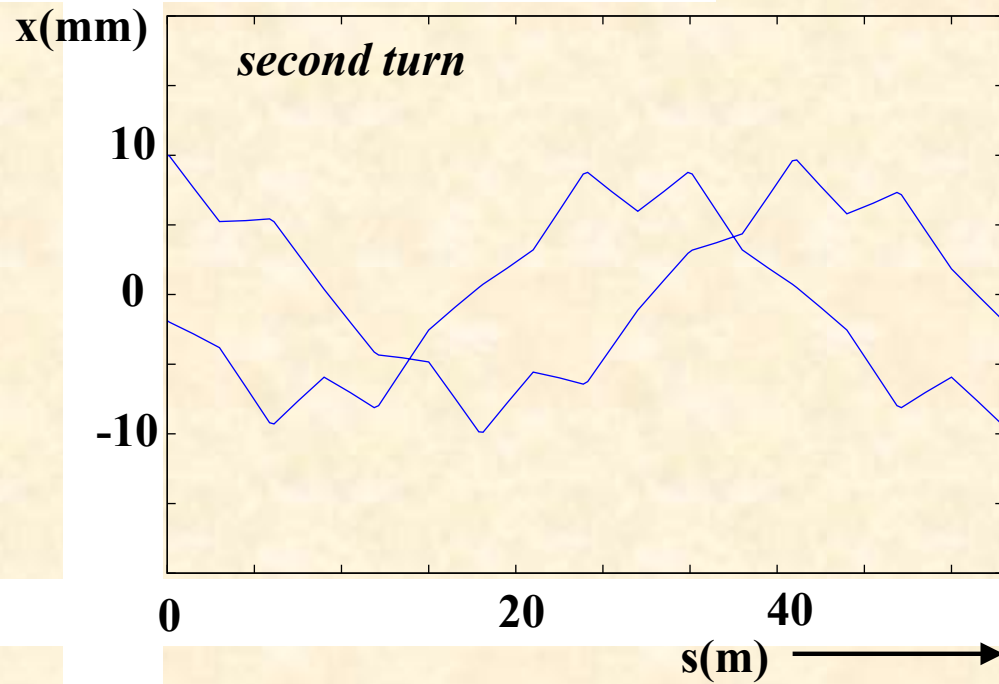
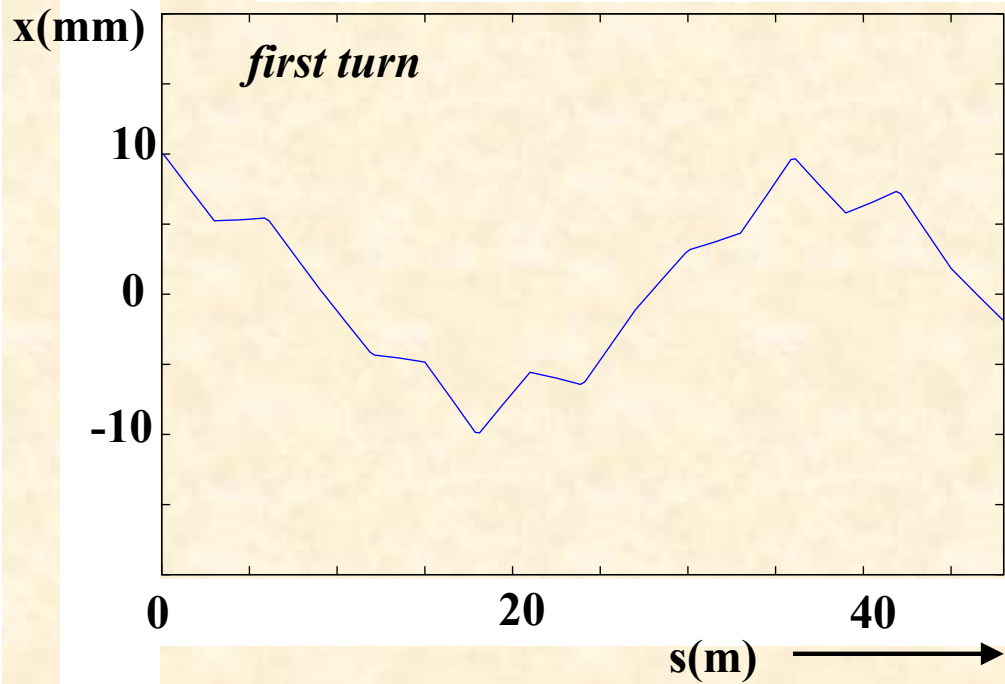
The Symphony of the Beam ???

What would happen if we would connect the tune signal to a loud speaker ?

Multiple Choice:

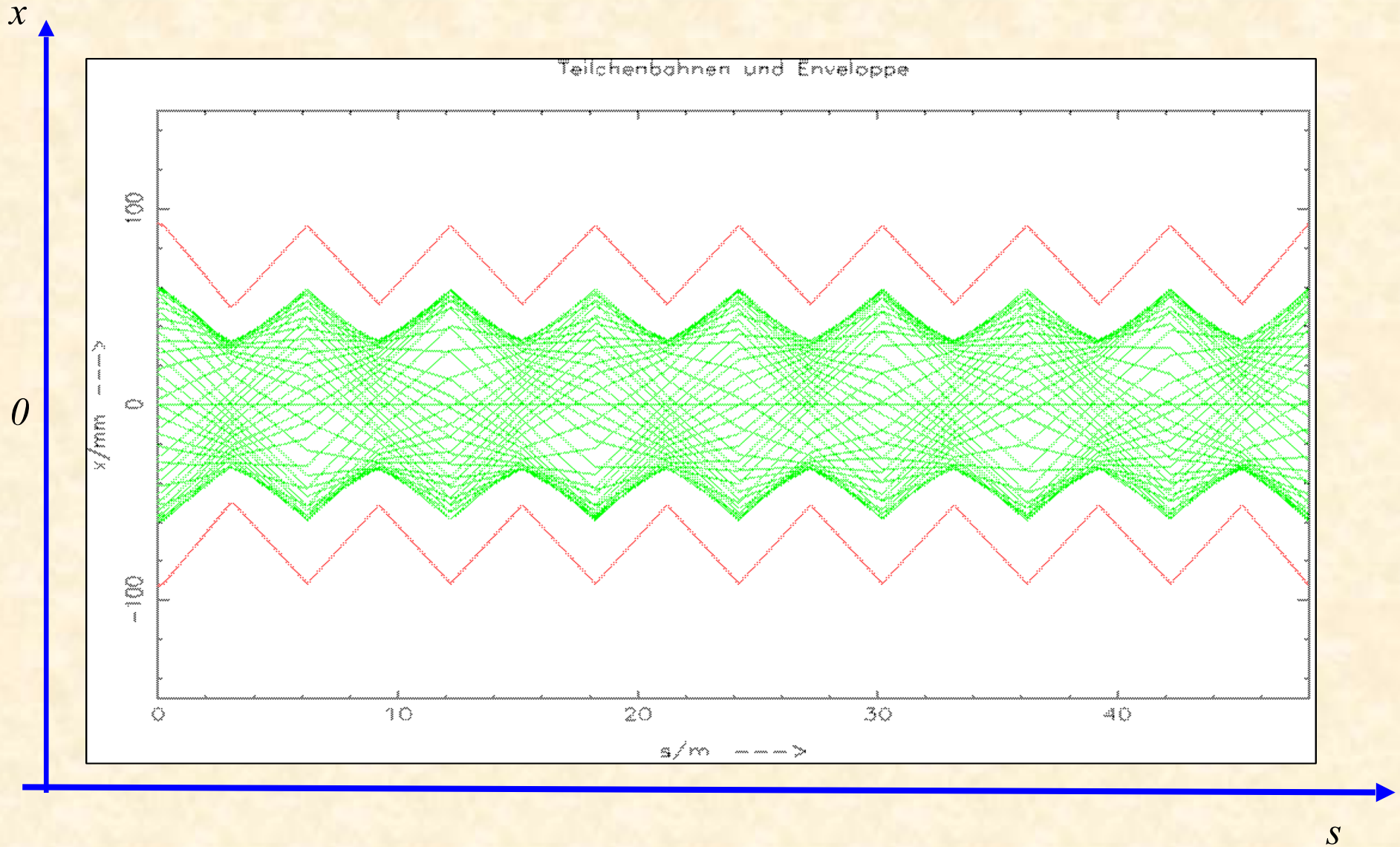
- a) **Nothing, in vacuum we cannot hear sound.**
- b) **we would not hear anything, as the frequency is well beyond our ear's sensitivity.**
- c) **we would hear "the sound of silence", LOL**
- d) **well ...**

Question: what will happen, if the particle performs a second turn ?



Question: what will happen, if the particle performs a second turn ?

... or a third one or ... 10^{10} turns



11.) Hill's Equation:

Astronomer Hill:

*differential equation for motions with periodic focusing properties
„Hill's equation“*

Example: particle motion with
periodic coefficient



equation of motion: $x''(s) + K(s) \cdot x(s) = 0$ *Hill's equation“*

*restoring force \neq const,
 $k(s)$ = depending on the position s
 $k(s+L) = k(s)$, periodic function*

*we expect a kind of quasi harmonic
oscillation: amplitude & phase will depend
on the position s in the ring.*

Hill's Equation:

Hill's equation: the origins

ON THE PART OF THE
MOTION OF THE LUNAR PERIGEE
WHICH IS A FUNCTION OF THE
MEAN MOTIONS OF THE SUN AND MOON
BY
G. W. HILL
in WASHINGTON.

Hill's original paper on orbital mechanics (1886)

12.) The Beta Function

General solution of Hill's equation (... Floquet's theorem, see e.g. cern-94-01)

Ansatz: $x(s) = \sqrt{\varepsilon} \sqrt{\beta(s)} \cdot \cos(\psi(s) + \phi)$ (i)

$\varepsilon, \Phi =$ integration **constants** determined by initial conditions

$\beta(s)$ **periodic function** given by **focusing properties** of the lattice \leftrightarrow quadrupoles

$$\beta(s + L) = \beta(s)$$

Inserting (i) into the equation of motion ...

$$\psi(s) = \int_0^s \frac{ds}{\beta(s)}$$

$\Psi(s) =$ „**phase advance**“ of the oscillation between point „0“ and „s“ in the lattice.
For one complete revolution: number of oscillations per turn „**Tune**“

$$Q_y = \frac{1}{2\pi} \oint \frac{ds}{\beta(s)}$$

The Beta Function

Amplitude of a particle trajectory:

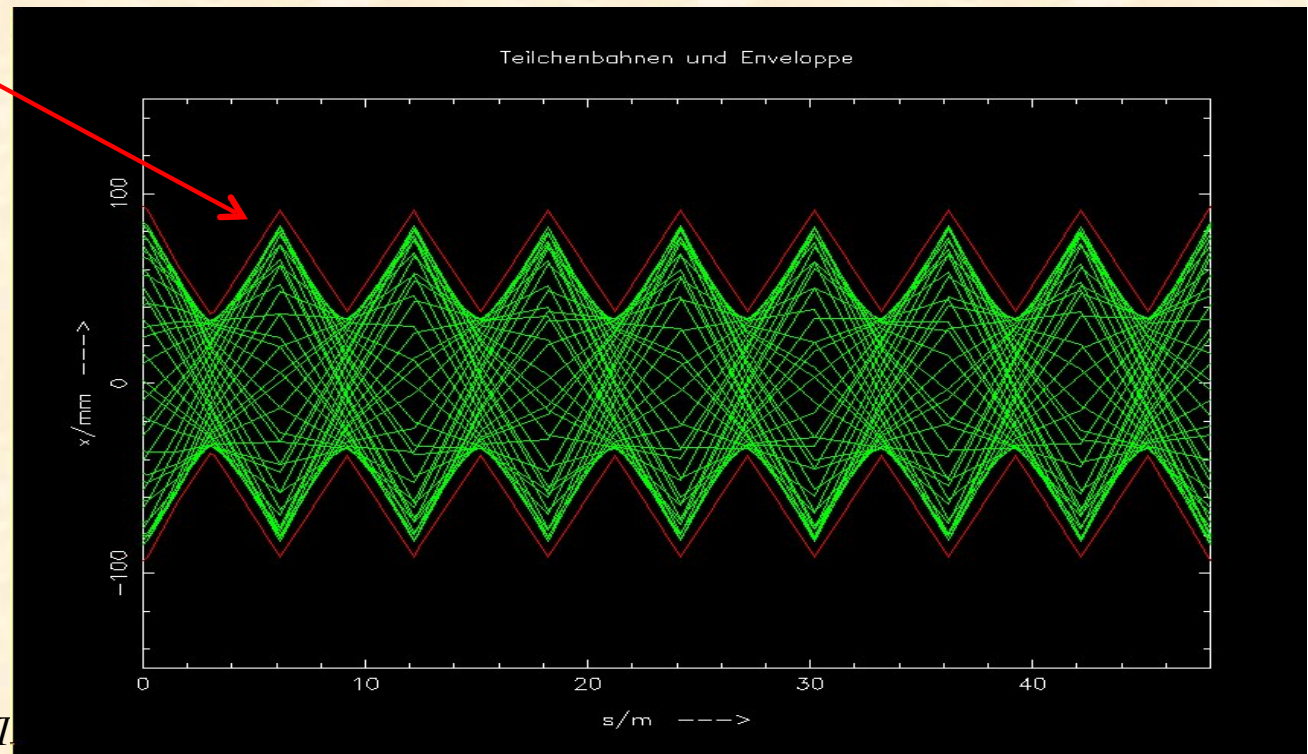
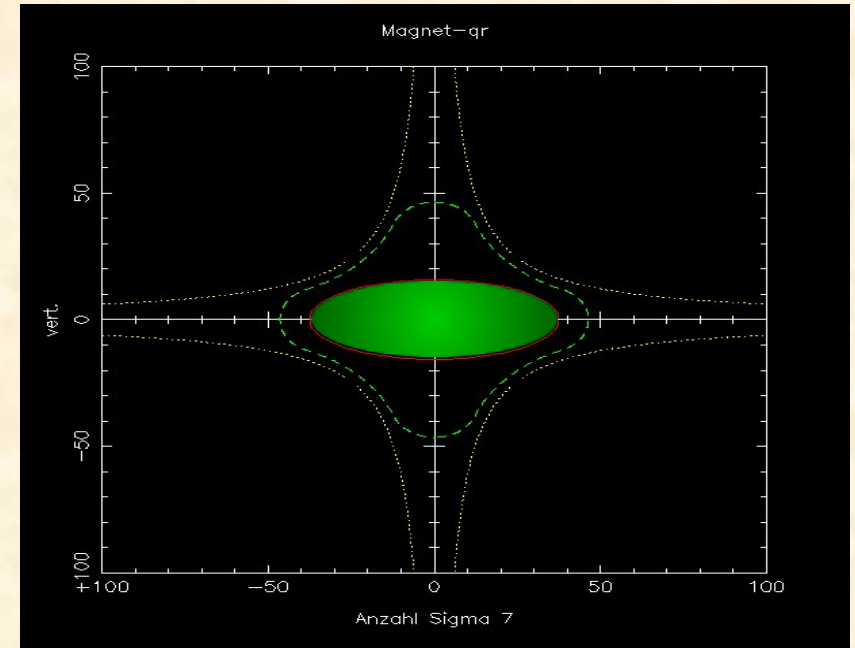
$$x(s) = \sqrt{\varepsilon} \cdot \sqrt{\beta(s)} \cos(\psi(s) + \phi)$$

Maximum size of a particle amplitude

$$\hat{x}(s) = \sqrt{\varepsilon} \sqrt{\beta(s)}$$

β determines the beam size
(... the envelope of all particle trajectories at a given position “s” in the storage ring.

It reflects the periodicity of the magnet structure.



13.) Beam Emittance and Phase Space Ellipse

general solution of Hill equation

$$\left\{ \begin{array}{l} (1) \quad \mathbf{x}(s) = \sqrt{\varepsilon} \sqrt{\beta(s)} \cos(\psi(s) + \phi) \\ (2) \quad \mathbf{x}'(s) = -\frac{\sqrt{\varepsilon}}{\sqrt{\beta(s)}} \left\{ \alpha(s) \cos(\psi(s) + \phi) + \sin(\psi(s) + \phi) \right\} \end{array} \right.$$

from (1) we get

$$\cos(\psi(s) + \phi) = \frac{\mathbf{x}(s)}{\sqrt{\varepsilon} \sqrt{\beta(s)}}$$

$$\alpha(s) = -\frac{1}{2} \beta'(s)$$

$$\gamma(s) = \frac{1 + \alpha(s)^2}{\beta(s)}$$

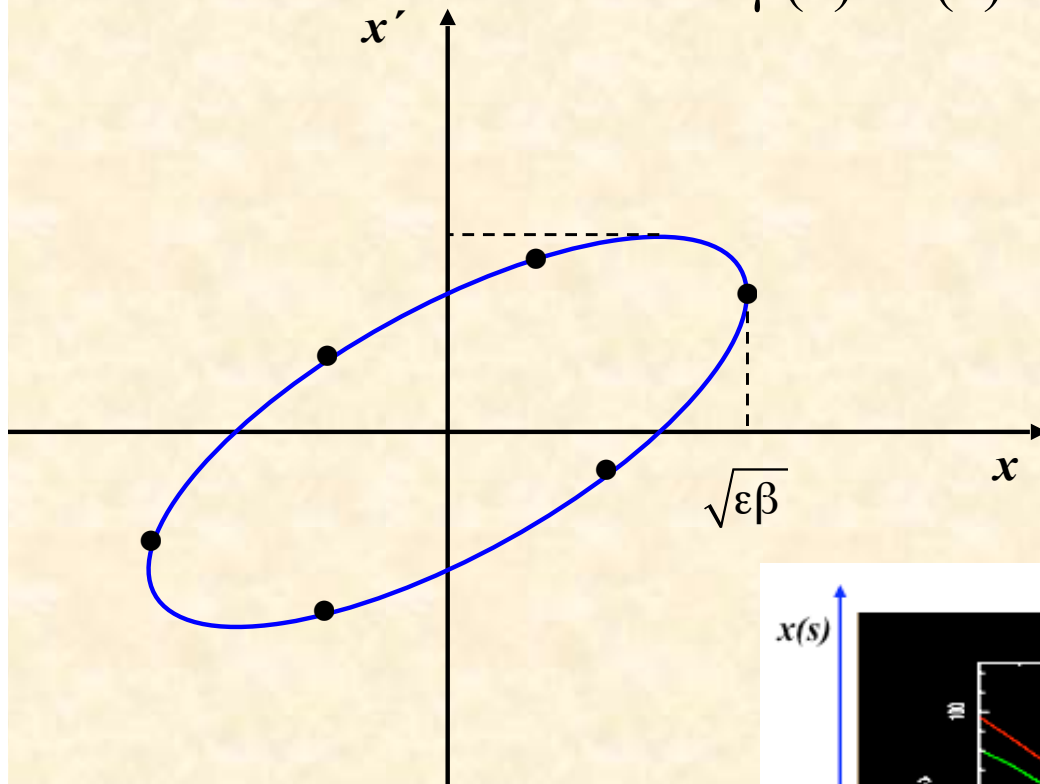
Insert into (2) and solve for ε

$$\varepsilon = \gamma(s) \mathbf{x}^2(s) + 2\alpha(s) \mathbf{x}(s) \mathbf{x}'(s) + \beta(s) \mathbf{x}'^2(s)$$

- * ε is a **constant** of the motion ... it is independent of „s“
- * parametric representation of an **ellipse** in the $x \ x'$ space
- * shape and orientation of ellipse are given by α, β, γ

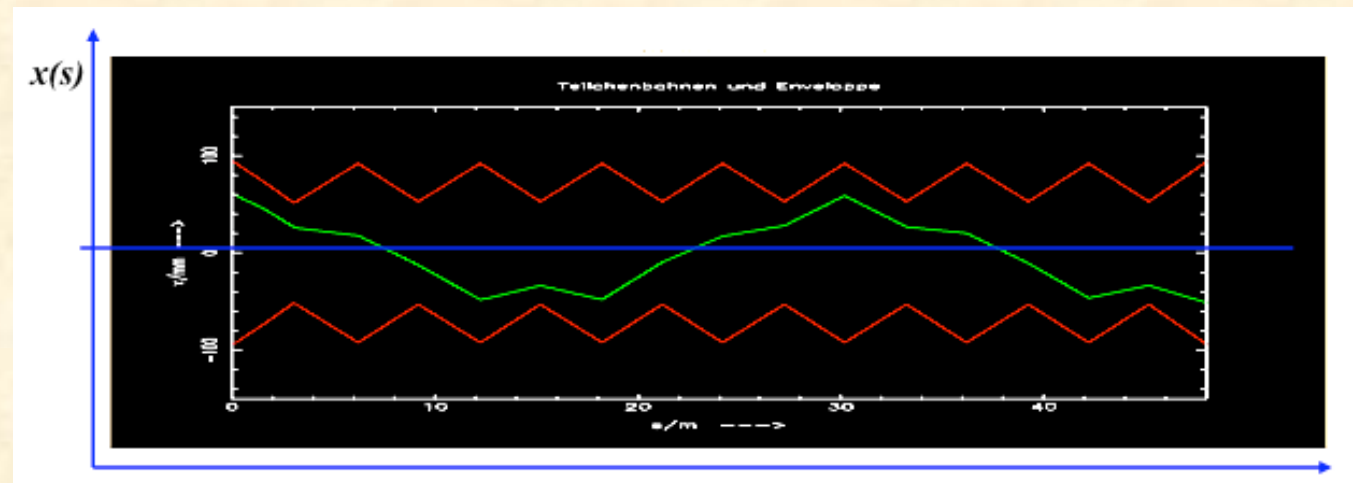
Beam Emittance and Phase Space Ellipse

$$\varepsilon = \gamma(s) x^2(s) + 2\alpha(s)x(s)x'(s) + \beta(s) x'^2(s)$$



Liouville: in reasonable storage rings area in phase space is constant.

$$A = \pi * \varepsilon = \text{const}$$



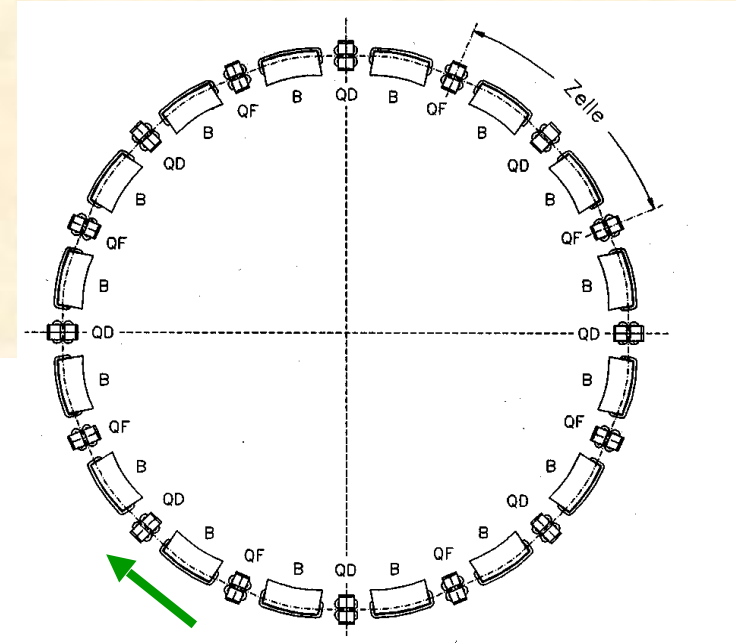
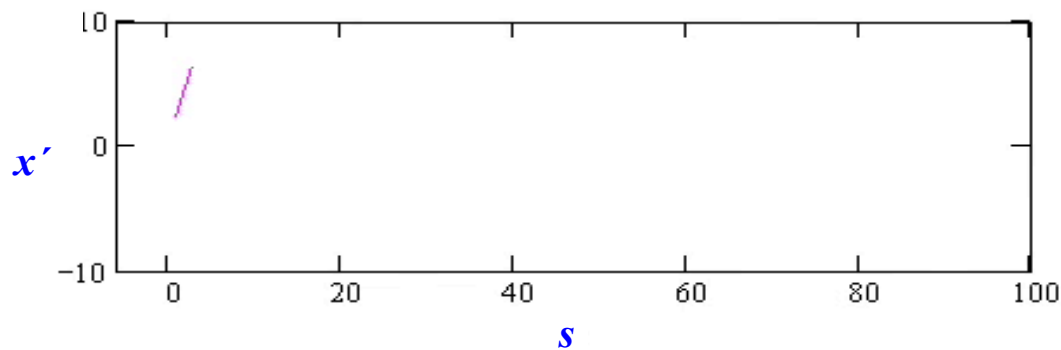
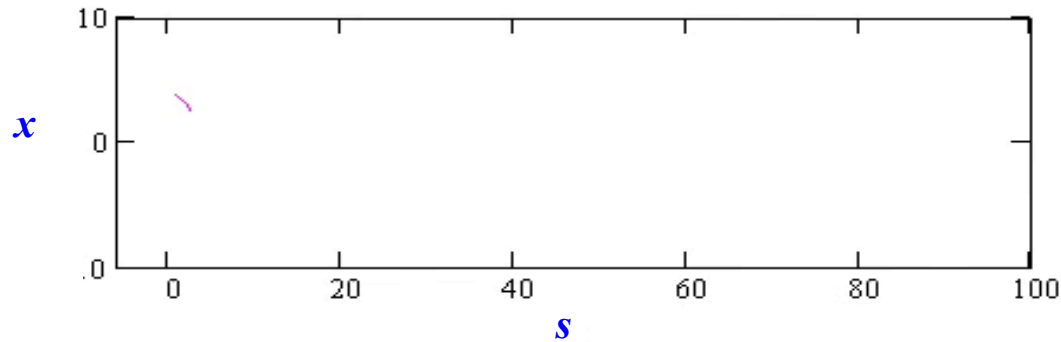
ε beam emittance = *woozilycity* of the particle ensemble, *intrinsic beam parameter*, cannot be changed by the foc. properties.

Scientificquely speaking: area covered in transverse x, x' phase space ... and it is constant !!!

Particle Tracking in a Storage Ring

Calculate x , x' for each linear accelerator element according to matrix formalism

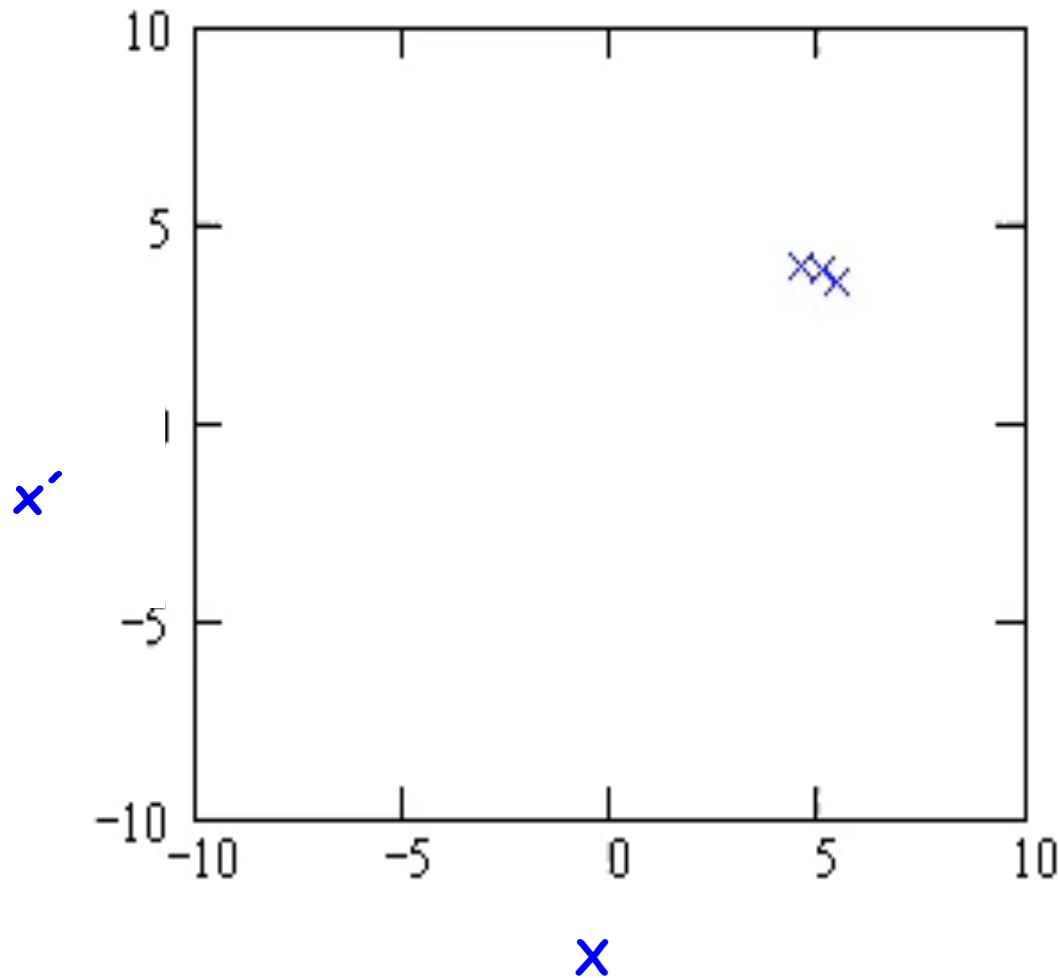
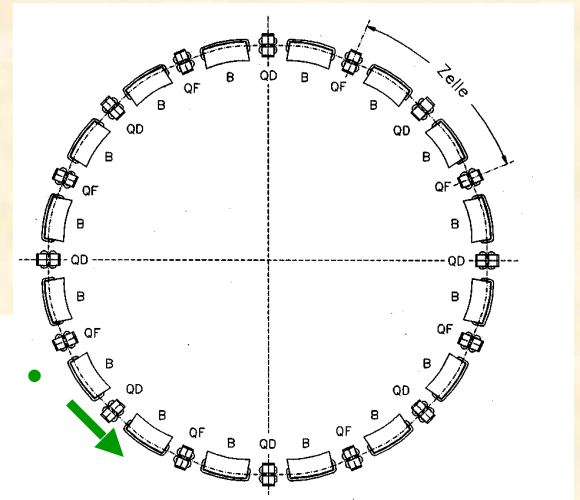
plot x , x' as a function of „ s “



$$\begin{pmatrix} x \\ x' \end{pmatrix}_{s1} = M_{\text{turn}} * \begin{pmatrix} x \\ x' \end{pmatrix}_{s0}$$

... and now the ellipse:

note for each turn x, x' at a given position „ s_1 “ and plot in the phase space diagram

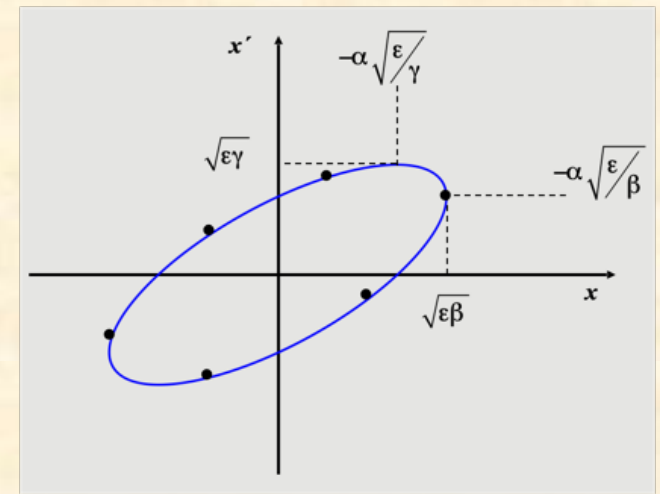


Particle Tracking in a Storage Ring
Calculate x, x' for each accelerator element according to matrix formalism and plot x, x' at a given position „ s “ turn by turn in the phase space diagram

Phase Space Ellipse

particle trajectory: $x(s) = \sqrt{\varepsilon} \sqrt{\beta(s)} \cos \{ \psi(s) + \phi \}$

max. Amplitude: $\hat{x}(s) = \sqrt{\varepsilon\beta}$ \longrightarrow x' at that position ...?



... put $\hat{x}(s)$ into $\varepsilon = \gamma(s) x^2(s) + 2\alpha(s)x(s)x'(s) + \beta(s) x'^2(s)$ and solve for x'

$$\varepsilon = \gamma \cdot \varepsilon\beta + 2\alpha \sqrt{\varepsilon\beta} \cdot x' + \beta x'^2$$

\longrightarrow $x' = -\alpha \cdot \sqrt{\varepsilon / \beta}$

* A high β -function means a large beam size and a small beam divergence. !
 ... et vice versa !!!

* In the middle of a quadrupole $\beta = \text{maximum}$,
 $\alpha = \text{zero}$ } $x' = 0$

... and the ellipse is flat

Phase Space Ellipse

$$\varepsilon = \gamma(s) x^2(s) + 2\alpha(s)x(s)x'(s) + \beta(s) x'^2(s)$$

$$\alpha(s) = -\frac{1}{2}\beta'(s)$$

$$\gamma(s) = \frac{1+\alpha(s)^2}{\beta(s)}$$

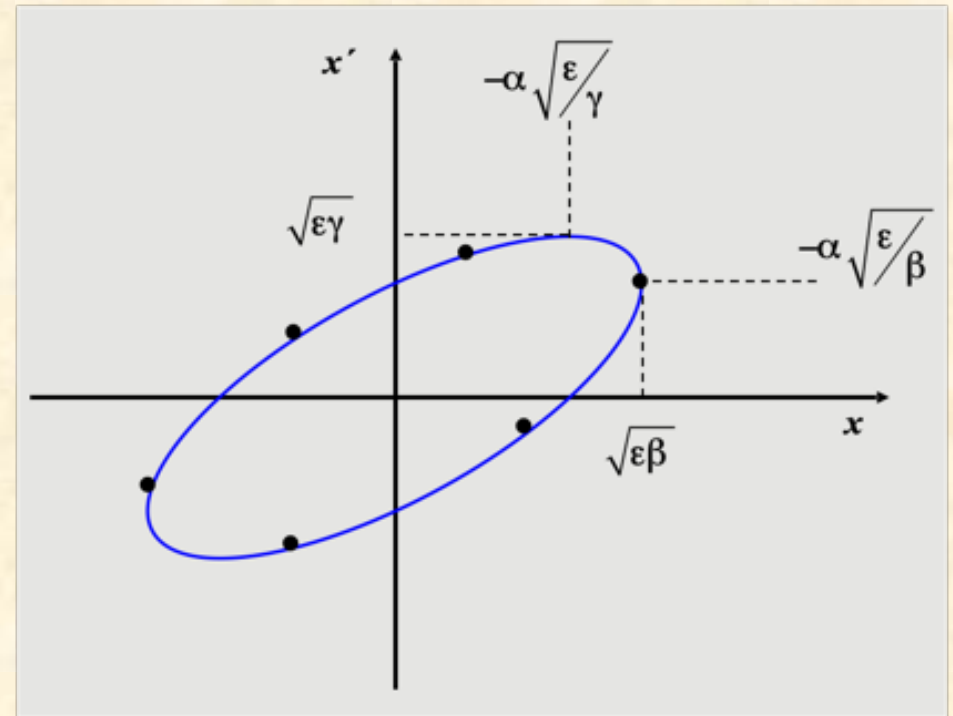
$$\longrightarrow \varepsilon = \frac{x^2}{\beta} + \frac{\alpha^2 x^2}{\beta} + 2\alpha \cdot xx' + \beta \cdot x'^2$$

... solve for x' $x'_{1,2} = \frac{-\alpha \cdot x \pm \sqrt{\varepsilon\beta - x^2}}{\beta}$

... and determine \hat{x}' via: $\frac{dx'}{dx} = 0$

$$\longrightarrow \hat{x}' = \sqrt{\varepsilon\gamma}$$

$$\longrightarrow \hat{x} = \pm\alpha \sqrt{\frac{\varepsilon}{\gamma}}$$



shape and orientation of the phase space ellipse depend on the Twiss parameters $\beta \alpha \gamma$

Liouville states that phase space is conserved.

*Primarily, this refers to 6-dimensional phase space
($x-x'$, $y-y'$ and $s-dp/p$).*

*When the component phase spaces are uncoupled,
the phase space is conserved within the 2- dimensional
and/or 4-dimensional spaces.*

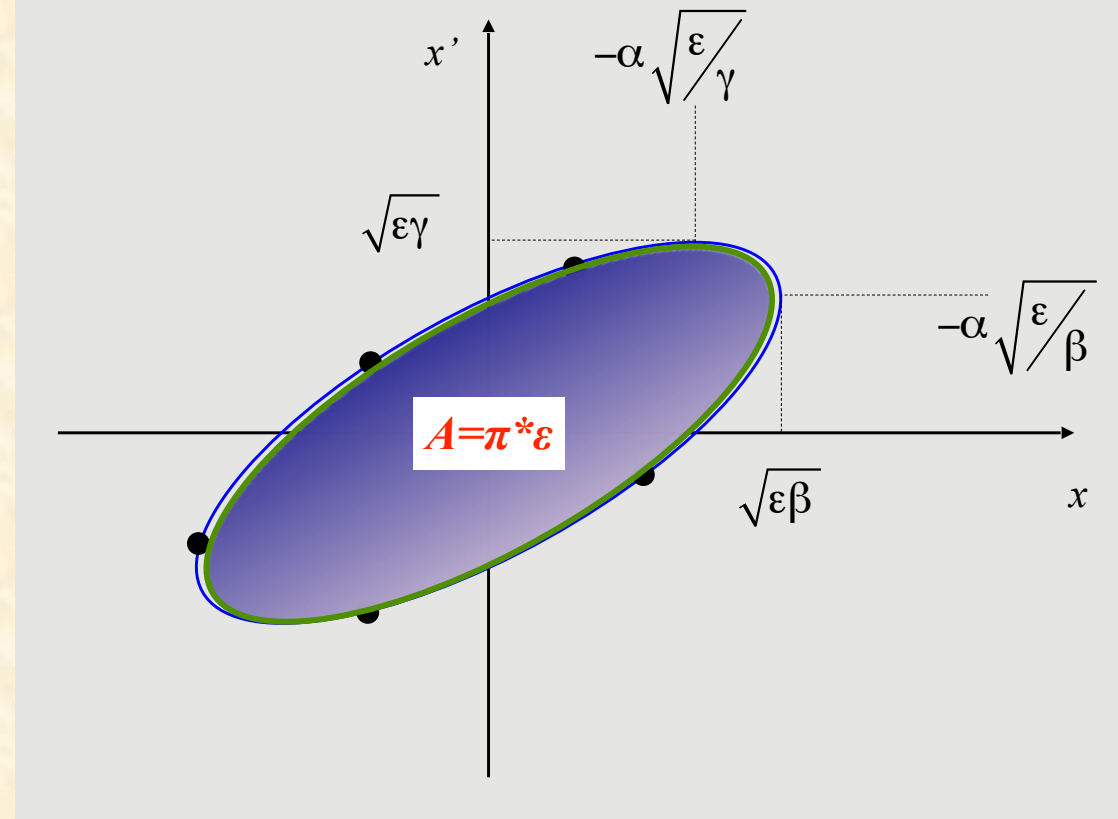
*The invariant of the motion in the uncoupled $x-x'$ or $y-y'$ spaces
is another way of saying the phase space is conserved.*

*Phase space is not conserved if ions change, e.g. by stripping
or nuclear fragmentation, or if non-Hamiltonian forces appear
e.g. scattering or photon emission.*

(Phil Bryant)

Phase Space Area & Emittance

shape and orientation of the phase space ellipse depend on the Optics parameters α , β , γ



The emittance of a beam is related to the phase-space area that it occupies and is therefore related to the motion invariants of the constituent ions.

A practical definition of emittance requires a choice for the limiting ellipse that defines the phase-space area of the beam.

Usually this is related to some number of standard deviations of the beam distribution, for example “the 1-sigma emittance is ...”.

(Phil Bryant)

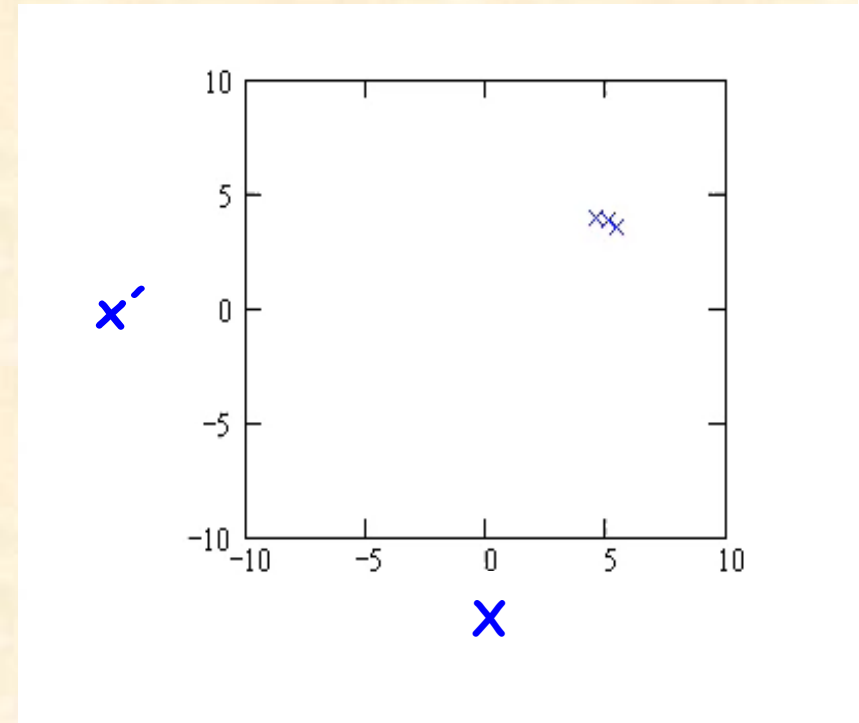
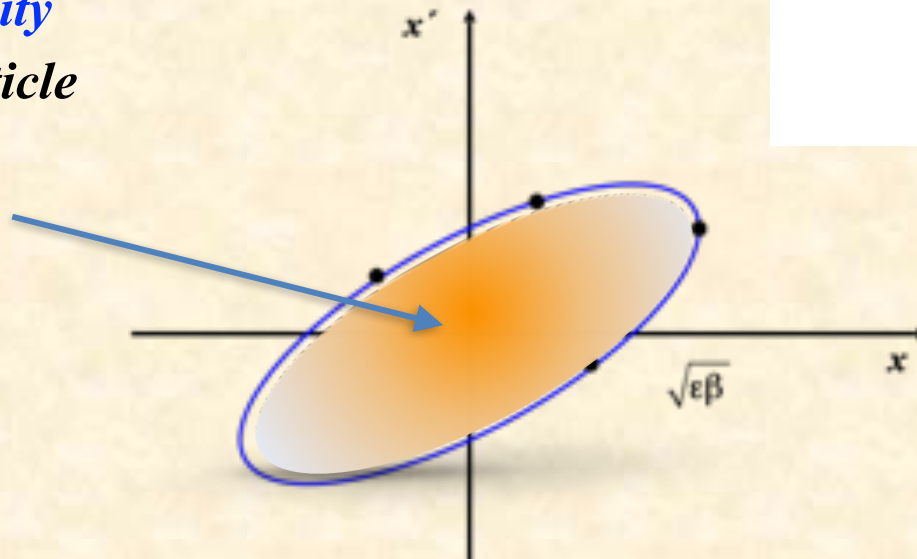
14.) Theorem of Liouville

... and now the ellipse:

note for each turn x, x' at a given position „s“ and plot in the phase space diagram

under the influence of conservative forces, the particle kinematics will always follow an ellipse in phase space x, x' phase space volume = constant

We use the area of that beam-ellipse as quality attribute for the particle ensemble: $A = \varepsilon \pi$



Time for a blue Slide ...

Why do we do that ?

—> *the beam size is given by two parameters:
β function - focusing properties
ε as intrinsic beam quality*

—> *beam size:* $\sigma = \sqrt{\varepsilon \cdot \beta}$

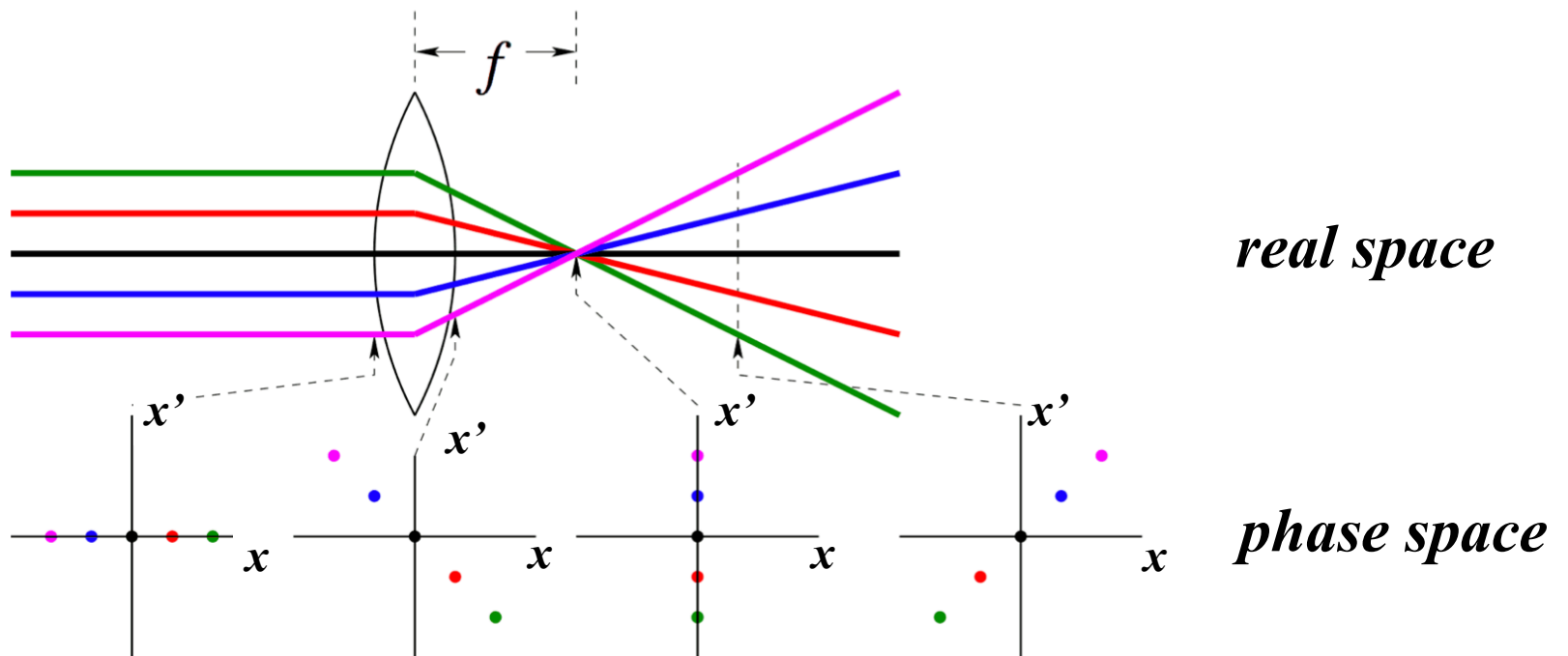
—> *the stability of the phase space ellipse, ε,
tells us about the stability
of the particle oscillation, which is ...
... “the lifetime” of the beam.*

—> *the size of the ellipse tells us about the particle density,
... which is the beam quality in collision.*

Phase Space & Real Space

... don't worry: it takes some time to fully find your way in both worlds.

Particle trajectories:

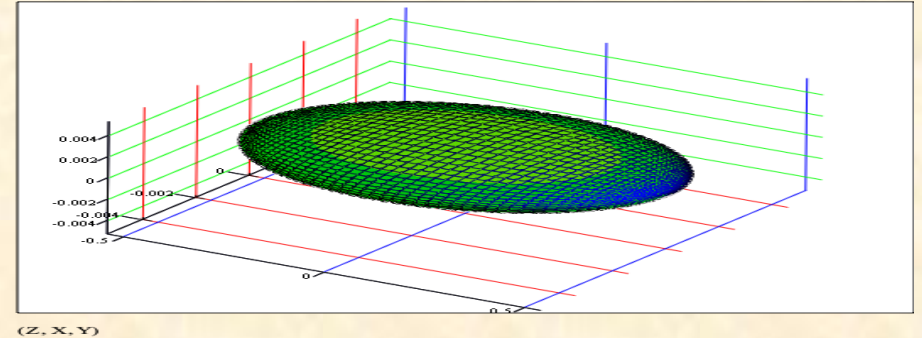
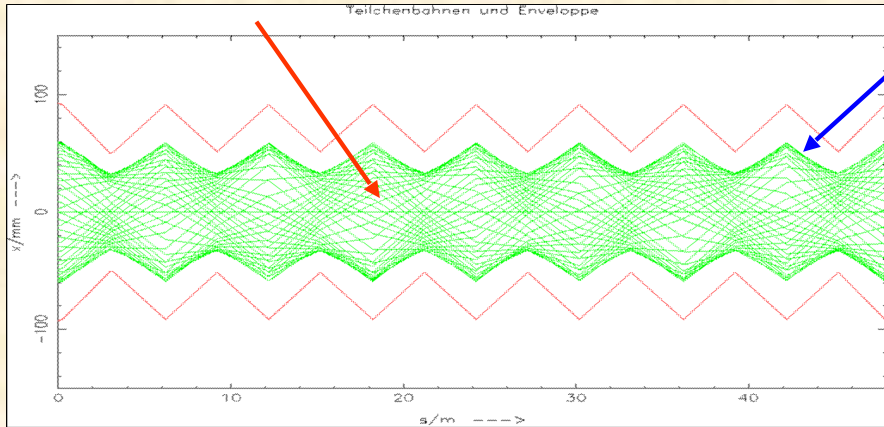


A real Beam:

particle bunch

$$x(s) = \sqrt{\epsilon} \sqrt{\beta(s)} \cdot \cos(\Psi(s) + \phi)$$

$$\hat{x}(s) = \sqrt{\epsilon} \sqrt{\beta(s)}$$



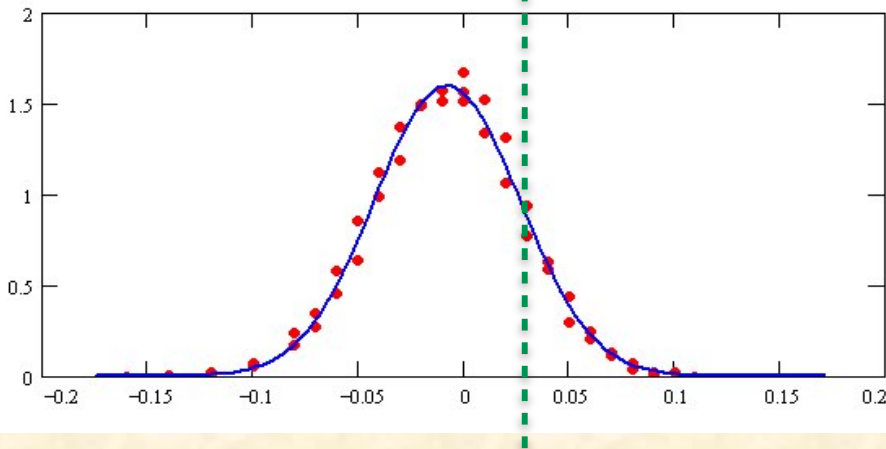
single particle trajectories, $N \approx 10^{11}$ per bunch

Gauß Particle Distribution:
$$\rho(\mathbf{x}) = \frac{N \cdot e}{\sqrt{2\pi}\sigma_x} \cdot e^{-\frac{1}{2}\frac{x^2}{\sigma_x^2}}$$

particle at distance 1σ from centre \leftrightarrow 68.3 % of all beam particles

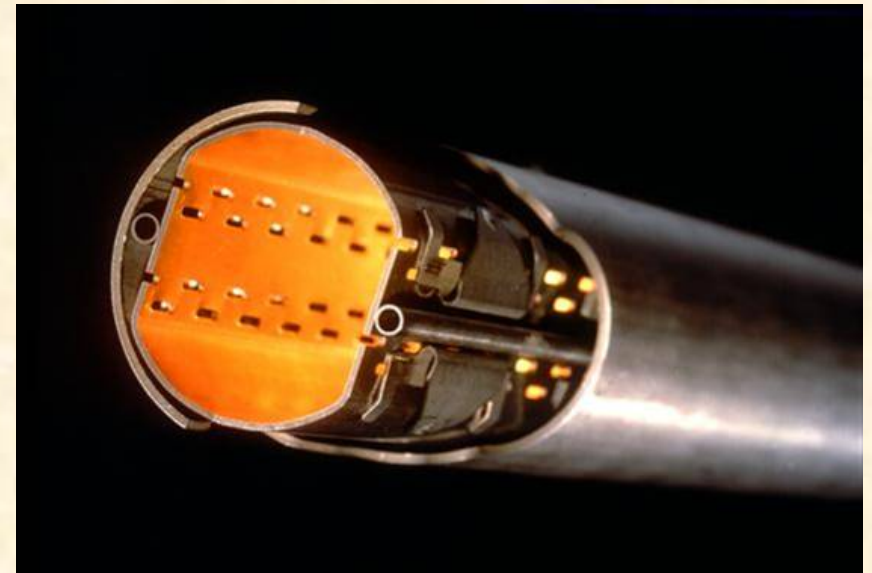
vertical:

$$\sigma_{\text{fit}} = 24.376 \cdot \mu\text{m}$$



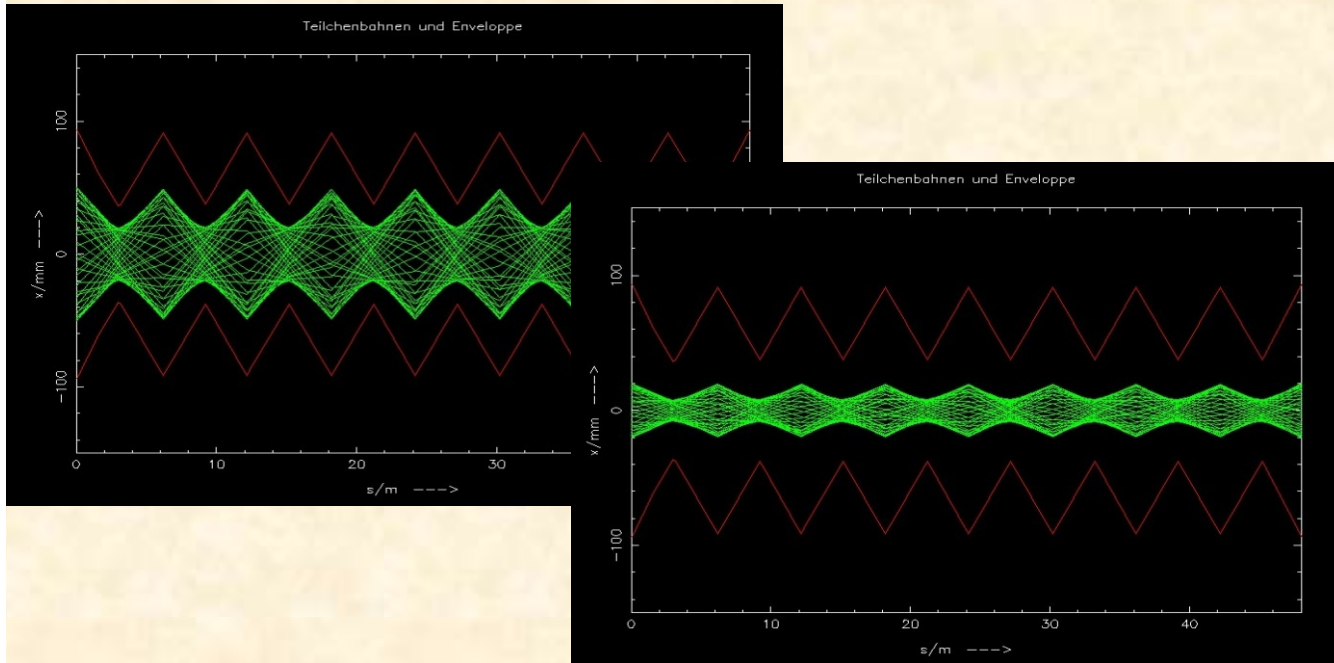
LHC:

$$\sigma = \sqrt{\epsilon \cdot \beta} = \sqrt{5 \cdot 10^{-10} \text{ m rad} \cdot 180 \text{ m}} = 0.3 \text{ mm}$$



aperture requirements: $r_0 \geq 10 \cdot \sigma$

Emittance and Beam Size:



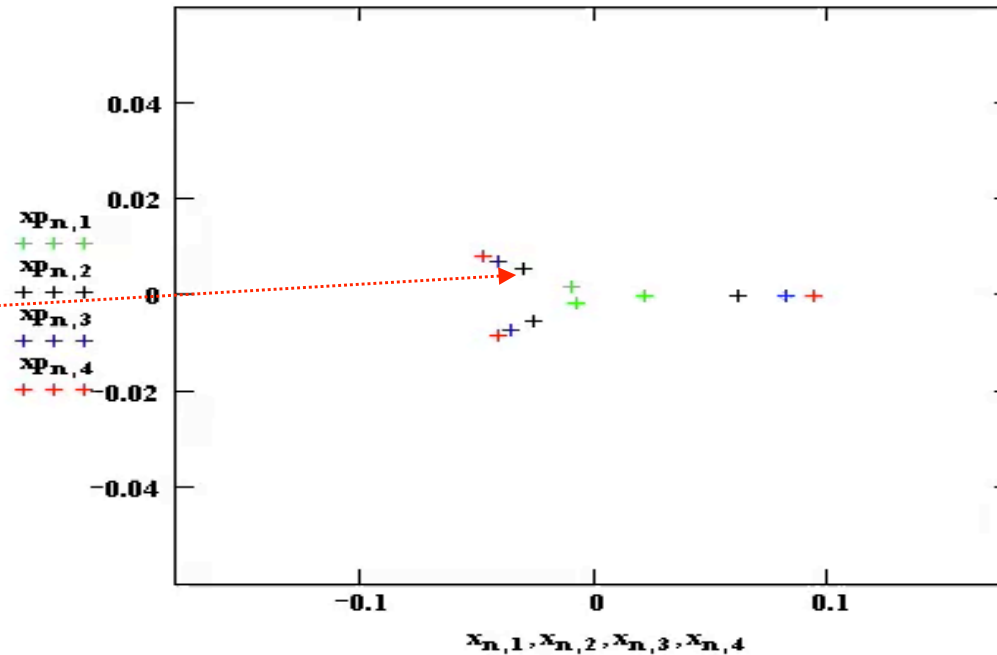
Example: LHC

beam parameters in the arc

$$\beta(x) \approx 180 \text{ m}$$

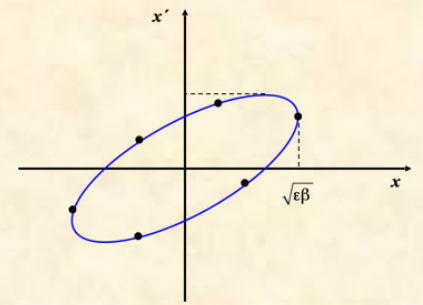
$$\varepsilon \approx 5 \cdot 10^{-10} \text{ rad} \cdot \text{m} \quad (\Leftrightarrow 1\sigma)$$

$$\sigma = \sqrt{\varepsilon\beta} \approx 0.3 \text{ mm}$$

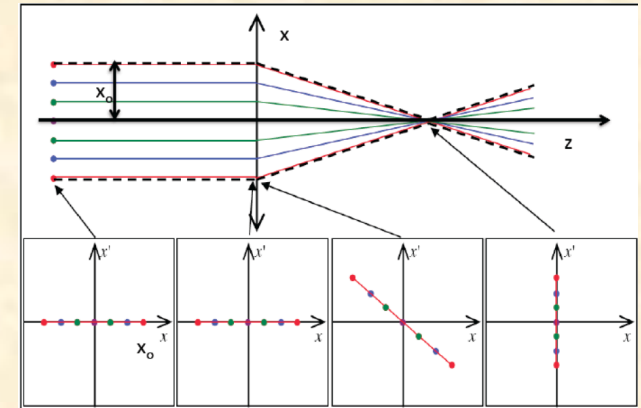


15.) Statistical Definition of Emittance:

The emittance is *the* quality parameter of the particle distribution



*the ideal case ... that never really exists ...
laminar (“LASER like) beam*



the real case ... the non-laminar (“real”) beam

Maxwell distribution:

source temperature “T”

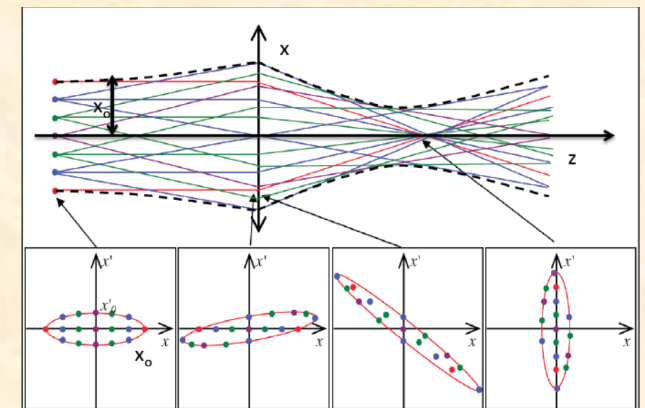
kinetic energy per degree of freedom:

$$E_{kin} = \frac{1}{2}kT$$

transverse momentum of the particles:

$$\frac{1}{2}mv_x^2 = \frac{p_x^2}{2m} = \frac{1}{2}kT \quad \longrightarrow \quad \sqrt{\langle p_x^2 \rangle} = \sqrt{mkT}$$

the particles have an intrinsic (transverse) momentum distribution



Statistical Definition of Emittance:

The r.m.s. emittance is a statistical definition of the amount of phase space covered by a beam. If the beam is centred, (symmetric situation) ($\langle x \rangle = \langle x' \rangle = 0$) we can write:

$$\epsilon_{rms} = \frac{1}{N} \sqrt{\Sigma x^2 \Sigma x'^2 - (\Sigma x x')^2}$$

If we really refer to the actual particle distribution our emittance definition is much more precise.

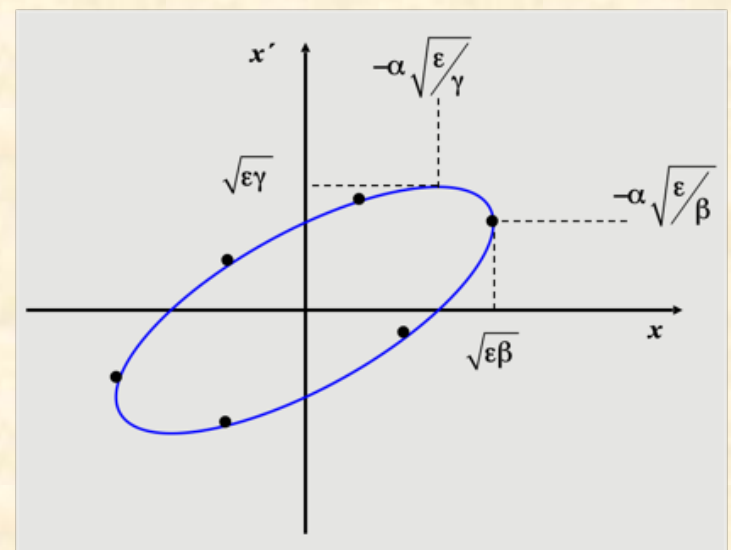
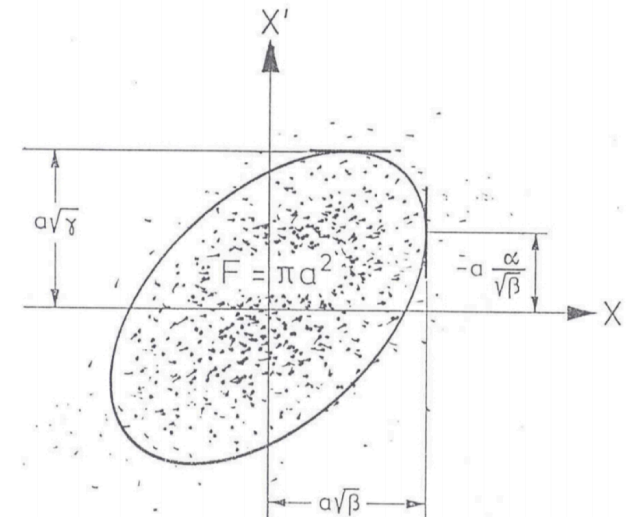
We can translate into our Twiss language via:

$$\gamma_x \cdot \epsilon_{rms} = \langle x'^2 \rangle = \sigma_{x'}^2$$

$$\beta_x \cdot \epsilon_{rms} = \langle x^2 \rangle = \sigma_x^2$$

$$\alpha_x \cdot \epsilon_{rms} = \langle x x' \rangle$$

The beam is composed of particles distributed in phase space.



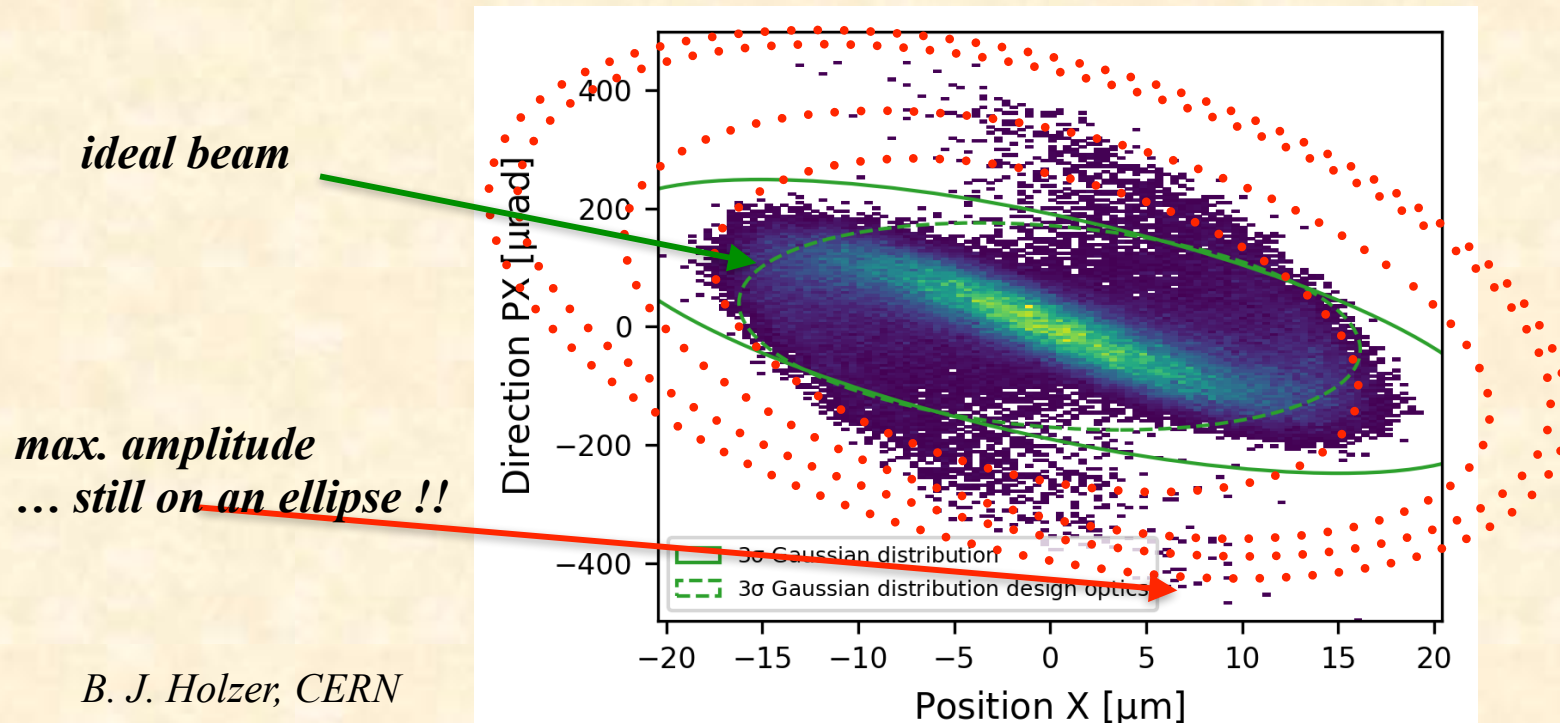
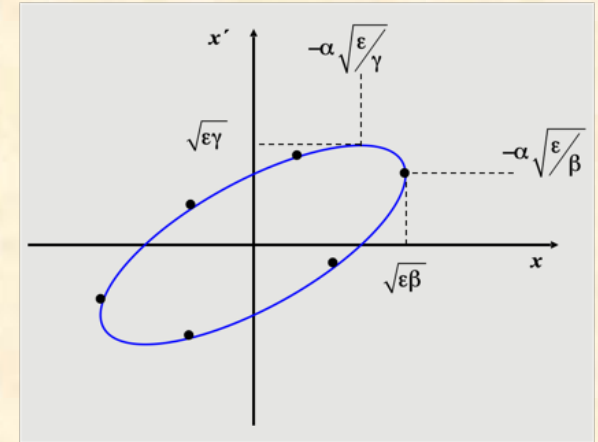
Emittance Dilution:

As soon as we inject the beam into an accelerator lattice, it is *the actual Twiss parameters*, that define the phase space ellipse in its shape and orientation.

We should *optimise the α , β , γ to fit as much as possible to the actual distribution.*

And we should *keep ϵ as small as possible.*

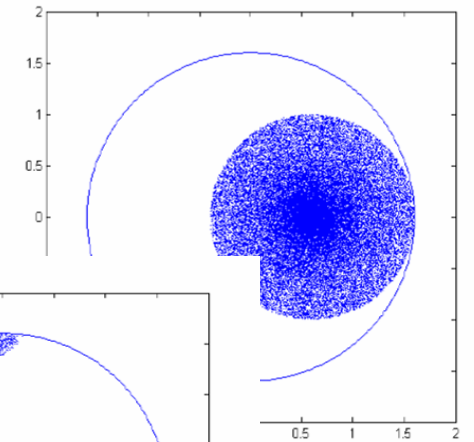
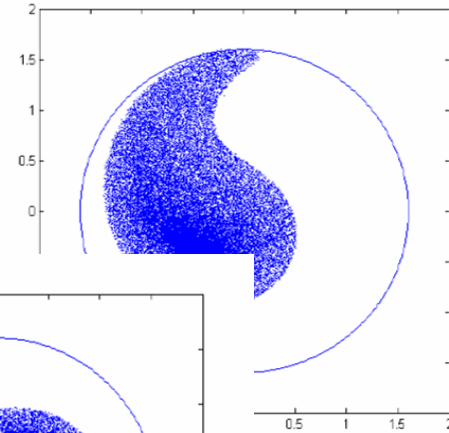
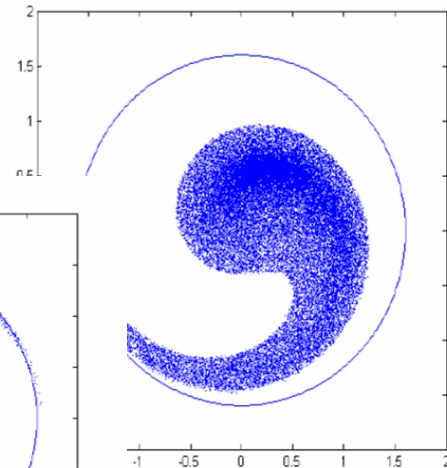
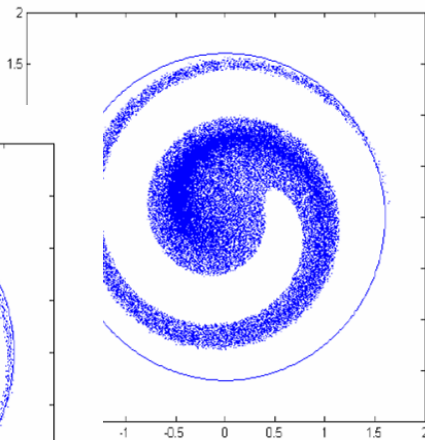
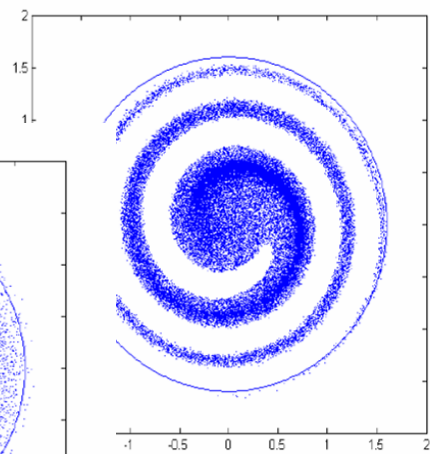
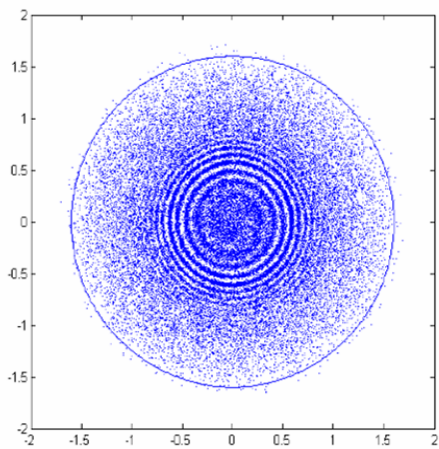
In the synchrotron *each single particle will follow its phase space ellipse*, that is defined by the ring optics.



Filamentation

Non-linear effects (e.g. magnetic field multipoles) distort the harmonic oscillation and lead to amplitude dependent effects in the particle motion in phase space.

Over many turns, a non-ideal phase-space distribution is smeared out and transformed into an emittance increase.



16.) Transfer Matrix M ... *yes we had the topic already*

*general solution
of Hill's equation*

$$\left\{ \begin{array}{l} x(s) = \sqrt{\varepsilon} \sqrt{\beta(s)} \cos \{ \psi(s) + \phi \} \\ x'(s) = \frac{-\sqrt{\varepsilon}}{\sqrt{\beta(s)}} \left[\alpha(s) \cos \{ \psi(s) + \phi \} + \sin \{ \psi(s) + \phi \} \right] \end{array} \right.$$

remember the trigonometrical gymnastics: $\sin(a + b) = \dots$ etc

$$x(s) = \sqrt{\varepsilon} \sqrt{\beta_s} (\cos \psi_s \cos \phi - \sin \psi_s \sin \phi)$$

$$x'(s) = \frac{-\sqrt{\varepsilon}}{\sqrt{\beta_s}} \left[\alpha_s \cos \psi_s \cos \phi - \alpha_s \sin \psi_s \sin \phi + \sin \psi_s \cos \phi + \cos \psi_s \sin \phi \right]$$

starting at point $s(0) = s_0$, where we put $\Psi(0) = 0$

$$\cos \phi = \frac{x_0}{\sqrt{\varepsilon \beta_0}},$$

$$\sin \phi = -\frac{1}{\sqrt{\varepsilon}} \left(x'_0 \sqrt{\beta_0} + \frac{\alpha_0 x_0}{\sqrt{\beta_0}} \right)$$

inserting above ...

$$\underline{x(s)} = \sqrt{\frac{\beta_s}{\beta_0}} \{ \cos\psi_s + \alpha_0 \sin\psi_s \} \underline{x_0} + \{ \sqrt{\beta_s \beta_0} \sin\psi_s \} \underline{x'_0}$$

$$\underline{x'(s)} = \frac{1}{\sqrt{\beta_s \beta_0}} \{ (\alpha_0 - \alpha_s) \cos\psi_s - (1 + \alpha_0 \alpha_s) \sin\psi_s \} \underline{x_0} + \sqrt{\frac{\beta_0}{\beta_s}} \{ \cos\psi_s - \alpha_s \sin\psi_s \} \underline{x'_0}$$

which can be expressed ... for convenience ... *in matrix form* $\begin{pmatrix} x \\ x' \end{pmatrix}_s = M \begin{pmatrix} x \\ x' \end{pmatrix}_0$

$$M = \begin{pmatrix} \sqrt{\frac{\beta_s}{\beta_0}} (\cos\psi_s + \alpha_0 \sin\psi_s) & \sqrt{\beta_s \beta_0} \sin\psi_s \\ \frac{(\alpha_0 - \alpha_s) \cos\psi_s - (1 + \alpha_0 \alpha_s) \sin\psi_s}{\sqrt{\beta_s \beta_0}} & \sqrt{\frac{\beta_0}{\beta_s}} (\cos\psi_s - \alpha_s \sin\psi_s) \end{pmatrix}$$

* we can calculate *the single particle trajectories* between two locations in the ring, if we know the $\alpha \beta \gamma$ at these positions.

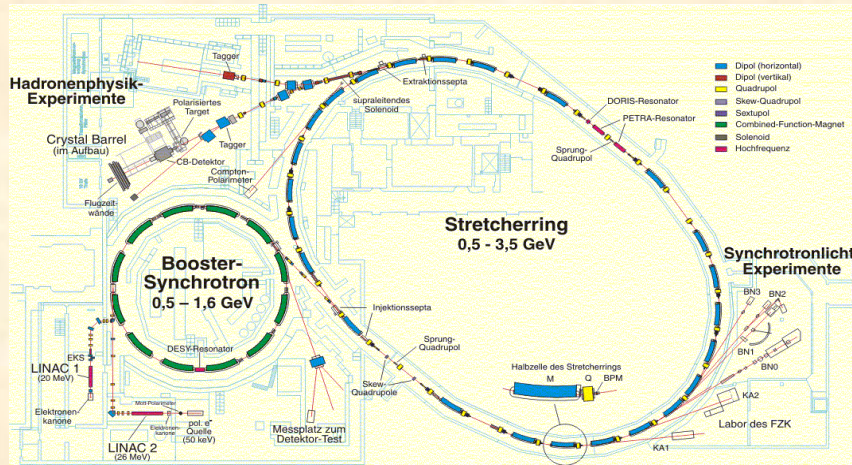
* and nothing but the $\alpha \beta \gamma$ at these positions.

* ... !

17.) Periodic Lattices

transfer matrix for particle trajectories as a function of the lattice parameters

$$M = \begin{pmatrix} \sqrt{\frac{\beta_s}{\beta_0}} (\cos\psi_s + \alpha_0 \sin\psi_s) & \sqrt{\beta_s \beta_0} \sin\psi_s \\ \frac{(\alpha_0 - \alpha_s) \cos\psi_s - (1 + \alpha_0 \alpha_s) \sin\psi_s}{\sqrt{\beta_s \beta_0}} & \sqrt{\frac{\beta_0}{\beta_s}} (\cos\psi_s - \alpha_s \sin\psi_s) \end{pmatrix}$$



ELSA Electron Storage Ring

„This rather formidable looking matrix simplifies considerably if we consider one complete turn ...“

One Turn Matrix

$$M(s) = \begin{pmatrix} \cos\psi_{turn} + \alpha_s \sin\psi_{turn} & \beta_s \sin\psi_{turn} \\ -\gamma_s \sin\psi_{turn} & \cos\psi_{turn} - \alpha_s \sin\psi_{turn} \end{pmatrix}$$

$$\psi_{turn} = \int_s^{s+L} \frac{ds}{\beta(s)}$$

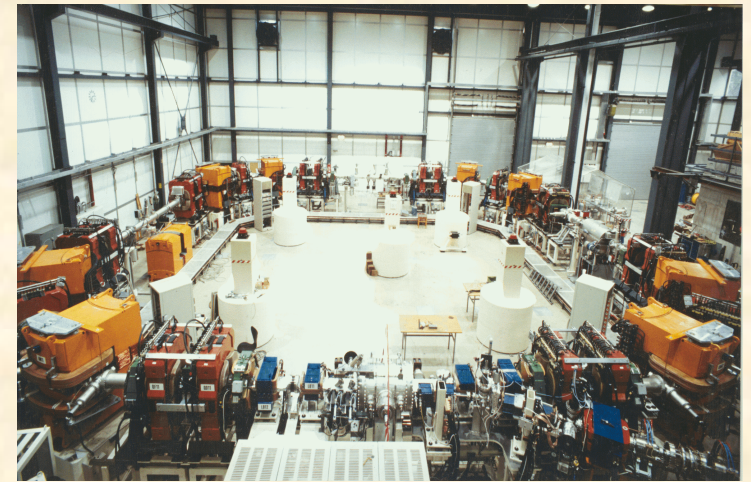
ψ_{turn} = *phase advance per period*

Tune: Phase advance per turn in units of 2π

$$Q = \frac{1}{2\pi} * \oint \frac{ds}{\beta(s)}$$

Stability Criterion:

Question: *what will happen, if we do not make too many mistakes and your **particle performs one complete turn** ?*



Matrix for 1 turn:

$$M = \begin{pmatrix} \cos\psi_{turn} + \alpha_s \sin\psi_{turn} & \beta_s \sin\psi_{turn} \\ -\gamma_s \sin\psi_{turn} & \cos\psi_{turn} - \alpha_s \sin\psi_{turn} \end{pmatrix} = \underbrace{\cos\psi}_{\mathbf{I}} \cdot \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} + \underbrace{\sin\psi}_{\mathbf{J}} \begin{pmatrix} \alpha & \beta \\ -\gamma & -\alpha \end{pmatrix}$$

Matrix for N turns:

$$M^N = (\mathbf{I} \cos\psi + \mathbf{J} \sin\psi)^N = \mathbf{I} \cos N\psi + \mathbf{J} \sin N\psi$$

The motion for N turns remains bounded, if the elements of M^N remain bounded

$$\psi = \text{real} \quad \Leftrightarrow \quad |\cos\psi| < 1 \quad \Leftrightarrow \quad |\text{Trace}(M)| < 2$$

stability criterion proof for the disbelieving colleagues !!

$$\text{Matrix for 1 turn: } M = \begin{pmatrix} \cos\psi_{turn} + \alpha_s \sin\psi_{turn} & \beta_s \sin\psi_{turn} \\ -\gamma_s \sin\psi_{turn} & \cos\psi_{turn} - \alpha_s \sin\psi_{turn} \end{pmatrix} = \underbrace{\cos\psi}_{\mathbf{I}} \cdot \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} + \underbrace{\sin\psi}_{\mathbf{J}} \begin{pmatrix} \alpha & \beta \\ -\gamma & -\alpha \end{pmatrix}$$

Matrix for 2 turns:

$$\begin{aligned} M^2 &= (\mathbf{I} * \cos\psi_1 + \mathbf{J} * \sin\psi_1) * (\mathbf{I} * \cos\psi_2 + \mathbf{J} * \sin\psi_2) \\ &= \mathbf{I}^2 * \cos\psi_1 \cos\psi_2 + \mathbf{I}\mathbf{J} * \cos\psi_1 \sin\psi_2 + \mathbf{J}\mathbf{I} * \sin\psi_1 \cos\psi_2 + \mathbf{J}^2 \sin\psi_1 \sin\psi_2 \end{aligned}$$

now ...

$$\mathbf{I}^2 = \mathbf{I}$$

$$\mathbf{I} * \mathbf{J} = \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} * \begin{pmatrix} \alpha & \beta \\ -\gamma & -\alpha \end{pmatrix} = \begin{pmatrix} \alpha & \beta \\ -\gamma & -\alpha \end{pmatrix}$$

$$\mathbf{J} * \mathbf{I} = \begin{pmatrix} \alpha & \beta \\ -\gamma & -\alpha \end{pmatrix} * \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} = \begin{pmatrix} \alpha & \beta \\ -\gamma & -\alpha \end{pmatrix}$$

$$\left. \begin{array}{l} \mathbf{I} * \mathbf{J} = \mathbf{J} * \mathbf{I} \end{array} \right\}$$

$$\mathbf{J}^2 = \begin{pmatrix} \alpha & \beta \\ -\gamma & -\alpha \end{pmatrix} * \begin{pmatrix} \alpha & \beta \\ -\gamma & -\alpha \end{pmatrix} = \begin{pmatrix} \alpha^2 - \gamma\beta & \alpha\beta - \beta\alpha \\ -\gamma\alpha + \alpha\gamma & \alpha^2 - \gamma\beta \end{pmatrix} = \begin{pmatrix} -1 & 0 \\ 0 & -1 \end{pmatrix} = -\mathbf{I}$$

$$M^2 = \mathbf{I} * \cos(\psi_1 + \psi_2) + \mathbf{J} * \sin(\psi_1 + \psi_2)$$

$$M^2 = \mathbf{I} * \cos(2\psi) + \mathbf{J} * \sin(2\psi)$$

18.) Transformation of α, β, γ

consider two positions in the storage ring: s_0, s

$$\begin{pmatrix} x \\ x' \end{pmatrix}_s = M \cdot \begin{pmatrix} x \\ x' \end{pmatrix}_{s_0} \quad \text{where ... } M = M_{QF} \cdot M_{QD} \cdot M_B \cdot M_{Drift} \cdot M_{QF} \cdot \dots$$

for a single element, e.g. ...

$$M_{\text{foc}} = \begin{pmatrix} \cos(\sqrt{|K|} \cdot l_q) & \frac{1}{\sqrt{|K|}} \sin(\sqrt{|K|} \cdot l_q) \\ -\sqrt{|K|} \cdot \sin(\sqrt{|K|} \cdot l_q) & \cos(\sqrt{|K|} \cdot l_q) \end{pmatrix} = \begin{pmatrix} C & S \\ C' & S' \end{pmatrix}$$

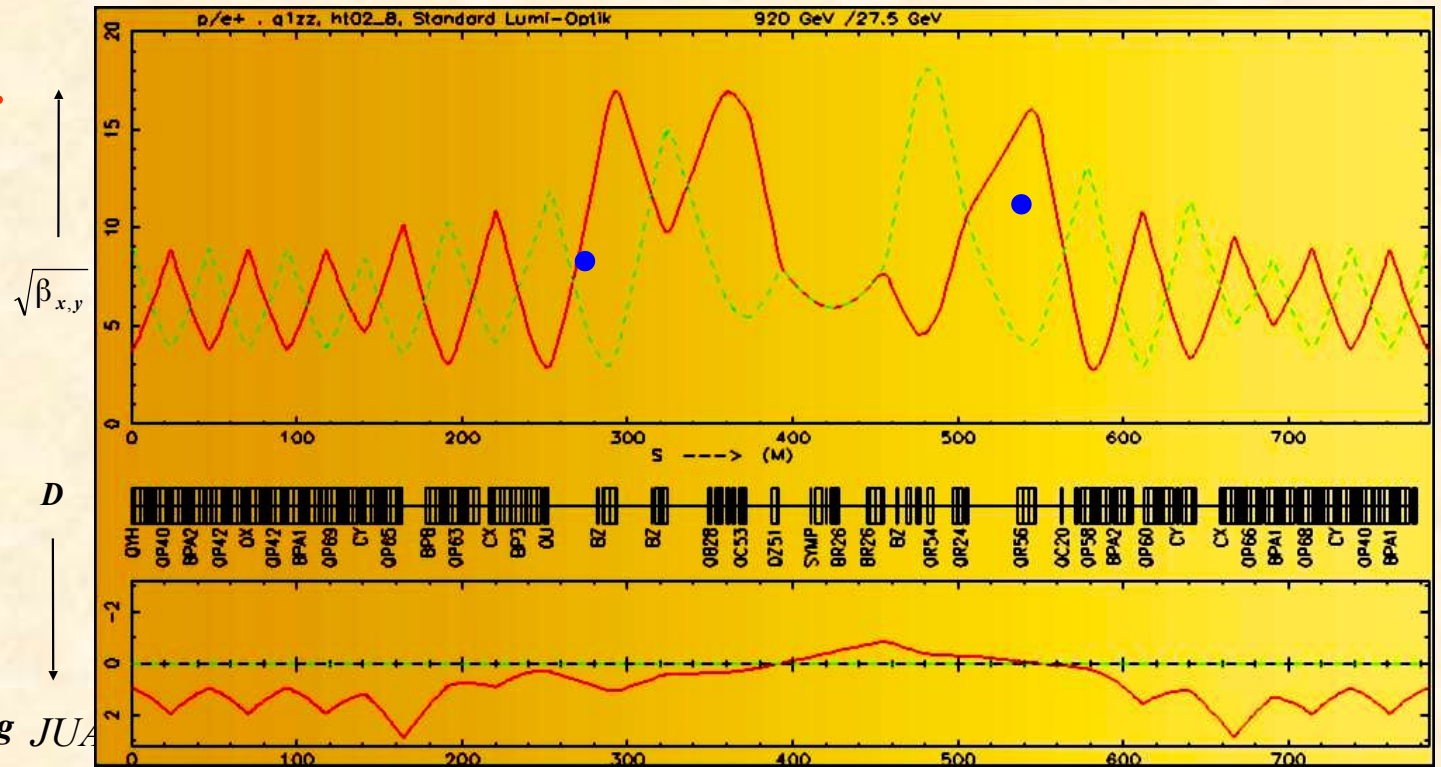
for a sequence of elements ...

$$M_{\text{seq}} = \begin{pmatrix} m_{11} & m_{12} \\ m_{21} & m_{22} \end{pmatrix} = \begin{pmatrix} C & S \\ C' & S' \end{pmatrix}$$

since $\epsilon = \text{const}$ (Liouville):

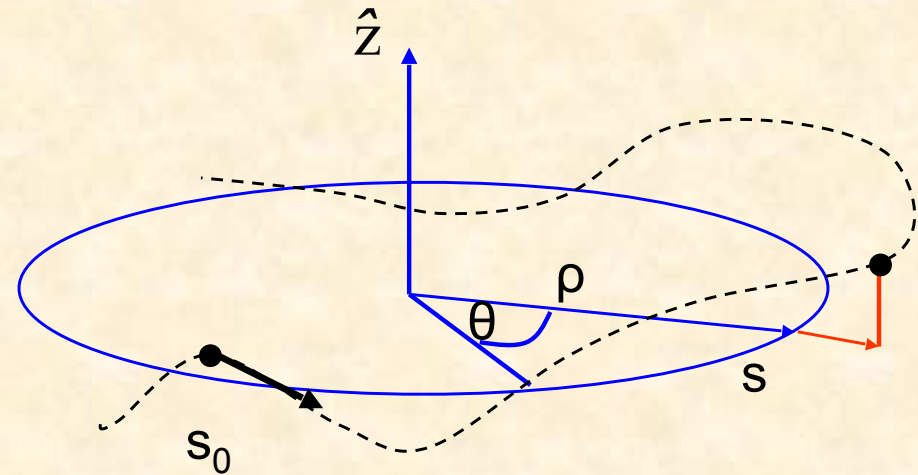
$$\epsilon = \beta x'^2 + 2\alpha x x' + \gamma x^2$$

$$\epsilon = \beta_0 x_0'^2 + 2\alpha_0 x_0 x_0' + \gamma_0 x_0^2$$



express x_0, x'_0 as a function of x, x' .

$$\begin{pmatrix} x \\ x' \end{pmatrix}_s = M \cdot \begin{pmatrix} x \\ x' \end{pmatrix}_{s_0}$$



... remember $W = CS' - SC' = 1$

$$\begin{pmatrix} x \\ x' \end{pmatrix}_0 = M^{-1} \cdot \begin{pmatrix} x \\ x' \end{pmatrix}_s$$

$$M^{-1} = \frac{1}{\det(M)} \begin{pmatrix} m_{22} & -m_{12} \\ -m_{21} & m_{11} \end{pmatrix} = \begin{pmatrix} S' & -S \\ -C' & C \end{pmatrix}$$

$\det(M) = 1$



$$\begin{aligned} x_0 &= S'x - Sx' \\ x'_0 &= -C'x + Cx' \end{aligned}$$

inserting into ε

$$\varepsilon = \beta x'^2 + 2\alpha xx' + \gamma x^2$$

$$\varepsilon = \beta_0 (Cx' - C'x)^2 + 2\alpha_0 (S'x - Sx')(Cx' - C'x) + \gamma_0 (S'x - Sx')^2$$

sort via x, x' and compare the coefficients to get

$$\beta(s) = C^2\beta_0 - 2SC\alpha_0 + S^2\gamma_0$$

$$\alpha(s) = -CC'\beta_0 + (SC' + S'C)\alpha_0 - SS'\gamma_0$$

$$\gamma(s) = C'^2\beta_0 - 2S'C'\alpha_0 + S'^2\gamma_0$$

in matrix notation:

$$\begin{pmatrix} \beta \\ \alpha \\ \gamma \end{pmatrix}_s = \begin{pmatrix} C^2 & -2SC & S^2 \\ -CC' & SC' + CS' & -SS' \\ C'^2 & -2S'C' & S'^2 \end{pmatrix} \cdot \begin{pmatrix} \beta_0 \\ \alpha_0 \\ \gamma_0 \end{pmatrix}$$



- 1.) *this expression is important*
- 2.) *given the Twiss parameters α, β, γ at any point in the lattice we can transform them and calculate their values at any other point in the ring.*
- 3.) *the transfer matrix is given by the focusing properties of the lattice elements, the elements of M are just those that we used to calculate single particle trajectories.*
- 4.) *go back to point 1.)*

Résumé:

equation of motion:
$$\mathbf{x}''(s) + \mathbf{K}(s) \mathbf{x}(s) = 0, \quad K = 1/\rho^2 - k$$

general solution of Hill's equation:
$$x(s) = \sqrt{\varepsilon} \sqrt{\beta(s)} \cdot \cos(\psi(s) + \phi)$$

phase advance & tune:
$$\psi_{12}(s) = \int_{s_1}^{s_2} \frac{1}{\beta(s)} ds, \quad Q(s) = \frac{1}{2\pi} \oint \frac{1}{\beta(s)} ds$$

emittance:
$$\varepsilon = \gamma(s) \mathbf{x}^2(s) + 2\alpha(s) \mathbf{x}(s) \mathbf{x}'(s) + \beta(s) \mathbf{x}'^2(s)$$

transfer matrix from $s_1 \longrightarrow s_2$:
$$M = \begin{pmatrix} \sqrt{\frac{\beta_s}{\beta_0}} (\cos \psi_s + \alpha_0 \sin \psi_s) & \sqrt{\beta_s \beta_0} \sin \psi_s \\ \frac{(\alpha_0 - \alpha_s) \cos \psi_s - (1 + \alpha_0 \alpha_s) \sin \psi_s}{\sqrt{\beta_s \beta_0}} & \sqrt{\frac{\beta_0}{\beta_s}} (\cos \psi_s - \alpha_s \sin \psi_s) \end{pmatrix}$$

matrix for 1 turn:

$$M(s) = \begin{pmatrix} \cos \psi_{turn} + \alpha_s \sin \psi_{turn} & \beta_s \sin \psi_{turn} \\ -\gamma_s \sin \psi_{turn} & \cos \psi_{turn} - \alpha_s \sin \psi_{turn} \end{pmatrix}$$

stability criterion:

$$|\text{Trace}(M)| < 2$$