



**Gábor Kasza**

10th Day of Femtoscopy

Gyöngyös, 29/10/2024

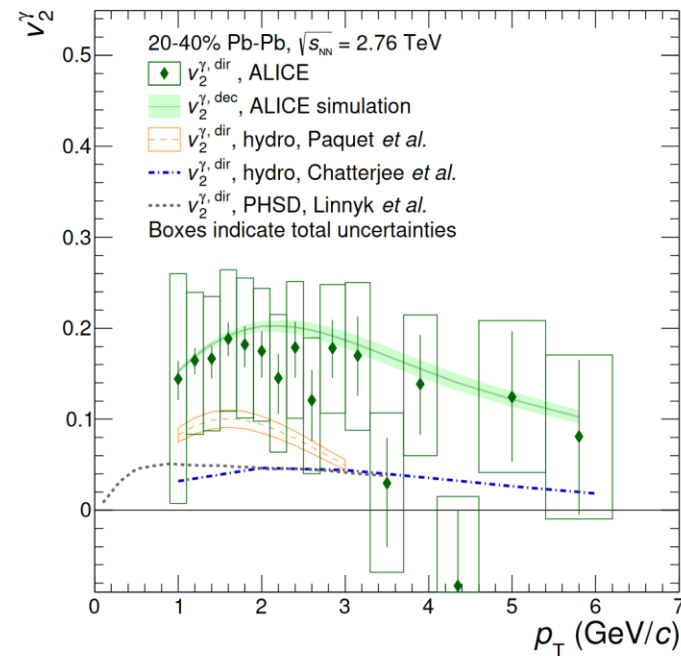
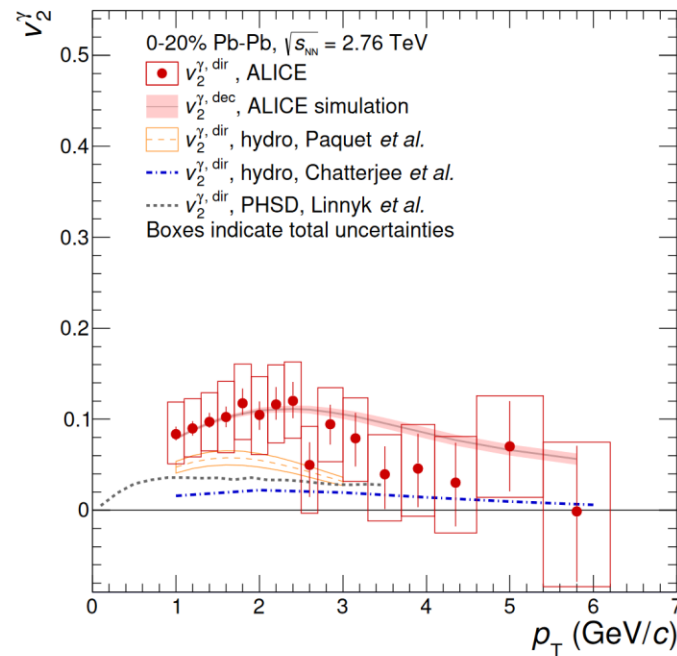
# Description of direct photon spectra by analytic hydrodynamics

**MATE**

 **WIGNER**

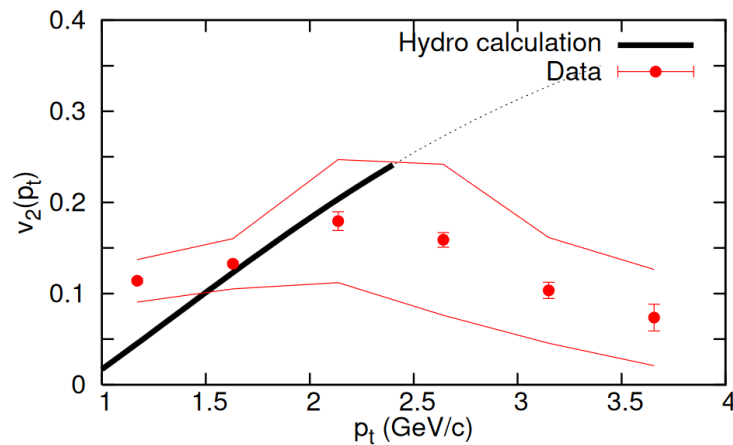
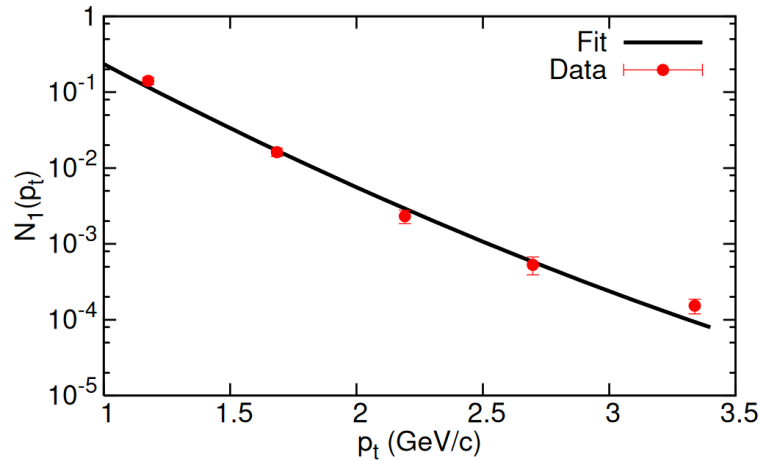
# Motivation

- *Direct photon puzzle*: the measured  $v_2$  of direct photons is of the same order of magnitude as for hadrons.
- $v_2$  cannot be described simultaneously with direct photon spectra using the theoretical models known so far.



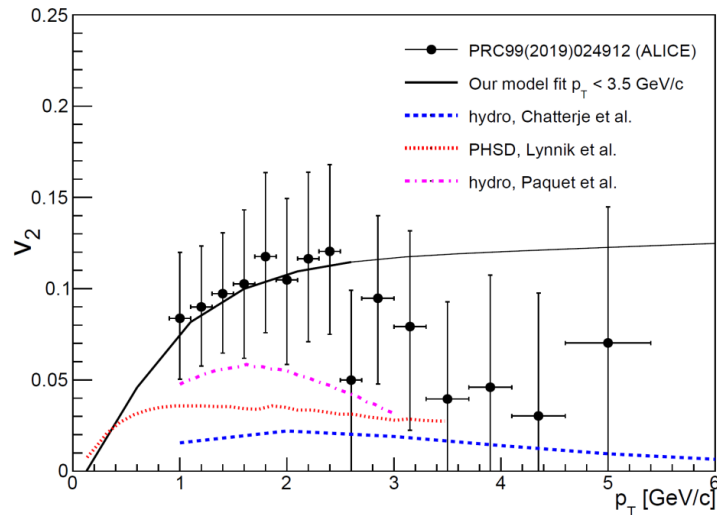
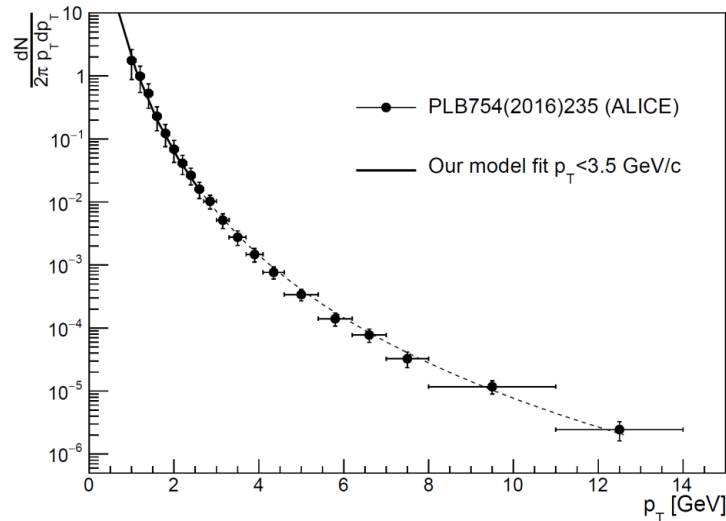
*Phys.Lett.B* 789 (2019)  
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## Earlier success of analytic hydro



- Based on the Csörgő-Csernai-Hama-Kodama solution of relativistic hydro.  
*Acta Phys.Hung.A 21 (2004) 73-84*
- No acceleration, but 1+3d.
- Gaussian temperature profile.
- Analytic calculation* of spectrum and  $v_2$  using second-order saddle-point approximation.
- Fitted to PHENIX Au+Au @ 200 GeV data.
- $T_0 > 507 \pm 12$  MeV

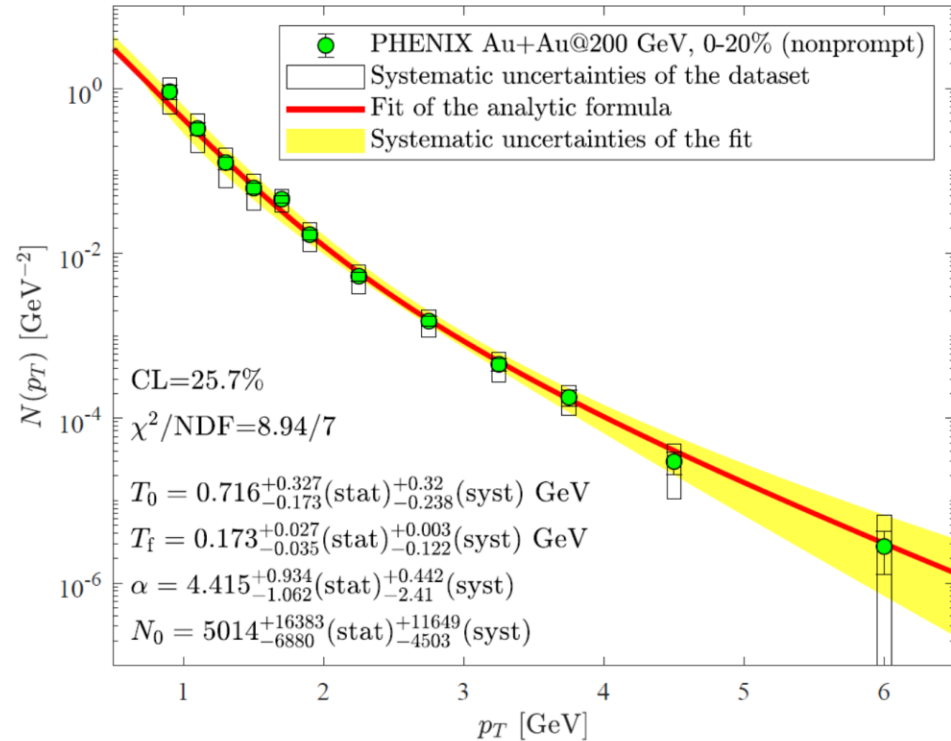
# Recent successes of analytic hydro 1



- Same model (based on CCHK solution, Gaussian temperature, 1+3d, no acceleration), but:
- Numeric calculation* of observables to avoid analytic approximations.
- Fitted to ALICE Pb+Pb @ 2.76 GeV data.
- v<sub>2</sub> and spectrum were fitted simultaneously.*
- $T_0 = 418 \pm 31$  MeV

S. Lökös and G. K.: submitted to EPJA

## Recent successes of analytic hydro 2



- *Scaling behaviour of data* has been found.

- Based on the relativistic hydrodynamic solution of Csörgő, Kasza, Csanád and Jiang.  
*T. Csörgő, G.K., M. Csanád, Z. Jiang: Universe 4 (2018) 6, 69*
- *Locally accelerating velocity field*, inhomogeneous temperature, but only 1+1d.
- *Analytic calculation* of spectrum using saddle-point approximation.
- Fitted to the non-prompt component of PHENIX Au+Au @ 200 GeV data.
- *Non-prompt component*: dominated by hydrodynamic evolution.

## ► New 1+1d model with generalized EoS

- Same 1+1d model (based on the CKCJ solution, accelerating velocity field, inhomogeneous temperature), but:
- Generalized for *a broadened class of EoS that contains lQCD EoS*.
- The spectrum is embedded to the 1+3d space, but  $v_2$  *cannot be calculated*.
- The spectrum has a
  - low temperature component ( $T < T_c$  let's call it hadronic component),
  - high temperature component ( $T > T_c$  let's call it QGP component).

## Equation of State

- EoS:  $\mu=0$ ,  $\varepsilon=\kappa(T)p$
- *Requiring temperature inhomogeneity*: strong constraint for  $\kappa(T)$ :

$$\frac{d}{dT} \left[ \frac{\kappa(T)T}{1 + \kappa(T)} \right] = \frac{c_Q}{1 + \kappa(T)} \quad T > T_c, \kappa(T)=\kappa_Q(T)$$

$$\frac{d}{dT} \left[ \frac{\kappa(T)T}{1 + \kappa(T)} \right] = \frac{c_H}{1 + \kappa(T)} \quad T < T_c, \kappa(T)=\kappa_H(T)$$

- Solutions for  $\kappa(T)$ :

$$\kappa_Q(T) = \frac{c_Q \left( \frac{T}{T_c} \right)^{1+c_Q} + \frac{\kappa_c - c_Q}{\kappa_c + 1}}{\left( \frac{T}{T_c} \right)^{1+c_Q} - \frac{\kappa_c - c_Q}{\kappa_c + 1}} \longrightarrow c_Q = \kappa(T \gg T_c)$$

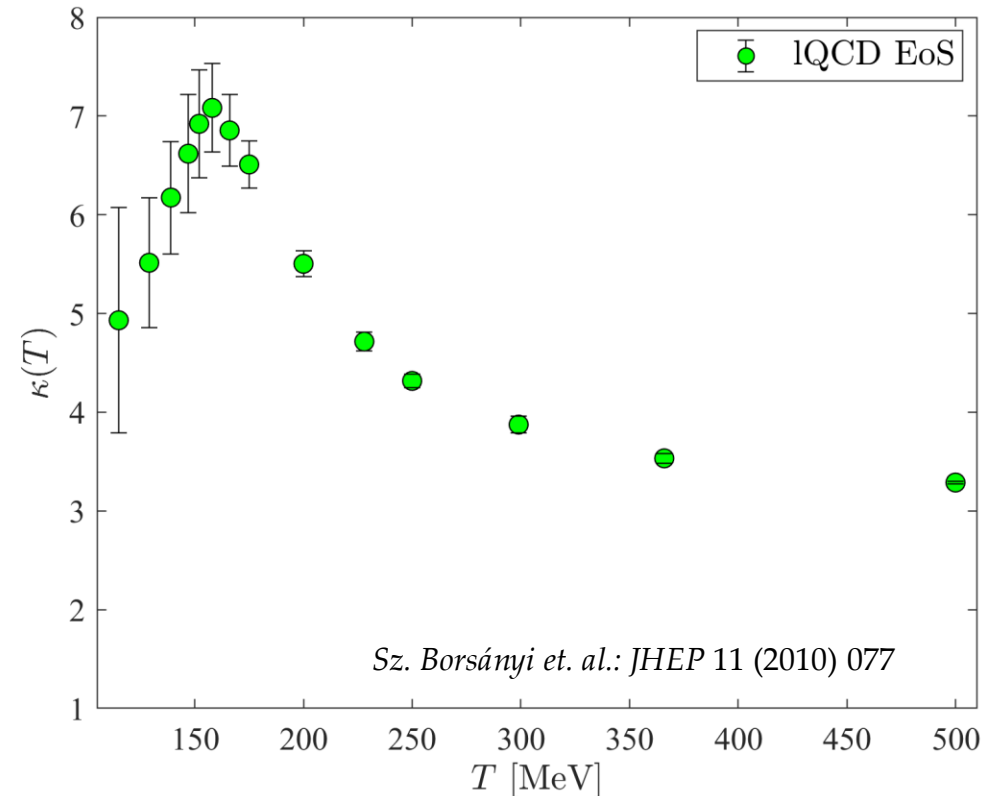
$$c_H \text{ controls the peak of } \kappa(T) \longleftarrow \kappa_H(T) = \frac{c_H \left( \frac{T}{T_f} \right)^{1+c_H} - \frac{c_H - \kappa_f}{\kappa_f + 1}}{\left( \frac{T}{T_f} \right)^{1+c_H} + \frac{c_H - \kappa_f}{\kappa_f + 1}}$$

## Parametrization of EoS

- $\kappa_Q(T)$  and  $\kappa_H(T)$  are matched at  $T_c$ :

$$\kappa(T) = \Theta(T - T_c)\kappa_H(T) + \Theta(T_c - T)\kappa_Q(T)$$

- Main goal: *mimic the lQCD EoS*
- $T_f = 124$  MeV: extracted from slopes of hadronic  $p_T$  spectra
- $\kappa_f = \kappa(T_f) \approx 5.5$
- $T_c \approx 160$  MeV
- $\kappa_c = \kappa(T_c) \approx 7$
- $\kappa(T \gg T_c) \approx 3.25$





# Temperature and Source Function

- Analytic solutions for the temperature:

$$T(\tau, \eta_z) = T_0 \left( \frac{\tau_0}{\tau} \right)^{\frac{\lambda}{c_H}} \left[ 1 + \frac{c_H - 1}{\lambda - 1} \sinh^2 (\Omega - \eta_z) \right]^{-\frac{\lambda}{2c_H}} \quad T < T_c$$

$$T(\tau, \eta_z) = T_0 \left( \frac{\tau_0}{\tau} \right)^{\frac{\lambda}{c_Q}} \left[ 1 + \frac{c_Q - 1}{\lambda - 1} \sinh^2 (\Omega - \eta_z) \right]^{-\frac{\lambda}{2c_Q}} \quad T > T_c$$

- Source function:

$$S(x^\mu, p^\mu) d^4x = \frac{g}{(2\pi\hbar)^3} \frac{H(\tau)}{\tau_R} \frac{p_\mu d\Sigma^\mu}{\exp\left(\frac{p^\mu u_\mu}{T}\right) - 1}$$

- $H(\tau)$ : opacity for photons  
(step-function, photons do not participate in strong interaction)

# ▶ Spectrum

- Two-component direct photon spectrum:

$$N(p_T) = \frac{dN}{2\pi p_T dp_T} = \underbrace{\int_{\tau_c}^{\tau_f} N(p_T, \tau, c_H) d\tau}_{\text{Low- } T \text{ (hadronic) component}} + \underbrace{\int_{\tau_0}^{\tau_c} N(p_T, \tau, c_Q) d\tau}_{\text{High- } T \text{ (QGP) component}}$$

*Low- T (hadronic) component    High- T (QGP) component*

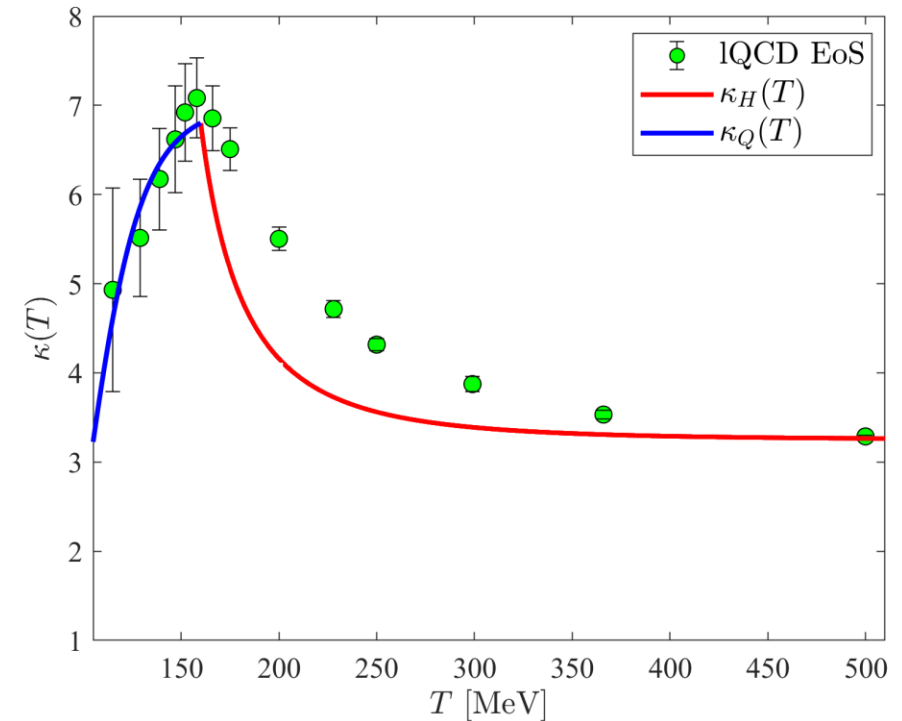
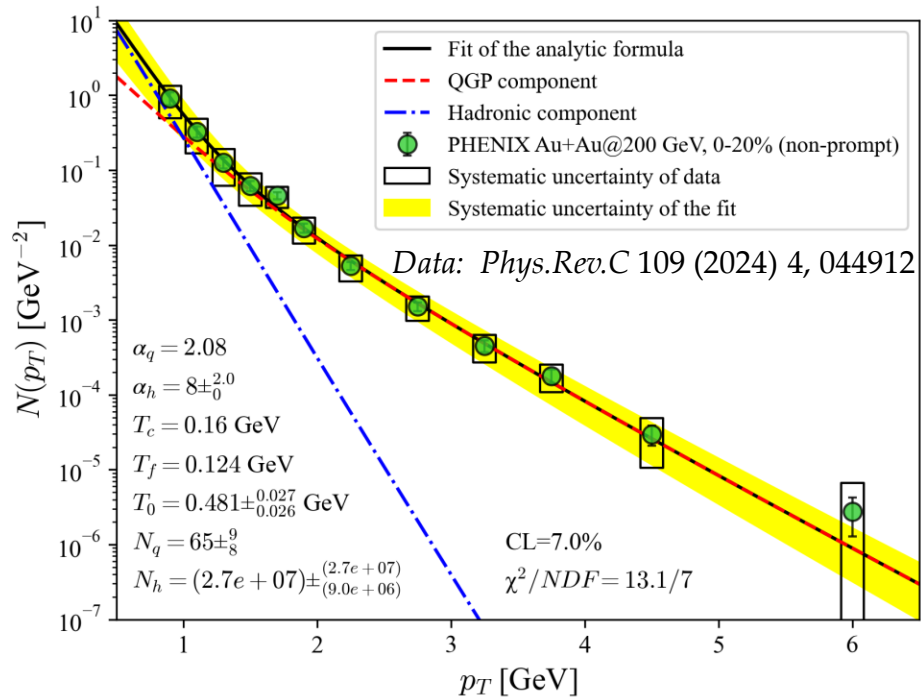
- Low-T component** (where  $\alpha_H = 2c_H/\lambda - 3$  and  $\lambda$  is fixed):

$$\left. \frac{d^2 N_H}{2\pi p_T dp_T dy} \right|_{y=0} = N_{0,H} \frac{2\alpha_H}{3\pi^{3/2}} \left[ \frac{1}{T_f^{\alpha_H}} - \frac{1}{T_c^{\alpha_H}} \right]^{-1} p_T^{-\alpha_H - 2} \Gamma \left( \alpha_H + \frac{5}{2}, \frac{p_T}{T} \right) \Big|_{T=T_f}^{T=T_c}$$

- High-T component** (where  $\alpha_Q = 2c_Q/\lambda - 3$  and  $\lambda$  fixed):

$$\left. \frac{d^2 N_Q}{2\pi p_T dp_T dy} \right|_{y=0} = N_{0,Q} \frac{2\alpha_Q}{3\pi^{3/2}} \left[ \frac{1}{T_c^{\alpha_Q}} - \frac{1}{T_0^{\alpha_Q}} \right]^{-1} p_T^{-\alpha_Q - 2} \Gamma \left( \alpha_Q + \frac{5}{2}, \frac{p_T}{T} \right) \Big|_{T=T_c}^{T=T_0}$$

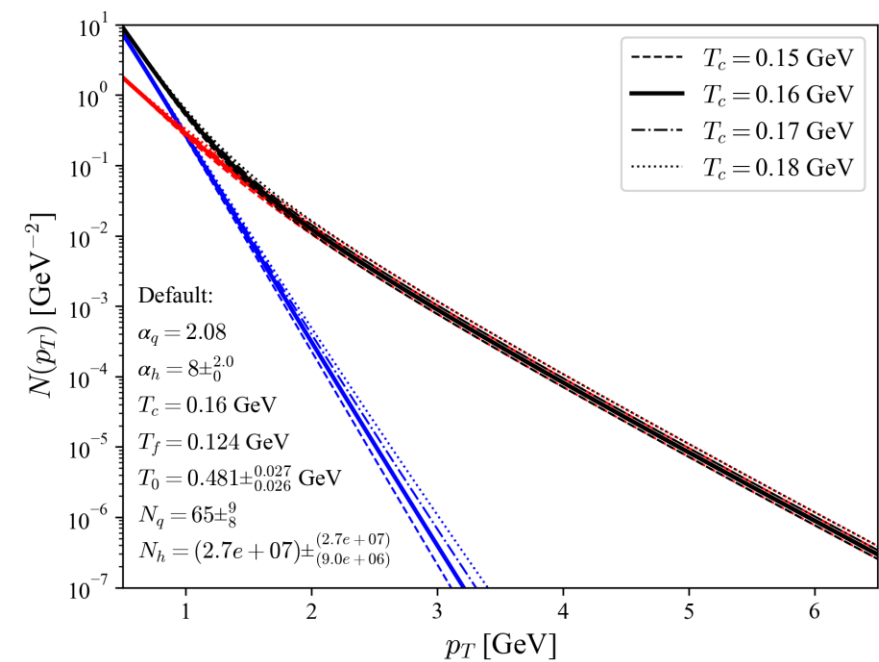
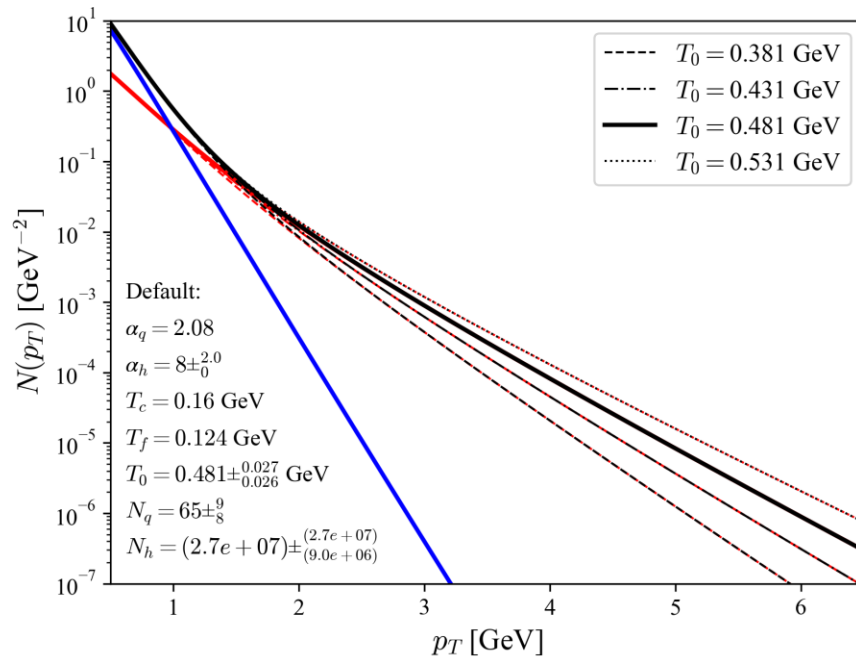
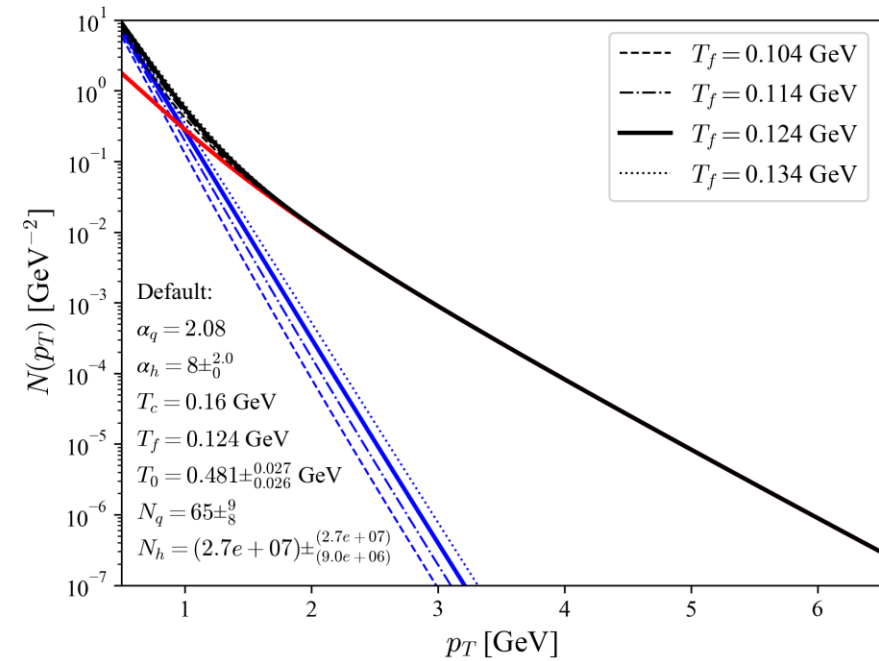
# Results



- EoS constrained by hydro eqs.: *qualitatively similar to 1QCD EoS*.
- Two component spectrum: *quantitatively good description of data* (CL=7.0%).
- Realistic value is obtained for the initial temperature.

# Results

- Varying  $T_f$  (upper panel)
- Varying  $T_0$  (lower left panel)
- Varying  $T_c$  (lower right panel)



# Open questions

- Where simple analytic hydrodynamic models succeed, more complex models struggle.
  - *Hydro ensures the thermalization of the system.*
- Will the 1+3d hydrodynamic model be successful in describing the PHENIX data as well? *According to our preliminary results, yes.*
- Will the 1+1d two-component hydrodynamic model be successful in describing the ALICE data as well?
- Perfect fluid models works well.
  - *Is the effect of viscosity negligible?*

# Open questions

- Where simple analytic hydrodynamic models succeed, more complex models struggle.
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- Perfect fluid models works well.
  - *Is the effect of viscosity negligible?*

***Thank you for your attention!***

# ▾ Backup slides

# Simple 1+3d model (L&K)

- Source function:

$$S(x, p)d^4x = \frac{g}{(2\pi\hbar)^3} \frac{\Theta(\tau - \tau_0) - \Theta(\tau - \tau_f)}{\tau_R} p^\mu u_\mu \exp\left(\frac{p^\mu u_\mu}{T}\right) dt d^3x$$

- Using the CCHK solution:

$$u_\mu = \gamma(1, \mathbf{v}) = \gamma\left(1, \frac{\mathbf{r}}{t}\right)$$

*Hubble-type velocity field*

$$T(\tau, s) = T_f \left(\frac{\tau_f}{\tau}\right)^{3/\kappa} \mathcal{T}(s)$$

*Inhomogeneous temperature profile*

$$\mathcal{T}(s) = \exp(-bs/2)$$

*Scale function is chosen to be Gaussian*

- Scale variable:

$$s = \frac{r^2}{R^2} (1 + \epsilon_2 \cos(2\alpha)) + \frac{r_z^2}{Z^2} \begin{matrix} \longrightarrow \epsilon_2 = \frac{Y - X}{Y + X} \\ \longrightarrow \frac{1}{R^2} = \left(\frac{1}{X^2} + \frac{1}{Y^2}\right) \end{matrix}$$



## Simple 1+3d model: observables

- Invariant transverse momentum spectrum:

$$N_1(p_T, \phi) = E \left. \frac{d^3 N}{dp^3} \right|_{p_z=0} = \left. \frac{d^3 N}{d\phi dp_T dy} \right|_{y=0} = \int S(t, r, \alpha, r_z, p_T, \phi) dt r d\alpha dr dr_z$$

$$\left. \frac{d^2 N}{p_T dp_T dy} \right|_{y=0} = \int_0^{2\pi} d\phi N_1(p_T, \phi)$$

- Elliptic flow:

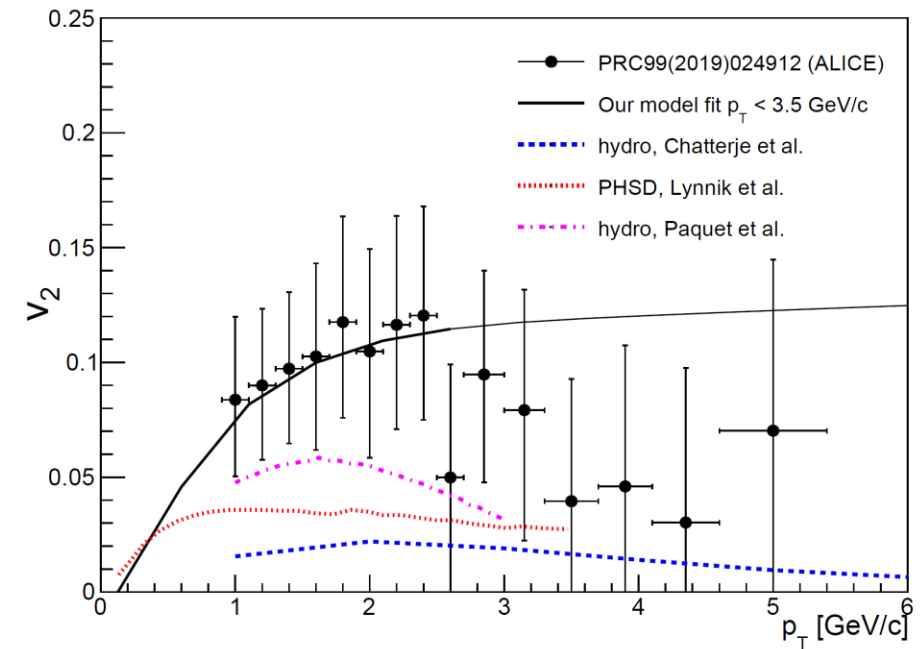
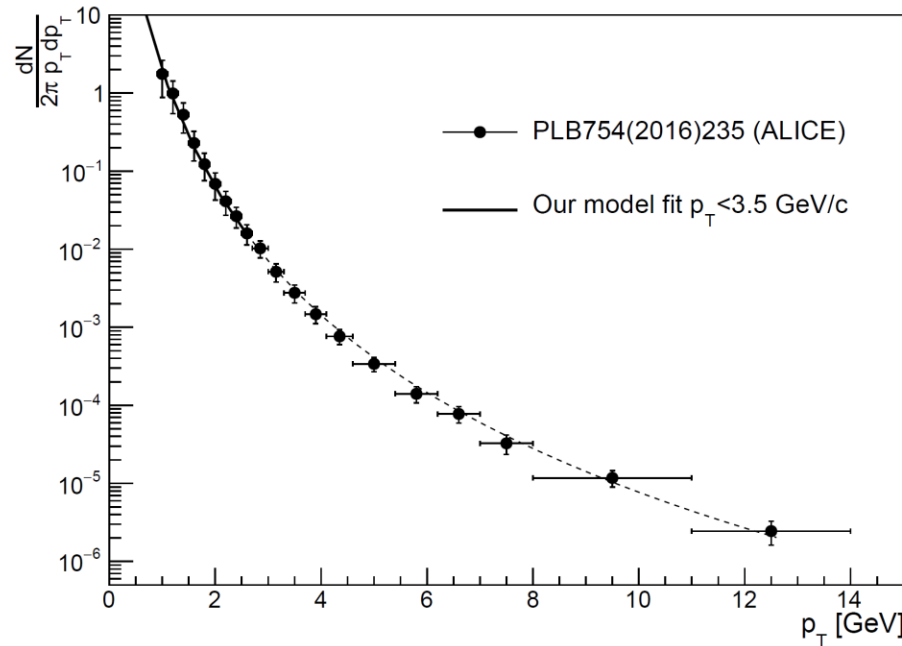
$$v_2(p_T) = \frac{\int_0^{2\pi} d\phi \cos(2\phi) N_1(p_T, \phi)}{\int_0^{2\pi} d\phi N_1(p_T, \phi)}$$

# Simultaneous fit

- Dataset: ALICE Pb+Pb@2.76 TeV, 0-20%

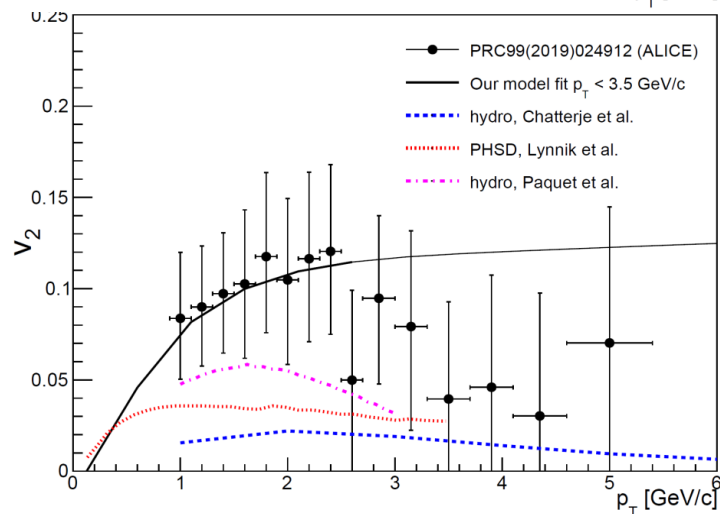
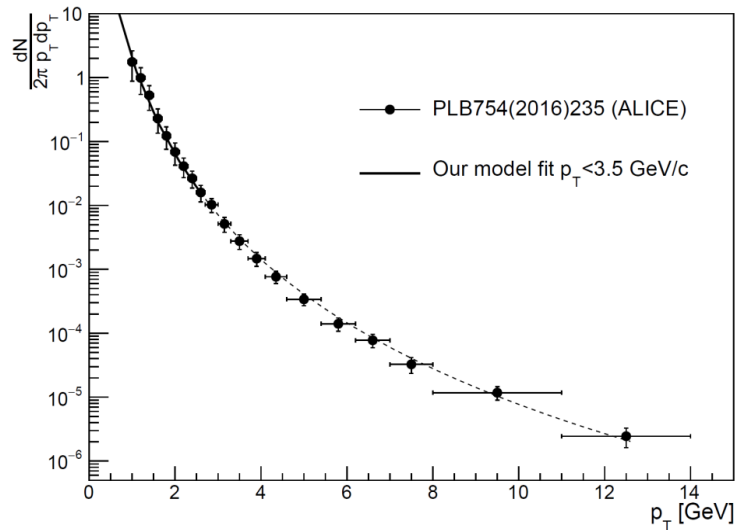
*Phys.Lett.B* 754 (2016) 235-248

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- Our model works better than the more complex microscopic models.
- Problem: why our model works on the whole  $p_T$ -range?

# Simultaneous fit



- $T_f = 0.123081 \pm 0.00766099$  GeV (limited)
- $\tau_f = 5.51272 \pm 0.484373$  fm/c (limited)
- $(dR/dt)/\sqrt{b} = 1.9$  (fixed)
- $(dZ/dt)/\sqrt{b} = 1.2$  (fixed)
- $\epsilon_2 = 0.271011 \pm 0.0120792$  (limited)
- $\kappa = 4.16755 \pm 0.074314$  (limited)
- normalization =  $0.0642342 \pm 0.0180748$