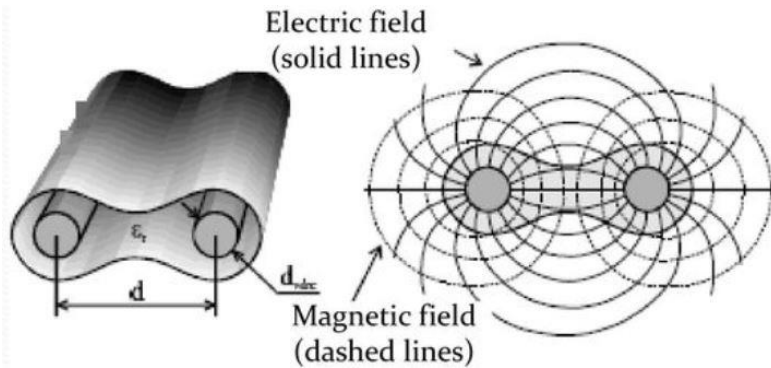


RF Engineering Transmission Lines II & Pillbox Resonator I

Christine Völlinger (CERN) & Manfred Wendt (BNL)

- Wave patterns in guided wave systems depend on the number of conductors used.

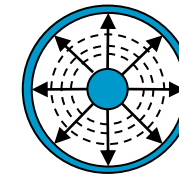
“Open” two-wire system (TEM)



TEM-propagation = 2 conductor system

Source: Ma, *Electromagnetic Waves and Applications Part III, univ. lecture*

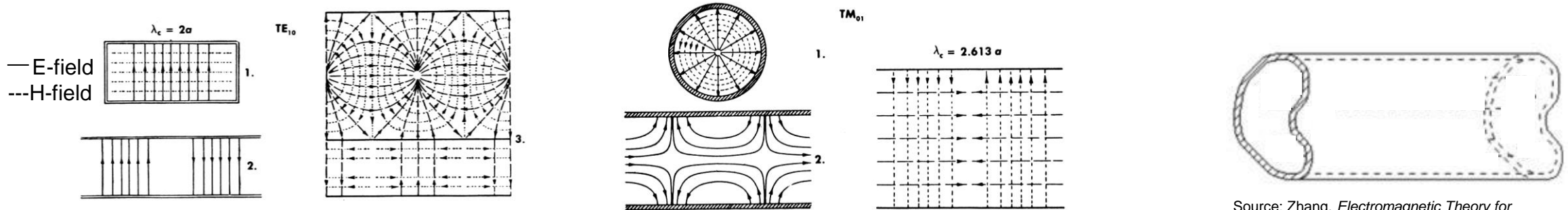
Coaxial line (TEM)



— E-field --- H-field

Picture: Coaxial loads for the SPS 200 MHz cavity © CERN

Uniform waveguides rectangular, round, random cross-section (all non-TEM)



Source: Saad, *Microwave Engineers Handbook, vol. 1, Artech House*

Source: Zhang, *Electromagnetic Theory for Microwaves and Optoelectronics, 2nd ed., Springer*

Rectangular waveguide



Circular waveguide



Arbitrarily shaped cross-section waveguide



Source: Zhang, *Electromagnetic Theory for Microwaves and Optoelectronics*, 2nd ed., Springer

- Waveguides exist in different shapes. They are just hollow metallic tubes with uniform cross-sections.
- The metallic waveguide is a completely enclosed system without any radiation losses.
- Further, losses due to a dielectric layer, as we have seen for the coaxial line do not exist, equally the inner conductor of the coaxial line is eliminated.
- Waveguides present a one-conductor system, so *no TEM-mode propagation* is possible. Propagation happens in TE-mode (*Transverse-Electric*) and TM-mode (*Transverse-Magnetic*). These modes have a lower frequency limit, the so-called cut-off frequency below which propagation is not possible. The cut-off frequency depends on the dimensions of the waveguide cross-section.

Rectangular waveguide



Circular waveguide

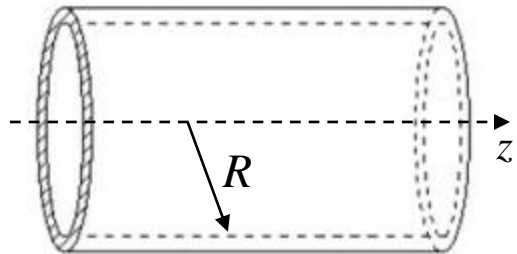


Arbitrarily shaped cross-section waveguide



Source: Zhang, *Electromagnetic Theory for Microwaves and Optoelectronics*, 2nd ed., Springer

- Waves are following the waveguide shape, even if it is bend.
- For this, the propagating waves (formerly unbound) now need to fulfil boundary conditions on the walls of the waveguide.
- Remember? From EM-field description, we know that:
 - on *perfect conducting surfaces*, the tangential electric field has to vanish to fulfill the boundary condition. We speak of an “electric wall”: $E_{\text{tangential}} = 0$.
 - the walls of a waveguide can approximated as a perfect conductor (exception: loss calculation).



- Like all transmission lines, also the circular waveguide is characterized by a propagation constant, an attenuation constant and a characteristic impedance.
- These quantities are derived by field theory analysis (see lecture of A. Mostacci). Propagation is usually assumed in z -direction.
- For a perfect conducting tube (no resistive attenuation) with radius R , we obtain propagation constants for the TE-mode and the TM-mode:

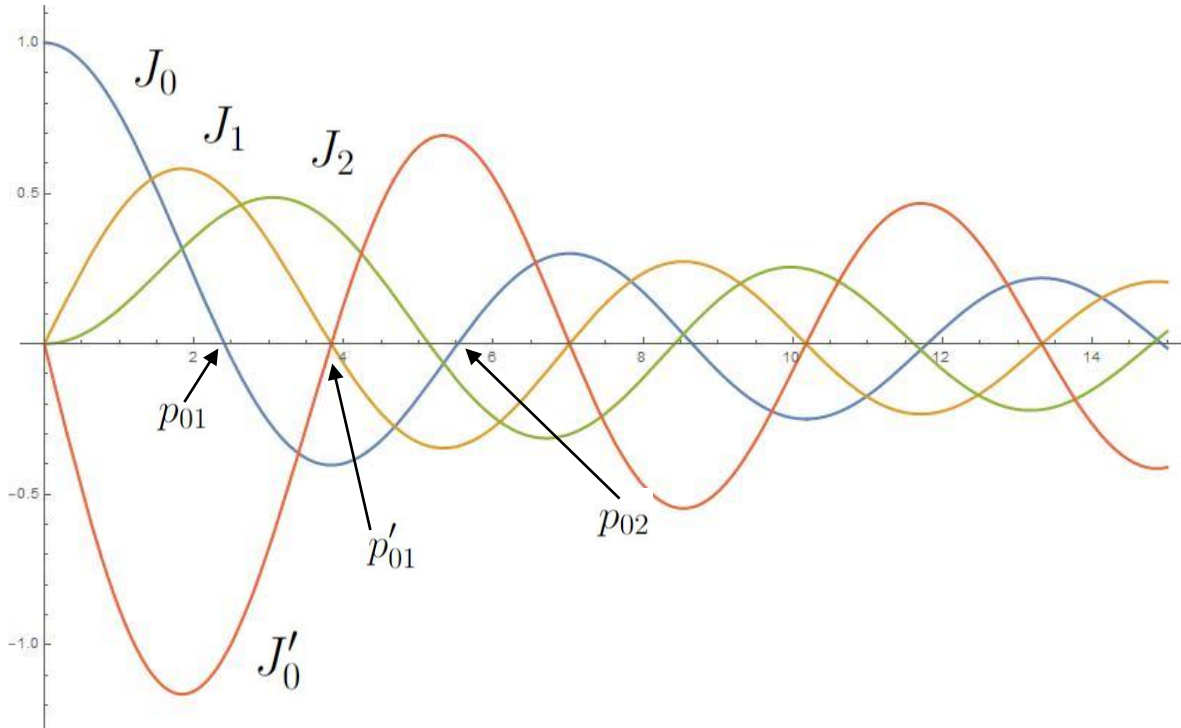
$$\beta_{nm} = \sqrt{\omega^2 \epsilon \mu - \left(\frac{p'_{nm} \text{ OR } p_{nm}}{R} \right)^2}$$

$p_{nm} \rightarrow$ Roots of the Bessel-function for TM-mode $J_n(p_{nm}) = 0$
 $p'_{nm} \rightarrow$ Roots of the derivative of the B.F. for TE-mode $J'_n(p'_{nm}) = 0$

- This leads to the so-called cut-off frequencies for the different modes:

$$f_{c,nm} = \frac{1}{2\pi \sqrt{\mu \epsilon}} \left(\frac{p'_{nm} \text{ OR } p_{nm}}{R} \right)$$

Recall the Bessel function and the derivative of the Bessel function with their roots:



$p_{01} = 1^{\text{st}}$ root of Bessel function of first type J_0

$p_{02} = 2^{\text{nd}}$ root of Bessel function of first type J_0

$p'_{01} = 1^{\text{st}}$ root of derivative of Bessel function of first type J'_0

TABLE 3.3 Values of p'_{nm} for TE Modes of a Circular Waveguide

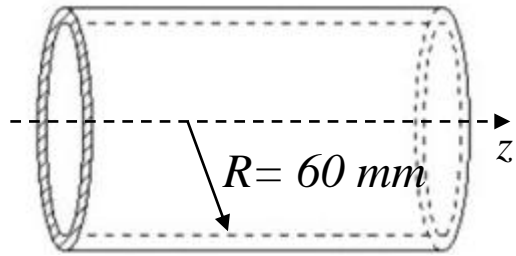
n	p'_{n1}	p'_{n2}	p'_{n3}
0	3.832	7.016	10.174
1	1.841	5.331	8.536
2	3.054	6.706	9.970

TABLE 3.4 Values of p_{nm} for TM Modes of a Circular Waveguide

n	p_{n1}	p_{n2}	p_{n3}
0	2.405	5.520	8.654
1	3.832	7.016	10.174
2	5.135	8.417	11.620

Source: Pozar, *Microwave Engineering*, 4th ed., Wiley

Let's see an example.



→ Calculating cut-off frequencies for a circular waveguide with $R = 60 \text{ mm}$.

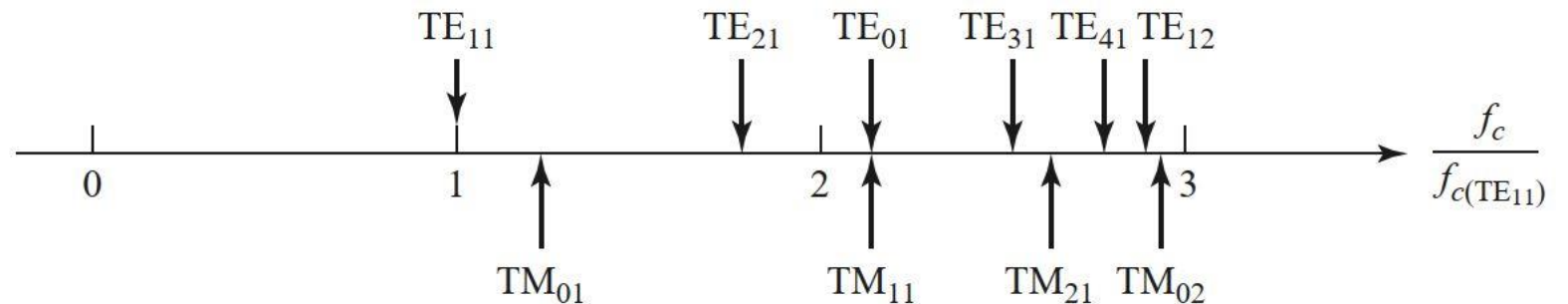
$$f_c(\text{TE}_{11}) = \frac{1}{2\pi\sqrt{\mu_0\epsilon_0}} \frac{p'_{11}}{R} = \frac{2.9 \cdot 10^8}{2\pi} \frac{1.841}{0.06} = 1.416 \text{ GHz}$$

$$f_c(\text{TM}_{01}) = \frac{1}{2\pi\sqrt{\mu_0\epsilon_0}} \frac{p_{01}}{R} = \frac{2.9 \cdot 10^8}{2\pi} \frac{2.405}{0.06} = 1.850 \text{ GHz}$$

Frequencies where these modes start propagating.

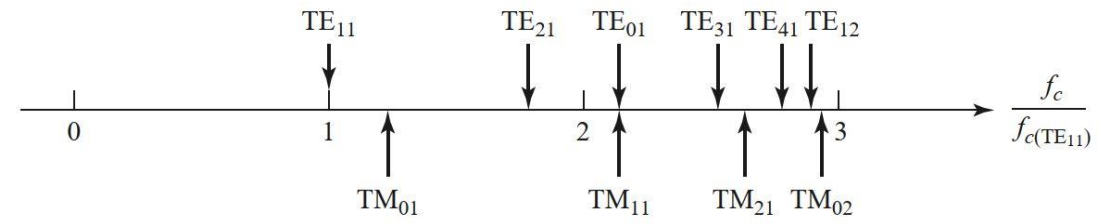
$$f_{c,nm} = \frac{1}{2\pi\sqrt{\mu\epsilon}} \left(\frac{p'_{nm} \text{ (TE-mode) OR } p_{nm} \text{ (TM-mode)}}{R} \right)$$

Chart of cut-off frequencies relative to f_c of TE_{11} -mode:



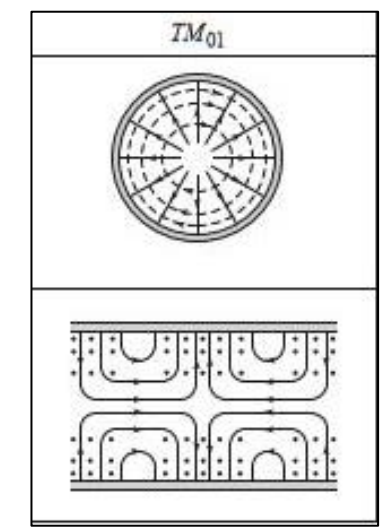
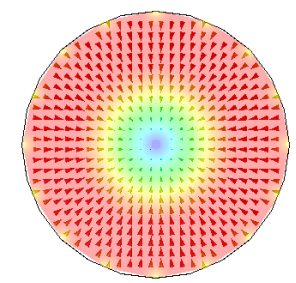
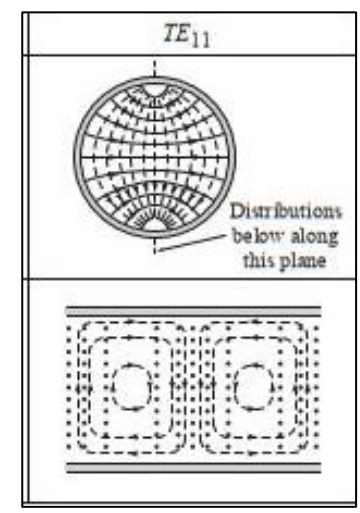
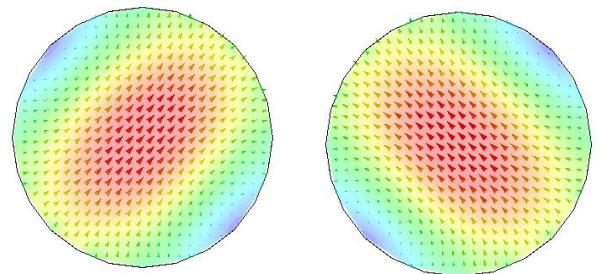
Source: Pozar, *Microwave engineering*, 4th ed., Wiley

Waveguide modes and field patterns



1st mode is TE_{11} . Electric field is transverse.
Mode has 2 polarisations (orientations of the electrical field).

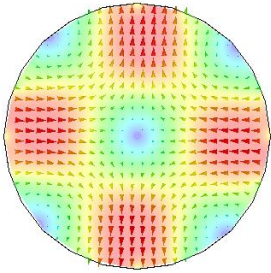
- 2nd mode is TM_{01} . Magnetic field is transverse.



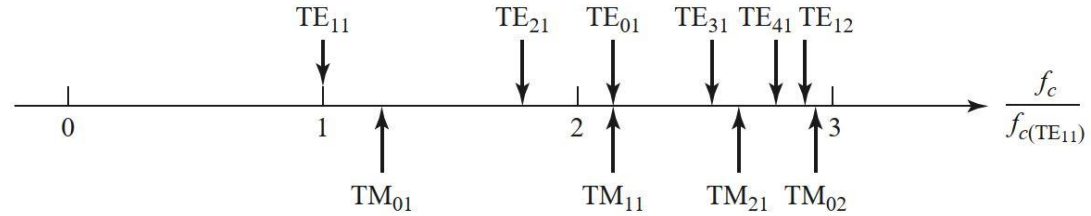
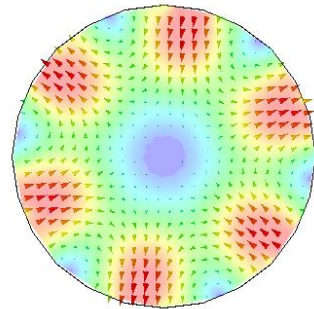
Simulaton pictures: courtesy E. Jensen, source pictures and chart: Pozar, *Microwave engineering*, 4th ed., Wiley

Waveguide modes and field patterns

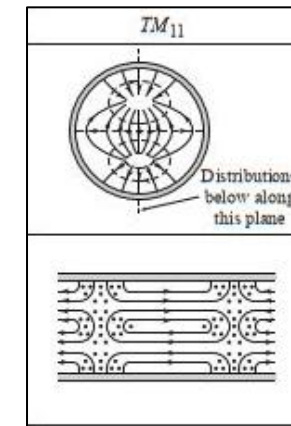
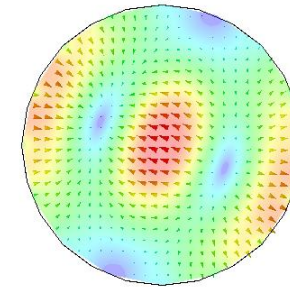
- 3rd mode is TE_{21} .
Electric field is transverse



- 6th mode is TE_{31} .
Electric field is transverse.



- 4th mode is TM_{11} . Magnetic field is transverse.



Further field plots can e.g. be found here: Lee, Lee and Chuang, *Plot of Modal Field Distribution in Rectangular and Circular Waveguides*. IEEE, Trans. Microwave Theory and Techniques, Vol. MTT-33, no. 3, March 1985

What is actually happening if the wave frequency in a waveguide is below cut-off?

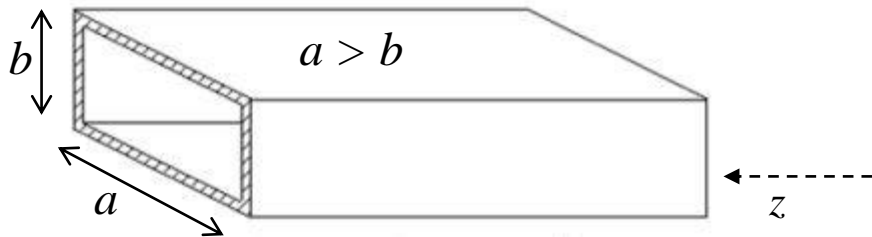
Remember? Propagation constant is defined: $\gamma = \alpha + j\beta$

$$\beta_{nm} = \sqrt{\omega^2 \epsilon \mu - \underbrace{\left(\frac{p'_{nm} \text{ OR } p_{nm}}{R} \right)^2}_{\epsilon \mu \omega_c^2}} \quad \Rightarrow \quad \beta^2 = \mu \epsilon (\omega^2 - \omega_c^2)$$

Now for frequencies below cut off, β would become complex: $\beta = j \cdot \sqrt{\mu \epsilon} \sqrt{\omega_c^2 - \omega^2}$

Then, the propagation constant γ becomes fully real, and the wave is damped even in the loss-less case ($\alpha=0$), since the propagation is following: $e^{-\beta z}$

This is called an evanescent wave = a wave that cannot propagate.



Source: Zhang, *Electromagnetic Theory for Microwaves and Optoelectronics*, 2nd ed., Springer

- All field quantities are derived by field theory analysis (see lecture of A. Mostacci).
- Wave propagation is usually assumed in z -direction.
- The rectangular waveguide is characterized by a propagation constant, attenuation constant and a characteristic impedance.

- We distinguish again between TE-modes and TM-modes. TEM-modes do not exist in a waveguide.
- We will see that the hollow rectangular waveguide can propagate TE- and TM-modes, and that these modes have cut-off frequencies below which propagation is not possible, similar to the TE- and TM-modes of the circular waveguide.

The *derivation* of the cut-off frequency is straightforward, please refer to textbooks, if you are interested. The field components are obtained by the separation method.

TE-modes

- TE-modes are characterized by a zero *electric* field in propagation direction. The electric field is in the transverse plane, only.
- Each mode has a cut-off frequency:

$$f_{c,mn} = \frac{1}{\sqrt{\mu\epsilon}} \sqrt{\left(\frac{m}{a}\right)^2 + \left(\frac{n}{b}\right)^2}$$

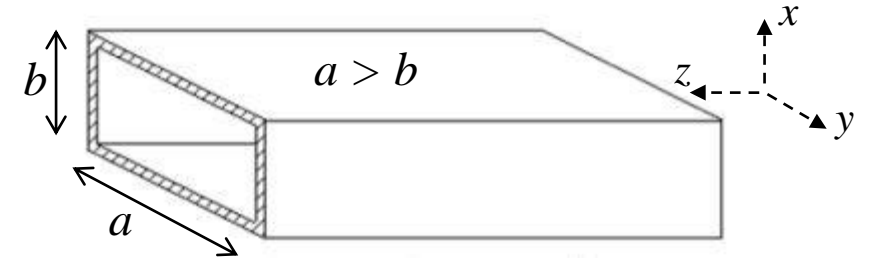
$$\omega_{c,mn} = \frac{1}{\sqrt{\mu\epsilon}} \sqrt{\left(\frac{m\pi}{a}\right)^2 + \left(\frac{n\pi}{b}\right)^2}$$

- TE₁₀-mode is the mode with the lowest cut-off frequency: $\omega_{c,10} = \frac{1}{2a\sqrt{\mu\epsilon}}$

The mode with the *lowest cut-off frequency* is called *the dominant mode*.

Note that if $n=m=0$, all field components become zero, there is no TE₀₀ mode.

A waveguide is called *overmoded*, if more than one mode is propagating in the waveguide at the same time.



Source: Zhang, *Electromagnetic Theory for Microwaves and Optoelectronics*, 2nd ed., Springer

TM-modes

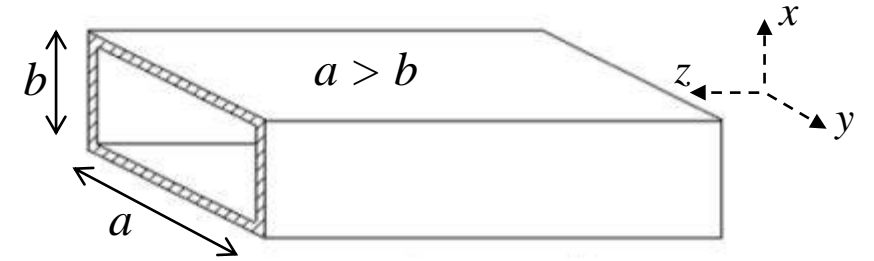
- TM-modes are characterized by a zero *magnetic* field in propagation direction. The magnetic field is in the transverse plane, only.
- Also in this case, each mode has a cut-off frequency (identical to TE):

$$f_{c,mn} = \frac{1}{\sqrt{\mu\epsilon}} \sqrt{\left(\frac{m}{a}\right)^2 + \left(\frac{n}{b}\right)^2} \quad \omega_{c,mn} = \frac{1}{\sqrt{\mu\epsilon}} \sqrt{\left(\frac{m\pi}{a}\right)^2 + \left(\frac{n\pi}{b}\right)^2}$$

- TM_{11} has as cut-off frequency: $\omega_{c,11} = \frac{1}{\sqrt{\mu\epsilon}} \sqrt{\left(\frac{\pi}{a}\right)^2 + \left(\frac{\pi}{b}\right)^2}$

There is no TM_{00} -mode, neither do TM_{01} or TM_{10} exist. For $n=m=0$, all field components become zero.

Note that TE_{10} -mode remains the dominant mode, it is the mode with the lowest cut-off frequency.



Source: Zhang, *Electromagnetic Theory for Microwaves and Optoelectronics*, 2nd ed., Springer

TM-modes

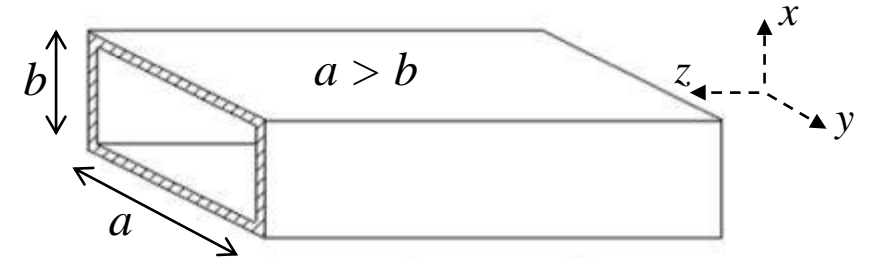
- TM-modes are characterized by a zero *magnetic* field in propagation direction. The magnetic field is in the transverse plane, only.
- Also in this case, each mode has a cut-off frequency (identical to TE):

$$f_{c,mn} = \frac{1}{\sqrt{\mu\epsilon}} \sqrt{\left(\frac{m}{a}\right)^2 + \left(\frac{n}{b}\right)^2} \quad \omega_{c,mn} = \frac{1}{\sqrt{\mu\epsilon}} \sqrt{\left(\frac{m\pi}{a}\right)^2 + \left(\frac{n\pi}{b}\right)^2}$$

- TM_{11} has as cut-off frequency: $\omega_{c,11} = \frac{1}{\sqrt{\mu\epsilon}} \sqrt{\left(\frac{\pi}{a}\right)^2 + \left(\frac{\pi}{b}\right)^2}$

There is no TM_{00} -mode, neither do TM_{01} or TM_{10} exist. For $n=m=0$, all field components become zero.

Note that TE_{10} -mode remains the dominant mode, it is the mode with the lowest cut-off frequency.



Source: Zhang, *Electromagnetic Theory for Microwaves and Optoelectronics*, 2nd ed., Springer

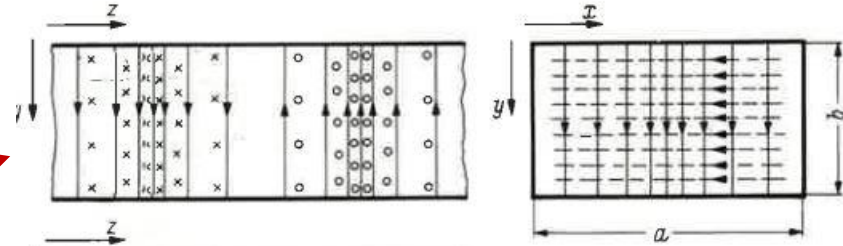
Rectangular Waveguides (4/6)

Field pattern for TE_{01} -mode

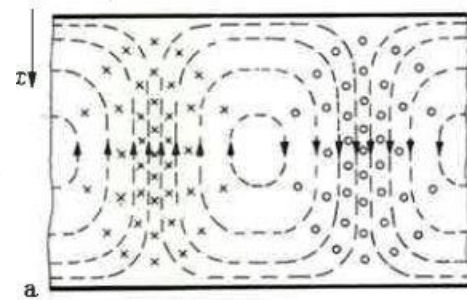
What do we actually see?

Side view narrow
(y/z-plane)

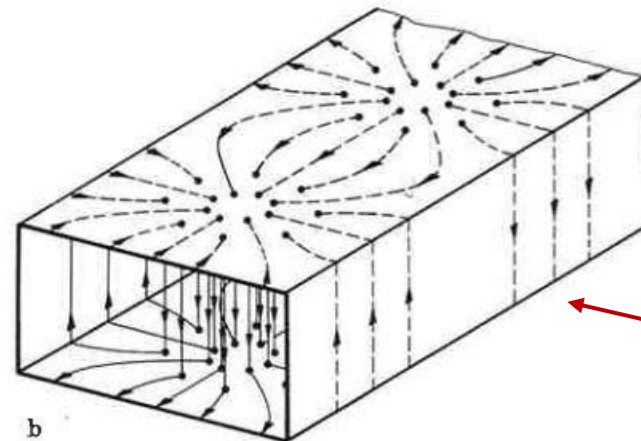
Side view broad
(x/z-plane)



View from front
(x/y-plane)



E-field ---H-field

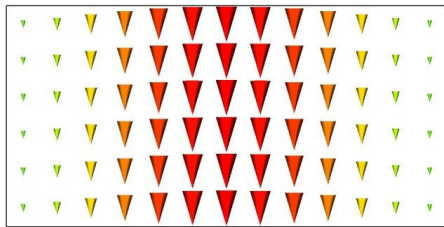


Wall currents and displacement currents inside the waveguide.

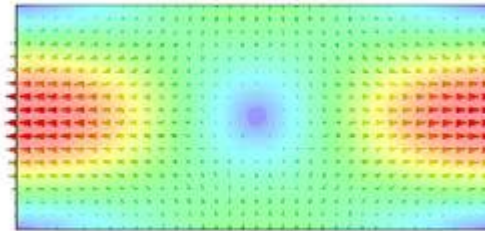
source: Zinke/Brunswig, *Hochfrequenztechnik*, 5th ed., Springer

Fundamental and higher waveguide modes and field patterns for TE-modes

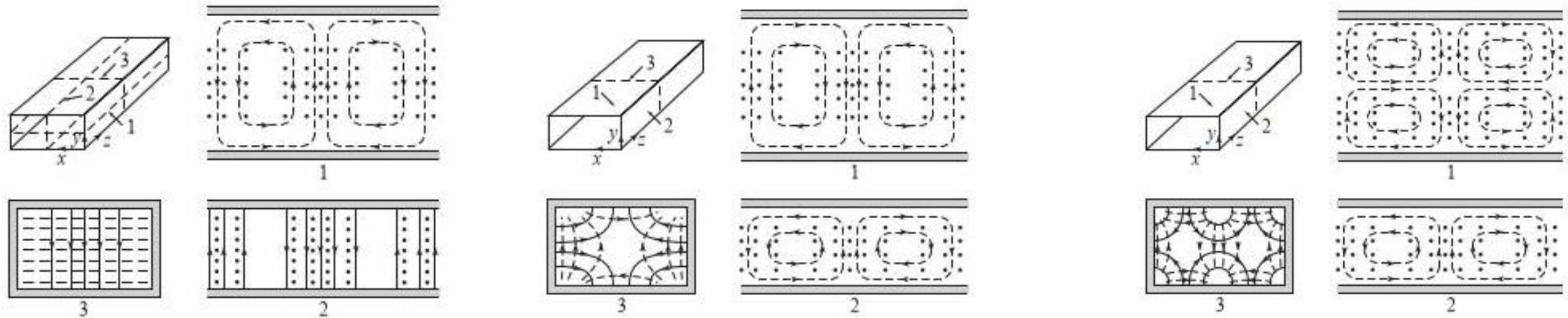
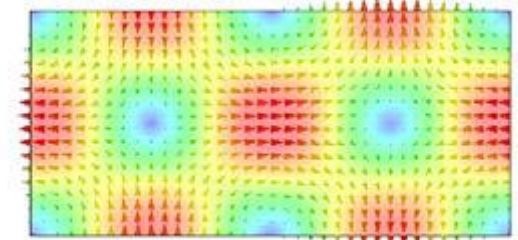
TE₁₀-mode



TE₁₁-mode



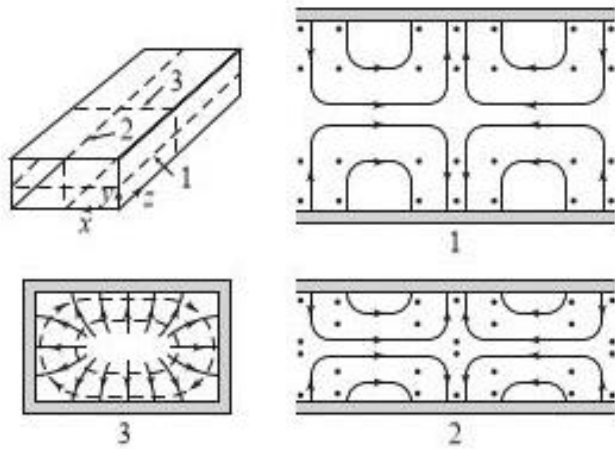
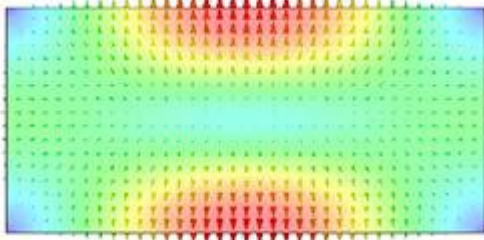
TE₂₁-mode



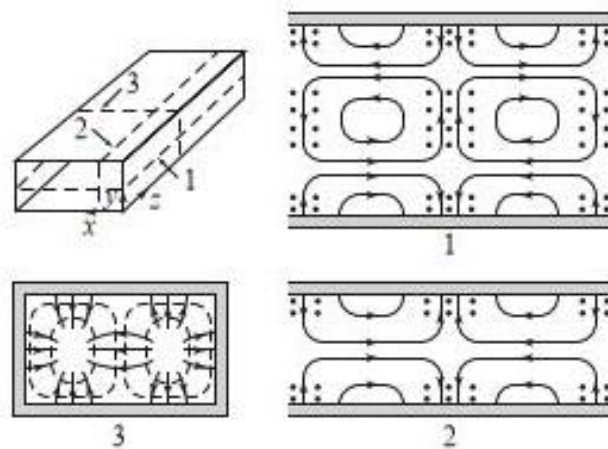
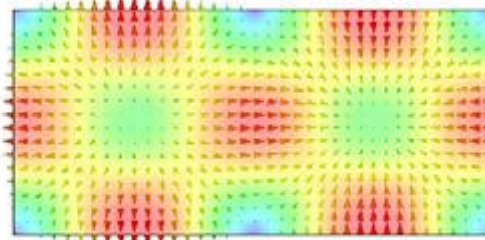
Simulaton pictures: courtesy E. Jensen, field pattern source: Pozar, *Microwave engineering*, 4th ed., Wiley

Fundamental and higher waveguide modes and field patterns for TM-modes

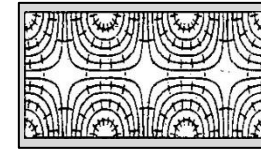
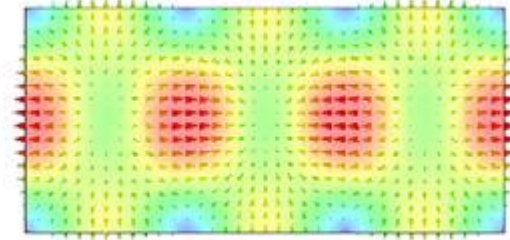
TM₁₁-mode



TM₂₁-mode



TM₃₁-mode



Simulaton pictures: courtesy E. Jensen, field pattern source: Pozar, *Microwave engineering*, 4th ed., Wiley, TM31-mode from Lee et al., MTT-33, 1985

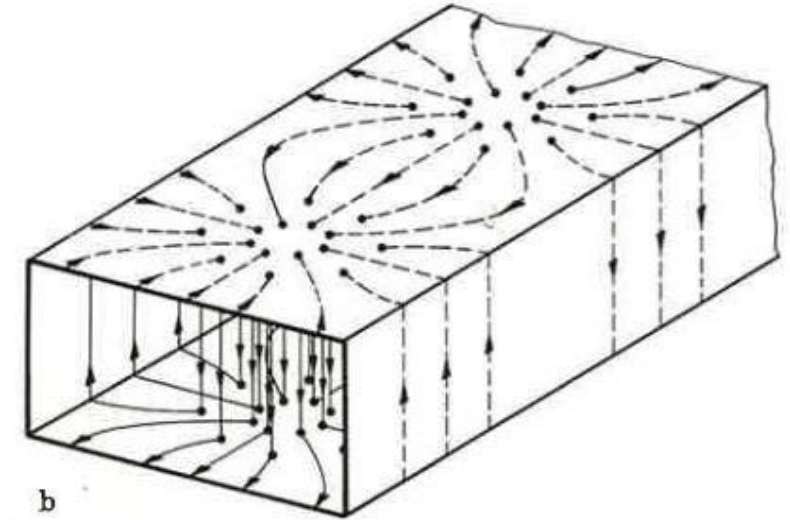
Amplifiers of the 200 MHz travelling wave system in the CERN SPS and their coaxial power lines leading to the underground installation (feeder lines)



Quiz!
Test your knowledge.

- **Waveguides are used for TEM signal transport**

- true
- false

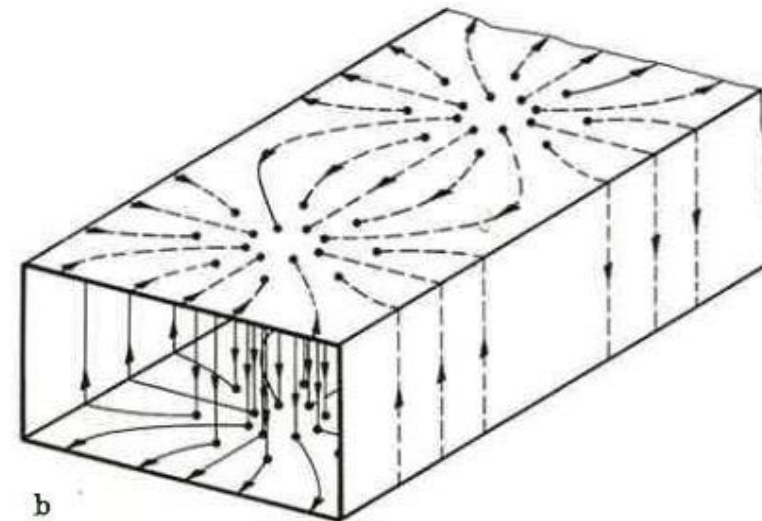


- **In a waveguide, lossy wall-currents exist in...?**

- circular waveguides due to the shape of their outer conductor
- rectangular waveguides due to the shape of their outer conductor
- rectangular waveguides only, but not in circular waveguides and not in elliptical waveguides
- all waveguides

- **Waveguides are used for TEM signal transport**

- true
- false (correct answer)

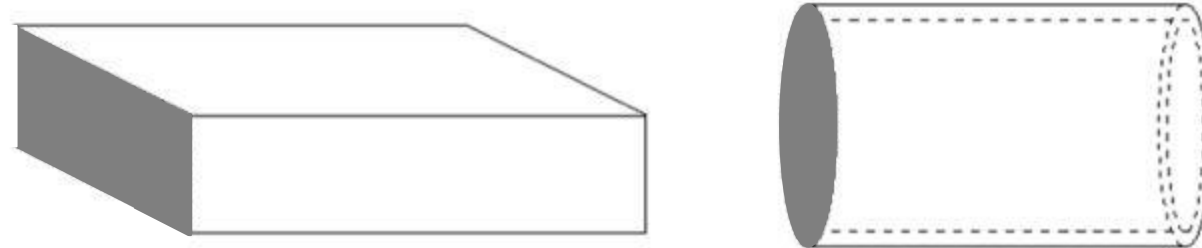


- **In a waveguide, lossy wall-currents exist in...?**

- circular waveguides due to the shape of their outer conductor
- rectangular waveguides due to the shape of their outer conductor
- rectangular waveguides only, but not in circular waveguides and not in elliptical waveguides
- all waveguides (correct answer)

Microwave resonators can be built from waveguides by closing the open ends.

We call this an *RF-resonator* or a *cavity*.



Source: Zhang, *Electromagnetic Theory for Microwaves and Optoelectronics*, 2nd ed., Springer

- A cavity stores electric and magnetic energy inside the hollow body.
- The frequency of the electromagnetic field resonance depends on the cavity dimensions.
- Same as with the rectangular and the circular waveguide, the electromagnetic field has to fulfil the boundary conditions on the resonator walls. The field builds up in resonant modes.
- Cavities that are used for the acceleration of particles in our accelerators are mostly of a cylindrical and flat shape. This is why we call them pillbox cavities.
- We will later see how resonances can be excited in cavities by feeding RF-signal in, and how such a cavity can be used in accelerators for particle acceleration.

- Resonators are classified by their quality factor Q .
- The quality factor (or Q -value) can be used as a measure of “how well the cavity is resonating” (more details will come).

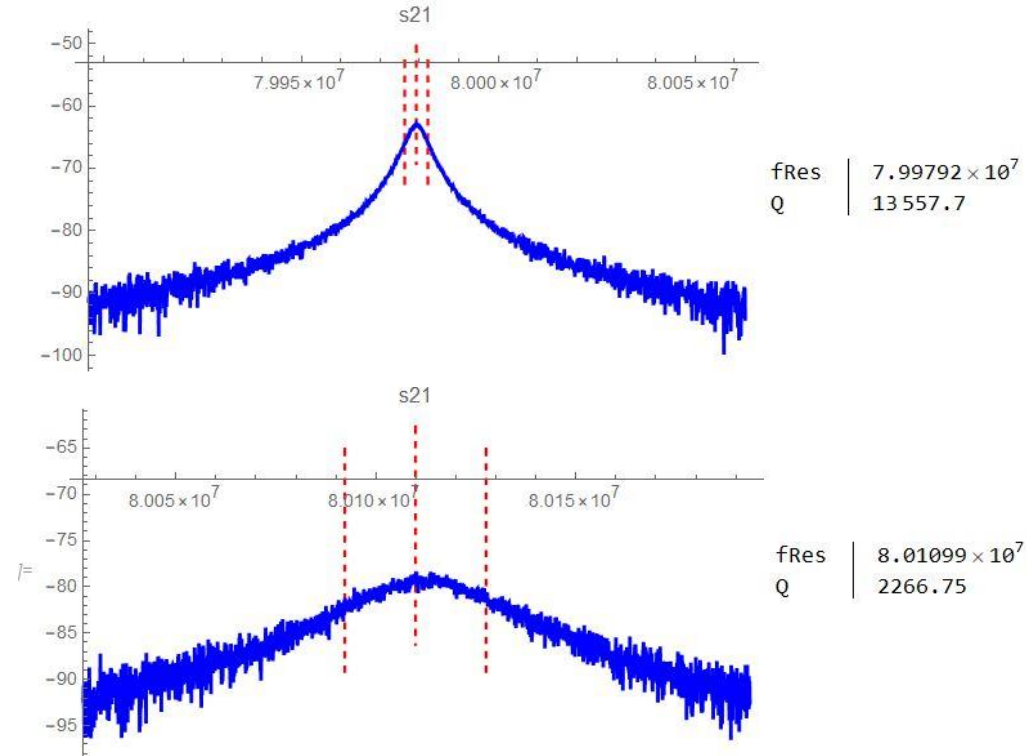
define $Q_0 = 2\pi \frac{\text{energy stored in the resonator}}{\text{energy dissipated in the resonator}}$

$Q_0 = \text{unloaded } Q$

- High Q -value is desired in accelerating cavities; Q -value is one of the accelerator efficiency figures-of-merit.
- Q -value reduces e.g. due to the power dissipated in the metallic walls or other loss mechanisms.
- The connection of the cavity resonator to the outer world will reduce its Q -value as well (*we say “it loads the cavity with an additional loss mechanism”*).
- If the cavity is filled with a dielectric, this will produce additional losses → we don't have this in accelerators.



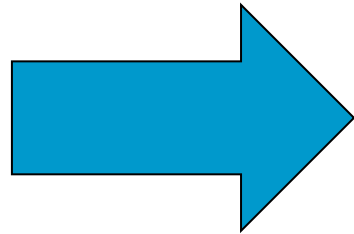
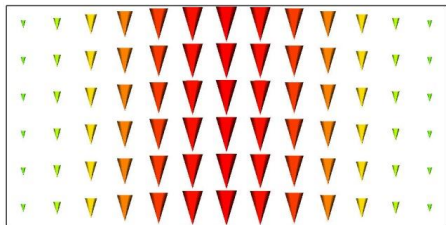
Example: measurements taken on the PS 80MHz pillbox cavity



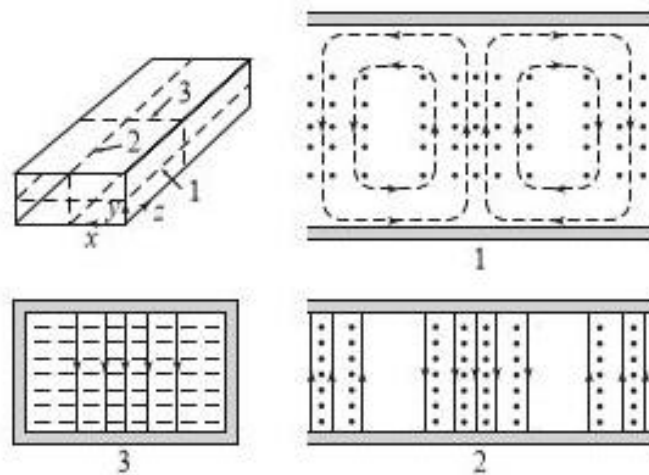
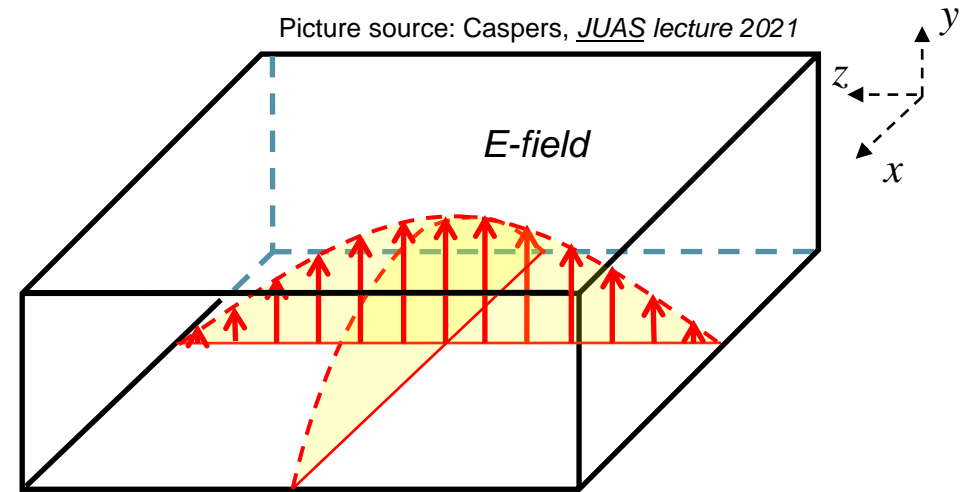
- We tested this cavity for Q-deterioration and shifting of its fundamental mode (~80 MHz).
- Q-values obtained from 3-dB-measurement (see dashed red lines, method will be introduced later).
→ Plot on the top has a higher Q-value than the plot below.

From the field pattern of the TE₁₀-mode in the rectangular waveguide → imagine that the resonant field in a rectangular resonator will build up in a similar way.

Waveguide TE₁₀-mode



Picture source: Caspers, *JUAS lecture 2021*



- The field is building up in a standing wave pattern, using sine and cosine functions.
- From the boundary condition on the resonator wall, the electric field has to be zero, thus we can see multiples of electric field maxima.
- The *mode indices* are counting the maxima of the field along one axis (we see either a TE_{xyz} or a TM_{xyz}- mode).
- The mode shown is a TE₁₀₁-mode.

Source: Pozar, *Microwave engineering*, 4th ed., Wiley

Formulae for a rectangular resonator
(for derivation see lecture of A. Mostacci or textbooks).

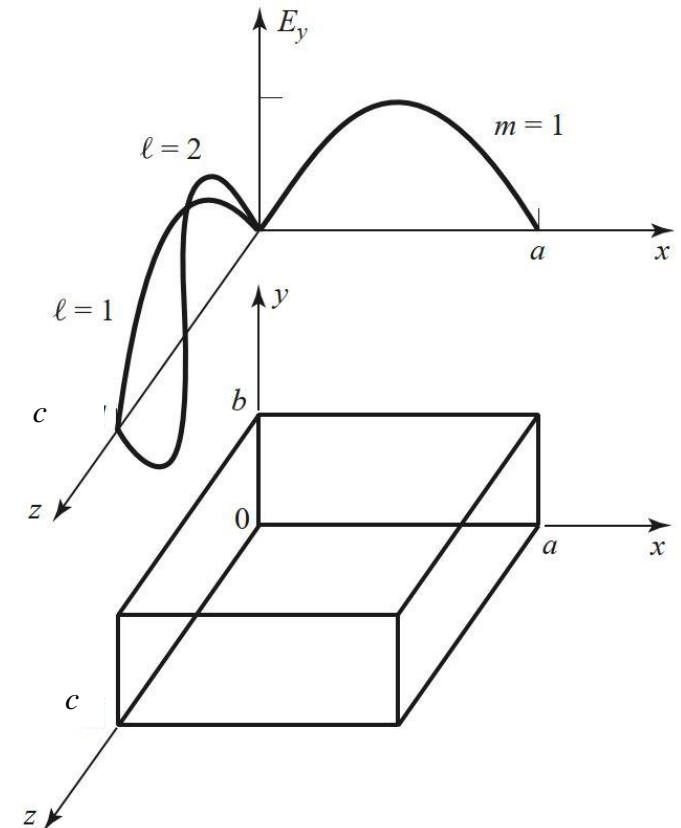
Source: Pozar, *Microwave Engineering*, 4th ed., Wiley

- resonant wavelength:
$$\lambda_0 = \frac{2}{\sqrt{\left(\frac{m}{a}\right)^2 + \left(\frac{n}{b}\right)^2 + \left(\frac{l}{c}\right)^2}}$$

- Resonant frequencies for TE_{mnl} -mode and TM_{nml} -modes:

$$f_{mnl} = \frac{1}{2\pi\sqrt{\mu\epsilon}} \sqrt{\left(\frac{m\pi}{a}\right)^2 + \left(\frac{n\pi}{b}\right)^2 + \left(\frac{l\pi}{c}\right)^2}$$

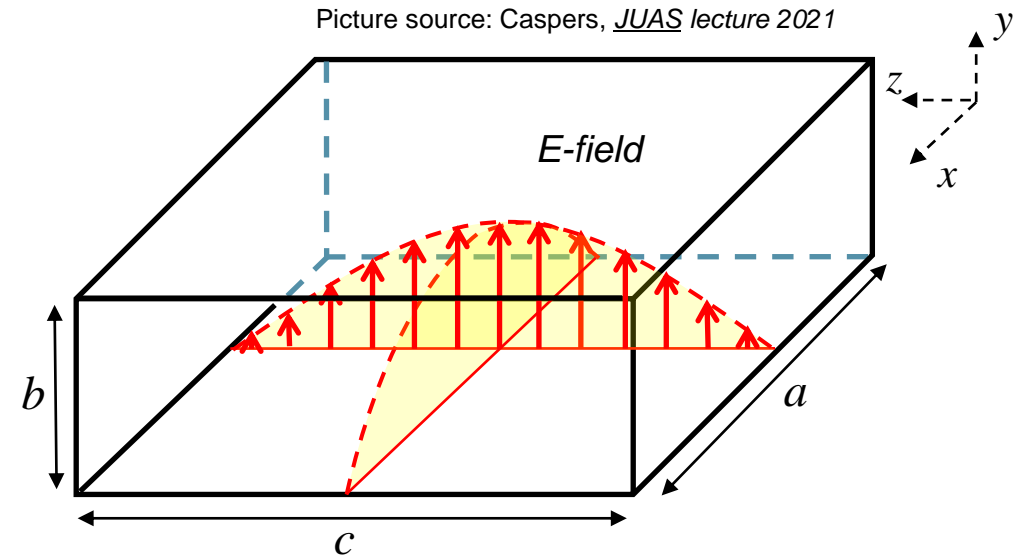
- Q-value for TE_{101} :
$$Q_{TE101} = \frac{\lambda_0 b}{\delta} \frac{(a^2 + b^2)^{3/2}}{2c^3(a + 2b) + a^3(c + 2b)}$$



TE_{101} -mode is the lowest resonant mode, hence also the dominant mode in this resonator.

For the case $a=c$, the formulae simplifies strongly.

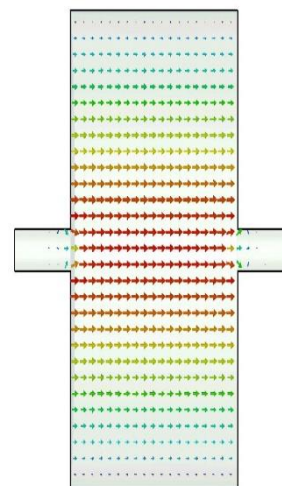
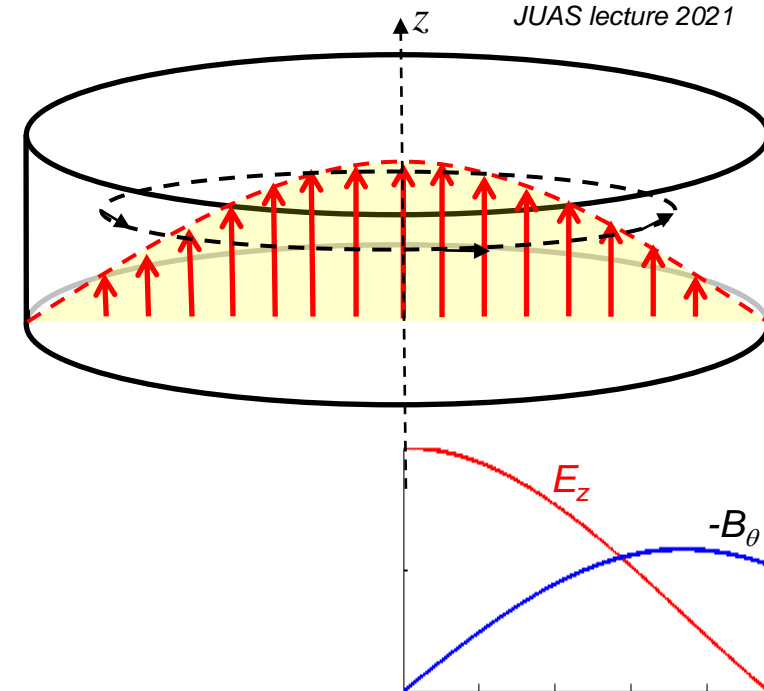
- resonant wavelength: $\lambda_0 = \sqrt{2}a$
- Q-value for TE_{101} : $Q_0 = \frac{1}{\delta} \frac{ab}{a + 2b}$



A cylindrical cavity resonator can be seen as a section of circular waveguide shorted at both ends.

- From the TE_{11} -mode (dominant mode for circular waveguide)
 - TE_{111} -mode is the dominant TE-mode for a cylindrical cavity.
 - TM_{010} -mode is the dominant TM-mode for a cylindrical cavity.
 - TM_{010} -mode mode is used for acceleration as it has a large electric field along the z-axis.
 - Electric field is described by Besselfunction J_0 .
 - TM_{010} -mode and pillbox mode indices?

Picture source: Caspers, JUAS lecture 2021



Pillbox cavity is flat compared to a “long” cylindrical cavity. Also, beam-pipe connection is needed.

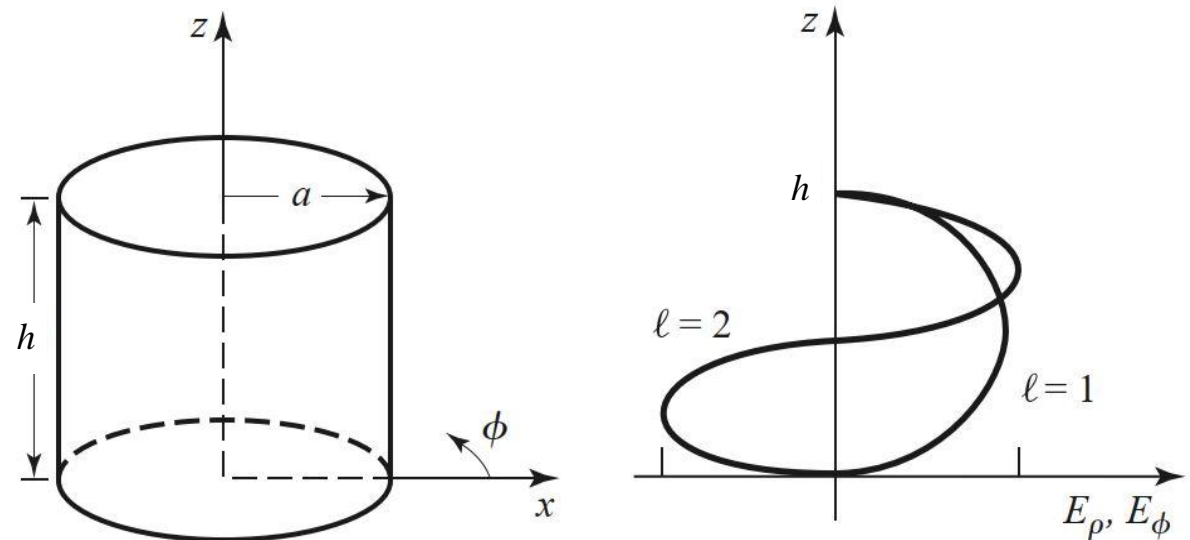
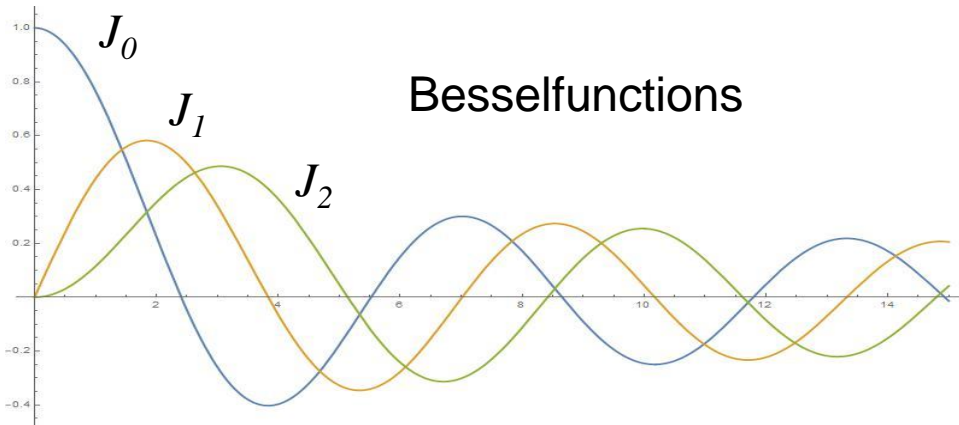
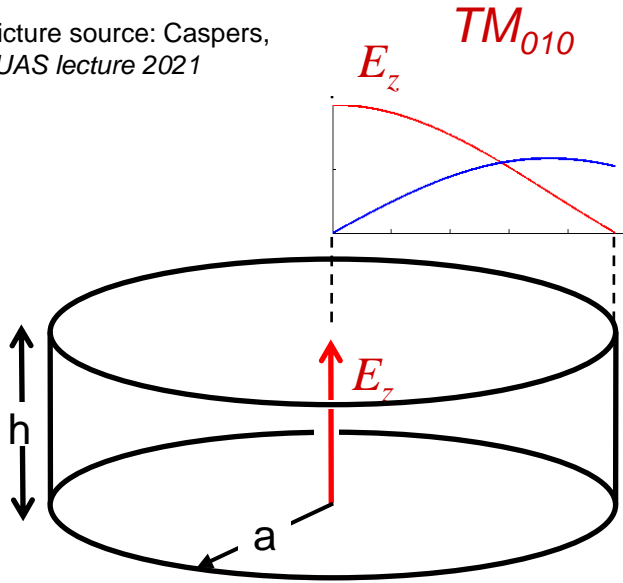
Simulation: P. Kramer

Pillbox Cavity Resonators (2/6)

Picture source: Caspers, JUAS lecture 2021

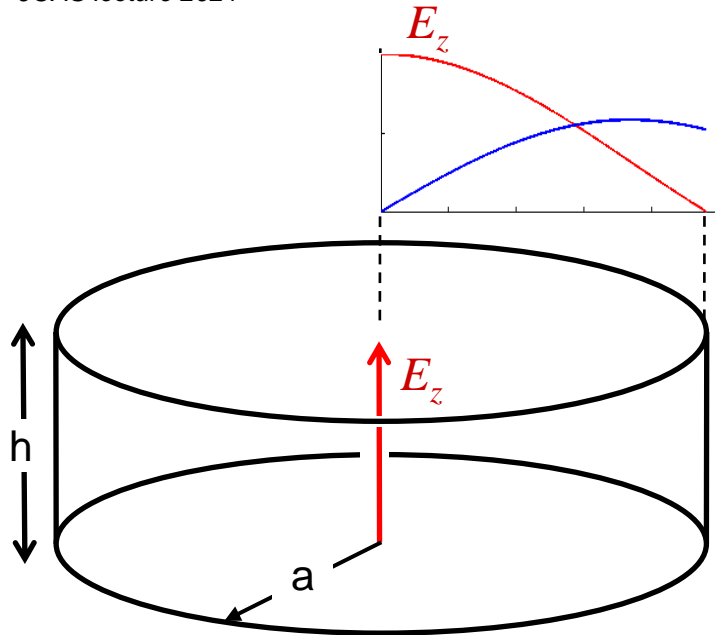
Mode indices of a pillbox cavity are as follows:

- The electric field on the cavity's axis following the Besselfunction J_0
- The *first mode-index gives the order of the Bessel function.*
- The *second mode-index indicates the root of the Bessel function* which is the number of its zero-crossings.
- The *third mode-index is the number of half waves (maxima) along the z-axis.*



Source: Pozar, Microwave engineering, 4th ed., Wiley

Picture source: Caspers,
JUAS lecture 2021



Formulae for a circular cavity resonator (for derivation see lecture of A. Mostacci or textbooks).

- Resonant frequencies for TE_{nml}-mode:

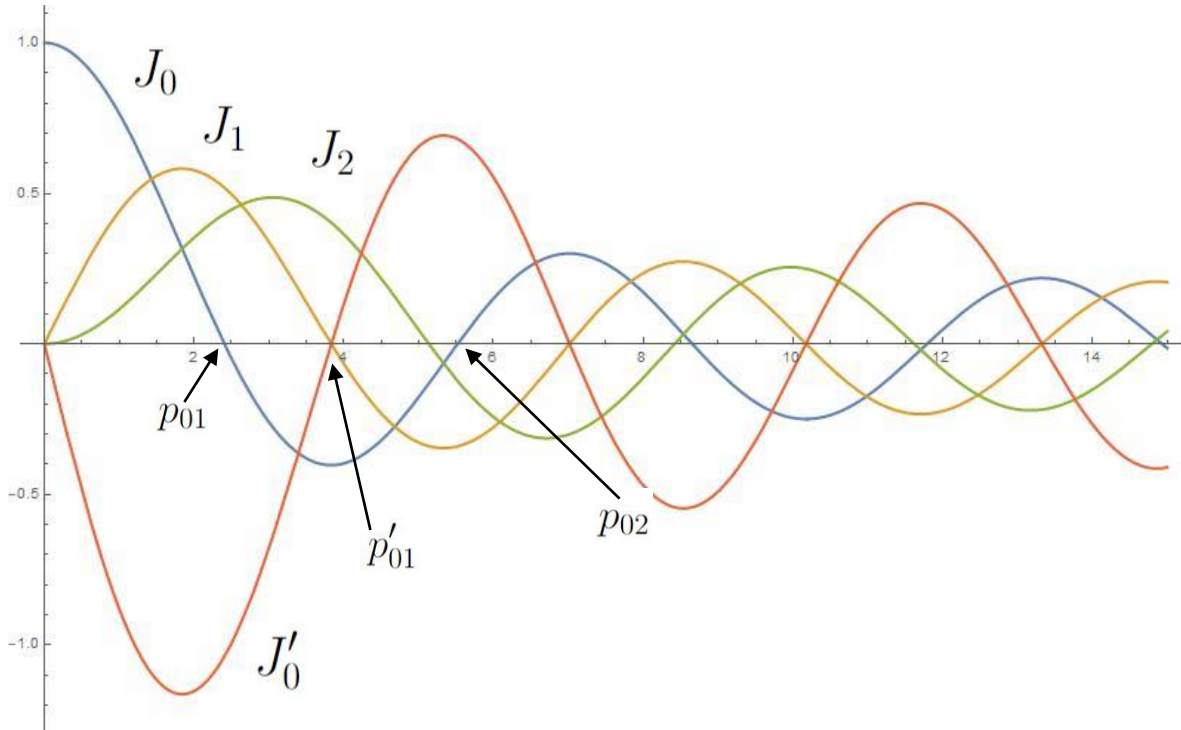
$$f_{\text{TE},nml} = \frac{c_0}{2\pi} \sqrt{\left(\frac{p'_{nm}}{a}\right)^2 + \left(\frac{l\pi}{h}\right)^2}$$

- Resonant frequency for TM_{nml}-modes:

$$f_{\text{TM},nml} = \frac{c_0}{2\pi} \sqrt{\left(\frac{p_{nm}}{a}\right)^2 + \left(\frac{l\pi}{h}\right)^2}$$

Note that the TM₀₁₀- mode is *independent of the cavity height h* (until the TE₁₁₁ mode shows up, roughly at $2a/h=1$).

Recall the Bessel function and the derivative of the Bessel function with their roots:



$p_{01} = 1^{\text{st}}$ root of Bessel function of first type J_0

$p_{02} = 2^{\text{nd}}$ root of Bessel function of first type J_0

$p'_{01} = 1^{\text{st}}$ root of derivative of Bessel function of first type J'_0

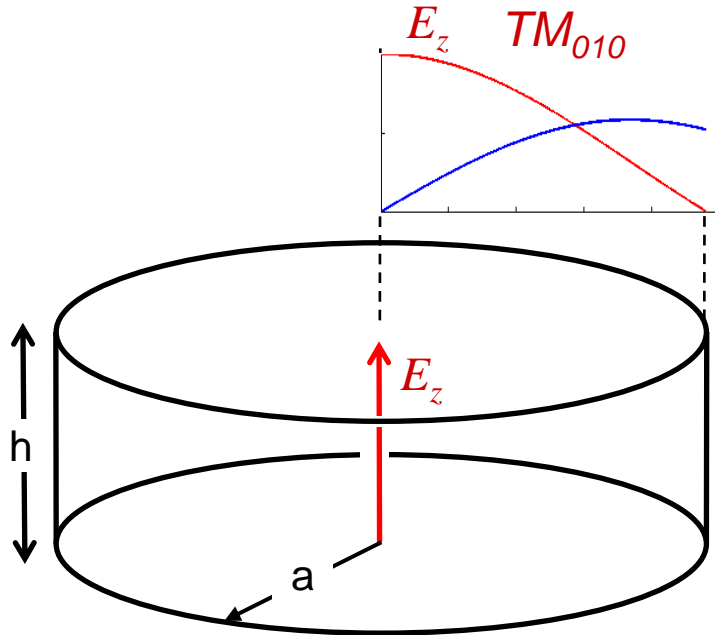
TABLE 3.3 Values of p'_{nm} for TE Modes of a Circular Waveguide

n	p'_{n1}	p'_{n2}	p'_{n3}
0	3.832	7.016	10.174
1	1.841	5.331	8.536
2	3.054	6.706	9.970

TABLE 3.4 Values of p_{nm} for TM Modes of a Circular Waveguide

n	p_{n1}	p_{n2}	p_{n3}
0	2.405	5.520	8.654
1	3.832	7.016	10.174
2	5.135	8.417	11.620

Source: Pozar, *Microwave Engineering*, 4th ed., Wiley



For the case of the TM_{010} -mode, we have no dependence on cavity height, so we get a much simpler formulae:

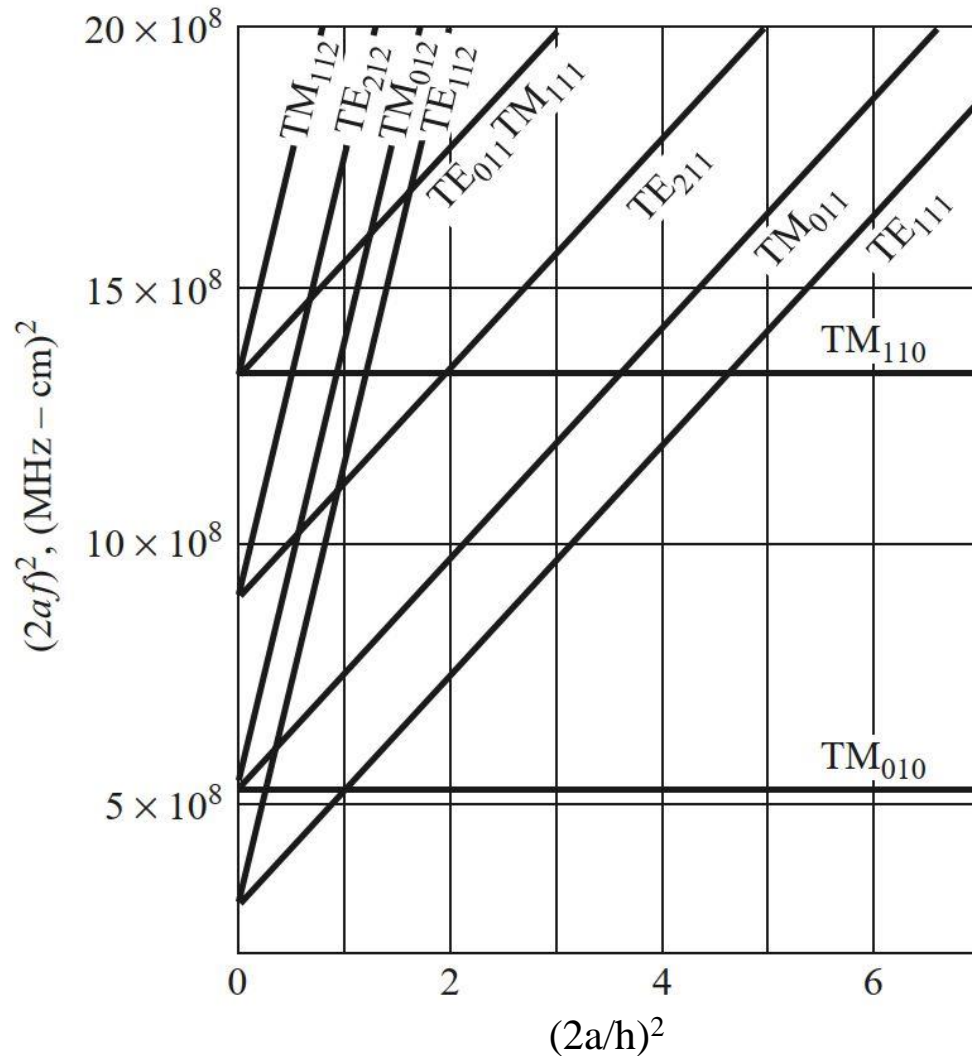
$$0.383 \lambda_{TM,010} = a$$

$$Q = \left(0.383 \frac{\lambda_{TM,010}}{\delta} \right) \left[1 + \left(0.383 \frac{\lambda_{TM,010}}{h} \right) \right]^{-1}$$



$$Q = \frac{0.383 \lambda_{TM,010}}{\delta} \left[1 + \frac{a}{h} \right]^{-1} = \frac{a}{\delta} \left[1 + \frac{a}{h} \right]^{-1}$$

Remember? Skinddepth is: $\delta = \sqrt{\frac{2}{\omega\sigma\mu}}$



Resonant chart for a general cylindrical cavity showing the excited modes as a function of cavity dimensions.

→ 1st TE-mode is TE111

→ 1st TM-mode is TM010, and shows up for ratios $(2a/h)^2 > 1$

Source: Pozar, *Microwave engineering*, 4th ed., Wiley

Thank you for your attention.

Let's have a break!

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